

# Heavy Neutral Leptons: Theoretical Overview

June 7<sup>th</sup>  
LHCP 2024  
Northeastern University  
Boston, USA



Juraj Klarić (*he/him*)



University of  
Zagreb

# The four Portals to Light New Physics



The four Horsemen (SM-apocalypse)  
A. Durer, MFA Boston

Four portals to Light New Physics:

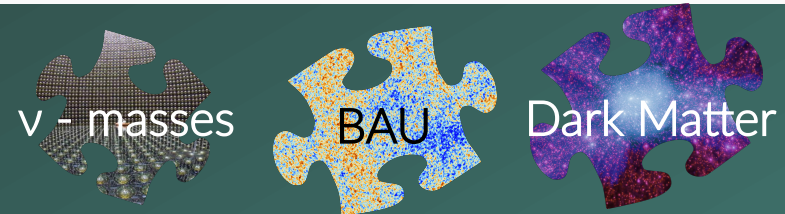
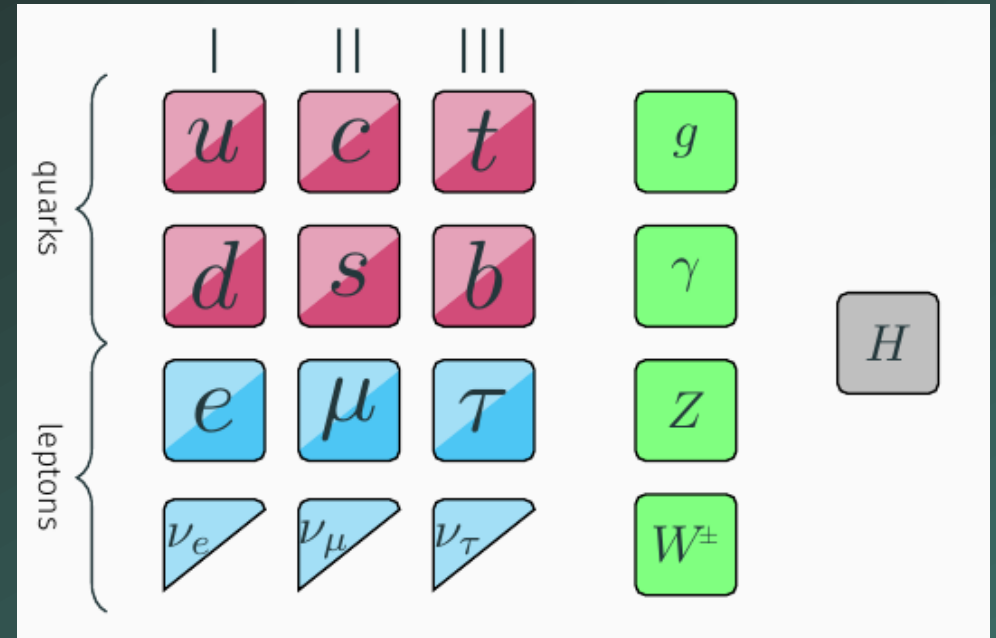
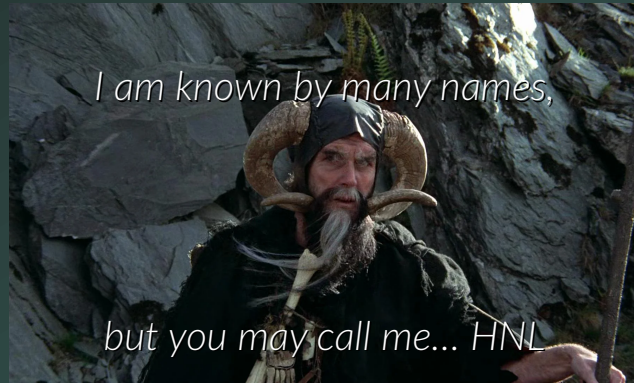
- Scalar (spin 0)
- Pseudo-scalar (spin 0)  
*Axions or  
Axion-Like-Particles (ALPs)*
- Vector (spin 1)  
*Dark Photon*
- Fermion (spin  $\frac{1}{2}$ )

*Heavy Neutral Leptons (HNLs)*

[See the Monday session for more light NP]

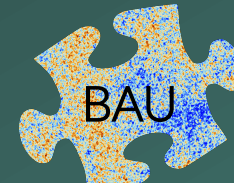
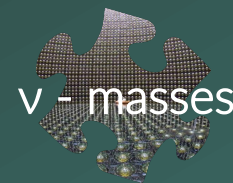
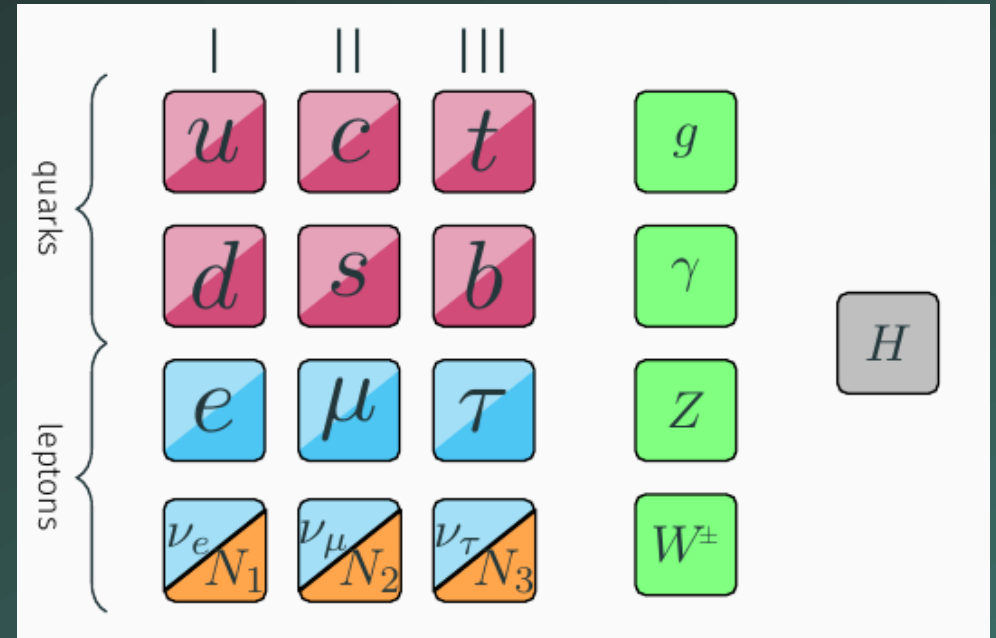
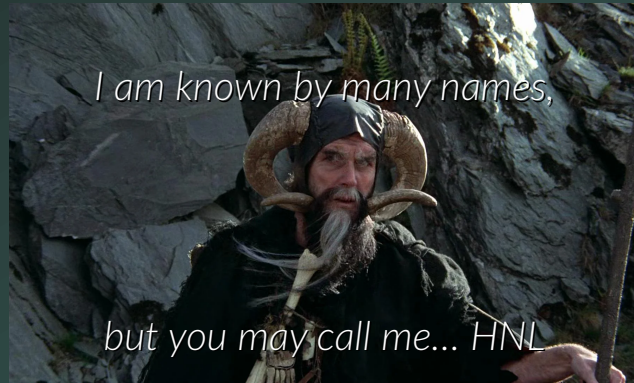
# Heavy Neutral Leptons

- Also known as:
  - Right-handed neutrinos
  - Heavy Majorana neutrinos
  - Heavy neutrinos



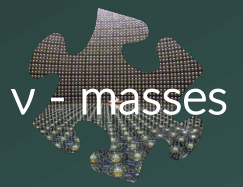
# Heavy Neutral Leptons

- Also known as:
  - Right-handed neutrinos
  - Heavy Majorana neutrinos
  - Heavy neutrinos



# The Seesaw Mechanism

## How *Heavy* are HNLs?



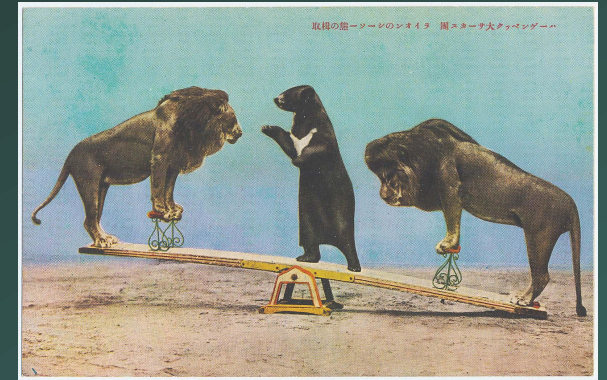
The Seesaw Lagrangian

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu}_L & \overline{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & 0 \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

$$m_\nu = m_D$$

Neutrino masses  
are extremely small

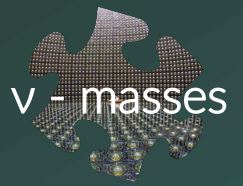
$$m_\nu \lesssim 0.8 \text{ eV}$$



Unknown Artist,  
Japanese, MFA Boston

# The Seesaw Mechanism

## How *Heavy* are HNLs?



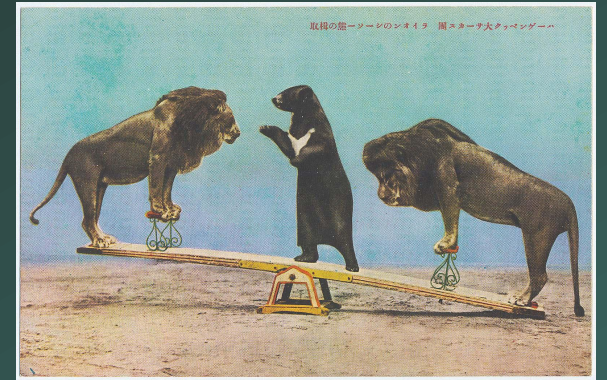
The Seesaw Lagrangian

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu}_L & \overline{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

$$m_\nu = -m_D M_M^{-1} m_D^T$$

Neutrino masses  
are extremely small

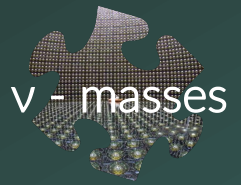
$$m_\nu \lesssim 0.8 \text{ eV}$$



Unknown Artist,  
Japanese, MFA Boston

# The Seesaw Mechanism

## How *Heavy* are HNLs?



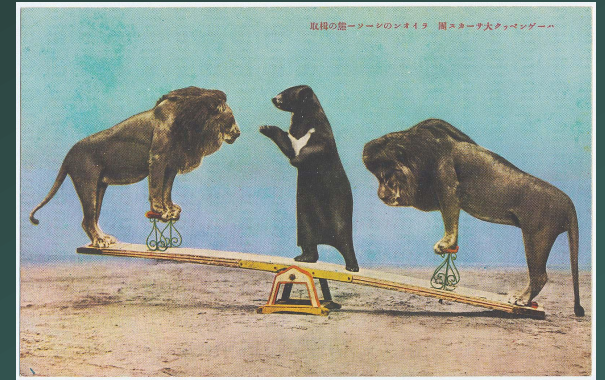
The Seesaw Lagrangian

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu}_L & \overline{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

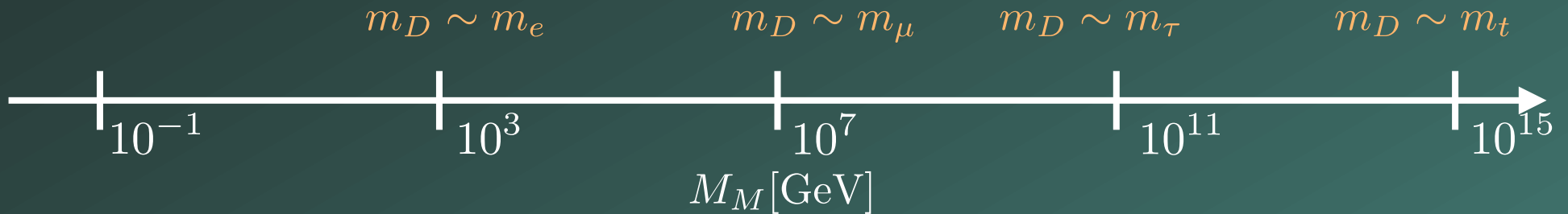
$$m_\nu = -m_D M_M^{-1} m_D^T$$

Neutrino masses  
are extremely small

$$m_\nu \lesssim 0.8 \text{ eV}$$

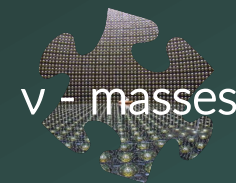


Unknown Artist,  
Japanese, MFA Boston



# The Seesaw Mechanism

## How *Heavy* are HNLs?



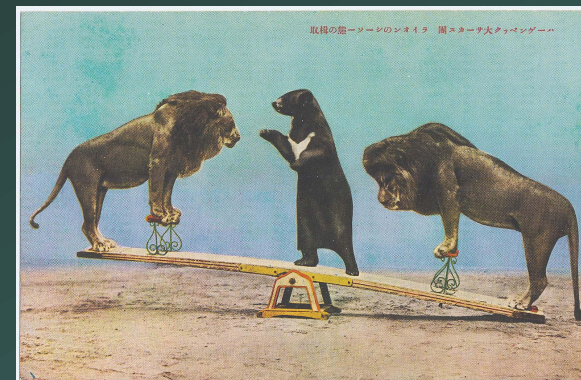
The Seesaw Lagrangian

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu}_L & \overline{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

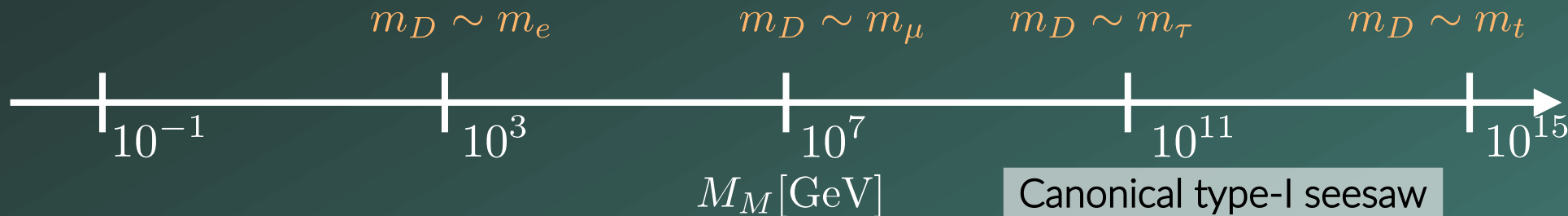
$$m_\nu = -m_D M_M^{-1} m_D^T$$

Neutrino masses  
are extremely small

$$m_\nu \lesssim 0.8 \text{ eV}$$



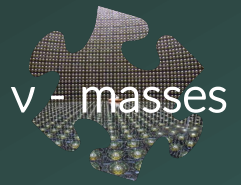
Unknown Artist,  
Japanese, MFA Boston





# The Seesaw Mechanism

## How *Heavy* are HNLs?



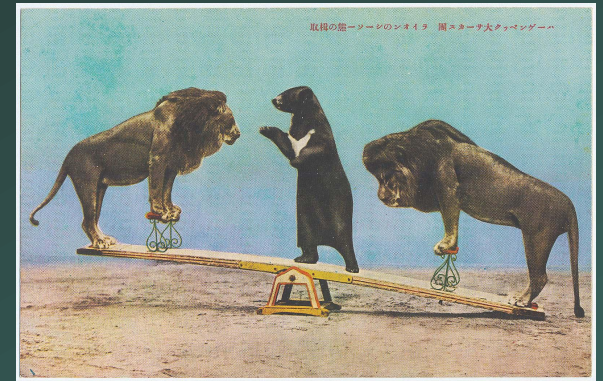
### The Seesaw Lagrangian

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu}_L & \overline{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

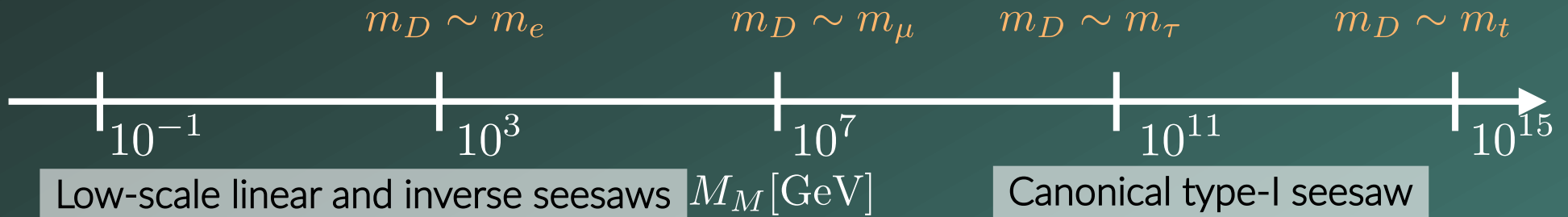
$$m_\nu = -m_D M_M^{-1} m_D^T$$

Neutrino masses  
are extremely small

$$m_\nu \lesssim 0.8 \text{ eV}$$

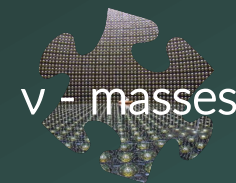


Unknown Artist,  
Japanese, MFA Boston



# The Seesaw Mechanism

## How *Heavy* are HNLs?



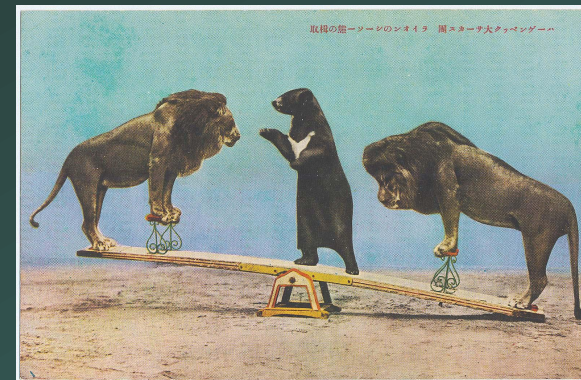
The Seesaw Lagrangian

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu}_L & \overline{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

$$m_\nu = -m_D M_M^{-1} m_D^T$$

Neutrino masses  
are extremely small

$$m_\nu \lesssim 0.8 \text{ eV}$$



Unknown Artist,  
Japanese, MFA Boston

Collider testable region  $m_D \sim m_e$

$m_D \sim m_\mu$

$m_D \sim m_\tau$

$m_D \sim m_t$



Low-scale linear and inverse seesaws  $M_M [\text{GeV}]$

Canonical type-I seesaw

# How do HNLs interact with the SM?

## Minimal Renormalizable Interactions

- The only **renormalizable** interactions are the **Yukawa** couplings:

$$\bar{L}_a \tilde{H} Y_{ai} \nu_{Ri}$$

# How do HNLs interact with the SM?

## Minimal Renormalizable Interactions

- The only **renormalizable** interactions are the **Yukawa** couplings:

$$\bar{L}_a \tilde{H} Y_{ai} \nu_{R_i}$$

$$\rightarrow \bar{\nu}_{L_a} \langle H \rangle Y_{ai} \nu_{R_i} = \bar{\nu}_L m_D \nu_R$$

## After Electroweak Symmetry Breaking

- HNL mass eigenstates are a **mixture** of  $\nu_R$  and  $\nu_L$

$$N_i \simeq \nu_{R_i} + \theta_{ia}^T \nu_{L_a}^c$$

# How do HNLs interact with the SM?

## Minimal Renormalizable Interactions

- The only **renormalizable** interactions are the **Yukawa** couplings:

$$\bar{L}_a \tilde{H} Y_{ai} \nu_{R_i}$$

$$\rightarrow \bar{\nu}_{L_a} \langle H \rangle Y_{ai} \nu_{R_i} = \bar{\nu}_L m_D \nu_R$$

## After Electroweak Symmetry Breaking

- HNL mass eigenstates are a **mixture** of  $\nu_R$  and  $\nu_L$

$$N_i \simeq \nu_{R_i} + \theta_{ia}^T \nu_{L_a}^c$$

## Beyond Minimal Interactions

- HNLs are a part of a **light hidden sector**, e.g. **LRSM**...
- Other particles present e.g.  $W_R$ ,  $Z'$

# How do HNLs interact with the SM?

## Minimal Renormalizable Interactions

- The only **renormalizable** interactions are the **Yukawa** couplings:

$$\bar{L}_a \tilde{H} Y_{ai} \nu_{R_i}$$

$$\rightarrow \bar{\nu}_{L_a} \langle H \rangle Y_{ai} \nu_{R_i} = \bar{\nu}_L m_D \nu_R$$

## After Electroweak Symmetry Breaking

- HNL mass eigenstates are a **mixture** of  $\nu_R$  and  $\nu_L$

$$N_i \simeq \nu_{R_i} + \theta_{ia}^T \nu_{L_a}^c$$

## Beyond Minimal Interactions

- HNLs are a part of a **light hidden sector**, e.g. **LRSM**...
- Other particles present e.g.  $W_R$ ,  $Z'$

## Heavy new mediators

- HNLs are the **only light new particles**
- Additional interactions are described by effective interactions
  - e.g. the  $\nu_R$  **SMEFT**

# How do HNLs interact with the SM?

## Minimal Renormalizable Interactions

- The only **renormalizable** interactions are the **Yukawa** couplings:

$$\bar{L}_a \tilde{H} Y_{ai} \nu_{R_i}$$

$$\rightarrow \bar{\nu}_{L_a} \langle H \rangle Y_{ai} \nu_{R_i} = \bar{\nu}_L m_D \nu_R$$

## After Electroweak Symmetry Breaking

- HNL mass eigenstates are a **mixture** of  $\nu_R$  and  $\nu_L$

$$N_i \simeq \nu_{R_i} + \theta_{ia}^T \nu_{L_a}^c$$

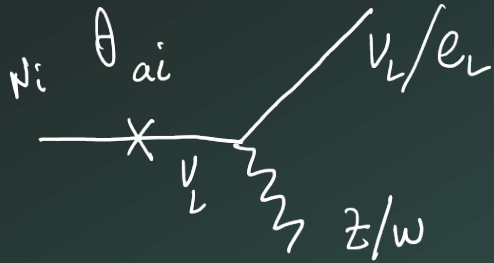
## Beyond Minimal Interactions

- HNLs are a part of a **light hidden sector**, e.g. **LRSM**...
- Other particles present e.g.  $W_R$ ,  $Z'$

## Heavy new mediators

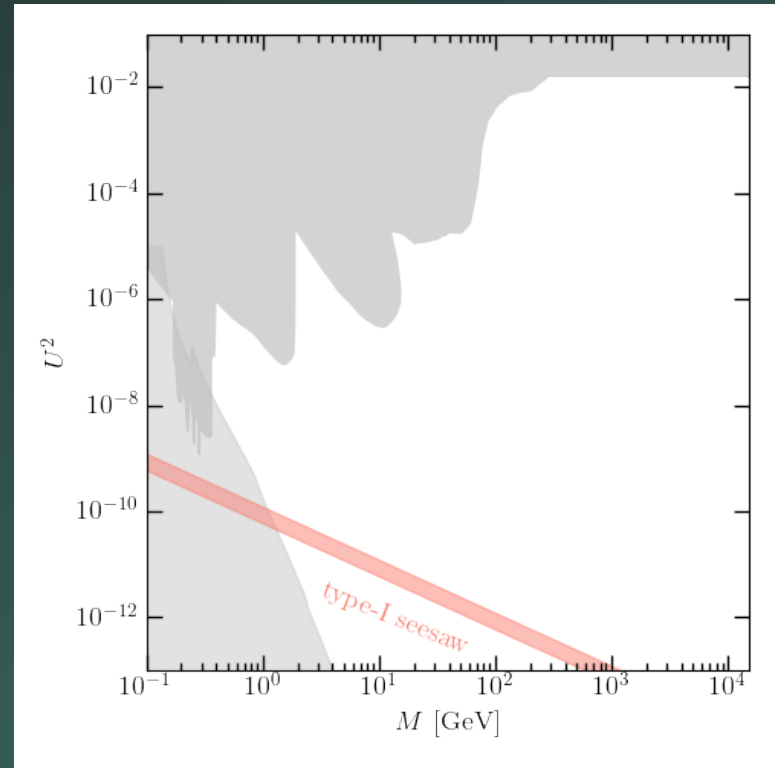
- HNLs are the **only light new particles**
- Additional interactions are described by effective interactions
  - e.g. the  $\nu_R$  **SMEFT**

# The parameter space of the minimal low-scale seesaw model



$$U_{ai}^2 \equiv |\theta_{ai}|^2 = |(m_D M_M^{-1})_{ai}|^2$$

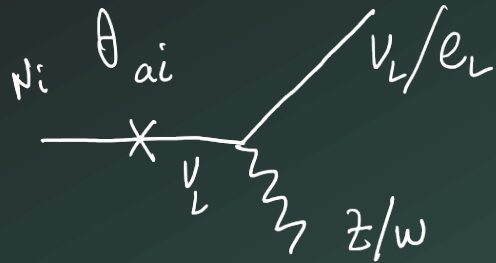
$$U^2 = \sum_{a,i} U_{ai}^2 \quad U^2 \gtrsim m_\nu / M$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

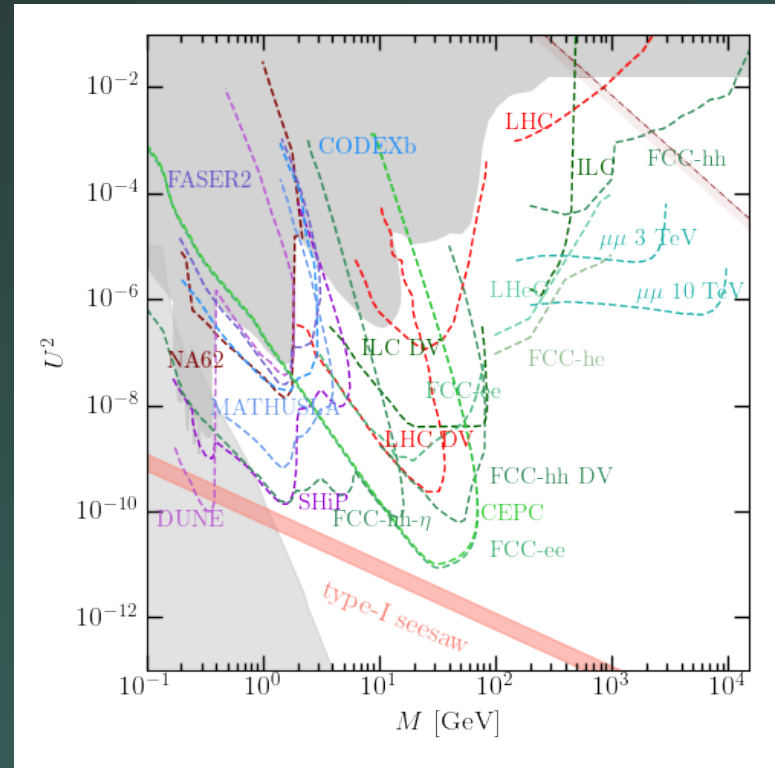


# The parameter space of the minimal low-scale seesaw model



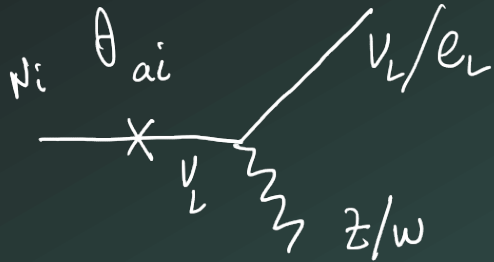
$$U_{ai}^2 \equiv |\theta_{ai}|^2 = |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2 \quad U^2 \gtrsim m_\nu / M$$



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

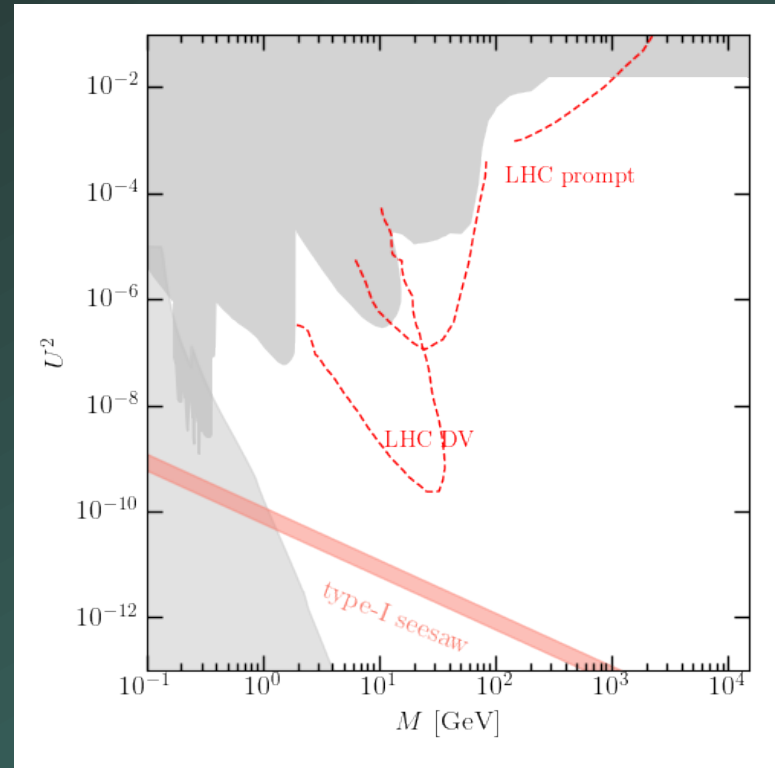
# The parameter space of the minimal low-scale seesaw model



$$U_{ai}^2 \equiv |\theta_{ai}|^2 = |(m_D M_M^{-1})_{ai}|^2$$

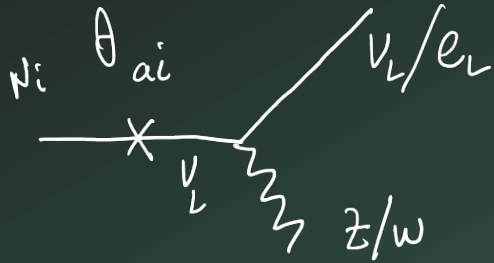
$$U^2 = \sum_{a,i} U_{ai}^2 \quad U^2 \gtrsim m_\nu / M$$

[see the talks by Frattari, Portales, Schaarschmidt, Knolle, Lunerti, Henry and Bird]



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

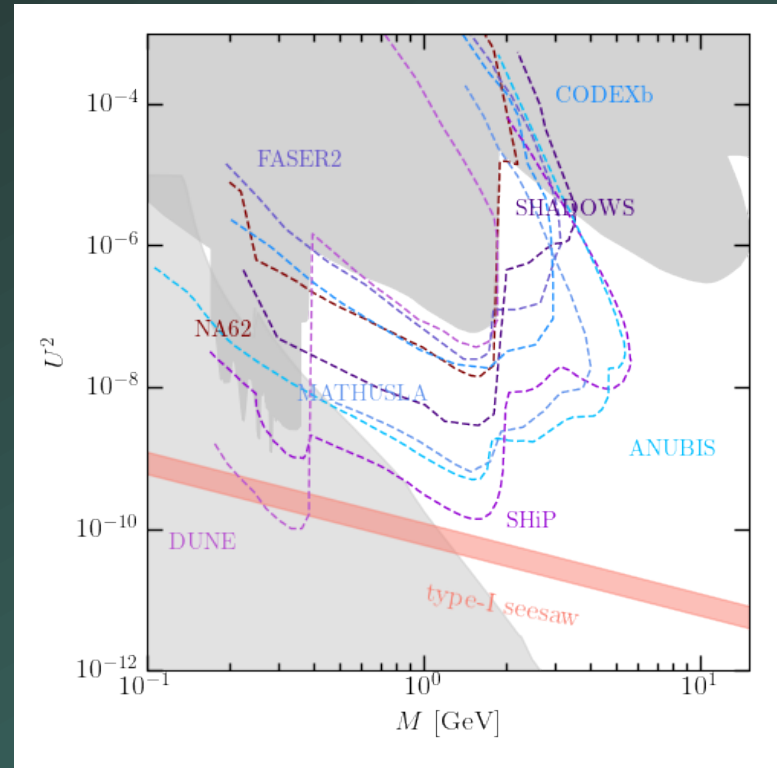
# The parameter space of the minimal low-scale seesaw model



$$U_{ai}^2 \equiv |\theta_{ai}|^2 = |(m_D M_M^{-1})_{ai}|^2$$

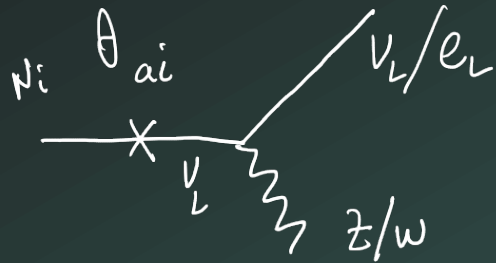
$$U^2 = \sum_{a,i} U_{ai}^2 \quad U^2 \gtrsim m_\nu / M$$

[see the session on future facilities]



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

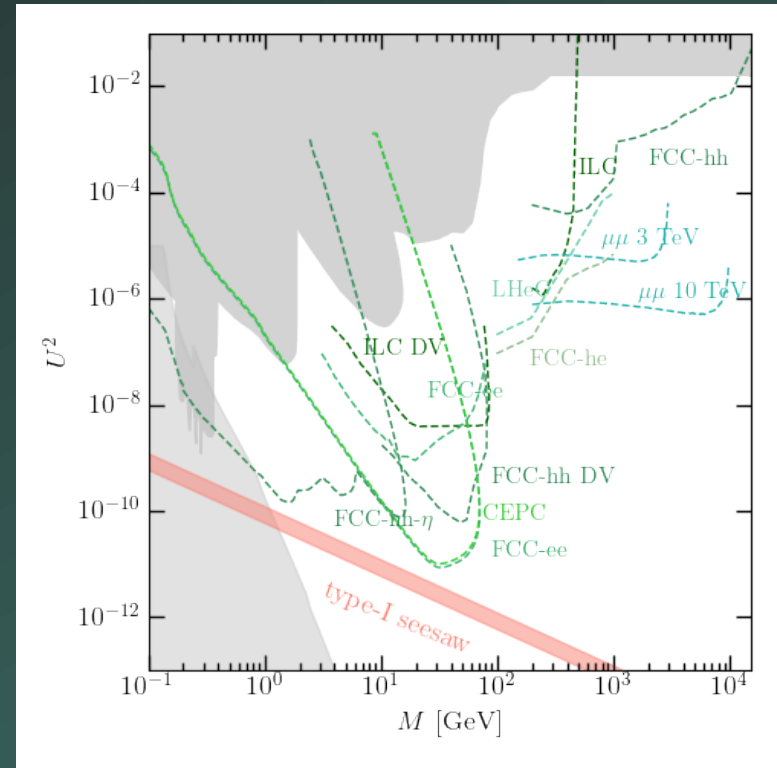
# The parameter space of the minimal low-scale seesaw model



$$U_{ai}^2 \equiv |\theta_{ai}|^2 = |(m_D M_M^{-1})_{ai}|^2$$

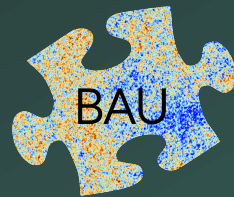
$$U^2 = \sum_{a,i} U_{ai}^2 \quad U^2 \gtrsim m_\nu / M$$

[see the session on future facilities]



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

# HNLs and the $B_{\text{aryon}}A_{\text{symmetry}}$ of the $U_{\text{niverse}}$



- Three [Sakharov '67] conditions:

1. Baryon number violation

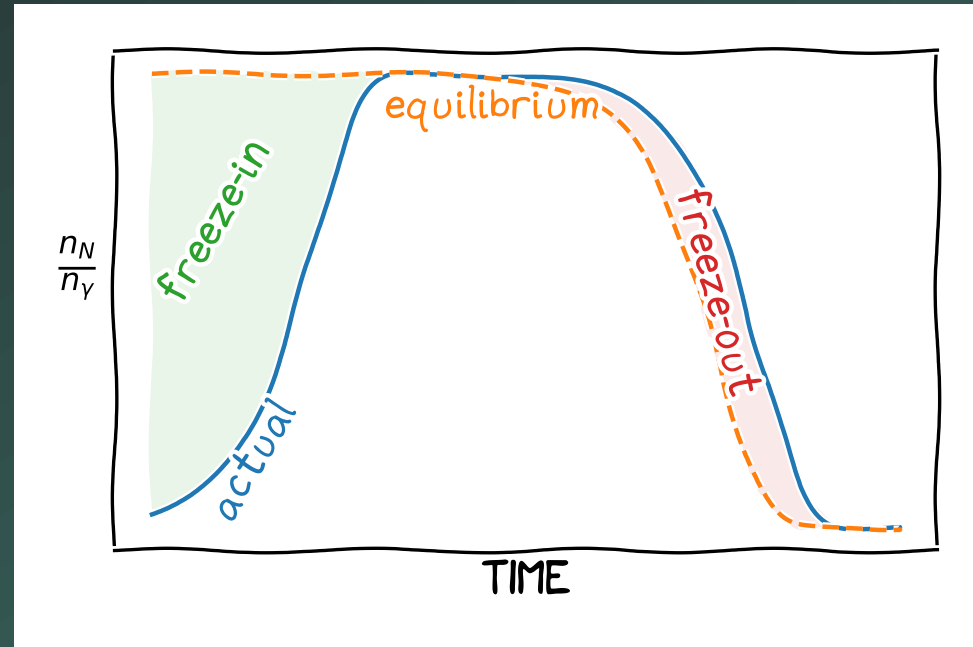
- sphaleron processes ✓

2. C and CP violation

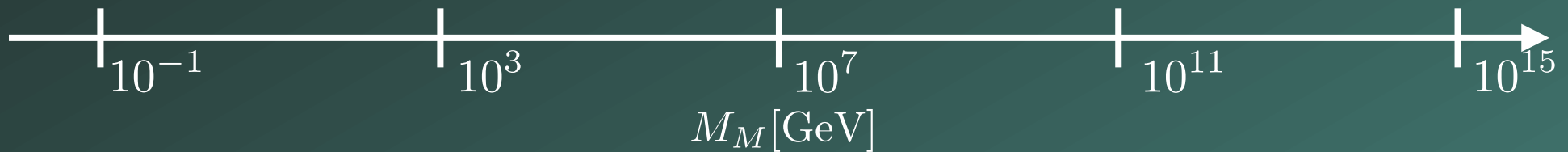
- HNL decays and oscillations ✓

3. Deviation from equilibrium

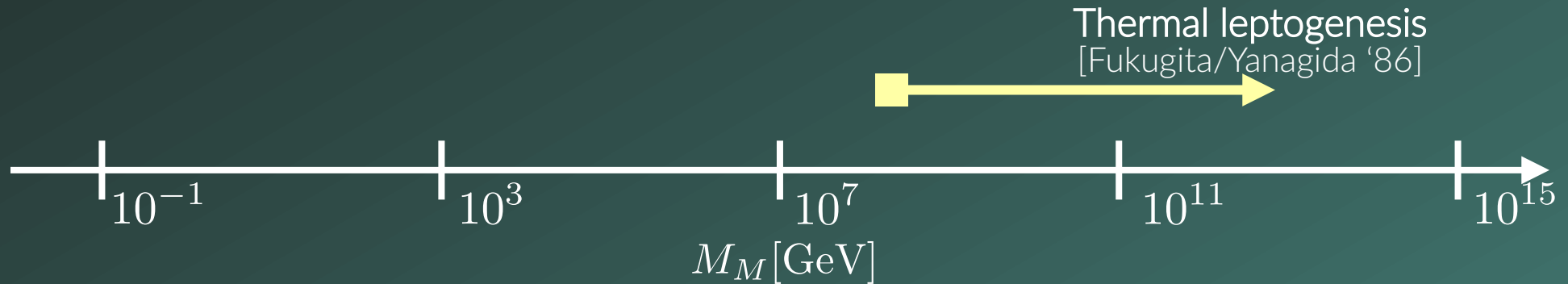
- freeze-in and freeze-out of HNLs ✓



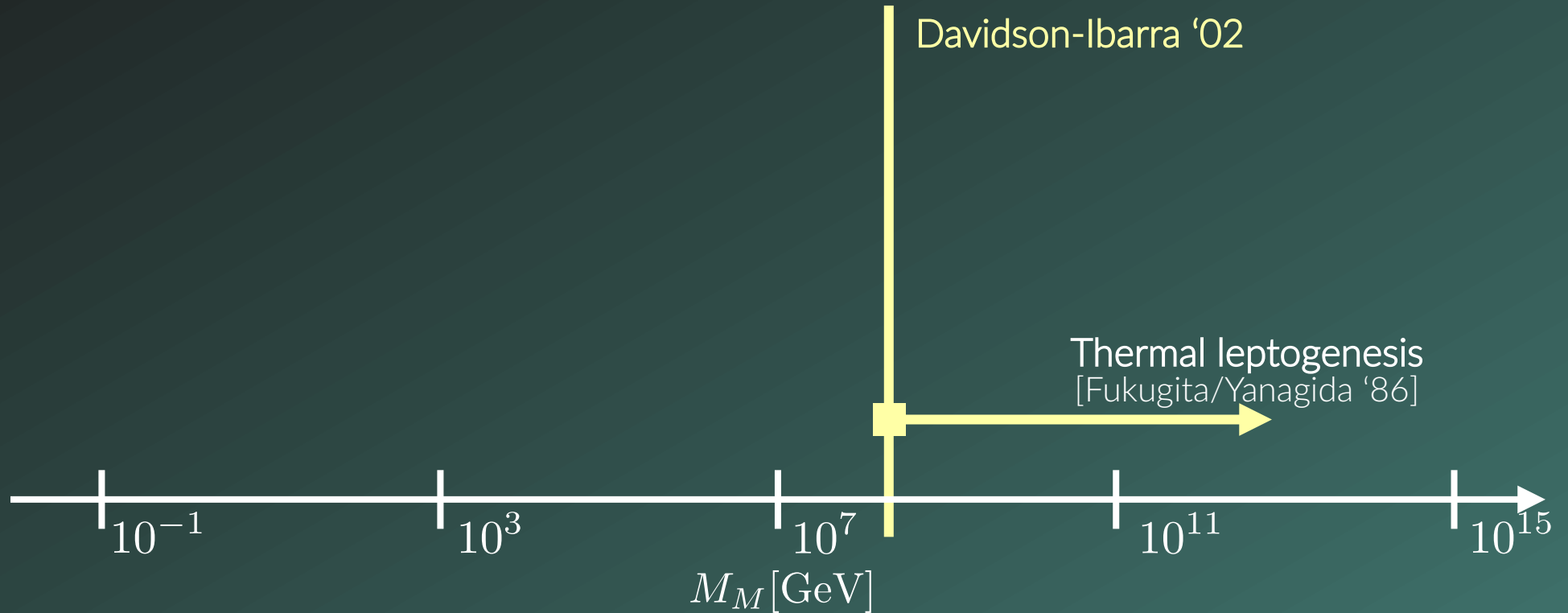
What does leptogenesis tell us about HNL masses?



# What does leptogenesis tell us about HNL masses?

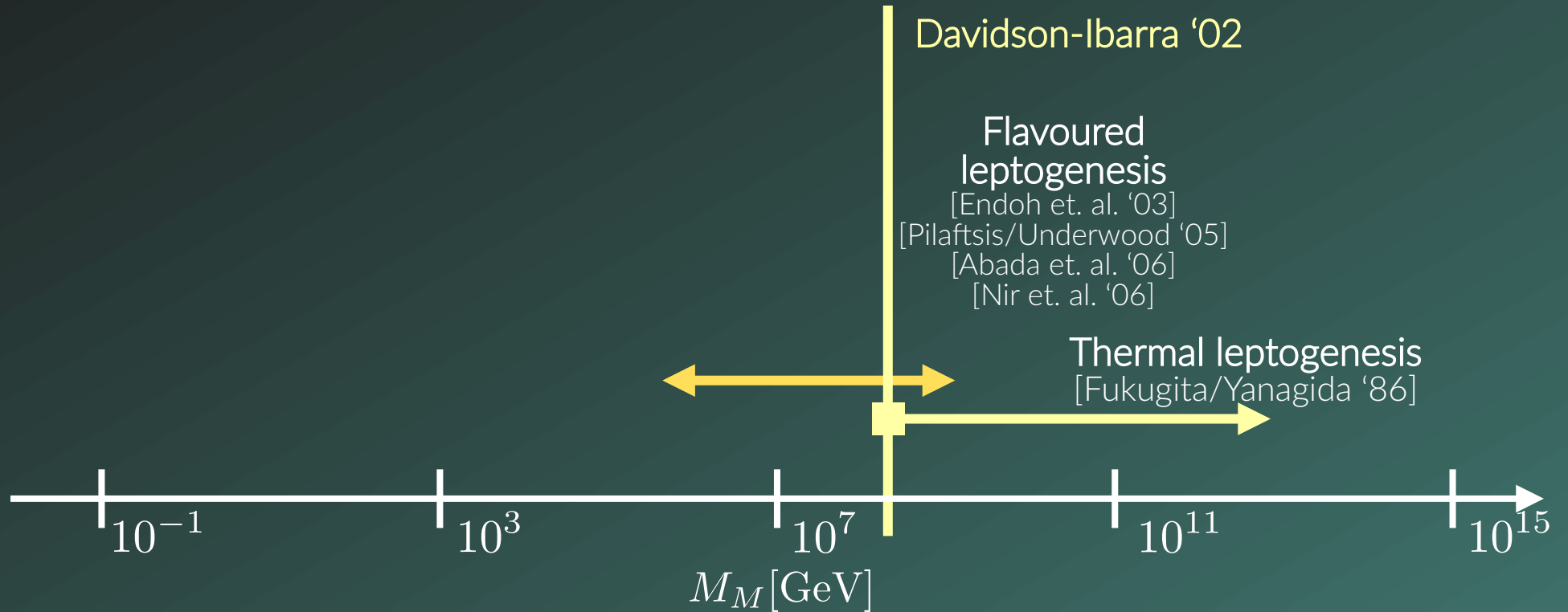


# What does leptogenesis tell us about HNL masses?

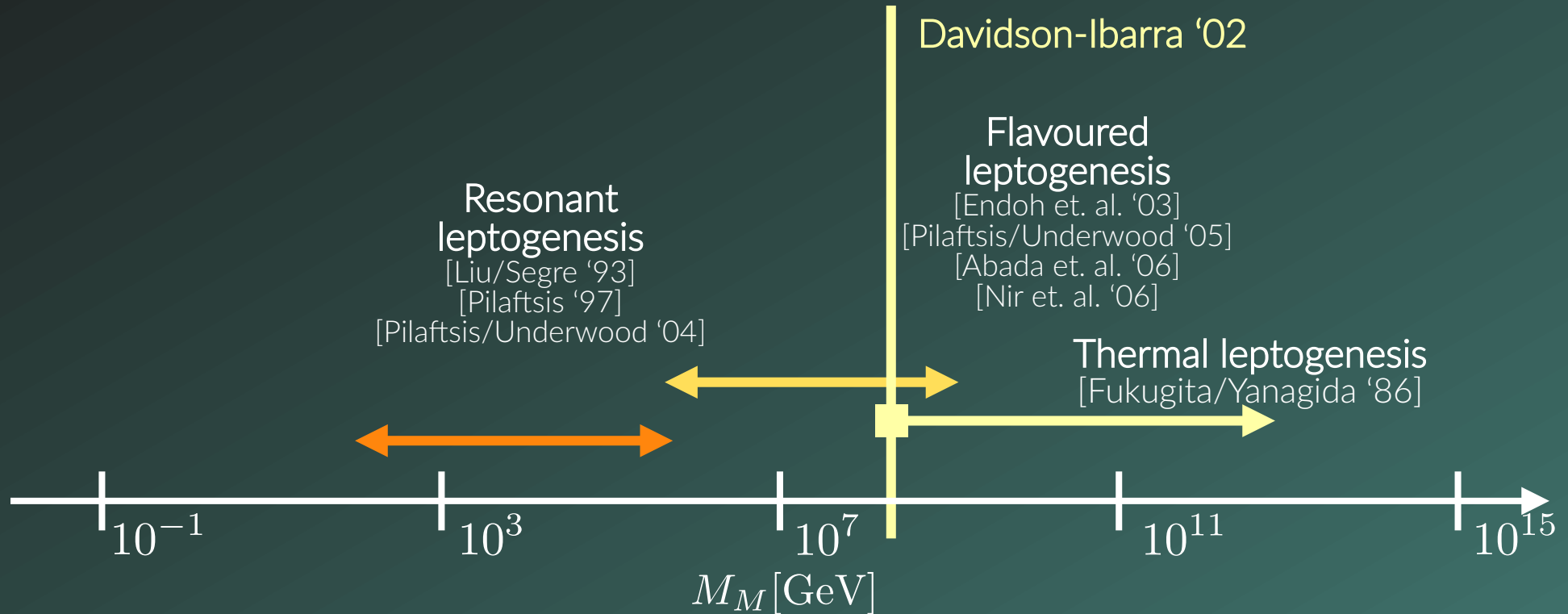




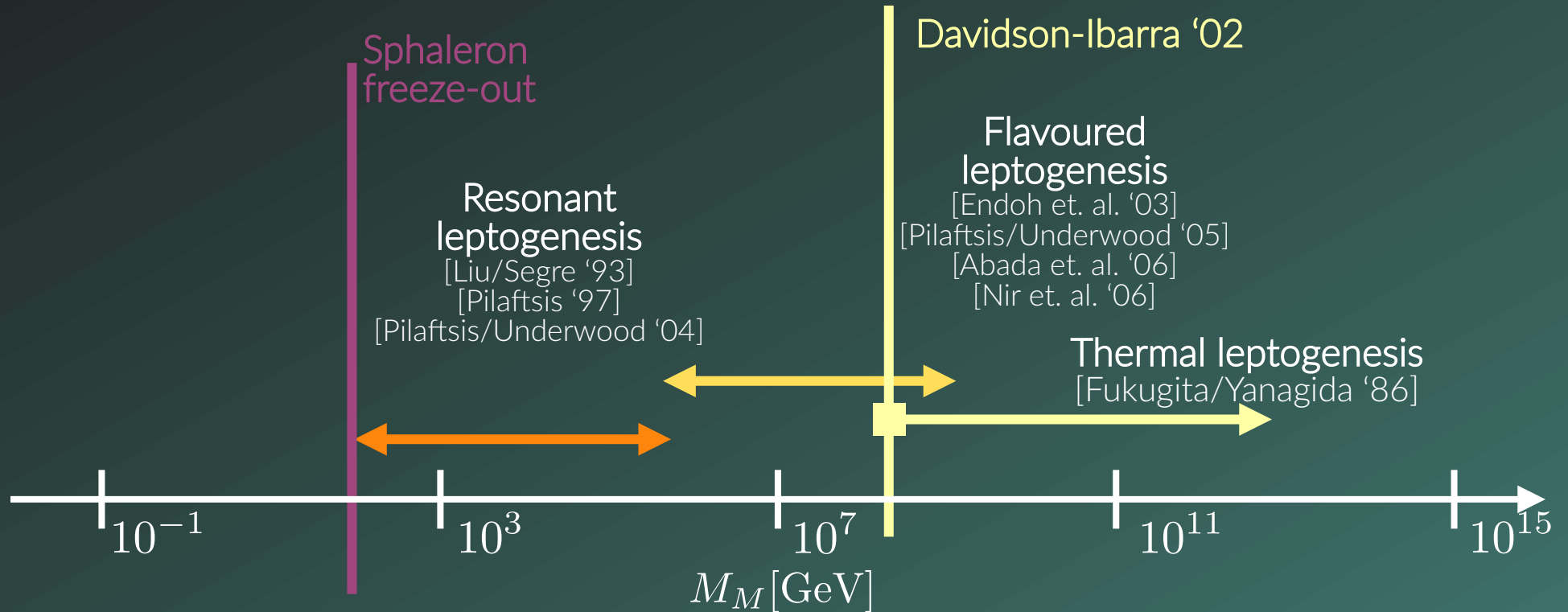
# What does leptogenesis tell us about HNL masses?



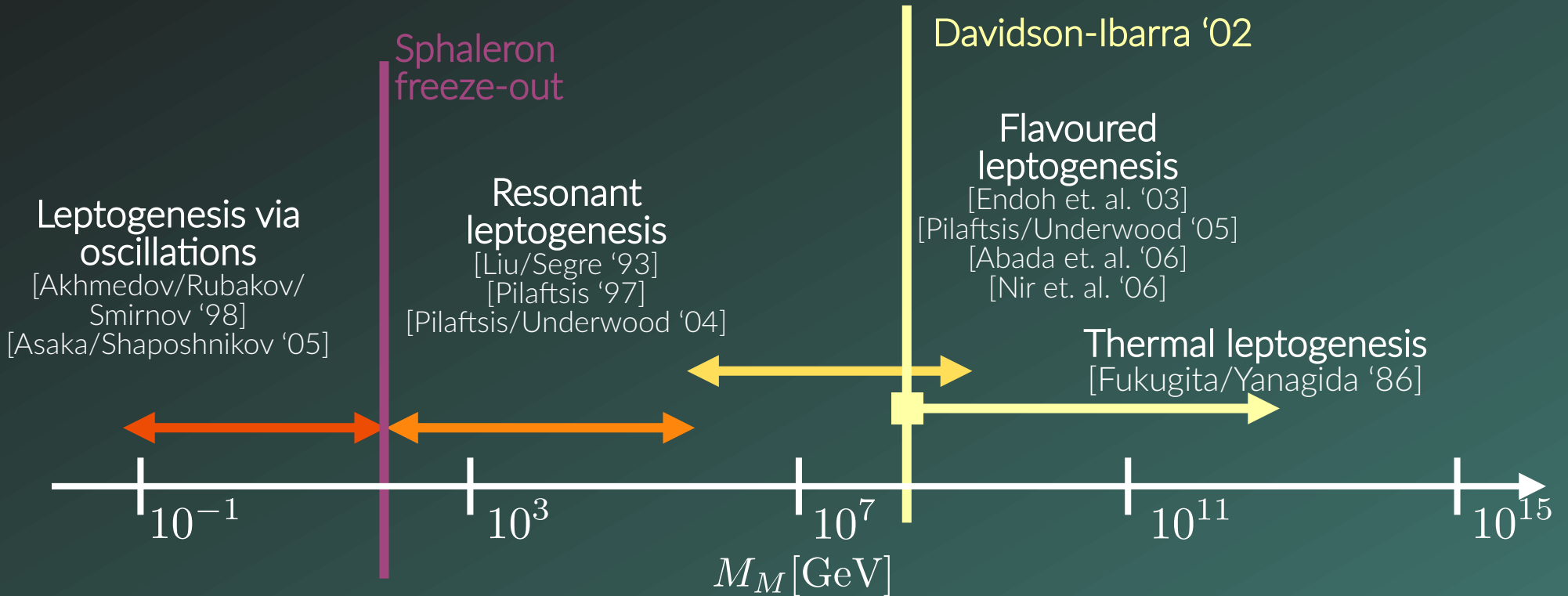
# What does leptogenesis tell us about HNL masses?



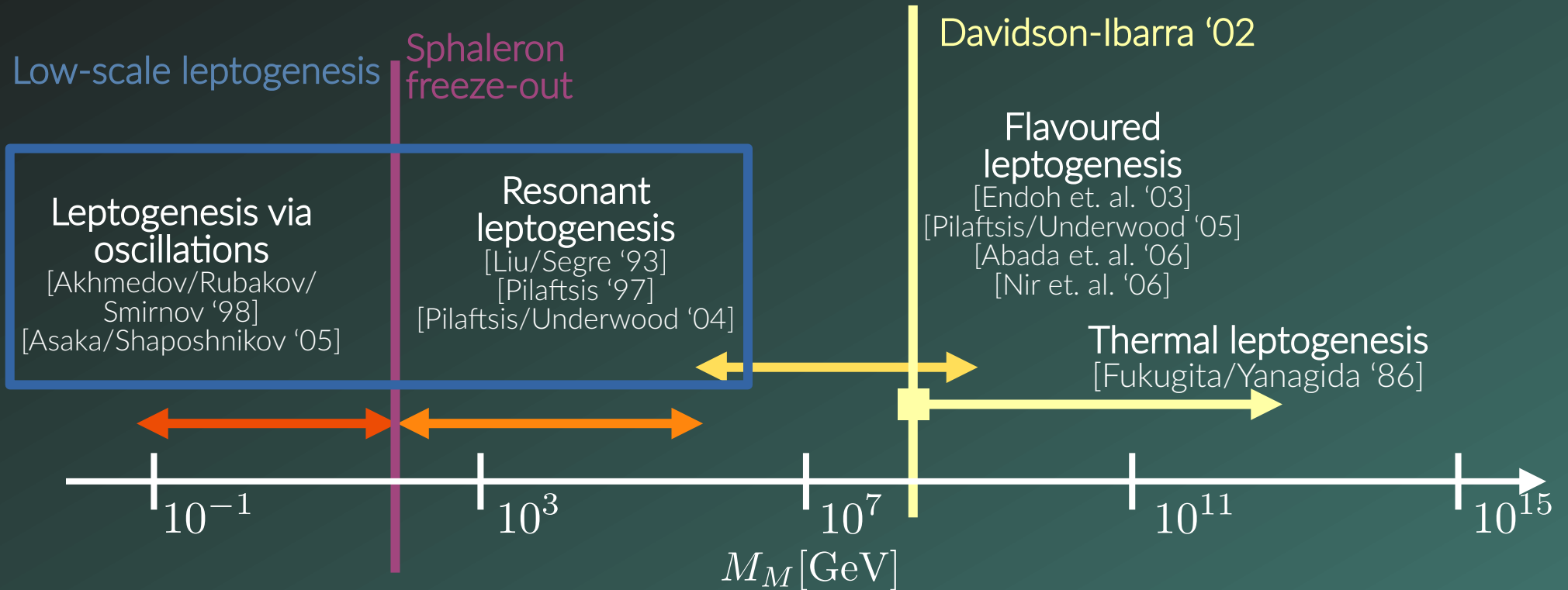
# What does leptogenesis tell us about HNL masses?



# What does leptogenesis tell us about HNL masses?

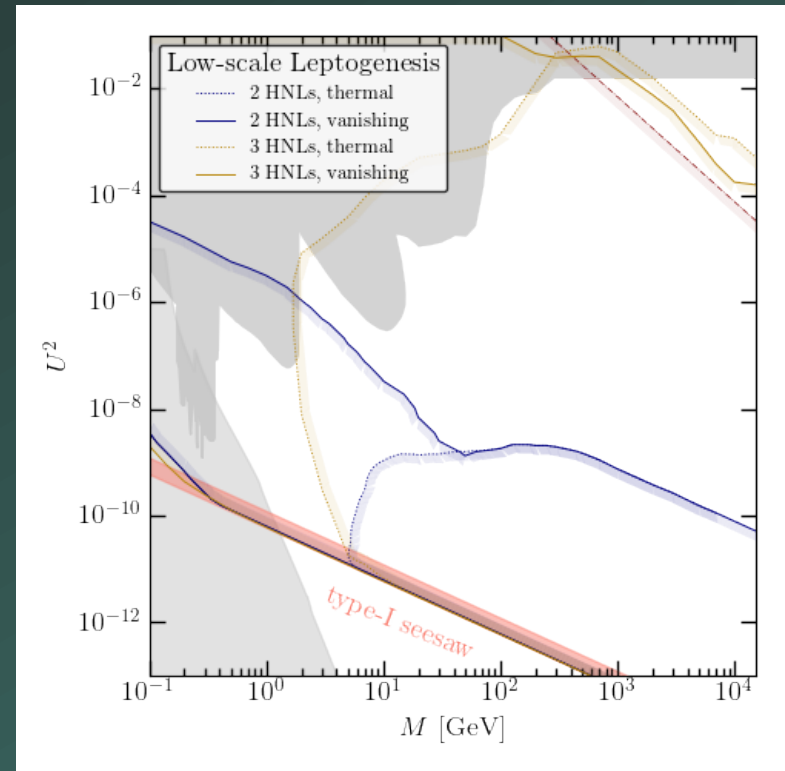


# What does leptogenesis tell us about HNL masses?



# The parameter space of leptogenesis

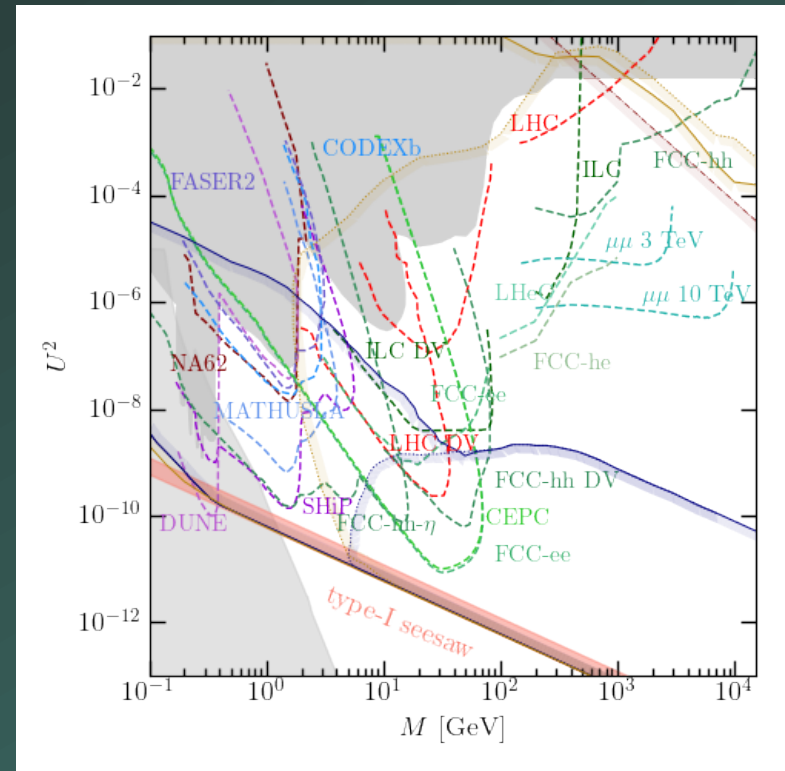
- With 2 HNLs leptogenesis is possible for *all masses above 100 MeV*
- Leptogenesis is possible in the *entire experimentally accessible parameter space* for 3 HNLs
- Both **vanishing** (no additional interactions) and **thermal** (high-scale additional interactions) leptogenesises possible
- Leptogenesis **within reach** of HL-LHC
- High complementarity between **colliders** and **dedicated LLP** searches



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]  
[leptogenesis bounds from JK/Timiryasov/Shaposhnikov 2103.16545  
and Drewes/Georis/JK 2106.16226 ]

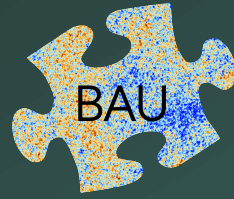
# The parameter space of leptogenesis

- With 2 HNLs leptogenesis is possible for *all masses above 100 MeV*
- Leptogenesis is possible in the *entire experimentally accessible parameter space* for 3 HNLs
- Both **vanishing** (no additional interactions) and **thermal** (high-scale additional interactions) leptogenesis possible
- Leptogenesis **within reach** of HL-LHC
- High complementarity between **colliders** and **dedicated LLP** searches



[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]  
[leptogenesis bounds from JK/Timiryasov/Shaposhnikov 2103.16545  
and Drewes/Georis/JK 2106.16226 ]

# Sakharov conditions at colliders



Can we test the Sakharov conditions at colliders?

## 1. Baryon number violation

- sphaleron processes

## 2. C and CP violation

- HNL decays and oscillations

## 3. Deviation from equilibrium

- freeze-in and freeze-out of HNLs

- Sphalerons processes are extremely suppressed at  $T \approx 0$  GeV
- In the SM  $B-L$  is conserved: **lepton number violation**
- CP violation in HNL decays is challenging – but we could measure **CPV in  $\nu$  – oscillations**
- The equilibration rate of HNLs is directly probed by measuring their **couplings & branching ratios**



# Deviation from Equilibrium: HNL branching ratios

- HNL branching ratios are highly constrained by the measured parameters in the minimal model (2 HNLs)

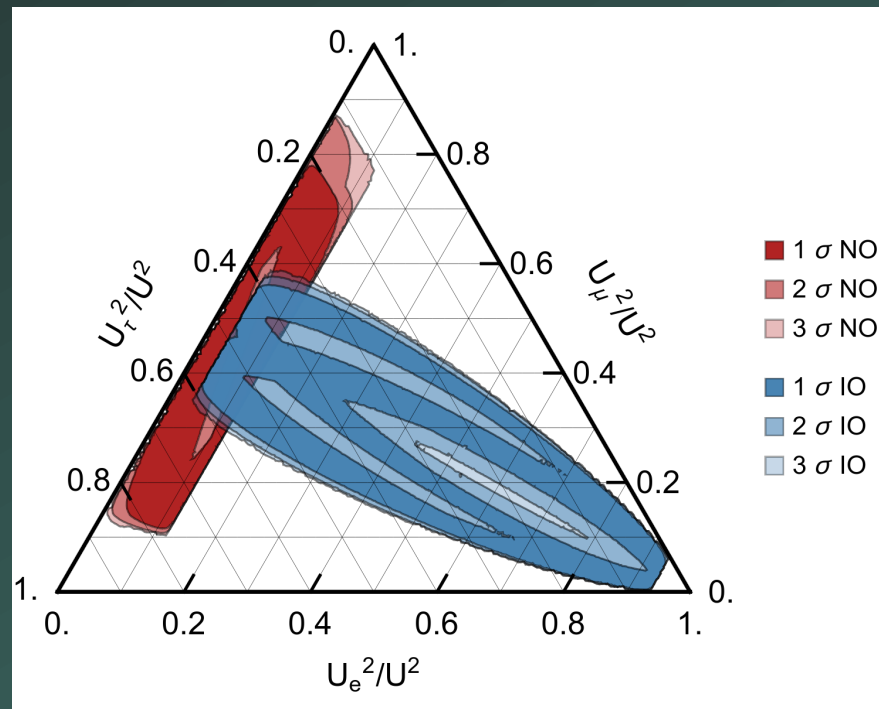
[Snowmass white paper 2203.08039]

- Leptogenesis imposes further constraints on the branching ratios

[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK 1710.03744]

- Branching ratios become even more predictive when combined with Flavor and CP symmetries

[Drewes/Georis/Hagedorn/JK 24xx.xxxx]



# Deviation from Equilibrium: HNL branching ratios

- HNL branching ratios are highly constrained by the measured parameters in the minimal model (2 HNLs)

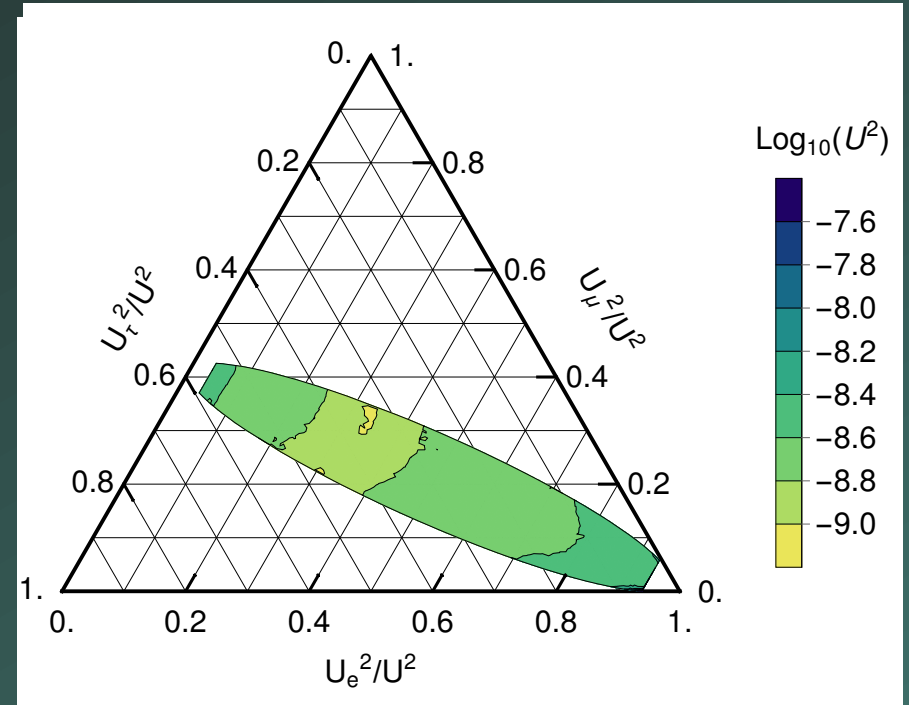
[Snowmass white paper 2203.08039]

- Leptogenesis imposes further constraints on the branching ratios

[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK 1710.03744]

- Branching ratios become even more predictive when combined with Flavor and CP symmetries

[Drewes/Georis/Hagedorn/JK 24xx.xxxx]



# Deviation from Equilibrium: HNL branching ratios

- HNL branching ratios are highly constrained by the measured parameters in the minimal model (2 HNLs)

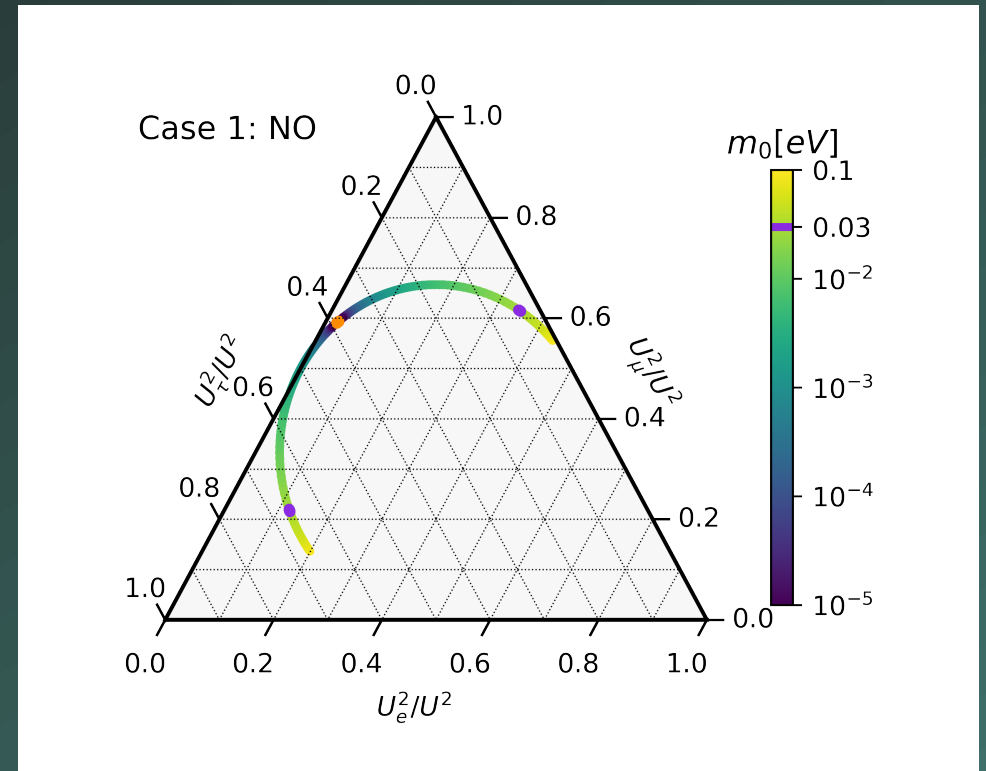
[Snowmass white paper 2203.08039]

- Leptogenesis imposes further constraints on the branching ratios

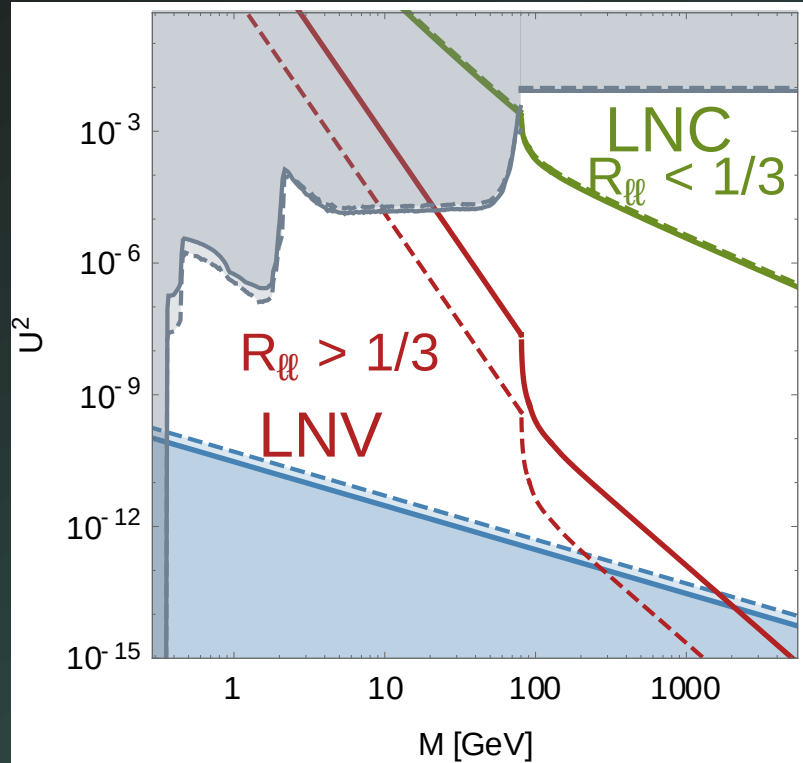
[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK 1710.03744]

- Branching ratios become even more predictive when combined with Flavor and CP symmetries

[Drewes/Georis/Hagedorn/JK 24xx.xxxx]



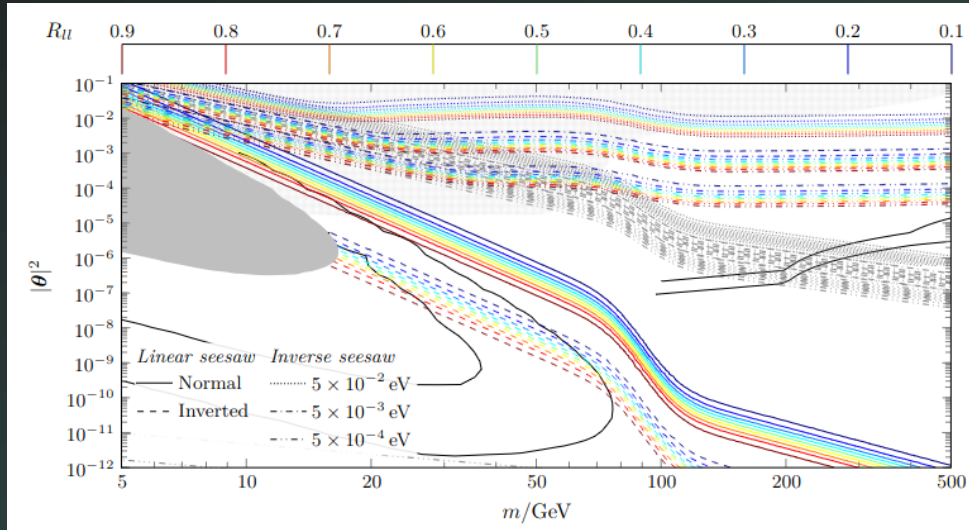
# Lepton Number Violation



[Drewes/Klose/JK 1907.13034]

- In low-scale seesaw mechanisms HNLs preserve approximate  $B-L$  symmetry
- LNC only for  $\Gamma \gg \Delta M$
- Absence of fine tuning implies lower bound on HNL mass splitting  $\Delta M > \Delta m_\nu$
- Prompt decays can be sensitive to decoherence effects
- For *tiny* mass splittings HNL oscillations between LNV and LNC processes possible

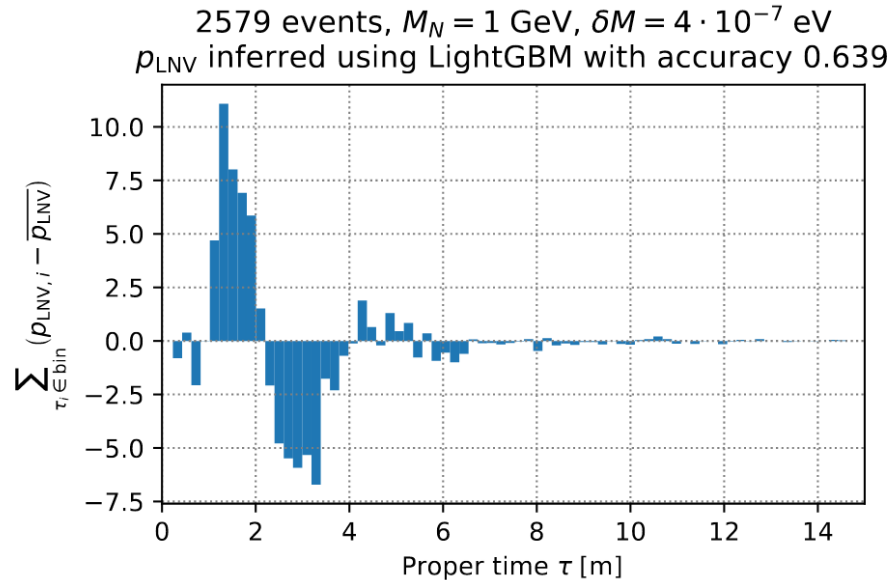
# Lepton Number Violation



[Antusch/Hajer/Roskopp 2308.07297]

- In low-scale seesaw mechanisms HNLs preserve approximate  $B-L$  symmetry
- LNC only for  $\Gamma \gg \Delta M$
- Absence of fine tuning implies lower bound on HNL mass splitting  $\Delta M > \Delta m_\nu$
- Prompt decays can be sensitive to decoherence effects
- For *tiny* mass splittings HNL oscillations between LNV and LNC processes possible

# Lepton Number Violation

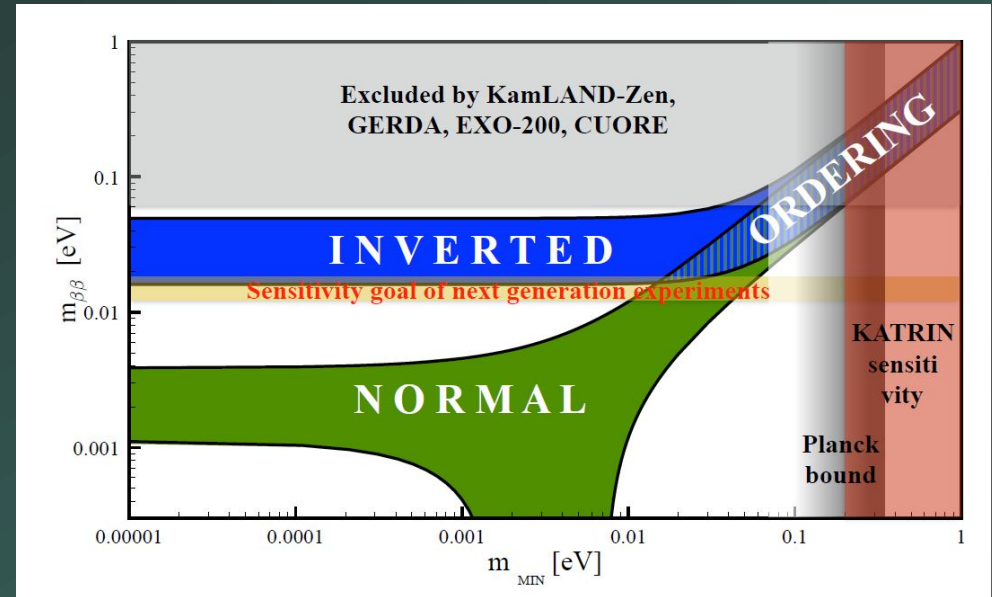


[Tastet/Timiryasov 1912.05520]

- In low-scale seesaw mechanisms HNLs preserve approximate  $B-L$  symmetry
- LNC only for  $\Gamma \gg \Delta M$
- Absence of fine tuning implies lower bound on HNL mass splitting  $\Delta M > \Delta m_\nu$
- Prompt decays can be sensitive to decoherence effects
- For *tiny* mass splittings HNL oscillations between LNV and LNC processes possible

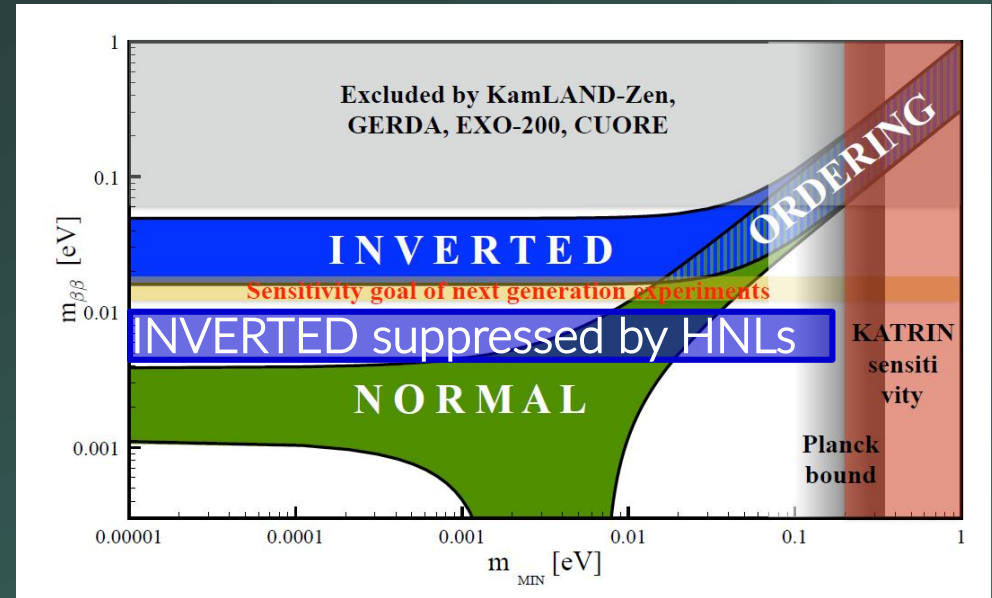
# Indirect LNV: neutrinoless double $\beta$ decay

- Smoking gun signature of LNV
- GeV-scale HNLs can modify the naive expectation
- Suppression of the IO signal implies light HNLs
- Target region testable with HL-LHC, SHiP & FCC-ee



# Indirect LNV: neutrinoless double $\beta$ decay

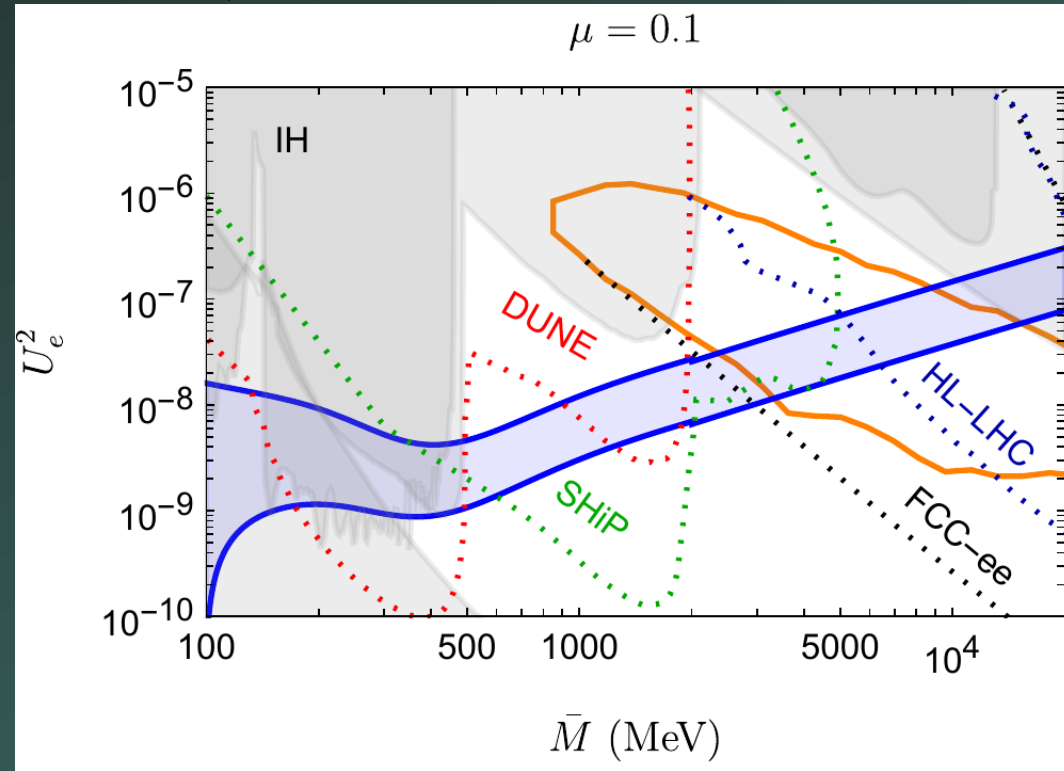
- Smoking gun signature of LNV
- GeV-scale HNLs can modify the naive expectation
- Suppression of the IO signal implies light HNLs
- Target region testable with HL-LHC, SHiP & FCC-ee





# Indirect LNV: neutrinoless double $\beta$ decay

- Smoking gun signature of LNV
- GeV-scale HNLs can modify the naive expectation
- Suppression of the IO signal implies light HNLs
- Target region testable with HL-LHC, SHiP & FCC-ee



[Fig from de Vries, Drewes, Georis, JK, Plakkot 24XX.XXXX]  
[figures from 1910.04688]

# Conclusions

HNLs can offer a minimal solution to the origins of **neutrino masses** and the **baryon asymmetry of the Universe**

- the existence of HNLs is already being tested at existing experiments
- excellent synergy between **high-energy** and **high-intensity** experiments!
- leptogenesis is a viable baryogenesis mechanism for all HNL masses above  $O(100)$  MeV scale
- Indirect probes can lead to **clear target regions for HNLs**
- HNLs have a very rich phenomenology

# Additional References

- **Canonical seesaw:** [Minkowski '77] [Gell-Mann/Ramond/Slansky '79] [Mohapatra/Senjanović '80] [Yanagida '79] [Schechter/Valle '80]
- **Low-scale seesaws:** [Mohapatra '93] [Mohapatra/Valle '86] [Bernabeu/Santamaria/Vidal/Mendez/Valle '86] [Gavela/Hambye/Hernandez/Hernandez '09] [Branco/Grimus/Lavoura '89] [Malinsky/Romao/Lavoura '89]
- **Leptogenesis:** [Fukugita/Yanagida '86]
- **Flavoured leptogenesis:** [Endoh et. al. '03] [Pilaftsis/Underwood '05] [Abada et. al. '06] [Nir et. al. '06]
- **Resonant Leptogenesis:** [Liu/Segre '93, Pilaftsis '97, Pilaftsis/Underwood '04;'05]
- **Leptogenesis via neutrino oscillations:** [Akhmedov/Rubakov/Smirnov '98, Asaka/Shaposhnikov '05]