

# **CPV and lifetime measurements from CMS**

Enrico Lusiani (University & INFN PD), on behalf of the CMS Collaboration

Large Hadron Collider Physics Conference 2024

# Introduction and outline

- The LHC allows to test SM predictions and look for new physics in many ways
  - Direct evidence of new particles and signatures
  - Indirect evidence from SM deviations in rare decays and precision measurements
- Three precision measurements are presented in this talk:
  - Search for CP violation in D<sup>0</sup> → K<sub>S</sub>K<sub>S</sub>
  - Measurement of the B<sub>s</sub> → J/ψ K<sub>s</sub> effective lifetime
  - Measurement of the time-dependent CP violation in B<sub>s</sub> mesons

More b-physics results from CMS can be found here

# Search for CP violation in D<sup>0</sup> → K<sub>s</sub>K<sub>s</sub>

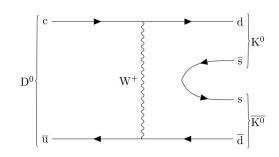
arXiv:2405.11606

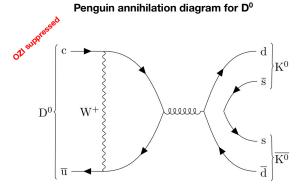
### W exchange diagram for D<sup>0</sup>

- CP violation in the up-quark sector is not studied as well as in the down-quark one
  - Expected to be suppressed by the GIM mechanism and CKM element size
- Observation of a significant CPV → hints of BSM physics
  - First observation of CPV in D mesons in 2019 by LHCb with  $D^0 \rightarrow K^+K^-$  and  $D^0 \rightarrow \pi^+\pi^-$  decays [PRL122(2019)211803]
- Presented here: measurement of the direct CPV in
   D<sup>0</sup> → K<sub>s</sub>K<sub>s</sub> decays

$$A_{CP} = \frac{\Gamma(D^0 \to K_S^0 K_S^0) - \Gamma(\overline{D}^0 \to K_S^0 K_S^0)}{\Gamma(D^0 \to K_S^0 K_S^0) + \Gamma(\overline{D}^0 \to K_S^0 K_S^0)}$$

• From theory, CPV in  $D^0 \rightarrow K_S K_S$  could be as large as O(1%)[PRD92(2015)054036]





# **Measurement strategy**

- Use  $D^0$  from  $D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \overline{D}{}^0 \pi^-$ , so that the pion charge tags the  $D^0$  flavor
- This introduces an additional asymmetry due to the D\*+/D\*- differences in the measurement

$$A_{raw} = \frac{N(D^{0}) - N(\overline{D}^{0})}{N(D^{0}) + N(\overline{D}^{0})}$$

$$A_{prod} = \frac{\sigma_{pp \to D^{*+}X} - \sigma_{pp \to D^{*-}X}}{\sigma_{pp \to D^{*+}X} + \sigma_{pp \to D^{*-}X}}$$

$$A_{det} \approx \frac{\epsilon_{\pi^{+}} - \epsilon_{\pi^{-}}}{\epsilon_{\pi^{+}} + \epsilon_{\pi^{-}}}$$

- Need a reference channel:  $\Delta A_{CP} = A_{CP}(D^0 \rightarrow K_S K_S) A_{CP}(D^0 \rightarrow K_S \pi^+\pi^-)$ 
  - $\circ$  Reference channel is very similar in kinematics and topology ightharpoonup and  $A_{\text{det}}$  cancel out
  - $\circ$  CPV in D<sup>0</sup> ightharpoonup  $K_{_{\rm S}}$   $\pi^+\pi^-$  already measured consistent with zero [PRD86(2012)032007]

$$\Delta A_{CP} = A_{raw}(D^0 o K_S K_S) - A_{raw}(D^0 o K_S \, \pi^+ \pi^-)$$

Signal channel  $R_S^0$ Reference channel  $R_S^0$ 

# **A**<sub>CP</sub> extraction

To extract the CP asymmetry a **2D maximum-likelihood fit is** performed on the invariant mass of the **D**\*\* and **D**<sup>0</sup>

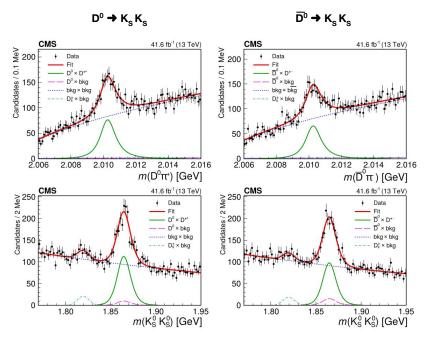
- Fit is done simultaneously on the D\*+ and D\*- samples with only the yields left to float
- Main fit components (signal channel):
  - o D<sup>0</sup> x D\*+, the signal component
  - $\circ$  D<sup>0</sup> x *bkg*, real D<sup>0</sup> but fake D\*+
  - o bkg x bkg, background in both dimensions
- Notable selections:  $m(\pi^+\pi^-) \in PDG \pm 20 \text{ MeV}$ ,  $m(K_{\varsigma}K_{\varsigma}) \in [1.7,2.0] \text{ GeV}$ , displaced by >9(2) $\sigma$  in xyz(xy)
- Background suppression: fit alternative topologies, select based on vertex probabilities
- Yields:

### Reference channel

Pion charge	N
$\pi^+$	$944800\pm3500$
$\pi^-$	$930150\pm3400$

### Signal channel

Pion charge	N	
$\pi^+$	$1095 \pm 46$	
$\pi^-$	$951 \pm 44$	



### Systematic uncertainties

Source	Uncertainty, %
$m(\mathrm{D}\pi^\pm)$ signal model	0.10
$m(\mathrm{D}\pi^\pm)$ background model	0.02
$m(K_S^0K_S^0)$ signal model	0.04
$m(K_S^0K_S^0)$ background model	0.02
$m(K_S^0K_S^0)$ fit range	0.04
Reweighting	0.09
$\Delta A_{CP}$ in MC	0.13
Total	0.20

### Results and outlook

Putting everything together, ΔA<sub>CP</sub> is measured

$$\Delta A_{CP} = 6.3 \pm 3.0 \, (\text{stat}) \pm 0.2 \, (\text{syst}) \, \%$$

• Using the world-average value of  $A_{CP}(K_S \pi^+\pi^-) = (-0.1 \pm 0.8)\%$ ,  $A_{CP}(K_S K_S)$  is found to be

$$A_{CP}(D^0 o K^0_S\,K^0_S) = 6.2\pm3.0\, ext{(stat)}\pm0.2\, ext{(syst)}\pm0.8(A_{CP}(K^0_S\,\pi^+\pi^-))\,\%$$

- Consistent with no CP violation at  $2\sigma$ , with LHCb [PRD104(2021)L031102] [(-3.1 ± 1.3)%] at  $2.7\sigma$  and Belle [PRL119(2017)171801] [(0.0 ± 1.5)%] at  $1.8\sigma$
- This is the first CMS study of CP violation in the charm sector, paving the way for future measurements using
  - More data
  - Refined techniques
  - Different channels

# Measurement of the $B_s \rightarrow J/\psi K_s$ effective lifetime

CMS PAS BPH-22-001

# **Motivations**

 B<sub>s</sub> mesons are produced in flavor eigenstates, but propagate as mass ones, which, if no CPV in the mixing, coincide with CP eigenstates

$$B_s^H o \mathsf{CP} \; \mathsf{odd} \qquad B_s^L o \mathsf{CP} \; \mathsf{even}$$

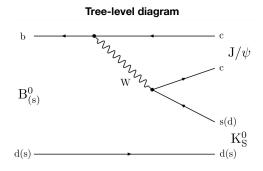
If the two eigenstates have different lifetimes (as for the B<sub>s</sub>), we can relate the mass eigenstate rate asymmetry A<sub>λr</sub>, with the CPV observable λ

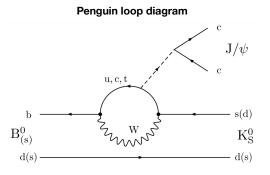
$$A_{\Delta\Gamma} = \frac{R_H - R_L}{R_H + R_L} = \frac{-2\,\mathcal{R}(\lambda)}{1 + |\lambda|^2}$$

R<sub>H</sub> and R<sub>L</sub>: coefficients in the untagged decay rate

$$\Gamma(B_s \to f) + \Gamma(\overline{B}_s \to f) = R_H e^{-\Gamma_H t} + R_L e^{-\Gamma_L t}$$

- This presentation is about a measurement of the B<sub>s</sub> effective lifetime  $\tau$  in the CP-odd final state J/ $\psi$  K<sub>s</sub> performed with the CMS Run 2 data set
- This process is related to  $B^0 \rightarrow J/\psi K_S$  via U-spin flavor symmetry
  - $\circ$  A<sub>AC</sub> can be used to determine penguin contributions to the measurement of sin(2 $\beta$ )
  - The measurement can also probe the CKM angle γ





### The effective lifetime

The effective lifetime is defined as the expected value of the untagged decay rate

$$\tau(J/\psi \, K_S) \equiv \frac{\int_0^\infty t(\Gamma_{B_s \to J/\psi K_S} + \Gamma_{\overline{B}_s \to J/\psi K_S}) dt}{\int_0^\infty (\Gamma_{B_s \to J/\psi K_S} + \Gamma_{\overline{B}_s \to J/\psi K_S}) dt} = \frac{\tau_{B_s}}{1 - y_s^2} \left(\frac{1 + 2A_{\Delta\Gamma} y_s + y_s^2}{1 + A_{\Delta\Gamma} y_s}\right)$$
Normalized decay width difference  $y_s = \tau_{B_s} \Delta\Gamma/2$ 

Using the latest measurements and assuming the SM  $(A_{\Lambda\Gamma} = 0.94 \pm 0.07, \tau_{Bs} = 1.520 \pm 0.005 \text{ ps}, \Delta\Gamma = 0.084 \pm 0.005 \text{ ps}^{-1})$ 

$$\tau(J/\psi K_S)|_{SM}$$
 = 1.62 ± 0.02 ps

- Available measurement from LHCb:  $\tau(J/\psi K_s) = 1.75 \pm 0.14$  ps [Nucl.Phys.B(2013)873]
- In this analysis the decay time is measured in the transverse plane as

$$t = \frac{L_{xy} \cdot M_{B_S}}{p_T}$$

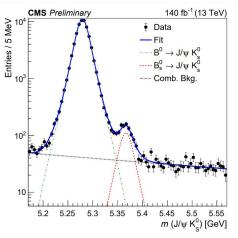
### Fit and results

- The effective lifetime is measured with a 2D UML fit to the invariant mass and proper decay time
- Background sources
  - $B^0 \rightarrow J/\psi K_s$ : irreducible, treated as a control channel
  - Combinatorial: suppressed with dedicated BDT selection
  - B<sup>0</sup> → J/ψ K\*<sup>0</sup>: negligible
  - $\circ$  J/ $\psi$   $\Lambda^0$ : suppressed with constraints on the decay kinematics
- Result (using 727 ± 35 B<sub>s</sub> signal candidates)

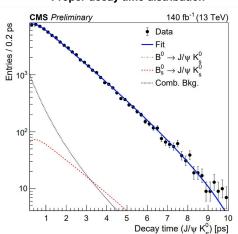
$$au(J/\psi\, K_{\mathcal{S}})^{\mathit{eff}}$$
 = 1.59  $\pm$  0.07 (stat)  $\pm$  0.03 (syst) ps

- The control channel's effective lifetime is found to be in good agreement with the world-average value
- The measured  $B_s \rightarrow J/\psi K_s$  effective lifetime is in agreement with the SM prediction and compatible with the previous LHCb results at 2.1 $\sigma$
- This is the most precise measurement of this quantity to date

### Invariant mass distribution



### Proper decay time distribution



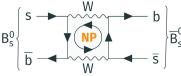
# Measurement of the time-dependent CP violation in B<sub>s</sub> mesons

CMS PAS BPH-23-004

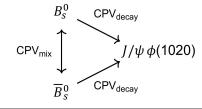
**Dataset**: 2017-18 (96 fb<sup>-1</sup>)

# **Motivations**

- B<sub>s</sub> mesons decays allow us to study the time-dependent
   CP violation generated by the interference between direct decays and flavor mixing
  - CPV in the interference is possible even if there is no CPV in decay and mixing
- The weak phase φ<sub>s</sub> is the main CPV observable
  - Predicted by the SM to be  $\phi_s \approx -2\beta_s = -37 \pm 1 \text{ mrad } (\text{ICKMfitter}, \text{UTfit})$  $\beta_s \rightarrow \text{ angle of the } B_s \text{ unit. triangle}$
- New physics can change the value of φ<sub>s</sub> up to ~100% via new particles contributing to the flavor oscillations [RMP88(2016)045002]

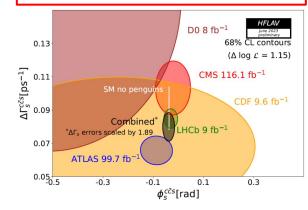


- ullet  $\phi_{\rm s}$  has been **first measured** by the **Tevatron** experiments D0 and CDF
- At LHC φ<sub>s</sub> has been measured several times by ATLAS, LHCb, and CMS
- This presentation is about the latest CMS results with the *golden* channel  $B_s \rightarrow J/\psi \ \phi(1020) \rightarrow \mu^+\mu^- K^+K^-$



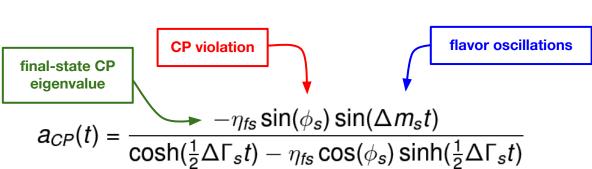
$$\Gamma\left(B_{S_{(w,\overline{B}_{S}^{0})}}^{0} \to f\right)(t) \stackrel{?}{\neq} \Gamma\left(\overline{B}_{S_{(w,\overline{B}_{S}^{0})}}^{0} \to f\right)(t)$$

 $\begin{aligned} a_{\text{CP}}(t) &\propto \Gamma_{\overline{\mathsf{B}}_{\mathsf{S}} \to \mathsf{f}}(t) - \Gamma_{\mathsf{B}_{\mathsf{S}} \to \mathsf{f}}(t) \\ &\propto -\eta_{\mathsf{fS}} \sin(\phi_{\mathsf{S}}) \sin(\Delta m_{\mathsf{S}} t) \end{aligned}$ 



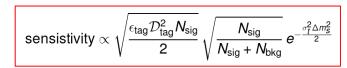
From: [Jevtic and Li, CERN seminar (2023)]

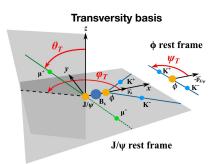
# A time-, flavor- and angular-dependent measurement



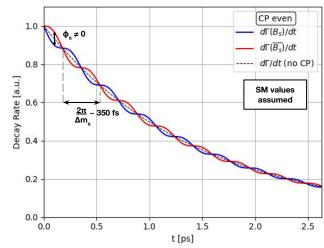
### **Core ingredients**

- Time-dependent angular analysis to separate the CP eigenstates ("transversity basis" used)
- Time-dependent flavor analysis to resolve the B<sub>s</sub> mixing oscillations (T ~ 350 fs, CMS  $\sigma_{t}$  ~ 65 fs)





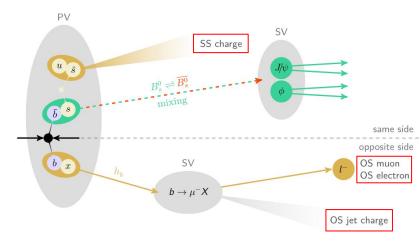
### Decay rate for a CP-even final state



# Flavor tagging overview

- A cutting-edge flavor tagging framework has been engineered to extract the best possible results from data
- Four DNN-based algorithms are used, divided into two main categories
  - Opposite side (OS): exploits decay products of the other B hadron in the event
    - **1.** OS muon: leverages  $b \rightarrow \mu^{-}X$  decays
    - **2. OS electron**: leverages  $b \rightarrow e^{-}X$  decays
    - **3. OS jet**: capitalizes on charge asymmetries in the OS *b*-jet
  - o Same side (SS): exploits the B<sub>s</sub> fragmentation
    - **4. SS tagger**: leverages charge asymmetries in the B<sub>s</sub> fragmentation

### Schematic representation of a generic event



### Useful definitions

$$\xi_{tag} = \begin{cases} +1 & \text{for } B_s \\ -1 & \text{for } \overline{B}_s \end{cases}$$

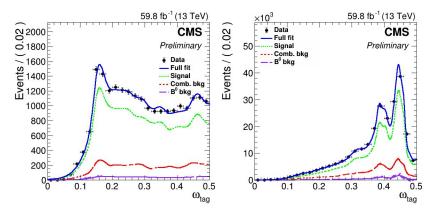
$$0 & \text{if no tagging decision is made}$$

$$\epsilon_{tag} = \frac{N_{tag}}{N_{tot}}, \quad \omega_{tag} = \frac{N_{mistag}}{N_{tag}}, \quad \mathcal{D}_{tag} = 1 - 2\omega_{tag}, \quad P_{tag} = \epsilon_{tag}\mathcal{D}_{tag}^2$$

# Flavor tagging performance

- The SS and any one of the OS algorithms overlap in about 20% of the events
  - In these cases, the information is combined to further improve the tagging inference
- The combined flavor tagging framework achieves
   a tagging power of P<sub>tag</sub> = 5.6% when applied to the
   B<sub>s</sub> data sample
  - Among the highest ever recorded at LHC
  - x3~4 improvement with respect to prev. CMS results
- This is the first CMS implementation of the OS jet and same-side tagging techniques
  - SS accounts for half of the performance

### $\omega_{tag}$ distribution in the *muon-tagging* trigger category (left) and the *standard* one (right) for 2018 data



Flavor tagging performance (mutually exclusive categories)

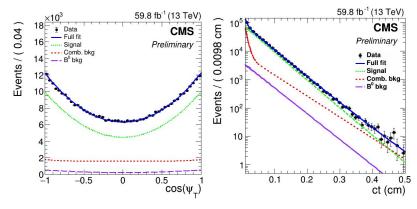
Category	$\varepsilon_{\mathrm{tag}}$ [%]	$\mathcal{D}_{ ext{eff}}^2$	$P_{\mathrm{tag}}$ [%]
Only OS muon	$6.07 \pm 0.05$	0.212	$1.29 \pm 0.07$
Only OS electron	$2.72 \pm 0.02$	0.079	$0.214 \pm 0.004$
Only OS jet	$5.16 \pm 0.03$	0.045	$0.235 \pm 0.003$
Only SS	$33.12 \pm 0.07$	0.080	$2.64 \pm 0.01$
SS + OS muon	$0.62 \pm 0.01$	0.202	$0.125 \pm 0.003$
SS + OS electron	$2.77 \pm 0.02$	0.150	$0.416 \pm 0.005$
SS + OS jet	$5.40 \pm 0.03$	0.124	$0.671 \pm 0.006$
Total	$55.9 \pm 0.1$	0.100	$5.59 \pm 0.02$

# Fit strategy

- The physics parameters are extracted with unbinned multidimensional extended maximum-likelihood (UML) fit
  - Physics parameters:  $\phi_s$ ,  $|\lambda|$ ,  $\Delta\Gamma_s$ ,  $\Gamma_s$ ,  $\Delta m_s$ ,  $|A_0|^2$ ,  $|A_\perp|^2$ ,  $|A_S|^2$ ,  $\delta_{\parallel}$ ,  $\delta_{\perp}$ ,  $\delta_{S\perp}$
  - $\circ$  Observables:  $m_{Bs}$ , t,  $\sigma_{t}$ ,  $\cos \theta_{T}$ ,  $\cos \psi_{T}$ ,  $\phi_{T}$ ,  $\omega_{tag}$
- Fit model

$$\frac{P(t,\sigma_t,\Theta,\xi_{tag},\omega_{tag},m\mid\alpha)}{|\epsilon(t)|} = \left[ \boxed{\Gamma(t,\Theta,\xi_{tag},\omega_{tag}\mid\alpha)} \otimes \boxed{G(t\mid\sigma_t)} \right] \cdot \underbrace{\epsilon(\Theta)} \cdot P(\sigma_t)P(m)P(\omega_{tag}) + P_{bkg}(...)$$

- Analytical decay rate
- Time resolution (extracted from prompt background)
- Angular efficiency (extracted from MC)
- **Time efficiency** (from  $B^0 \rightarrow J/\psi K^*$  events in data)
  - Implemented as reweighting
- Backgrounds sources:
  - 75%: combinatorial
  - $\sim$  25%: B<sup>0</sup>  $\rightarrow$  J/ψ K\*  $\rightarrow$  μμ Kπ
  - $\circ$  negligible: Λ<sub>b</sub> → J/ψ Λ<sup>0</sup> → μμ Kp (treated as systematic uncertainty)
- The statistical uncertainties and fit bias are estimated with 1300 bootstrap distributions



# Results

### Fit results

Parameter	Fit value	Stat. uncer.	Syst. uncer.
$\phi_s$ [mrad ]	-73	$\pm 23$	±7
$\Delta\Gamma_s$ [ps <sup>-1</sup> ]	0.0761	$\pm 0.0043$	$\pm 0.0019$
$\Gamma_s$ [ps <sup>-1</sup> ]	0.6613	$\pm 0.0015$	$\pm 0.0028$
$\Delta m_s  [\hbar \mathrm{ps}^{-1}]$	17.757	$\pm 0.035$	$\pm 0.017$
$ \lambda $	1.011	$\pm 0.014$	$\pm 0.012$
$ A_0 ^2$	0.5300	$\pm 0.0016$	$\pm 0.0044$
$ A_{\perp} ^2$	0.2409	$\pm 0.0021$	$\pm 0.0030$
$ A_{\rm S} ^2$	0.0067	$\pm 0.0033$	$\pm 0.0009$
$\delta_{\parallel}$	3.145	$\pm 0.074$	$\pm 0.025$
$\delta_{\perp}^{"}$	2.931	$\pm 0.089$	$\pm 0.050$
$\delta_{\mathrm{S}\perp}$	0.48	$\pm 0.15$	$\pm 0.05$

•  $\phi_s$  and  $\Delta\Gamma_s$  are found in agreement with the SM

$$\phi_s^{SM} \simeq -37 \pm 1 \; \mathrm{mrad} \qquad \Delta \Gamma_s^{SM} = 0.091 \pm 0.013 \; \mathrm{ps^{-1}}$$

•  $\Gamma_s$  and  $\Delta m_s$  are consistent with the latest world averages

$$\Gamma_s^{WA} = 0.6573 \pm 0.0023 \text{ ps}^{-1}$$
  $\Delta m_s^{WA} = 17.765 \pm 0.006 \text{ } \hbar \text{ps}^{-1}$ 

- $|\lambda|$  is consistent with no direct CPV ( $|\lambda| = 1$ )
- This measurement utilizes the largest ever effective statistics  $N_{Bs} \cdot P_{taq}$  for a single  $\phi_s$  measurement (~27.5k)
  - ο The precision on  $φ_s$  is comparable with the world's most precise single measurement by LHCb ( $φ_s$  = -39 ± 22 (stat) ± 6 (syst) mrad) [PRL132(2024)051802]
  - $\circ$  This is the most precise single measurement of  $\Delta \Gamma_s$  to date in this channel

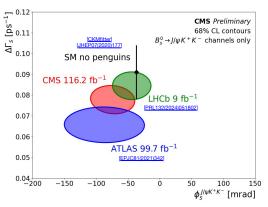
# **Combination with 8 TeV results**

 These results supersede <u>PLB816(2021)136188</u> and are further combined with those obtained CMS at 8 TeV [PLB757(2016)97], yielding

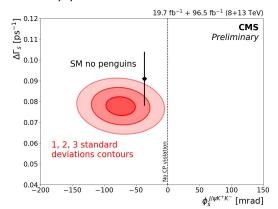
$$\phi_{s}$$
 =  $-74 \pm 23$  [mrad]  $\Delta\Gamma_{s}$  = 0.0780  $\pm$  0.0045 [ps $^{-1}$ ]

- Due to the high difference in statistical power between the two results the sensitivity gain is small
- The combined value for the weak phase  $\phi_s$  is consistent with the SM prediction, the latest world average, and with zero (no CPV) at 3.2 s.d.
  - This is the first evidence of CPV in  $B_s \rightarrow J/\psi K^+K^-$  decays
- These results helps to further constrain possible BSM effects in the B<sub>s</sub> system

### Comparison with other LHC experiments



### 1, 2, 3 standard deviations contours



# **Outlook**

# **Summary and outlook**

- This presentation showed three recent CMS results on the physics of CP violation
  - CP violation in D<sup>0</sup> → K<sub>s</sub>K<sub>s</sub>
    - First CMS results on CP violation in the charm sector
  - Effective lifetime measurement in the CP-odd decay B<sub>s</sub> → J/ψ K<sub>s</sub>
    - Most precise determination of  $\tau_{eff}(B_s \rightarrow J/\psi K_s)$
  - Measurement of the time-dependent CP violation in B<sub>κ</sub> → J/ψ φ
    - First evidence of CP violation in B<sub>s</sub> → J/ψ K<sup>+</sup>K<sup>-</sup>
- CMS recent contributions in flavor physics prove that it can be one of the leading actors in several key areas of study, such as rare decays and CP violation
- Thanks to the advancements in trigger strategies and flavor tagging techniques, CMS is able to compete in measurements for which the detector was not designed
- Run 3 will provide unique opportunities thanks of a revamped trigger strategy, which will lead
  to the collection of an unprecedented amount of data suitable for flavor physics studies

Stay tuned in the future for other exciting CMS results!

# **Backup**

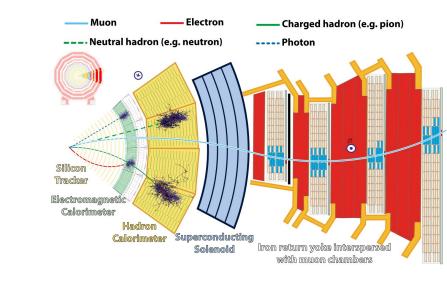
# The CMS detector

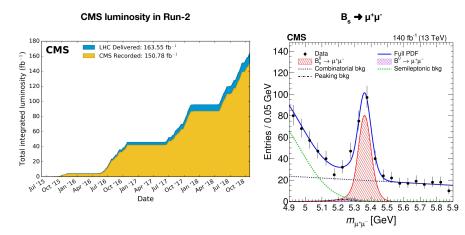
# CMS is a general purpose detector able to perform a vast range of physics studies, including flavor physics

- Excellent tracking system able to reconstruct vertices with high decay time resolution
   (e.g., σ, ~ 65 fs for B<sub>s</sub> → J/ψ φ) up to |η| < 2.5</li>
  - Complementary to LHCb (2 <  $|\eta|$  < 5)
- + Enormous amount of data collected
  - ~ 7.5 · 10<sup>13</sup> bb pairs produced at Point 5 during Run 2 (geometric acceptance not considered)
- High pile up  $N_{PV} \sim 40$  (in Run 2)
- No reliable hadronic particle identification available

### Some CMS flavor physics highlights from recent years

- B<sub>s</sub>  $\rightarrow \mu^+\mu^-$  (world's most precise) [PLB842(2023)137955]
- $\eta \rightarrow \mu^+\mu^-\mu^+\mu^-$  observation [PRL131(2023)091903]
- f<sub>s</sub>/f<sub>II</sub> measurements [PRL131(2023)121901]
- Triple J/ψ production observation [Nat.Phys.19(2023)338]
- R(K) LFU test [BPH-22-005]
- R(J/ψ) LFU test [BPH-22-012]



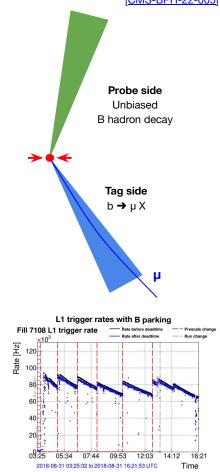


# Backup

- CP violation in D<sup>0</sup> -

# The CMS B parking dataset

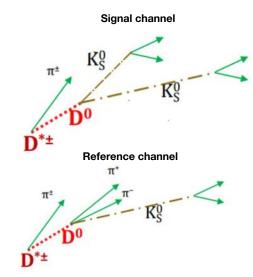
- Designed to allow CMS to perform B physics measurements on difficult/impossible to trigger final states (e.g. fully hadronic final states)
- Achieved with a set of single muon triggers (tags) with different thresholds in p<sub>T</sub> and impact parameter
  - Luminosity decreases during a run → less restrictive triggers enabled
    - Maximises the available trigger bandwidth
  - Events are parked for later reconstruction
  - Very high purity of ~80%
- No impact on the standard CMS physics programme
- 10 billion unbiased B hadron decays collected in 2018 (L<sub>int</sub> ~ 41 fb<sup>-1</sup>)

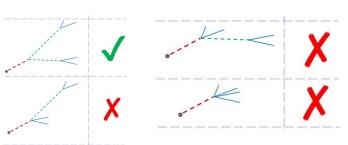


# **Event selection**

- First,  $K_s \rightarrow \pi^+\pi^-$  are reconstructed fitting the  $\pi$  tracks to a common vertex
  - $| m(\pi^+\pi^-) m(K_s^{\text{w.a.}}) | < 20 \text{ MeV}, p_{\tau}(K_s) > 2.2(1.0) \text{ GeV}$
- In the signal channel, two K<sub>S</sub> candidates are required and fitted to a common vertex to form D<sup>0</sup> → K<sub>S</sub>K<sub>S</sub> candidates
  - 1.7 GeV < m(K<sub>S</sub>K<sub>S</sub>) < 2.0 GeV</li>
  - $K_s$  displacement in xyz from the D<sup>0</sup> vertex >9 $\sigma$  and >7 $\sigma$
  - o D<sup>0</sup> displacement in xyz (xy) from the PV >9 $\sigma$  (>2 $\sigma$ )
- In the **reference channel**, two track with  $p_T > 0.6$  GeV are used to form the  $D^0 \rightarrow K_S \pi^+ \pi^-$  candidate
  - 1.823 < m(K<sub>S</sub>π⁺π⁻) < 1.908 GeV</li>
- **Finally**, an additional track with -1.2 <  $|\eta|$  < 1.2 and  $p_T$  > 0.36 GeV is added to form  $D^{*+} \rightarrow D^0 \pi^+$  candidates
  - $\qquad \qquad m(D^0 \, \pi^+) = m(D^0 \pi^+) m(D^0) + m_{PDG}(D^0)$

 Background suppression: several fits corresponding to incorrect topologies are performed and vertex probabilities requirements are imposed





# **Selection**

Table 1: Optimized selection criteria in the signal channel  $D^0 \to K^0_S K^0_S.$ 

Variable	Requirement
$p_{\rm T}$ of tagging pion from ${\rm D}^{*\pm} \to {\rm D}\pi^{\pm}$	> 0.35 GeV
$\eta$ of tagging pion from $D^{*\pm} \to D\pi^{\pm}$	$-1.2 < \eta < 1.2$
$p_{\mathrm{T}}(\mathrm{K}^0_{\mathrm{S}})$	> 2.2  GeV and $> 1.0  GeV$
$P_{vtx}(\mathrm{D}\pi^{\pm})$	> 5%
$P_{vtx}(K_S^0K_S^0)$	> 1%
$P_{vtx}(\pi^+\pi^-)$ for $K_S^0 \to \pi^+\pi^-$	> 1%
$D^0$ vertex displacement from the PV in $xy$	$> 2 \mathrm{s.d.}$
$D^0$ vertex displacement from the PV in $xyz$	> 9 s.d.
$K_S^0$ vertex displacement from the $D^0$ vertex in $xyz$	> 9  s.d. and $> 7  s.d.$
angle between $D^0$ momentum and displacement from PV in $xyz$	< 0.205 rad
angle between $D^0$ momentum and displacement from PV in $xy$	< 0.237 rad
angle between $D^0$ momentum and displacement from BX in $xy$	< 0.237 rad

# **Backup**

- B<sub>s</sub> → J/ψ K<sub>s</sub> effective lifetime -

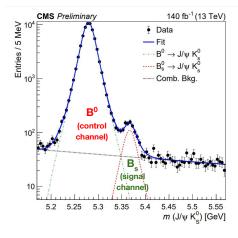
# **Event selection and efficiency**

- **Trigger**:  $J/\psi \rightarrow \mu^+\mu^-$  candidate with  $p_{\tau} > 20$  (25) GeV for 2016 (2017-18)
- Offline K<sub>s</sub> → π<sup>+</sup>π<sup>-</sup> selection:
  - $\circ$  Displaced by >15 $\sigma$  from the beamspot and >5 $\sigma$  from the B<sub>s</sub> vertex
  - Invariant mass within 70 MeV from world-average value
- Background sources
  - $\land \land \Rightarrow p\pi^-$ : suppressed with constraints on the decay kinematics
  - $\circ$  B<sup>0</sup>  $\rightarrow$  J/ψ K<sub>s</sub>: irreducible, treated as a control channel
  - $B^0 \rightarrow J/\psi K^{*0}$ : negligible
  - Combinatorial: suppressed with dedicated BDT selection
- **Time efficiency** is measured in simulations for B<sub>s</sub> and B<sup>0</sup> (control channel)

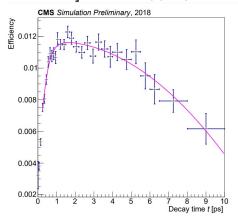
$$\epsilon(t) = \frac{t_{reco}}{t_{gen} \otimes \delta(t)}$$

Modeled with a combination of polynomials and logistic functions

### Invariant mass distribution



### B<sub>a</sub> time efficiency (2018)



# Backup

- CP violation in B<sub>s</sub> -

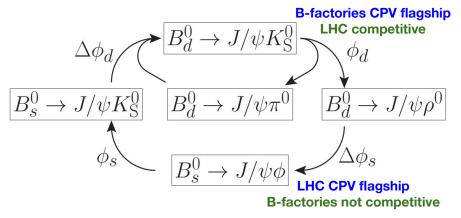
# **Penguin contributions**

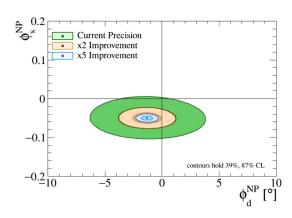
Assuming this is negligible

Penguin pollutions are expected to be small for B<sub>s</sub>, but they are not well constrained

$$\Delta\phi_s^{\mathsf{penguin}} pprox 3 \pm 10 \; \mathsf{mrad}$$

Analysis of penguin and NP contributions is possible using Cabibbo-favored control channels





# **Decay time and its resolution**

 The time dependence of the decay rate is parametrized with the proper decay length ct, measured in the transverse plane as

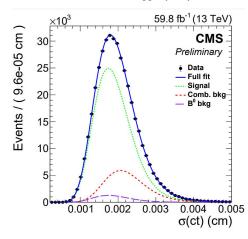
$$ct = c \cdot \frac{m_{Bs}^{w.a.} \cdot L_{xy}}{p_T}$$
 with  $L_{xy} \equiv ||\overline{r}_{xy}(SV) - \overline{r}_{xy}(PV)||$ 

- Its uncertainty is obtained by fully propagating the uncertainties in L<sub>xv</sub> and p<sub>T</sub>
  - The uncertainty on L<sub>xy</sub> dominates for most of the ct spectrum, with σ(p<sub>τ</sub>) taking over at high values (ct ≥ 3 mm)
- The ct uncertainty is calibrated in a prompt data sample of  $B_s \rightarrow J/\psi \phi$ , obtained by removing the displacement requirement in the *muon-tagging* data sets
  - Modeled with two gaussians to obtain the effective dilution and resolution

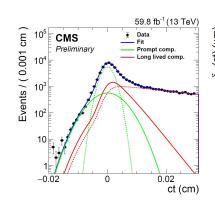
$$\delta_{\text{eff}} = \sqrt{\frac{-2 \ln \mathcal{D}}{\Delta m_s^2}}$$
 with  $\mathcal{D} = \sum_{i=1}^2 f_i \exp\left(-\frac{\sigma_i \Delta m_s^2}{2}\right)$ 

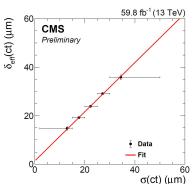
• Excellent agreement found, with corrections ~5%

### Proper decay length uncertainty distribution for the standard trigger (2018)



### Time resolution calibration for 2018 data





# **Acceptance and efficiency effects**

- The efficiency in selecting and reconstructing the B<sub>s</sub> candidates is not independent of the decay time and angular observables
  - To properly fit the decay rate model an efficiency parametrization is needed

### Time efficiency

- Modeled in the  $B^0 \rightarrow J/\psi K^{*0}$  data control channel with corrections from simulations
- Ultimately parametrized with Bernstein's polynomials

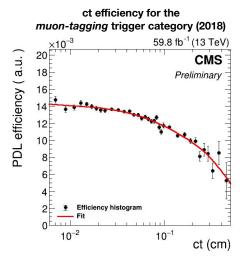
$$\varepsilon_{B^0}^{\mathrm{data}}(ct) = \frac{N_{B^0}(ct)}{e^{-\Gamma_d^{\mathrm{w.a.}}} \otimes P_{B^0}(\sigma_{ct})} \qquad \qquad \varepsilon_{B_s}^{\mathrm{data}}(ct) = \varepsilon_{B^0}^{\mathrm{data}}(ct) \cdot \frac{\varepsilon_{B_s}^{\mathrm{MC}}(ct)}{\varepsilon_{B^0}^{\mathrm{MC}}(ct)}$$

### Angular efficiency

- Estimated with KDE distributions in simulated events
- The simulated data samples are corrected to match the data
  - An iterative procedure is used to simultaneously correct the kinematics of the final state particles and the differences in the physics parameters set in the MC with respect to what measured in the data

# standard trigger category (2018) x10<sup>-3</sup> 59.8 fb<sup>-1</sup> (13 TeV) CMS Preliminary 14 10 8 6 4 2 Efficiency histogram 10<sup>-2</sup> 10<sup>-1</sup> ct (cm)

ct efficiency for the



# **Decay rate model**

### Flavor tag decision

(flips c, and d, signs)

Mistag probability

terms for SM φ

-S

 $D\cos(\delta_{\perp} - \delta_{\parallel})$ 

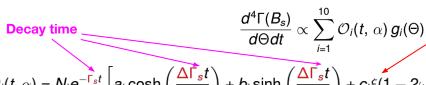
 $-S\cos(\delta_{\parallel}-\delta_{0})$ 

 $D\cos(\delta_{\perp}-\delta_0)$ 

 $D\sin(\delta_{\parallel}-\delta_{\rm S})$ 

 $S\sin(\delta_1 - \delta_S)$ 

 $D\sin(\delta_0 - \delta_S)$ 



# $\mathcal{O}_{i}(t,\alpha) = N_{i}e^{-\Gamma_{s}t}\left[a_{i}\cosh\left(\frac{\Delta\Gamma_{s}t}{2}\right) + b_{i}\sinh\left(\frac{\Delta\Gamma_{s}t}{2}\right) + c_{i}\xi(1-2\omega)\cos(\Delta m_{s}t) + d_{i}\xi(1-2\omega)\sin(\Delta m_{s}t)\right]$

### **Angular** variables

i	$g_i(\theta_T, \psi_T, \varphi_T)$
1	$2\cos^2\psi_T(1-\sin^2\theta_T\cos^2\varphi_T)$
2	$\sin^2 \psi_T (1 - \sin^2 \theta_T \sin^2 \varphi_T)$
3	$\sin^2\psi_T\sin^2\theta_T$
4	$-\sin^2\psi_T\sin2 heta_T\sinarphi_T$
5	$\frac{1}{\sqrt{2}}\sin 2\psi_T\sin^2\theta_T\sin 2\varphi_T$
6	$\frac{1}{\sqrt{2}}\sin 2\psi_T\sin 2\theta_T\cos \varphi_T$
7	$\frac{2}{3}(1-\sin^2\theta_T\cos^2\varphi_T)$
8	$\frac{1}{3}\sqrt{6}\sin\psi_T\sin^2\theta_T\sin2\varphi_T$
9	$\frac{1}{3}\sqrt{6}\sin\psi_T\sin 2\theta_T\cos\varphi_T$
10	$\frac{4}{3}\sqrt{3}\cos\psi_T(1-\sin^2\theta_T\cos^2\varphi_T)$

$_T\cos^2\varphi_T)$
$\sin^2 \varphi_T$
$\theta_T$
$\sin \varphi_T$
$\sin 2\varphi_T$
$T \cos \varphi_T$
$s^2 \varphi_T$ )

$$\frac{2}{3}(1-\sin^2\theta_T\cos^2\varphi_T)$$

$$\frac{1}{3}\sqrt{6}\sin\psi_T\sin^2\theta_T\sin 2\varphi_T$$

$$\frac{1}{3}\sqrt{6}\sin\psi_T\sin 2\theta_T\cos\varphi_T$$

$$\frac{1}{3}\sqrt{3}\cos\psi_T(1-\sin^2\theta_T\cos^2\varphi_T)$$

2		
2		

 $|A_0(0)|^2$ 

 $|A_{||}(0)|^2$  $|A_{\perp}(0)|^2$  $|A_{\parallel}(0)||A_{\perp}(0)|$ 

 $|A_0(0)||A_{\parallel}(0)|$ 

 $|A_0(0)||A_{\perp}(0)|$ 

 $|A_{\rm S}(0)|^2$ 

 $k_{SP}|A_{S}(0)||A_{\parallel}(0)|$ 

 $k_{SP} |A_{S}(0)| |A_{\perp}(0)|$ 

 $k_{SP}|A_{S}(0)||A_{0}(0)|$ 

Sensitive to direct CPV

$$S = -\frac{2|\lambda|\sin\phi_s}{1+|\lambda|^2}$$

Sensitive to φ ~ 0

 $C\sin(\delta_{\perp} - \delta_{\parallel})$ 

 $\cos(\delta_{\parallel} - \delta_{0})$ 

 $C\sin(\delta_{\perp}-\delta_0)$ 

 $C\cos(\delta_{\parallel} - \delta_{\varsigma})$ 

 $\sin(\delta_1 - \delta_S)$ 

 $C\cos(\delta_0 - \delta_S)$ 

$$D = -\frac{2|\lambda|\cos\phi_s}{1+|\lambda|^2}$$

 $-D\sin(\delta_{\perp}-\delta_{S})$   $C\sin(\delta_{\perp}-\delta_{S})$ 

 $S\cos(\delta_{\parallel}-\delta_{\parallel})$ 

 $D\cos(\delta_{\parallel}-\delta_{0})$ 

 $S\cos(\delta_1-\delta_0)$ 

 $S\sin(\delta_{\parallel}-\delta_{\rm S})$ 

 $S\sin(\delta_0 - \delta_S)$ 

Sensitive to  $\phi_{\rm s} \sim \pi/2$ 

 $\sin(\delta_{\perp} - \delta_{\parallel})$ 

 $C\cos(\delta_{\parallel}-\delta_{0})$ 

 $\sin(\delta_{\perp} - \delta_0)$ 

 $\cos(\delta_{\parallel} - \delta_{\rm S})$ 

 $\cos(\delta_0 - \delta_S)$ 

### **Conventions**

- $\begin{aligned} |A_{/\!/}|^2 &= |A_0|^2 |A_{\perp}|^2 \\ \delta_0 &= 0 \\ \delta_{S\perp} &= \delta_S \delta_{\perp} \end{aligned}$

- $\Delta\Gamma_{\rm s} > 0$

### **Physics parameters**

- $\phi_s$ ,  $|\lambda|$

### S-P wave effective coupling k<sub>ep</sub> ≈ 0.54

- Introduced since m(K+K-) is not fitted
- Evaluated from the S- and P-wave lineshape interference

# **Trigger strategy**

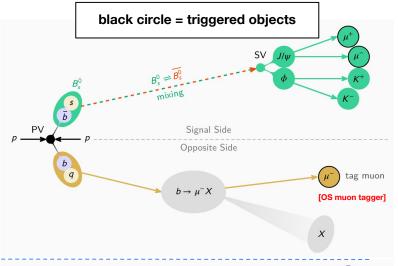
### **Muon-tagging** trigger

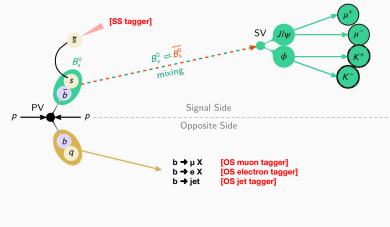
- $J/\psi \rightarrow \mu^+\mu^-$  candidate plus an additional muon (for tagging)
- ≈50 000 signal candidates
- Used for time resolution modeling
- Tagging algorithms deployed: OS-muon
  - $\circ$   $P_{tag} \sim 10\%$  (muon at trigger level enhance tagging efficiency)

### Standard trigger

- Displaced J/ψ → μ<sup>+</sup>μ<sup>-</sup> candidate + φ(1020) → K<sup>+</sup>K<sup>-</sup>
- ≈450 000 signal candidates
- Tagging algorithms deployed: OS-muon, OS-electron, OS-jet, Same Side
  - P<sub>taq</sub> ~ **5%**







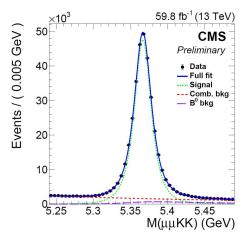
### **Dataset and selection**

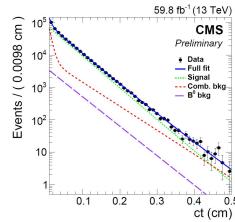
- **Dataset**: L<sub>int</sub> = 96 fb<sup>-1</sup> collected in 2017-2018
  - Why no 2016 data? Very different data set (old inner tracker detector with worse time resolution and different trigger menu)
- Signal candidates: 491 270 ± 950
- Notable selection requirements:

Variable	Requirement
ct ( <i>muon-tagging</i> HLT)	$>$ 60 $\mu$ m
ct (standard HLT)	> 100 $\mu$ m
$ct/\sigma_ct$ (standard HLT)	> 3
$ m(K^+K^-) - m_{\phi(1020)} $	< 10 MeV
$ m(\mu^+\mu^-) - m_{J/\psi} $	< 150 MeV

- To avoid overlaps, events that pass both trigger category selections are placed only in the muon-tagging one
  - This depletes the standard trigger category of OS muons
- The PV of choice is the closest in 3D to the line that passes through the SV and parallel to the B<sub>g</sub> momentum

### Invariant mass and proper decay length distributions for the standard trigger (2018)





### Flavor, neural networks, and probabilities

- The tagging inference logic differs between algorithms
  - Lepton taggers (OS muon, OS electron)

0

- Lepton charge  $\rightarrow \xi_{tag}$ ; DNN score  $\rightarrow \omega_{tag}$  (DNN trained for correct-tag vs mistag)

  OS  $\ell^- \rightarrow$  OS  $b \xrightarrow{tag}$  signal  $B_s$ OS  $\ell^+ \rightarrow$  OS  $\overline{b} \xrightarrow{tag}$  signal  $\overline{B}_s$ DNN score

  Charge-based taggers (OS jet, SS)

  DNN score  $\rightarrow$  Prob(B<sub>s</sub>)  $\rightarrow \xi_{tag}$ ,  $\omega_{tag}$  (DNN trained for B<sub>s</sub> vs  $\overline{B}_s$ )  $s_{DNN} > 0.5 + \epsilon \xrightarrow{tag}$  signal  $B_s$  with  $\omega_{tag} = 1 \overline{s_{DNN}}$   $s_{DNN} < 0.5 \epsilon \xrightarrow{tag}$  signal  $\overline{B}_s$  with  $\omega_{tag} = 1 \overline{s_{DNN}}$ 
  - $\bullet$   $\bullet$  is used to remove events with  $\omega_{tag} \sim 50\%$
- The algorithms are optimized and trained in simulated events and calibrated in data with self-tagging
   B⁺ → J/ψ K⁺ decays
  - $\circ$   $\;$  The calibration is performed by comparing  $\omega_{tag}$  predicted by the DNN and the one measured in data

## **Calibration strategy (and other tricks)**

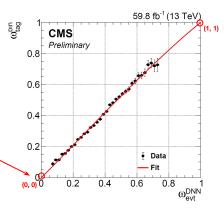
- A multi-pronged strategy has been devised to improve the ω<sub>tag</sub> estimation and suppress systematic effects
  - All models are constructed from the start as probability estimators, i.e. score~ω<sub>tag</sub>
    - Loss function: cross-entropy, which is the likelihood for the probability P(true class | score)
    - Output layer: Sigmoid function, which normalizes the output to a probability distribution
  - 2. All DNNs are calibrated with the *Platt scaling*, which ensures that the calibrated score is still a probability
    - The Platt scaling is a linear calibration of the score before the last sigmoid layer
  - 3. In calibrating the charge-based taggers (which provide a probability for  $B_s$  vs  $\overline{B}_s$ ):
    - A. The output is symmetrized due to the initial LHC charge imbalance

$$s_{DNN}^{sym}(x) = \frac{s_{DNN}(x) + [1 - s_{DNN}(\overline{x})]}{2}$$

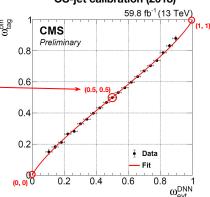
B. The symmetry is explicitly forced in the calibration function by removing the constant term

This strategy cancels almost all the systematic effects associated with flavor tagging

### OS-Muon calibration (muon-tagging trigger 2018)



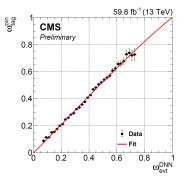
#### OS-jet calibration (2018)



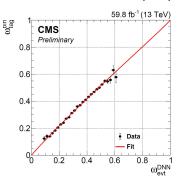
## **OS-lepton tagging**

- OS-lepton tagging techniques search for b → ℓ<sup>-</sup>X decays of the other B hadron in the event
- The charge of the lepton is used as tagging feature and a fully connected DNN is used to estimate the mistag probability
- Lepton selection
  - Loose kinematic cuts
  - Separated from the signal B meson
  - MVA discriminator against fakes
  - OS-electrons are searched only if no OS-muon is found in the event (explicit orthogonality)
- Mistag estimation
  - Fully connected DNN with ReLU activation and dropout
  - Inputs: lepton kinematics and surrounding activity
- Trained on simulated B<sub>s</sub> → J/ψ φ(1020) events and calibrated in B<sup>+</sup> → J/ψ K<sup>+</sup> data

#### OS-Muon calibration (muon-tagging trigger 2018)



#### **OS-Electron calibration (2018)**

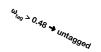


#### Schematic DNN model representation

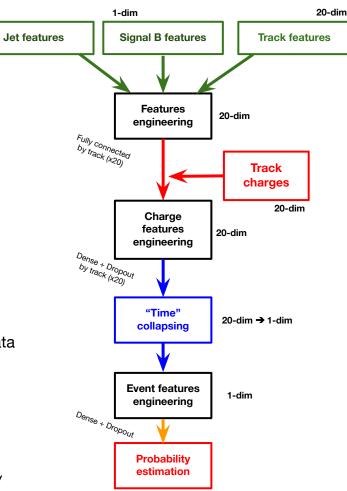
# **OS-jet tagging**

- The OS-jet algorithm exploits charge asymmetries in the jet structure and is based on a DNN called <u>DeepJetCharge</u>
  - Inputs: features from signal B meson, OS jet and its constituents
    - NB: The only flavor asymmetry is in the charges
  - Based on the DeepSets architecture <u>ref</u>
- Jet selection
  - No OS-lepton candidate
  - At least 2 tracks with |IP<sub>2</sub>| < 1 cm
  - Separated from the signal B meson
  - jet b-tagging discriminator
- Additional nearby tracks are used due to the poor jet clustering performance in the kinematic region of interest (p<sub>⊤</sub> < 20 GeV)</li>
- Trained on simulated  $B_s \rightarrow J/\psi \phi$  events and calibrated in  $B^+ \rightarrow J/\psi K^+$  data
- The trained network produces the probability of signal B meson containing a b̄ quark (i.e. being a B)
- The score is finally used to compute both  $\boldsymbol{\xi}_{tag}$  and  $\boldsymbol{\omega}_{tag}$

$$s_{DNN} > 0.52 \xrightarrow{\text{tag}} \text{ signal } B_s \text{ with } \omega_{tag} = 1 - s_{DNN} \ s_{DNN} < 0.48 \xrightarrow{\text{tag}} \text{ signal } \overline{B_s} \text{ with } \omega_{tag} = s_{DNN}$$



1-dim



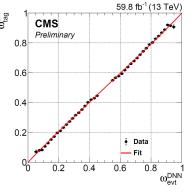
## **SS** tagger

- The SS tagger consists of a DNN (DeepSSTagger), derived from DeepJetCharge, able to probe the fragmentation products of a B meson and exploit tracks with high flavor correlation
- DeepSSTagger uses the kinematic information from up to 20 tracks (ordered by |IP<sub>z</sub>|) around the reconstructed B meson
- Track selection
  - $\circ$  ΔR(trk, B) < 0.8,  $|IP_{z}(PV)|$  < 0.4 cm,  $|IP_{xv}(PV)|/\sigma_{dxv}$  < 1
  - Overlap with signal and OS is carefully avoided with geometrical cuts and vetos
- Trained on an equal-weight mixture of  $B_s \rightarrow J/\psi \phi$  and  $B^+ \rightarrow J/\psi K^+$  to make the model invariant for  $B_s \leftrightarrow B^+$  for calibration purposes
  - Calibration directly in B<sub>s</sub> was found to be not feasible in CMS
    - Tested:  $B_{\epsilon}$  →  $D_{\epsilon}^{-}\pi^{+}$  (not enough stat.) and  $B_{\epsilon}^{**}$  →  $B^{+(*)}K^{-}$  (too much uncer. from  $B^{0**}$  bkg)
  - The trained network produces the probability of signal B meson containing a negatively charged quark alongside the b quark (i.e., being a B<sub>s</sub> or B<sup>-</sup>)

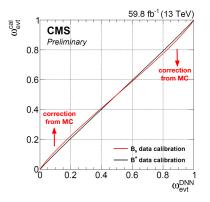
### Calibration

- The SS is calibrated B<sup>+</sup> → J/ψ K<sup>+</sup> data, with residual differences ~10% corrected with simulations
- $\circ$  Events with  $\omega_{\text{\tiny tag}} > 0.46$  are removed before the calibration and assumed untagged

### Same-side tagger calibration (B<sup>+</sup> data 2018)



### Comparison between Same-side tagger B<sup>+</sup> and B<sub>c</sub> calibrations (2018)

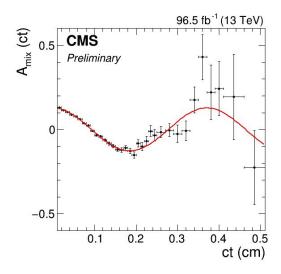


# Tagging validation with B<sup>0</sup> events

- The flavor tagging framework is validated in the B<sup>0</sup> → J/ψ K\*<sup>0</sup> control channel (~2M events)
- The time-dependent **mixing asymmetry** is measured to extract the flavor mixing oscillation frequency  $\Delta m_d$  with a precision of ~1% (comparable with BaBar and Belle)
  - Excellent agreement with world-averages is observed
    - No bias in mixing frequency measurements
- Study performed also in each tagging category (see backup)
- The time-integrated mixing is also measured for each tagger and their dependency on the expected tagging dilution is compared
  - The dependency between the measured A<sub>mix</sub> and the estimated D<sub>tag</sub> is found to be well described by a linear relationship, indicating that all four techniques behave in the same predictable way

#### B<sup>0</sup> flavor mixing asymmetry

$$A_{mix}(ct) = \frac{N_{unmix}(ct) - N_{mix}(ct)}{N_{unmix}(ct) + N_{mix}(ct)}$$



### Offline selection

### Requirements common between the two HLTs

- 5.24 < m(μμΚΚ) < 5.49 GeV</li>
- $p_{\tau}(B_s) > 9.5 \text{ GeV}$
- Vertex probability > 2%
- $\sigma(ct) < 50 \, \mu m$
- $|\eta(\mu)| < 2.4$
- $|\eta(K)| < 2.5$
- |m(μμ) m(J/ψ<sup>PDG</sup>)| < 150 MeV</li>
- $|m(KK) m(\phi(1020)^{PDG})| < 10 \text{ MeV}$

### Requirements specific to the muon-tagging HLT

- $p_{T}(\mu) > 3.5 \text{ GeV}$
- $p_{T}(K) > 1.15 \text{ GeV}$
- ct > 60 μm

### Requirements specific to the standard HLT

- muon-tagging trigger vetoed
- $p_{\tau}(\mu) > 4 \text{ GeV}$
- $p_T(K) > 0.9 \text{ GeV}$
- $p_{T}(\mu\mu) > 6.9 \text{ GeV}$
- ct > 100  $\mu$ m, ct/ $\sigma$ (ct) > 3

- Selection requirement optimized with the a genetic algorithm to maximize  $S/\sqrt{S+B}$
- To avoid overlaps, the muon-tagging trigger is vetoed in the standard trigger category
- The **PV** of choice is the closest in 3D to the line that passes through the SV and parallel to the B<sub>s</sub> momentum

## **OS-lepton taggers selection**

### **OS Muon**

### Requirements

- $\circ$  p<sub>T</sub> > 2 GeV
- $\circ$   $|\eta| < 2.4$
- |d<sub>x</sub>(PV)| < 1 cm</p>
- $\circ$   $\Delta R(B_s) > 0.4$
- Discriminators vs fakes

# Deployed in both trigger categories

- Dense DNN for  $\omega_{tag}$  estimation
  - Inputs: kinematics, IP, surrounding activity

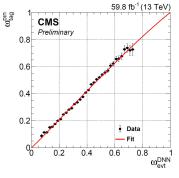
#### OS electron

### Requirements

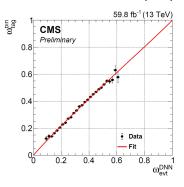
- No OS muon selected in the event
- $\circ$  p<sub>T</sub> > 2.5 GeV
- $\circ$   $|\eta| < 2.4$
- $\circ$   $|d_{z}(PV)| < 0.2 \text{ cm}$
- $|d_{xy}(PV)| < 0.08 \text{ cm}$
- $\circ$   $\Delta R(B_s) > 0.4$
- Discriminators vs fakes

- Deployed only in the standard trigger category
- Dense DNN for  $\omega_{tag}$  estimation
  - Inputs: kinematics, IP, surrounding activity

#### OS-Muon calibration (muon-tagging HLT 2018)



#### OS-Electron calibration (2018)



## **Taggers combination**

### Overlap logic

Overlap	OS muon	OS electron	OS jet	SS
OS muon		Χ	Χ	$\checkmark$
OS electron	X		X	1
OS jet	X	X		<b>√</b>
SŚ	✓	$\checkmark$	$\checkmark$	

### Tag decision combination

$$\xi(\xi_1, \xi_2, \omega_1, \omega_2) = \begin{cases} \xi_1 & \text{if } \omega_1 < \omega_2 \\ \xi_2 & \text{if } \omega_2 < \omega_1 \end{cases}$$

### Mistag combination

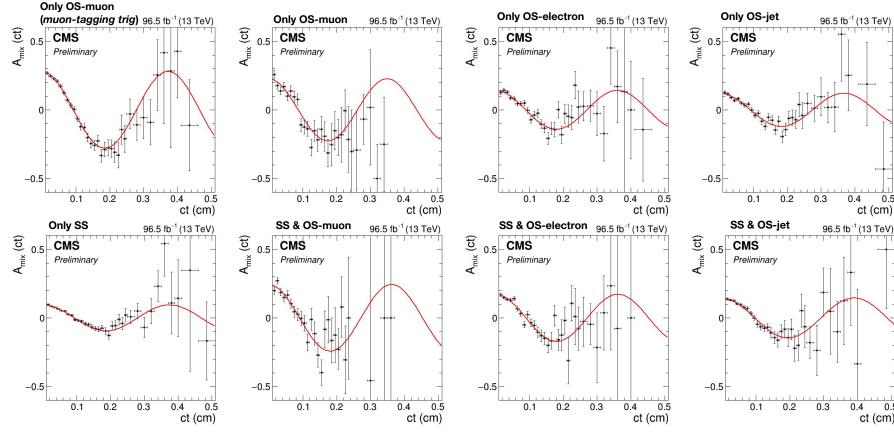
$$p(\overline{b}) = \prod_{i=1}^{2} \left( \frac{1 - \xi_i}{2} + \xi_i (1 - \omega_i) \right) \qquad p(b) = \prod_{i=1}^{2} \left( \frac{1 + \xi_i}{2} - \xi_i (1 - \omega_i) \right)$$

$$P(\overline{b}) = \frac{p(\overline{b})}{p(\overline{b}) + p(b)} \qquad P(b) = \frac{p(b)}{p(\overline{b}) + p(b)}$$

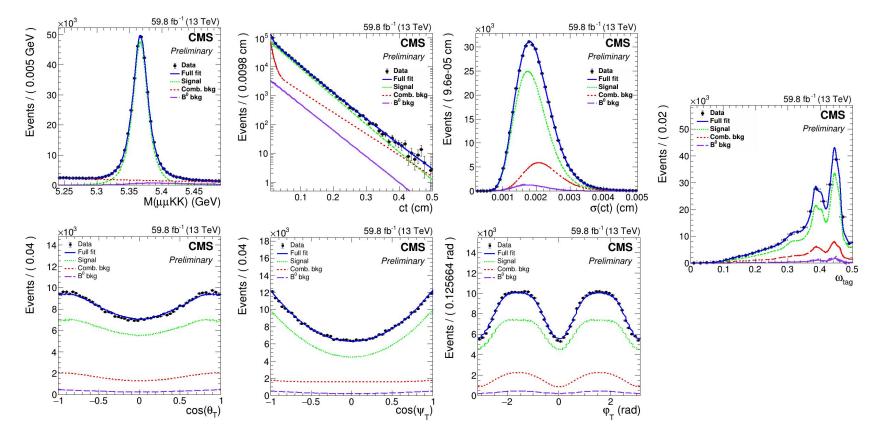
# Mixing asymmetry for different tagging categories

All, but the first, refers to the standard trigger category

ct (cm)



## Fit projections (standard trigger category 2018)



## Systematic uncertainty classification

- Type-I: unaccounted uncertainties
  - Account for the finite statistics of simulated/control samples and uncertainties in calibrations and efficiency
  - Always propagated to the final results
  - Evaluated with two procedures
    - <u>Type-I full</u>: obtained by sampling the samples/parameters of interest ~100 times, repeating the fit each time, and taking the RMS of the results as uncertainty
    - 2. <u>Type-I simple</u>: obtained by sampling a parameter only two times at  $\pm 1\sigma_{\text{stat}}$
- Type-II: method and model assumptions
  - Account for possible bias induced by the assumptions made in the fit model and the analysis methods
  - Evaluated only if a **significant** bias is observed while testing an alternative (good) hypothesis
  - $\circ$  A significant bias for a parameter V is defined as a difference  $\Delta$  in the fit results of **more than 20% of its \sigma\_{stat}**
  - $\circ$  In these cases, **half of the bias** is taken as uncertainty, assuming that the *true* bias is uniformly distributed between 0 and  $\Delta$

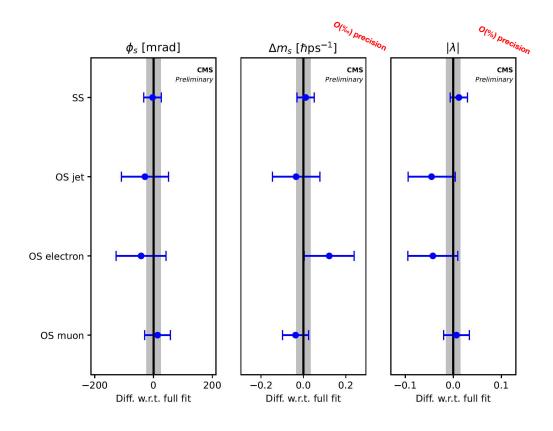
The fit bias does not fall into either of these two categories

## **Systematic uncertainty overview**

	$\phi_s$	$\Delta\Gamma_s$	$\Gamma_s$	$\Delta m_s$	$ \lambda $	$ A_0 ^2$	$ A_{\perp} ^2$	$ A_{\rm S} ^2$	$\delta_{\parallel}$	$\delta_{\perp}$	$\delta_{\mathrm{S}\perp}$
	[mrad]	$[ps^{-1}]$	$[ps^{-1}]$	$[\hbar ps^{-1}]$					[rad]	[rad]	[rad]
Statistical uncertainty	23	0.0043	0.0015	0.035	0.014	0.0016	0.0021	0.0033	0.074	0.089	0.15
Model bias	4	0.0011	0.0002	0.004	0.006	0.0012	0.0022	0.0006	0.015	0.017	0.03
Flavor tagging	4	$< 10^{-4}$	0.0005	0.007	0.002	$< 10^{-4}$	$< 10^{-4}$	0.0006	0.012	0.016	0.03
Angular efficiency	4	0.0002	$< 10^{-4}$	0.015	0.011	0.0042	0.0019	0.0001	0.017	0.044	0.02
Time efficiency	< 1	0.0014	0.0026	$< 10^{-3}$	$< 10^{-3}$	0.0004	0.0005	$< 10^{-4}$	0.001	0.002	$< 10^{-2}$
Time resolution	< 1	$< 10^{-4}$	$< 10^{-4}$	$< 10^{-3}$	$< 10^{-3}$	$< 10^{-4}$	$< 10^{-4}$	$< 10^{-4}$	$< 10^{-3}$	0.001	$< 10^{-3}$
Model assumptions	_	0.0005	0.0006	-	-	-	-	_	_	-	-
B <sup>0</sup> background	< 1	0.0002	0.0003	$< 10^{-3}$	$< 10^{-3}$	$< 10^{-4}$	$< 10^{-4}$	$< 10^{-4}$	$< 10^{-3}$	$< 10^{-3}$	$< 10^{-2}$
$\Lambda_b^0$ background			0.0004	7/	12 <u></u> -1	0.0004	0.0003	9 <u></u> 0	9 <u></u>	9 <u></u>	9 <u></u> 1
S-P wave interference	< 1	$< 10^{-4}$	$< 10^{-4}$	$< 10^{-3}$	$< 10^{-3}$		$< 10^{-4}$	$< 10^{-4}$	$< 10^{-3}$		
$P(\sigma_{ct})$ uncertainty	< 1	0.0002	0.0003	$< 10^{-3}$	$< 10^{-3}$	0.0001	0.0001	$< 10^{-4}$	$< 10^{-3}$	$< 10^{-3}$	$< 10^{-2}$
Total systematic uncertainty	7	0.0019	0.0028	0.017	0.012	0.0044	0.0030	0.0009	0.025	0.050	0.05

- Model bias, flavor tagging, and angular efficiency are found to be the leading systematic sources for  $\phi_{\text{\tiny s}}$
- The measurement is still heavily statistically limited for  $\phi_s$

## Validation: fit with individual tagging techniques



- To check the consistency and stability of the tagging framework, the fit to data is repeated with only one tagging algorithm deployed at a time
  - The grey area represents the result and statistical uncertainty of the full fit
  - Only flavor-sensitive parameters are presented
- Excellent agreement between the various tagging techniques

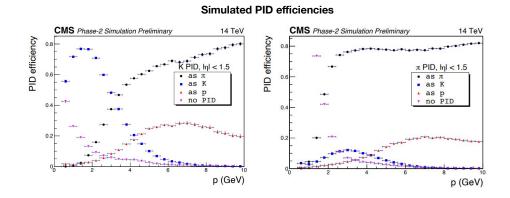
# Comparison with theory and world averages

Parameter	Measured value	World-average value	Theory prediction	-
$\phi_s \text{ [mrad ]}$	$-73 \pm 24$	$-49 \pm 19$	$-37 \pm 1$	[CKMfitter, UTfit]
$\Delta\Gamma_s \; [\mathrm{ps}^{-1}]$	$0.0761 \pm 0.0047$	$0.084 \pm 0.005$	$0.091 \pm 0.013$	[Lenz & Tetlalmatzi-Xolocotzi]
$\Gamma_s  [\mathrm{ps}^{-1}]$	$0.6613 \pm 0.0032$	$0.6573 \pm 0.0023$	-	
$\Delta m_s \ [\hbar \mathrm{ps}^{-1}]$	$17.757 \pm 0.039$	$17.765 \pm 0.006$	$18.77 \pm 0.86$	[Lenz & Tetlalmatzi-Xolocotzi]
$ \lambda $	$1.011 \pm 0.018$	$1.001 \pm 0.018$	1	
$ A_0 ^2$	$0.5300 \pm 0.0047$	$0.520 \pm 0.003$	_	
$ A_{\perp} ^2$	$0.2409 \pm 0.0037$	$0.253 \pm 0.006$	_	
$ A_S ^2$	$0.0067 \pm 0.0034$	$0.030 \pm 0.005$	_	
$\delta_{  }$	$3.145 \pm 0.078$	$3.18 \pm 0.06$	<u></u> 0	
$\delta_{\perp}$	$2.931 \pm 0.102$	$3.08 \pm 0.12$	_	
$\delta_{S\perp}$	$0.48 \pm 0.16$	$0.23 \pm 0.05$	, <del>,</del> ,	_

Reference: CERN-CMS-DP-2022-025

## Flavor tagging in Phase-2 with MTD

- The MTD (Mip Timing Detector) provides time information of charged tracks at its surface
- The reconstruction algorithm utilizes compatible times of tracks from a vertex to offer time-of-flight based particle identification (PID) as a natural byproduct
- Same-side tagging could utilize charge correlation between the s-quark in the B<sub>s</sub> and a nearby soft kaon for flavor tagging
- The PID from MTD, when integrated in the Phase-2 extrapolation of this analysis, shows a significant improvement of the tagging performances



### Relative gain in P<sub>tag</sub> (only SS)

PID scenario	Gains in P <sub>tag</sub>			
MC truth (perfect PID < 3 GeV)	+66%			
PID with $\sigma_{BTL} = 40 \text{ ps}$	+24%			
PID with $\sigma_{BTL} = 70 \text{ ps}$	+14%			