

Non-resonant searches at the TeV scale

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N.OF-H-INV-SIT

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Introduction

 Many unresolved questions that the Standard Model cannot answer

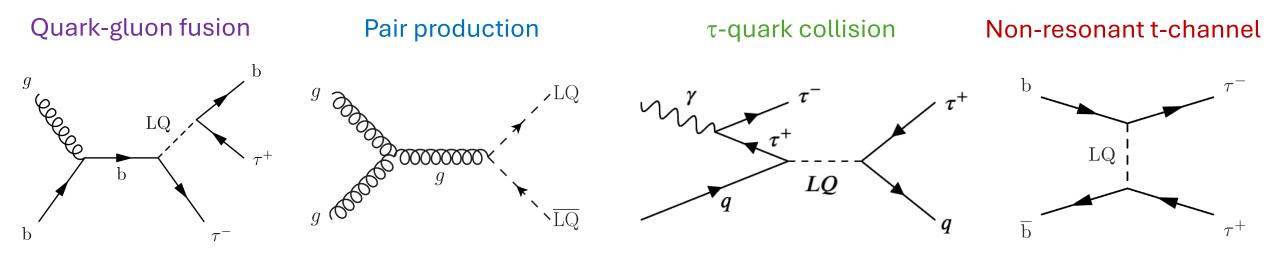
 (eg. dark matter, Higgs mass fine tuning, baryogenesis, neutrino masses, ...)



- Wide range of new physics searches are conducted at the LHC, such as for new resonances (Louis' talk), Extended Higgs sector (Shigeki's talk), Long-lived particles (Guglielmo's talk), precision measurements (Andrew's talk)
- Searches that are covered in this talk:
 - New symmetries giving rise to new particles, eg. Leptoquarks, Supersymmetry, Heavy Neutrinos
 - Extra dimensions leading to broad excesses or periodic signals (clockwork), or Quantum Black Holes
- Can only present a few results ^{(CMS Results}, ATLAS Results)
- Non-resonant Higgs boson pair production not covered in this talk (> John's talk tomorrow)

Leptoquarks

- In the SM there are striking similarities between quark and lepton families, but no explanation for it
- Could there be a deeper symmetry between leptons and quarks?
- Featured in models like Grand Unified Theories, compositeness, technicolor models, superstrings, R-violating supersymmetry, ...
- Hypothetical color-triplet bosons carrying both a baryon and a lepton number, with fractional charge
- Decay to a lepton (or neutrino) and a quark
- Parameters: λ_{LQ} (coupling), β (BR of LQ to charged lepton+quark), LQ mass, κ (vector LQs, <u>hep-ph/9610408</u>)
- Many production modes, for example:



Leptoquarks decaying to μ + b

↓LQ

g 000000000-

000000000

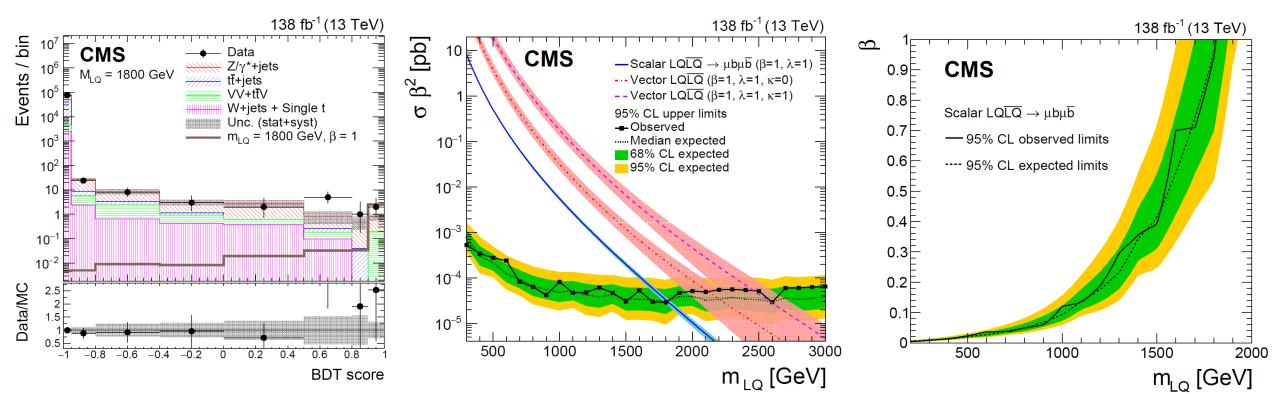
LO

١<u>LQ</u>

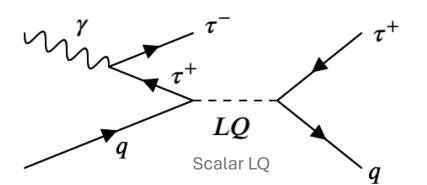
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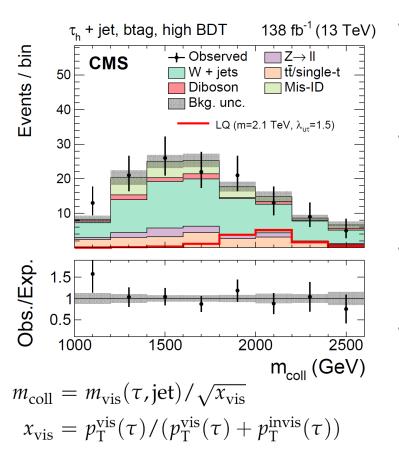
- Only LQ pair production considered
- Final state: Two muons (p_T > 53 GeV, m_{µµ}>250 GeV), two jets (p_T > 50 GeV) among which at least one is b-tagged
- Two control regions for background estimation (Z, tt; VV, ttV)
- BDTs trained on kinematic variables of the muons and jets to enrich a selection in signal events
- BDT cut optimized for each LQ mass hypothesis. **No excess found, but strong limits obtained!**



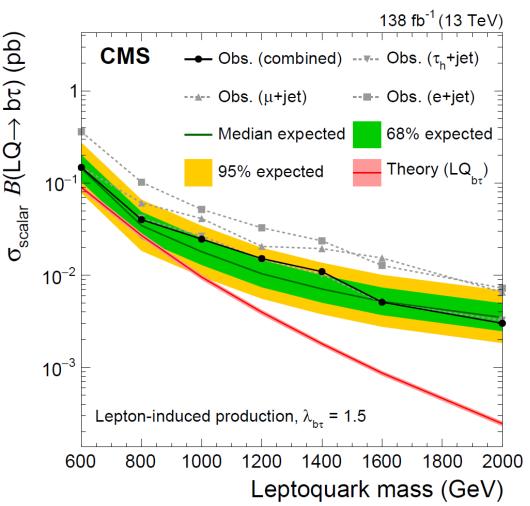
Leptoquarks via τ + quark collisions 2308.06143



- Novel production mode, made possible by advancements in the lepton and photon density functions of the proton
- Exploring three channels: hadronic τ + jet, e + jet, μ + jet, further separated if jet is b-tagged or not (different LQ decays)

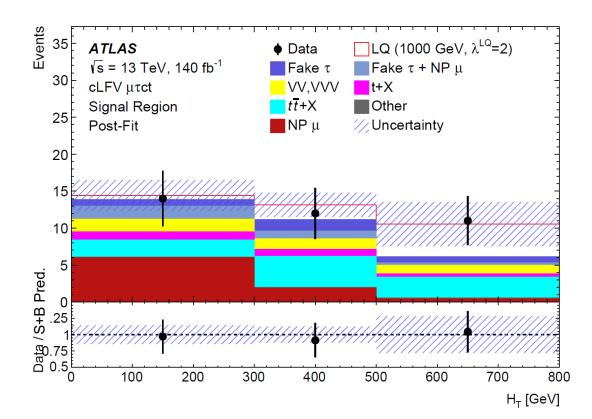


- BDTs trained in each
 channel to separate signal
 from background
- Categories based on the BDT score → 7 signalenriched categories
- Final discriminant: collinear mass
- W+jets background normalization estimated from control region that is fitted simultaneously



Leptoquarks via charged-lepton-flavor violation 2403.06742 6/22

- SM-predicted charged-lepton-flavor violation (cLFV) rates well below experimental sensitivity
- Leptoquarks are a candidate for introducing BSM cLFV interactions, ie. in the Scalar Leptoquark Model S₁
- Model introduces couplings between all up-type quarks and all charged $\lambda_{ki} \in \begin{pmatrix} \lambda_{t\tau} & \lambda_{c\tau} & \lambda_{u\tau} \\ \lambda_{t\mu} & \lambda_{c\mu} & \lambda_{u\mu} \\ \lambda_{te} & \lambda_{ce} & \lambda_{ue} \end{pmatrix} \equiv \lambda^{LQ} \begin{pmatrix} 10 & 1 & 0.1 \\ 1 & 0.1 & 0.01 \\ 0.1 & 0.01 & 0.001 \end{pmatrix}$



Analysis:

Final state: 2 µ (same-sign), 1 τ_{had}, ≥ 1 jet (including exactly one b-jet)

u, c

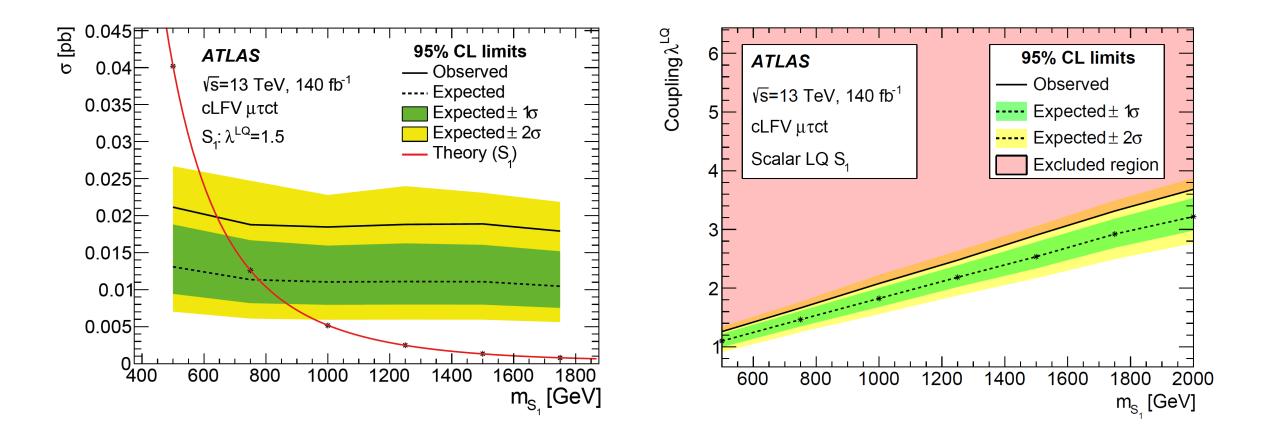
 S_1

u, c

- Final discriminant: scalar sum of the lepton and jet transverse momenta $\rm H_{T}$
- Slight excess at high H_T bins (~ 1.6 σ) here shown for m_{S1}=1 TeV and λ^{LQ} =2.0
- Fit includes SR and a CR enriched in ttbar+ μ events

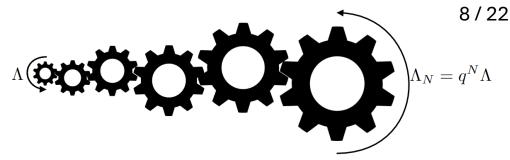
Leptoquarks via charged-lepton-flavor violation 2403.06742 7/22

- Excess almost not dependent on the S₁ mass, cross section limits almost flat
- LQ excluded for masses below 620 GeV, assuming λ_{LQ} = 1.5
- Exclusion of coupling values of λ_{LO} = 1.3 to 3.7 for LQ masses between 0.5 and 2.0 TeV

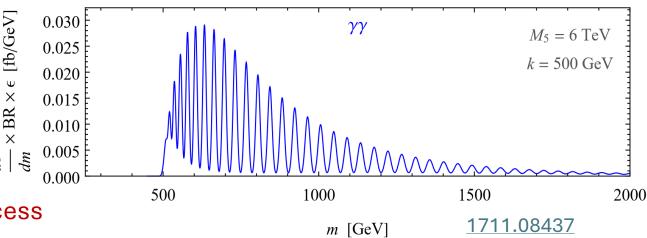


Periodic Signals (Clockwork)

 Clockwork mechanism can generate large hierarchies with only O(1) couplings and N fields <u>1610.07962</u>

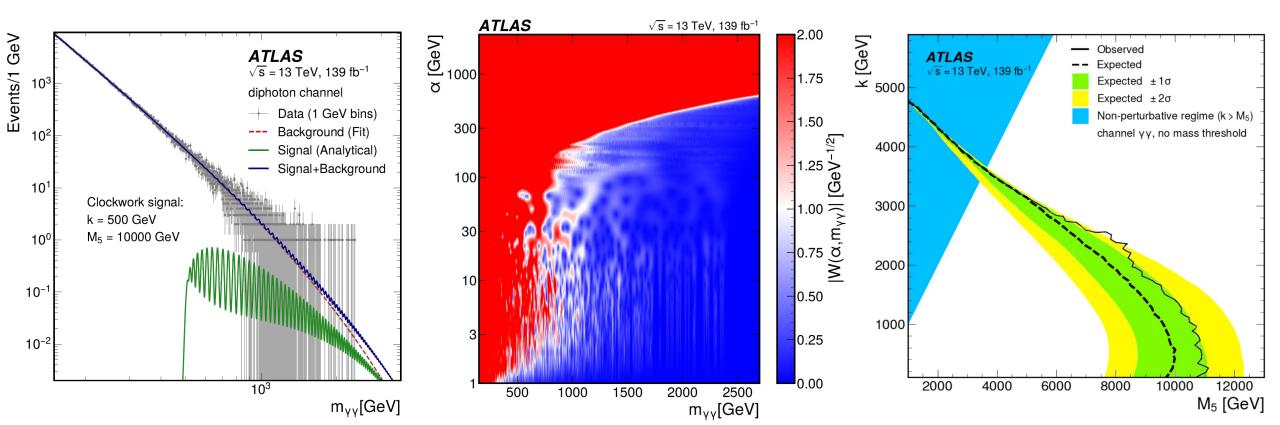


- This concept offers **solutions to hierarchy problems**, such as Higgs boson mass naturalness, why gravity is so weak, or why dark matter is cosmologically stable
- The clockwork gravity model assumes extra dimensions and predicts a narrowly-spaced spectrum of resonances in mass, such as towers of Kaluza-Klein (KK) gravitons → periodic signal
- Mass spacing between signals is a few percent at the onset an falls below 1% at high masses
- **Parameters:** k onset of the KK graviton spectrum, M₅ five-dimensional reduced Planck mass
- Signal cross section roughly scales as σ ~ M_5^{-3}
- Resolution effects can wash out the periodic structure and result in a broad signal shape especially at high masses
- In the continuum limit, the KK spacing is so tight [¬] to that it cannot be resolved anymore → one broad excess



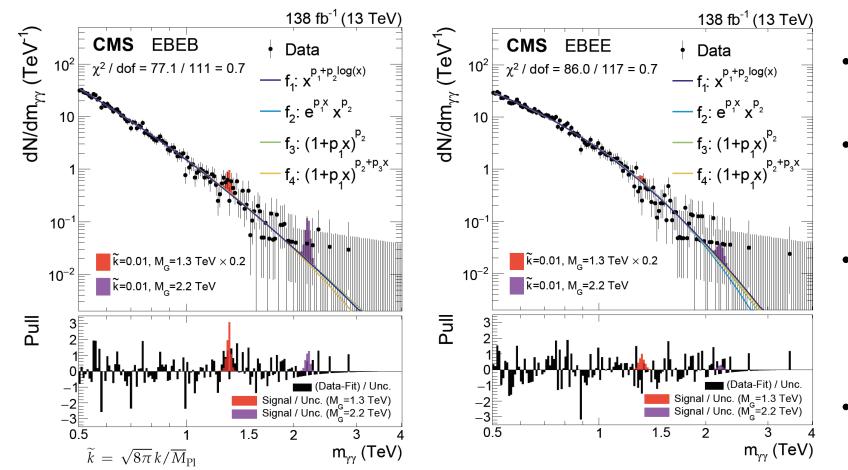
Periodic Signals (Clockwork)

- Search conducted in the invariant mass spectra of **di-photon and di-electron final states**
- Continuous wavelet transform (CWT) used to transform the mass spectra into scalograms ("image") of mass vs. frequency. Scaling parameter α either dilates or compresses the signal.
 It is inversely proportional to the frequency (i.e. if α is large, then the signal is "stretched").
- **Convolutional neural networks** or **autoencoders** used to search for anomalies ("islands") in the scalogram, but no signal detected in the dataset



Broad Resonances in Diphoton 2405.09320

- Arkani-Hamed, Dimopoulos, Dvali (ADD) model predicts n compactified extra dimensions. Kaluza-Klein modes of the graviton are tightly spaced and result in a broad, non-resonant excess. Interference with the background is considered.
- In the **continuum limit of the clockwork framework**, an infinite tower of very narrow KK graviton modes also leads to a **continuous excess** in the diphoton mass spectrum. No interference with the background.



- Diphoton trigger p_T > 60-70 GeV. Offline p_T cut > 125 GeV
- Two categories:
 - EBEB (both γ in the barrel)
 - EBEE (one barrel, one endcap)
- Background model:
 - Sherpa MC reweighted to NNLO
 - Data-driven estimate of misidentified jets
- Fit to **binned m_{γγ} distribution**

Broad Resonances in Diphoton 2405.09320

ADD model:

Lower limits on the mass scale M_s [TeV] for three theory conventions:

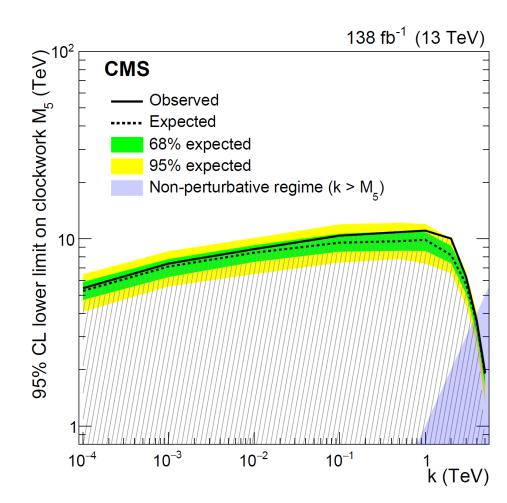
- GRW: Giudice-Rattazzi-Wells <u>hep-ph/9811291</u>
- Hewett: <u>hep-ph/9811356</u>
- HLZ: Han-Lykken-Zhang <u>hep-ph/9811350</u>

Signal:	GRW	Hewett		HLZ				
		negative	positive	$n_{\rm ED} = 3$	$n_{\rm ED} = 4$	$n_{\rm ED} = 5$	$n_{\rm ED} = 6$	$n_{\rm ED} = 7$
Expected:	$8.7\substack{+0.7 \\ -0.6}$	$7.3^{+0.3}_{-0.3}$	$7.8\substack{+0.6 \\ -0.5}$	$10.3\substack{+0.8 \\ -0.7}$	$8.7\substack{+0.7 \\ -0.6}$	$7.9\substack{+0.6 \\ -0.5}$	$7.3\substack{+0.6 \\ -0.5}$	$6.9^{+0.6}_{-0.5}$
Observed:	9.3	7.1	8.3	11.1	9.3	8.4	7.8	7.4

M_s is the ultraviolet cutoff parameter for the virtual graviton exchange

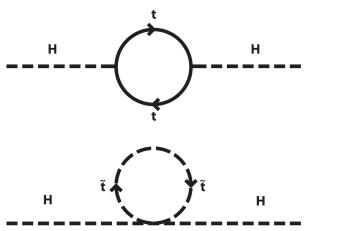
Continous clockwork gravity model:

Values of M_5 excluded for k values in the range of 0.2 – 2000 GeV. Strongest exclusion of $M_5 < 11$ TeV for k=1 TeV

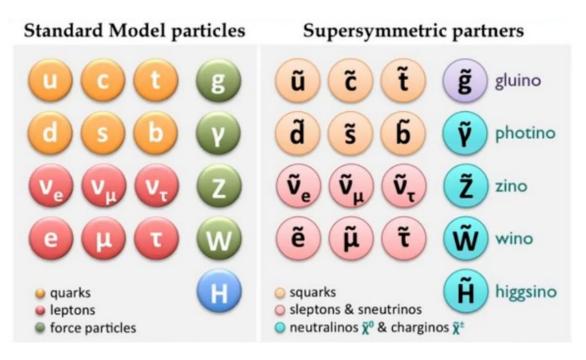


Supersymmetry

- Fundamental symmetry between fermions (half spin) and bosons (integer spin)
- Can stabilize the Higgs mass up to the Planck scale
- Provides dark matter candidates
- Gauge couplings unify at high energies
- Preserves baryon asymmetry



Sfermion loops cancel divergent corrections to the Higgs boson mass



R-parity:

 $P_{R} = (-1)^{3(B-L)+2s}$

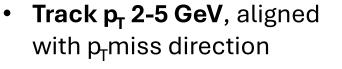
+1 for SM, -1 for SUSY partners

In R-violation scenarios the LSP is not stable

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Compressed Higgsinos with mildly-displaced tracks 2404.01996 ^{13/22}

- Higgsinos are the SUSY partners to the Higgs bosons. Mass eigenstates: $\tilde{\chi}_1^{\pm}$, $\tilde{\chi}_2^0$, $\tilde{\chi}_1^0$
 - Higgsino masses are connected to the EWSB, favored to be at EW scale @(100 GeV) even if SUSY mass scale is very high (natural SUSY)
 - Small mass-splitting $\Delta m(\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^0) \approx 0.3-1$ GeV
 - Very compressed topology, difficult! Strongest bounds came from <u>LEP</u>, up to now.
 - $\widetilde{\chi}_1^{\pm}$ flight length is $\mathcal{O}(0.1-1 \text{ mm})$.
 - Decays dominantly to charged pions with a **mildly-displaced low-p_T track**.

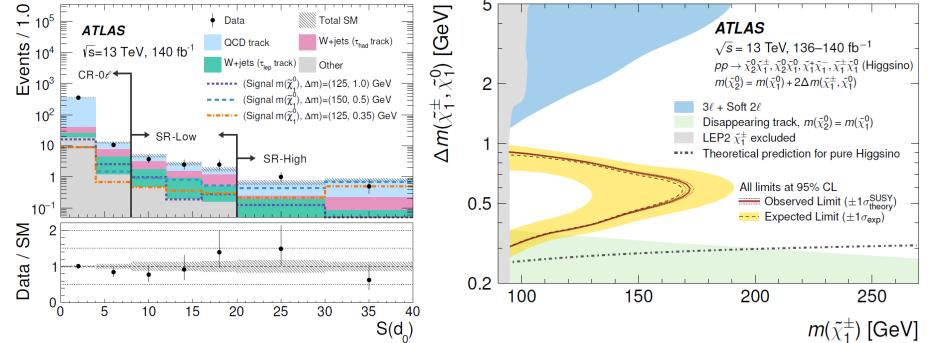


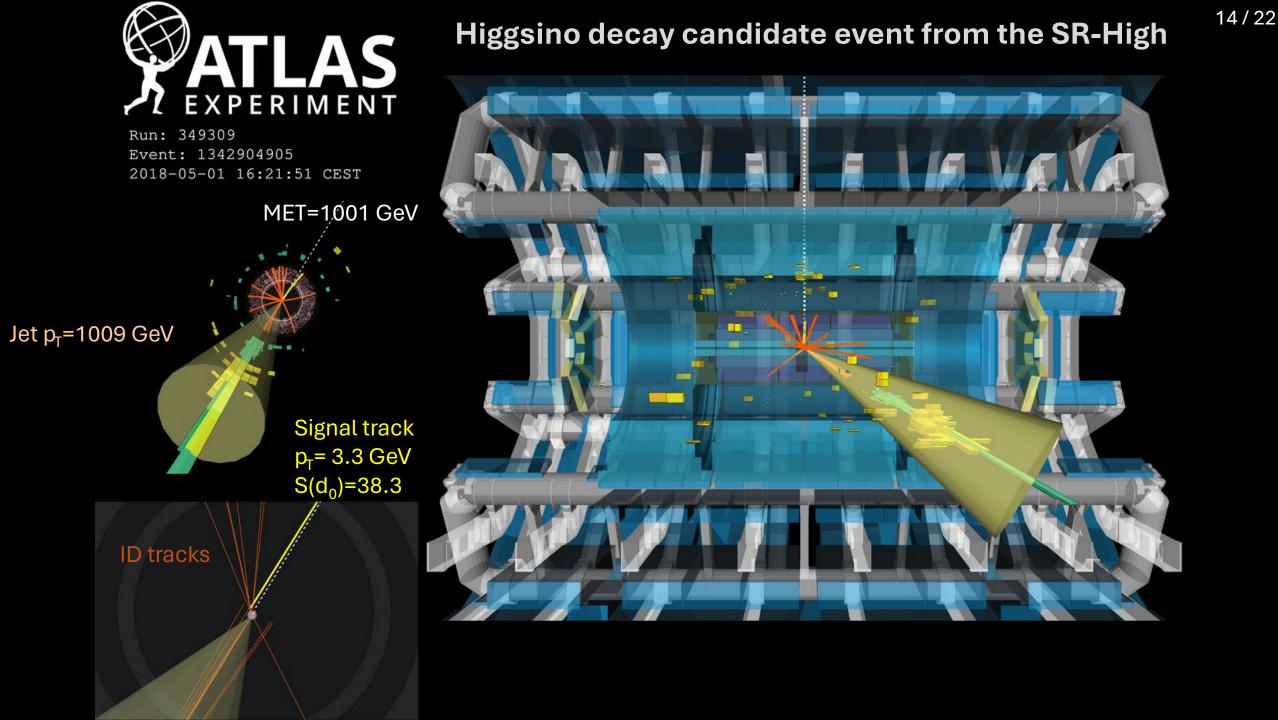
jet

- Must have a signal in first layer of the inner detector
- High p_T jet (> 250 GeV)
- MET > 600 GeV

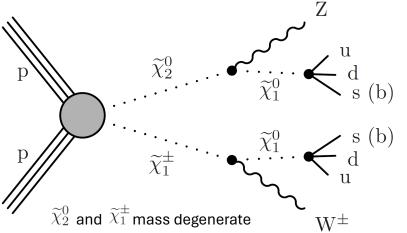
p

Final discriminant:
 S(d0), transverse impact
 parameter significance



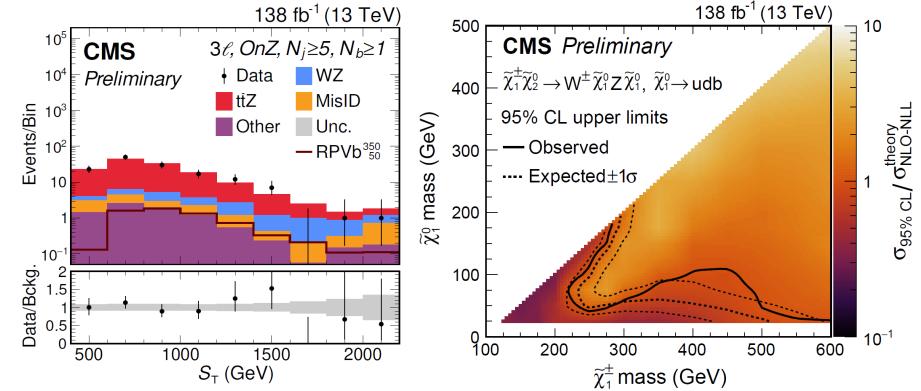


Hadronic R-parity violating SUSY decays CMS-PAS-SUS-23-015



- Lepton pair consistent with Z mass ("OnZ")
- Effective mass S_T:
 Scalar p_T sum of charged leptons, jets, and p_Tmiss
 → No combinatorics
- Wide excess expected
- Related RPV-SUSY search by ATLAS: 2401.16333

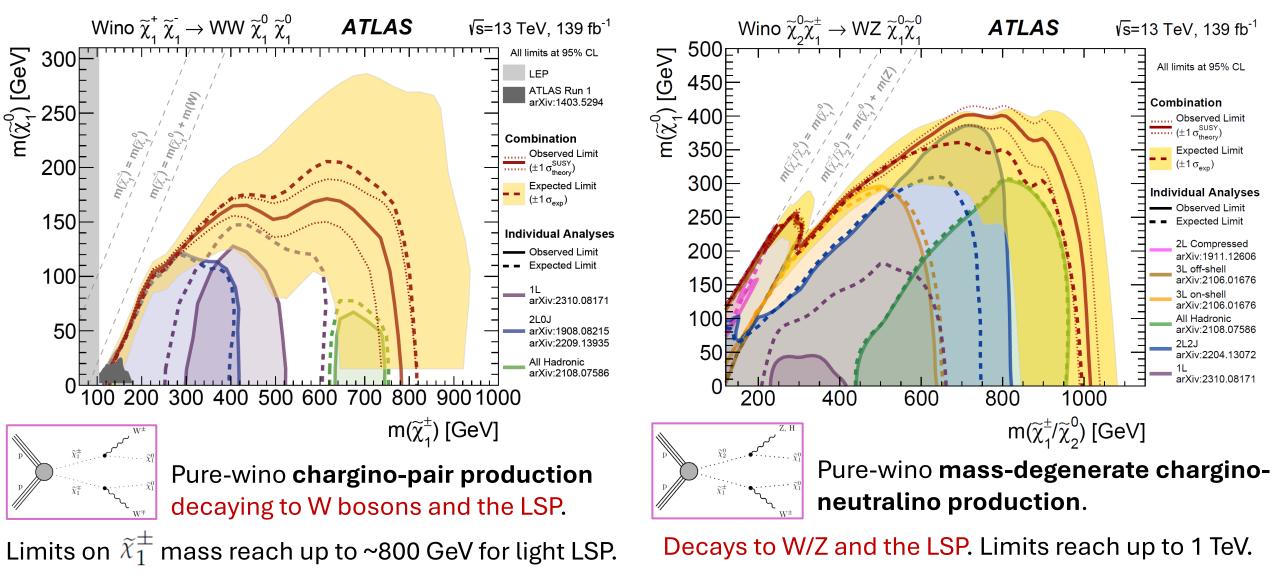
- Production of chargino-neutralino pair through electroweak process
 - Cascade decays to $\widetilde{\chi}_1^0$ and a vector boson (Z or W), followed by hadronic R-parity violating prompt decays of the LSP
 - $W \rightarrow lv, Z \rightarrow ll. LSP \rightarrow uds \text{ or } LSP \rightarrow udb \text{ (shown here)}$
 - Final state consists of **3 leptons and up to 6 jets**
 - Control regions with 4 or 5 leptons enriched in bkg



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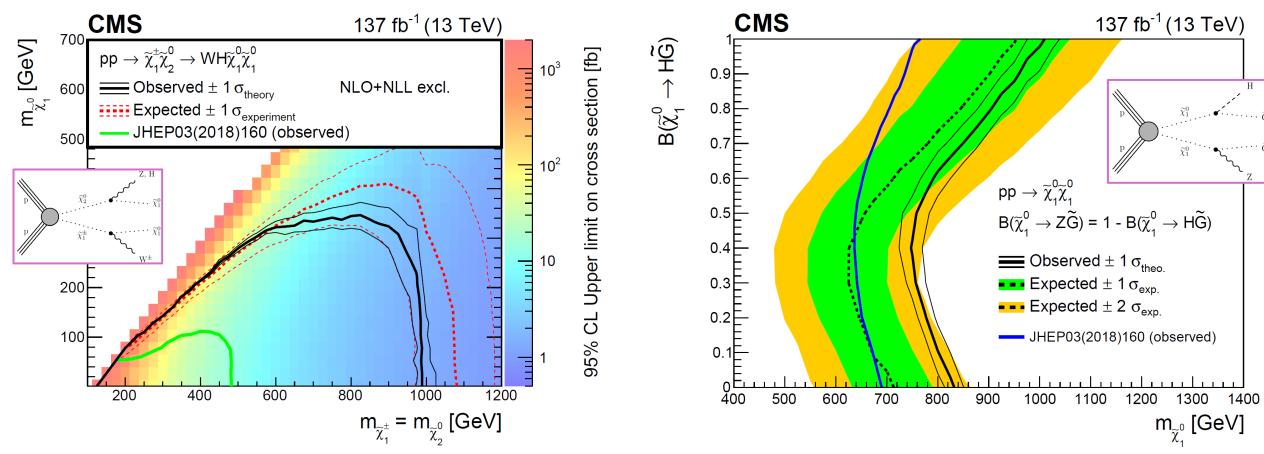
ATLAS combination of charginos and neutralinos 2402.08347

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Combinations extend sensitivity to SUSY production up to 100 GeV in (N)LSP masses, sensitivity to SUSY production cross-sections increased by up to 40%

CMS combination of charginos and neutralinos 2402.01888



Wino-bino model: **Mass-degenerate charginoneutralino production**, decays to W, H and the LSP.

1σ excess driven by Hadronic WW/WZ/WH search <u>2205.09597</u>

For LSP mass of 50 GeV, $\widetilde{\chi}_1^\pm$ below 990 GeV excluded.

Gauge mediated SUSY breaking: **Neutralino-pair production** with decays to H or Z and Gravitino

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Exclusion ranges between 750 GeV and 1 TeV depending on $\mathcal{B}(\widetilde{\chi}^0_1 \to H\widetilde{G}\,)$

Quantum Black Holes

- Predicted in **theories with low scale of quantum gravity** M_D (order of **1-10 TeV**)
- Models considered:
 - Large extra dimensions (Arkani-Hamed-Dimopoulos-Dvali, ADD)
 - Warped extra dimensions (Randall-Sundrum, RS1)

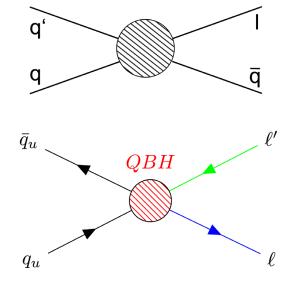
- Planck $M \sim M_{pl}$ hidden $v_{eff} = \exp(-ky) v$ U
 - Randall-Sundrum
- Global symmetries such as baryon or lepton number may not be conserved in strong-gravity interactions
- QBH production in pp collisions via **2-to-2 scattering** , for example:

 $uu \to \bar{d}\ell^+, \quad ud \to \bar{u}\ell^+, \quad \bar{d}\bar{d} \to d\ell^+$

• QBH decay to two-particle final states with large branching fraction (51-74% in ADD or RS1).

In contrary, semi-classical BH decay into multiparticle final states via Hawking radiation.

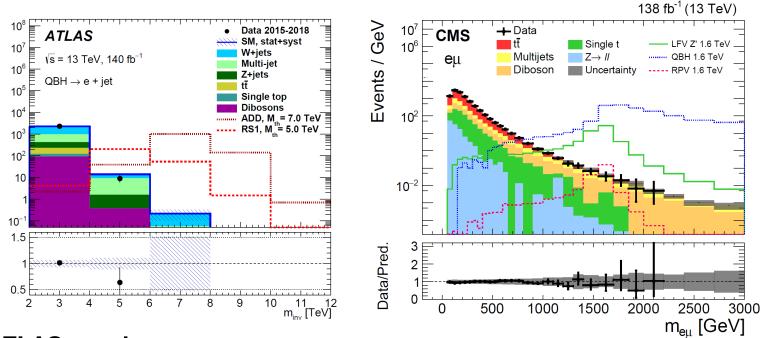
QBH



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Quantum Black Holes

<u>2307.14967</u> <u>2205.06709</u>



ATLAS result:

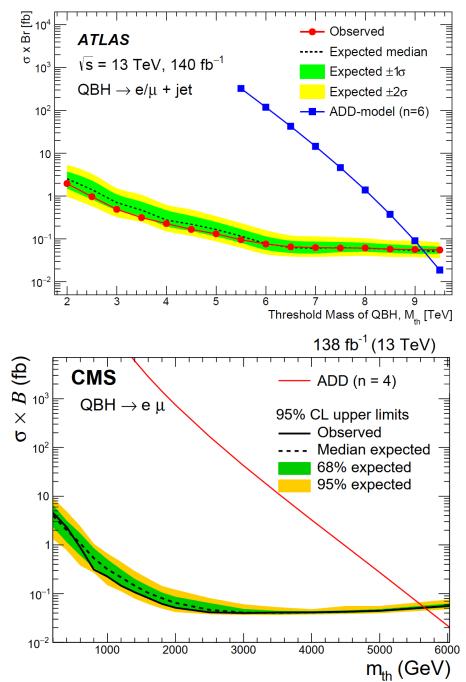
Events / 2.0 Te^v

Data / SM

- electron+jet and muon+jet final states
- Fit to SR and three CRs (W, Z, top) using m_{inv} of lepton and jet
- QBH in ADD model with n=6 excluded for M_{th}<9.2 TeV

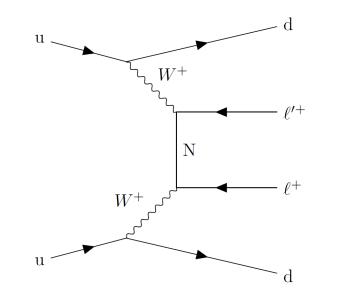
CMS result:

- $e\mu$, $e\tau$ and $\mu\tau$ channels, separate limits obtained
- Jets faking leptons is estimated from data in control samples
- Best channel: $e\mu$, QBH in ADD with n=4 excluded for $M_{th} < 5.6$ TeV



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Heavy Majorana Neutrinos

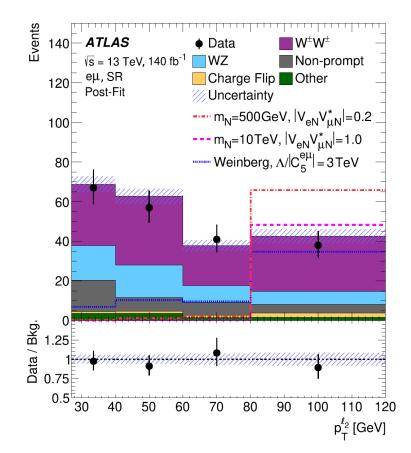


- VBF same-sign W boson scattering
- Final state: two leptons + two jets

- Neutrinos masses implied by the observation of neutrino oscillations
- Majorana nature means that these neutrinos are their own anti-particles
- Heavy leptons and extended scalar sectors present in Seesaw, Left-Right Symmetric or GUT models
- Phenomenological Type-I Seesaw model: Introduces a new heavy Majorana neutrino N, that generates the small neutrino masses m_v ~ @(vev²/m_N) (vev is the vacuum expectation value 246 GeV, m_N is the mass of the heavy Majorana neutrino)
- Heavy neutrino mass-mixes with SM neutrinos, **mixing parameters V**_{eN} and V_{μ N}
- Cross section ~ $|V_{eN} V_{\mu N}|^2$

Heavy Majorana Neutrinos 2403.15016

- Two channels: ee or eµ (+ two jets), that are then statistically combined
- Jets have a large rapidity gap ($|\Delta y|$ >2) and large invariant mass (m_{ii} > 500 GeV)
- Most discriminating variable: p_{T} of the subleading lepton
- WW and WZ control region for background estimation

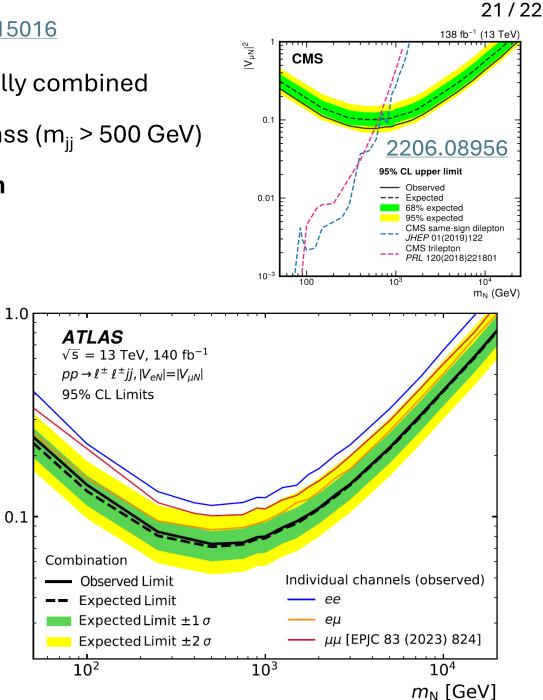


Limits set on the mixing parameter value as a function of the neutrino mass within 50 GeV to 20 TeV

1.0

 $|V_{\ell N}|^2$

No excess found



Summary

- Diverse set of searches presented with the full Run 2 dataset from ATLAS and CMS, non-resonant searches complementary to resonance and other searches.
- No significant excess found, stringent limits in various models set:
 - Vector LQ decays to μ +b for λ =1, β =1 and κ =1 excluded for masses up to 2.6 TeV
 - Scalar LQ excluded for masses below 620 GeV, assuming λ_{LQ} = 1.5 in cLFV interactions
 - M₅ up to 11 TeV excluded for k=1 TeV in extra-dimensional clockwork models
 - Natural SUSY Higgsinos with small mass-splitting excluded up to 170 GeV
 - Chargino/neutralino masses up to 1 TeV excluded in SUSY combinations
 - QBH in ADD model with n_D =4 (6) excluded for masses up to 5.6 (9.2) TeV
 - Neutrino mass-mixing parameter values excluded for heavy neutrinos up to 20 TeV
 - Searches with Run 3 data in progress! Some expected improvements:
 - Larger luminosity increases sensitivity for many searches, eg. EW SUSY, high-mass tails
 - Cross sections for TeV-scale processes profit from increased beam energy (e.g. QBH cross section at 9 TeV doubles, strong SUSY production also gains)

Backup

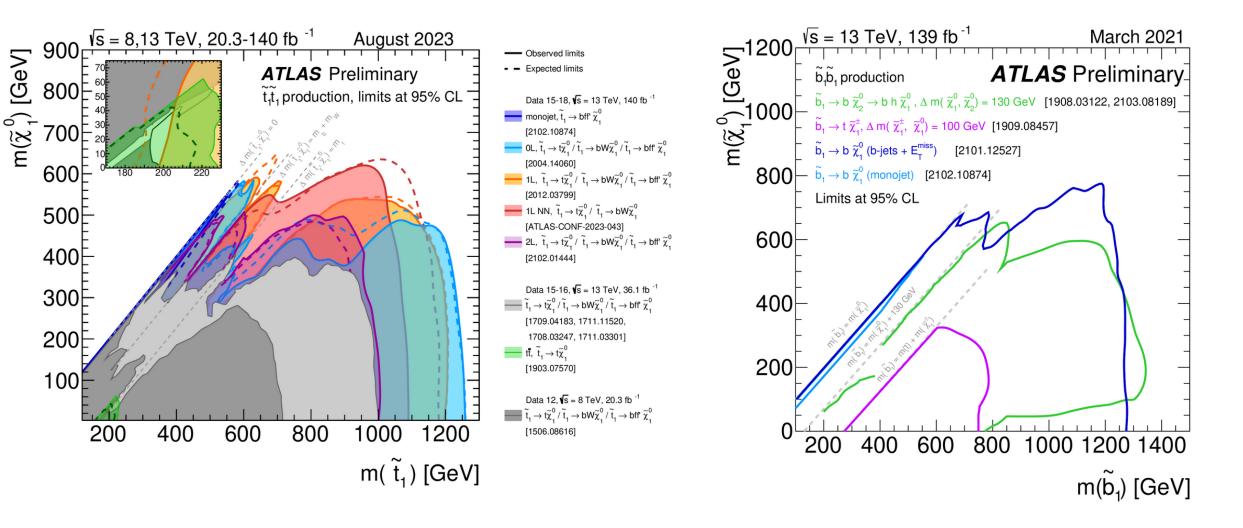
ATLAS Preliminary $\sqrt{s} = 13$ TeV

ATLAS SUSY Searches* - 95% CL Lower Limits August 2023

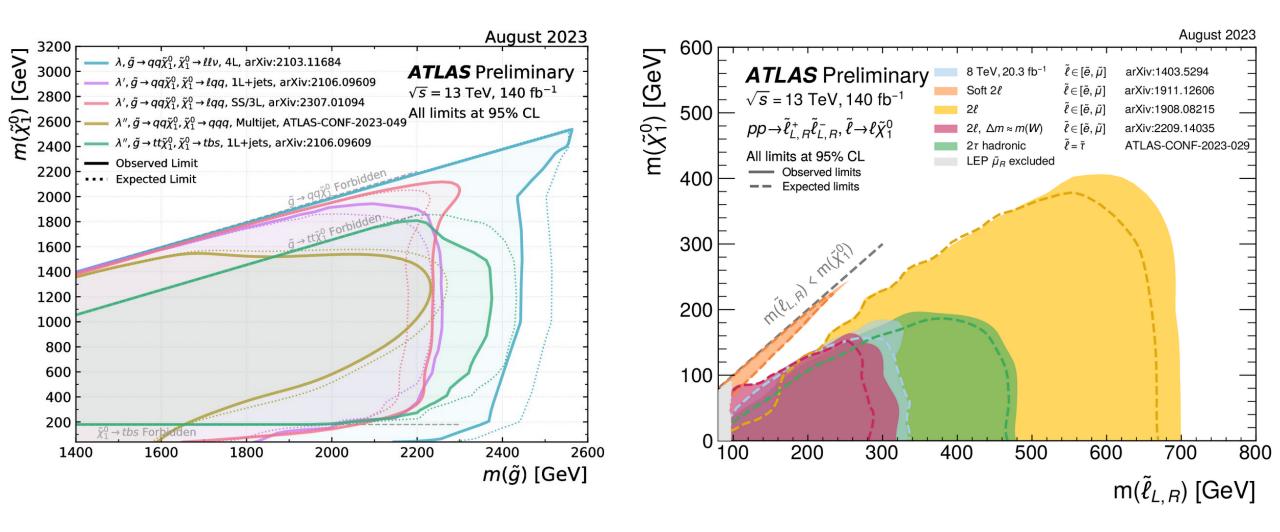
Model	Signature ∫	<i>L dt</i> [fb	¹] Mass limit		$\sqrt{s} = 13$ lev Reference
$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{k}_1^0$	0 e, μ 2-6 jets E_T^{miss} mono-jet 1-3 jets E_T^{miss}	140 140	<i>q</i>	0 1.85 m(ℓ ₁ ⁰)<400 GeV m(∂)-m(ℓ ₁ ⁰)=5 GeV	2010.14293 2102.10874
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{\chi}_1^0$	$0 e, \mu$ 2-6 jets E_T^{miss}	140	R Forbidde	2.3 m(\tilde{k}_{1}^{0})=0 GeV	2010.14293 2010.14293
$\begin{array}{c} \mathbf{x} \\ $	1 e, μ 2-6 jets $ee, \mu\mu$ 2 jets E_r^{miss}	140 140	a 1000000	2.2 m(ℓ ₁ ⁰)<600 GeV 2.2 m(ℓ ₁ ⁰)<700 GeV	2101.01629 2204.13072
$\tilde{g}_{\tilde{g}}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_{1}^{0}$	$\begin{array}{llllllllllllllllllllllllllllllllllll$	140 140	5 - 56 - 56 - 56 - 56 - 56 - 56 - 56 -	1.97 m(\tilde{v}_{1}^{0}) <600 GeV 1.15 m(\tilde{v}_{1}^{0})=200 GeV	2008.06032 2307.01094
$\tilde{\boldsymbol{g}}$ $\tilde{\boldsymbol{g}}, \tilde{\boldsymbol{g}} \rightarrow t \tilde{\boldsymbol{\mathcal{X}}}_1^0$	$\begin{array}{ccc} \text{0-1} \ e,\mu & \text{3} \ b & E_T^{\text{miss}} \\ \text{SS} \ e,\mu & \text{6 jets} \end{array}$	140 140	as, 23,	2.45 m(k ⁰ ₁)<500 GeV 1.25 m(k ¹ ₁)=300 GeV	2211.08028 1909.08457
$\tilde{b}_1 \tilde{b}_1$	$0 e, \mu$ $2 b$ E_T^{miss}	140	$b_1 \\ b_1 \\ 0.68$	1.255 m(\tilde{k}_{1}^{0})<400 GeV 10 GeV< $\Delta m(b, k_{1}^{0})$ >20 GeV	2101.12527 2101.12527
sy to $b_1b_1, b_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$ $\tilde{\tau}_1\tilde{\tau}_1, \tilde{\tau}_1 \rightarrow \tilde{\chi}_1^0$ $\tilde{\tau}_1\tilde{\tau}_1, \tilde{\tau}_1 \rightarrow \tilde{\chi}_1^0$	$\begin{array}{ccc} 0 \ e, \mu & 6 \ b & E_T^{miss} \\ 2 \ \tau & 2 \ b & E_T^{miss} \end{array}$	140 140	bi Forbidden bi 0.13-0.85	0.23-1.35 $\Delta m(\hat{\xi}_2^0, \hat{\xi}_1^0) = 130 \text{ GeV}, m(\hat{\xi}_1^0) = 100 \text{ GeV} \\ \Delta m(\hat{\xi}_2^0, \hat{\xi}_1^0) = 130 \text{ GeV}, m(\hat{\xi}_1^0) = 0 \text{ GeV} \end{cases}$	1908.03122 2103.08189
	0-1 $e, \mu \ge 1$ jet E_T^{miss} 1 e, μ 3 jets/1 b E_T^{miss}	140 140	ři ři Forbidden 1	1.25 m($\tilde{\chi}_{1}^{0}$)=1 GeV 05 m($\tilde{\chi}_{1}^{0}$)=500 GeV	2004.14060, 2012.03799 2012.03799, ATLAS-CONF-2023-043
\tilde{b} $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b v, \tilde{\tau}_1 \rightarrow \tau \tilde{G}$	1-2 τ 2 jets/1 b E_T^{miss}	140 36.1	ī Forbidden	1.4 m(ř ₁)=800 GeV	2108.07665 1805.01649
0 0	$0 e, \mu$ mono-jet E_T^{fniss}	140	ī ₁ 0.55	$m(\tilde{\xi}_{0}^{0})=0 \text{ GeV}$ $m(\tilde{r}_{1},\tilde{z})-m(\tilde{\xi}_{1}^{0})=5 \text{ GeV}$	2102.10874
$ \begin{split} \tilde{r}_1 \tilde{r}_1, \tilde{r}_1 \rightarrow t \tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h \tilde{\chi}_1^0 \\ \tilde{r}_2 \tilde{r}_2, \tilde{r}_2 \rightarrow \tilde{r}_1 + Z \end{split} $	$\begin{array}{ccc} \textbf{1-2} \ e, \mu & \textbf{1-4} \ b & E_T^{\text{miss}} \\ \textbf{3} \ e, \mu & \textbf{1} \ b & E_T^{\text{miss}} \end{array}$	140 140	ī1 0.06 ī2 Forbidden 0.86	7-1.18 m(ℓ ⁰ ₂)=500 GeV m(ℓ ⁰ ₁)=360 GeV, m(ℓ ₁)-m(ℓ ⁰ ₁)= 40 GeV	2006.05880 2006.05880
$\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via WZ	$\begin{array}{llllllllllllllllllllllllllllllllllll$	140 140	$\hat{\chi}_1^*/\tilde{\chi}_0^0$ 0.96 $\hat{\chi}_1^*/\tilde{\chi}_2^0$ 0.205	$m(\tilde{\chi}_1^0)=0$, wino-bino $m(\tilde{\chi}_1^1)\cdot m(\tilde{\chi}_1^0)=5$ GeV, wino-bino	2106.01676, 2108.07586 1911.12606
$\tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp}$ via <i>WW</i> $\tilde{\chi}_{1}^{\pm} \tilde{\chi}_{2}^{0}$ via <i>Wh</i>	$\begin{array}{ccc} 2 \ e, \mu & E_T^{\rm miss} \\ {\rm Multiple} \ \ell/{\rm jets} & E_T^{\rm miss} \\ 2 \ e, \mu & E_T^{\rm miss} \end{array}$	140 140	\tilde{x}_1^{\pm} 0.42 $\tilde{x}_1^{\pm}/\tilde{x}_2^{0}$ Forbidden 1	$m(\tilde{\chi}_1^0)=0$, wino-bino $m(\tilde{\chi}_1^0)=70$ GeV, wino-bino	1908.08215 2004.10894, 2108.07586
$\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{1}^{\mp}$ via $\tilde{\ell}_{r}/\tilde{\nu}$	$2 e, \mu$ E_T^{miss} 2τ E_T^{miss}	140 140	\tilde{X}_{1}^{\pm} 1.	$m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^{\pm})+m(\tilde{\chi}_1^{0}))$	1908.08215 ATLAS-CONF-2023-029
$\begin{array}{c} \overleftarrow{\boldsymbol{\Sigma}} \\ \overleftarrow{\boldsymbol{\Sigma}} \\ \overrightarrow{\boldsymbol{\Sigma}} \\ \overrightarrow{\boldsymbol{\Sigma}} \\ \overrightarrow{\boldsymbol{U}} \overrightarrow$	$\begin{array}{ccc} 2 \tau & E_T^{\text{miss}} \\ 2 e, \mu & 0 \text{ jets } E_T^{\text{miss}} \\ ee, \mu \mu & \geq 1 \text{ jet } E_T^{\text{miss}} \end{array}$	140 140 140	τ [τ̄ _R , τ̄ _{R,L}] 0.34 0.48 ζ 0.26 0.7	m(ℓ ₁ ⁰)=0 m(ℓ ₁ ⁰)=0 m(ℓ)-m(ℓ ₁ ⁰)=10 GeV	1908.08215 1911.12606
$\hat{H}\hat{H}, \hat{H} \rightarrow h\hat{G}/Z\hat{G}$	$0 e, \mu \ge 3 b E^{\text{miss}}_{\text{rest}}$	140 140	H 0.55	$\begin{array}{c} BR(\hat{k}_{d}^{0} \to h\tilde{G}) = 1 \\ BR(\hat{k}_{d}^{0} \to Z\tilde{G}) = 1 \end{array}$	To appear 2103.11684
	$e, \mu \ge 2$ large jets E_T $e, \mu \ge 2$ large jets E_T $e, \mu \ge 2$ jets E_T	140 140	H 0.45-0.93 H 0.77	$BR(\tilde{x}_1^0 \to Z\tilde{G})=1$ $BR(\tilde{x}_1^0 \to Z\tilde{G})=BR(\tilde{x}_1^0 \to h\tilde{G})=0.5$	2108.07586 2204.13072
Direct $\tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-}$ prod., long-lived $\tilde{\chi}_{1}^{\pm}$	Disapp. trk 1 jet E_T^{miss}	140		Pure Wino	2201.02472
	-			Pure higgsino	2201.02472
Stable \tilde{g} R-hadron Stable \tilde{g} R-hadron Hetastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$ $\tilde{\ell}\tilde{\ell}, \tilde{\ell} \rightarrow \ell\tilde{G}$	pixel dE/dx E_T^{miss} pixel dE/dx E_T^{miss} Displ. lep E_T^{miss}	140 140	\tilde{g} \tilde{g} [$\tau(\tilde{g}) = 10 \text{ ns}$]	2.05 2.2 m(\tilde{x}_1^0)=100 GeV	2205.06013 2205.06013
	Displ. lep E_T^{miss}	140	ē,μ 0.34	$\tau(\tilde{\ell}) = 0.1 \text{ ns}$ $\tau(\tilde{\ell}) = 0.1 \text{ ns}$	2011.07812 2011.07812
	pixel dE/dx E_T^{miss}	140	τ 0.34 τ 0.36	$\tau(\tilde{\ell}) = 10 \text{ ns}$	2205.06013
$ \begin{split} \tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp} / \tilde{\chi}_{1}^{0} , \tilde{\chi}_{1}^{\pm} \rightarrow Z\ell \rightarrow \ell\ell\ell \\ \tilde{\chi}_{1}^{\pm} \tilde{\chi}_{1}^{\mp} / \tilde{\chi}_{2}^{0} \rightarrow WW / Z\ell\ell\ell\ell\nu\nu \end{split} $	$3 e, \mu$ $4 e, \mu$ 0 jets E_T^{miss}	140 140	$\hat{x}_{1}^{+}/\hat{x}_{0}^{0}$ [BR(Z _T)=1, BR(Z _e)=1] 0.625 1 $\hat{x}_{1}^{+}/\hat{x}_{0}^{+}$ [$\lambda_{133} \neq 0, \lambda_{12k} \neq 0$] 0.95	05 Pure Wino 1.55 m(𝔅 ⁰)=200 GeV	2011.10543 2103.11684
$\begin{array}{c} \tilde{\chi}_{1}^{1}\chi_{1}^{1}\chi_{2}^{0} & \tilde{\chi}_{1}^{0}\chi_{1}^{0}\chi_{1}^{0} \\ \tilde{\chi}_{1}^{0},\tilde{\chi}_{1}^{0}\chi_{1}^{0}\chi_{1}^{0} & ds \\ \tilde{\chi}_{1},\tilde{\tau}\rightarrow\tilde{\chi}_{1}^{0},\tilde{\chi}_{1}^{0}\rightarrow ds \\ \tilde{\chi}_{1},\tilde{\tau}\rightarrow\tilde{\chi}_{1}^{1},\tilde{\chi}_{1}^{\pm}\rightarrow bbs \end{array}$	≥8 jets	140	$\tilde{g} = [m(\tilde{x}_1^0) = 50 \text{ GeV}, 1250 \text{ GeV}]$	1.6 2.25 Large X''_112	To appear
$\begin{array}{c} \overbrace{ii}^{i}, i \rightarrow i \tilde{X}_{1}^{0}, \tilde{X}_{1}^{0} \rightarrow i b s \\ ii, i \rightarrow b \tilde{X}_{1}^{i}, \tilde{X}_{1}^{\pm} \rightarrow b b s \end{array}$	Multiple $\geq 4b$	36.1 140	i [J] ³ ₃₂₃ =2e-4, 1e-2] 0.55 1 i Forbidden 0.95	05 m($\tilde{\chi}_1^0$)=200 GeV, bino-like m($\tilde{\chi}_1^+$)=500 GeV	ATLAS-CONF-2018-003 2010.01015
$\overbrace{i_1}^{m} \overbrace{i_1 \to bs}^{m} i_1 \to bs$	2 jets + 2 b	36.7	<i>ī</i> ₁ [<i>qq, bs</i>] 0.42 0.61	mat/=300 GBV	1710.07171
$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e,μ 2 b 1 μ DV	36.1 136	$\frac{\tilde{t}_1}{\tilde{t}_1}$ [10-10< λ'_{23k} < 10-8, 30-10< λ'_{23k} < 30-9] 1.	0.4-1.45 BR($\tilde{i}_1 \rightarrow be/b\mu$)>20% BR($\tilde{i}_1 \rightarrow q\mu$)=100%, cos θ_i =1	1710.05544 2003.11956
$\tilde{\chi}_1^{\pm}/\tilde{\chi}_2^0/\tilde{\chi}_1^0, \tilde{\chi}_{1,2}^0 \rightarrow tbs, \tilde{\chi}_1^{+} \rightarrow bbs$	$1-2 e, \mu \ge 6$ jets	140	\bar{x}_1^0 0.2-0.32	Pure higgsino	2106.09609
*Only a selection of the available m phenomena is shown. Many of the	ass limits on new states or	1	0 ⁻¹	1 Mass scale [TeV]	

phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

SUSY Summary Plots

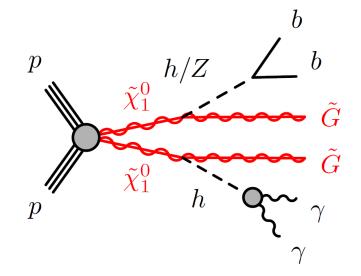


SUSY Summary Plots

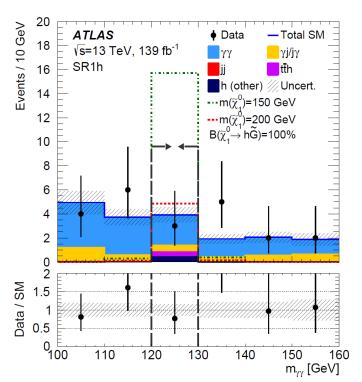


Pair-produced Higgsinos 2404.01996

•

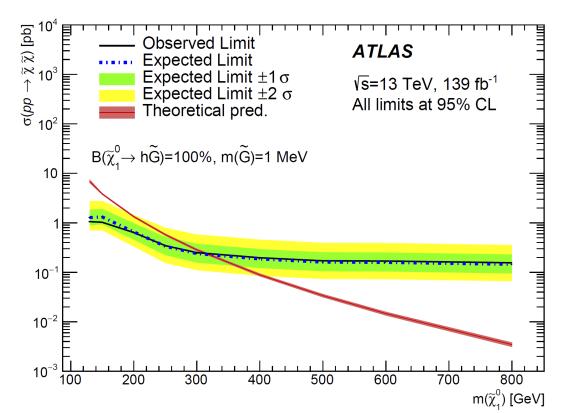


- Neutralinos $ilde{\chi}^0_1$ decay to a light Gravitino $ilde{G}$ and a Higgs or Z boson
- Neutralino masses at EW scale, even if SUSY mass scale is very high
- Final state: Missing energy, two photons and two b-jets
- Selection requires 120 GeV < $m_{\gamma\gamma}$ < 130 GeV



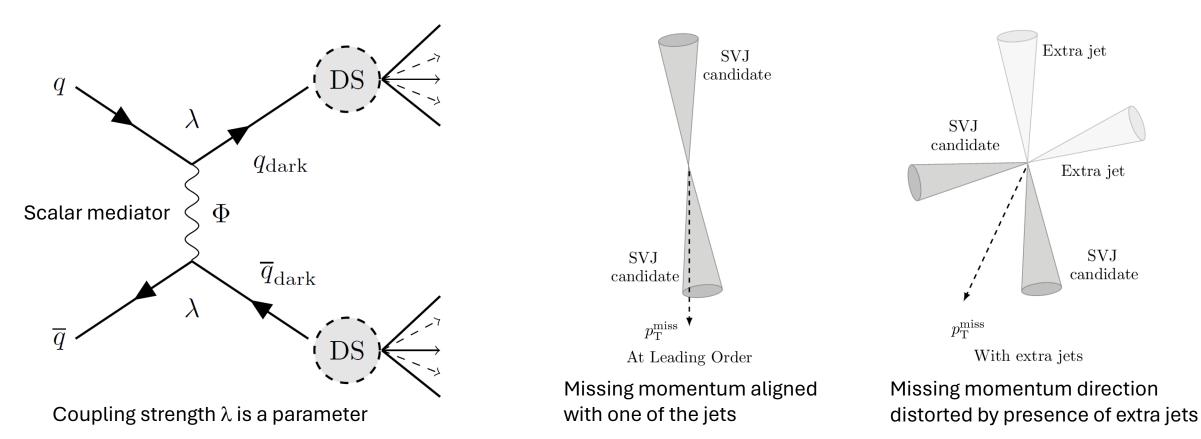
Three signal regions:

- <mark>SR1h:</mark> MET < 100 GeV 100 GeV < m_{bb} <140 GeV
- <mark>SR1Z:</mark> MET < 100 GeV 60 GeV < m_{bb} <100 GeV
- SR2 (targets heavy $\tilde{\chi}_1^0$): MET > 100 GeV 35 GeV < m_{bb} <145 GeV



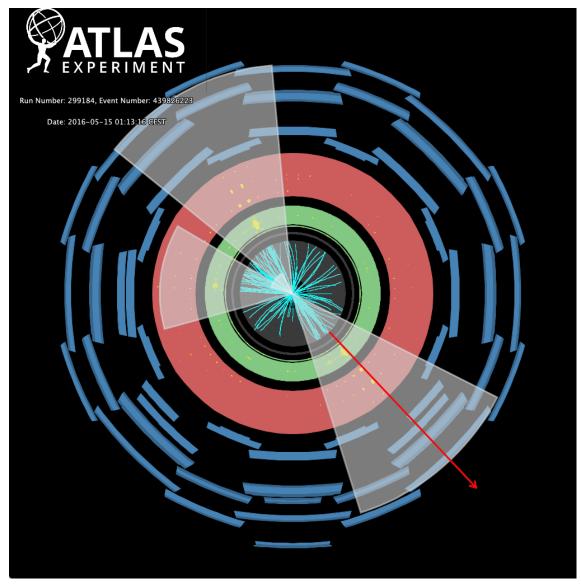
Semi-visible Jets

- Strongly-coupled dark sector featuring a dark shower (DS) including dark hadrons
- These dark hadrons may or may not be stable:
 - If all dark hadrons are stable \rightarrow missing energy signature.
 - If all dark hadrons decay to SM particles → multijet signature
- "Semi-visible" means the fraction of stable hadrons (R_{inv}) has intermediate values



Semi-visible Jets

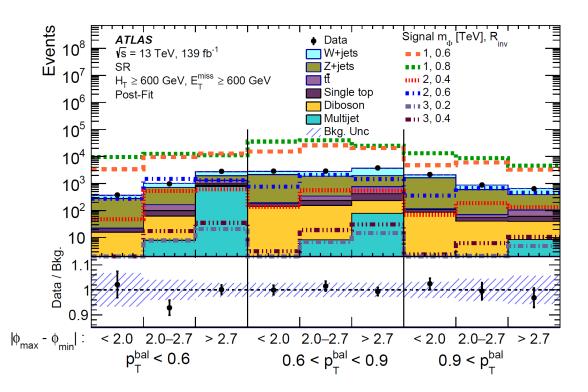
2305.18037



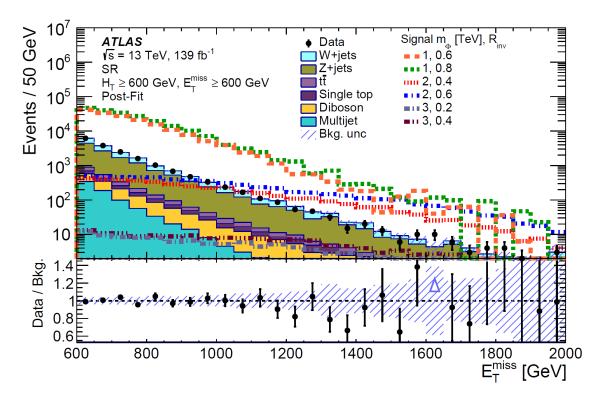
2305.18037

• Missing E_T trigger with online MET >70-110 GeV

- Events with at least two jets, leading jet $p_T > 250$ GeV, one jet close to the missing p_T direction
- Events with at leptons or at least two b-jets vetoed in the SR. Three control regions with leptons used for background estimation.
- In the SR: MET > 600 GeV, H_T > 600 GeV
- p_T balance $p_T^{\text{bal}} = \frac{|\vec{p_T}(j_1) + \vec{p_T}(j_2)|}{|\vec{p_T}(j_1)| + |\vec{p_T}(j_2)|}$ and azimuthal separation between jets for final discrimination and fit

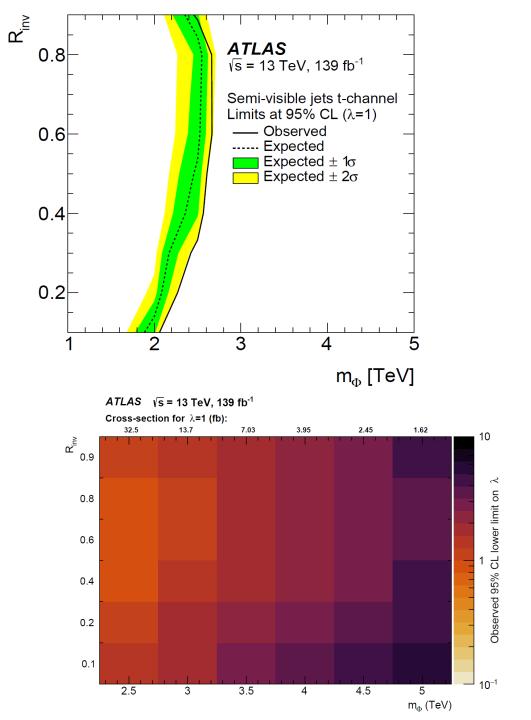


Semi-visible Jets



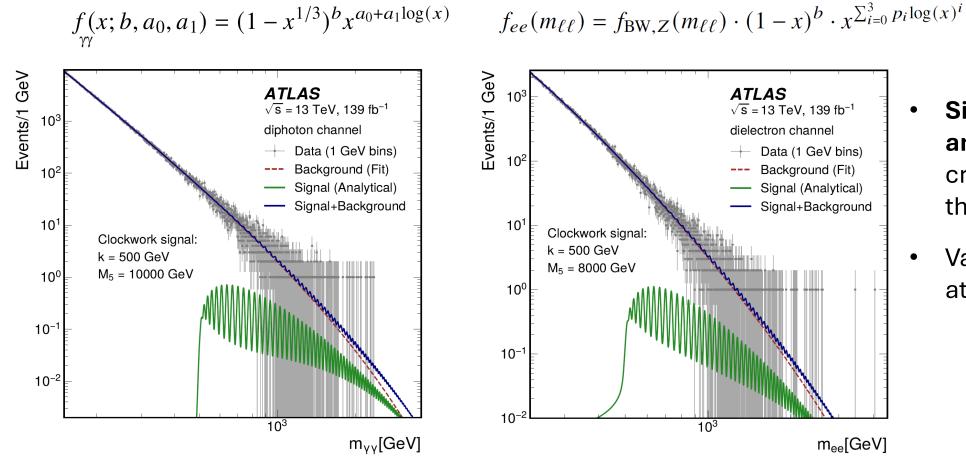
2305.18037

- No excess found
- Upper limits on the mediator mass between
 2.4 2.7 TeV, depending on R_{inv}
- Limits on coupling strength λ set (for m_{Φ} > 2.5 TeV)



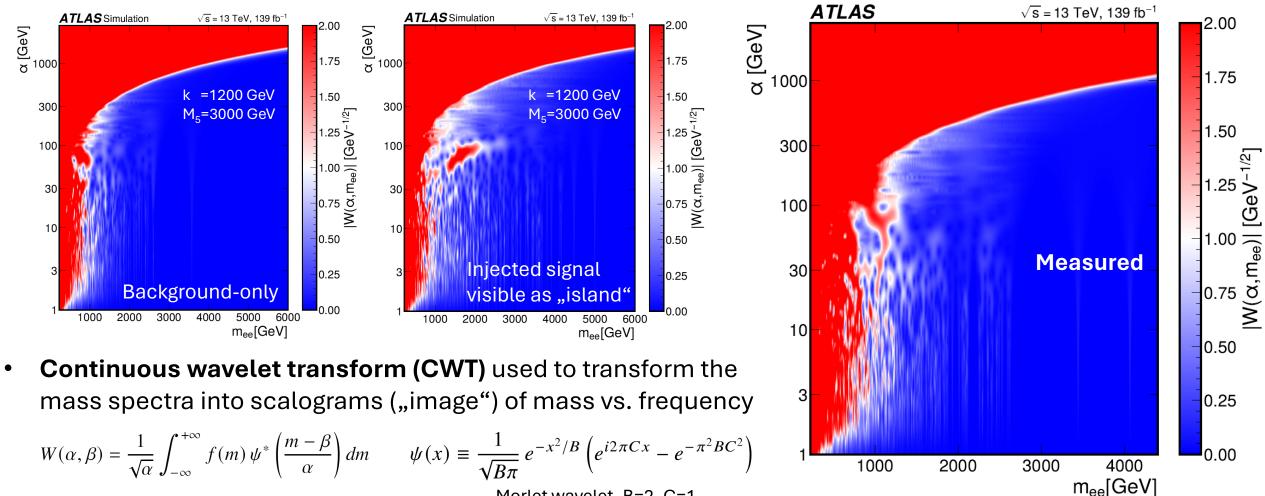
Periodic Signals (Clockwork): Analysis 2305.10894

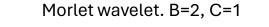
- Search conducted in **di-photon and di-electron final states**
- Di-photon: Tight ID, photon $E_T > 25$ GeV, $E_T/m_{\gamma\gamma} > 0.35$ (0.25), $m_{\gamma\gamma} > 150$ GeV
- Di-electron: Medium ID, $E_T > 30$ GeV, $m_{ee} > 225$ GeV
- Background modelled with **fit function** determined from simulated mass distributions:



- Signals also modelled with analytical functions created for any point in the k-M₅ plane
- Validated with MC samples at specific parameter points

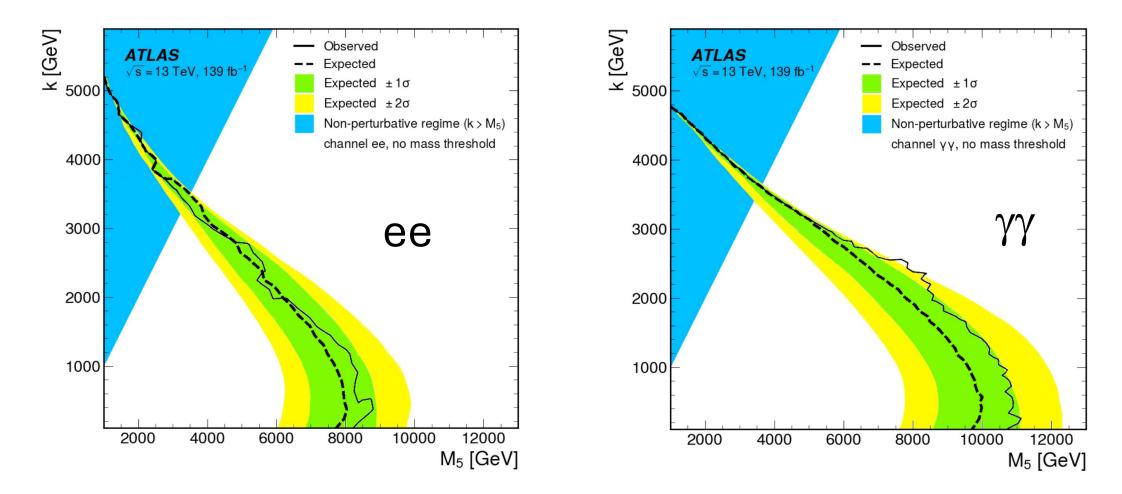
Periodic Signals (Clockwork): Analysis 2305.10894





- Scaling parameter α either dilates or compresses the signal. It is inversely proportional to the frequency (meaning if α is large, then the signal is "stretched")
- Convolutional neural networks or Autoencoder used to search for anomalies ("islands") in the scalogram

Periodic Signals (Clockwork): Results



2305.10894

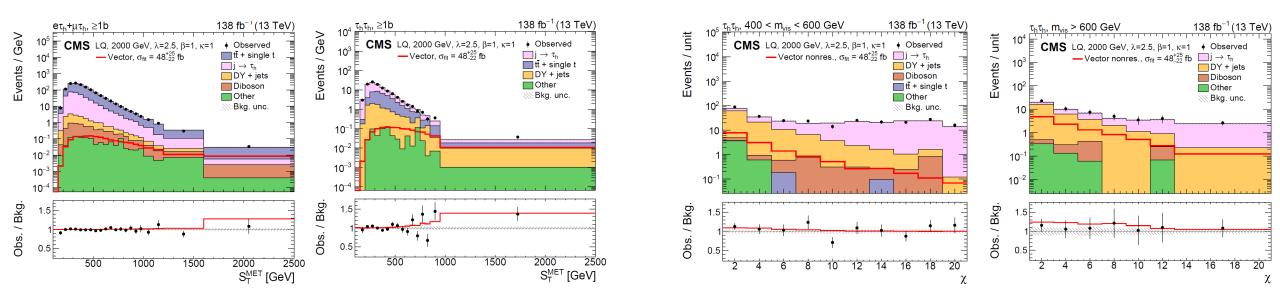
- No significant excess found, largest significance 1.5 σ in the dielectron channel for **model-independent** AE analysis
- Model-dependent limits in k-M₅ plane derived, obtained from the classifier NN analysis.
- Maximum excluded M_5 value is 11 TeV (8 TeV) for the $\gamma\gamma$ (ee) channel.

Leptoquarks decaying to τ + b: 2308.07826

- Produced singly resonant or nonresonant, and pair-produced. Non-resonant dominates cross section for large values of λ , scales with λ^4
- Analysis selects events with two leptonicaly or hadronicaly decaying taus ($e\tau_{had}$, $\mu\tau_{had}$, τ_{had} , $e\mu$, $\mu\mu$)
- Resonant: 0 or ≥1 b-jet, non-resonant: veto of high-pT jets
- Final discriminants: for non-resonant the angular separation of the taus $\chi = \exp(|\Delta \eta|)$, for resonant the scalar $p_T \text{ sum } S_T^{MET} = p_T^{-1} + p_T^{-2} + p_T^{-3} + p_T^{-miss}$

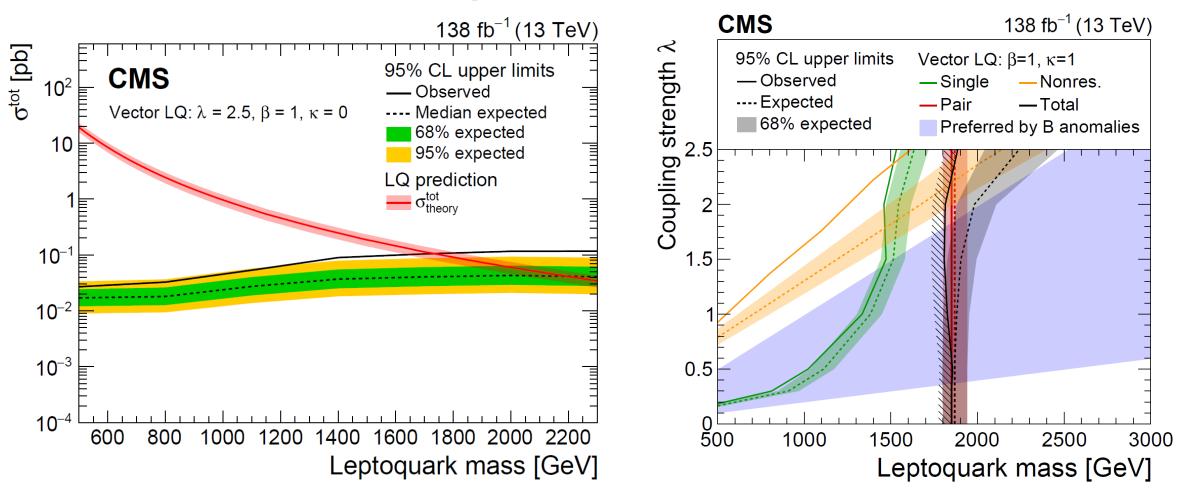
Resonant

Non-resonant



Leptoquarks decaying to τ + b: 2308.07826

(Result from 2023)



Vector LQ for λ =2.5, β =1 and κ =1 excluded for masses below 1.7 TeV

Broad excess found

2.8 σ local significance for 2 TeV LQ mass at λ =2.5

Pair-produced Higgsinos 2404.01996

