Precise SM measurements as BSM probes: a new purpose for the W mass measurement

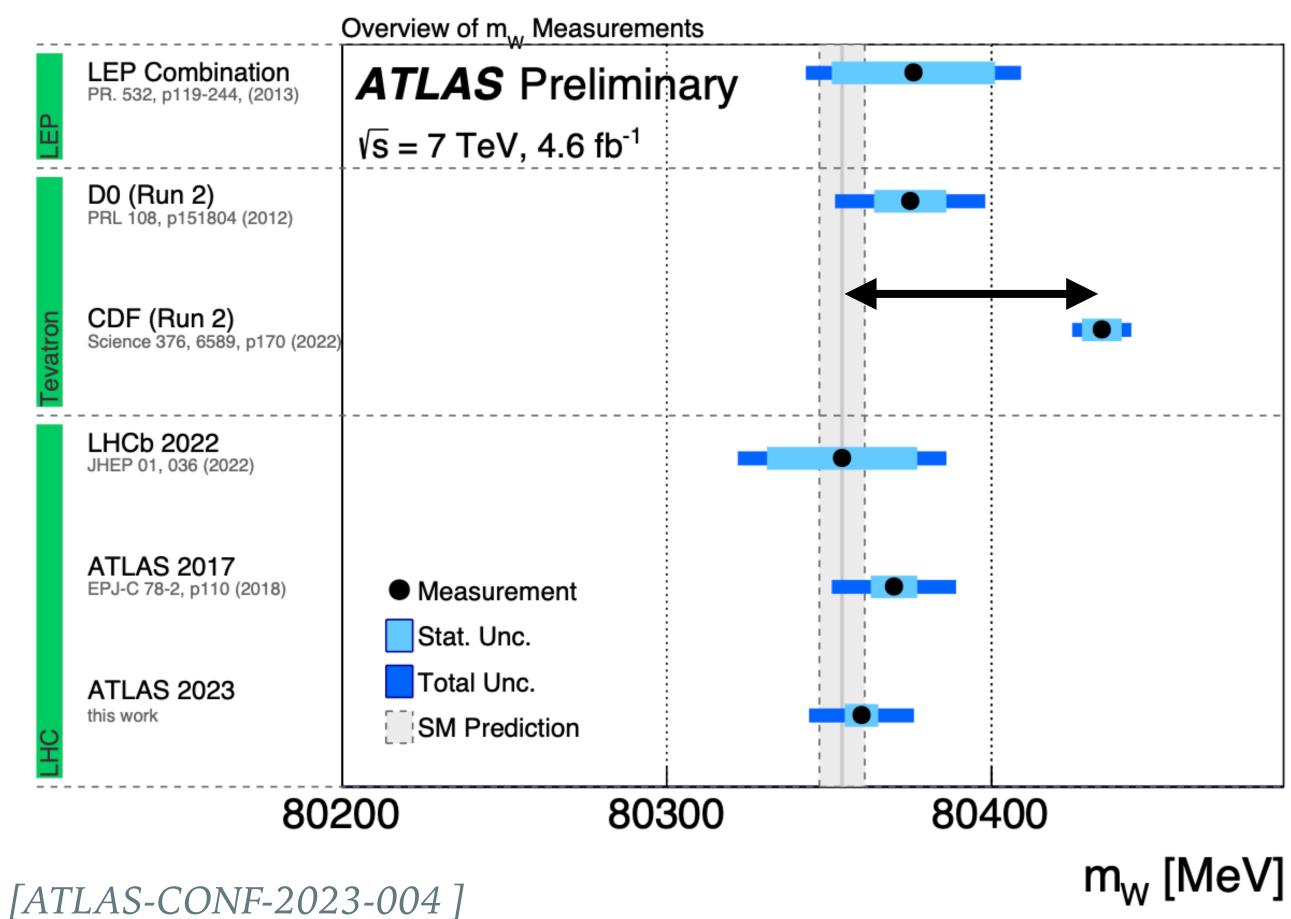
Lorenzo Ricci



Wednesday 17th May

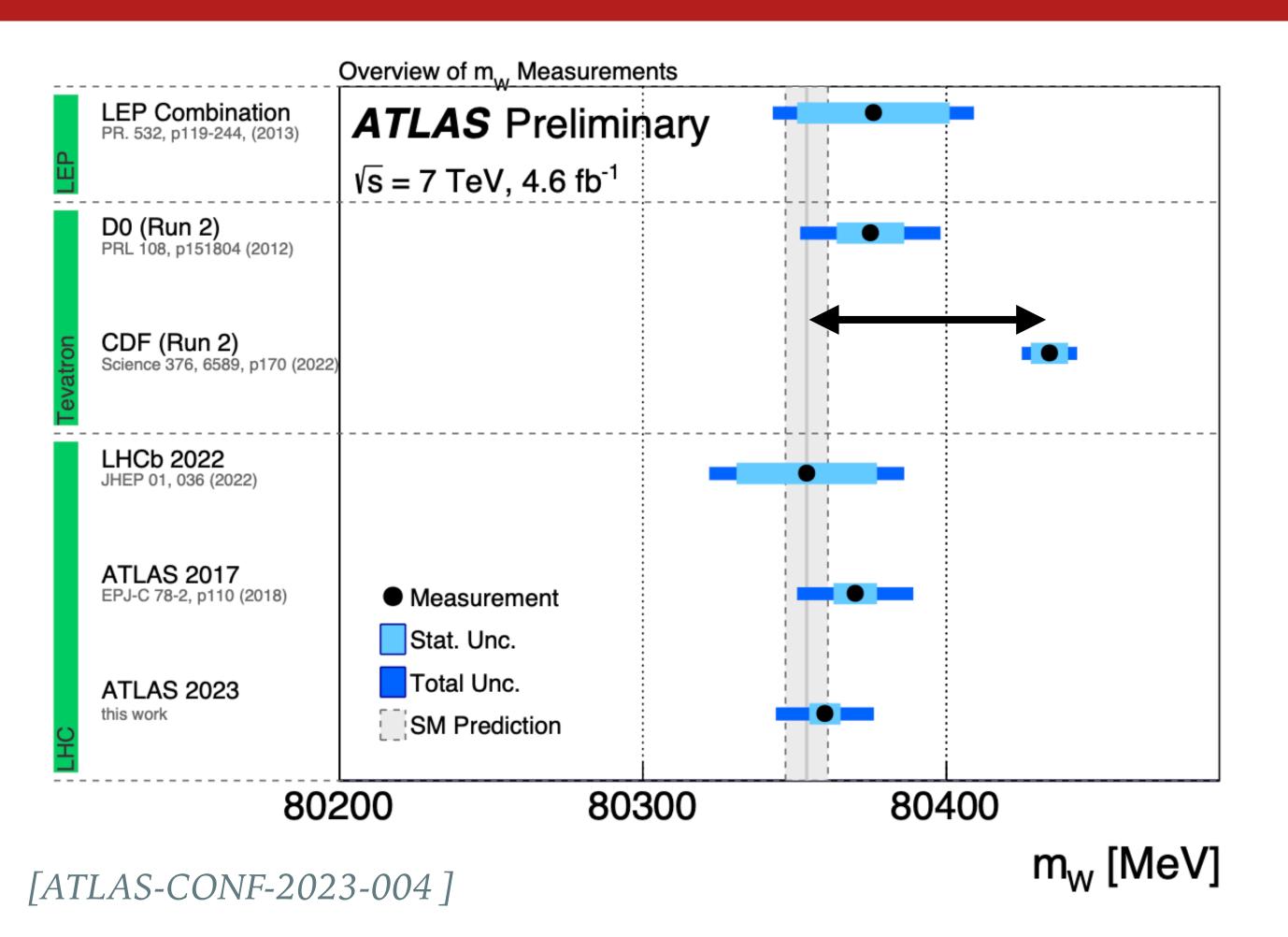
Mitchell 2023

Based on [K. Agashe, S. Airen, R. Franceschini, D. Kim, LR, D. Sathyan (in preparation)]



ATLAS-CONF-2023-004 J

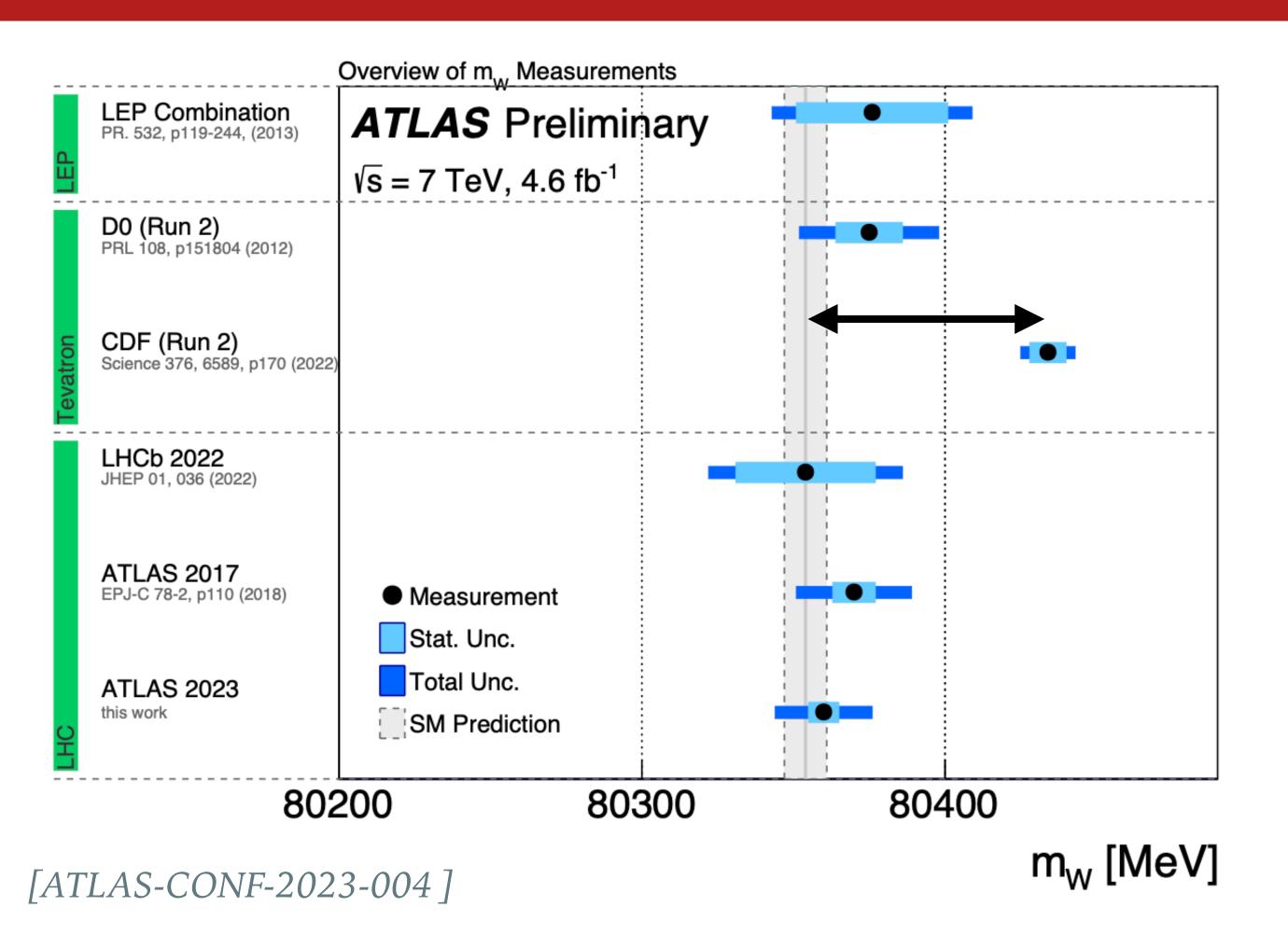
$$m_W^{SM}(G_F, m_Z, \alpha_{em}, \ldots) \neq m_W^{meas}$$



Can BSM explain the W mass anomaly?

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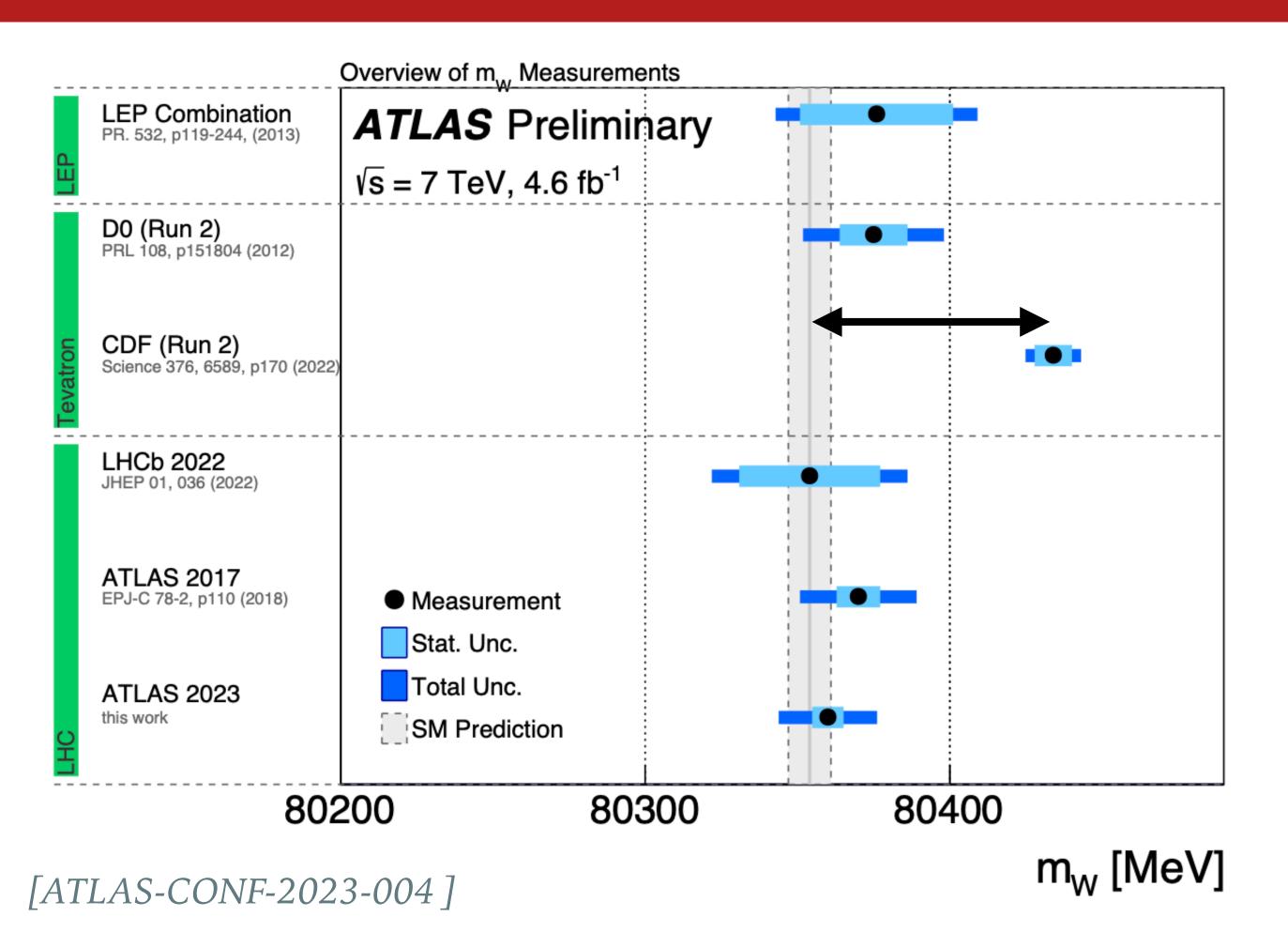


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[2204.04191, 2204.04514, 2204.04834, 2204.05285, 2204.05283, 2204.05085, 2204.05302, 2204.05260, 2204.05284, 2204.05296, 2204.05267, 2204.06327, 2204.05975, 2204.05992, 2204.05942, 2204.05760, 2204.05965, 2204.05728, 2204.06485,...]

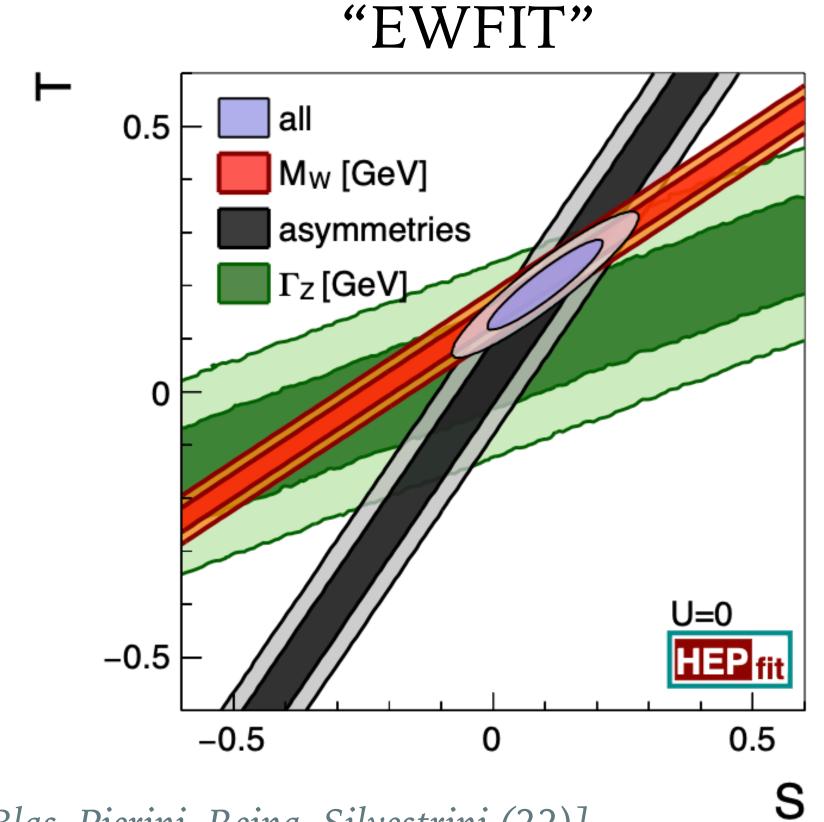
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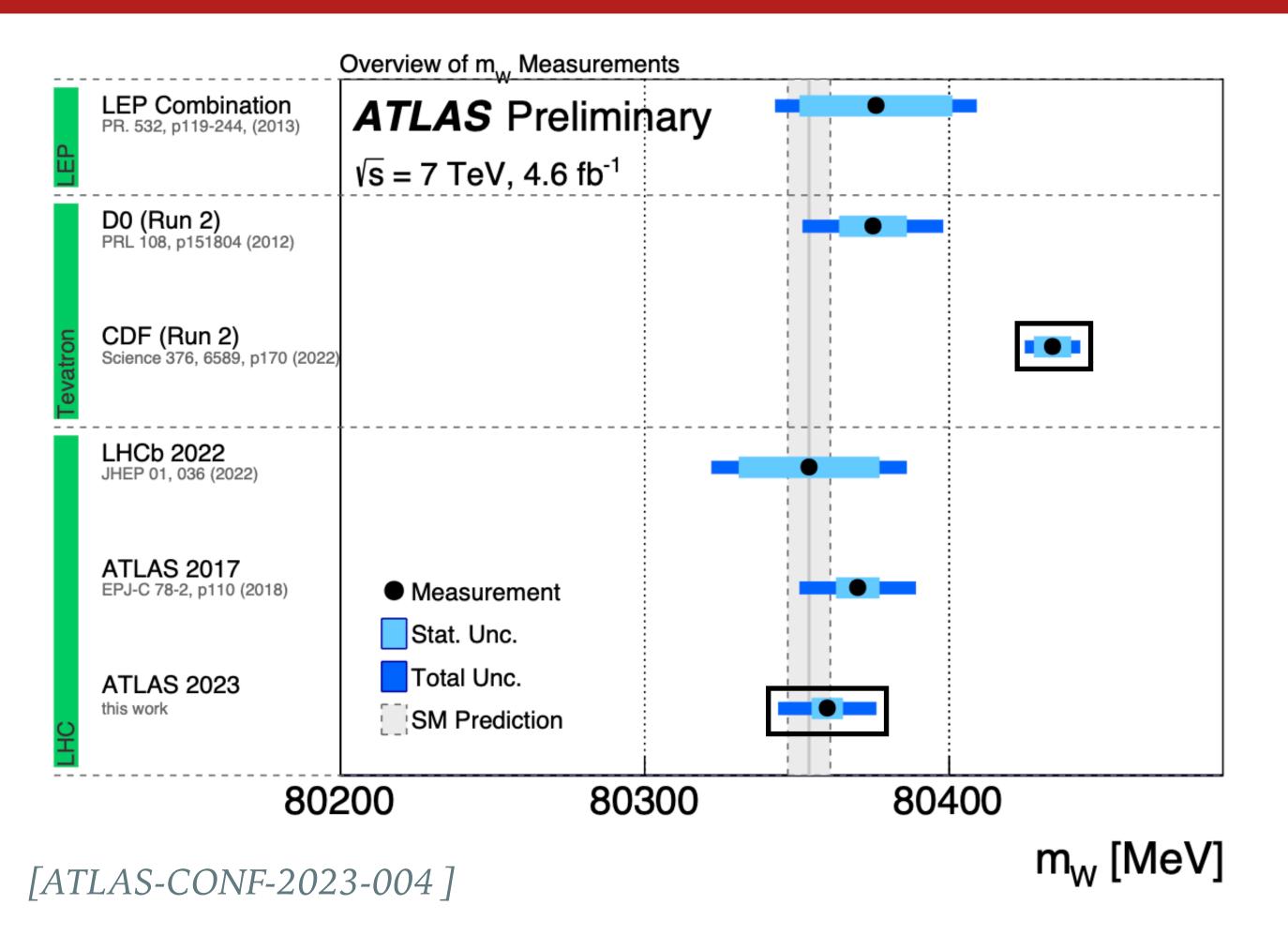


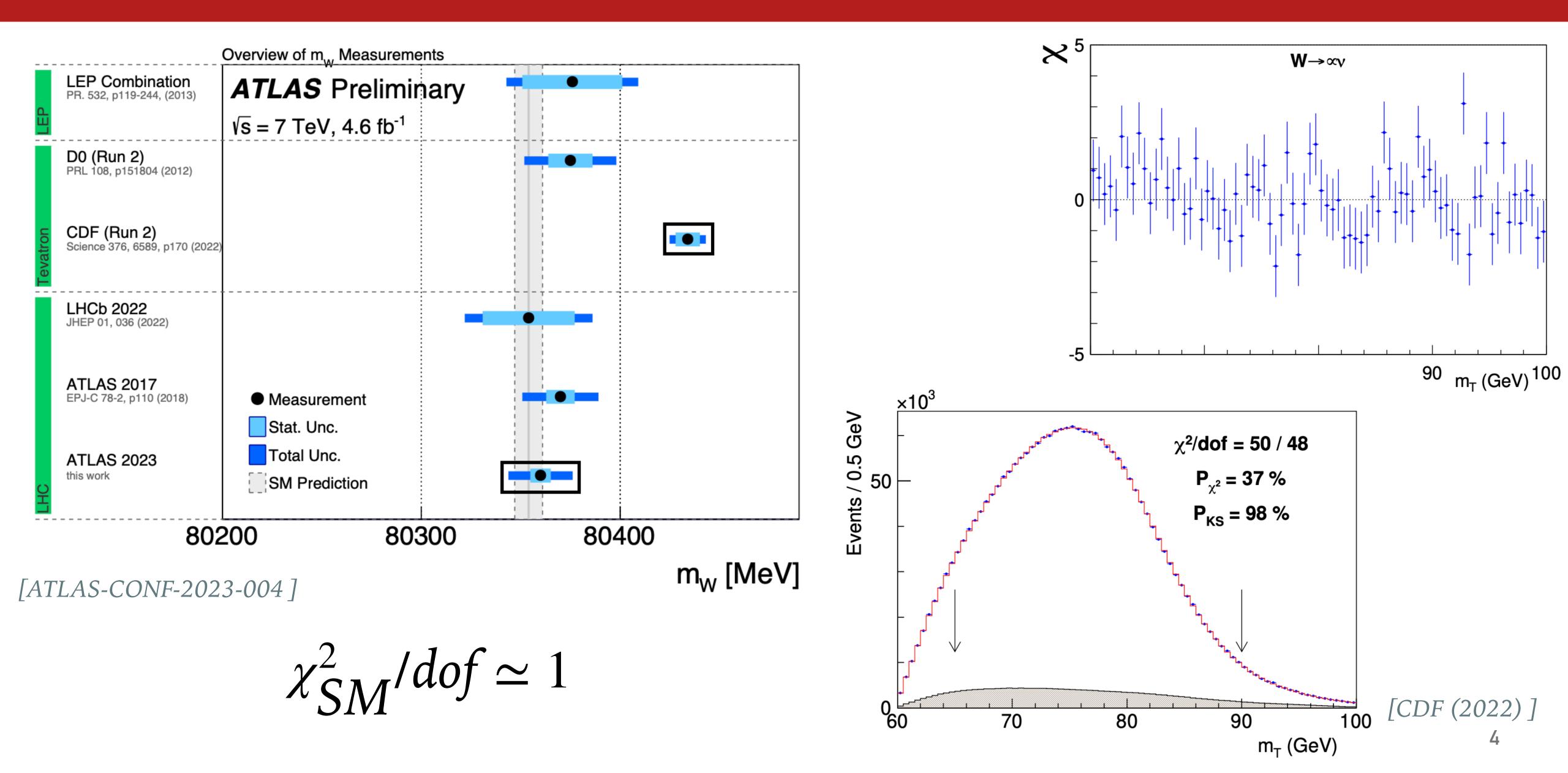
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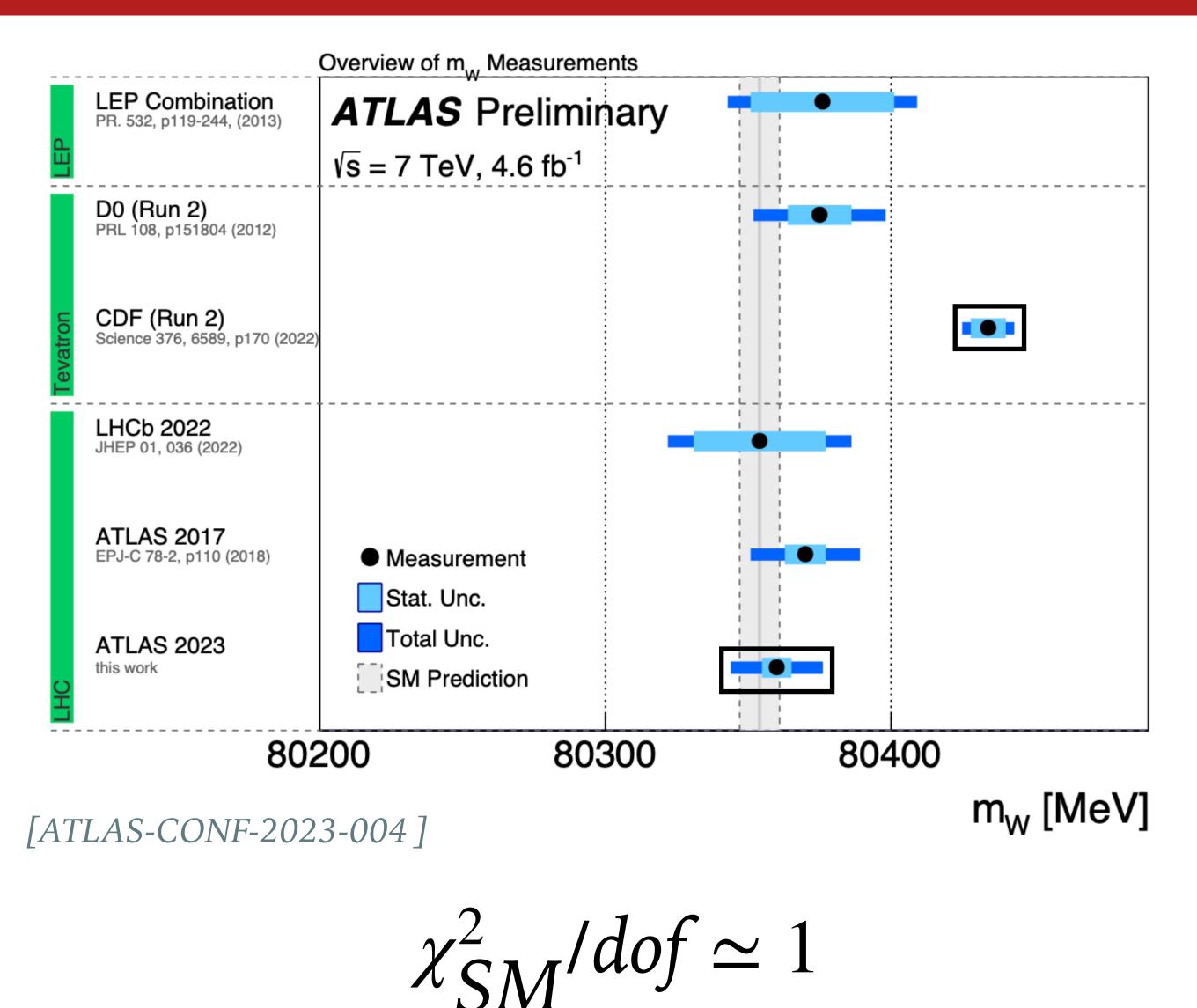
Can BSM explain the W mass anomaly?

$$m_W^{BSM}(G_F, m_Z, \alpha_{em}, S, T, \ldots) = m_W^{meas}$$



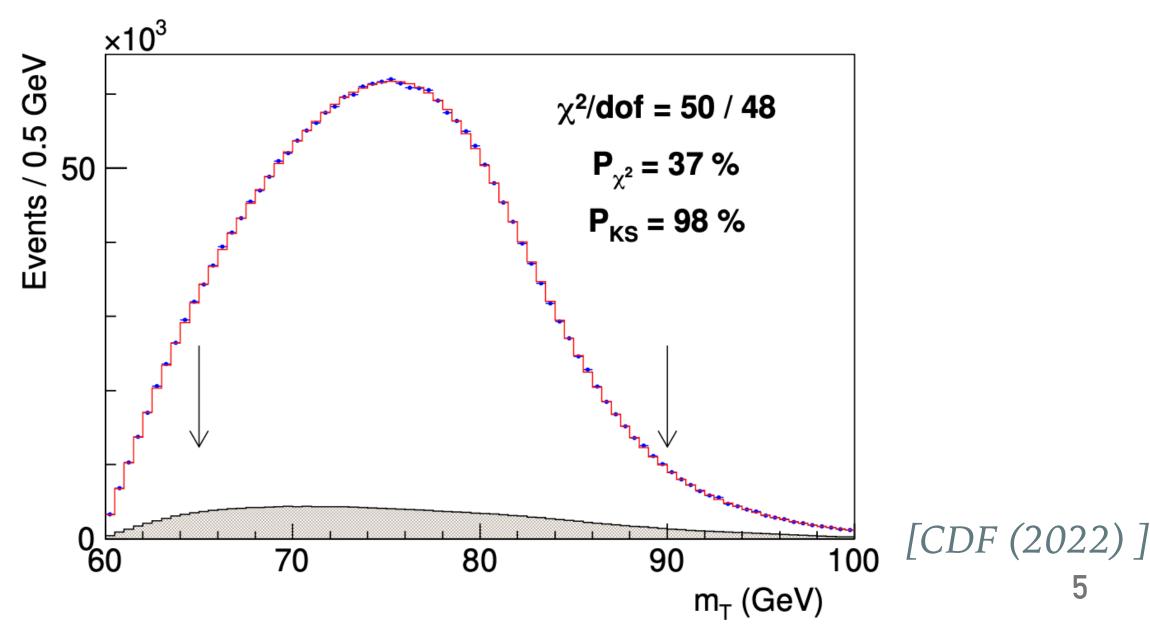


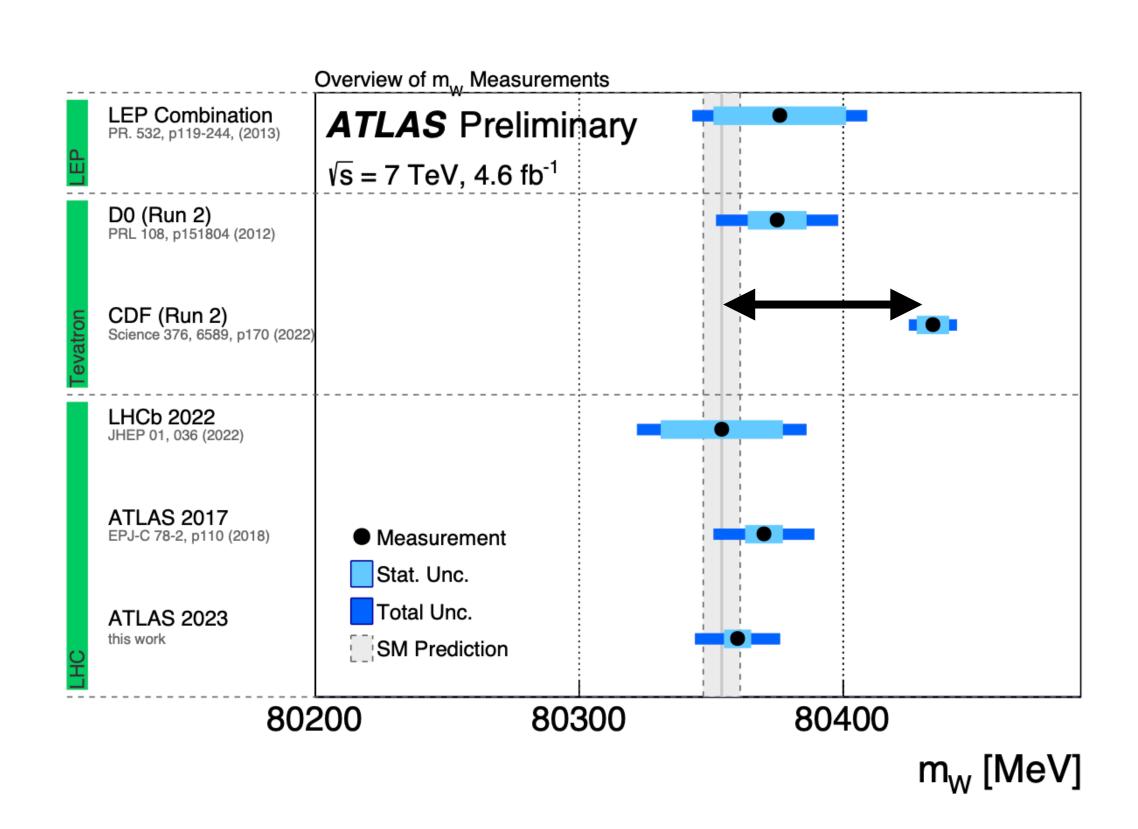




Can the W mass measurement *alone* constrain BSM?

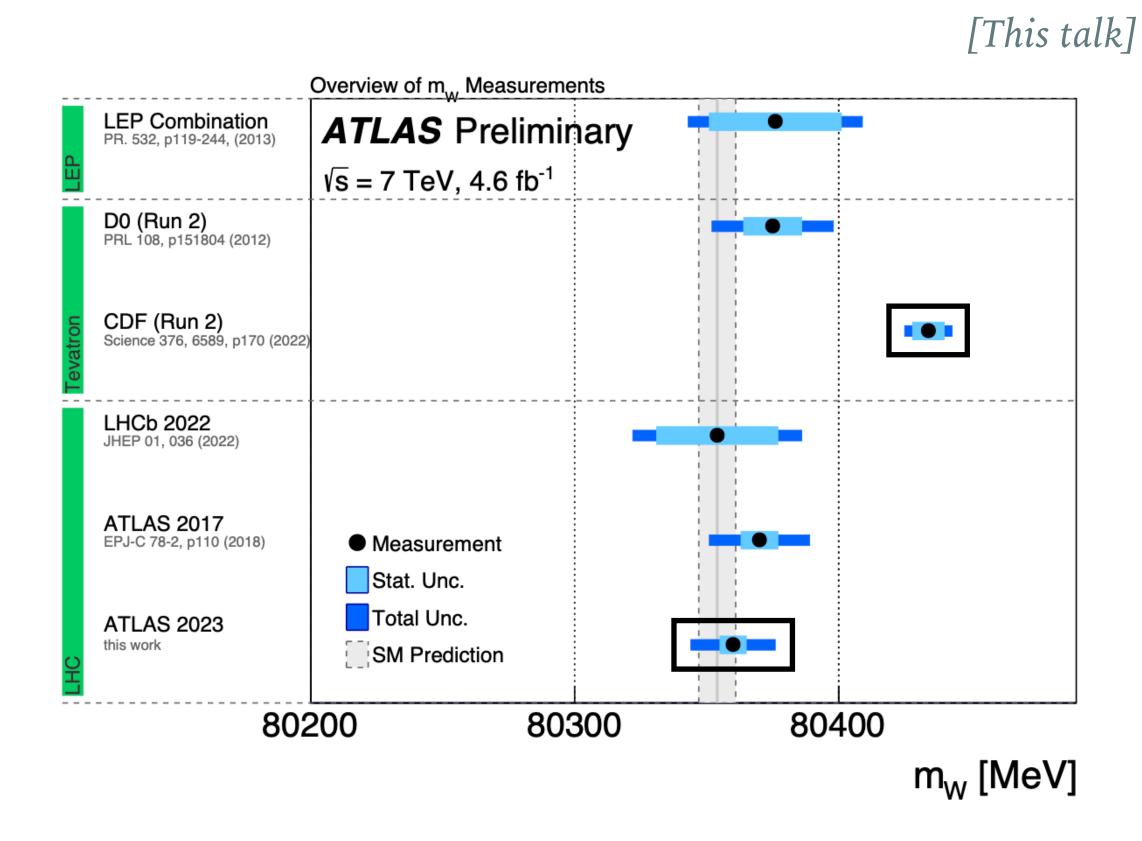
Independently from the EW fit and from any anomaly





Can BSM explain the W mass anomaly?

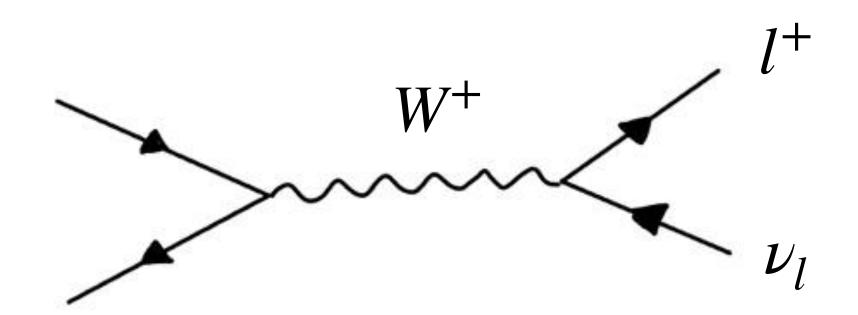
Relies on m_W prediction and measured value



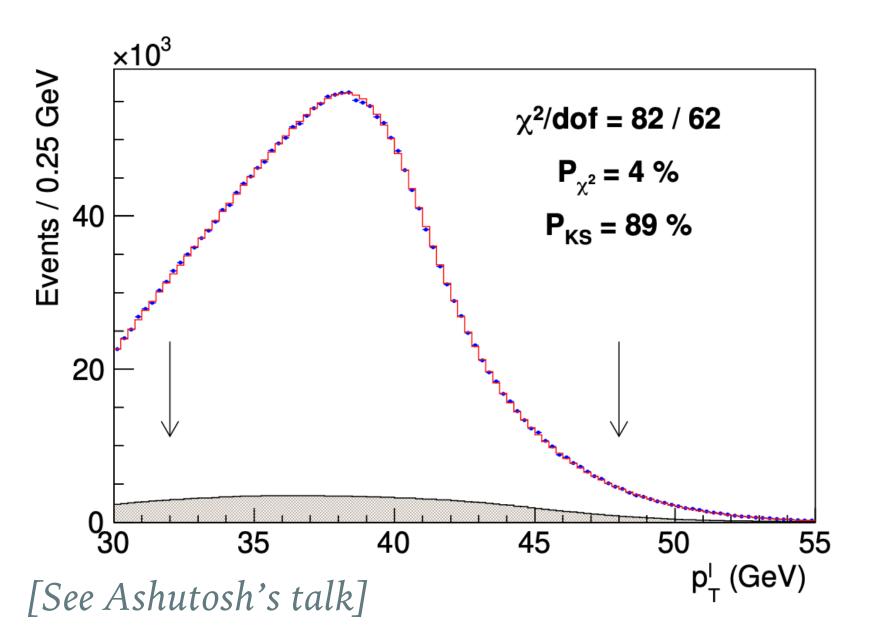
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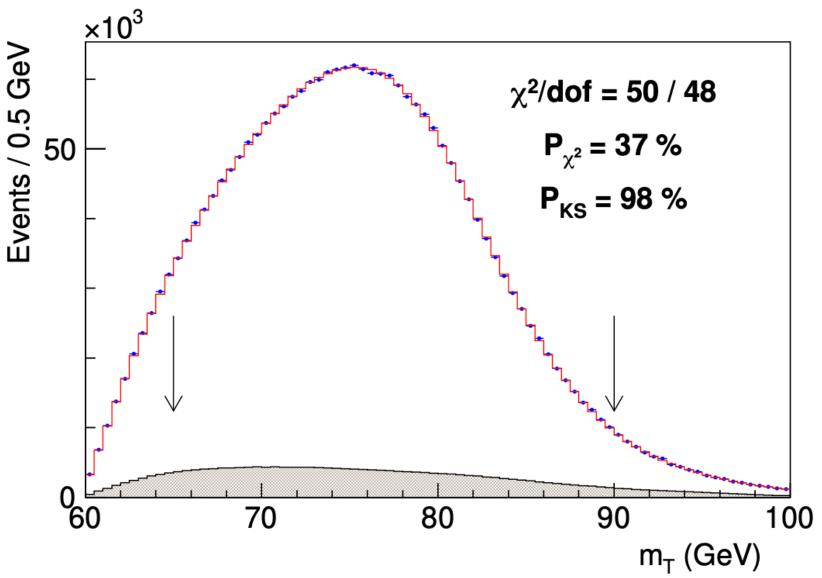
Relies only on the m_W measurement,

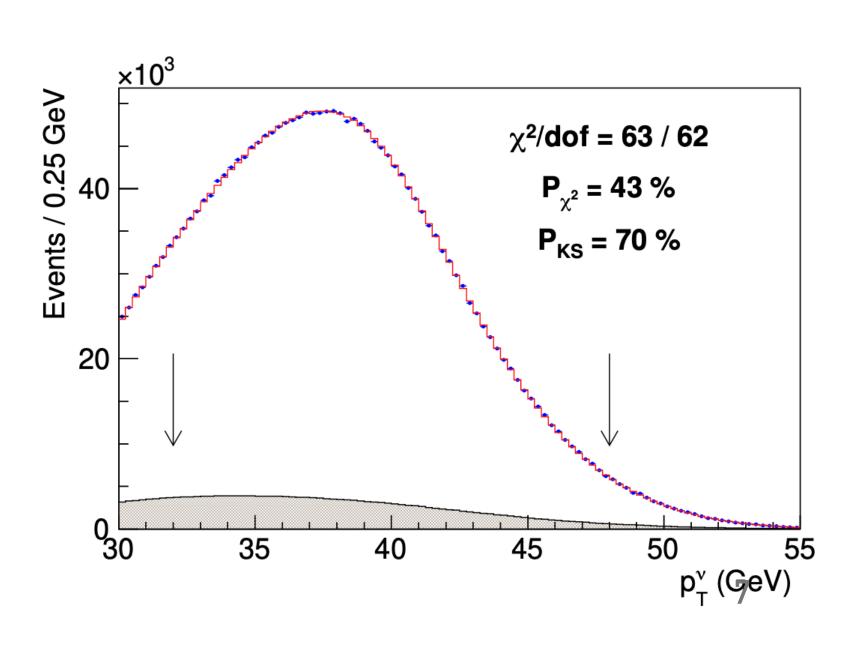
New physics in the W measurement



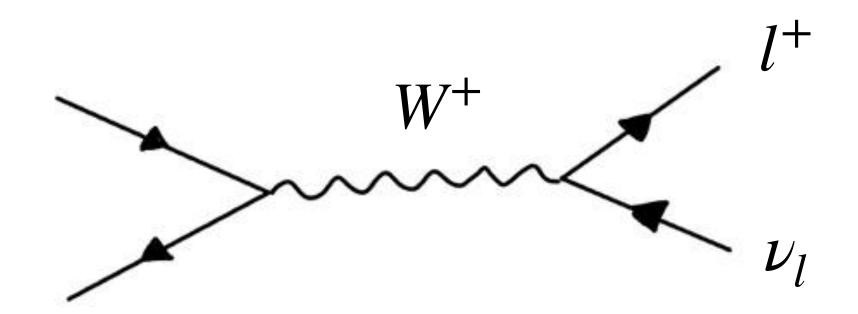
 m_W is extracted from the three kinematic distributions

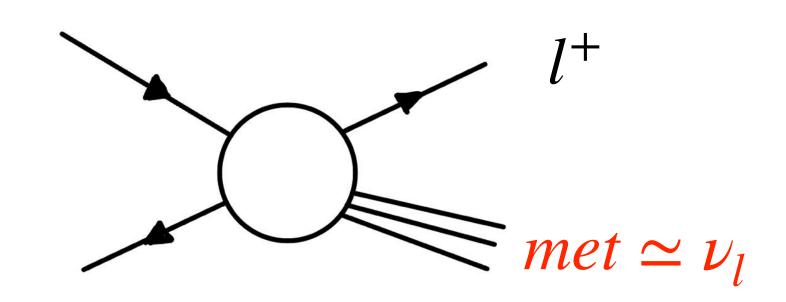






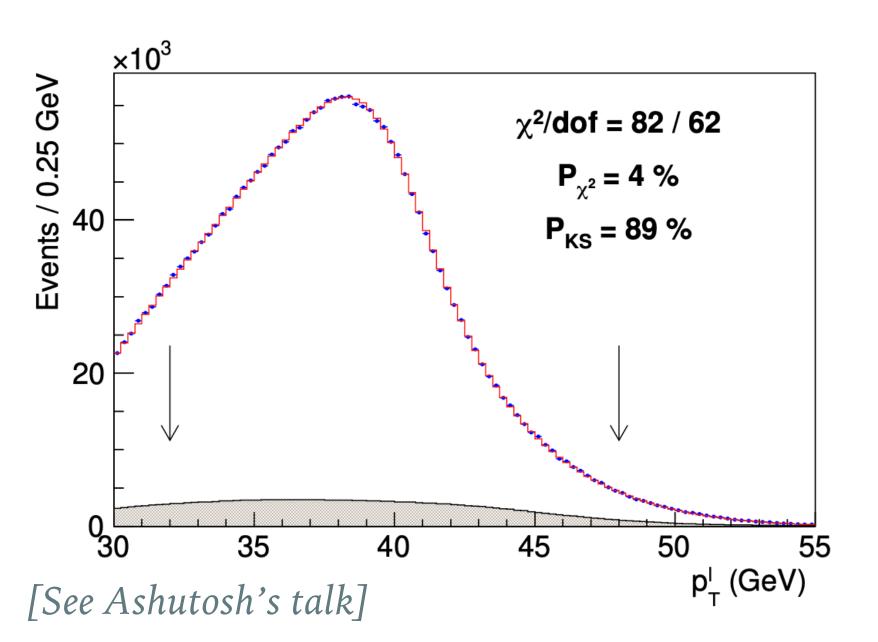
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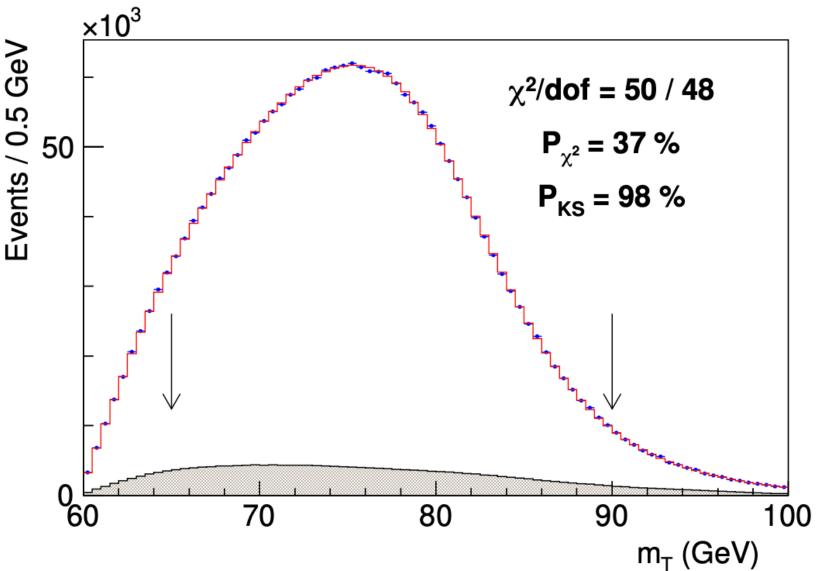


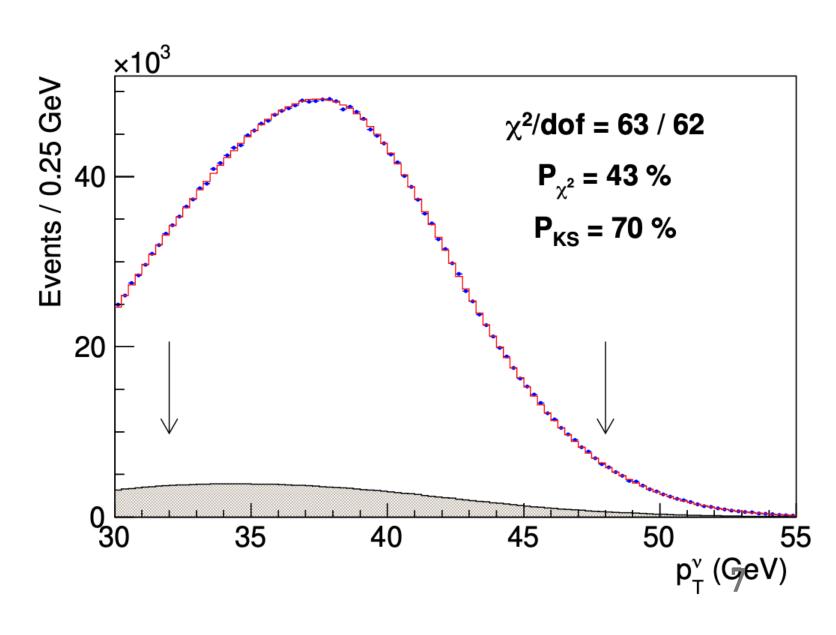


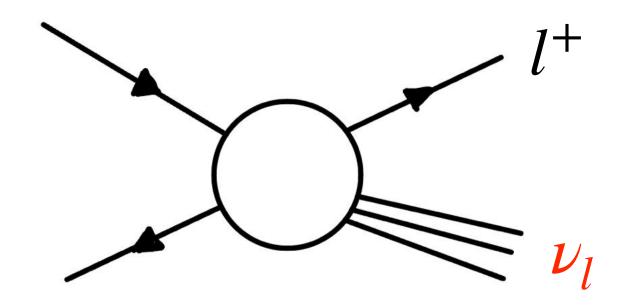
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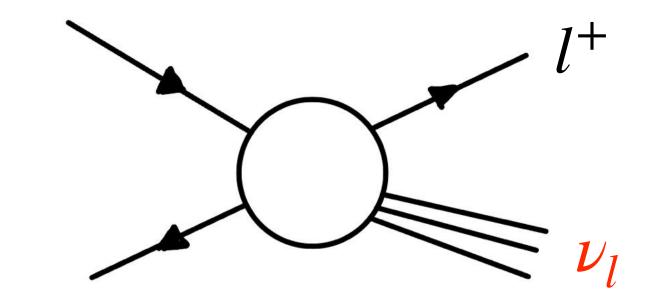
Any l + met process can contaminate the W sample and distort the distributions



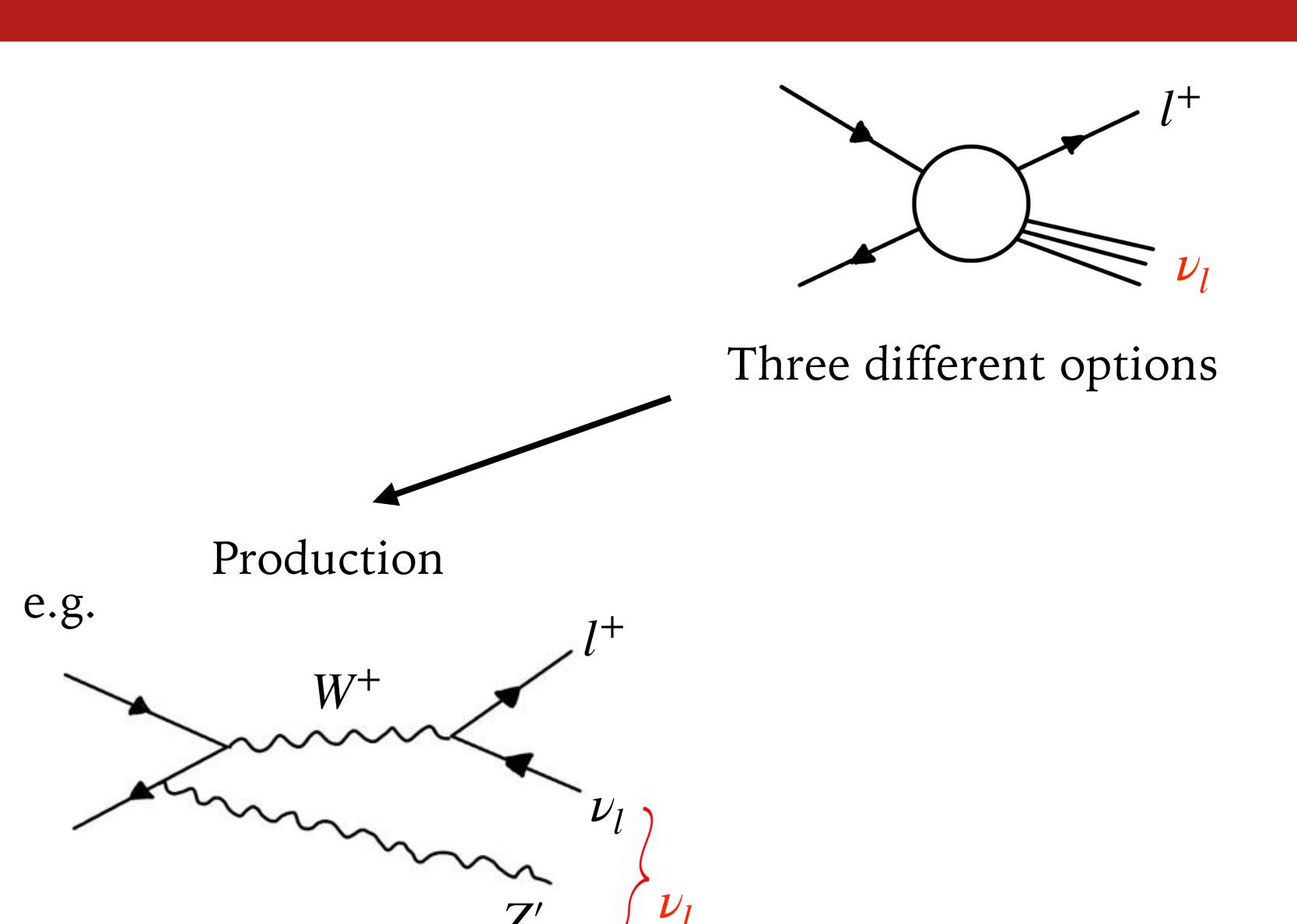


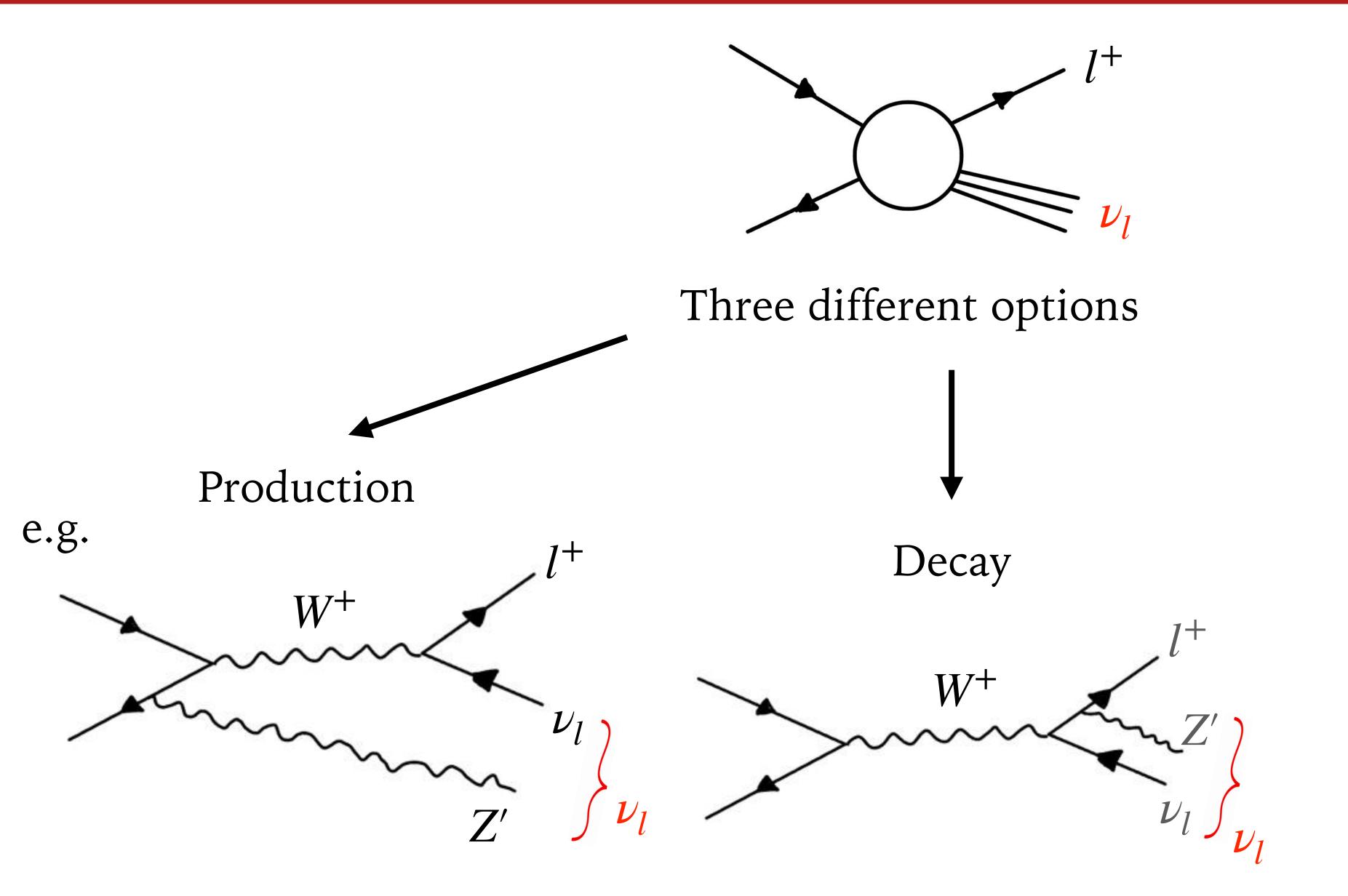


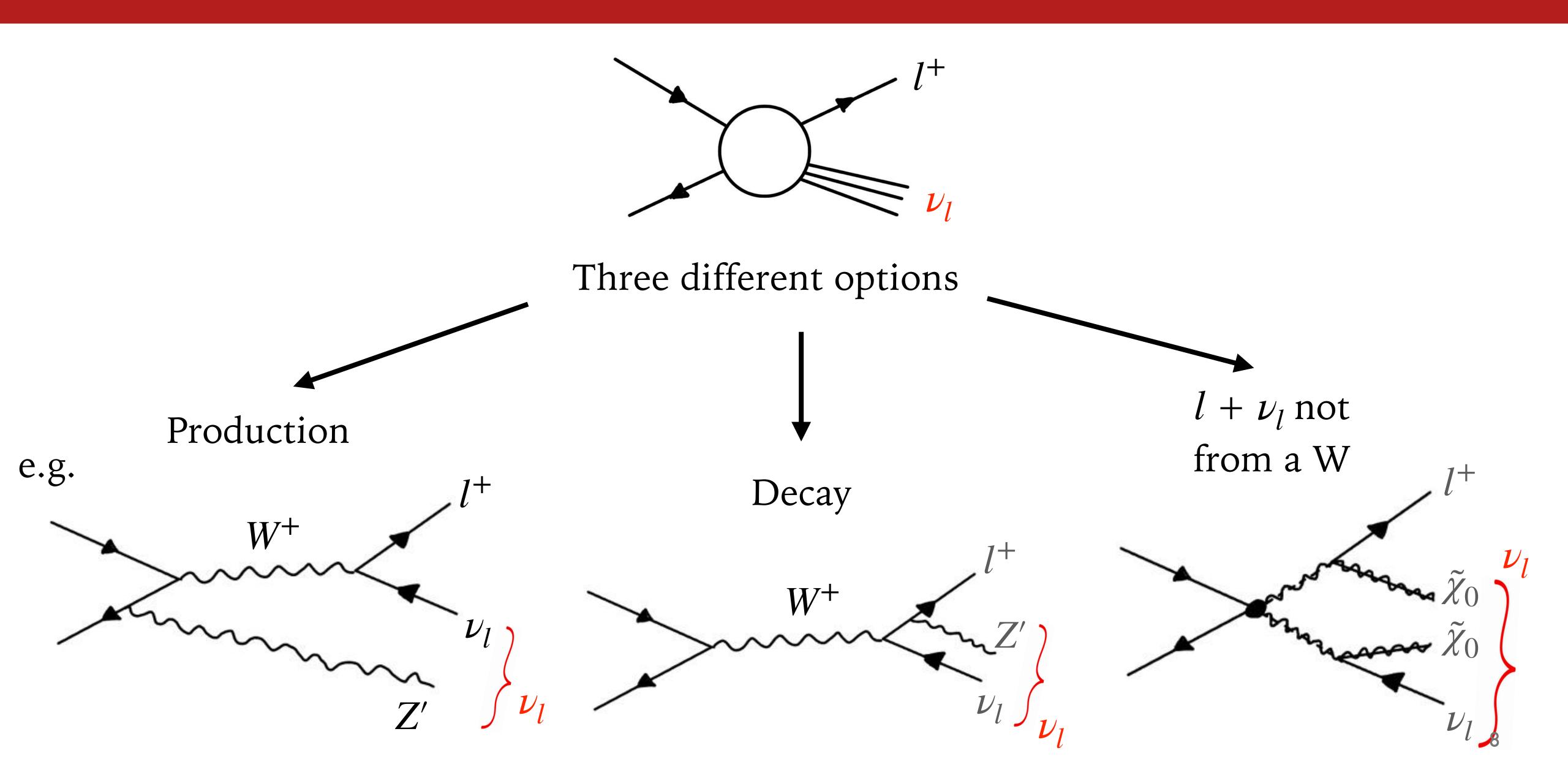


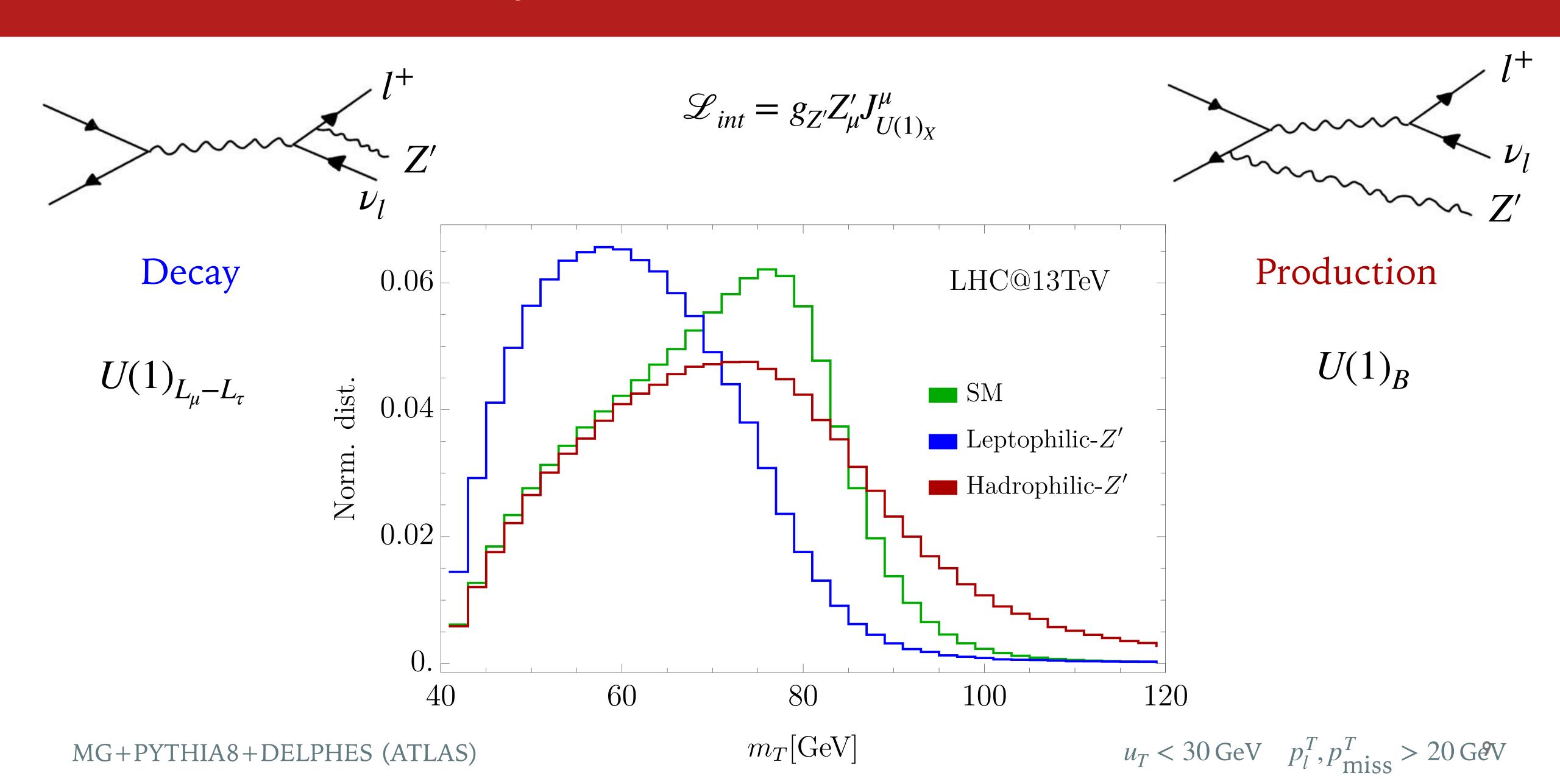


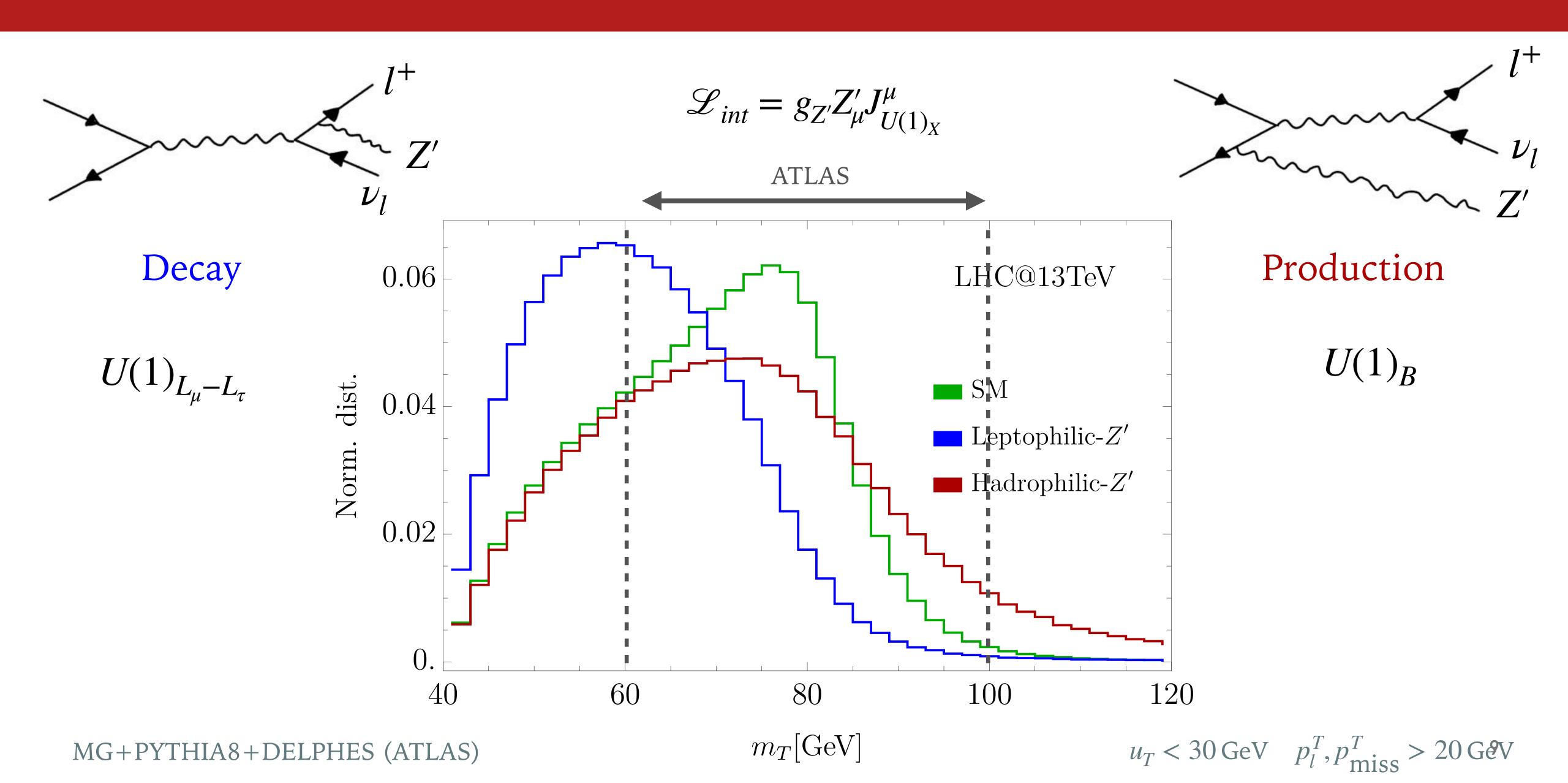
Three different options



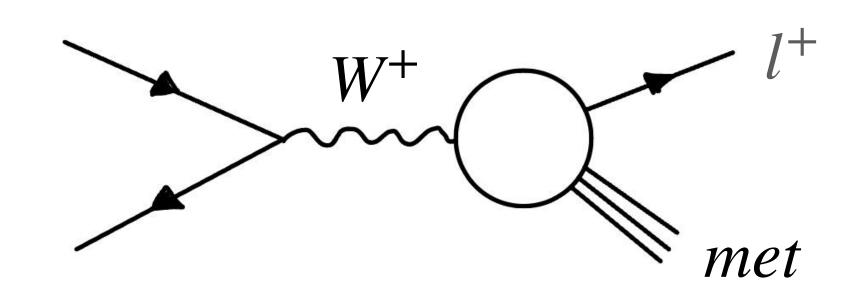


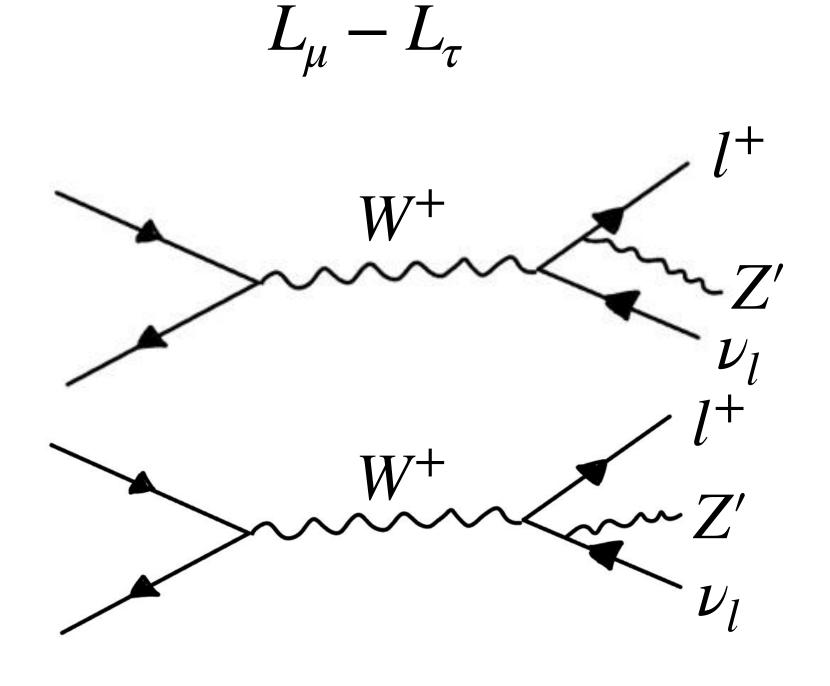




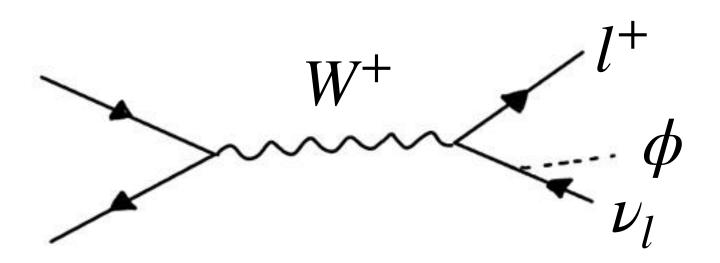


New physics in the W decay



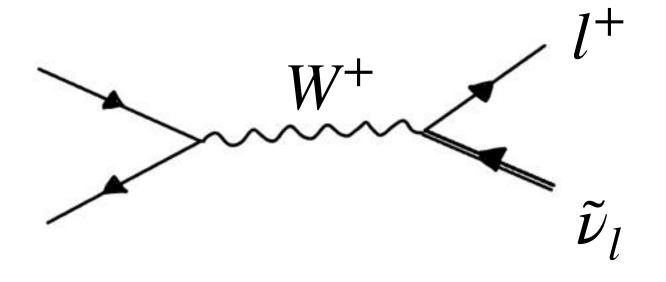


"Neutrinophilic" scalar



[1802.00009, 1910.0113]

Dirac neutrino portal

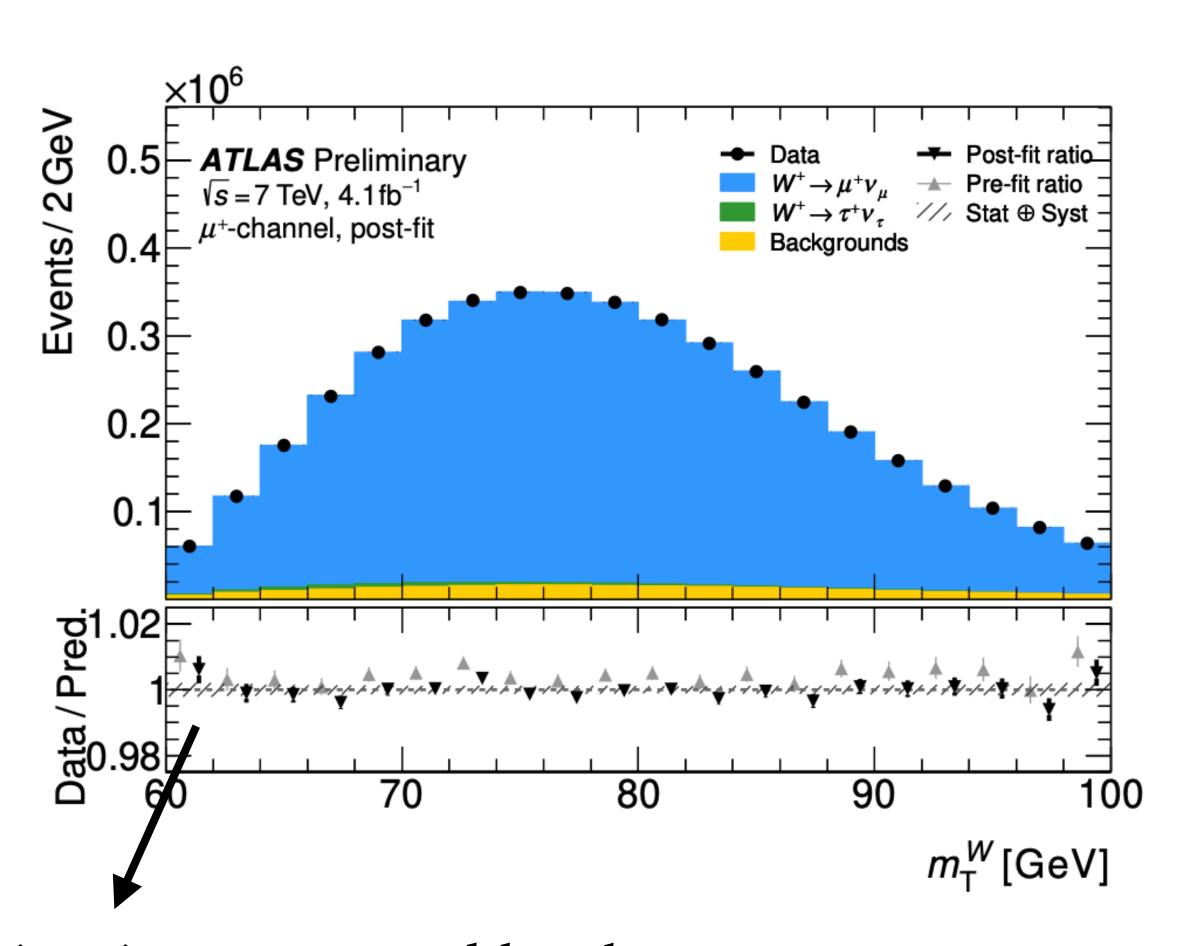


New physics in the W decay: $U(1)_{L_u-L_\tau} Z'$

$$\mathcal{L} = -\frac{1}{4}F^{'2} + \frac{1}{2}m_Z^{'2}Z^{'2} + g_{Z'}Z'_{\mu}J^{\mu}_{L_{\mu}-L_{\tau}}$$

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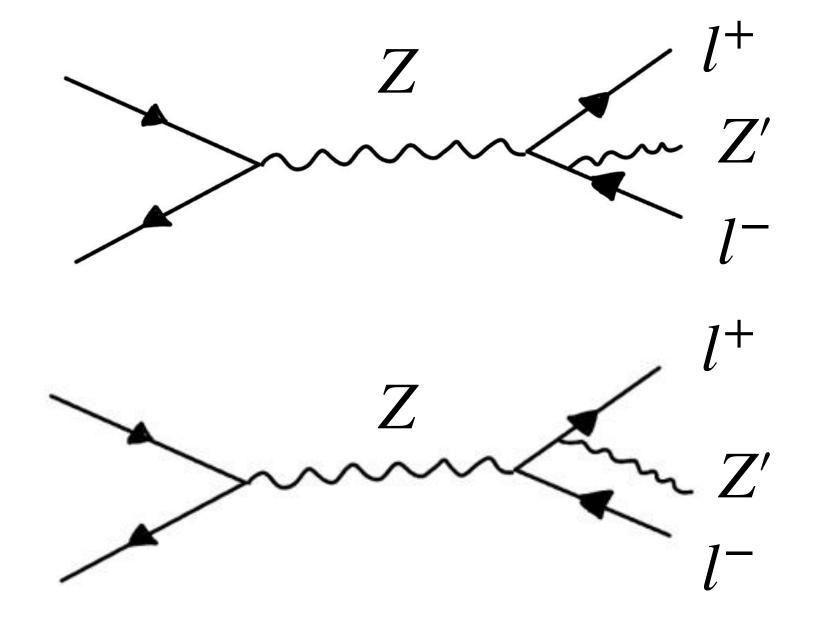


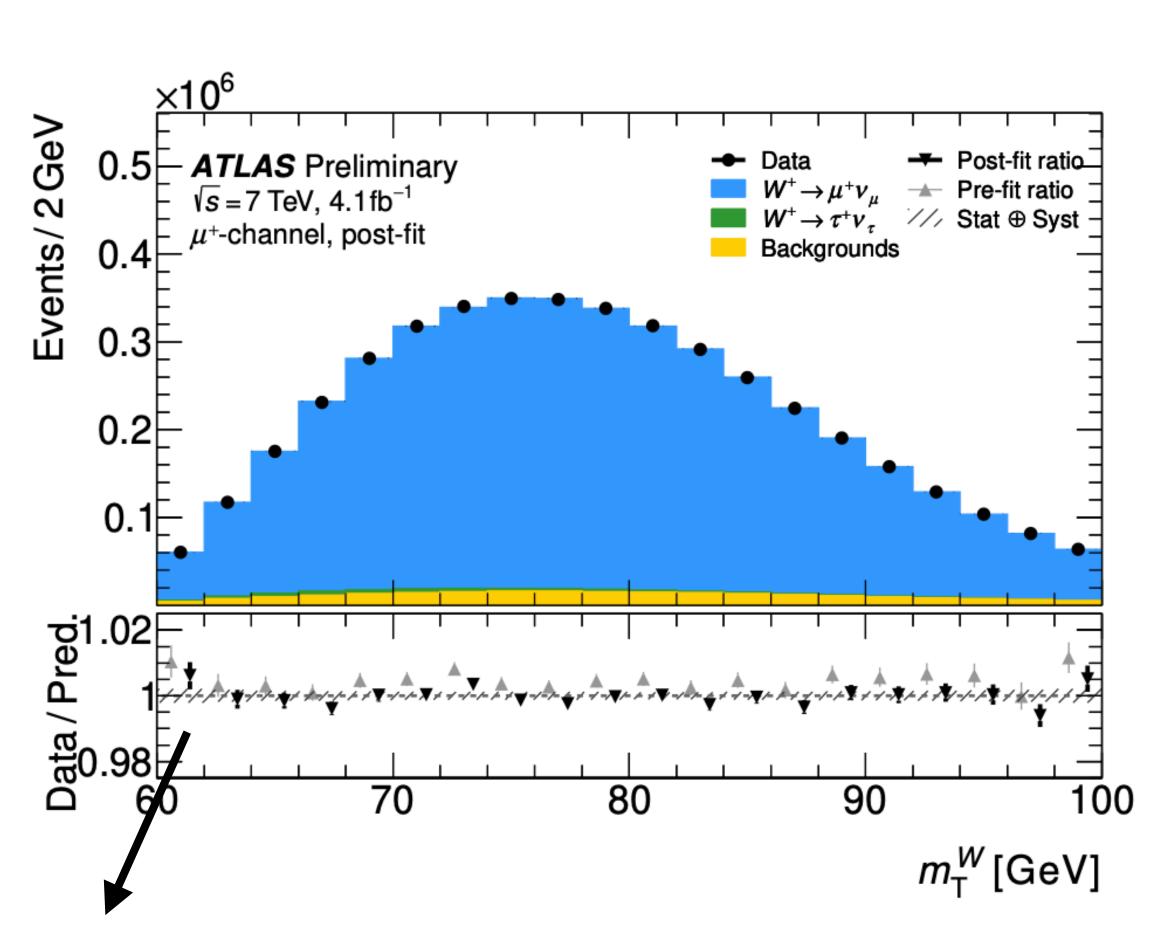
This precision is guaranteed by the data driven approach (Z calibration)

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It also modifies the Z decay!



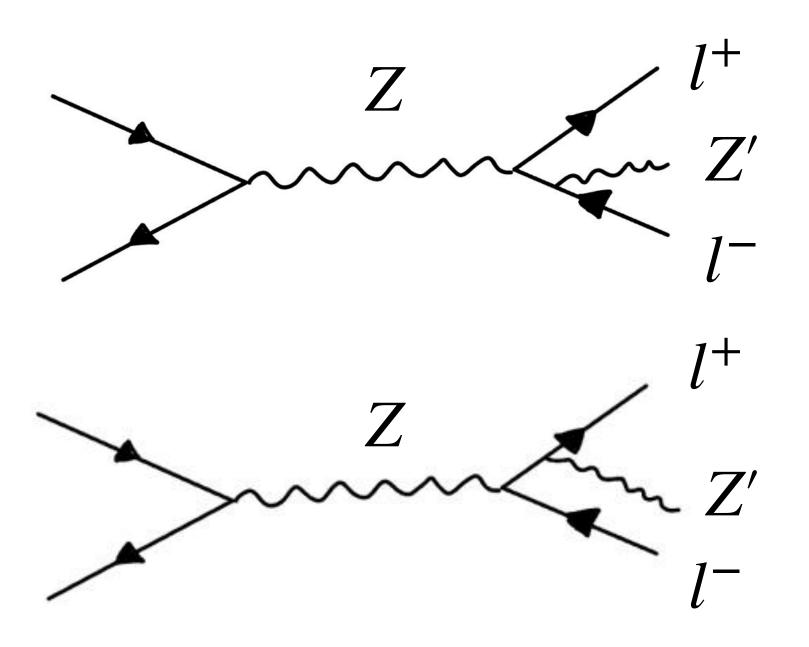


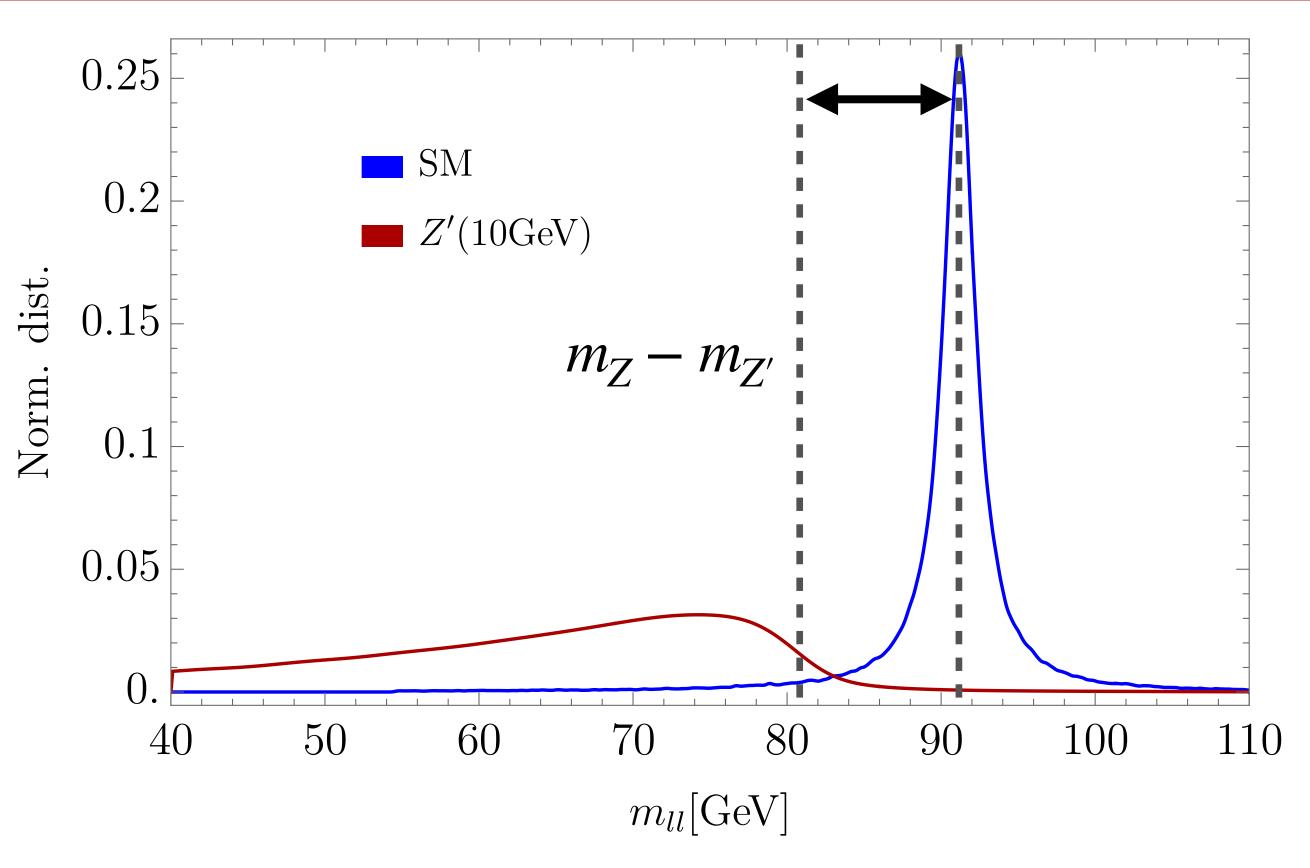
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New physics in the W decay: $U(1)_{L_u-L_ au}$ Z'

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It also modifies the Z decay!





We can veto on these events as far as $m_{Z'} > \Gamma_Z$

$$80 < m_{ll}[GeV] < 100$$
 ATLAS $66 < m_{ll}[GeV] < 116$ CDF

New physics in the W decay: $U(1)_{L_{\mu}-L_{ au}} Z'$

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- The $\chi^2 \equiv \chi^2(m_W^{data}; m_W, g_{Z'}, m_{Z'})$ depends on three parameters $(m_W, g_{Z'}, m_{Z'})$
- •In principle m_W must float
- •In practice, it is very hard for BSM to modify the best fit value of m_W

New physics in the W decay: $U(1)_{L_u-L_ au} \, Z'$

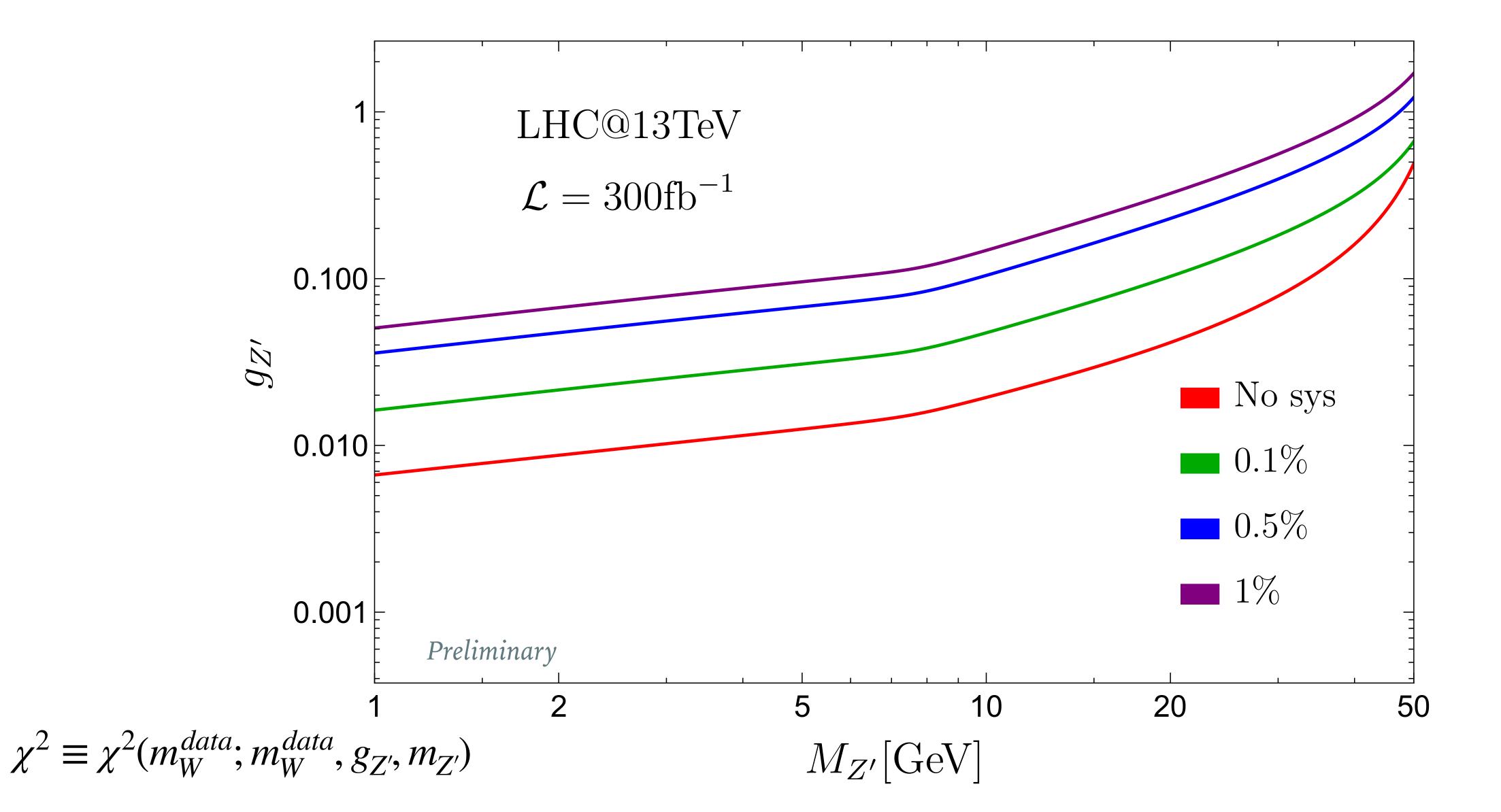
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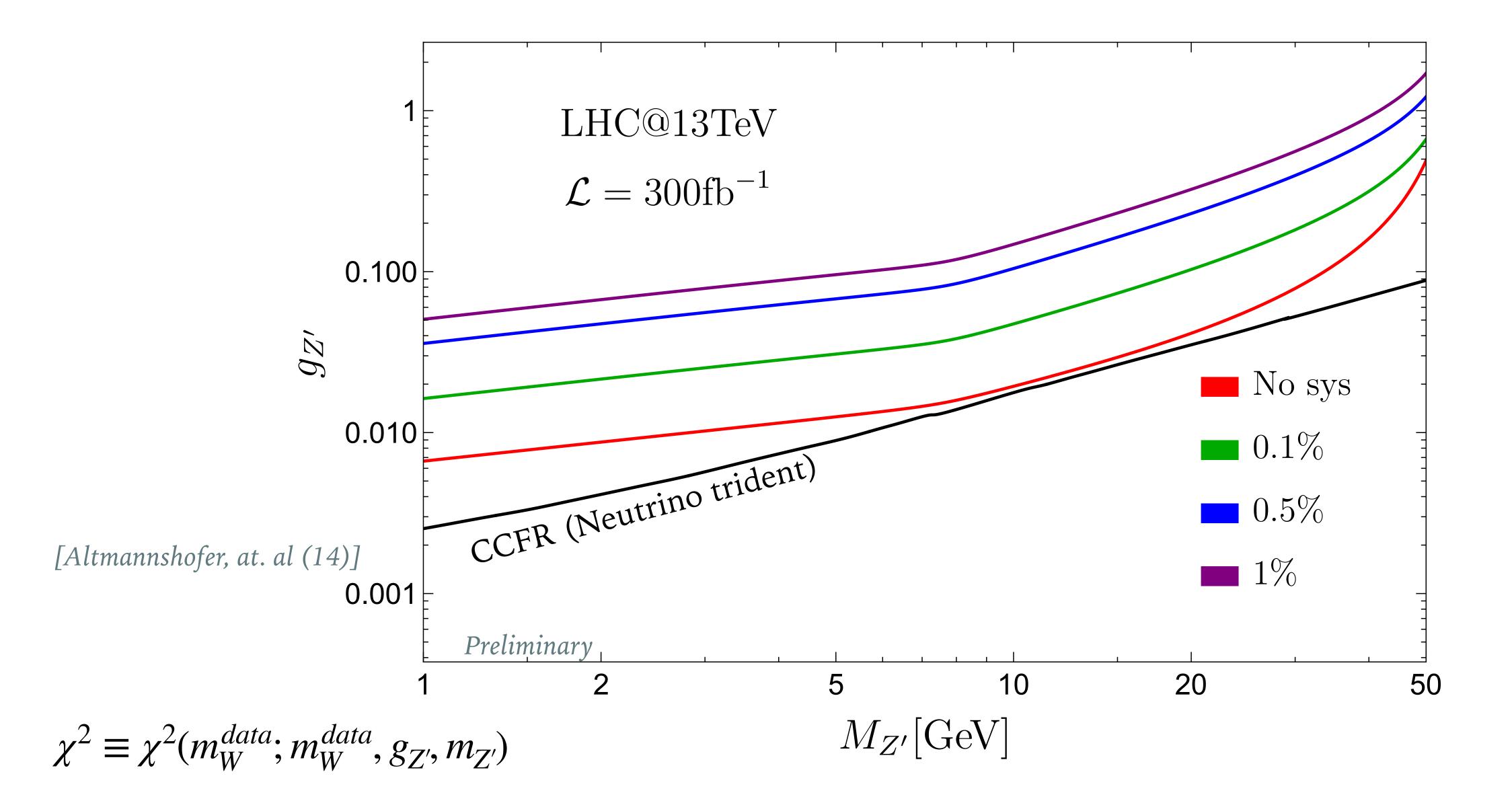
OUR ANALYSIS ----

- •MG+PYTHIA8+DELPHES(ATLAS)
- • $u_T < 30 \,\text{GeV}$ $p_l^T, p_{\text{miss}}^T > 20 \,\text{GeV}$
- $\bullet \chi^2 = \chi_{p_l}^2 + \chi_{m_T}^2$
- •Bins size: $p^T \rightarrow 0.5 \text{ GeV}$, $m_T \rightarrow 1 \text{ GeV}$
- Number of events normalised to the data

New physics in the W decay: $U(1)_{L_{\mu}-L_{\tau}}$ Z'



New physics in the W decay: $U(1)_{L_{\mu}-L_{\tau}}$ Z'



New physics in the W decay: Neutrinophilic scalar

$$\mathcal{L}_{int} = \frac{\tilde{\lambda}}{\Lambda^2} (LH)(LH)\phi$$

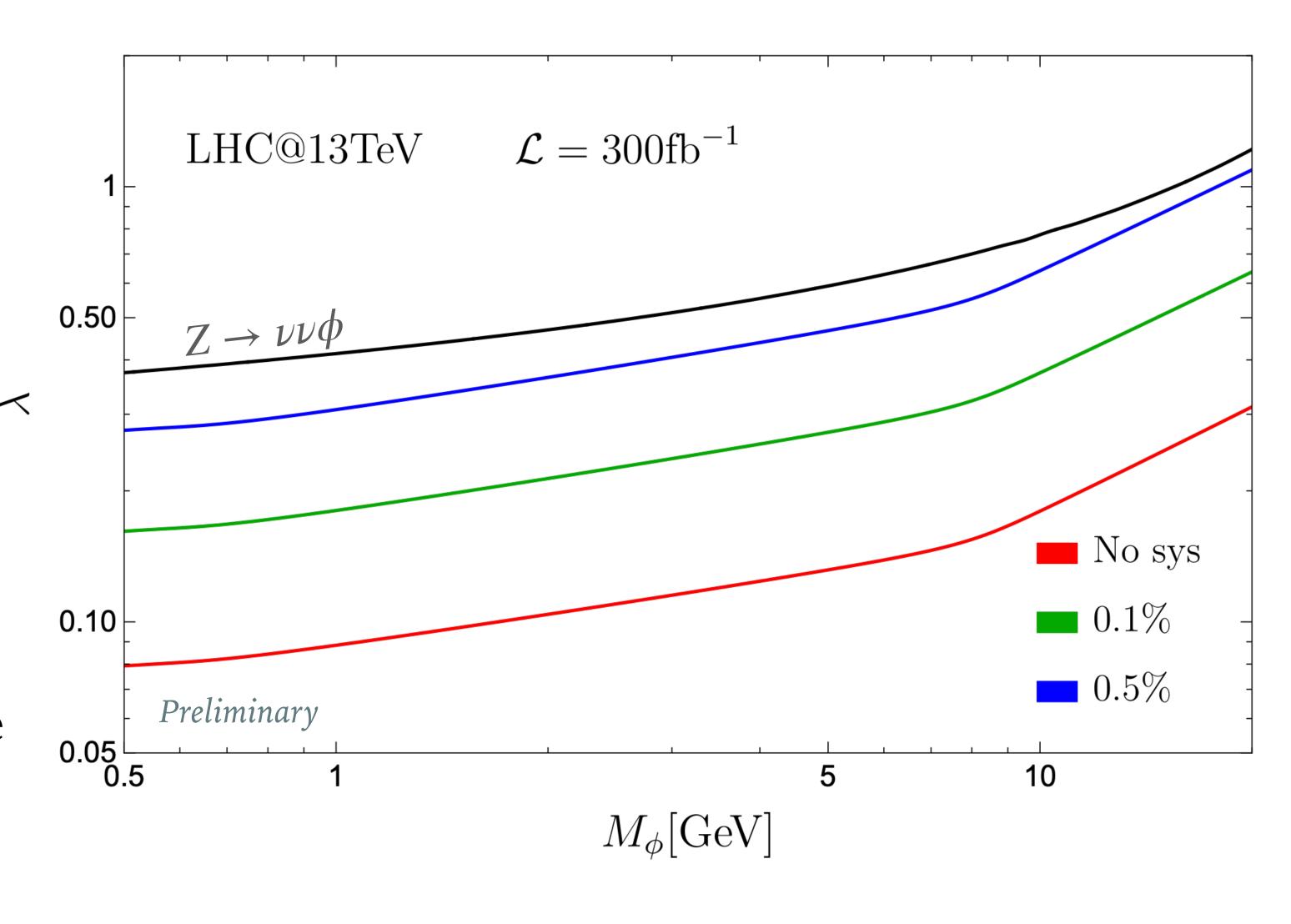
$$\downarrow < H >$$

$$\mathcal{L}_{int} = \frac{1}{2} \lambda \nu \nu \phi$$

$$l^+$$

$$\nu_l$$

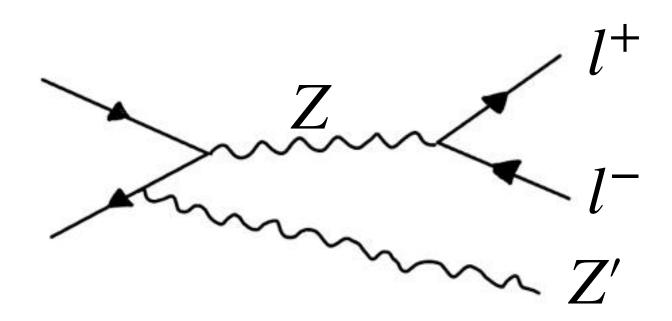
- No new effects in the Z sample
- We don't need to worry about calibration



$$\mathcal{L} = -\frac{1}{4}F^{2} + \frac{1}{2}m_{Z}^{2}Z^{2} + g_{Z'}Z'_{\mu}J^{\mu}_{B}$$

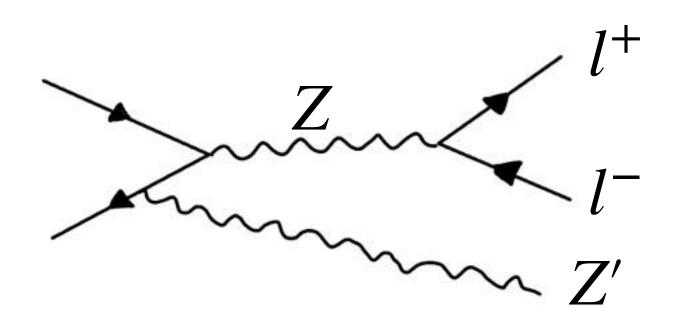
$$\mathcal{L} = -\frac{1}{4}F^{2} + \frac{1}{2}m_{Z}^{2}Z^{2} + g_{Z'}Z'_{\mu}J^{\mu}_{B}$$

• New physics can modify the Z sample (and so the calibration)



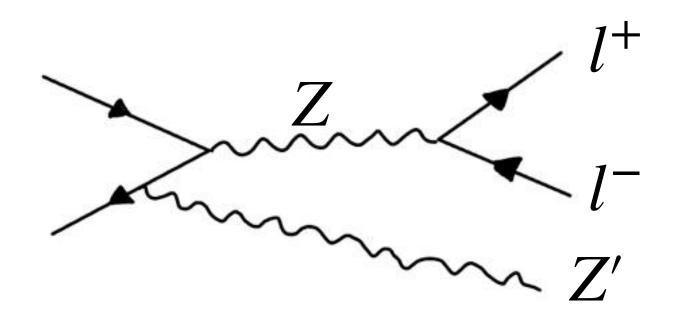
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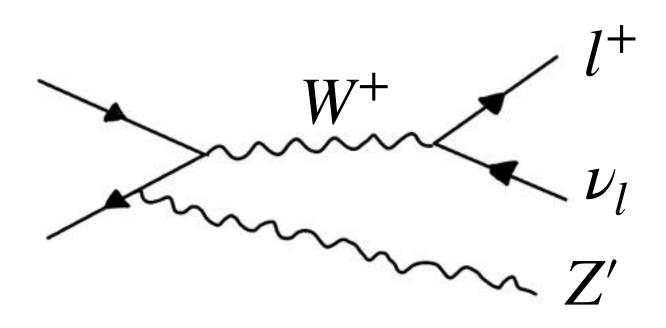
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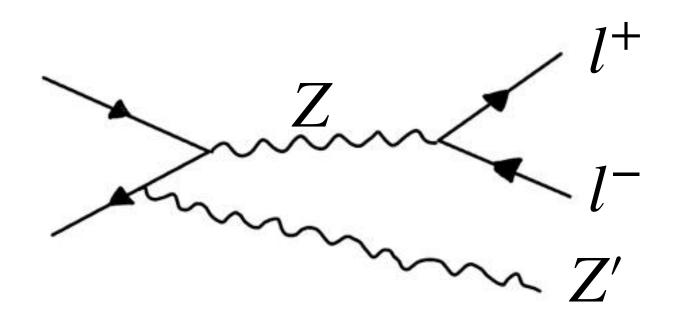


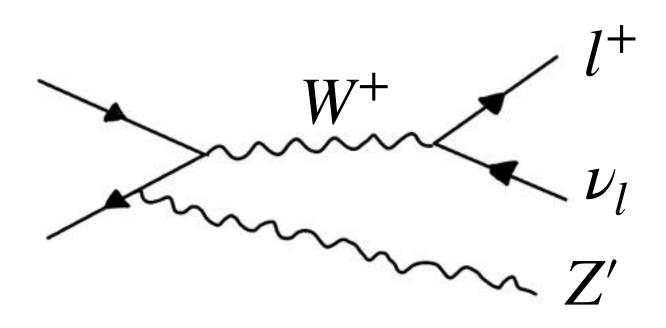


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$$p_{Z,l^+}^T \to p_{W,l^+}^T(pred) = p_{W,l^+}^T(meas)$$

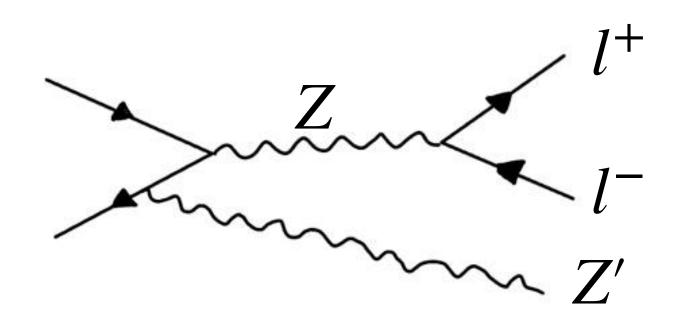


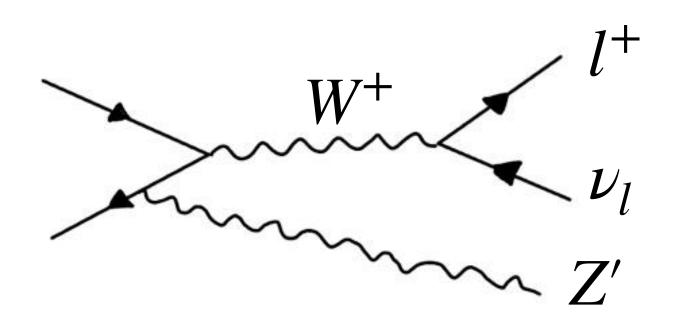


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$$\begin{aligned} p_{Z,l^{+}}^{T} &\to p_{W,l^{+}}^{T}(pred) = p_{W,l^{+}}^{T}(meas) \\ p_{Z,l^{-}}^{T} &\to p_{W,\nu_{l}}^{T}(pred) = p_{W,\nu_{l}+Z'}^{T}(meas) = p_{W,l^{+}}^{T}(meas) \end{aligned}$$

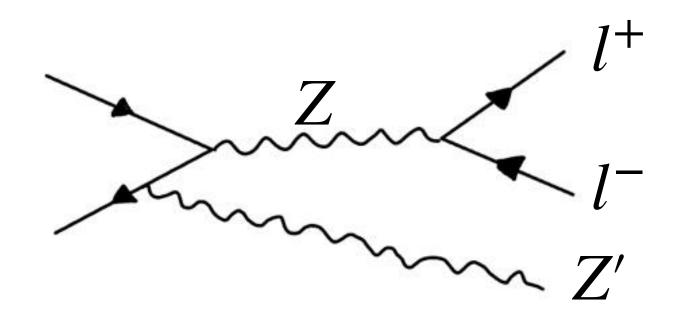


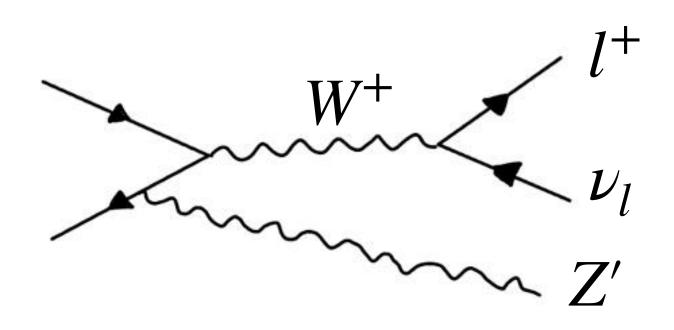


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$$\begin{split} p_{Z,l^+}^T &\to p_{W,l^+}^T(pred) = p_{W,l^+}^T(meas) \\ p_{Z,l^-}^T &\to p_{W,\nu_l}^T(pred) = p_{W,\nu_l+Z'}^T(meas) = p_{W,l^+}^T(meas) \\ m_T^Z &\to m_T^Z(pred) \neq m_T^Z(meas) \end{split}$$





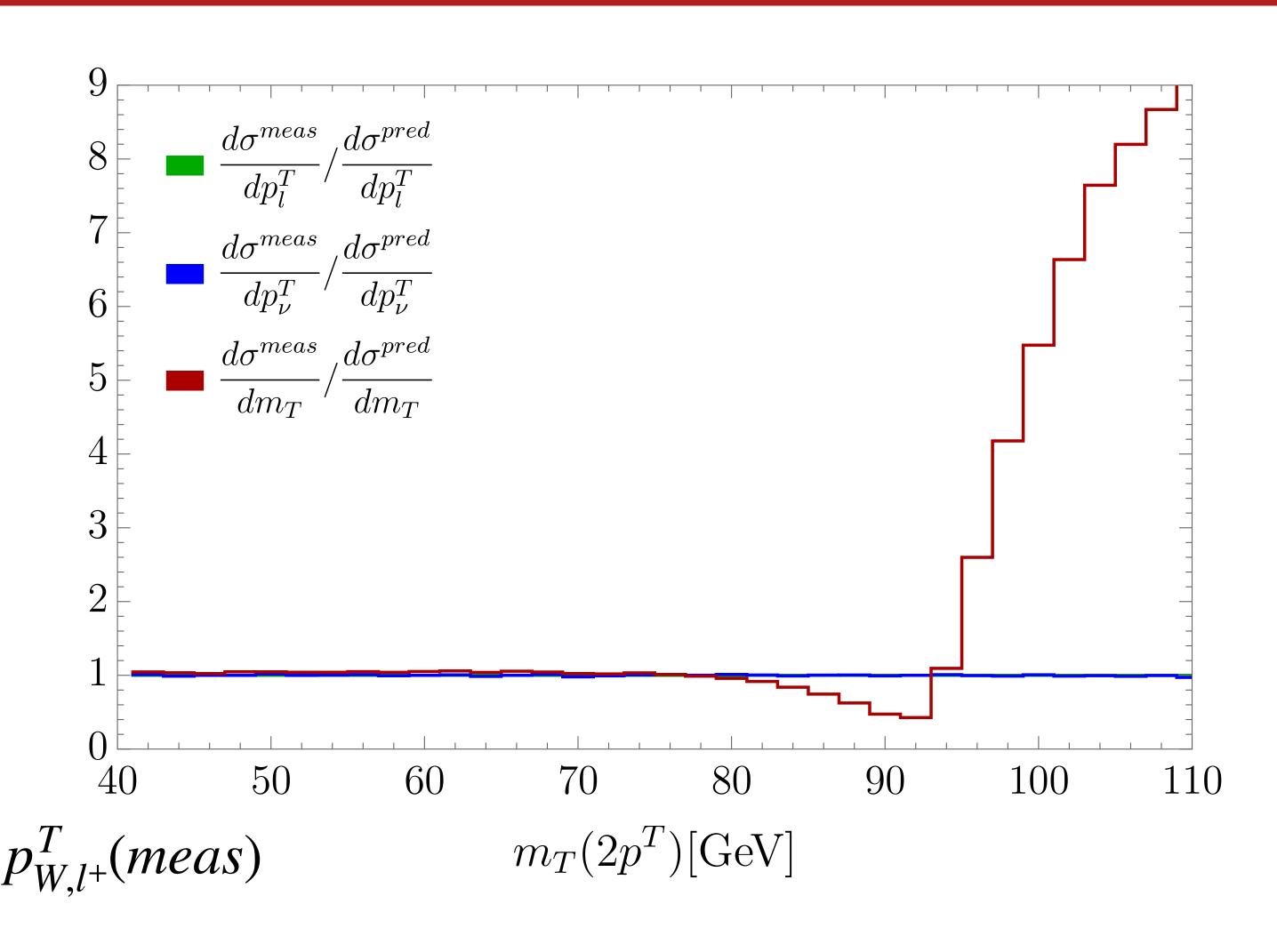
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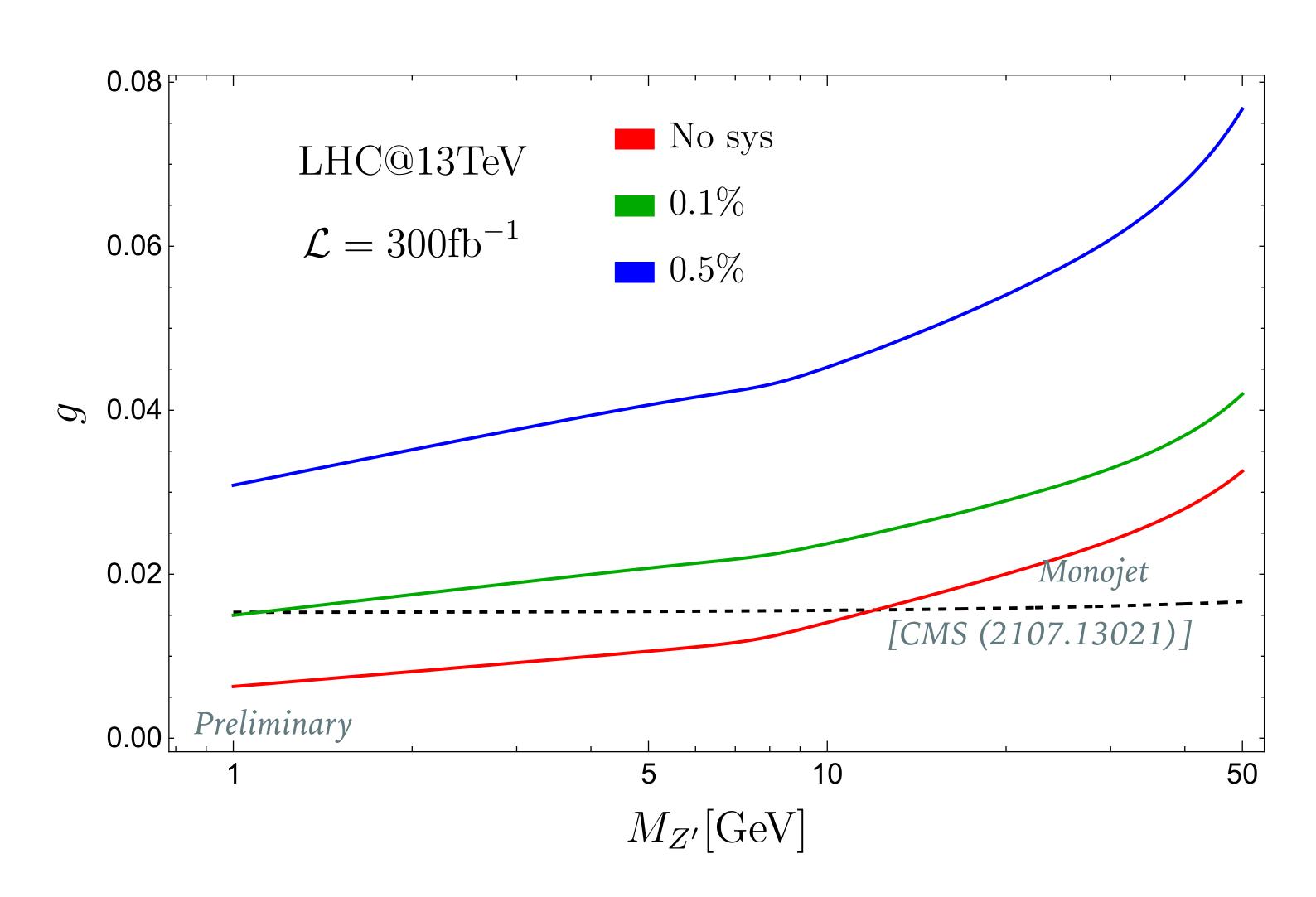
$$m_{T}^{Z} \rightarrow m_{T}^{Z}(pred) \neq m_{T}^{Z}(meas)$$



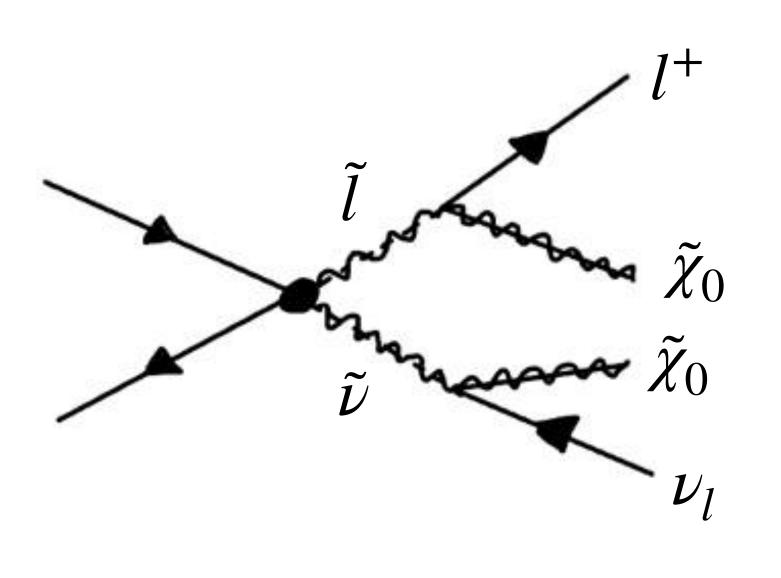
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- New physics can modify the Z sample (and so the calibration)
- A possible way out is to veto on Z+met

➤ *To cancel the gauge anomaly we need an anomalon sector and this may lead to additional constraints

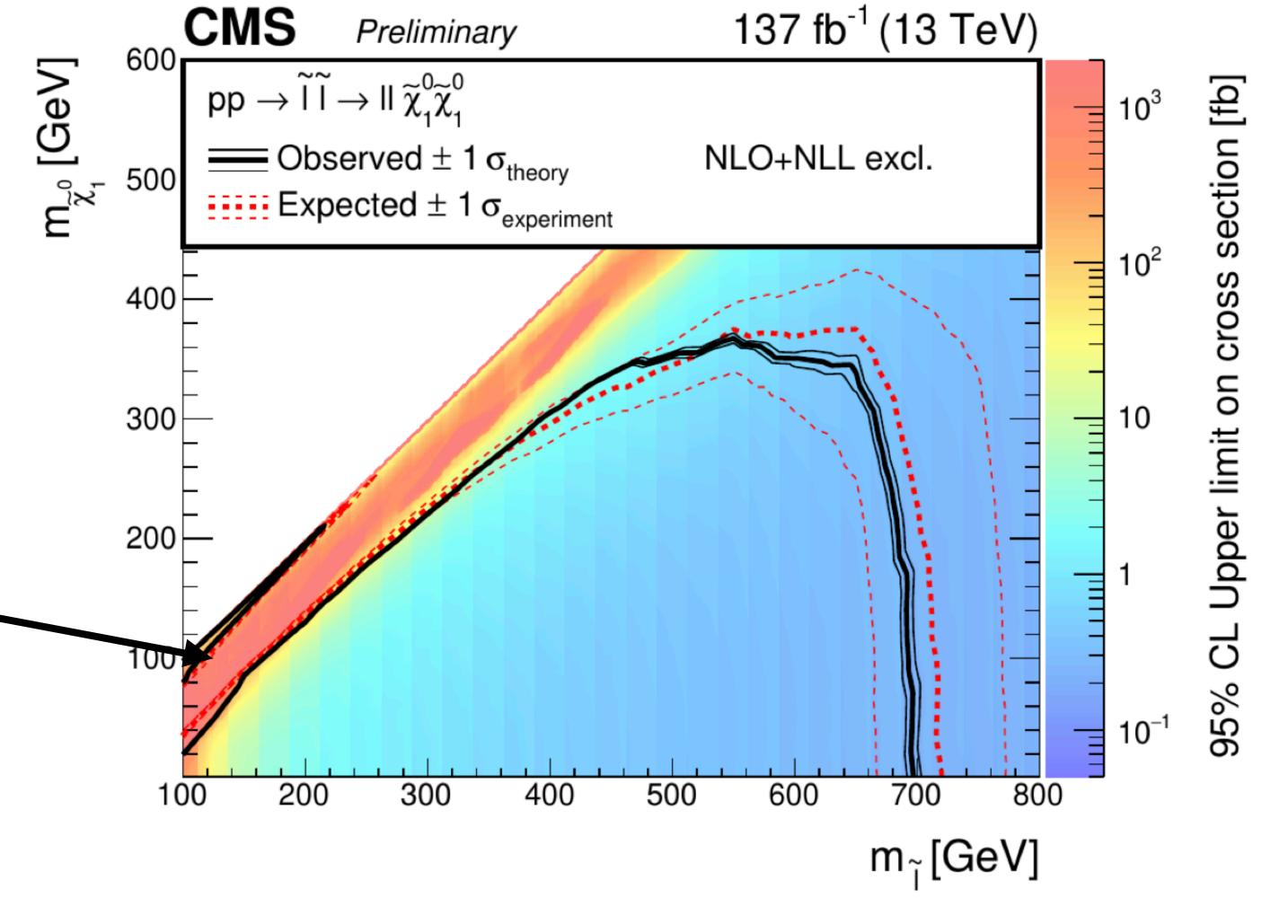


l+met not from a W



This region is not covered by "standard" searches

We can go in the SM region, where NP contaminates the W sample



CMS PAS SUS-21-008

Conclusion & Outlook

- ➤ Invisibly decaying new physics is an irreducible background for the W mass measurement
- ➤ Suited kinematic cuts allows to veto BSM from Z decay (crucial for the data driven approach)
- ➤ We performed a comprehensive study of the relevant BSM phenomenology and found competitive constraints

Conclusion & Outlook

- ➤ Invisibly decaying new physics is an irreducible background for the W mass measurement
- ➤ Suited kinematic cuts allows to veto BSM from Z decay (crucial for the data driven approach)
- ➤ We performed a comprehensive study of the relevant BSM phenomenology and found competitive constraints
- ➤ Precise measurements of the SM parameters are highly non-trivial test of the SM and powerful probes for New Physics:
 - ➤ Indirectly, through the EW fit
 - ➤ Directly, through precision tests of the SM distributions



Thank you!