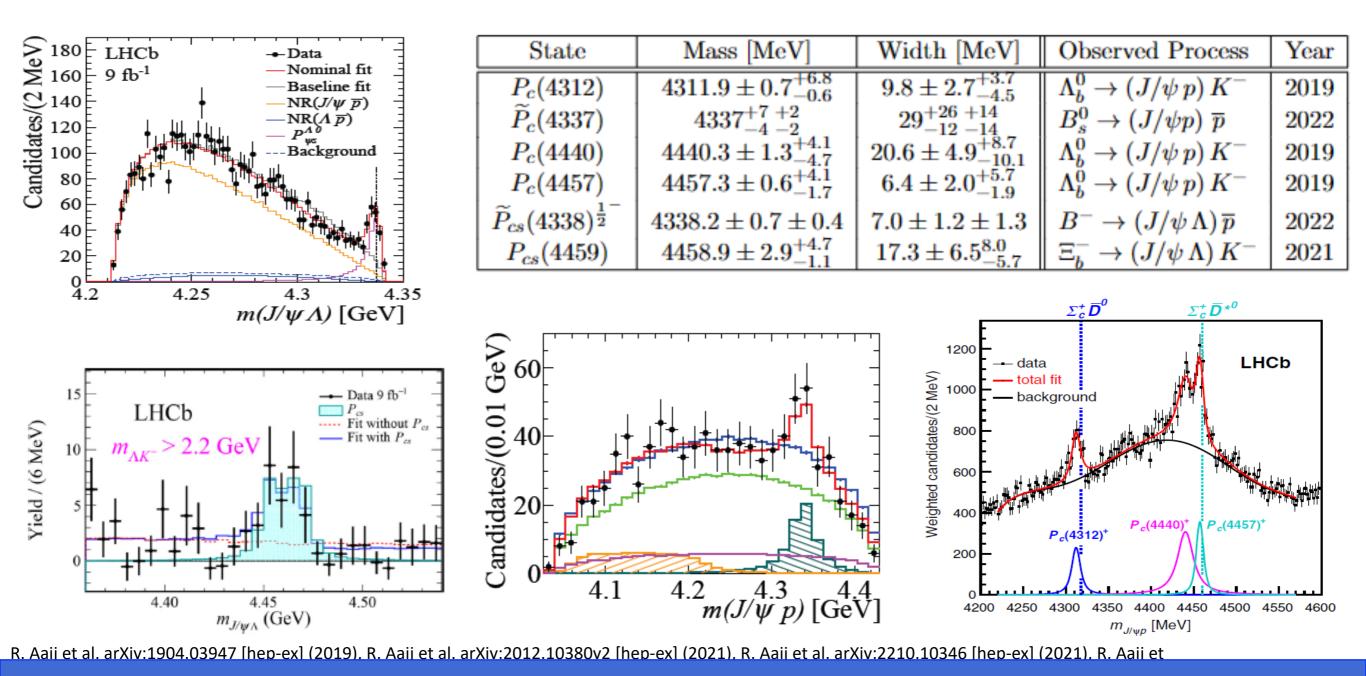
A MODEL OF PENTAQUARKS

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Current Experimental Situation



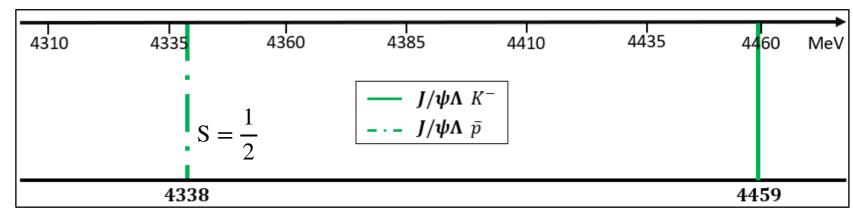
08/04/2024

Slide by D. Germani

K^- and \bar{p} associated production

- P_c^+ [*ccuud*] We divide the spectrum according to strangeness 4335 4360 4385 4435 4460 4310 4410 MeV content. $J/\psi p K^-$ Data suggests two different - **Ι/ψp** p type of production: in 4337 4312/ 4440 4457 association with a K^- or with an \bar{p} . P_{cs}^0 [ccuds]
- Pentaquark seems to appear in triplet

Can we build a model to account for these properties?

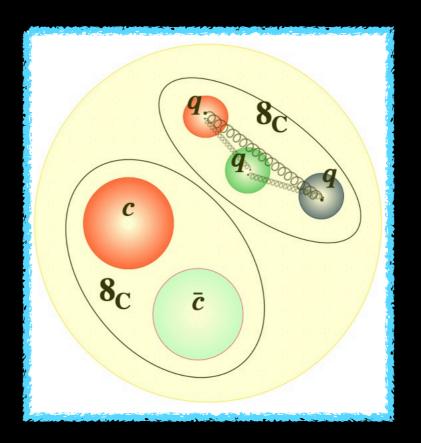


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BARYO-CHARMONIUM

Require Fermi Statistics for the three identical light quarks. The *energy values* of these three quarks depend on their *total spin:* this is *exchange interaction* (not there for spins higher than 1/2).



Under exchange the spin wf ψ_{spin} of two particles gets $(-1)^{2s+S}$. Statistics requires $\Psi(=\psi_{spin}\psi_{coord})$ to get $(-1)^{2s}$ so that $\psi_{coord.} \sim (-1)^S \Rightarrow L$ is even(odd) if S is even(odd). If we include color-flavor: $\Psi = \psi_{col-flav} \psi_{spin} \psi_{coord}$

Ψ	$\psi_{\rm col-flav}$	$\psi_{\rm coord}$	$\psi_{ m spin}$
А	S	S	А
A	S	A	S
A	А	А	А
А	А	S	S

On the use of Fermi statistics for pentaquarks see L. Maiani, ADP, V. Riquer Eur. Phys. J. C83 (2023) 5, 378

EXCHANGE INTERACTION

$$\psi_{\text{coord}}(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{\sqrt{2}} \left(\phi_1(\mathbf{r}_1) \phi_2(\mathbf{r}_2) \pm \phi_1(\mathbf{r}_2) \phi_2(\mathbf{r}_1) \right)$$

$$\langle U \rangle_{\psi_{\text{coord}}} = \int \psi_{\text{coord}}^*(\mathbf{r}_1, \mathbf{r}_2) U(\mathbf{r}_1 - \mathbf{r}_2) \psi_{\text{coord}}(\mathbf{r}_1, \mathbf{r}_2) d^3r_1 d^3r_2$$

 $\langle U \rangle_{\psi_{\text{coord}}} = C \pm J$

Color-flavor symmetric case:

Symmetric in coordinates i.e. $+ \Rightarrow S = 0$ in a pair of fermions (S = 0,1) $\Rightarrow S(S+1) - 1 = -1 \Rightarrow$

$$+J = \langle V \rangle_{\psi_{\text{spin}}} = \left\langle -\sum_{\text{pairs}} J_{ab} \underbrace{\left(\frac{1}{2} + 2\mathbf{S}_a \cdot \mathbf{S}_b\right)}_{S(S+1)-1}\right\rangle$$

EXCHANGE INTERACTION

Coordinates anti-symm. i.e. $\rightarrow S = 1$ in a pair of fermions (S = 0,1) $\Rightarrow S(S+1) - 1 = +1 \Rightarrow$

$$-J = \langle V \rangle_{\psi_{\text{spin}}} = \left\langle -\sum_{\text{pairs}} J_{ab} \left(\frac{1}{2} + 2\mathbf{S}_a \cdot \mathbf{S}_b \right) \right\rangle$$
$$\underbrace{\int_{S(S+1)-1}}_{S(S+1)-1} V_{ab} = \left[(\phi_a(\mathbf{r}_a)\phi_b(\mathbf{r}_b))^* U(\mathbf{r}_a - \mathbf{r}_b) (\phi_a(\mathbf{r}_b)\phi_b(\mathbf{r}_a)) d^3r_a d^3r_$$

Determine the exchange splitting of the energy levels of a system of three quarks (treating the interaction as a perturbation)

rb,

$$\begin{split} \Delta E_{1/2} &= \pm \sqrt{J_{12}^2 + J_{13}^2 + J_{23}^2 - J_{12}J_{13} - J_{12}J_{23} - J_{13}J_{23}} \\ \Delta E_{3/2} &= -J_{12} - J_{13} - J_{23} \end{split}$$

In the case the <u>color-flavor is anti-symmetric</u>, change the sign in the definition of the exchange potential V.

Let a, b, c be flavor indices and α, β, γ color indices. Requiring both color and flavor to be in the adjoint representation we can form the tensor

$$A^{abc}_{\alpha\beta\gamma} = \psi^{[a}_{\alpha}\psi^{b]}_{\beta}\psi^{c}_{\gamma} - \psi^{[a}_{\beta}\psi^{b]}_{\gamma}\psi^{c}_{\alpha} = \psi^{a}_{(\alpha}\psi^{b}_{\beta)}\psi^{c}_{\gamma} - \psi^{a}_{(\beta}\psi^{b}_{\gamma)}\psi^{c}_{\alpha}$$

anti-symmetric under the exchange of any two quarks provided that

$$\psi^a_\alpha \psi^b_\beta = -\,\psi^b_\beta \psi^a_\alpha$$

COLOR-FLAVOR

Similarly we can form the symmetric color-flavor tensor

$$S^{abc}_{\alpha\beta\gamma} = \eta^{[a}_{\alpha}\eta^{b]}_{\beta}\eta^{c}_{\gamma} - \eta^{[a}_{\beta}\eta^{b]}_{\gamma}\eta^{c}_{\alpha} = \eta^{a}_{[\alpha}\eta^{b}_{\beta]}\eta^{c}_{\gamma} - \eta^{a}_{[\beta}\eta^{b}_{\gamma]}\eta^{c}_{\alpha}$$

This is symmetric under the exchange of any two quarks, provided that

 $\eta^a_\alpha \eta^b_\beta = + \eta^b_\beta \eta^a_\alpha$

This allows to make the baryon octet: assume that tensor S represents the flavor octet with spin indices α, β, γ . Assign additional color indices and anti-symmetrize in color

$$\eta^{ai}_{\alpha}\eta^{bj}_{\beta} = \eta^{bi}_{\beta}\eta^{aj}_{\alpha} = -\eta^{bj}_{\beta}\eta^{ai}_{\alpha}$$

so that we can use $\eta \rightarrow \psi$. Spin-flavor (S) color (A)

$$\mathcal{S}^{abc}_{\alpha\beta\gamma} = \epsilon_{ijk}(\psi^{ai}_{[\alpha}\psi^{bj}_{\beta]}\psi^{ck}_{\gamma} + \psi^{ai}_{[\gamma}\psi^{bj}_{\beta]}\psi^{ck}_{\alpha}) = B^{acb}_{\alpha\gamma\beta}$$



We assign

 $S^{abc}_{\alpha\beta\gamma} \longrightarrow P$

for $(J/\psi p)$ or $(J/\psi \Lambda)$ pentaquarks produced in association with a kaon K^- and

 $A^{abc}_{\alpha\beta\gamma}\longrightarrow\widetilde{P}$

 $(J/\psi p)$ or $(J/\psi \Lambda)$ pentaquarks produced in association with an anti-proton \bar{p} ; in this case change the sign in the definition of the exchange potential V.

P pentaquarks of the $(J/\psi p)$ kind

Let us consider the case qqq = uud in the color-flavor configuration S.

$S_{\alpha\beta\gamma}^{121} = u_{[\alpha}d_{\beta]}u_{\gamma} + u_{[\gamma}d_{\beta]}u_{\alpha}$

since the *uu* cannot be anti-symmetric in flavor space. The pentaquark contains the light quarks u_{α} , d_{β} and u_{γ} with the *u* quarks, u_{α} and u_{γ} , in a color symmetric (repulsive) representation and *ud* in color attractive, anti-symmetric pairings. Therefore we require

$$J_{S}^{uu} = -\frac{1}{2}J_{A}^{ud} \qquad (C_{6} - 2C_{3}) = -\frac{1}{2}(C_{3} - 2C_{3})$$
$$= J_{S}^{uu} > 0 \qquad J_{13} = J_{23} = J_{A}^{ud} < 0$$

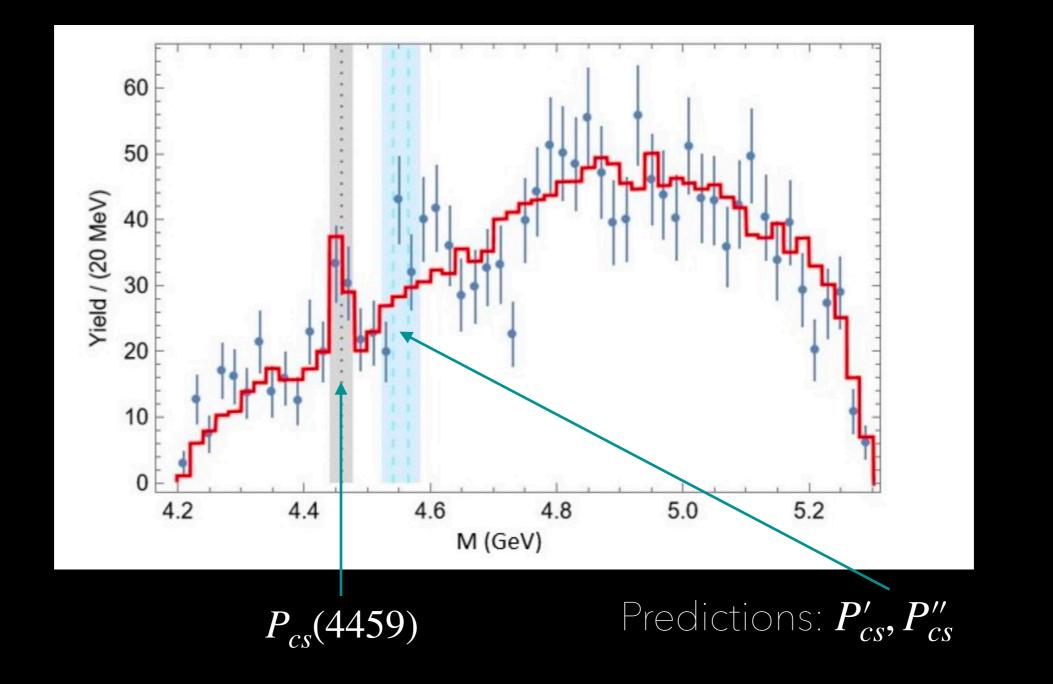
Assume that the ordering in mass corresponds to the lower one being spin 1/2 and the higher two being 3/2 and 1/2 respectively. Then we have to solve the simultaneous equations

$$\begin{split} M_{P_c}(4457) - M_{P_c}(4312) &= 2 \left| J_S^{uu} - J_A^{ud} \right| \\ M_{P_c}(4440) - \frac{1}{2} (M_{P_c}(4312) + M_{P_c}(4457)) &= -J_S^{uu} - 2J_A^{ud} \\ & \text{A solution can be found} \\ J_S^{uu} &= 29.9^{+2.5}_{-2.8} \text{ MeV} \qquad J_A^{ud} &= -42.8^{+2.4}_{-1.6} \text{ MeV} \\ & \frac{J_S^{uu}}{J_u^{ud}} &= -0.7 \pm 0.1 \end{split}$$

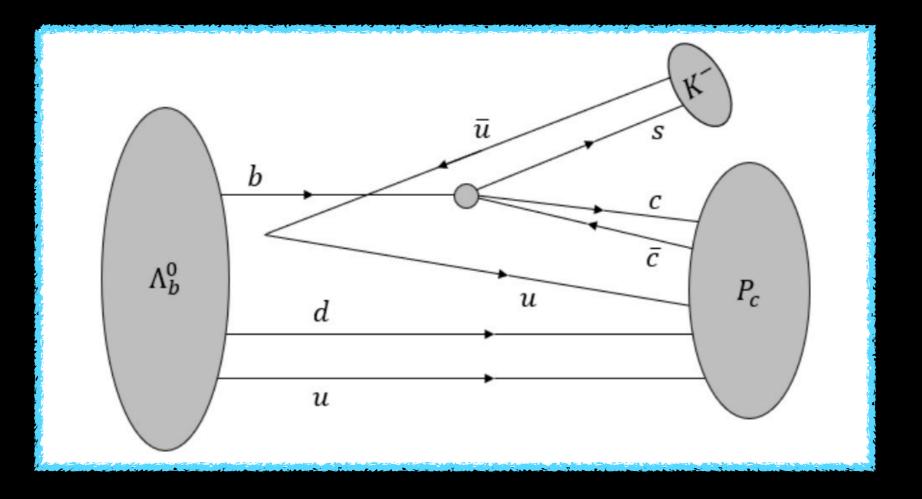
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P pentaquarks of the $(J/\psi\Lambda)$ kind

Using the same methods we can study the strange pentaquarks. One has been observed and two more are predicted.



\widetilde{P} pentaquarks of the $(J/\psi\Lambda,p)$ kind



Assuming that the color-flavor of the spectator ud pair is kept in the final state, the process $\Lambda_b^0 \to \tilde{P}_c K^-$ is expected to be suppressed, as suggested by observation: the $[ud]_0$ in Λ is in a CF-symmetric configuration as opposite to \tilde{P} as described by

$$A^{udu}_{\alpha\beta\gamma} = u_{(\alpha}d_{\beta)}u_{\gamma} - u_{(\gamma}d_{\beta)}u_{\alpha}$$

P and \widetilde{P} pentaquarks & conclusions

The spin ordering (1/2, 3/2, 1/2) is a main prediction.

	Mass [MeV]		Mass [MeV]
$P_{c}(4312)$	$(4311.9^{+7}_{-0.9})$	$P_{cs}(4459)$	$(4458.8^{+6}_{-3.1})$
$P_{c}(4440)$	(4440.0^{+4}_{-5})	\mathbf{P}_{cs}'	4541 ± 6
$P_c(4457)$	$(4457.3^{+7}_{-1.8})$	\mathbf{P}_{cs}''	4565 ± 6
$\widetilde{\mathbf{P}}_{\mathbf{c}}''$	4187 ± 7	$\widetilde{P}_{cs}(4338)$	(4338.2 ± 0.8)
$\widetilde{\mathbf{P}}_{\mathbf{c}}'$	4276 ± 12	$\widetilde{\mathbf{P}}'_{\mathbf{cs}}$	4387 ± 4
$\widetilde{P}_c(4337)$	$4332 \pm 7~(4337^{+7}_{-4}~^{+2}_{-2})$	$\widetilde{\mathbf{P}}_{\mathbf{cs}}''$	4435 ± 4