



### Time Independent CPV at LHCb

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### Introduction



- Charge-Parity (CP) symmetry = natural laws the same for (anti-)matter
  - C changes particles for anti-particles, e.g.,  $Q \rightarrow -Q$
  - P reverses spatial handedness,  $\vec{x} \rightarrow -\vec{x}$
  - CP symmetry apparently exact for electromagnetic and strong interactions
  - CP violation (CPV) seen in weak interactions
- CPV in the Standard Model (SM) comes from a single parameter in the Cabibbo-Kobayashi-Maskawa (CKM) mixing matrix, which determines the size of the couplings between quark flavors
- Not nearly large enough to explain observed (anti-)matter asymmetry! Prompts searches for sources Beyond the SM (BSM)

R.L. Workman et al. (Particle Data Group),

Sources of CPV



#### 1. CPV in decay ("direct")

R.L. Workman et al. (Particle Data Group), Prog. Theor. Exp. Phys. 2022, 083C01 (2022) and 2023 update

• Defined by different rates of charge-conjugated decay for hadron *M*:

$$\mathcal{A}_{f^{\pm}} \equiv \frac{\Gamma(M^- \to f^-) - \Gamma(M^+ \to f^+)}{\Gamma(M^- \to f^-) + \Gamma(M^+ \to f^+)}$$

- The only possible source of CPV for baryons and charged mesons
- 2. CPV in mixing
  - Comes from differences in mass and flavor eigenstates for neutral meson *M*:  $|M_{1,2}\rangle = p|M^0\rangle + q|\overline{M}^0\rangle, \quad |p|^2 + |q|^2 = 1$  $|q/p| \neq 1 \Rightarrow CPV$
  - The only source in charged-current semileptonic neutral-meson decays
- 3. CPV in interference between  $M^0 \to f$  and  $M^0 \to \overline{M}^0 \to f$

- Search for direct CP violation,  $\mathcal{A}^{CP}$
- Arises from interference between decay amplitudes
- Predicted to be small in SM, up to 1% for  $b \rightarrow c\bar{c}s$  and up to 5% for  $b \rightarrow c\bar{c}d$
- Also, measure two BF ratios,  $R(D^-D^0/D_s^-D^0) \& R(D^{*-}D^0/D^-D^0)$

$$R(D^{-}D^{0}/D_{s}^{-}D^{0}) \equiv \frac{\mathcal{B}(B^{-} \to D^{-}D^{0})}{\mathcal{B}(B^{-} \to D_{s}^{-}D^{0})} \frac{\mathcal{B}(D^{-} \to K^{+}\pi^{-}\pi^{-})}{\mathcal{B}(D_{s}^{-} \to K^{+}K^{-}\pi^{-})}$$
$$R(D^{*-}D^{0}/D^{-}D^{0}) \equiv \frac{\mathcal{B}(B^{-} \to D^{*-}D^{0})}{\mathcal{B}(B^{-} \to D^{-}D^{0})} \frac{\mathcal{B}(D^{*-} \to \overline{D}^{0}\pi^{-})\mathcal{B}(\overline{D}^{0} \to K^{+}\pi^{-})}{\mathcal{B}(D^{-} \to K^{+}\pi^{-}\pi^{-})}$$

*LHCb* 

202 [arXiv:2306.09945

 $\mathcal{A}^{CP} \equiv \frac{\Gamma\left(B^- \to D_{(s)}^{(*)-} D^{(*)0}\right) - \Gamma\left(B^+ \to D_{(s)}^{(*)+} \overline{D}^{(*)0}\right)}{\Gamma\left(B^- \to D_{(s)}^{(*)-} D^{(*)0}\right) + \Gamma\left(B^+ \to D_{(s)}^{(*)+} \overline{D}^{(*)0}\right)}$ 

Decay	World Average $\mathcal{A}^{CP}$ [%]			
$B^- \rightarrow D_s^-$	$D^0$	$-0.4 \pm 0.7$		
$B^- \rightarrow D_s^*$	$^{-}D^{0}$	-		
$B^- \rightarrow D_s^-$	$D^{*0}$	-		
$B^- \rightarrow D^-$	$D^0$	$1.6 \pm 2.5$		
$B^- \rightarrow D^-$	$D^{*0}$	$13 \pm 18$		
$B^- \rightarrow D^*$	$^{-}D^{0}$	$-6 \pm 13$		
$B^- \rightarrow D^*$	$^{-}D^{*0}$	$-15 \pm 11$		
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- 9 fb<sup>-1</sup> of data (Run 1 + 2)
- Reconstruct  $D_s^- D^0$ ,  $D^- D^0$ , or  $D^{*-} D^0$ final states
- Partially reconstruct intermediate  $D_s^{*-/0} \rightarrow D(\gamma/\pi^0)$  decays (contributes broad structure to invariant mass)



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Measurement of CP asymmetries and branching fraction ratios of  $B^-$  decays to two charm mesons

• 
$$\mathcal{A}^{CP} = \mathcal{A}_{raw} - \mathcal{A}_{P} - \mathcal{A}_{D}$$

- Calculate raw asymmetry (A<sub>raw</sub>) from yields and correct by the production and detection asymmetries (A<sub>P</sub> and A<sub>D</sub>) to get A<sup>CP</sup>
- *A*<sub>P</sub> and *A*<sub>D</sub> evaluated with kinematically-weighted calibration data
- Branching fraction ratios are measured for fully-reconstructed decays where we achieve high precision
  - Calculated using efficiency-corrected yields
  - Efficiency from data-corrected simulation

$$\mathcal{A}_{\rm raw} \equiv \frac{N\left(B^- \to D_{(s)}^{(*)-} D^{(*)0}\right) - N\left(B^+ \to D_{(s)}^{(*)+} \overline{D}^{(*)0}\right)}{N\left(B^- \to D_{(s)}^{(*)-} D^{(*)0}\right) + N\left(B^+ \to D_{(s)}^{(*)+} \overline{D}^{(*)0}\right)}$$

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$$\mathcal{A}_{\rm P} \equiv \frac{\sigma(B^-) - \sigma(B^+)}{\sigma(B^-) + \sigma(B^+)}$$

$$\mathcal{A}_{\mathrm{D}} \equiv \frac{\varepsilon \left( B^{-} \to D_{(s)}^{(*)-} D^{(*)0} \right) - \varepsilon \left( B^{+} \to D_{(s)}^{(*)+} \overline{D}^{(*)0} \right)}{\varepsilon \left( B^{-} \to D_{(s)}^{(*)-} D^{(*)0} \right) + \varepsilon \left( B^{+} \to D_{(s)}^{(*)+} \overline{D}^{(*)0} \right)}$$

$$R(D^{-}D^{0}/D_{s}^{-}D^{0}) = \frac{N(B^{-} \to D^{-}D^{0})}{N(B^{-} \to D_{s}^{-}D^{0})} \frac{\varepsilon(B^{-} \to D_{s}^{-}D^{0})}{\varepsilon(B^{-} \to D^{-}D^{0})}$$
$$R(D^{*-}D^{0}/D^{-}D^{0}) = \frac{N(B^{-} \to D^{*-}D^{0})}{N(B^{-} \to D^{-}D^{0})} \frac{\varepsilon(B^{-} \to D^{-}D^{0})}{\varepsilon(B^{-} \to D^{*-}D^{0})}$$

- $\mathcal{A}^{CP} = \mathcal{A}_{raw} \mathcal{A}_{P} \mathcal{A}_{D}$
- Fit  $m\left(D_{(s)}^{(*)-}D^{(*)0}\right)$  to extract  $\mathcal{A}_{\text{raw}}$ 
  - Asymmetry extracted from simultaneous fit to both charges
  - Background asymmetries constrained to improve resolution
- Extract  $\mathcal{A}_{\mathrm{P}}$  by kinematically weighting previous LHCb result from  $B^+ \to J/\psi K^+$
- Extract  $\mathcal{A}_D$  (dominated by  $K^-$  nuclear interaction) using independent, kinematically weighted calibration samples of  $D^+$ ,  $D^{*+}$ , and B decays





- $R(D^-D^0/D_s^-D^0) \& R(D^{*-}D^0/D^-D^0)$ 
  - Agree with world averages
  - Higher precision



- $R(D^-D^0/D_s^-D^0) \& R(D^{*-}D^0/D^-D^0)$ 
  - Agree with world averages
  - Higher precision
- $\mathcal{A}^{CP}$ 
  - No evidence of CP violation found
  - More precise than world averages
  - $\mathcal{A}^{CP}(B^- \to D^-_{(s)}D^0)$  agree with and supersede previous LHCb measurement
- Substantially improve knowledge of B<sup>-</sup> meson decays, helping to constrain BSM physics



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Search for CP violation in the phase space of  $D^0 \rightarrow K_S^0 K^{\pm} \pi^{\mp}$  decays with the energy test

- CPV in charm decays is expected to be small in the SM
- Observed for the first time in 2019 by LHCb in singly Cabibbo-suppressed (SCS) decay modes [1], unclear whether consistent with the SM
- $D^0 \rightarrow K_S^0 K^{\pm} \pi^{\mp}$  decays are dominated by SCS amplitudes, previously studied in amplitude analysis by LHCb [2], and contain other CP-sensitive amplitudes
- 1. Phys. Rev. Lett. 122 (2019) 211803, arXiv:1903.08726
- 2. Phys. Rev. D 93 (2016) 052018, arXiv:1509.06628



JHEP 03 (2024) 107 [arXiv:2310.19397]



Search for CP violation in the phase space of  $D^0 \rightarrow K_S^0 K^{\pm} \pi^{\mp}$  decays with the energy test

- The energy test quantifies whether two multi-dimensional datasets are consistent with the same underlying distribution
- Test statistic *T* near zero for no CPV, large for significant CPV
- Significance determined by comparing found *T*-value to null distribution
  - Repeatedly run the test with the  $D^0$ ,  $\overline{D}^0$  flavors randomly assigned = null dist.
  - The *p*-value = fraction of permutation samples with *T*-value > the found value



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Sum over  $n D^0$  pairs,  $\overline{n} \overline{D}{}^0$  pairs, and  $n, \overline{n} D^0 - \overline{D}{}^0$  pairs

$$\psi_{ij} = e^{-d_{ij}^2/2\delta^2}$$

$$d_{ij}^{2} = \left(s_{12,i} - s_{12,j}\right)^{2} + \left(s_{13,i} - s_{13,j}\right)^{2} + \left(s_{23,i} - s_{23,j}\right)^{2}$$

$$s_{12} = m^2 (K_S^0 K^{\pm}), s_{13} = m^2 (K_S^0 \pi^{\mp}), s_{23} = m^2 (K^{\pm} \pi^{\mp})$$

 $d_{ij}^2$  is the square of the Euclidean distance between candidates  $\delta$  sets the distance scale (optimized for max. sensitivity)





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T-value

### Conclusions



- Measured CPV in  $B^- \rightarrow D_{(s)}^{(*)-} D^{(*)0}$  and  $D^0 \rightarrow K_S^0 K^{\pm} \pi^{\mp}$  decays
  - JHEP 09 (2023) 202 [arXiv:2306.09945]
  - JHEP 03 (2024) 107 [arXiv:2310.19397]
- Found no evidence of CPV in these decays
- These measurements are statistically limited
  - We expect ≈5x greater integrated luminosity in Run 3+4
  - And we expect much larger hadron-trigger efficiency in Run 3+4
  - $\therefore$  we expect  $\geq 10x$  greater statistics in near future!



JHEP 09 (2023) 202 [arXiv:2306.09945]

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#### FIN



### BACKUP

- $\mathcal{A}_{\mathrm{D}} = \mathcal{A}_{K\pi} + \mathcal{A}_{\pi} + \mathcal{A}_{\mathrm{PID}} + \mathcal{A}_{\mathrm{TIS}} + \mathcal{A}_{\mathrm{TOS}}$
- Weight calibration samples to match signal kinematics
- $\mathcal{A}_{K\pi}$ 
  - Difference in raw asymmetry between  $D^+ \rightarrow K^- \pi^+ \pi^+$  and  $D^+ \rightarrow \overline{K}{}^0 \pi^+$
  - Corrected for  $\overline{K}^0$  asymmetry
- $\mathcal{A}_{\pi}$  from  $D^{*+} \to (D^0 \to K^- \pi^+ \pi^- \pi^+) \pi^+$
- $\mathcal{A}_{\text{PID}}$ 
  - Induced by tight particle identification (PID) requirements
  - From  $D^{*+} \rightarrow (D^0 \rightarrow K^- \pi^+)\pi^+$  without PID requirements
- $\mathcal{A}_{\text{TIS}}, \mathcal{A}_{\text{TOS}}$ 
  - Induced by hardware trigger requirements
  - From  $D^{*+} \to (D^0 \to K^- \pi^+) \pi^+$  and  $B \to \overline{D}{}^0 \mu^+ \nu_{\mu} X$

$$\mathcal{A}_{\mathrm{D}} \equiv \frac{\varepsilon \left( B^{-} \to D_{(s)}^{(*)-} D^{(*)0} \right) - \varepsilon \left( B^{+} \to D_{(s)}^{(*)+} \overline{D}^{(*)0} \right)}{\varepsilon \left( B^{-} \to D_{(s)}^{(*)-} D^{(*)0} \right) + \varepsilon \left( B^{+} \to D_{(s)}^{(*)+} \overline{D}^{(*)0} \right)}$$

Decay	$\mathcal{A}_{\mathrm{raw}}$	$\mathcal{A}_{ ext{P}}$	$\mathcal{A}_{\mathrm{D}}$
$D_s^- D^0$	$-1.3\pm0.2\pm0.1$	$-1.1\pm0.3\pm0.3$	$-0.7\pm0.2$
$D_{s}^{*-}D^{0}$	$-2.4 \pm 1.1 \pm 0.9$	$-1.1\pm0.4\pm0.3$	$-0.8\pm0.2$
$D_{s}^{-}D^{*0}$	$-0.8\pm0.8\pm0.4$	$-1.1\pm0.4\pm0.3$	$-0.8\pm0.2$
$D^{-}D^{0}$	$1.5\pm1.0\pm0.2$	$-1.1\pm0.4\pm0.3$	$0.1 \pm 0.2$
$D^{-}D^{*0}$	$-1.3 \pm 2.0 \pm 1.3$	$-1.1 \pm 0.4 \pm 0.3$	$0.1 \pm 0.2$
$D^{*-}D^{0}$	$2.4\pm1.6\pm0.2$	$-1.2\pm0.4\pm0.3$	$0.2 \pm 0.3$
$D^{*-}D^{*0}$	$1.3\pm2.1\pm1.6$	$-1.1\pm0.5\pm0.3$	$0.1\pm0.2$

Values of  $\mathcal{A}_{raw}$ ,  $\mathcal{A}_{P}$ , and  $\mathcal{A}_{D}$  in percent, averaged over all  $D^{0}$  decay modes and data-taking periods. The uncertainties on  $\mathcal{A}_{raw}$  are statistical and systematic, respectively. The first uncertainty on  $\mathcal{A}_{P}$  contains all sources of uncertainty except that on  $\mathcal{A}^{CP}(B^{+} \rightarrow J/\psi K^{+})$ , which is the second uncertainty.

- $\mathcal{A}_{\mathrm{D}} = \mathcal{A}_{K\pi} + \mathcal{A}_{\pi} + \mathcal{A}_{\mathrm{PID}} + \mathcal{A}_{\mathrm{TIS}} + \mathcal{A}_{\mathrm{TOS}}$
- Weight calibration samples to match signal kinematics
- $\mathcal{A}_{K\pi}$ 
  - Difference in raw asymmetry between  $D^+ \rightarrow K^- \pi^+ \pi^+$  and  $D^+ \rightarrow \overline{K}{}^0 \pi^+$
  - Corrected for  $\overline{K}^0$  asymmetry
- $\mathcal{A}_{\pi}$  from  $D^{*+} \to (D^0 \to K^- \pi^+ \pi^- \pi^+) \pi^+$
- $\mathcal{A}_{\text{PID}}$ 
  - Induced by tight particle identification (PID) requirements
  - From  $D^{*+} \rightarrow (D^0 \rightarrow K^- \pi^+)\pi^+$  without PID requirements
- $\mathcal{A}_{TIS}, \mathcal{A}_{TOS}$ 
  - Induced by hardware trigger requirements
  - From  $D^{*+} \to (D^0 \to K^- \pi^+) \pi^+$  and  $B \to \overline{D}{}^0 \mu^+ \nu_{\mu} X$

$$\mathcal{A}_{\mathrm{D}} \equiv \frac{\varepsilon \left( B^{-} \to D_{(s)}^{(*)-} D^{(*)0} \right) - \varepsilon \left( B^{+} \to D_{(s)}^{(*)+} \overline{D}^{(*)0} \right)}{\varepsilon \left( B^{-} \to D_{(s)}^{(*)-} D^{(*)0} \right) + \varepsilon \left( B^{+} \to D_{(s)}^{(*)+} \overline{D}^{(*)0} \right)}$$

Final state	$D_s^-$	$D^0$	$D^{-}$	$D^0$	$D^{*-}$	$^{-}D^{0}$
	Run 1	$\operatorname{Run}2$	$\operatorname{Run} 1$	$\operatorname{Run}2$	$\operatorname{Run} 1$	$\operatorname{Run}2$
$\mathcal{A}_{ ext{P}}$	0.42	0.43	0.41	0.43	0.48	0.48
$\mathcal{A}^{CP}(B^+ \to J/\psi K^+)$	0.30	0.30	0.30	0.30	0.30	0.30
$\mathcal{A}_{K\pi}$	0.28	0.11	0.04	0.04	0.10	0.00
$\mathcal{A}_{\pi}$	0.09	0.09	0.06	0.06	0.18	0.17
$\mathcal{A}_{ ext{PID}}$	0.29	0.03	0.25	0.11	0.55	0.10
$\mathcal{A}_{ ext{TIS}}$	0.08	0.10	0.08	0.10	0.09	0.11
$\mathcal{A}_{\mathrm{TOS}}$	0.01	0.03	0.01	0.02	0.01	0.01
Weighting	0.01	0.00	0.04	0.00	0.01	0.00
Part. rec. weighting	0.03	0.02	0.02	0.01	0.03	0.01
Total	0.67	0.55	0.58	0.55	0.82	0.61

Systematic uncertainties on the corrections for  $\mathcal{A}^{CP}$  in percent, averaged over all D0 decay modes.



$D^0$ decay mode	Run 1	Run 2	Run $1+2$
$K^{-}\pi^{+}$	$6.88 \pm 0.24 \pm 0.12$	$7.35 \pm 0.12 \pm 0.11$	$7.22 \pm 0.11 \pm 0.10$
$K^-\pi^+\pi^-\pi^+$	$6.93 \pm 0.38 \pm 0.23$	$7.40 \pm 0.18 \pm 0.15$	$7.30 \pm 0.16 \pm 0.14$
Combined	$6.89 \pm 0.20 \pm 0.12$	$7.36 \pm 0.10 \pm 0.10$	$7.25 \pm 0.09 \pm 0.09$

Values of  $R(D^-D^0/D_s^-D^0)/10^{-2}$  for each  $D^0$  decay mode, for Run 1 and Run 2 and the combined measurement. The first uncertainty is statistical and the second is systematic

$D^0$ decay mode	Run 1	Run 2	Run $1+2$
$K^-\pi^+$	$0.328 \pm 0.023 \pm 0.011$	$0.256 \pm 0.009 \pm 0.005$	$0.271 \pm 0.008 \pm 0.005$
$K^-\pi^+\pi^-\pi^+$	$0.316 \pm 0.033 \pm 0.015$	$0.272 \pm 0.012 \pm 0.008$	$0.278 \pm 0.012 \pm 0.007$
Combined	$0.324 \pm 0.019 \pm 0.010$	$0.262 \pm 0.007 \pm 0.005$	$0.271 \pm 0.007 \pm 0.005$

Values of  $R(D^{*-}D^{0}/D^{-}D^{0})$  for each  $D^{0}$  decay mode, for Run 1 and Run 2 and the combined measurement. The first uncertainty is statistical and the second is systematic

 $\begin{aligned} \mathcal{A}^{CP}(B^- \to D_s^- D^0) &= (+0.5 \pm 0.2 \pm 0.5 \pm 0.3)\% \\ \mathcal{A}^{CP}(B^- \to D_s^{*-} D^0) &= (-0.5 \pm 1.1 \pm 1.0 \pm 0.3)\% \\ \mathcal{A}^{CP}(B^- \to D_s^- D^{*0}) &= (+1.1 \pm 0.8 \pm 0.6 \pm 0.3)\% \\ \mathcal{A}^{CP}(B^- \to D^- D^0) &= (+2.5 \pm 1.0 \pm 0.4 \pm 0.3)\% \\ \mathcal{A}^{CP}(B^- \to D^- D^{*0}) &= (-0.2 \pm 2.0 \pm 1.4 \pm 0.3)\% \\ \mathcal{A}^{CP}(B^- \to D^{*-} D^0) &= (+3.3 \pm 1.6 \pm 0.6 \pm 0.3)\% \\ \mathcal{A}^{CP}(B^- \to D^{*-} D^{*0}) &= (+2.3 \pm 2.1 \pm 1.7 \pm 0.3)\% \end{aligned}$ 

Values of  $\mathcal{A}^{CP}$  where the first uncertainty is statistical, the second is systematic, and third is from  $\mathcal{A}^{CP}(B^+ \to J/\psi K^+)$ .