Hadronization in small and large systems

Andrea Beraudo

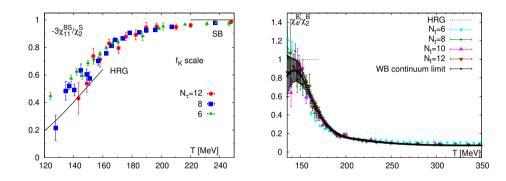
INFN - Sezione di Torino

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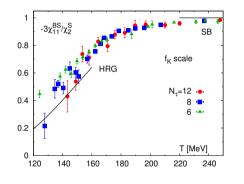


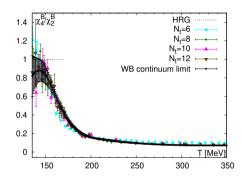
Premise: which are the carriers of conserved charges?



- In the QGP strangeness carried by quarks with |B|=1/3, PRD 86, 034509 (2012)
- $\chi_4^B/\chi_2^2 = B^2$, with |B|=1 (HG) or |B|=1/3 (QGP), PRL 111, 062005 (2013)

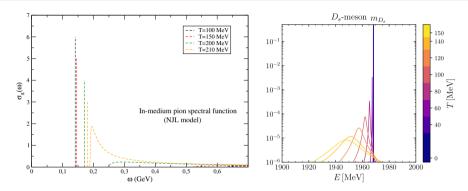
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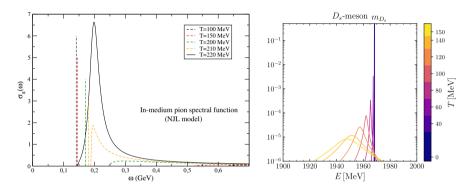
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One would expect a sharp change in the nature of these carriers... However, IQCD data show that also this change is very smooth!



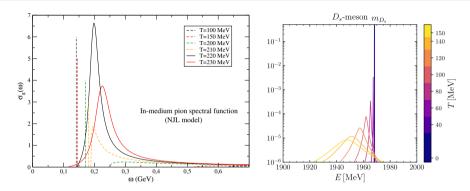
- At T=0 hadrons are stable eigenstates of H_{QCD}
- At $T \neq 0$ effective Lagrangians predict much richer structure of hadronic spectral functions (broadening, mass shift), both for light (NJL model) and heavy (non-linear chiral SU(3) model) hadrons¹

¹G. Montana et al., PLB 806 (2020)



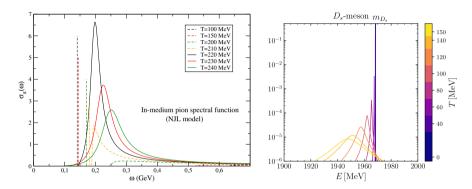
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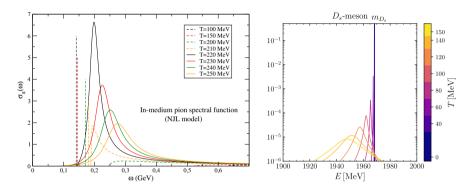
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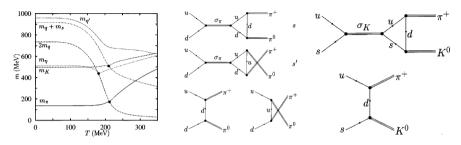
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Light sector: hadronization in effective chiral Lagrangians



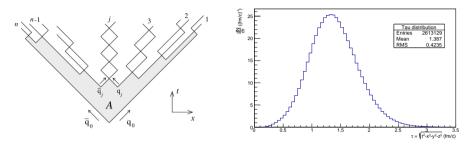
From $N_f = 2 + 1$ NJL Lagrangian (P. Rehberg et al, PRC 53 (1995) 410)

- Different Mott temperatures for the different hadrons, below which the decay channel $H \to q_1 + \overline{q}_2$ gets closed (non-universal hadronization temperature?)
- Hadronization, modeled as $q_1 + \overline{q}_2 \rightarrow H_1 + H_2$ process (exact four-momentum conservation \neq coalescence), takes time to occur:

$$\tau_{u}^{-1}(T_{H}) = \langle \sigma_{u\bar{u}}v \rangle \rho_{\bar{u}} + \langle \sigma_{u\bar{d}}v \rangle \rho_{\bar{d}} + \langle \sigma_{u\bar{s}}v \rangle \rho_{\bar{s}} \approx (2-3 \text{ fm/c})^{-1}$$

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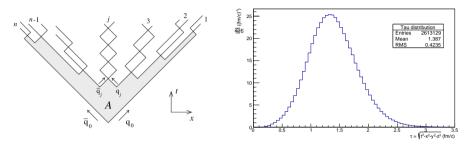
Hadronization does not occur suddenly



Also in string-fragmentation model (PYTHIA) primary hadron production takes time to occur, $\langle \tau \rangle \approx 1.3$ fm/c, however only recently model builders started investigating its implications (S. Ferreres-Solé and T. Sjostrand, EPJC 78 (2018) 11, 983)

- Probably irrelevant in e^+e^- collisions
- Important to consider if a dense medium (big or small), with its own time scales (lifetime, interaction rate, expansion rate...), is formed.

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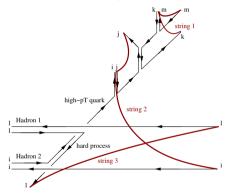
Notice that in most cases in HIC's hadronization is modeled as an instantaneous process (e.g. Cooper-Frye particlization or standard coalescence approaches)

Hadronization models: common features

Grouping colored partons into color-singlet structures: strings (PYTHIA), clusters (HERWIG), hadrons/resonances (coalescence/recombination).

Hadronization models: common features

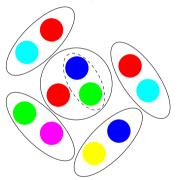
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• in "elementary collisions" (what is elementary?): from the hard process, shower stage, underlying event and beam remnants;

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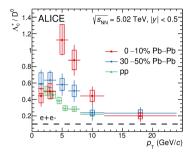
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- in "elementary collisions" (what is elementary?): from the hard process, shower stage, underlying event and beam remnants;
- in heavy-ion collisions (only?): from the hot medium produced in the collision.

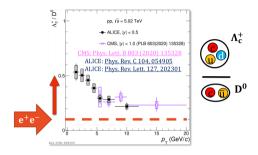
 NB Involved partons closer in space in this case and this has deep consequence!

Why the title of this talk?



• Strong enhancement of charmed baryon/meson ratio both in AA and pp collisions, incompatible with hadronization models tuned to reproduce e^+e^- data.

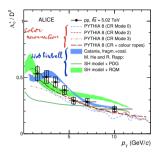
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• Strong enhancement of charmed baryon/meson ratio both in AA and pp collisions, incompatible with hadronization models tuned to reproduce e^+e^- data. Breaking of factorization of *inclusive* hadronic cross-sections in pp collisions

$$d\sigma_h \neq \sum_{a,b,X} f_a(x_1) f_b(x_2) \otimes d\hat{\sigma}_{ab \to c\bar{c}X} \otimes D_{c \to h_c}(z)$$

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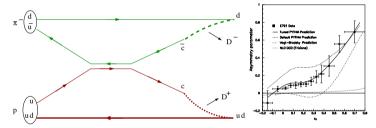
• Recent theory attempts to explain the data either based on Color Reconnection (CR) or on the formation of a small fireball: really different pictures?

Non-universality of hadronization: not the first observation

Breaking of factorization already observed in fixed target experiments at Fermilab and SPS (e.g. $\pi^- + p$ collisions) in the production of charmed hadrons sharing a valence (di-)quark with the beam or target remnant $(D^-/D^+, D_s^-/D_s^+, \Lambda_c^+/\overline{\Lambda}_c^-)$

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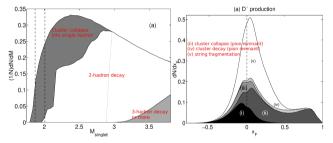


Second endpoint boosts the string along the direction of the beam-remnant (beam-drag effect), leading to an asymmetry in the rapidity distribution of D^-/D^+ mesons

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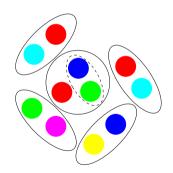


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Major contribution to asymmetry from collapse of a *very light cluster* into a single hadron (E. Norrbin and T. Sjostrand, PLB 442 (1998) 407 and EPJC 17 (2000) 137)!

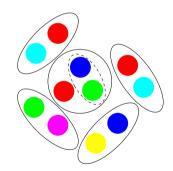
Local Color Neutralization (LCN): basic ideas



Both in AA and pp collisions a big/small deconfined fireball is formed. Around the QCD crossover temperature quarks undergoes recombination with the *closest* opposite color-charge (antiquark or diquark, favoring baryon production).

- Why? screening of color-interaction, minimization of energy stored in confining potential

Local Color Neutralization (LCN): basic ideas



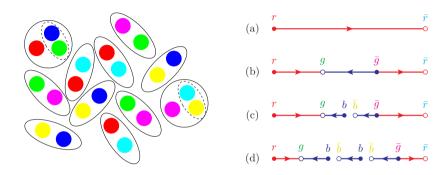
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- Why? screening of color-interaction, minimization of energy stored in confining potential
- Implication: recombination of particles from the same fluid cell
 Space-Momentum Correlation (SMC), recombined partons tend to share a common collective velocity

Color-singlet structures are thus formed, eventually undergoing decay into the final hadrons: $2 \to 1 \to N$ process, usually a charmed hadron plus a very soft particle

- Exact four-momentum conservation;
- No direct bound-state formation, hence no need to worry about overlap between the final hadron and the parent parton wave-functions

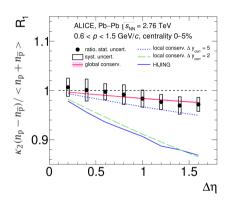
Implementation of global conservation laws

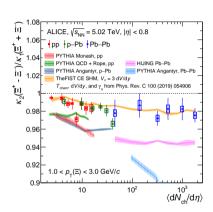


- In LCN and similar recombination approaches baryon number (and other charges as well) can be conserved over a very large volume;
- On the other hand in PYTHIA string-breaking (and possibly pop-corn) mechanism charge conservation occurs locally²

²L. Lonnblad and H. Shah, EPJC 83 (2023) 12, 1105

Global conservation of charges: experimental results

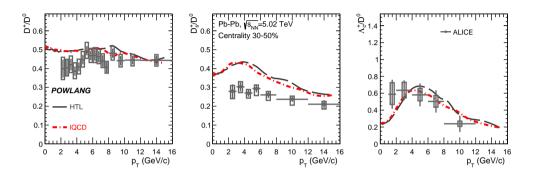




Normalized net-particle cumulants favour charge conservation (e.g. B, S) over a very big correlation volume³

³ALICE Coll. PLB 807 (2020) 135564 and 2405.19890

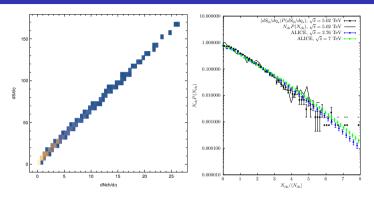
LCN: results in AA collisions



- Enhanced HF baryon-to-meson ratios up to intermediate p_T nicely reproduced, thanks to formation of *small invariant-mass* charm+diquark clusters⁴
- Smooth approach to e^+e^- limit $(\Lambda_c^+/D^0\approx 0.1)$ at high p_T : high- M_c clusters fragmented as Lund strings, as in the vacuum

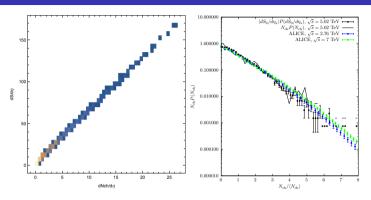
⁴A.B. et al., EPJC 82 (2022) 7, 607

Addressing pp collisions...



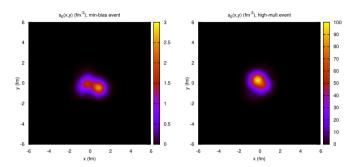
 EBE pp initial conditions generated with TrENTo and evolved with hydro codes (MUSIC and ECHO-QGP);

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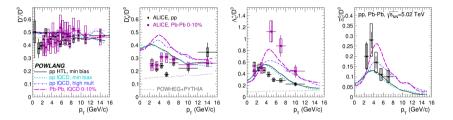
- EBE pp initial conditions generated with TrENTo and evolved with hydro codes (MUSIC and ECHO-QGP);
- Perfect correlation between initial entropy (dS/dy) and final particle multiplicity $(dN_{\rm ch}/d\eta)$, $S \approx 7.2N_{\rm ch}$. $P(N_{\rm ch})$ satisfying KNO scaling nicely reproduced;

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- Perfect correlation between initial entropy (dS/dy) and final particle multiplicity $(dN_{\rm ch}/d\eta)$, $S \approx 7.2N_{\rm ch}$. $P(N_{\rm ch})$ satisfying KNO scaling nicely reproduced;
- Samples of 10^3 minimum-bias ($\langle dS/dy \rangle_{\rm mb} \approx 37.6$, tuned to experimental $\langle dN_{\rm ch}/d\eta \rangle$) and high-multiplicity ($\langle dS/dy \rangle_{0-1\%} \approx 187.5$) events used to simulate HQ transport and hadronization.

Results in pp: particle ratios

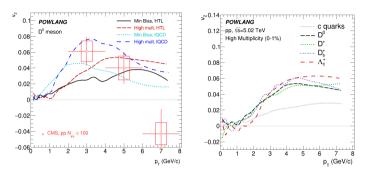


First results for particle ratios⁵:

- POWHEG+PYTHIA standalone strongly underpredicts baryon-to-meson ratio
- Enhancement of charmed baryon-to-meson ratio qualitatively reproduced if propagation+hadronization in a small QGP droplet is included
- Multiplicity dependence of radial-flow peak position (just a reshuffling of the momentum, without affecting the yields): $\langle u_{\perp} \rangle_{\rm pp}^{\rm mb} \approx 0.33$, $\langle u_{\perp} \rangle_{\rm pp}^{\rm hm} \approx 0.53$, $\langle u_{\perp} \rangle_{\rm PbPb}^{0-10\%} \approx 0.66$

⁵In collaboration with D. Pablos, A. De Pace, F. Prino et al., PRD 109 (2024) 1_□L011501 = ≥

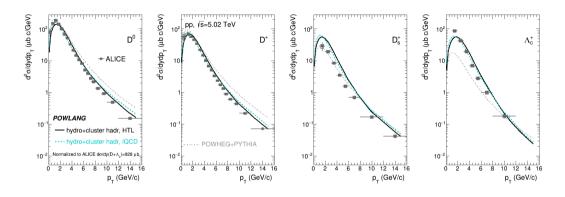
Results in pp: elliptic flow



Response to initial elliptic eccentricity ($\langle \epsilon_2 \rangle^{\mathrm{mb}} \approx \langle \epsilon_2 \rangle^{\mathrm{mh}} \approx 0.31$) \longrightarrow non-vanishing v_2 coefficient

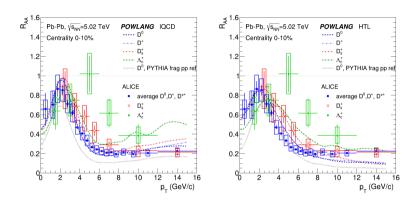
- Differences between minimum-bias and high-multiplicity results only due to longer time spent in the fireball ($\langle \tau_H \rangle^{\rm mb} \approx 1.95$ fm/c vs $\langle \tau_H \rangle^{\rm hm} \approx 2.92$ fm/c)
- Mass ordering at low p_T $(M_{qq} > M_q)$
- Sizable fraction of v_2 acquired at hadronization

Relevance to quantify nuclear effects

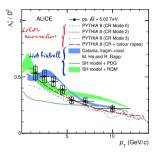


• Slope of the spectra in pp collisions better described including medium effects

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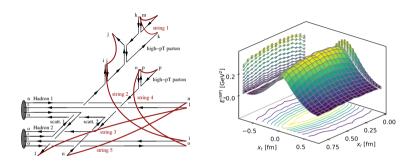


- Slope of the spectra in pp collisions better described including medium effects
- Inclusion of medium effects in minimum-bias pp benchmark fundamental to better describe charmed hadron $R_{\rm AA}$, both the radial-flow peak and the species dependence



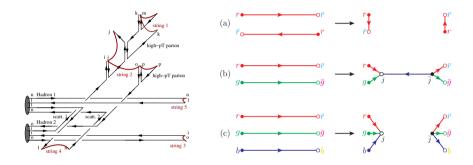
Charmed baryon enhancement in *pp* collisions can be accounted for *either* assuming the formation of a small fireball *or*, in PYTHIA, introducing the possibility of color-reconnection (CR).

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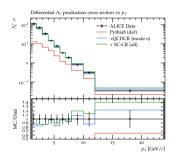
Charmed baryon enhancement in pp collisions can be accounted for either assuming the formation of a small fireball or, in PYTHIA, introducing the possibility of color-reconnection (CR). Strings have a finite thickness, in a dense environment they can overlap⁶

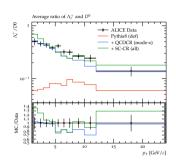
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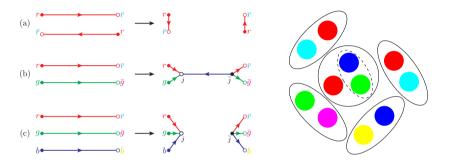
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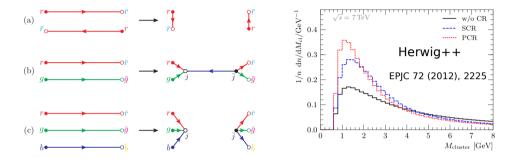
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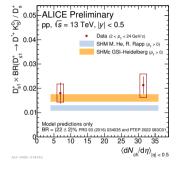
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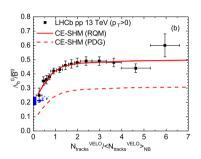


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HF statistical hadronization



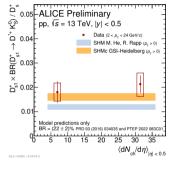


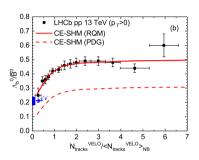
$$Z(\vec{Q}) = \int_{0}^{2\pi} \frac{d^{5}\phi}{(2\pi)^{5}} e^{i\vec{Q}\cdot\vec{\phi}} \exp\left[\sum_{j} \gamma_{s}^{N_{sj}} \gamma_{c}^{N_{cj}} \gamma_{b}^{N_{bj}} e^{-i\vec{q}_{j}\cdot\vec{\phi}} z_{j}\right], \quad \text{with} \quad z_{j} = (2J_{j} + 1) \frac{VT_{H}}{2\pi^{2}} m_{j}^{2} K_{2}\left(\frac{m_{j}}{T_{H}}\right)$$

Statistical description of HF production (Y. Dai, S. Zhao and M. He, 2402.03692) accounting for

• Deviation from chemical equilibrium (quark fugacities)

HF statistical hadronization



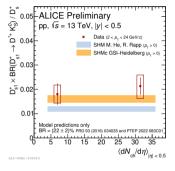


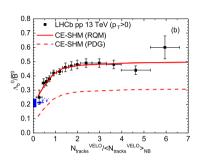
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Statistical description of HF production (Y. Dai, S. Zhao and M. He, 2402.03692) accounting for

- Deviation from chemical equilibrium (quark fugacities)
- Exact charge conservation within correlation volume V

HF statistical hadronization





$$Z(\vec{Q}) = \int_0^{2\pi} \frac{d^5 \phi}{(2\pi)^5} e^{i\vec{Q} \cdot \vec{\phi}} \exp\left[\sum_j \gamma_s^{N_{sj}} \gamma_c^{N_{cj}} \gamma_b^{N_{bj}} e^{-i\vec{q}_j \cdot \vec{\phi}} z_j\right], \quad \text{with} \quad z_j = (2J_j + 1) \frac{V T_H}{2\pi^2} m_j^2 K_2\left(\frac{m_j}{T_H}\right)$$

Statistical description of HF production (Y. Dai, S. Zhao and M. He, 2402.03692) accounting for

- Deviation from chemical equilibrium (quark fugacities)
- Exact charge conservation within correlation volume V
- Enlarged set of hadronic resonsance wrt PDG



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- A cross-fertilization between different communities will be welcome (lattice-QCD, QCD event generators, nuclear physics...) and necessary to achieve a deeper understanding of hadronization