

Multi-Photon Decay of the Higgs Boson at the LHC

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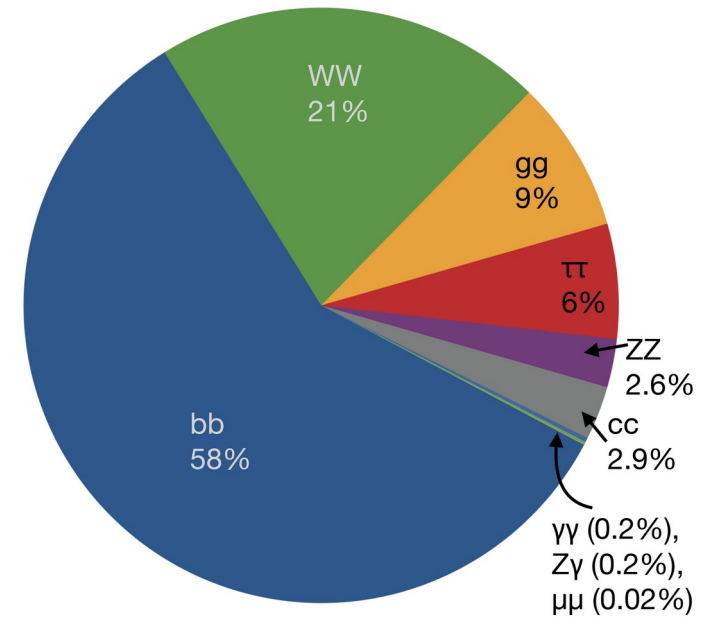
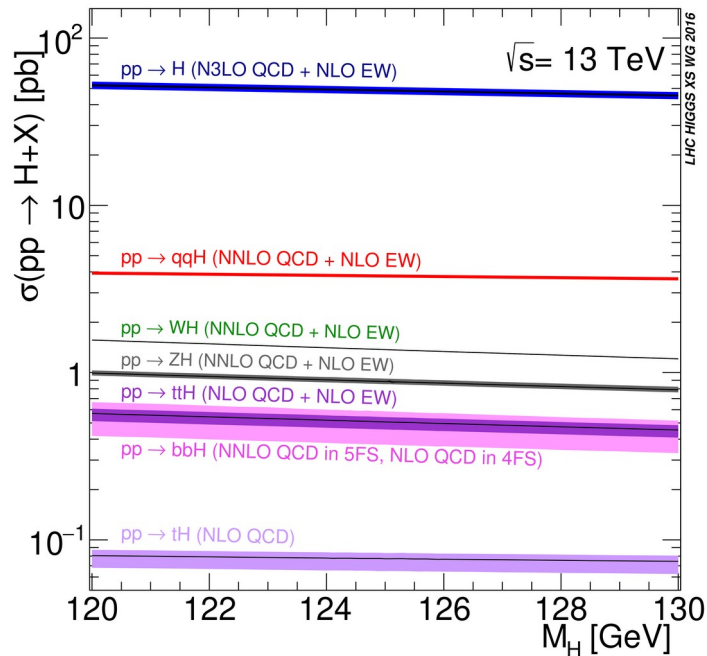
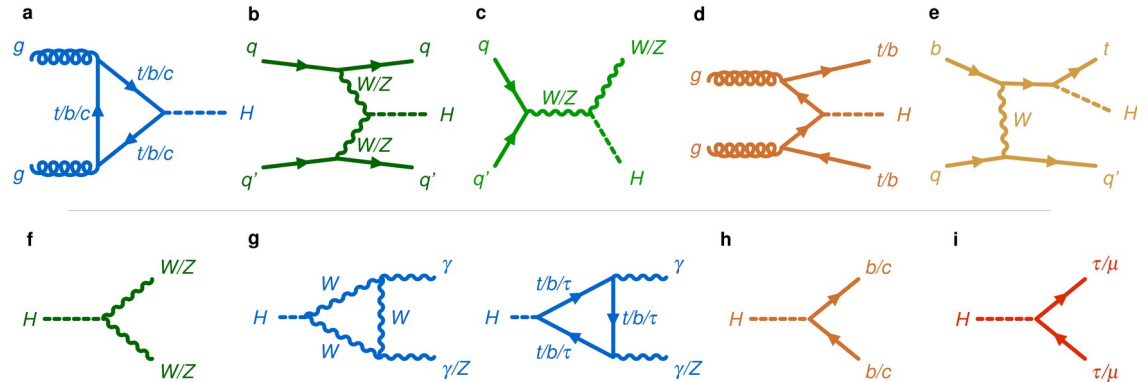
Hye-Sung Lee, Samuel D. Lane, **IML**, arXiv:2305.00013

Outline

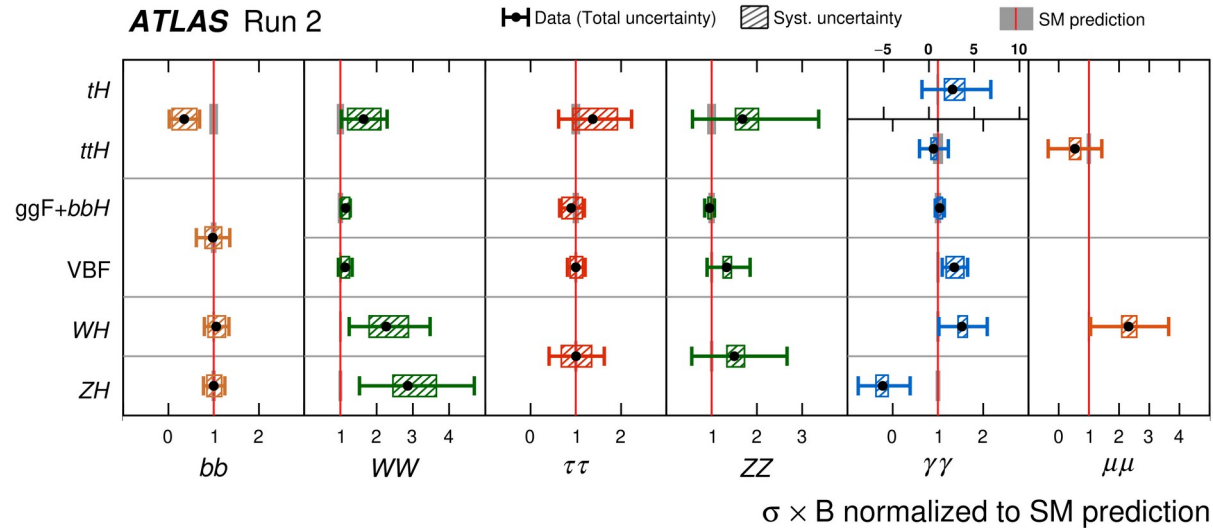
- Motivation.
- Introduction to dark axion model
 - Higgs to six photon decay.
- Classification of signals:
 - Events appearing as 2 photon, 4 photon, 6 photon, and non-isolated photon resonances.
- Detector efficiencies.
- Discussion of different signals and current constraints.
 - 4 photon signal: apparent Landau-Yang theorem violation.
- Conclude

Completely Predictive

- Once we know Higgs mass, can calculate production and decay rates to high precision.



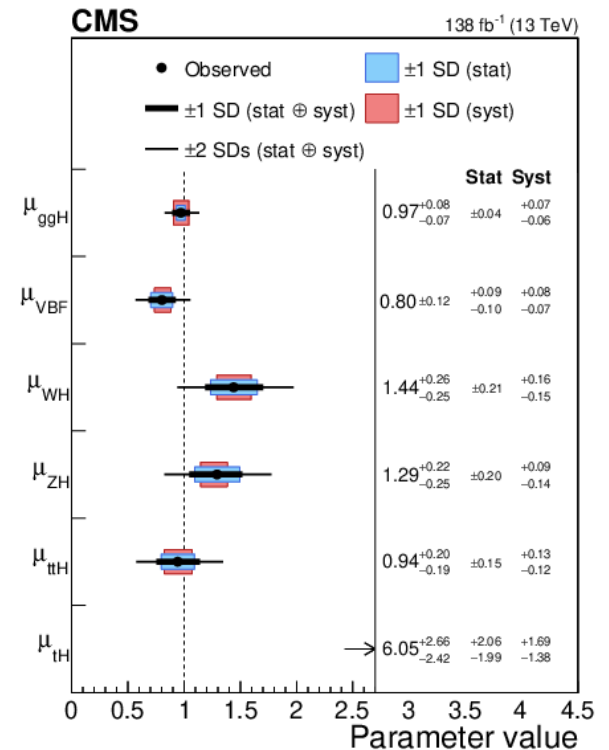
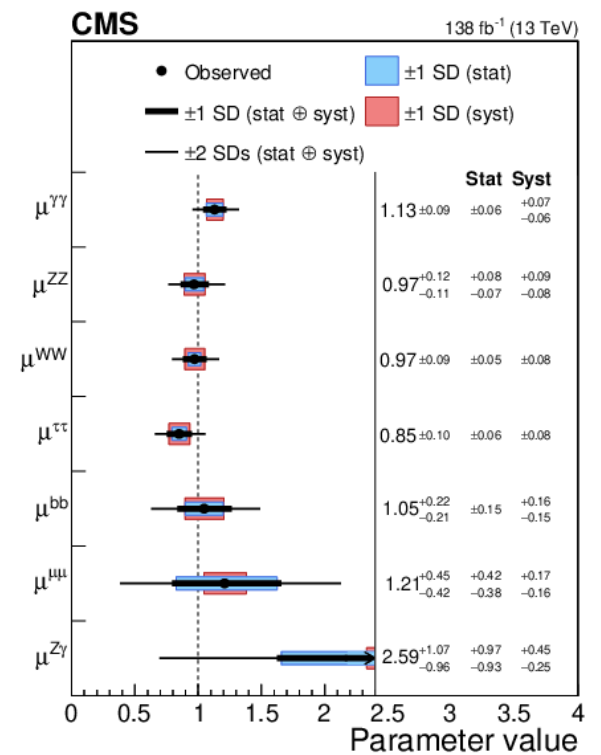
Where We're At



- Higgs precision shown in terms of signal strengths.

$$\mu_i^f = \frac{\sigma_{\text{exp}}(i \rightarrow h \rightarrow f)}{\sigma_{\text{SM}}(i \rightarrow h \rightarrow f)}$$

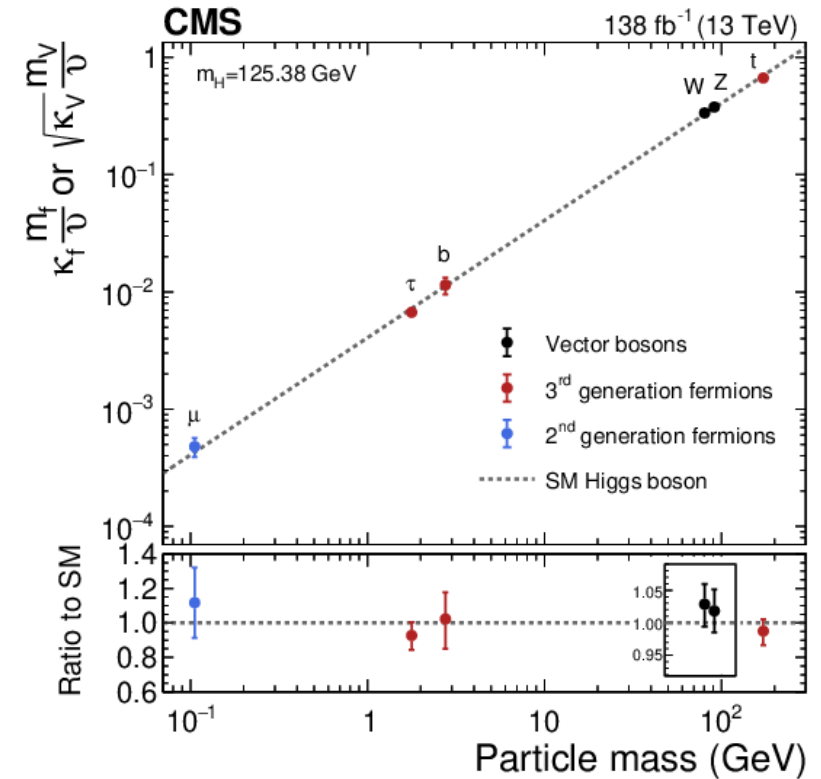
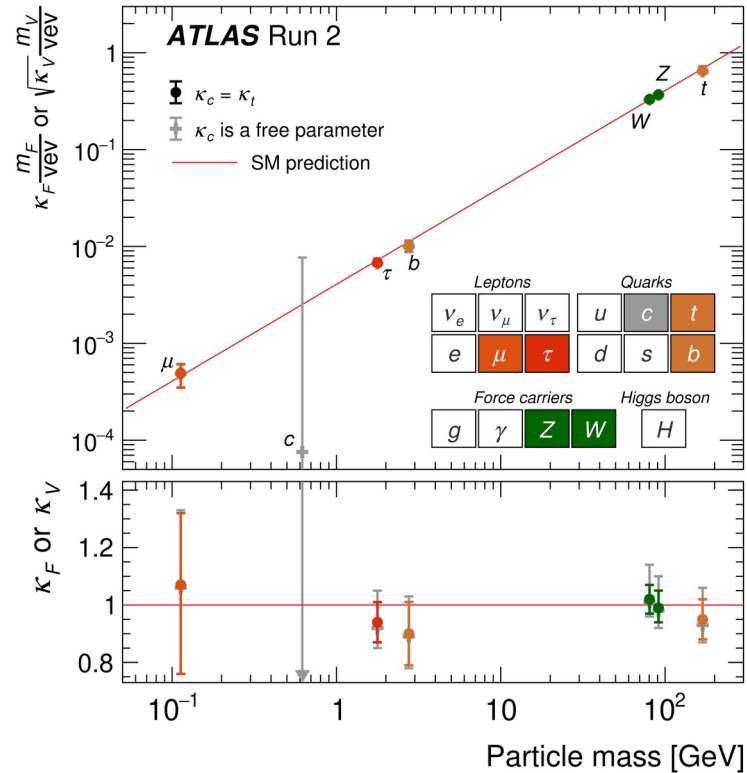
- Most measurements agree well with the SM.
- Already measuring many rates to order 10% accuracy.



Where We're At

CMS, Nature, arXiv:2207.00043

ATLAS, Nature, arXiv:2207.00092



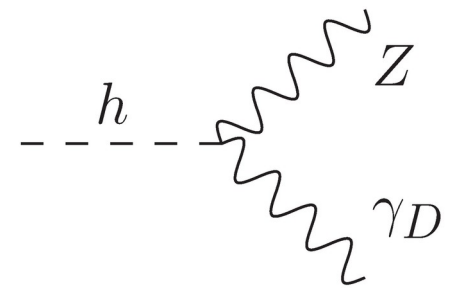
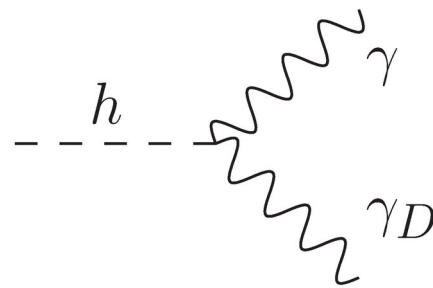
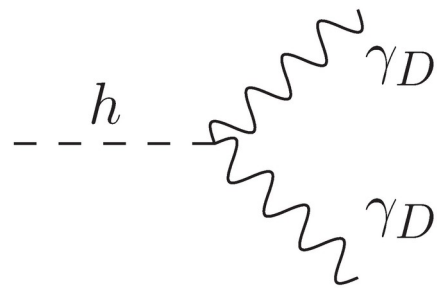
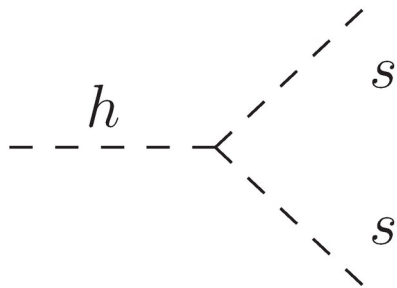
- Higgs boson associated with the source of EW symmetry breaking.
 - Higgs couplings should be proportional to masses (or squared masses)
 - Have verified that over many orders of magnitude in couplings and mass.

Where From Here?

- May still hope for deviations from SM physics.
- The Higgs width is very small: $\Gamma = 4.1 \text{ MeV}$
 - Generic partial width (ignoring longitudinal enhancements):

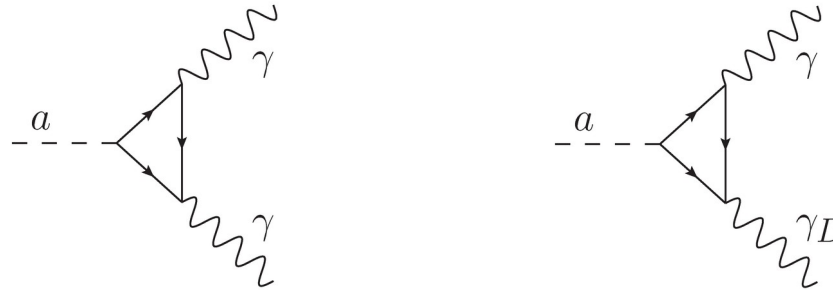
$$\Gamma \sim \frac{g^2}{8\pi} m_h \sim g^2 (5 \text{ GeV})$$

- Decays of Higgs into exotic states can be sensitive to small couplings.
- Many different types of portals (non-exhaustive list):
 - Scalar portal: $h s s$
 - Dark photon from scalar mixing: $h \gamma_D \gamma_D$
 - Dark photon vector boson mass mixing: $h Z \gamma_D$
 - Dark photon from loops: $h F_{D,\mu\nu} F_D^{\mu\nu}$ $h F_{D,\mu\nu} F^{\mu\nu}$
- Can give many different interesting Higgs decays:

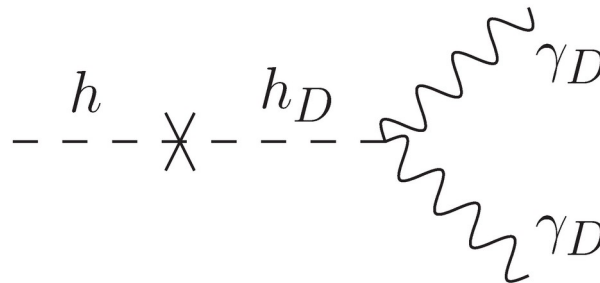


Dark Axion Portal

- Interested in exploring dark axion portal [Kaneta, Lee, Yun PRL118 \(2017\)](#)
 - Assume some heavy new physics that is charged under the SM electroweak gauge group and some new $U(1)_d$.
 - Introduce an axion like particle: a .
 - Can get interesting couplings between ALPs, photons, and dark photon:



- Assume new $U(1)_d$ is Higgsed, and SM Higgs can couple to dark photons via mixing?

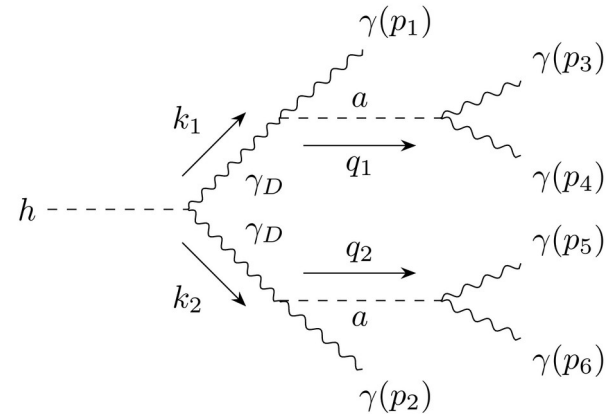
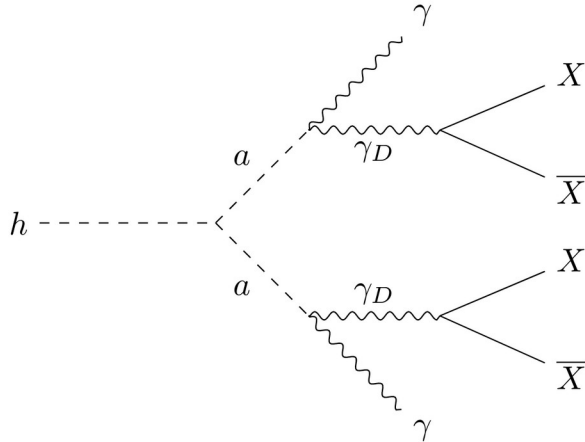


- Parameterize couplings as (and promptly forget about any UV completion):

$$L_{portal} = \frac{1}{4} G_{a\gamma\gamma} a F^{\mu\nu} \tilde{F}_{\mu\nu} + \frac{1}{2} G_{a\gamma\gamma_D} a F^{\mu\nu} \tilde{F}_{D\mu\nu} + \frac{1}{4} G_{a\gamma_D\gamma_D} a F_D^{\mu\nu} \tilde{F}_{D\mu\nu} + \frac{1}{2} \lambda_{h\gamma_D\gamma_D} h \gamma_D^\mu \gamma_{D\mu}$$

Dark Axion Portal

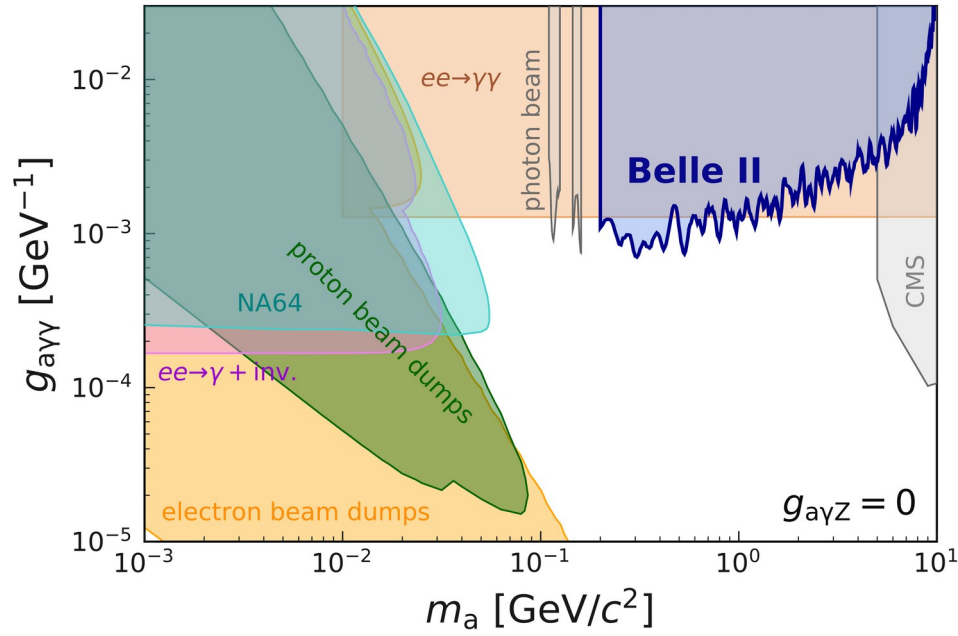
- Depending on mass ordering, can have many different interesting exotic decays:



- We will focus on six photon decay.
 - Will assume the dominant decays: $\gamma_D \rightarrow a \gamma$ $a \rightarrow \gamma \gamma$
- Depending on mass hierarchies, dark photon and ALP can be very boosted .
 - Many different combinations of collimated photons.
 - Will classify signals according to:
 - Photon-jets: when two photons would hit the same electromagnetic calorimeter cell.
[Draper, McKeen PRD85 \(2012\); Ellis, Roy, Scholtz PRD87 \(2013\), etc.](#)
 - Non-isolated photons: ξ -jets
[Chakraborty, Iyer, Roy, NPB932 \(2018\), Sheff, Steinberg, Wells PRD104 \(2021\), etc.](#)

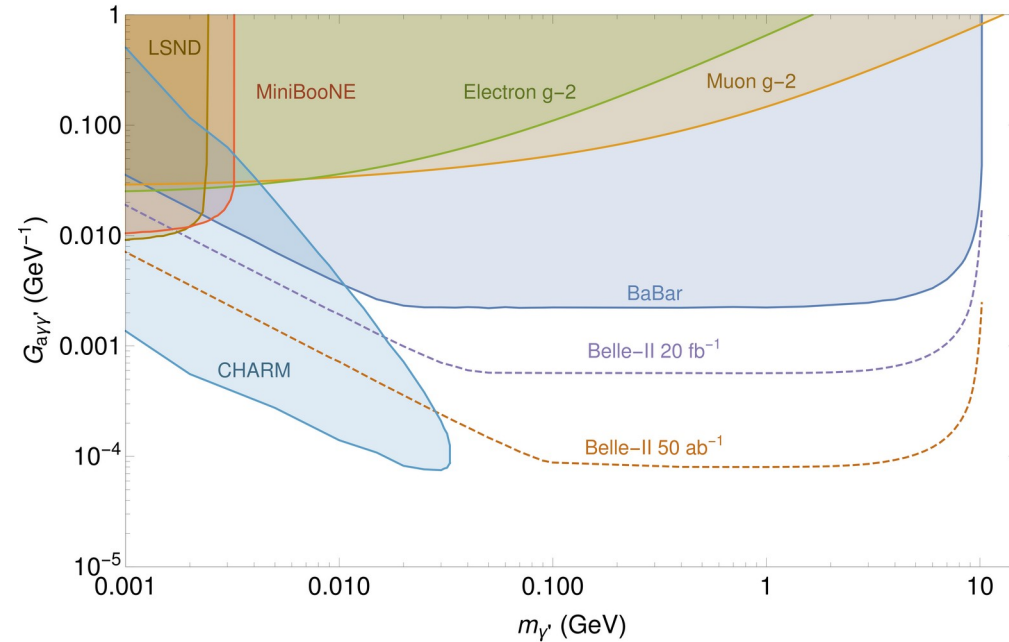
Constraints

ALP Constraints



Belle-II PRL125 (2020)

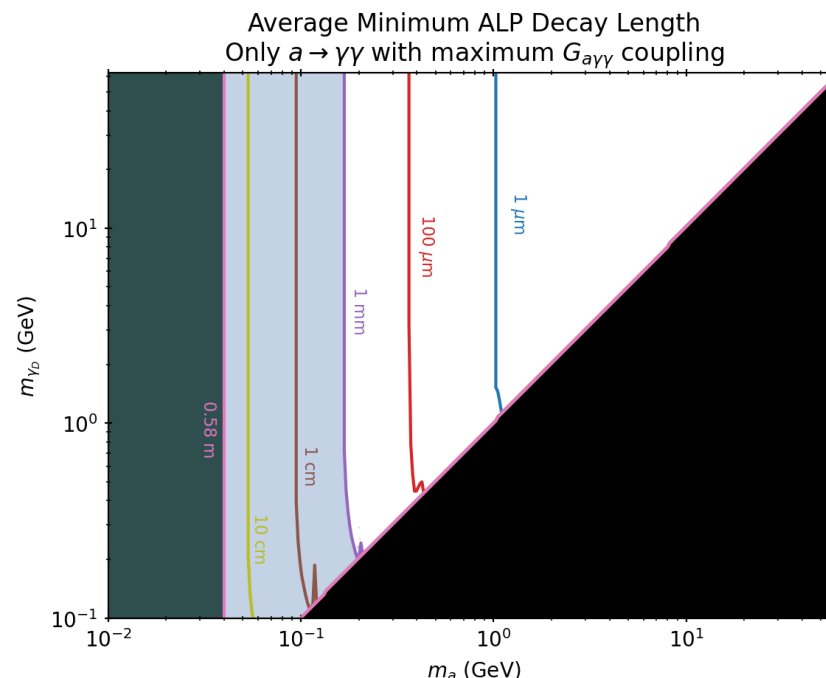
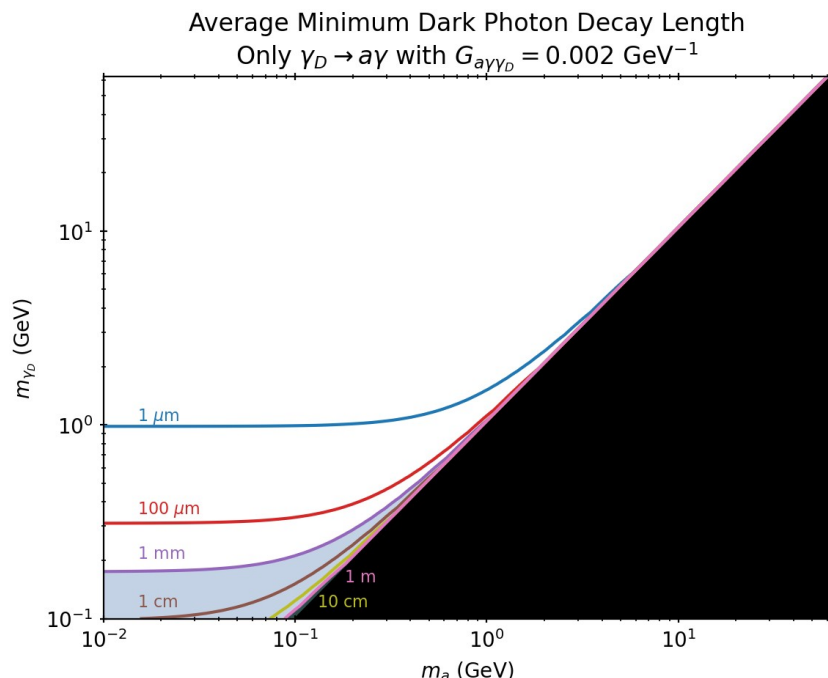
Dark ALP Constraints



DeNiverville, Lee, Seo, PRD98 (2018);
Deniverville, Lee, Lee PRD103 (2021)

- Very small couplings, can make dark photon or ALP long lived....

Decay Lengths for $h \rightarrow 2 \gamma_D \rightarrow 2 a 2 \gamma \rightarrow 6 \gamma$ Lee, Lane, IML, arXiv:2305.00013



- Average decay lengths of dark photon and ALP in Higgs rest frame.
 - Set couplings to the maximum allowed by current constraints.
 - Assume 100% branching ratios for: $\gamma_D \rightarrow a \gamma$ $a \rightarrow \gamma \gamma$
- For ALP masses below 40 MeV:
 - Beam dump/astrophysical/cosmological constraints make ALP long lived on LHC time scales.
- Will only consider ALPs with mass above 40 MeV.
 - Even then many ALPs/dark photons still decay within the tracker.
 - However decay products are photons, which don't leave tracks in the tracker.
 - Will not consider then as displaced vertices (there are subtleties about photon conversions...)
 - For a proposal to use photon conversions in the tracker to tag displaced ALP vertices see

Classification of Signatures

- Photon jets:

- Two photons would hit the same electromagnetic calorimeter cell.
- Defined as photons with angular separation:

$$\Delta R \leq 0.04$$

- After forming photon jets:
 - Collectively call remaining photons and photon jets as “observed photons”

- Intermediately separated photons:

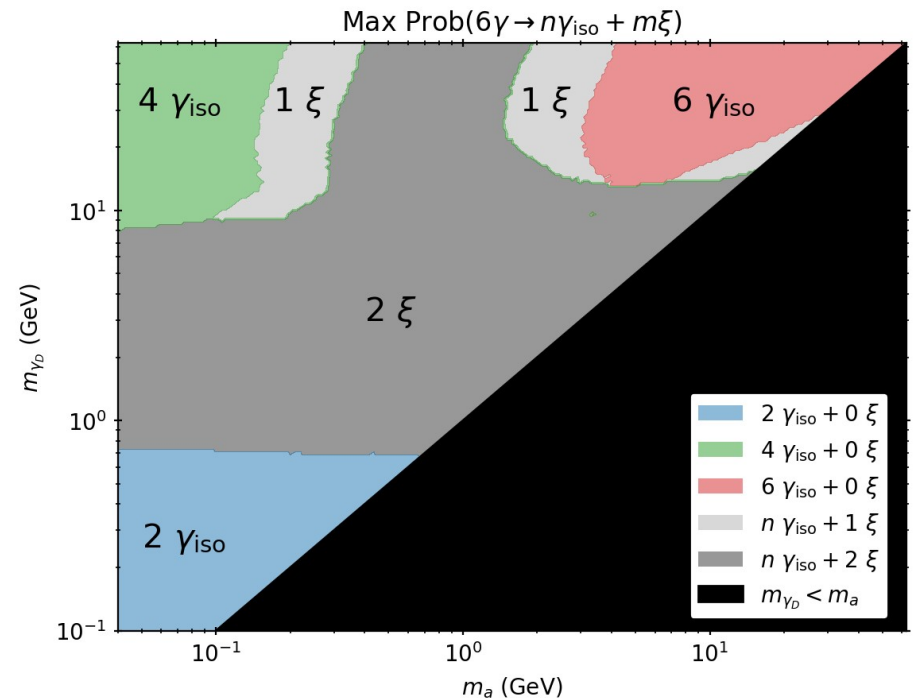
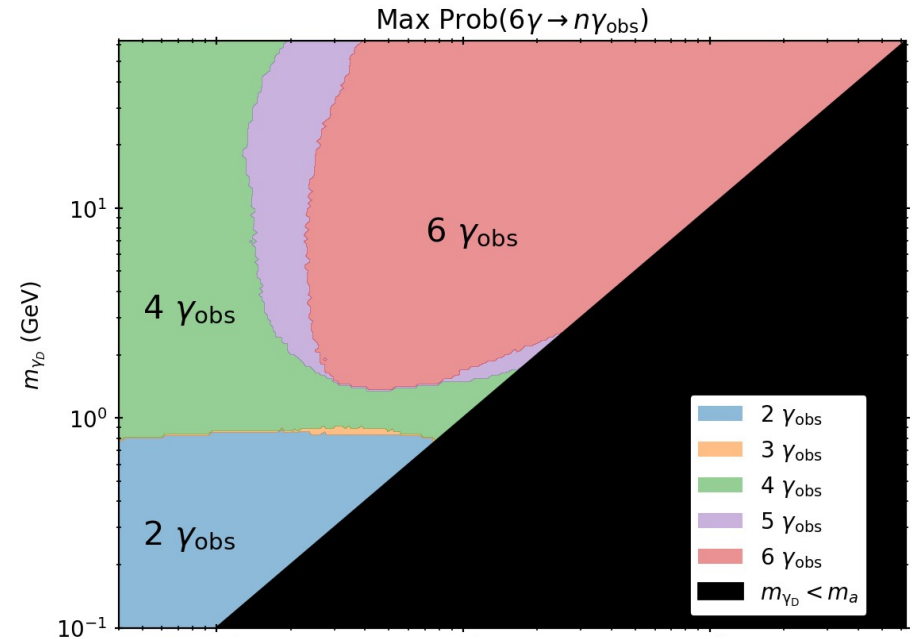
- Will label as ξ -jets.
- Define ξ -jets as observed photons with angular separation:

$$0.04 \leq \Delta R \leq 0.4$$

- After forming ξ -jets, any remaining photons or photon-jets are called “isolated photons”

- Angular separation defined as:

$$\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$$



- Between 6 and 4-photons:

- $m_a \ll m_{\gamma_D} \approx m_h/2$

- Dark photon at rest, ALP essentially massless, so approximate:

$$E_a \approx m_h/4$$

- Calculate angular separation of ALP decay products:

$$\Delta R_{\gamma\gamma} \approx 2 \frac{m_{\gamma_a}}{E_a} \approx 8 \frac{m_a}{m_h} \leq 0.04$$

- Gives constraint:

$$m_a \lesssim \frac{m_h}{200} \approx 0.625 \text{ GeV}$$

- $m_a \approx m_{\gamma_D}$

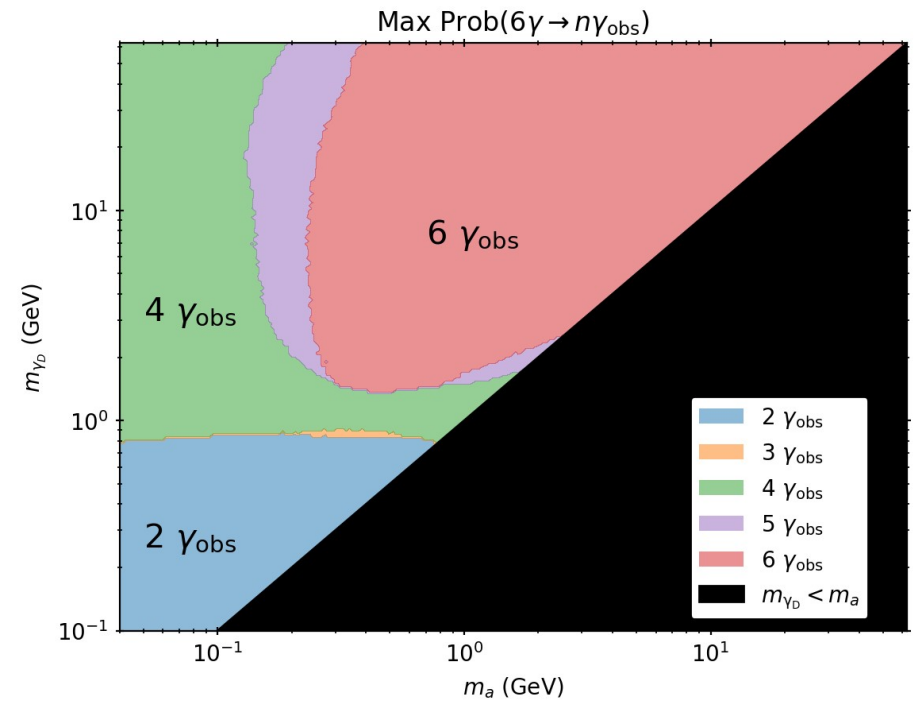
- ALP at rest in dark photon rest frame, photon from dark photon decay very soft.
 - Boosting to Higgs rest frame, ALP (roughly) inherits dark photon energy:

$$E_a \approx m_h/2.$$

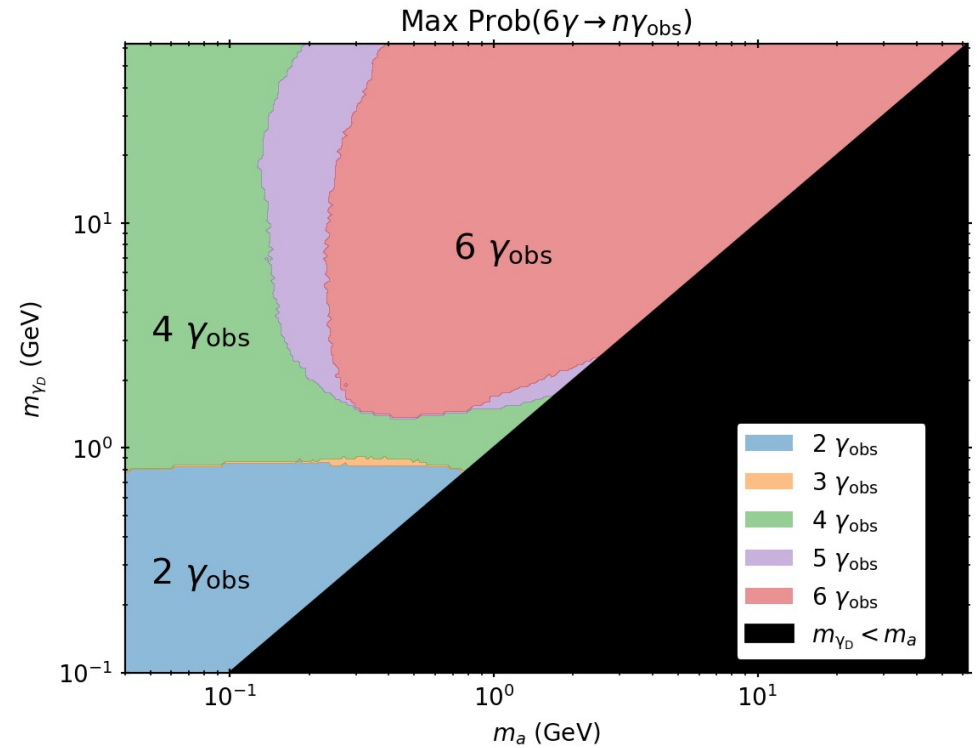
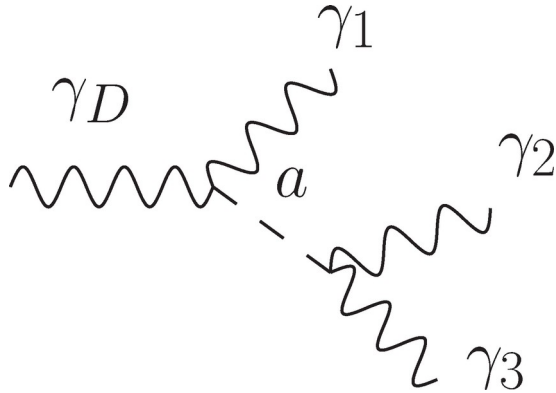
- Angular separation about axion decay products:

$$\Delta R_{\gamma\gamma} \approx 4 \frac{m_a}{m_h} \leq 0.04$$

- Gives constraint: $m_{\gamma_D} \approx m_a \lesssim \frac{m_h}{100} \approx 1.25 \text{ GeV}.$



- Between 4 and 2-photons:
 - Need to calculate angular separation between photons from dark photon decay.
 - Configuration where photons are most widely separated is when dark photon and ALP decay planes coincide:



- Assume $\Delta R_{\gamma_2 \gamma_3} \leq 0.04$
- To obtain $\Delta R_{\gamma_1 \gamma_3} \leq 0.04$ need angular separation between photon from dark photon and ALP:

$$\Delta R_{\gamma_1 a} \approx \frac{2 m_{\gamma_D}}{E_{\gamma_D}} \lesssim 0.02$$

- Gives a constraint on the dark photon mass: $m_{\gamma_D} \lesssim \frac{m_h}{200} \approx 0.625 \text{ GeV}$

Do the Signatures Survive in the Detector?

- Ensure photons land in barrel or endcap:
 - ATLAS: $|\eta| < 1.37$ or $1.52 < |\eta| < 2.5$
 - CMS: $|\eta| < 1.44$ or $1.57 < |\eta| < 2.5$
- Require some minimum transverse momentum on photons:

Channel	ATLAS p_T Requirements
1γ	$p_{1,T} > 150 \text{ GeV}$ [94]
2γ	$p_{1,T} > 35 \text{ GeV}$ and $p_{2,T} > 25 \text{ GeV}$ [24]
3γ	$p_{1,T} > 15 \text{ GeV}$, $p_{2,T} > 15 \text{ GeV}$, and $p_{3,T} > 15 \text{ GeV}$ [95]
4γ	$p_{1,T} > 30 \text{ GeV}$, $p_{2,T} > 18 \text{ GeV}$, $p_{3,T} > 15 \text{ GeV}$, and $p_{4,T} > 15 \text{ GeV}$ [95]
5γ	$p_{i,T} > 15 \text{ GeV}$ ($i = 1, 2, 3, 4, 5$)
6γ	$p_{i,T} > 15 \text{ GeV}$ ($i = 1, 2, 3, 4, 5, 6$)

Channel	CMS p_T Requirements
1γ	$p_{1,T} > 145 \text{ GeV}$ [98]
2γ	$p_{1,T} > 30 \text{ GeV}$ and $p_{2,T} > 18 \text{ GeV}$ [30]
3γ	$p_{1,T} > 15 \text{ GeV}$, $p_{2,T} > 15 \text{ GeV}$, and $p_{3,T} > 15 \text{ GeV}$ [95]
4γ	$p_{1,T} > 30 \text{ GeV}$, $p_{2,T} > 18 \text{ GeV}$, $p_{3,T} > 15 \text{ GeV}$, and $p_{4,T} > 15 \text{ GeV}$ [27]
5γ	$p_{i,T} > 15 \text{ GeV}$ ($i = 1, 2, 3, 4, 5$)
6γ	$p_{i,T} > 15 \text{ GeV}$ ($i = 1, 2, 3, 4, 5, 6$)

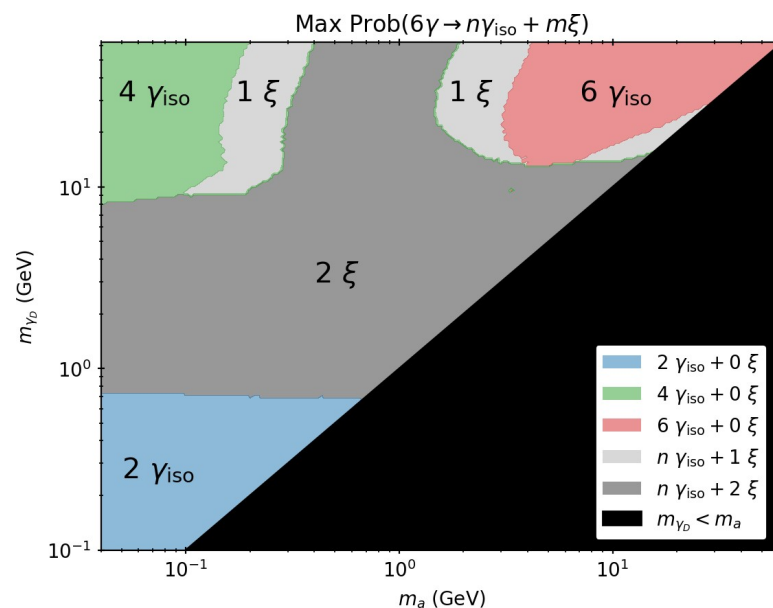
- No 6-photon triggers, so we estimated them choosing pTs such that photon identification is at 50-75%

Detector Efficiencies

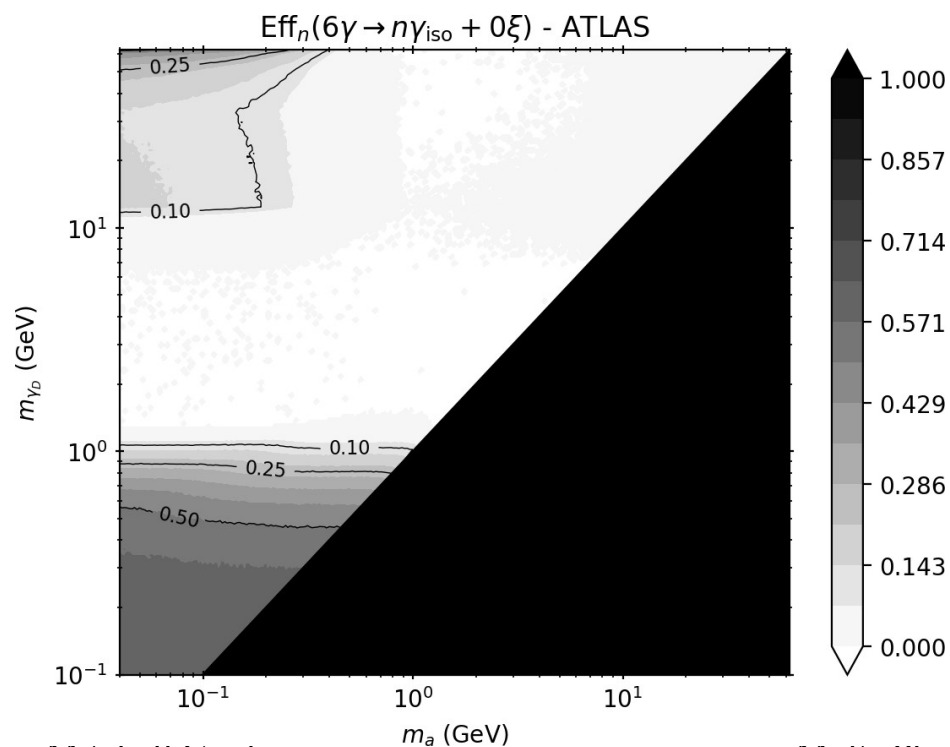
- Define a detector efficiency as percentage of signal events with n -isolated photons and m ξ -jets that pass a k -photon trigger:

$$\text{Eff}_k(h \rightarrow n\gamma_{\text{iso}} + m\xi)$$

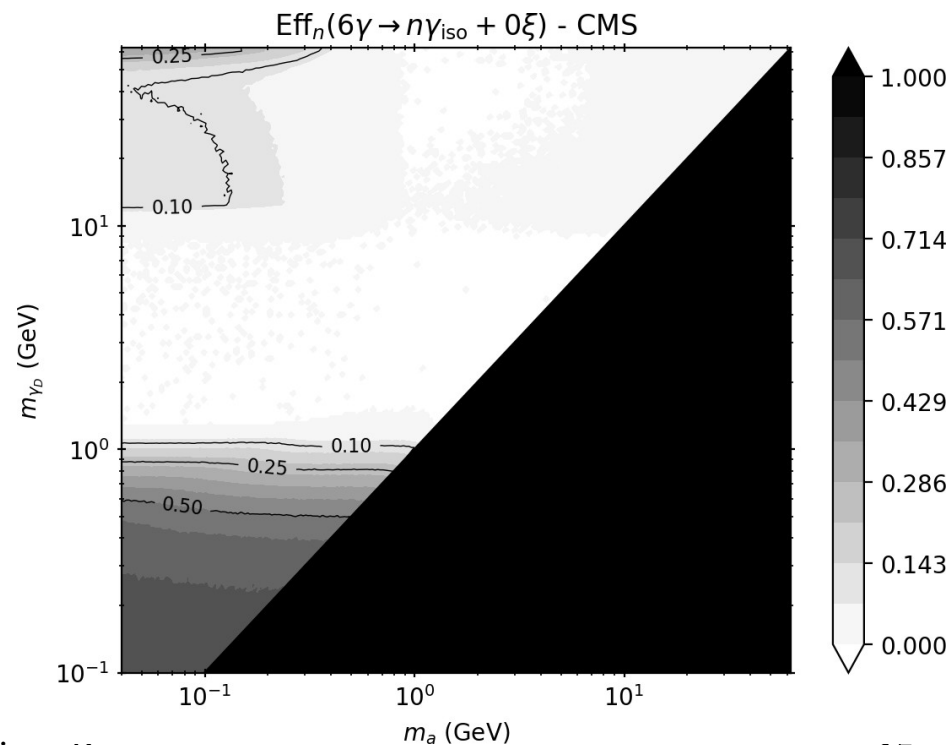
- n -isolated photon with no ξ -jet events passing n -photon triggers



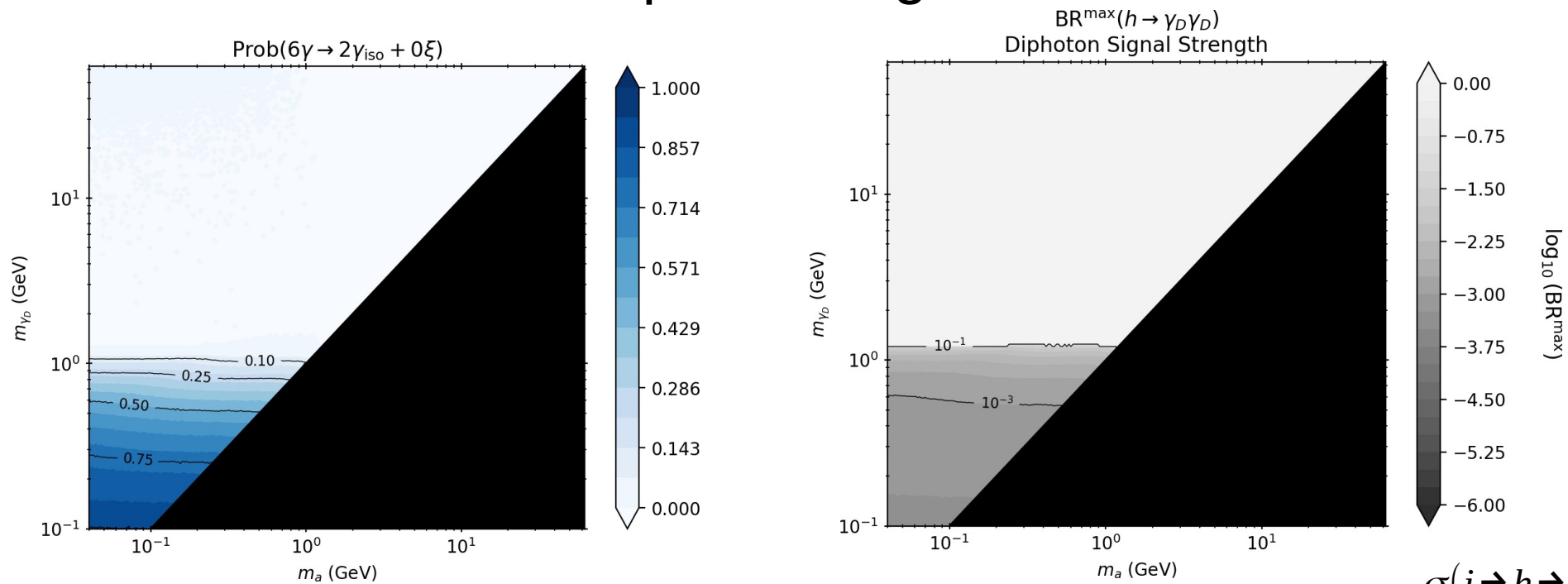
ATLAS



CMS



Di-photon signal



- Largest probability for di-photon signal at low dark photon masses.
- Signal strength measurements for ATLAS (CMS): [ATLAS 2207.00348](#), [CMS JHEP07 \(2021\)](#)

$$\mu = 1.04 \pm_{0.09}^{0.10} (1.12 \pm_{0.09}^{0.09})$$

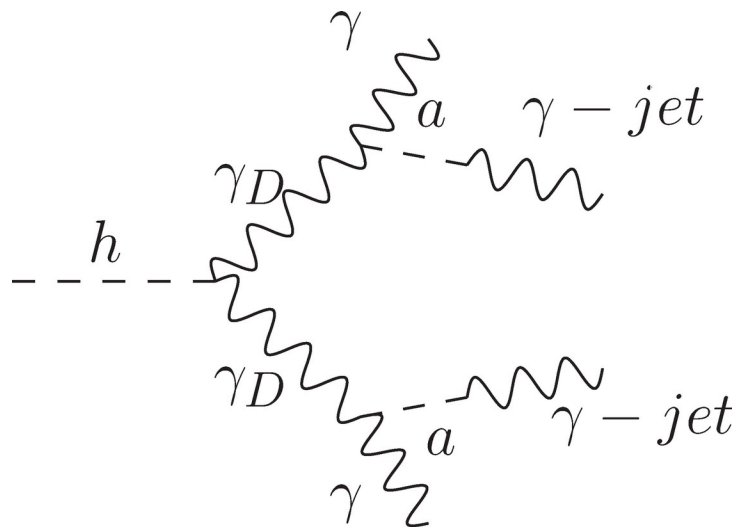
- BR into di-photons changed as:

$$\mu = \frac{1}{\text{BR}_{\text{SM}}(h \rightarrow \gamma\gamma)} \frac{\Gamma_{\text{SM}}(h \rightarrow \gamma\gamma) + \Gamma(h \rightarrow \gamma_D \gamma_D) \text{BR}^2(\gamma_D \rightarrow a\gamma) \text{BR}^2(a \rightarrow \gamma\gamma) \text{Prob}(6\gamma \rightarrow 2\gamma_{\text{iso}} + 0\xi)}{\Gamma_{\text{SM}}(h) + \Gamma(h \rightarrow \gamma_D \gamma_D)}$$

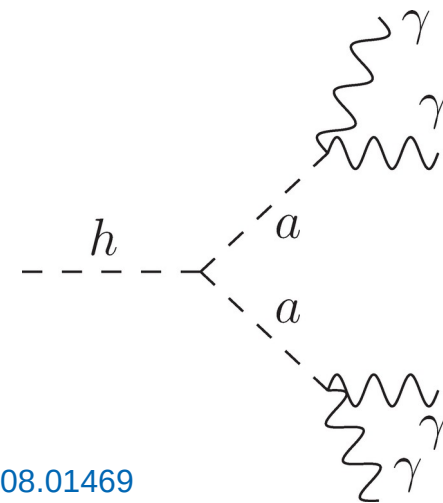
- Assume ALP and dark photon branching ratios as: $\text{BR}(\gamma_D \rightarrow a\gamma) = \text{BR}(a \rightarrow \gamma\gamma) = 1$
- From now on place limits on $\text{BR}(h \rightarrow \gamma_D \gamma_D)$

Four photon signal

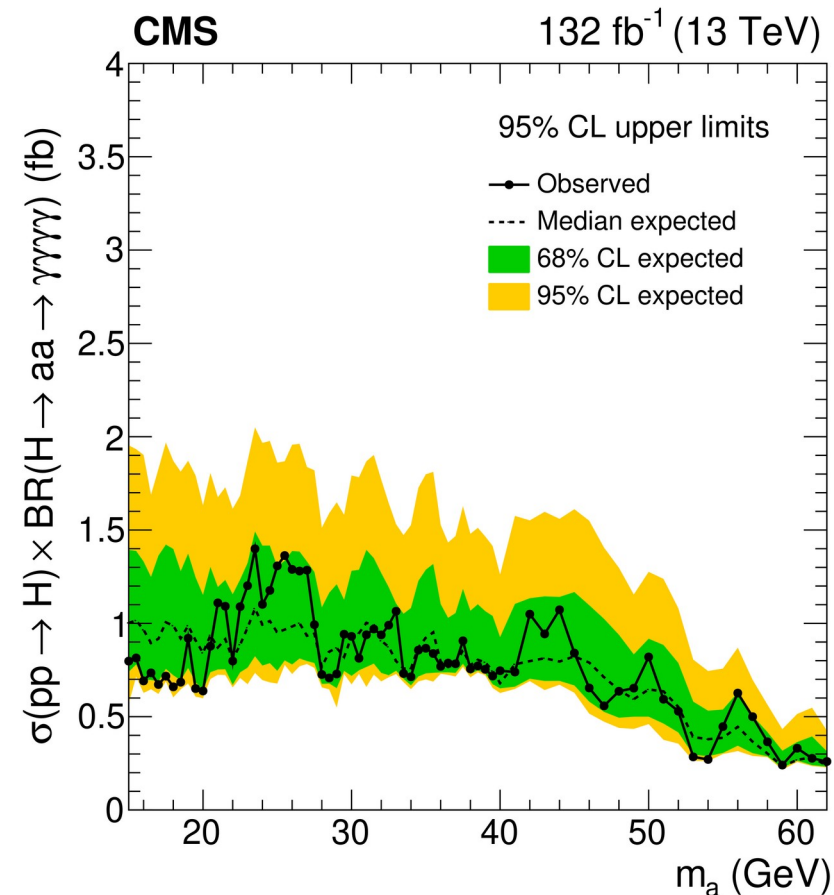
- Higgs to four photon has long been proposed [Draper, McKeen, PRD85 \(2012\)](#), [Curtin, et al PRD90 \(2014\)](#), etc.
- Unlike the typical scenario, our four isolated photons originate from a spin-1 particle:



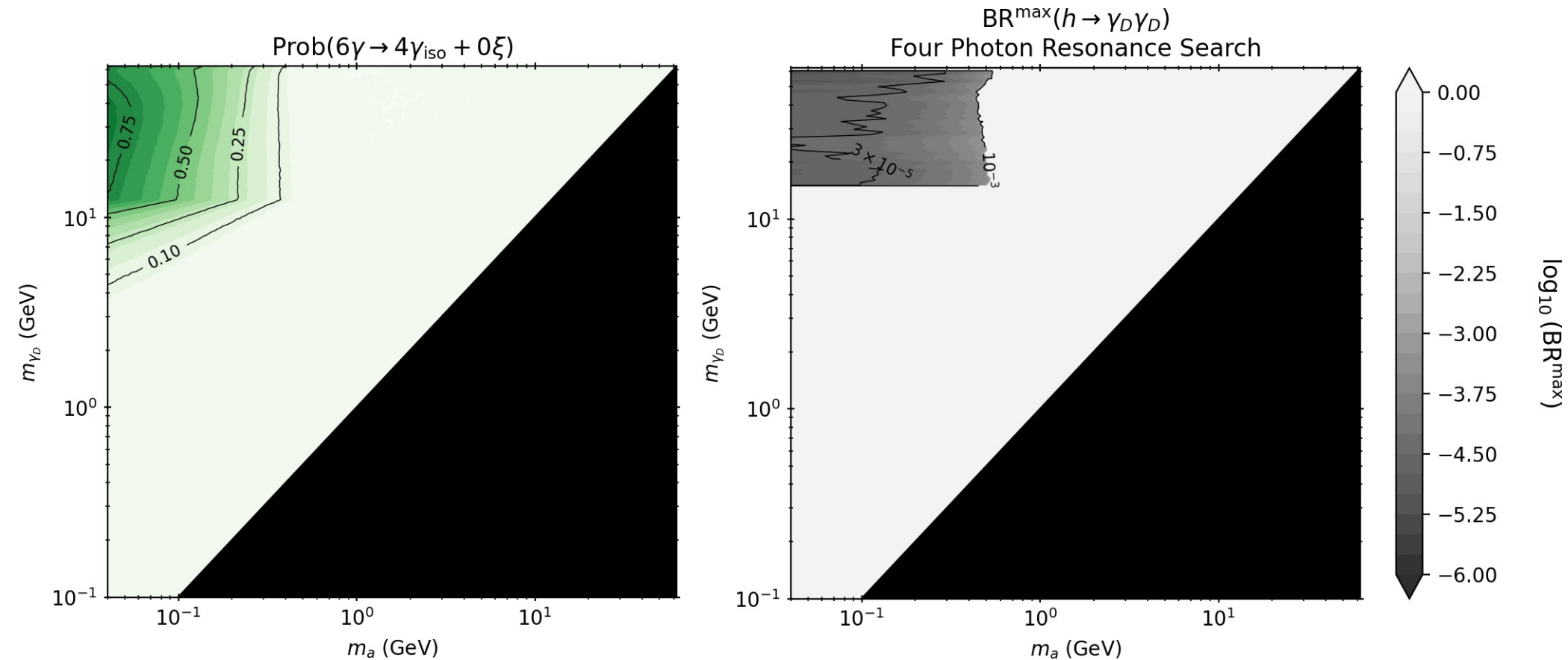
CMS 2208.01469



- An apparent violation of Landau-Yang theorem.
- Similar signals:
 - $Z \rightarrow \gamma a \rightarrow \gamma + \gamma - jet$ [Toro, Yavin PRD86 \(2012\)](#)
 - $Z' \rightarrow \gamma a \rightarrow \gamma + \gamma - jet$ [Chala, et al PLB755 \(2016\)](#)
- Detector resolutions are finite.

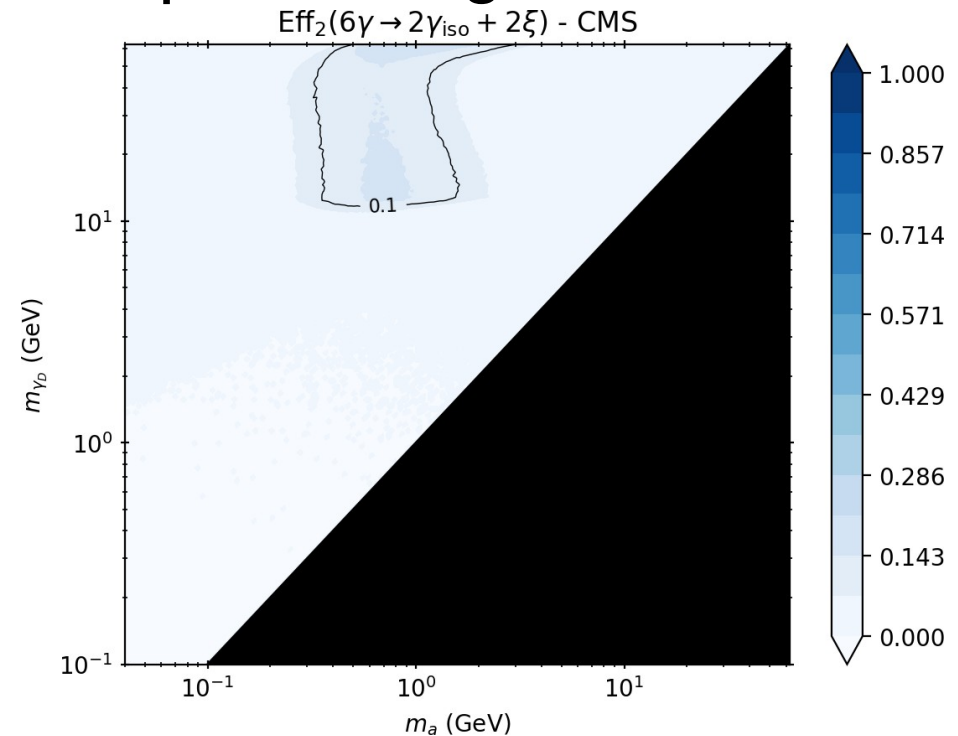
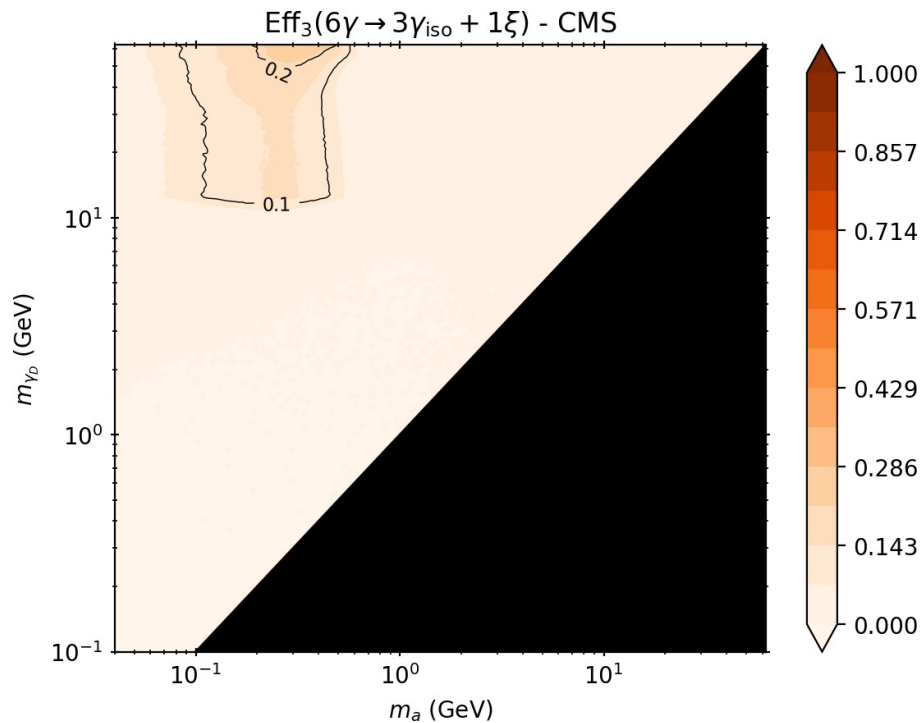


Four Photon Constraints



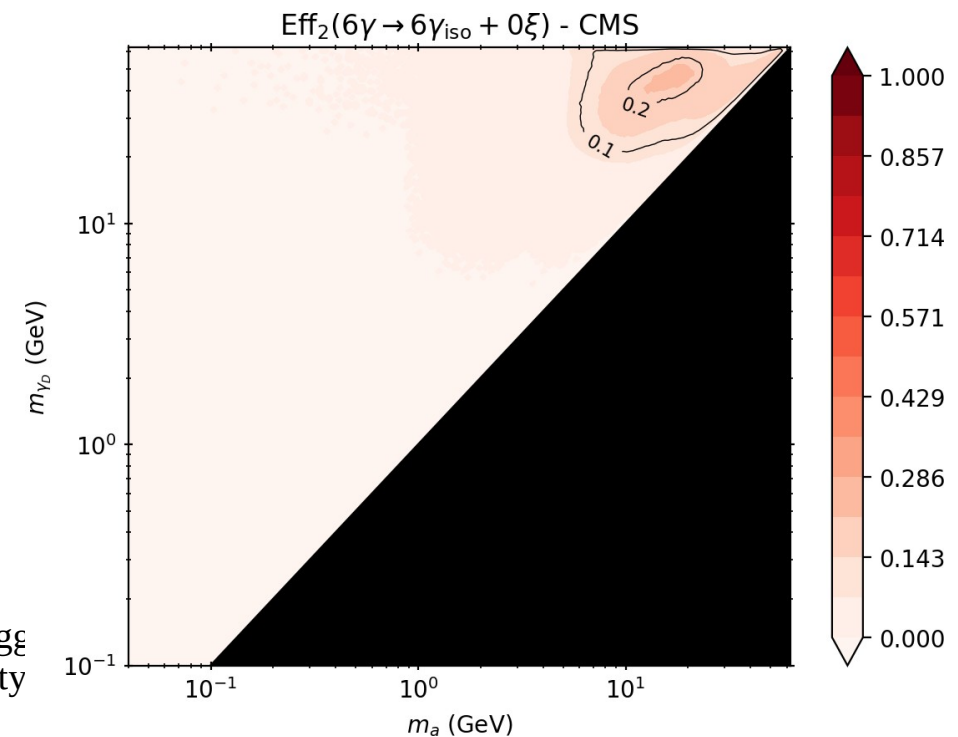
- Demand detector efficiencies greater than 1%
- Direct searches very constraining on this signal.
 - Limit branching ratio to less than 10^{-3} over relevant range.

Six photon and non-isolate photon signals

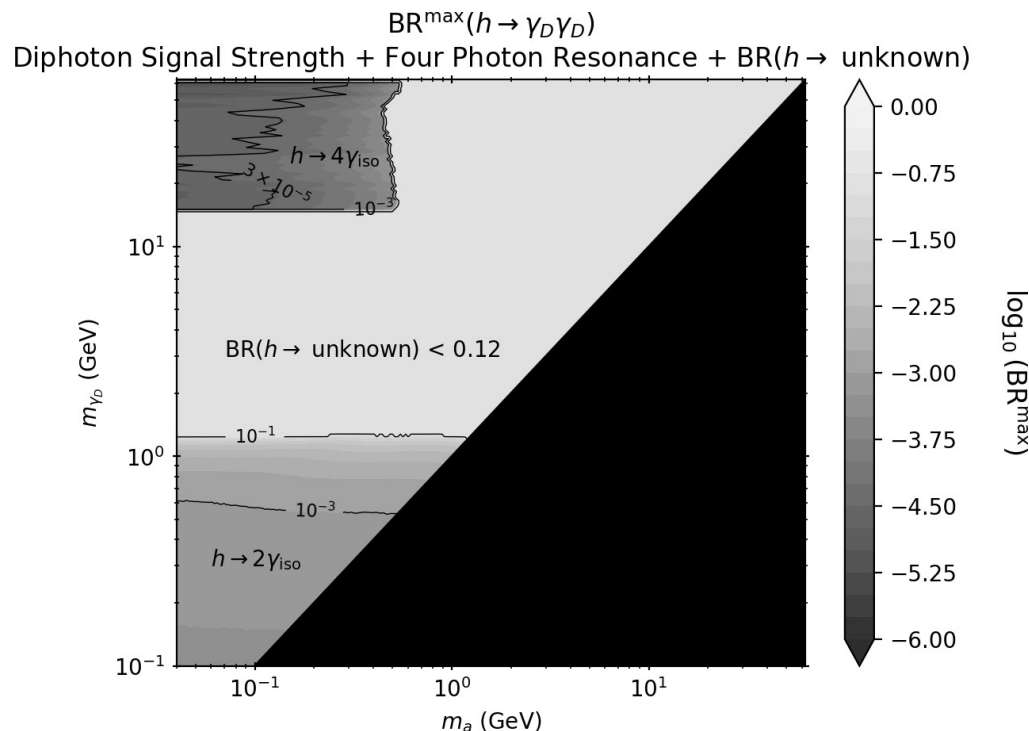


Lee, Lane, IML, arXiv:2305.00013

- 6-photon signal did not pass 6-photon triggers.
 - May be possible to lower trigger requirements.
- Showing the 2 and 3-photon triggers for 6-photon and ξ -jet signals
 - The photons that pass the trigger will not reconstruct the Higgs mass.
 - Would not survive a standard analysis.
 - However, it may be safe to “tape” and a reanalysis may find these signals.



Cumulative Result



- Six photon a ξ -jet signals will contribute to the total width of the Higgs.
 - In fits to Higgs precision measurements, assume there is some “undetermined” Higgs width into states not explicitly included in the fits.
 - Limit on the undetermined branching ratio [ATLAS, Nature, arXiv:2207.00092](#)

$$BR(h \rightarrow \text{unknown}) \leq 0.12$$

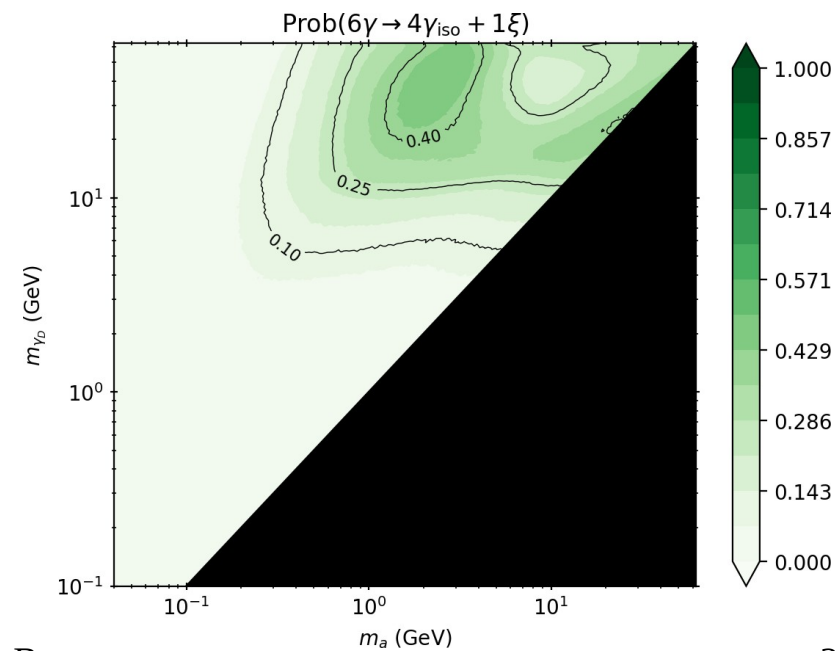
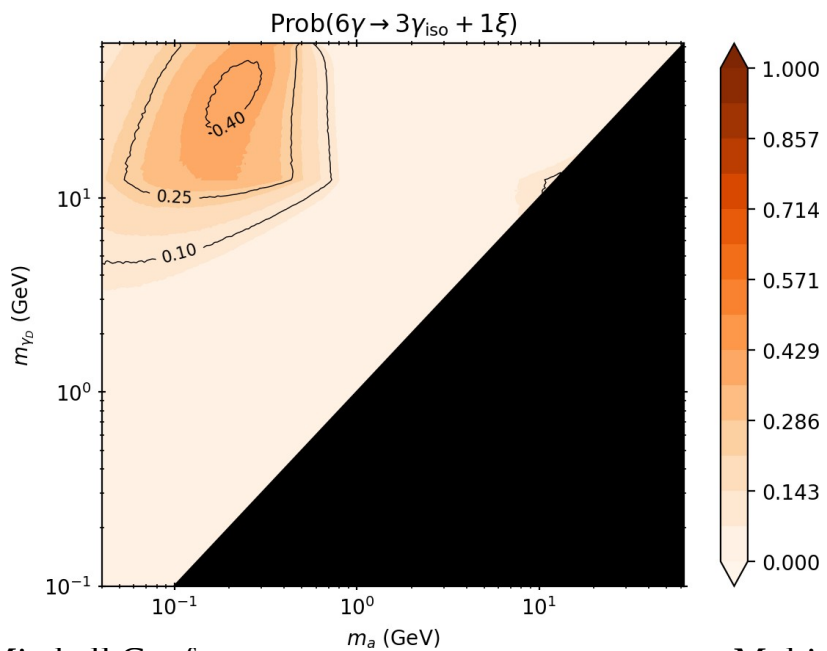
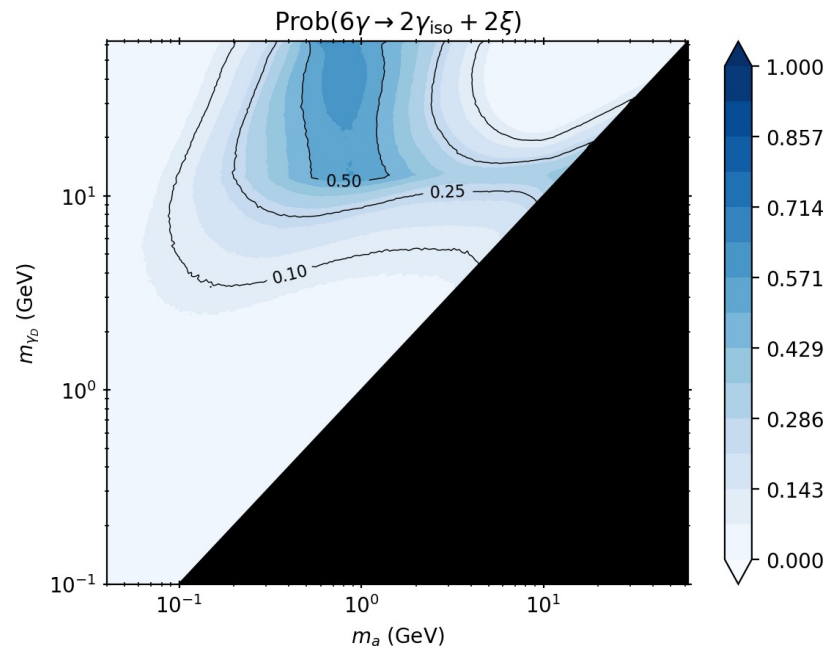
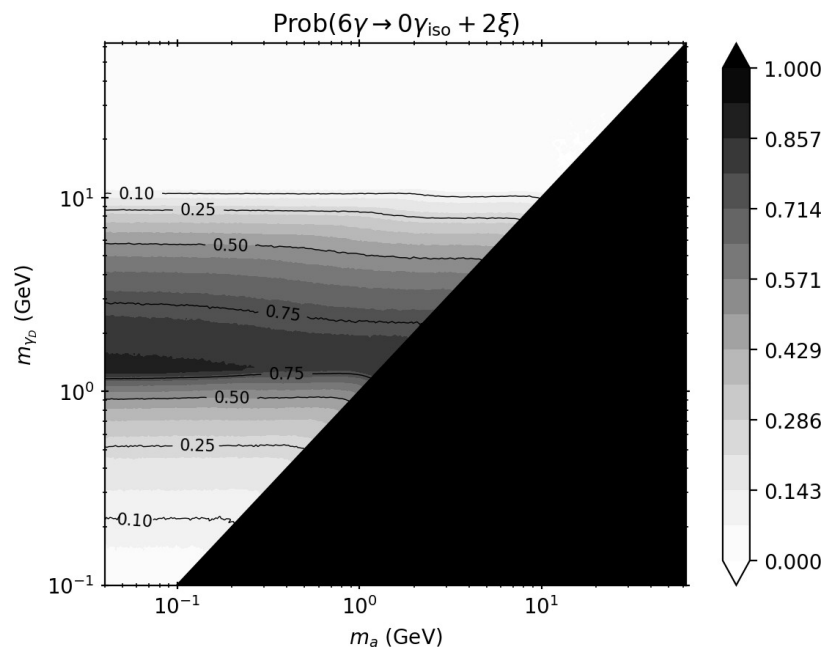
- Not relevant for two photon signal.
 - Changes di-photon partial width in addition to total width.
- Direct searches for four photons always better.

Conclusions

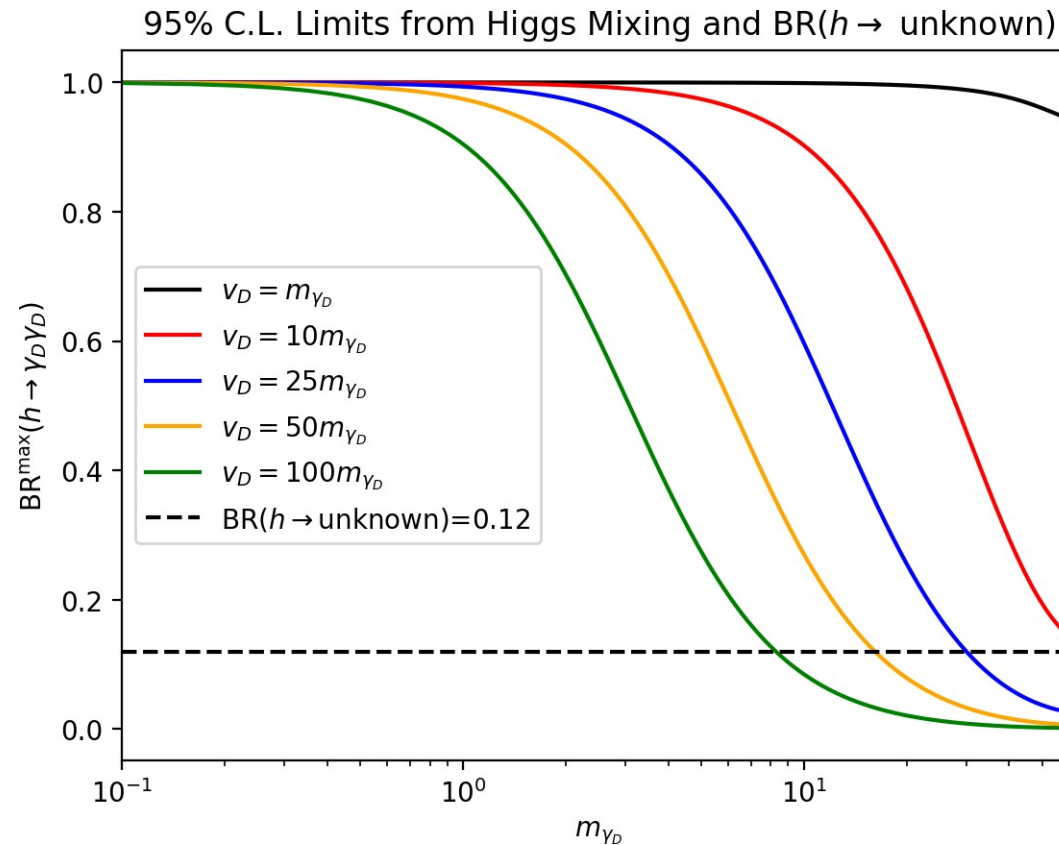
- In the past 11 years, Higgs physics has progressed greatly.
- May still have interesting signatures to search for.
 - May still expect deviations to SM predictions.
 - Due to small width, may still be sensitive decays through small couplings.
- Studied the dark axion model.
 - Have Higgs to six photon decay.
 - Can give rise to di-photon and four photon signatures.
 - Four-photon signature would be an apparent Landau-Yang theorem violation: dark photon decays into two photons.
 - More challenging signatures as well: six isolated photon and non-isolated photons.
 - Estimated a six photon trigger, however kills signal.
 - These events may have passed a two photon or three photon triggers.
 - Even if they pass the triggers, may not pass the standard analysis since the photons that pass the triggers would not reconstruct the Higgs mass.
 - However, these events may be save to “tape” and signal may be seen with a reanalysis.

Thank You

Non-Isolated Photon Signals



Higgs Mixing Constraints



- Higgs dark photon coupling from scalar mixing:

$$|\lambda_{h\gamma_D\gamma_D}^{\text{max}}| \sim |\sin \theta^{\text{max}}| \frac{m_{\gamma_D}^2}{v_D} \approx 0.1 \frac{m_{\gamma_D}^2}{v_D}$$