



Incoherent effects and long term beam behavior below the e-cloud instability threshold

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Thanks to:

G. Arduini, V. Baglin, E. Benedetto, O.E. Berrig, K. Li, K. Ohmi, E. Metral, G. Rumolo, E. Schaposhnikova, F. Zimmermann

Overview

Incoherent effect: Trapping or not Trapping
that is the problem !

The power of the (a) pinch

Modeling and estimates for LHC

Towards a systematic approach

Incoherent effects

it is caused by the repeated crossing of a resonance by a beam particle

The single particle tune modulation is caused by transverse-longitudinal coupling created by: chromaticity, space charge, pinched electron cloud

The difference between the source of the coupling arises from the type of amplitude dependent detuning they create

source\feature	detuning	amplitude dependence	driving resonances	Experiment verification
Chromaticity	weak	weak	no (?)	yes
Space Charge	strong	strong/decrease	yes/no	yes/ongoing
Electron Cloud	strong/weak	strong/decrease	yes	?
Nonlinearities	strong/weak	strong/increase	yes	yes

Resonance crossing

Basic phenomena → Separatrix crossing

Trapping into resonance



although the tune is modulated particles remain locked on the resonance



Large excursion of particles locked to compensate the modulation of the detuning

All possible intermediate dynamical regimes are possible



Trapping or not Trapping ?

Scattering by resonance



particle cross the separatrix but does not remain inside the island



Islands gives a kick to particle

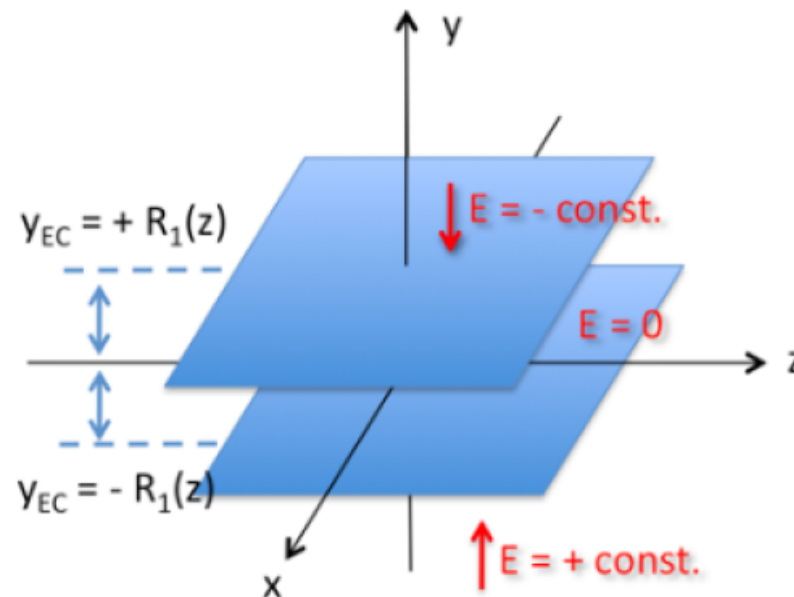
Modeling incoherent effect of EC

BEAM07

1D model of EC pinch

Electrons are in two planes moving apart according to the longitudinal position of a particle in a bunch

1 localized EC kick excite all structure resonances



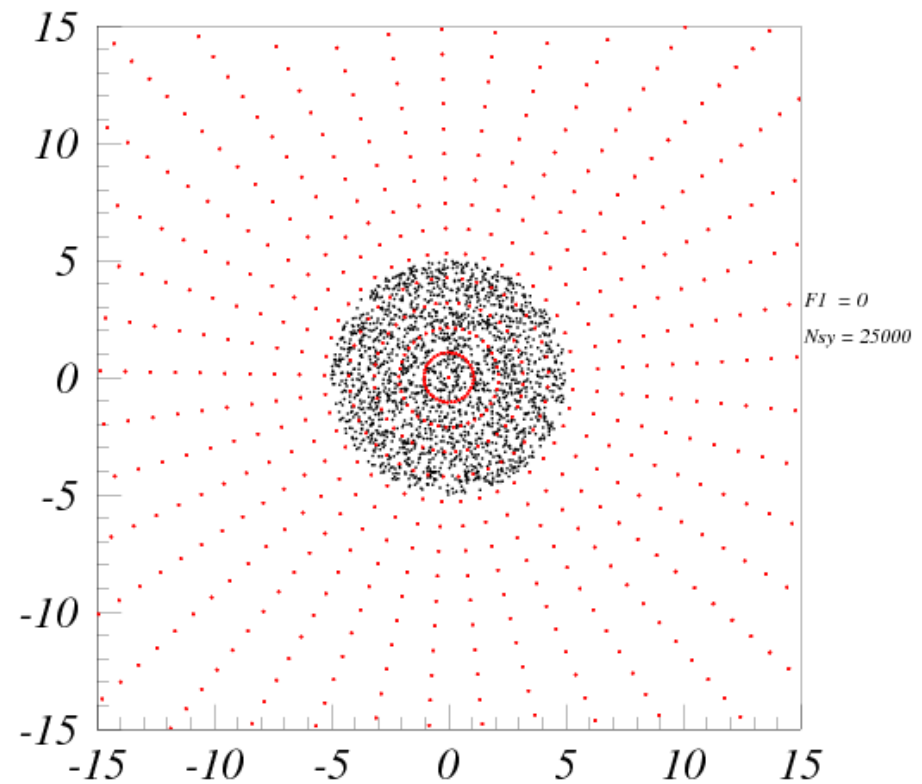


Crossing of the 4th order structure resonance

Trapping process for the 1D electron-cloud map

In red are show the orbit of the “frozen system”

This resonance crossing in NOT ADIABATIC

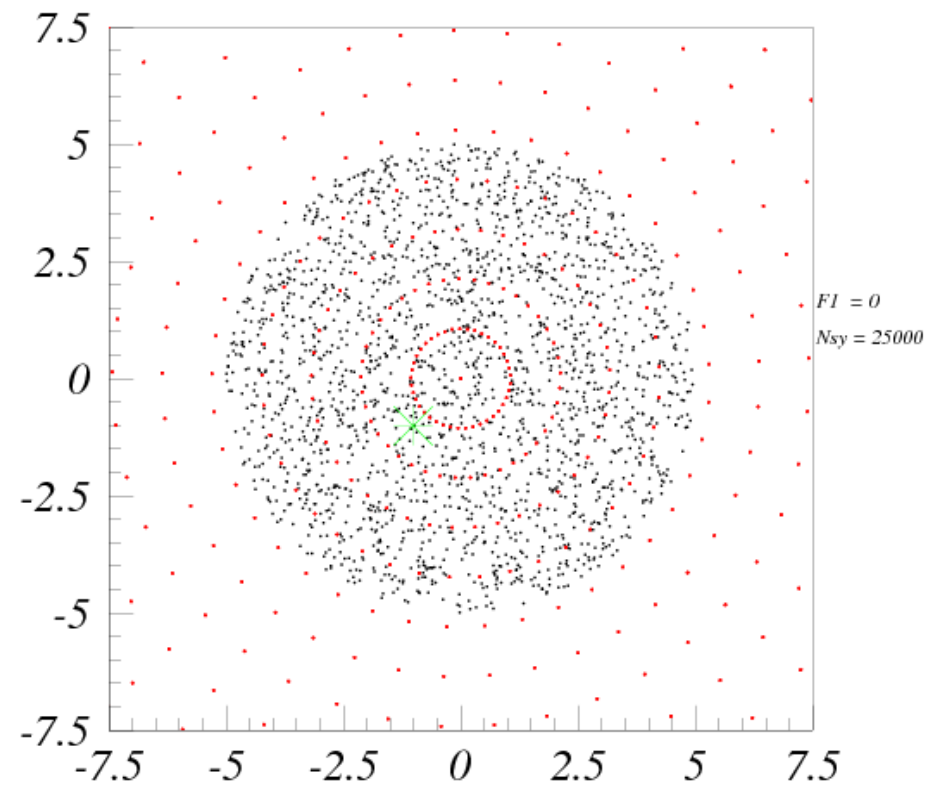


Sinusoidal Periodic Crossing: Period 25000 turns



The same periodic crossing
seen from the “frozen fixed point”

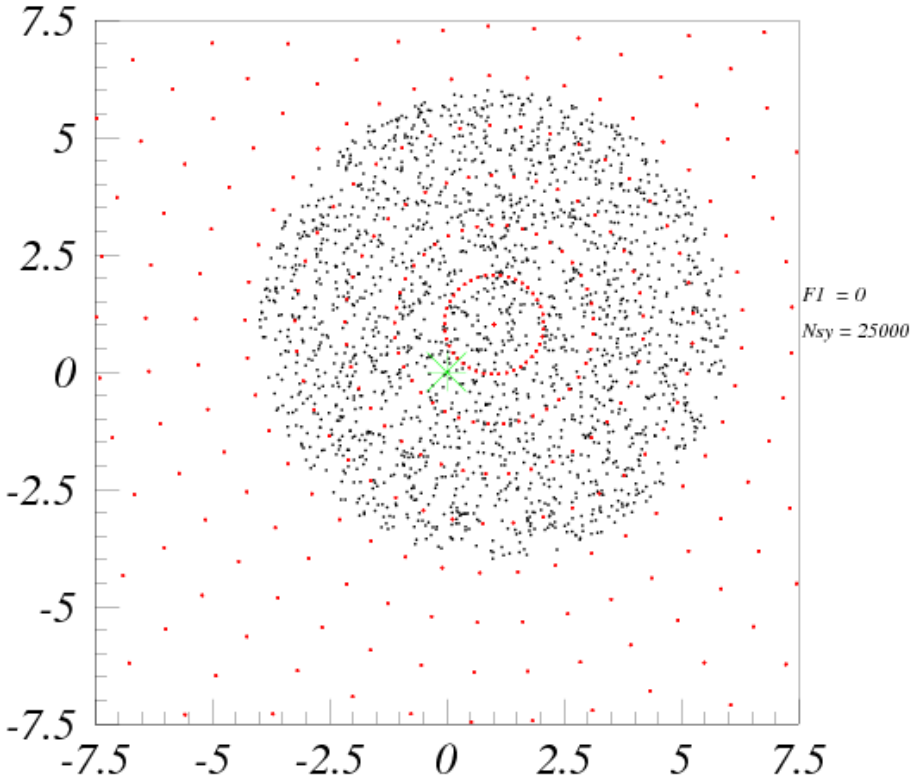
In green is shown where is the
real attraction point which exert
the trapping function



Sinusoidal Periodic Crossing: Period 25000 turns



The same beam dynamics seen from the reference frame of the “attraction point”



Sinusoidal Periodic Crossing: Period 25000 turns

Structure of EC-pinch

Scattering diffusive effect depends on the correct modeling
of the induced EC pinch structure

However

diffusive effect acts on time scale of 10^6 turns

An analytic modeling allow avoiding artificial emittance in long term tracking



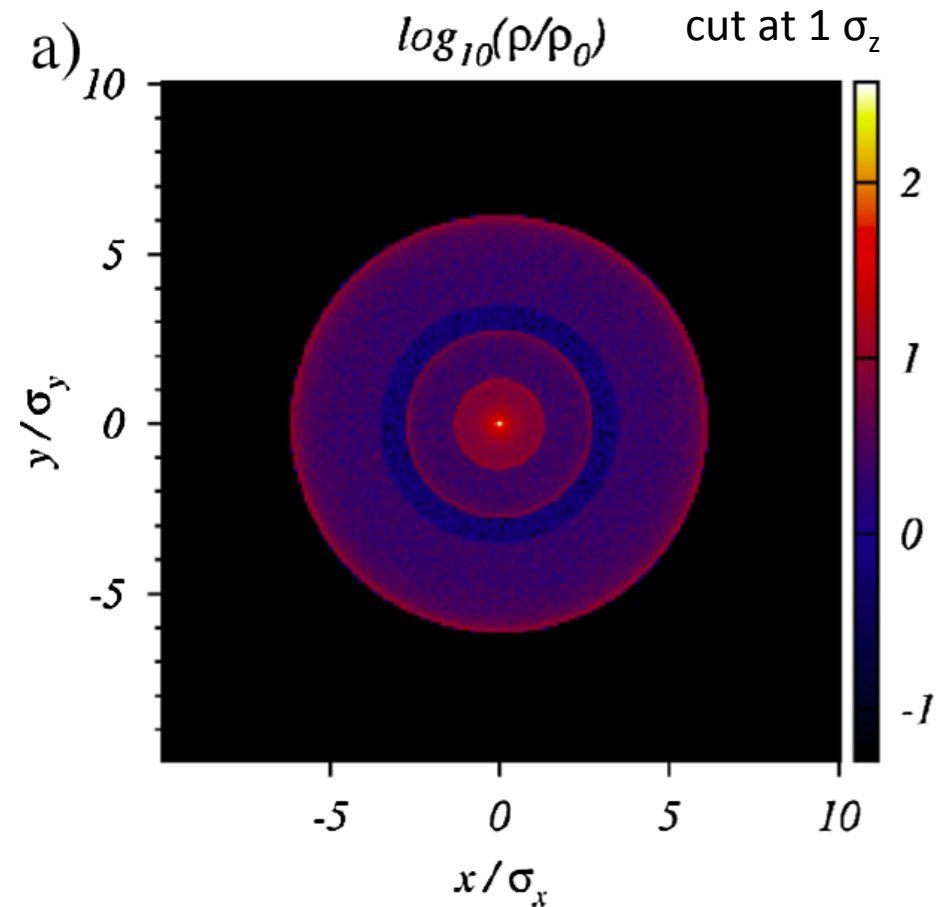
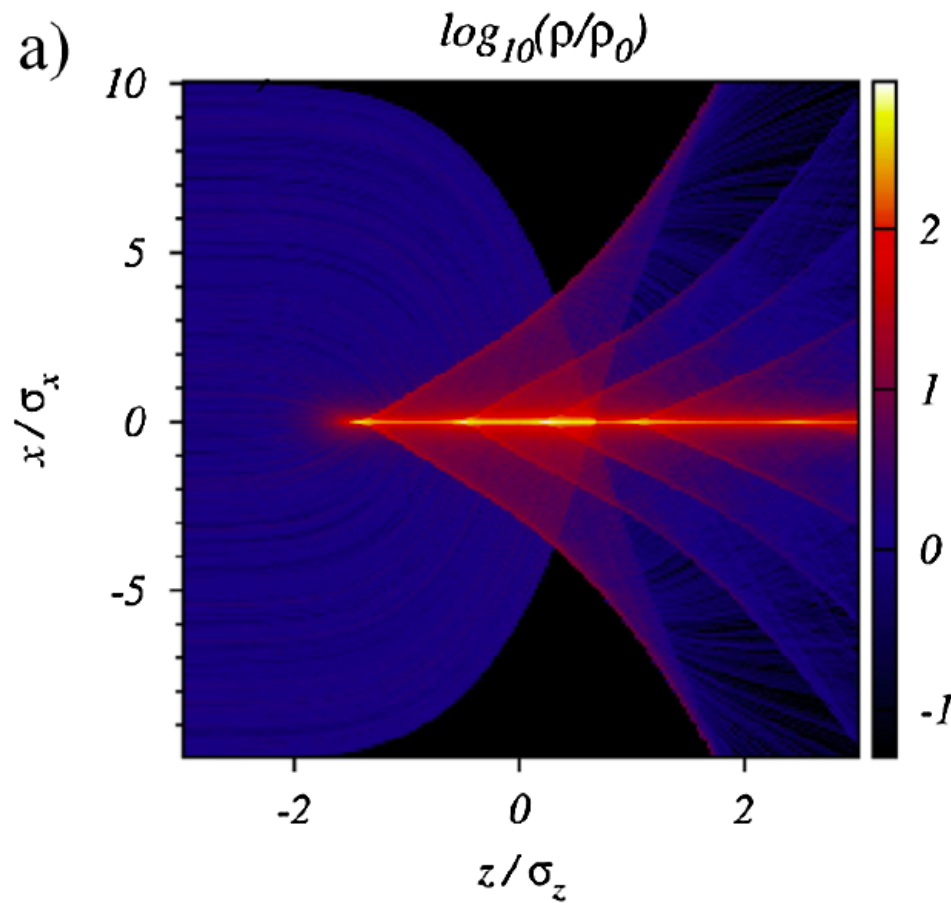
Closer look at the EC-pinch induced structure

EC-pinch in field free region

For nominal
LHC bunch

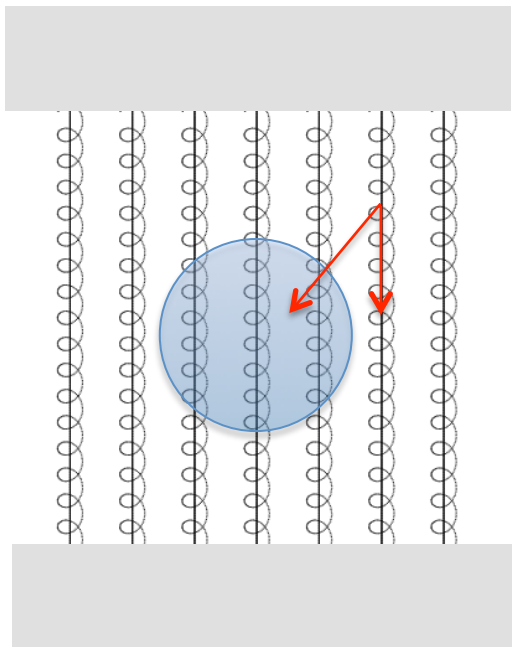
$N_b = 1.15 \times 10^{11}$, $\sigma_{x,y} = 0.88$ mm
 $\sigma_z = 11.4$ m. Initial electrons in a circle 10σ

No EC-self force

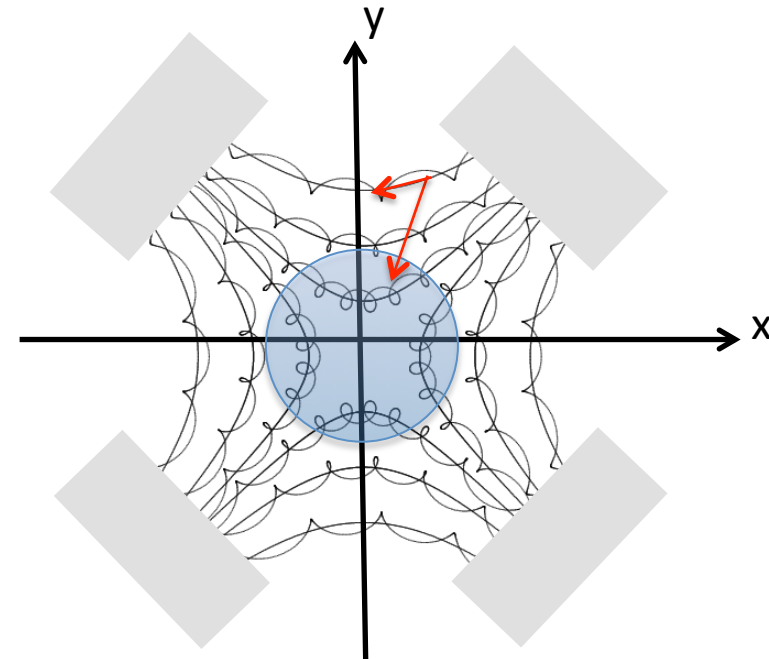


Strong field approximation

Dipole



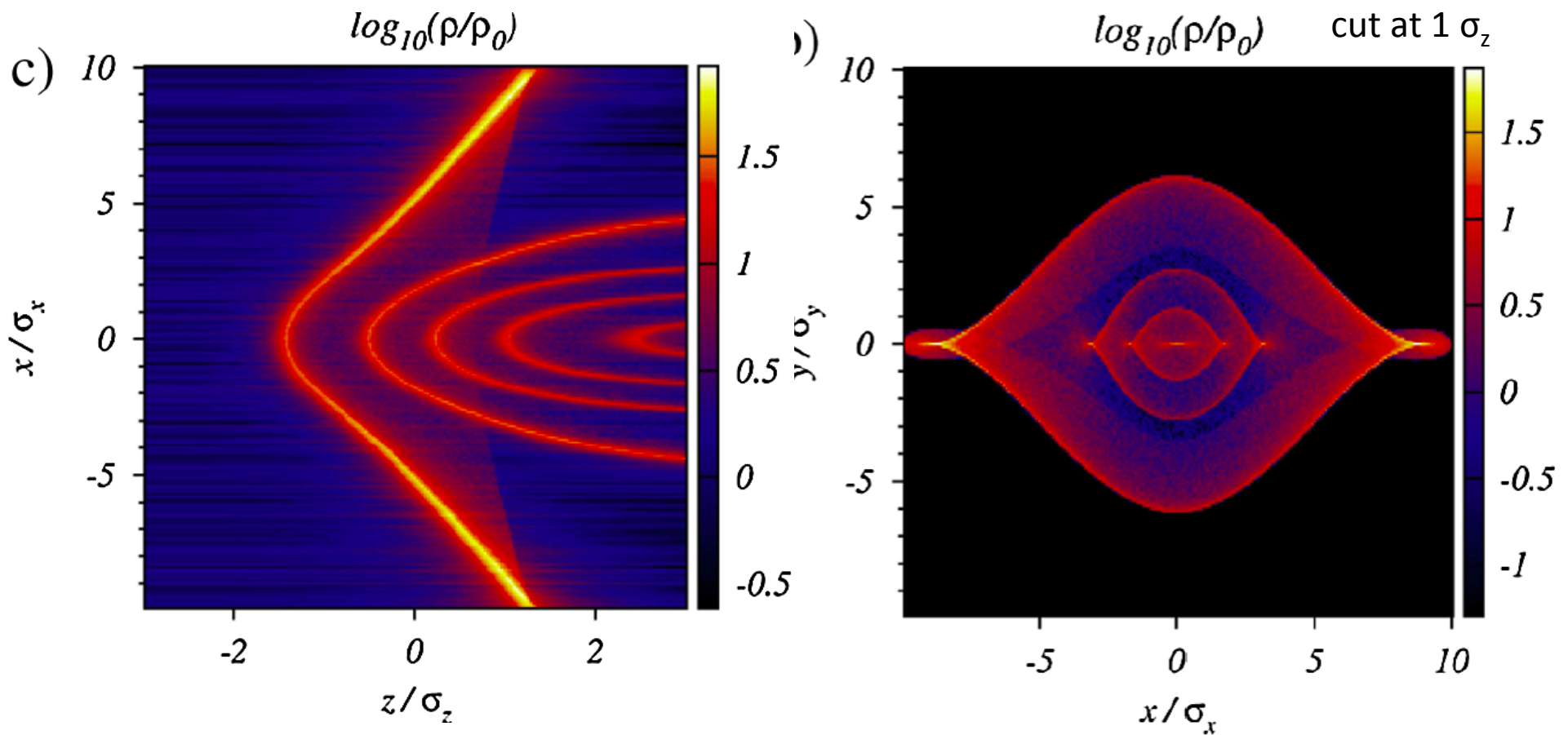
Quadrupoles



Force on the electrons is the projection of the force from the beam on the direction of the vector line field

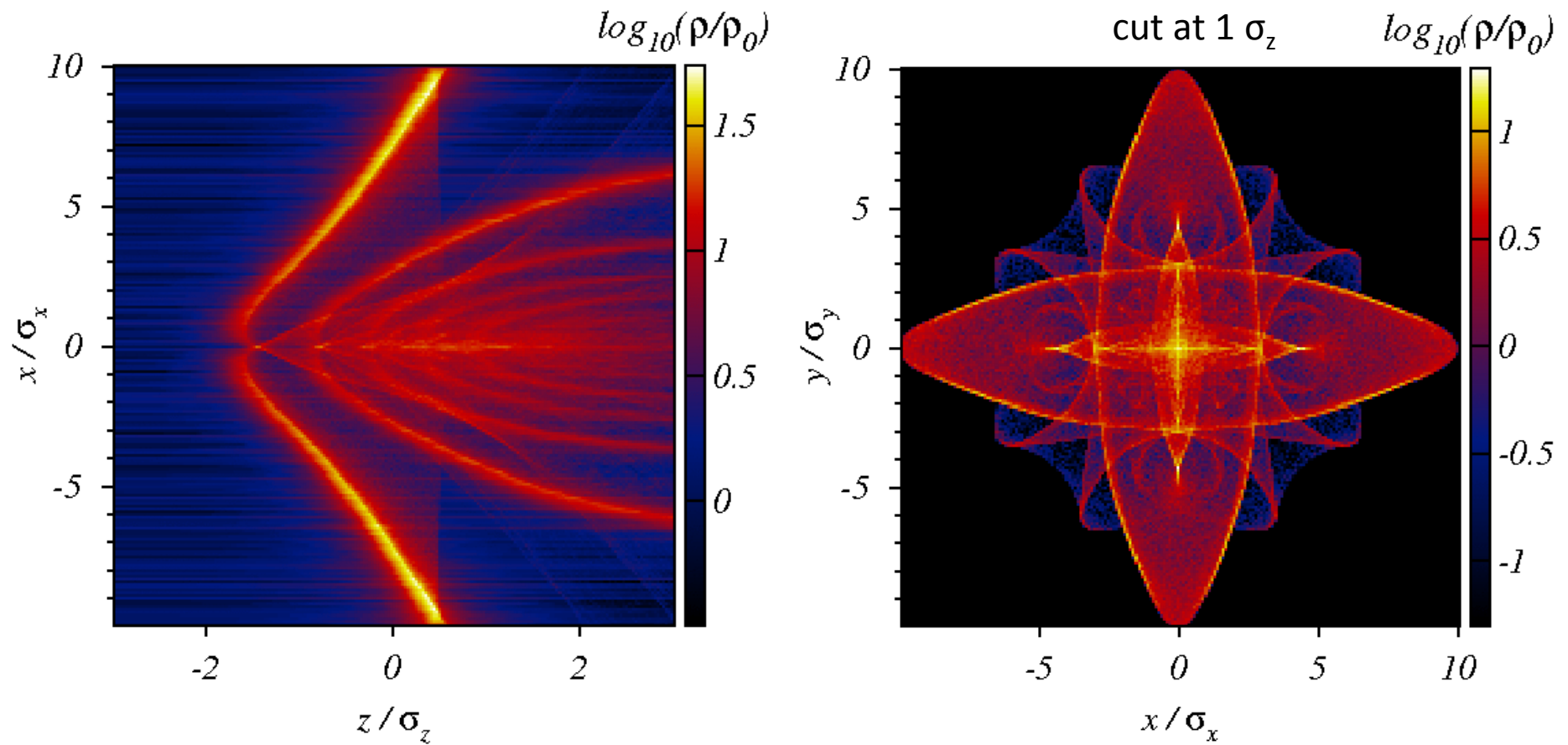
EC-pinch in dipoles

For nominal LHC bunch Based on the "strong field approximation"



EC-pinch in Quadrupoles

For nominal LHC bunch Strong B-field approximation



Approximations

Comments of K. Ohmi
and G. Rumolo

Electrons follow the magnetic field
This approximation is not valid at small radii

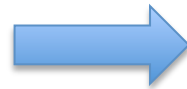
Comparison of quadrupoles with dipoles

$$\begin{array}{l} k_x = \frac{1}{B\rho} g \\ B_y = gx \end{array} \quad \longrightarrow \quad \frac{B_y}{B_0} = k_x \rho x$$

For $\rho \simeq 2.8 \text{ Km}$

$$\frac{B_{quad}}{B_0} = 0.023 \quad \text{at} \quad r = \sigma = 0.88 \text{ mm}$$

Strong B approximation not applicable



**Analysis of complete dynamics
is necessary**

EC-pinch in Quadrupole analysis

Complete equations
$$\frac{d}{dt} \left(m_e \gamma_e \frac{d\vec{R}}{dt} \right) = e\vec{E} + e \frac{d\vec{R}}{dt} \times \vec{B}$$

Assumptions of
2D bunch and
2D quadrupole

- $\vec{E} = (E_x, E_y, 0)$
- $\vec{B} = (B_x, B_y, 0) = (gy, gx, 0)$, quadrupole field.

Equation of
motion of
one electron

$$x = \sigma_r \tilde{x}, \quad y = \sigma_r \tilde{y}, \quad z = \sigma_r \tilde{z}.$$

$$\frac{d^2 x}{ds^2} = \frac{e}{c^2 m_e \gamma_e} E_x - k \frac{\gamma_p m_p}{\gamma_e m_e} \frac{dz}{ds} x$$

$$\frac{d^2 y}{ds^2} = \frac{e}{c^2 m_e \gamma_e} E_y + k \frac{\gamma_p m_p}{\gamma_e m_e} \frac{dz}{ds} y$$

$$\frac{d^2 z}{ds^2} = k \frac{\gamma_p m_p}{\gamma_e m_e} \left(\frac{dx}{ds} x - \frac{dy}{ds} y \right)$$



Including the force exerted by the protons the equations become

$$\frac{d^2 \tilde{x}}{ds^2} + k\sigma_r \frac{\gamma_p m_p}{\gamma_e m_e} \frac{d\tilde{z}}{ds} \tilde{x} + \omega_e^2(s) \tilde{x} = -\tilde{x} \frac{\omega_e^2(s)}{\tilde{r}^2} \left[2 \left(1 - e^{-\frac{\tilde{r}^2}{2}} \right) - \tilde{r}^2 \right]$$

$$\frac{d^2 \tilde{y}}{ds^2} - k\sigma_r \frac{\gamma_p m_p}{\gamma_e m_e} \frac{d\tilde{z}}{ds} \tilde{y} + \omega_e^2(s) \tilde{y} = -\tilde{y} \frac{\omega_e^2(s)}{\tilde{r}^2} \left[2 \left(1 - e^{-\frac{\tilde{r}^2}{2}} \right) - \tilde{r}^2 \right]$$

$$\frac{d^2 \tilde{z}}{ds^2} = k\sigma_r \frac{\gamma_p m_p}{\gamma_e m_e} \left(\frac{d\tilde{x}}{ds} \tilde{x} - \frac{d\tilde{y}}{ds} \tilde{y} \right)$$

Note: the effect of the quadrupole depends on $\frac{d\tilde{z}}{ds}$

But the change of kinetic energy is

$$\tilde{E}_k(\tilde{r}, \tilde{r}_0) = \int_{\tilde{r}}^{\tilde{r}_0} \frac{\omega_e^2(s)}{\tilde{r}} \left[2 \left(1 - e^{-\frac{\tilde{r}^2}{2}} \right) \right] d\tilde{r} = \omega_e^2(s) F(\tilde{r}, \tilde{r}_0)$$



Therefore certainly $\frac{1}{2} \left(\frac{d\tilde{z}}{ds} \right)^2 \leq \tilde{E}_k(0, \tilde{r}_0)$

which allows to estimate that for the electrons that start at rest

$$\beta_e \leq \sqrt{2} \sigma_r \omega_e(s) \sqrt{1.32 + 2 \ln(\tilde{r}_{pipe}/2)}$$

For LHC bunch $\rightarrow \omega_e(0) = 38.26 \text{ m}^{-1}$, $r_{pipe} = 0.025 \text{ m}$, and $\sigma_r = 0.88 \text{ mm}$

$$\rightarrow \beta_e \sim 0.1 \quad \rightarrow \gamma_e \simeq 1$$

Strong/Weak B threshold

The field B produced by the quadrupole is weak with respect to the field created by the protons if

$$\left(k\sigma_r \frac{\gamma_p m_p}{\gamma_e m_e}\right)^2 \left(\frac{d\tilde{z}}{dt}\right)^2 = < \frac{\omega_e^4(s)}{\tilde{r}^4} 4 \left(1 - e^{-\frac{\tilde{r}^2}{2}}\right)^2$$

But we can estimate the maximum $\frac{d\tilde{z}}{ds}$ from $\tilde{E}_k(\tilde{r}, \tilde{r}_0)$

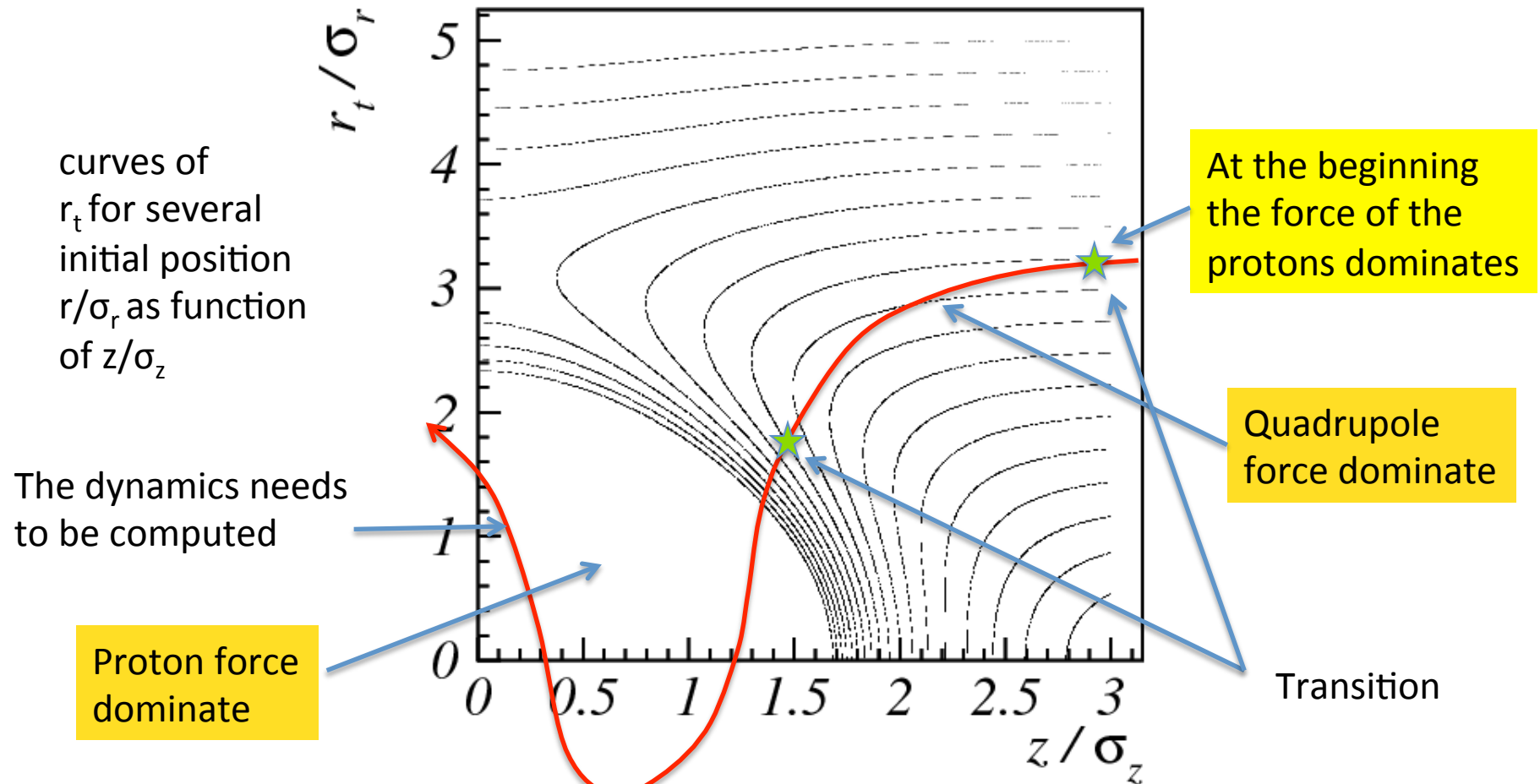
These two equations allow the definition of a threshold radius r_t such that all the electron $\tilde{r}_0 < \tilde{r}_t$ are dominated by a Weak B dynamics



There are region in which the pinch evolves very similar to a field free region and others in which is dominated by the quadrupole field

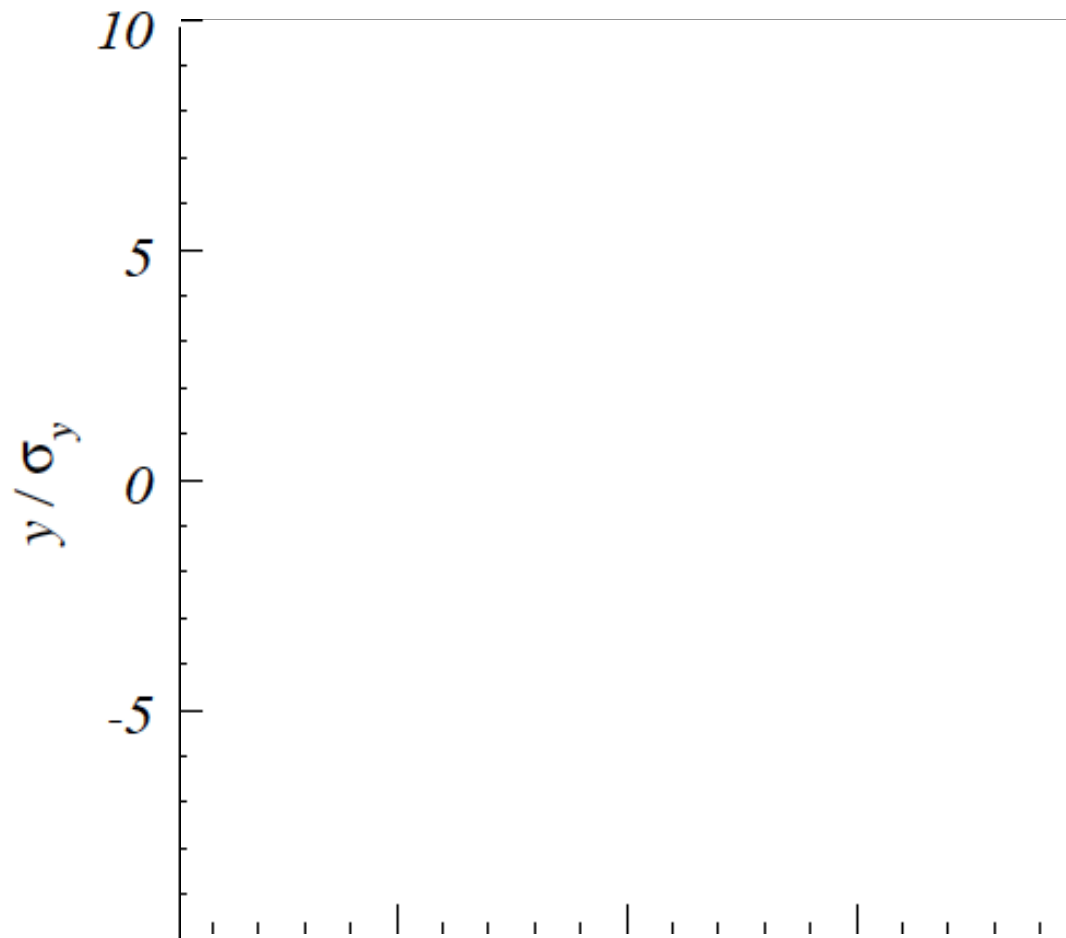
Example for the LHC injection energy

The transition of dominance of forces depends on the initial position of the electrons

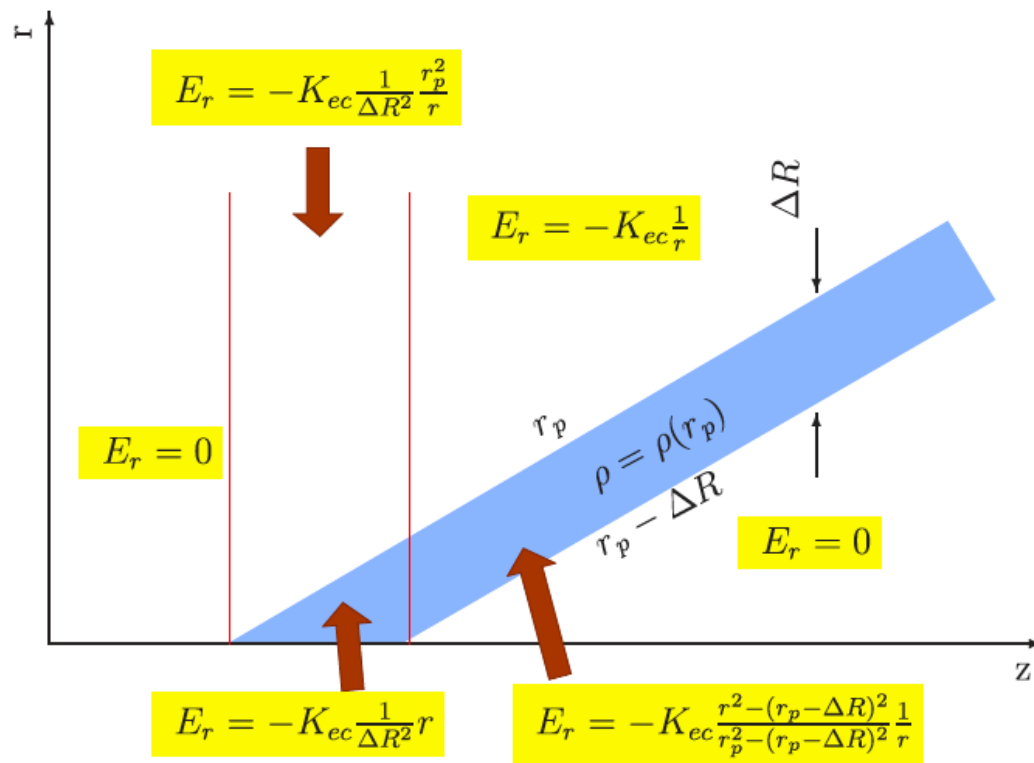


The full dynamics and transition strongly depends on the initial position of each particle

Pinch in Quadrupoles: strong B limit



Present modeling for long term tracking



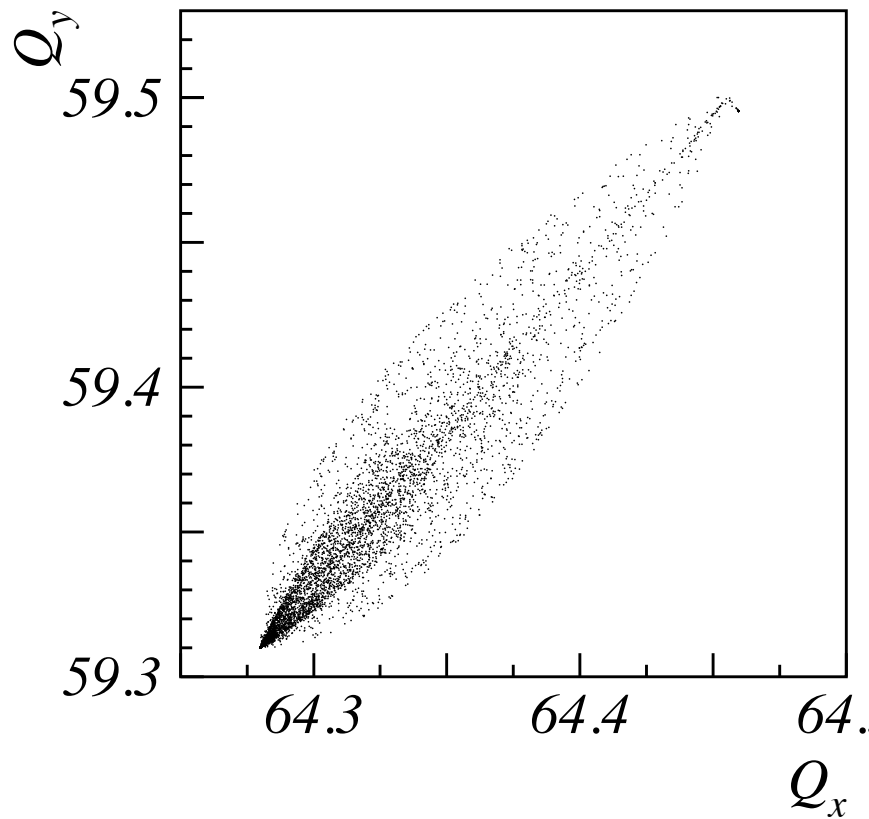
Input parameters

- 1) Maximum detuning
- 2) Ring slope
- 3) Number of rings

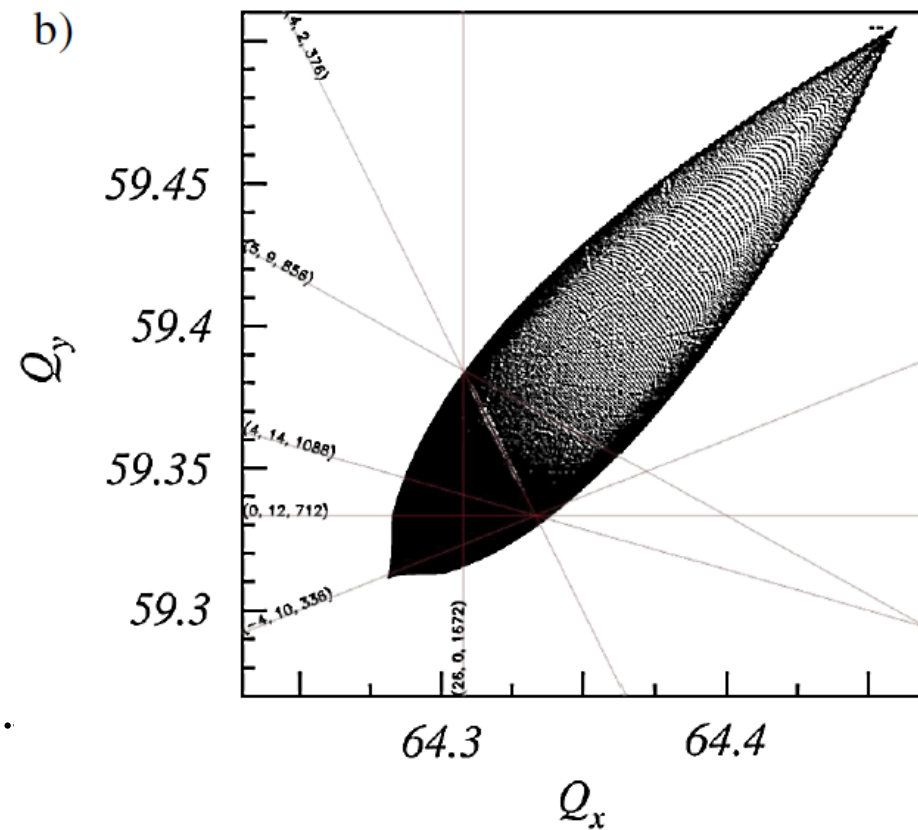
Tune footprint

Example for $DQ = 0.2$

Tune-footprint
of the full bunch



EC kicks create a web of high order
structure resonances



Estimates for LHC

Input parameter: the maximum incoherent tune spread
created by EC

Related to electron density

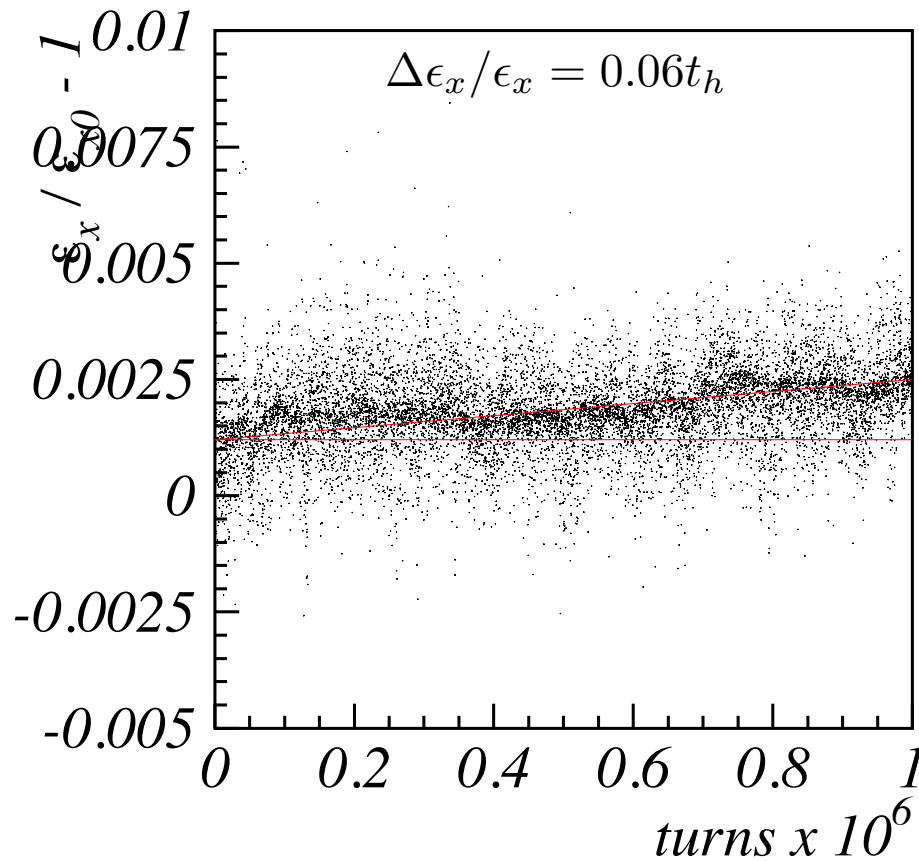
$$DQ_e = 0.1 \quad DQ_e = 0.02 \quad \rightarrow \text{K. Li}$$

$$Q_x = 64.28 \quad Q_y = 59.31 \quad (\text{injection})$$

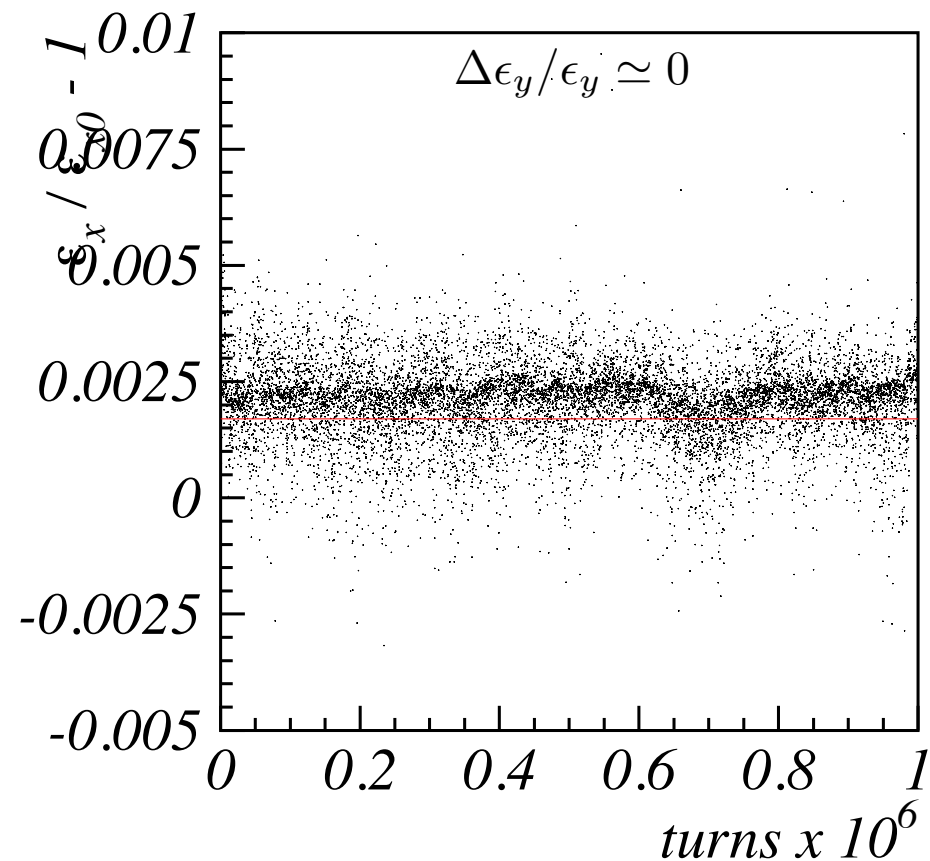
Pinch modeling: as field free \rightarrow conservative estimates

Estimates for LHC

EC Kicks all dipoles all quad,
but EC structure of free pinch



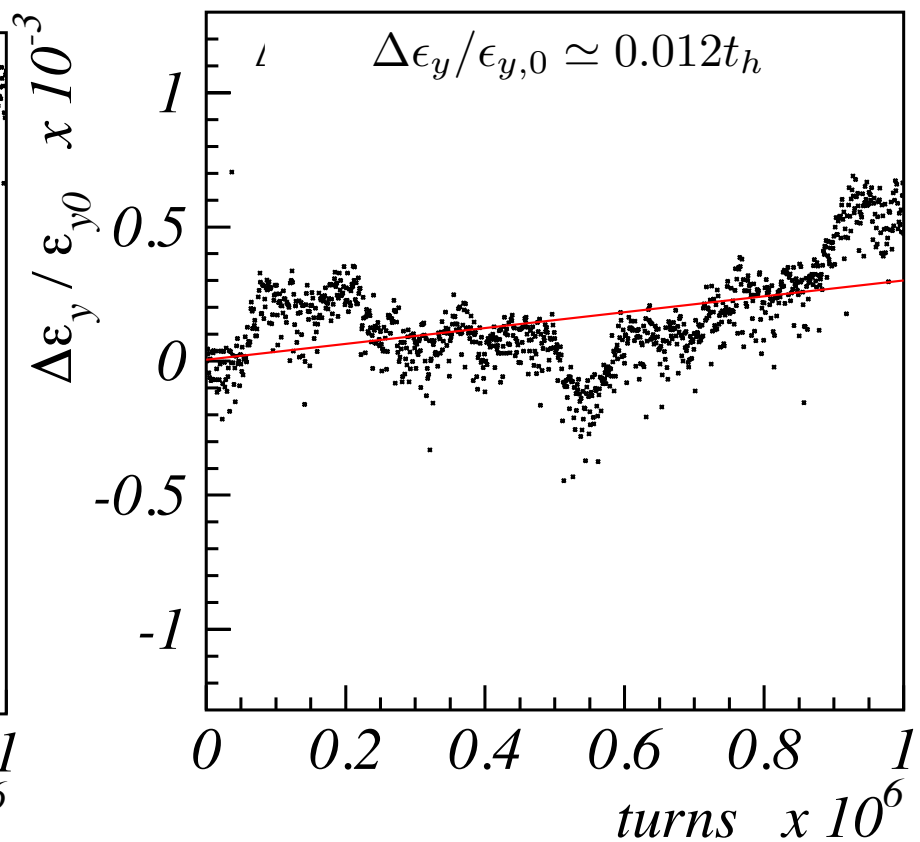
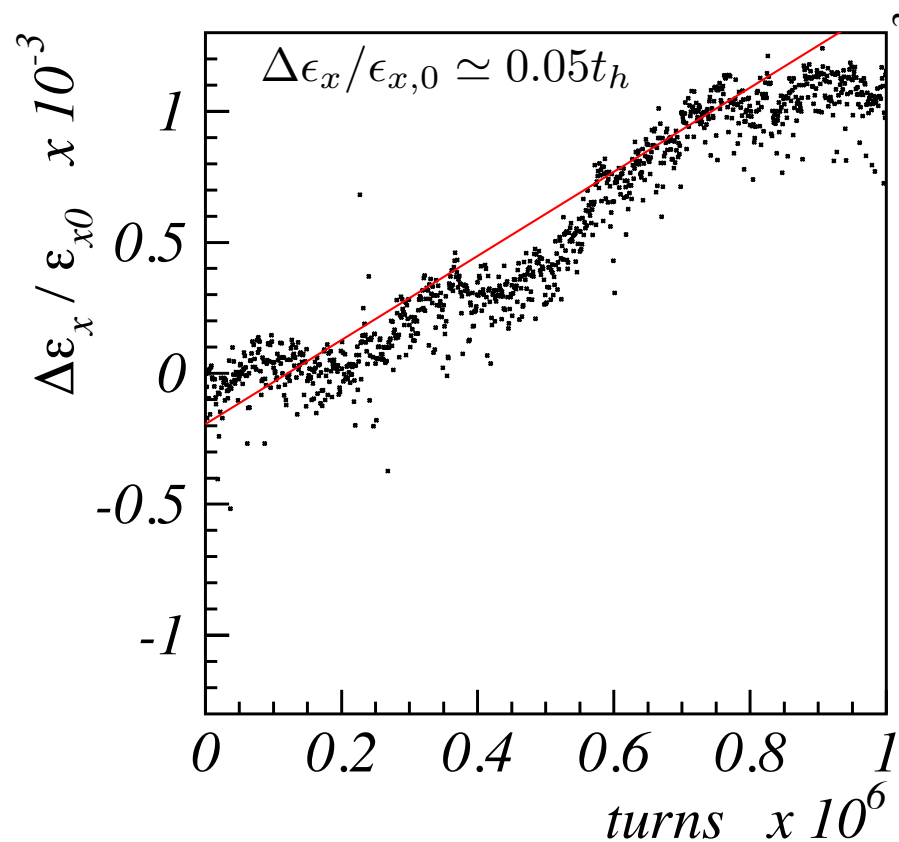
Simulation DQe = 0.1



Sensitivity to number of macro-particles

EC Kicks only in Dipoles,
but EC structure of free pinch

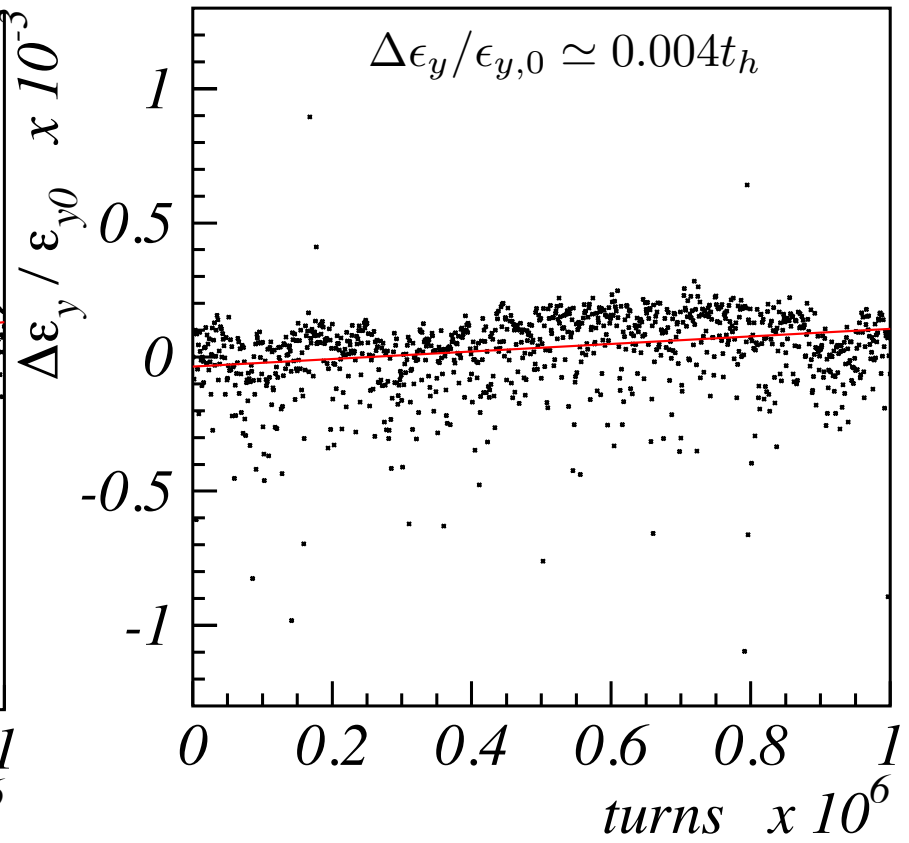
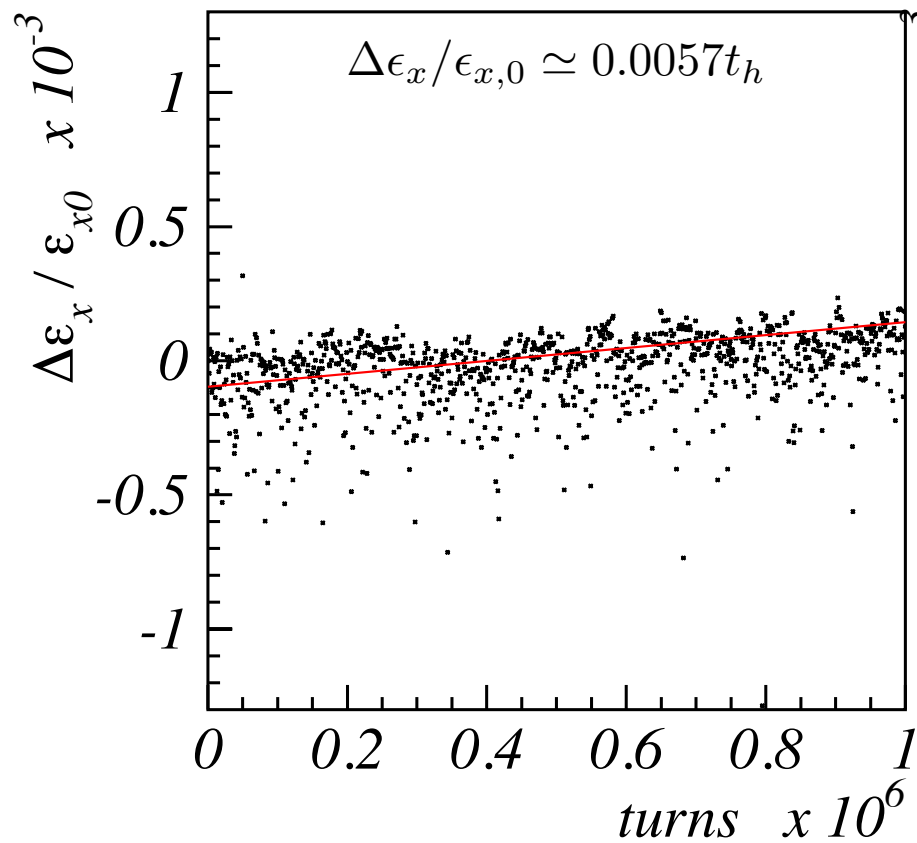
1152 dipoles DQe = 0.1





EC Kicks in all dipoles + Quadrupoles ,
but EC structure free field

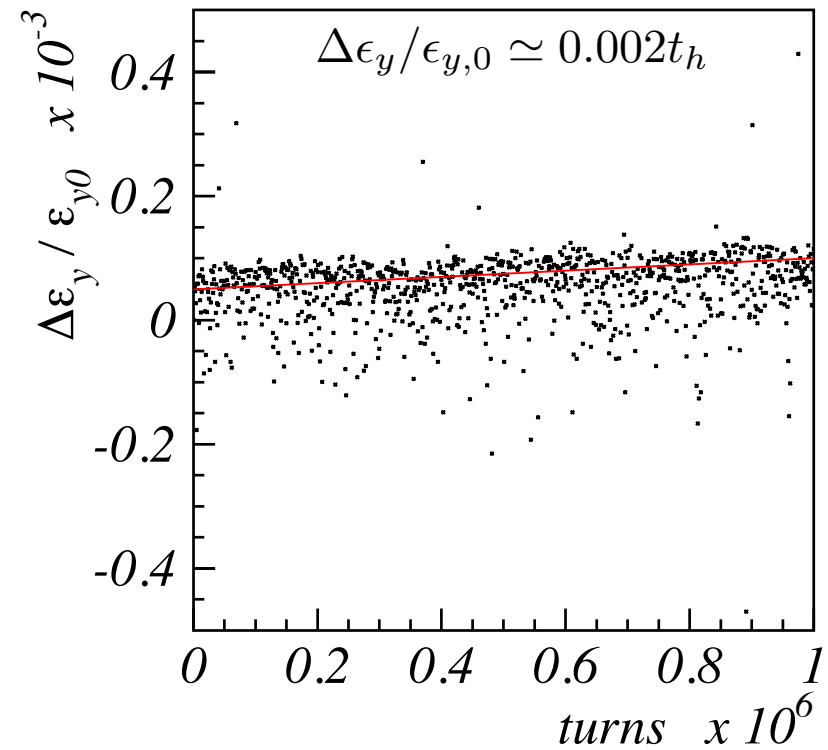
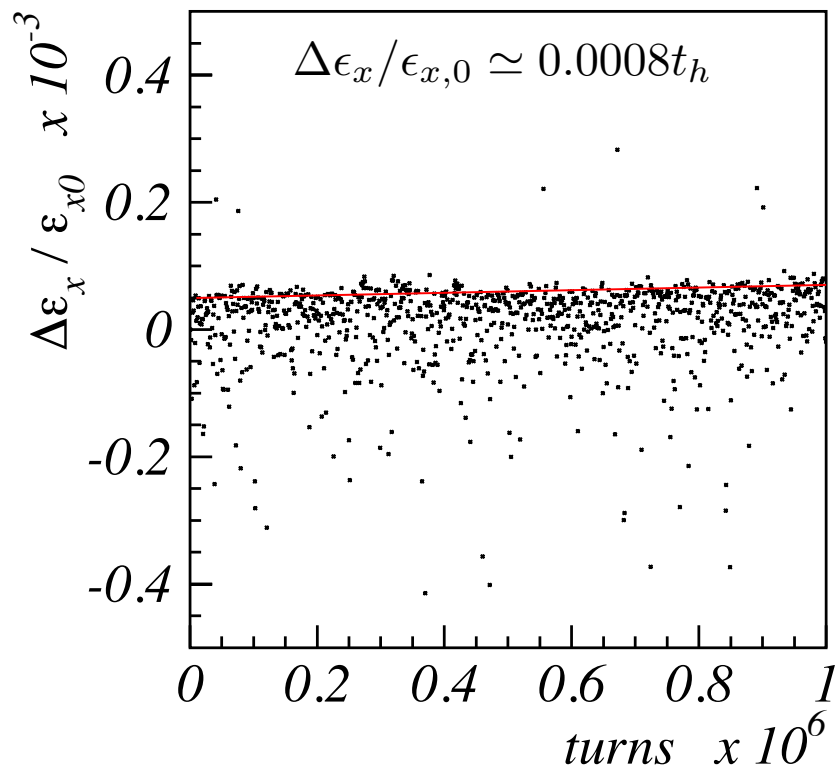
DQ = 0.02 Npar = 8000





EC Kicks all Dipoles,
EC structure of free pinch

DQ = 0.02 Npar = 19000



Present estimates

Issues affecting the results:

sensitivity to number of particles
sensitivity to EC-structure
sensitivity to EC kick random error
sensitivity to actual density of electrons
that is to maximum tune spread

Results sensitive to parameters: but emittance growth very slow

DQe	Dex/ex hour	Dey/ey hour	N macro- particles	Turns in simulation	Dipole arc 1152	Dipole	Quad
0.1	6%	~0%	8000	10^6	yes	yes	yes
0.1	5%	1.2%	1.6×10^4	10^6	yes	no	no
0.02	0.57%	0.4%	8000	10^6	yes	yes	yes
0.02	0.08%	0.2%	1.9×10^4	10^6	yes	yes	no

Approaching a Systematics

Difficult to have all information

EC kicks should take into account of

- 1) actual sizes of the beam (betax/betay)
- 2) pinch with actual electron extension (pipe vs. beam size)
- 3) realistic optics → correct phase advance among EC kicks → correct excitation of high order resonances
- 4) analytic modeling (acceptable approximation) of Coulomb field

Creating an element “EC Kick” which can be included into the accelerator structure for reaching the best modeling.

Help of O.E. Berrig, V. Baglin → location of solenoids, coated-uncoated sections



modeling the effect of random error on EC kicks → excitation of random resonances



Large overlap of topics between space charge and electron cloud incoherent effects. Development of models for analytic-diffusion based prediction is desirable to tracking simulations

Development of smart strategies for long term simulations (beam dynamics scaling)

**LHC issue: what will be the flexibility of LHC working point ???
over which range of Q_x, Q_y is it possible to move ???**