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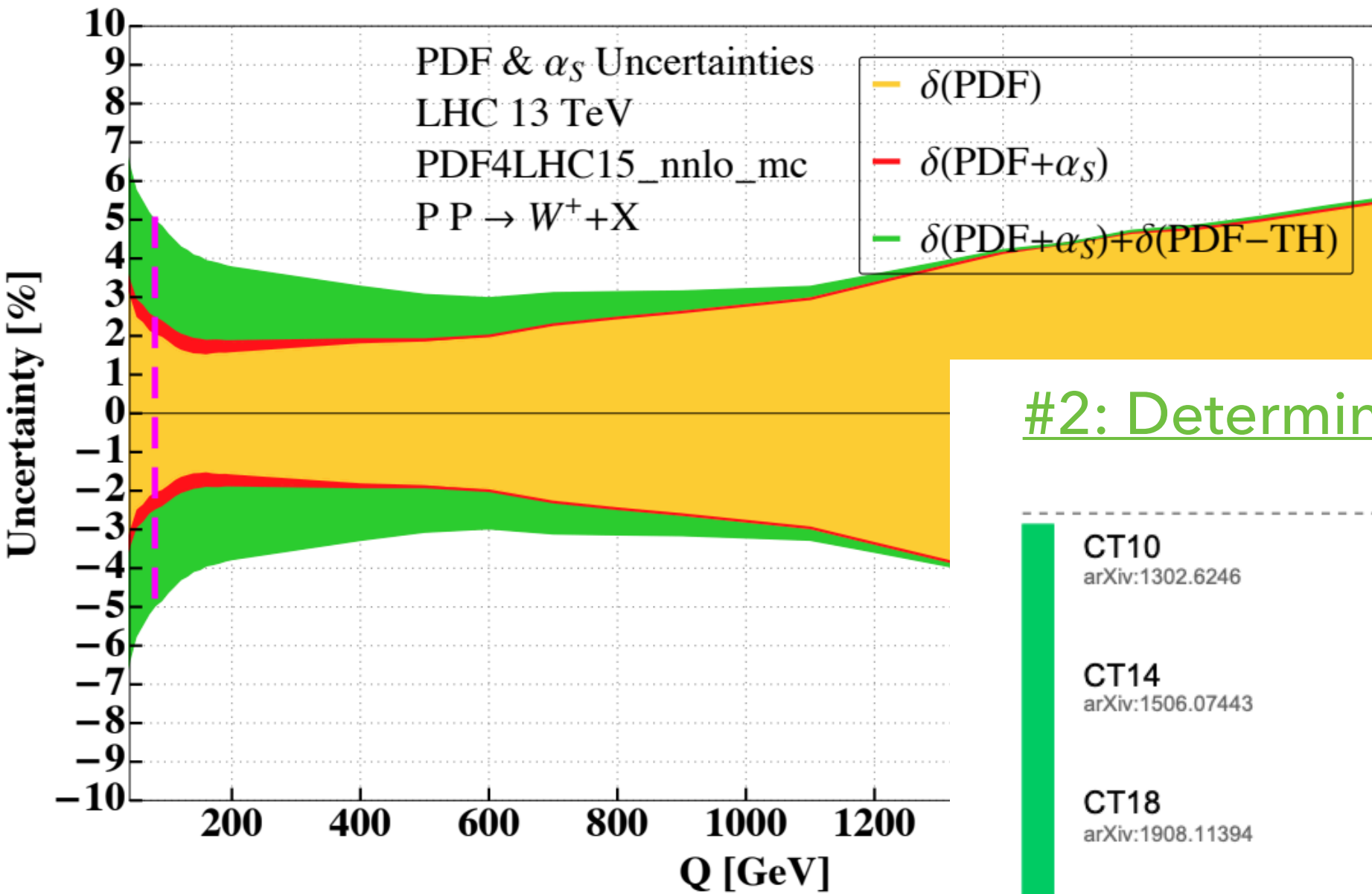


MARIA UBIALI  
UNIVERSITY OF CAMBRIDGE

# PDFS FOR PRECISION MEASUREMENTS AT THE LHC

# PDFS ARE KEY TO PRECISION PHYSICS

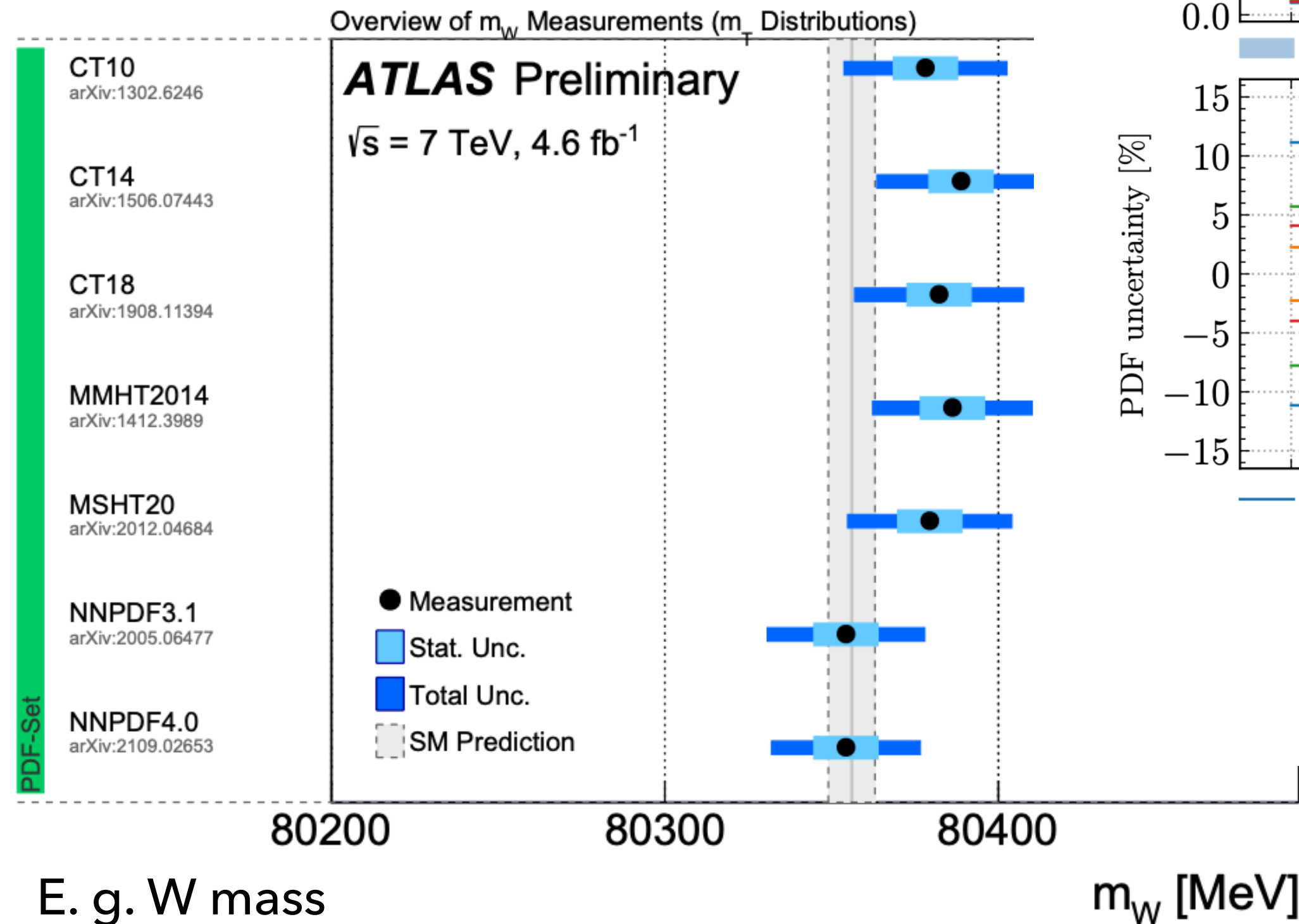
## #1: Theory uncertainty of SM predictions



E. g. Drell-Yan @ N3LO

C. Duhr et al, *JHEP* 11 (2020) 143

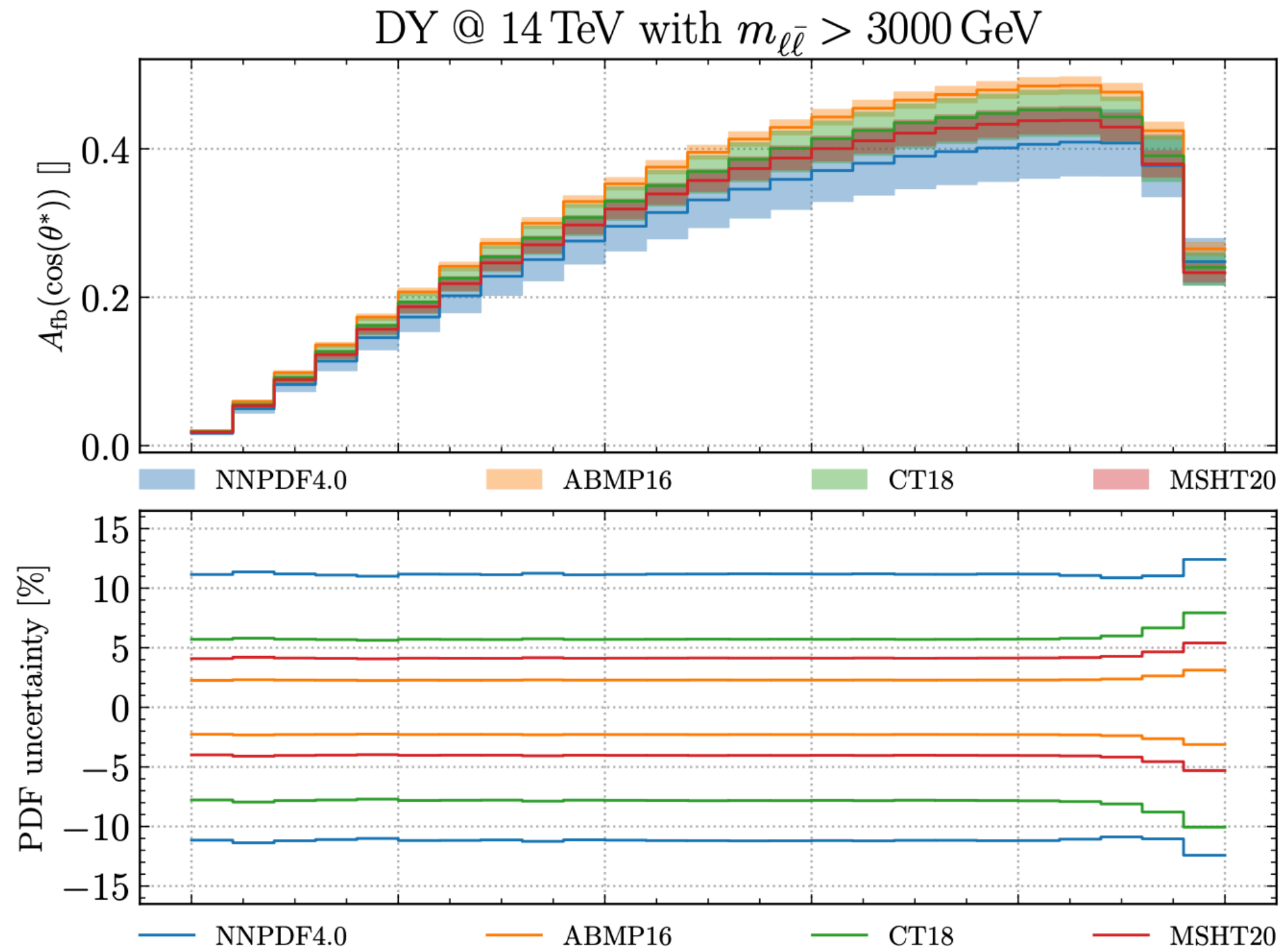
## #2: Determination of SM parameters



E. g. W mass

ATLAS-CONF-2023-004

## #3: NP searches



E. g. Heavy  $Z'$  search in DY Forward-Backward asymmetry

Ball et al arXiv: 2209.08115

# OUTLINE

- ➔ **Part I:** a short overview of recent developments in the determination of PDFs
- ➔ **Part II:** The precision vs accuracy challenge: a road-map to robustly test accuracy

# PART I: PROTON STRUCTURE AND PRECISION PHYSICS

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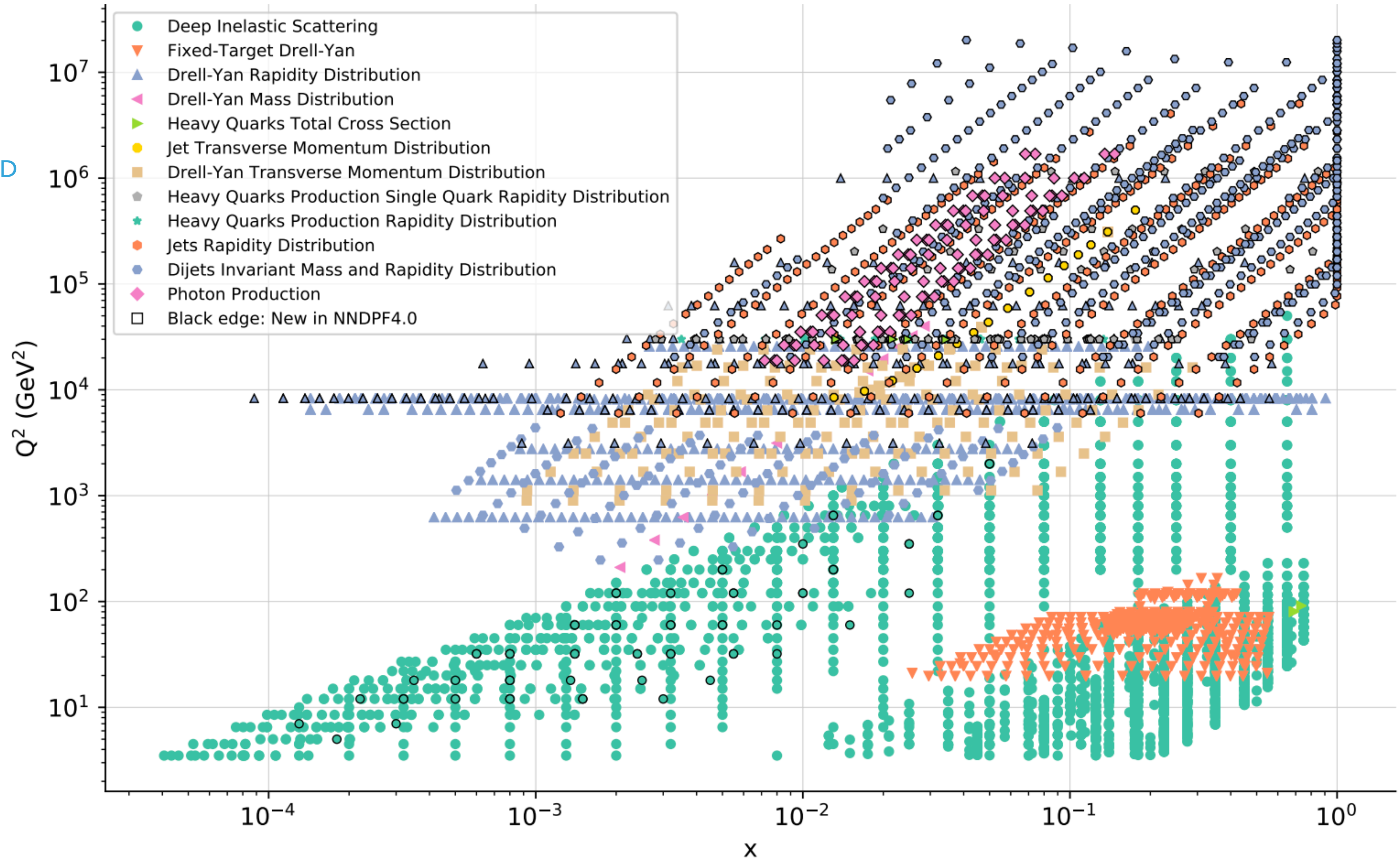
# PDF DETERMINATION

$$f_i(x, \mu)$$

Data  $\mu$  Perturbative QCD

- DGLAP evolution equations fully known up to NNLO and partially at N<sup>3</sup>LO
- Most theory predictions for processes in PDF fits known at NNLO (some at N<sup>3</sup>LO)
- Precision of most data of order few %
- Mostly from correlated systematic uncertainties

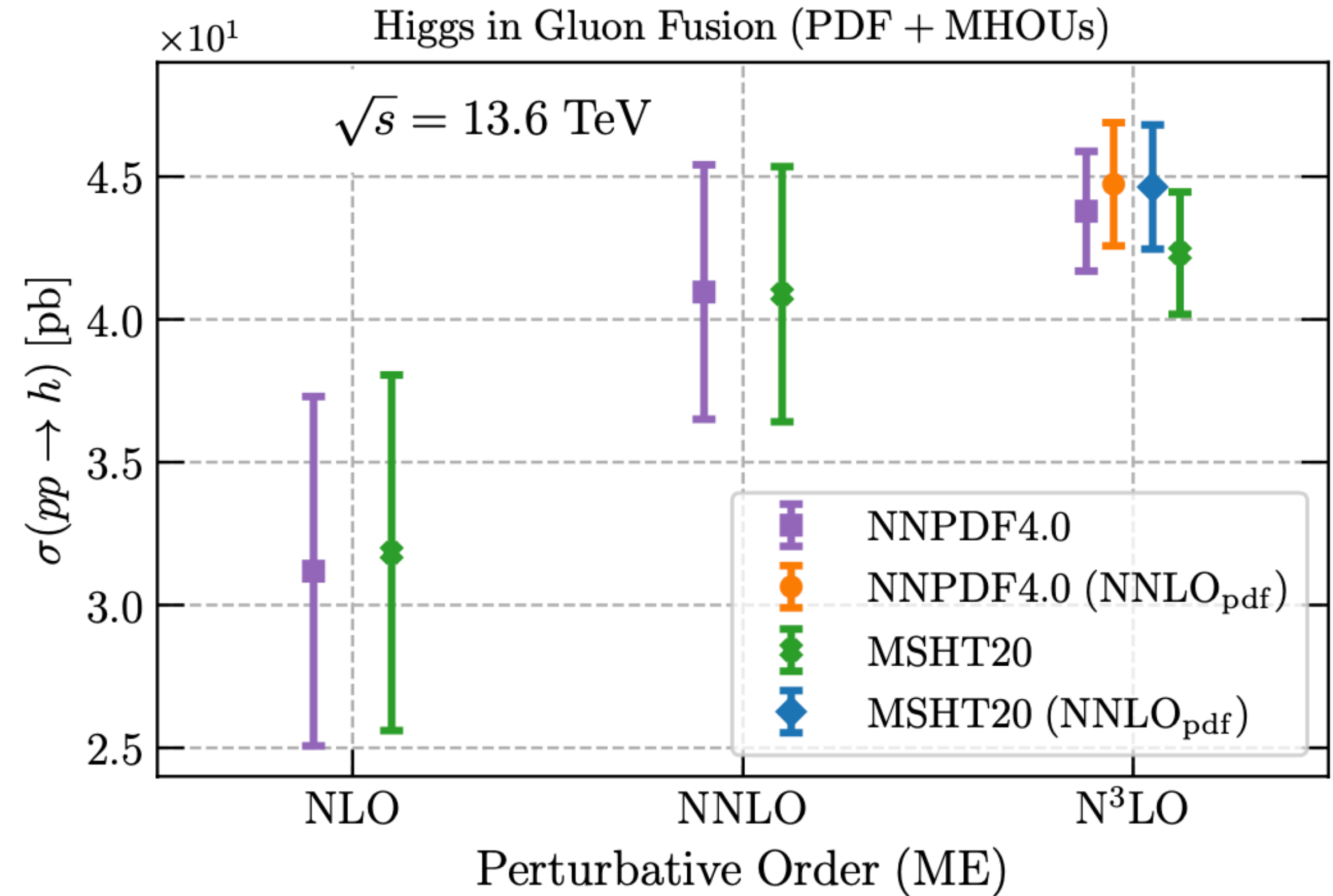
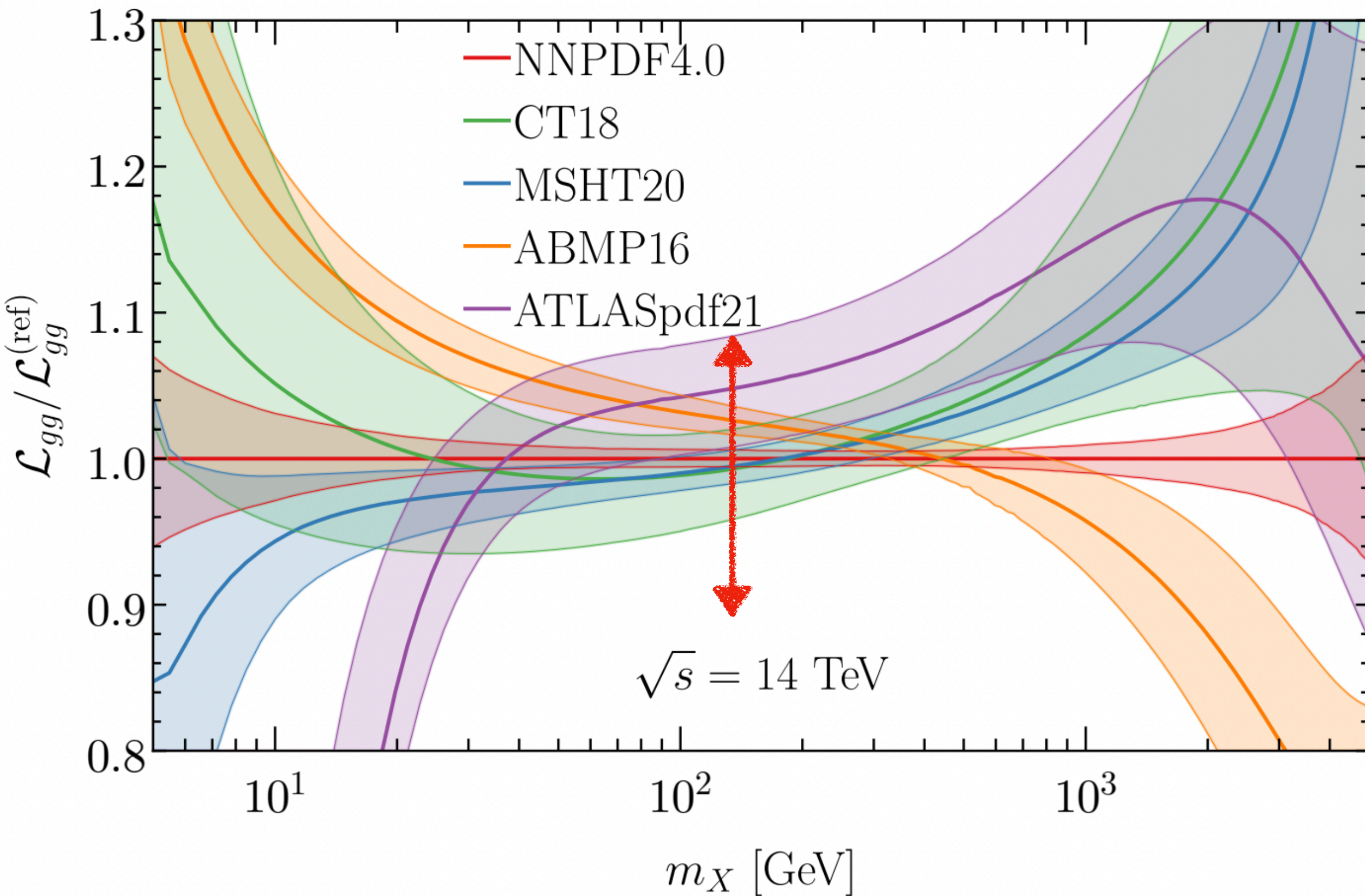
Ball et al, NNPDF4.0, arXiv:2109.02653



# WHERE WE STAND NOW

Ball et al 2402.18635

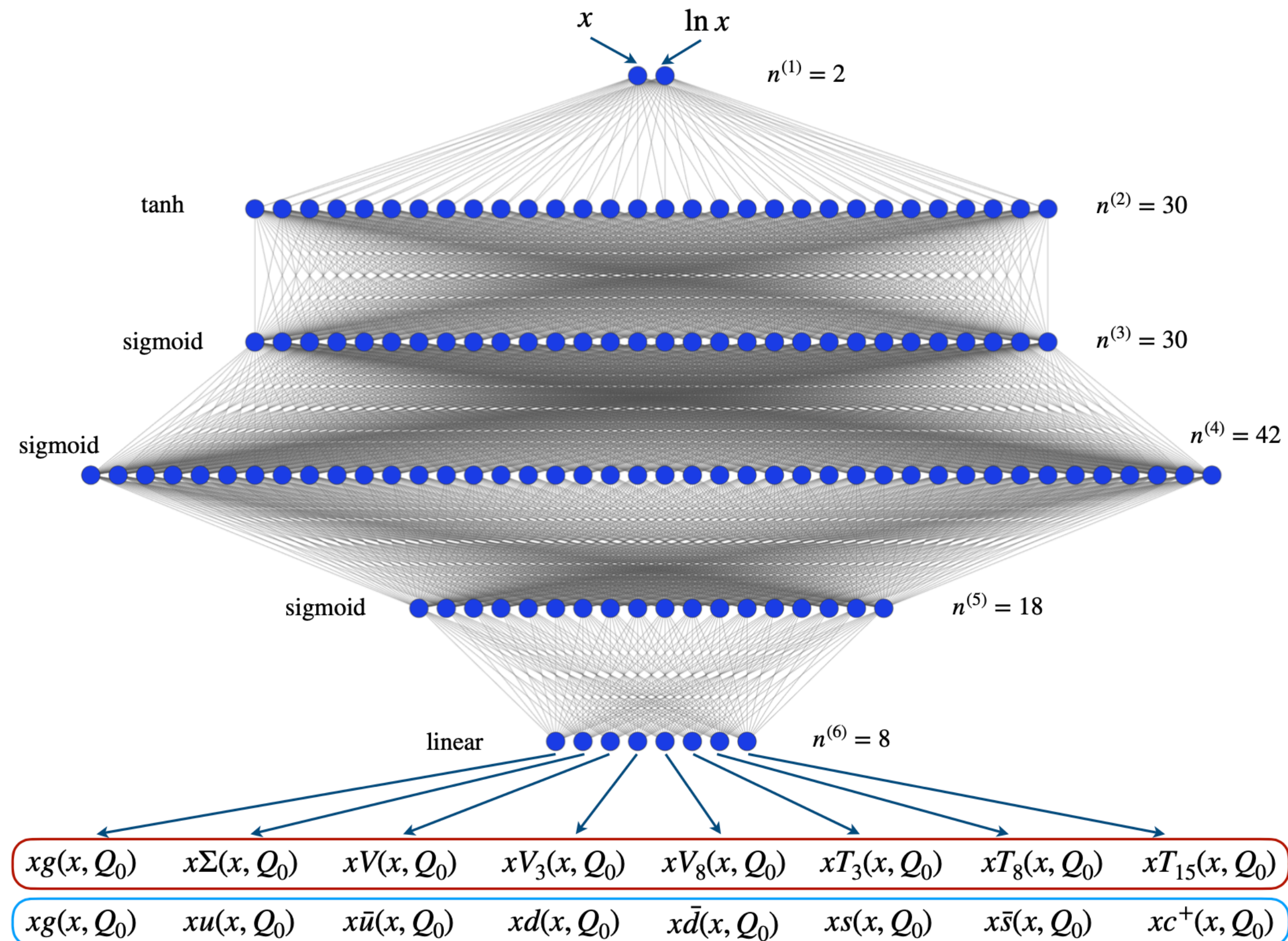
Snowmass 2021 white paper [arXiv:2203.13923]



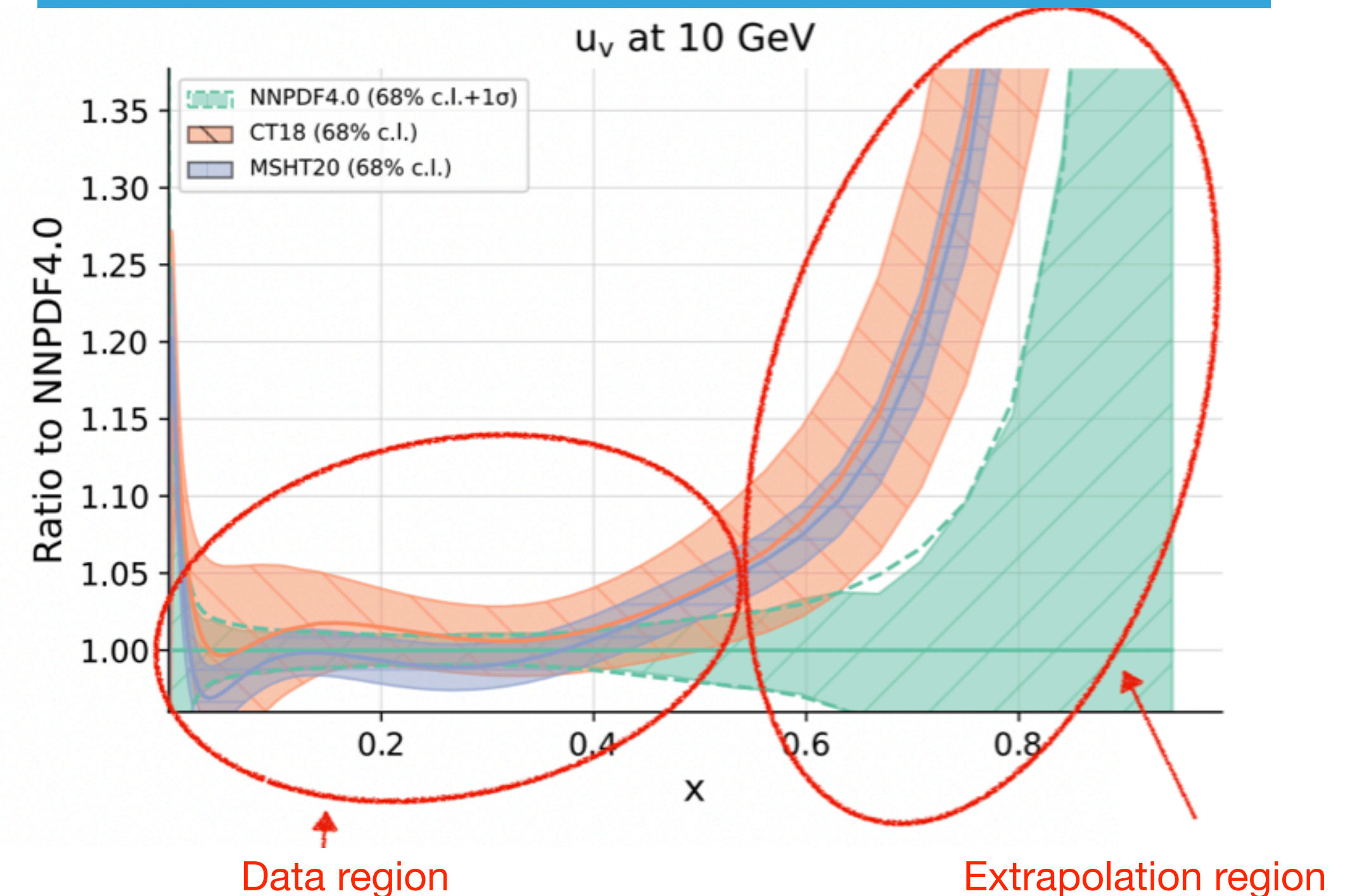
- **NNLO sets:** overall agreement around Higgs mass between global PDF sets with some shifts and large differences in PDF uncertainties due to differences in methodologies and datasets included.
- **Approx N3LO sets:** MSHT20aN3LO [arXiv:2207.04739] and NNPDF40aN3LO [arXiv:2402.18635] have been released

# NNPDF4.0: THE ROLE OF DATA AND METHODOLOGY

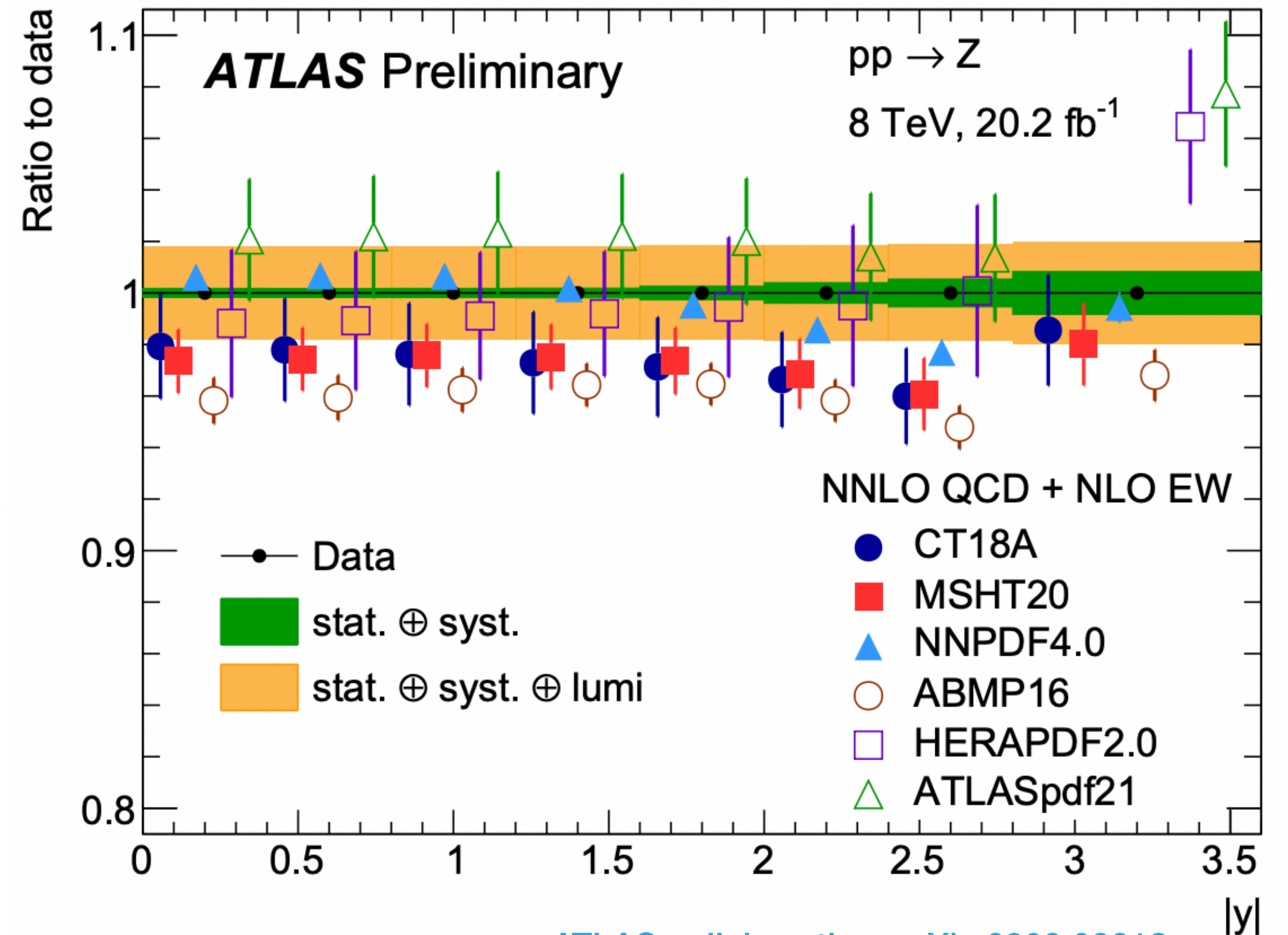
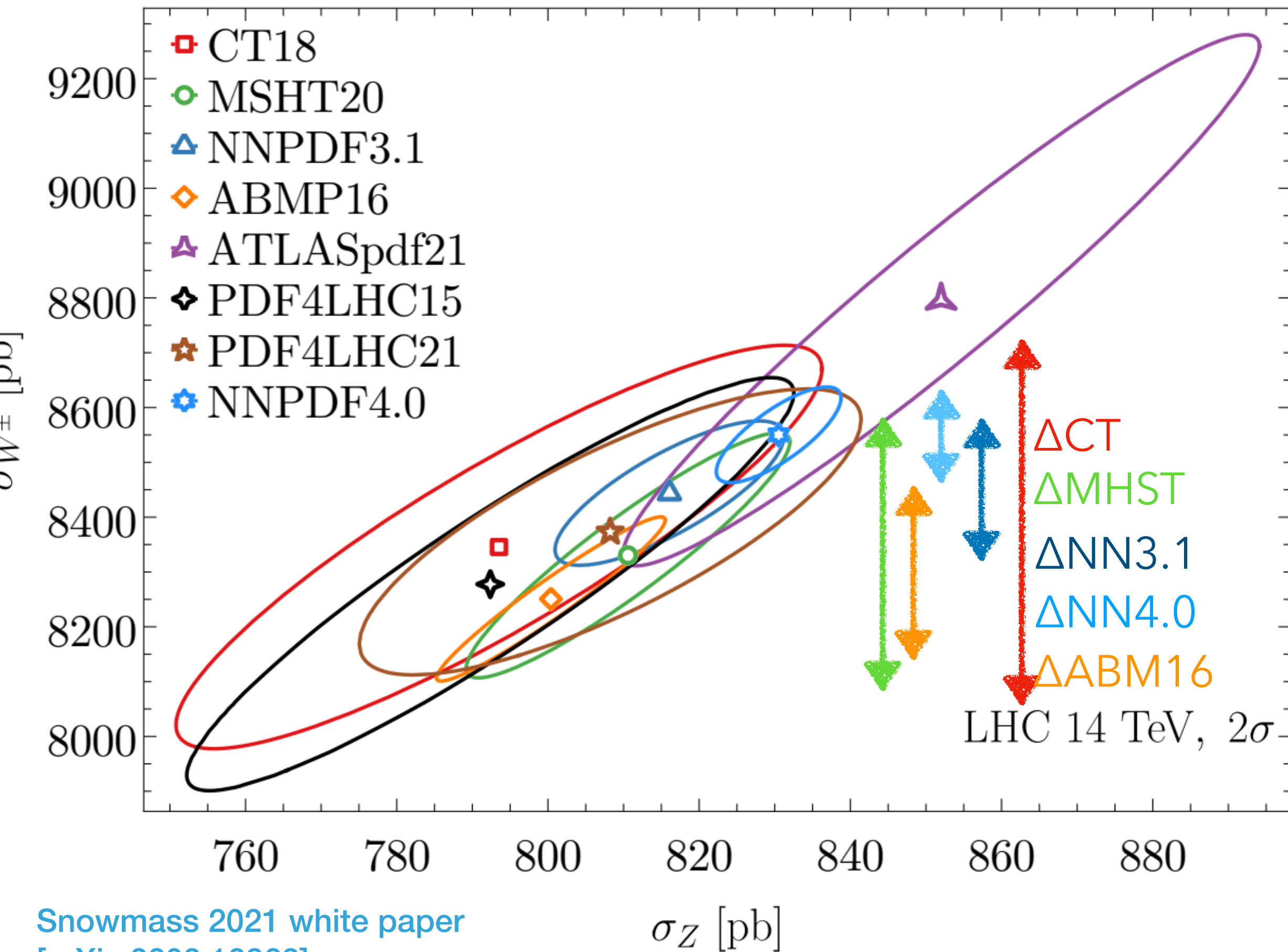
- NNPDF4.0 data: large portion of LHC data included, total of  $O(4500)$  points (shift in PDFs and some uncertainty reduction)
- NNPDF4.0 methodology: single neural network to parametrise 8 independent PDF combinations ( $g, u, d, s, u\sim, d\sim, s\sim, c=c\sim$ ) & optimisation strategy based on gradient descent rather than genetic algorithm (NNPDF3.1)
- Hyper-optimised methodology: scan of the hyper parameter space to find optimal minimisation settings (optimiser, initialiser, stopping patience, number of layers, learning rate, epochs, activation function) by minimising  $\chi^2_{\text{val}}$  [Carrazza et al, arXiv:1907.05075]



Thanks to superior methodology compared to NN31, NN40 displays smaller uncertainty than NN3.1,  $O(500)$  extra LHC data decrease uncertainties and shift PDFs.



# PHENOMENOLOGICAL IMPLICATIONS



ATLAS collaboration, arXiv:2309.09318

- NNPDF4.0 has the smallest uncertainty in the data region, towards 1% precision.
- Extremely precise predictions and small PDF uncertainties for electroweak observables: implications for extraction of SM precision parameters, like  $\alpha_s$  from Z  $p_T$  distribution and  $M_W$  from W  $m_T$  distributions [G. Bozzi's talk]
- Important to understand reason behind differences between different PDF sets & stress-test methodologies



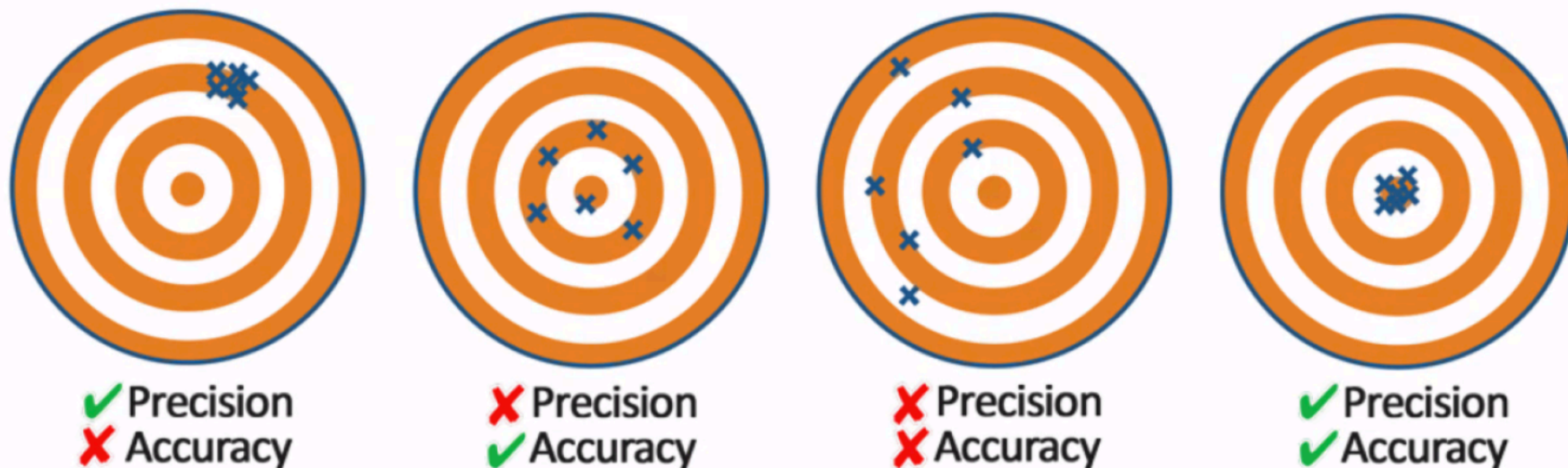
## **PART II: A ROADMAP TO TEST ACCURACY**

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# THE PRECISION VERSUS ACCURACY CHALLENGE

PDF4LHC21, 2203.05506

## PRECISION VS ACCURACY



Dataset	$N_{\text{pt}}$	$\chi^2/N_{\text{pt}}$		
		CT18	MSHT20	NNPDF3.1
BCDMS $F_2^p$	329/163 <sup>††</sup> /325 <sup>†</sup>	1.06	1.00	1.21
BCDMS $F_2^d$	246/151 <sup>††</sup> /244 <sup>†</sup>	1.06	0.88	1.10
NMC $F_2^d/F_2^p$	118/117 <sup>†</sup>	0.93	0.93	0.90
NuTeV dimuon $\nu + \bar{\nu}$	38+33	0.79	0.83	1.22
HERAI+II	1120	1.23	1.20	1.22
E866 $\sigma_{pd}/(2\sigma_{pp})$	15	1.24	0.80	0.43
LHCb 7 TeV & 8TeV $W,Z$	29+30	1.15	1.17	1.44
LHCb 8 TeV $Z \rightarrow ee$	17	1.35	1.43	1.57
ATLAS 7 TeV $W,Z$ (2016)	34	1.96	1.79	2.33
D0 $Z$ rapidity	28	0.56	0.58	0.62
CMS 7 TeV electron $A_{\text{ch}}$	11	1.47	1.52	0.76
ATLAS 7 TeV $W,Z$ (2011)	30	1.03	0.93	1.01
CMS 8TeV incl. jet	185/174 <sup>††</sup>	1.03	1.39	1.30
Total $N_{\text{pt}}$	—	2263	1991	2256
Total $\chi^2/N_{\text{pt}}$	—	1.14	1.15	1.20

### Challenges

- ▶ Inconsistency or tension in data of experimental origin (underestimate of systematics...)
- ▶ Deficiencies in fitting methodology (data-driven parametrisation change, optimisation issues, overfitting...)
- ▶ Inaccuracy in theoretical framework
  - ➔ Missing higher order uncertainties (QCD, EW)
  - ➔ Other corrections (nuclear, higher-twist, non-perturbative effects...)
- ▶ Fitting away possible BSM signals

# CLOSURE TEST: A TOOL TO TEST METHODOLOGY AND THEORY

- Closure tests for data region: imagine we knew the law of Nature  $f$ : is our fitting methodology able to reproduce it? Is the uncertainty faithful? Statistical validation of PDF uncertainties can be performed via closure tests.

L0 pseudo-data

$$y_0 = f + \eta$$

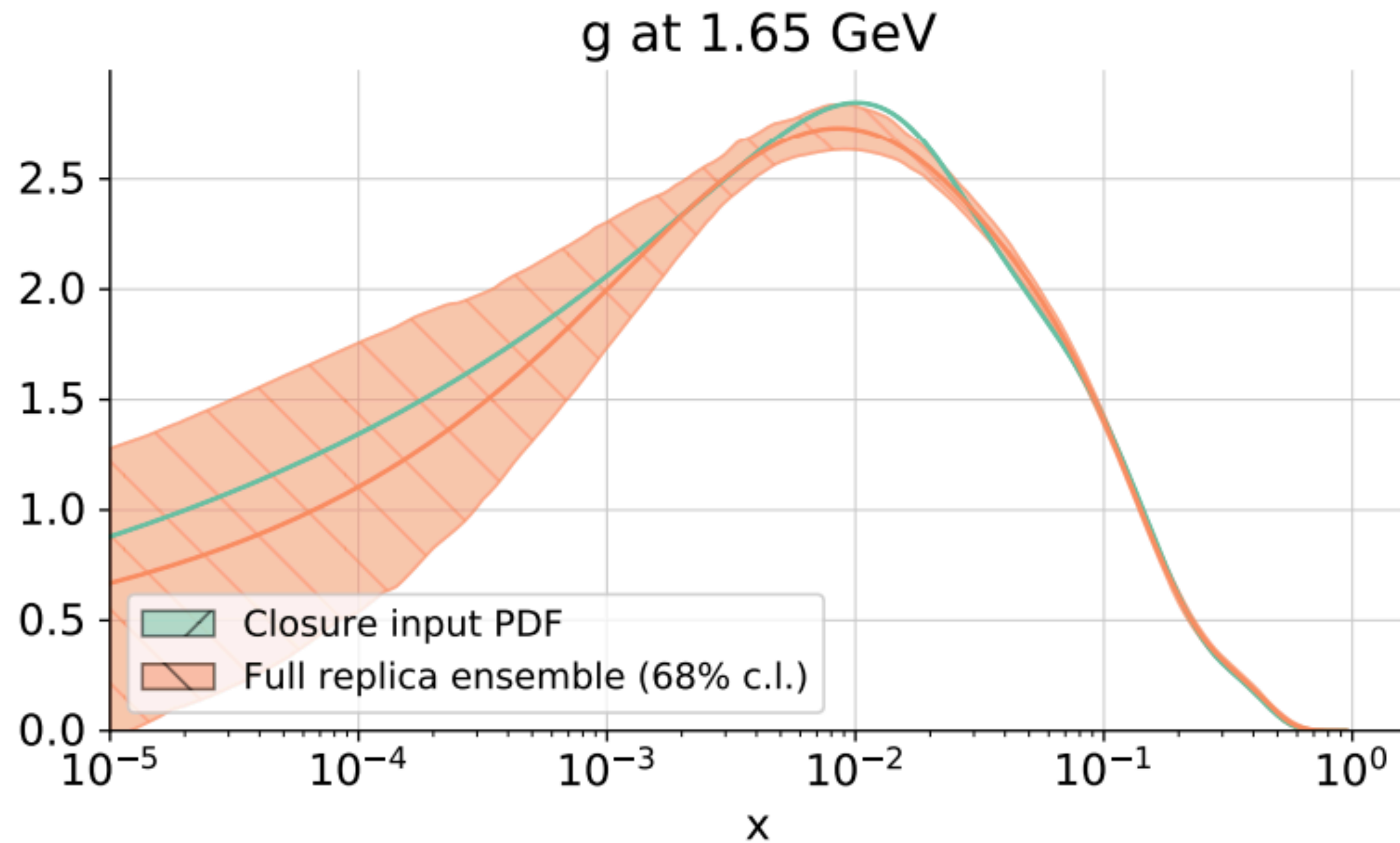
Experimental noise

Law of Nature

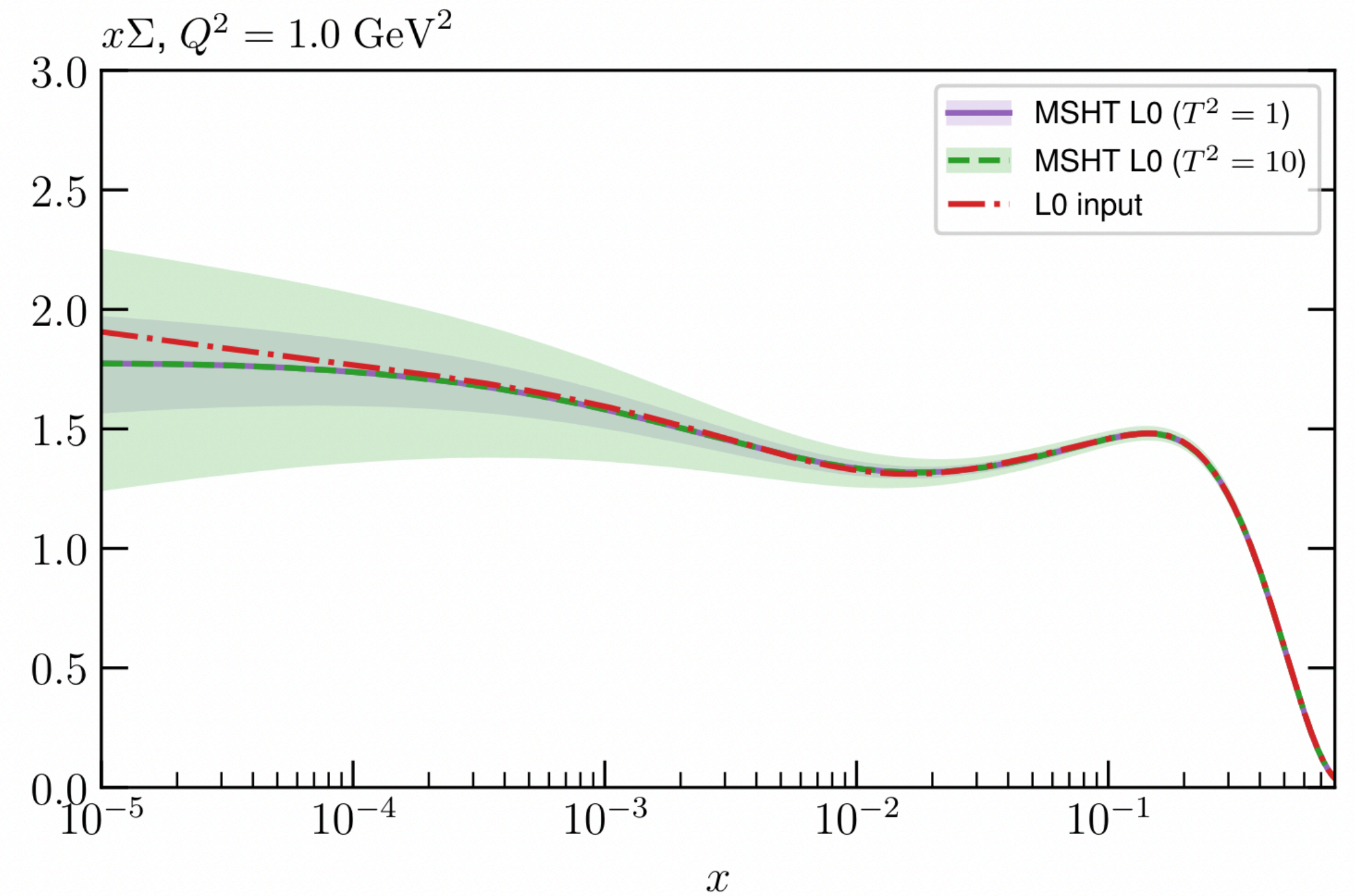
$$\hat{\sigma}_{\text{NNLO}} \otimes (f \otimes f) \text{ "true" PDF}$$

Test fitting methodology with consistent data (by construction)

Del Debbio et al, [arXiv: 2111.05787]



L. Harland Lang, DIS 2024



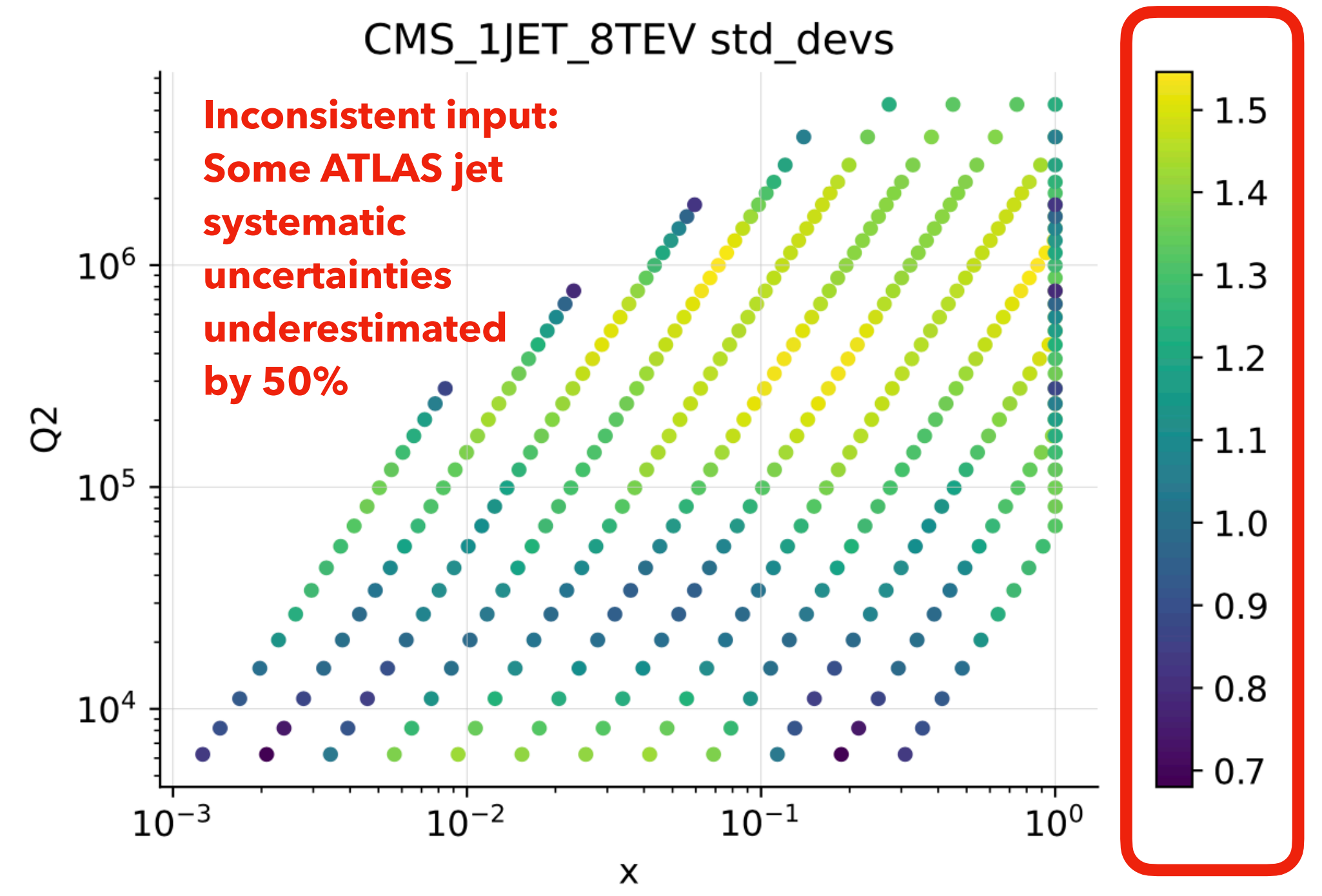
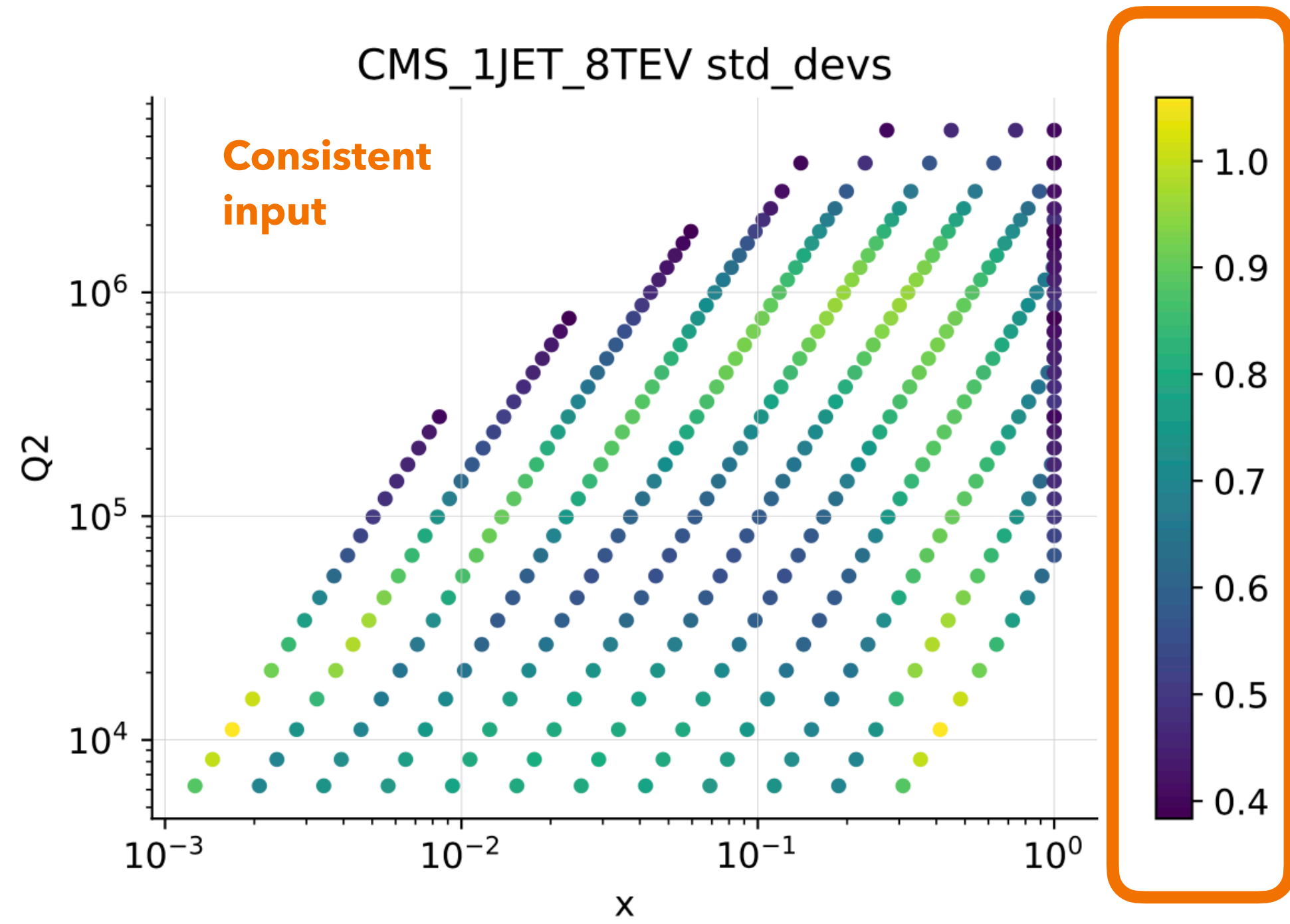
# CLOSURE TEST: A TOOL TO TEST METHODOLOGY AND THEORY

- Closure tests for data region: imagine we knew the law of Nature  $f$ : is our fitting methodology able to reproduce it? Is the uncertainty faithful? Statistical validation of PDF uncertainties can be performed via closure tests.
- What happens if experimentalists underestimated some systematics?

Statistical indicators of under ( $>1$ ) or over ( $<1$ ) estimated uncertainties allow to check effect of possible experimental inconsistencies on datasets included and on the PDFs

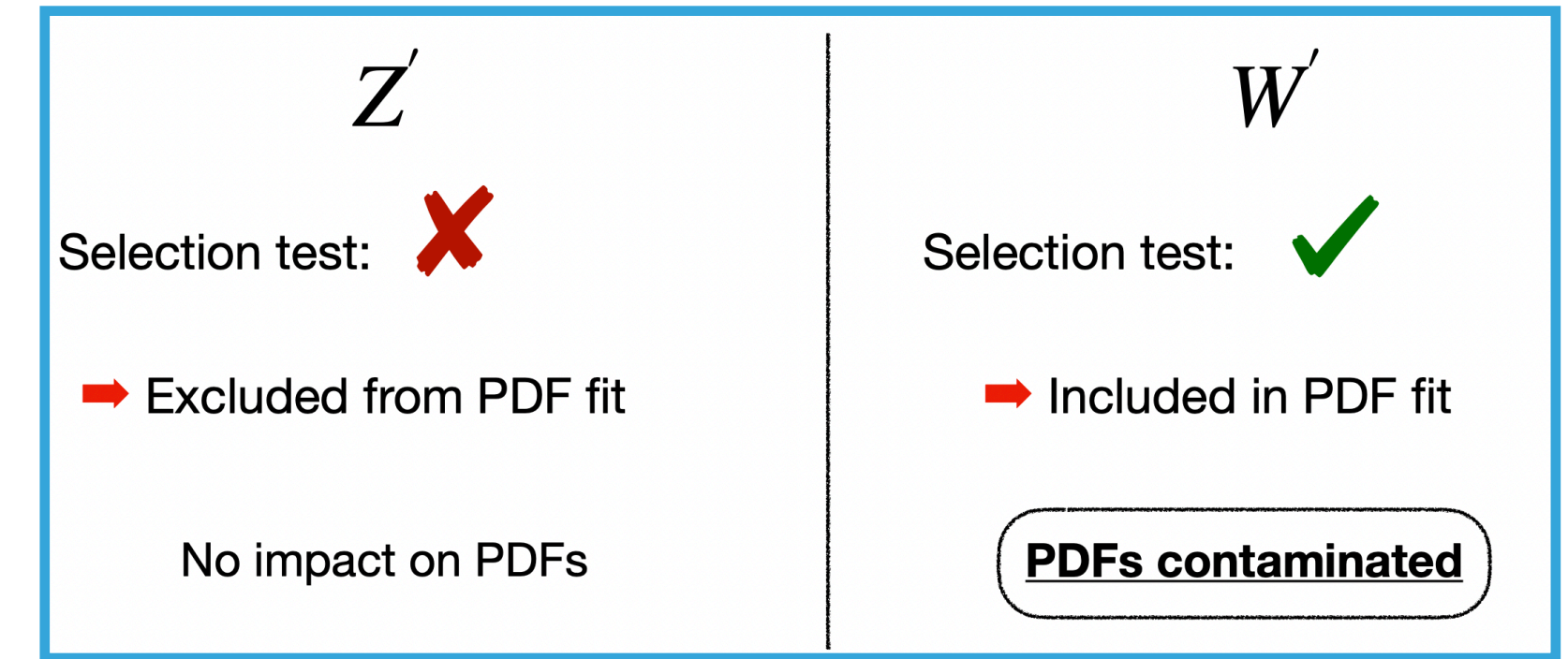
Test fitting methodology with inconsistent data

[Barontini, Costantini, De Crescenzo, Ubiali - in progress]

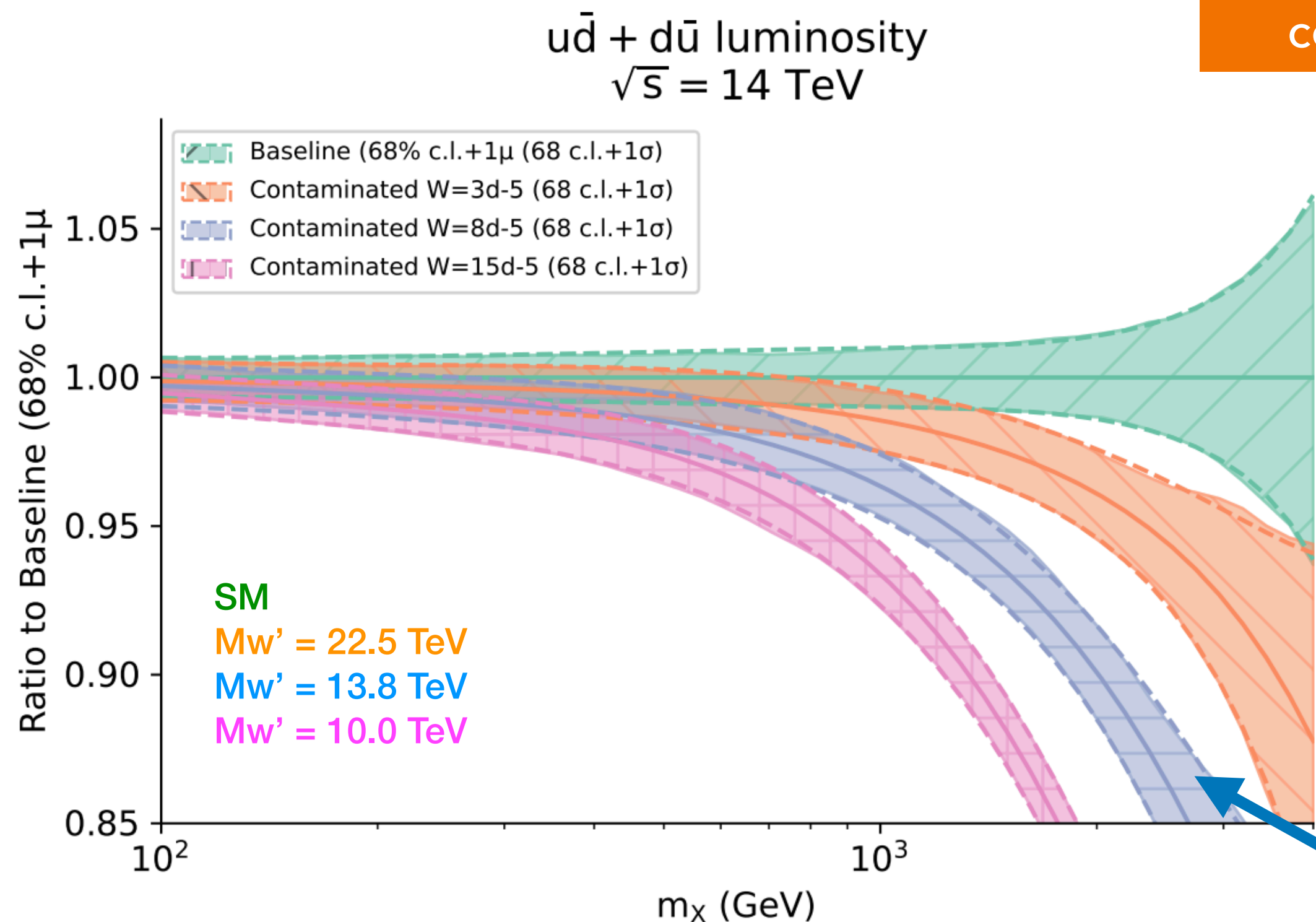


# CLOSURE TEST: A TOOL TO TEST METHODOLOGY AND THEORY

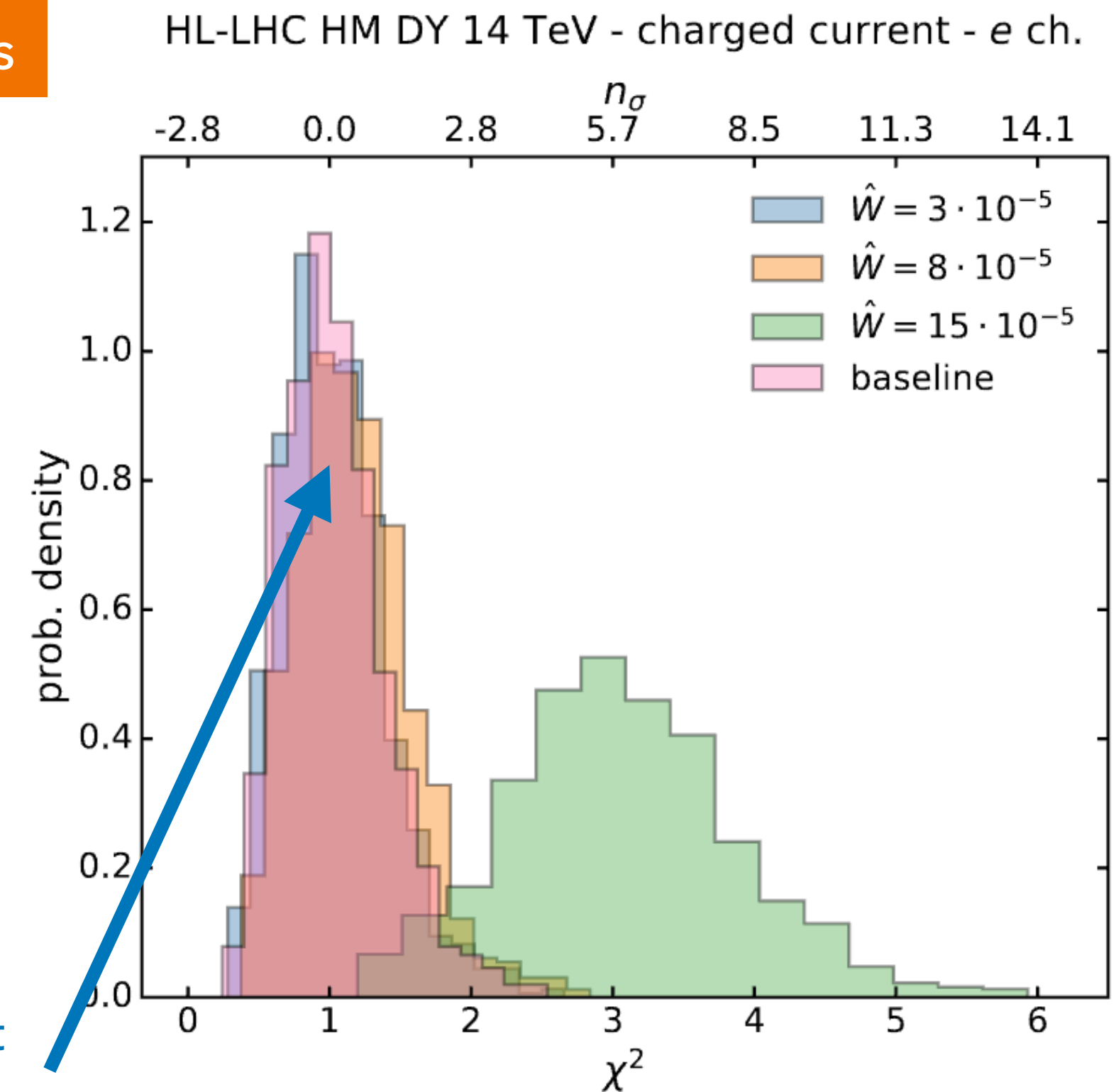
- Imagine that on top of the "true" PDFs one inject the "true" NP model in the pseudo-data
- Generate HL-LHC pseudo-data assuming "true" law of nature = "true" PDFs + "true" UV model
- Fit PDFs assuming SM
- Can PDFs absorb signs of new physics?



Test possible New Physics contamination in PDF fits



Max contamination allowed by global fit  
Without spoiling  $\chi^2$



# TEST GENERALISATION OF PDF AND EXTRAPOLATION

- Future tests help to discriminate among PDF sets [J. Cruz-Martinez et al, *Acta Phys.Polon.B* 52 (2021) 243 - on Run I ]
- [Chiefa, Costantini, Cruz-Martinez et al, in progress]: test all global PDF sets againsts new precise data from LHC Run I and Run II data & DIS HERA jets data.

How well do various PDF sets describe data that are not yet included in the fit?

- Remarks:

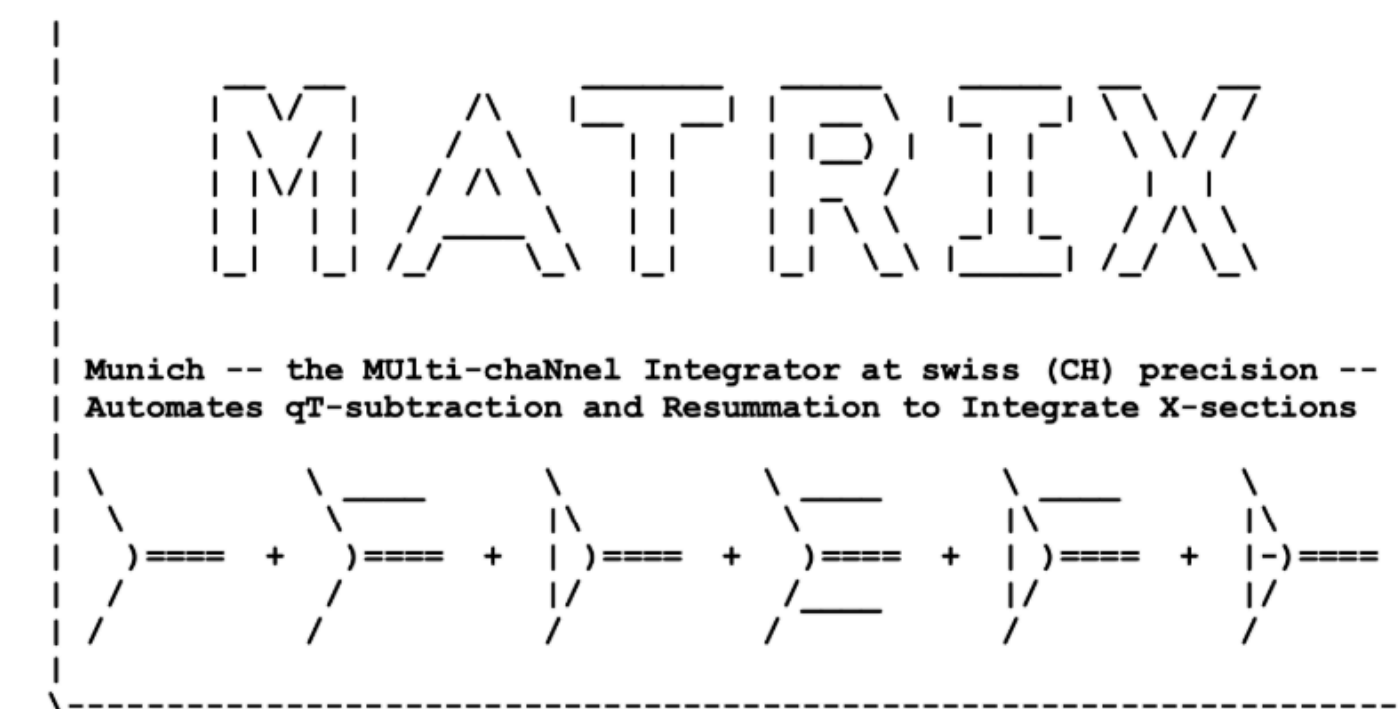
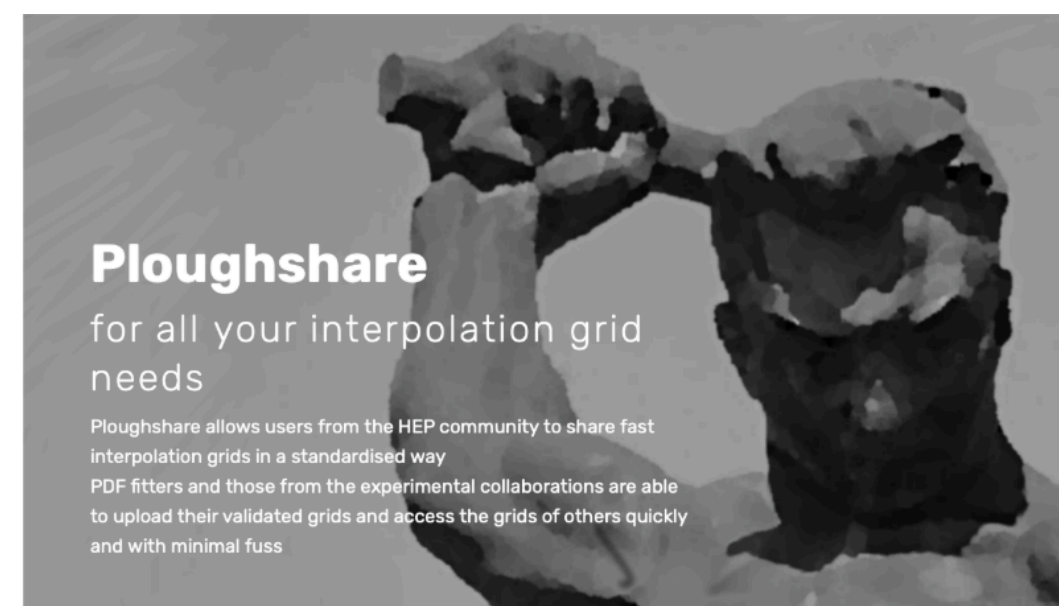
- ➔ All results are NNLO (no k-factor approximation) thanks to PineAPPL, NNLOJET, MATRIX and Ploughshare

PineAPPL

Carrazza et al: 2008.12789



Gehrmann-De Ridder et al: 1507.02850, 1605.04295...



Grazzini et al: 1711.06631

- ➔ NNLO PDF sets considered: PDF4LHC15, PDF4LHC21, ABMP16, CT18, CT18A, CT18Z, MSHT20, NNPDF3.1, NNPDF4.0
- ➔ The computation of  $\chi^2$  is always shown considering as uncertainties either (exp) or (exp + mho) or (exp + mho + pdf) and over the entire dataset / HEPdata entry

$$\chi^2 = \sum_{i,j=1}^{N_{\text{dat}}} \left( T_i^{(0)} - D_i \right) (\text{cov}^{-1})_{ij} \left( T_j^{(0)} - D_j \right)$$

$$\begin{aligned} \chi_{\text{exp}}^2 &\leftrightarrow \text{COV} = \text{COV}_{\text{exp}} \\ \chi_{\text{exp+mho}}^2 &\leftrightarrow \text{COV} = \text{COV}_{\text{exp}} + \text{COV}_{\text{mho}} \\ \chi_{\text{exp+mho+pdf}}^2 &\leftrightarrow \text{COV} = \text{COV}_{\text{exp}} + \text{COV}_{\text{mho}} + \text{COV}_{\text{pdf}} \end{aligned}$$

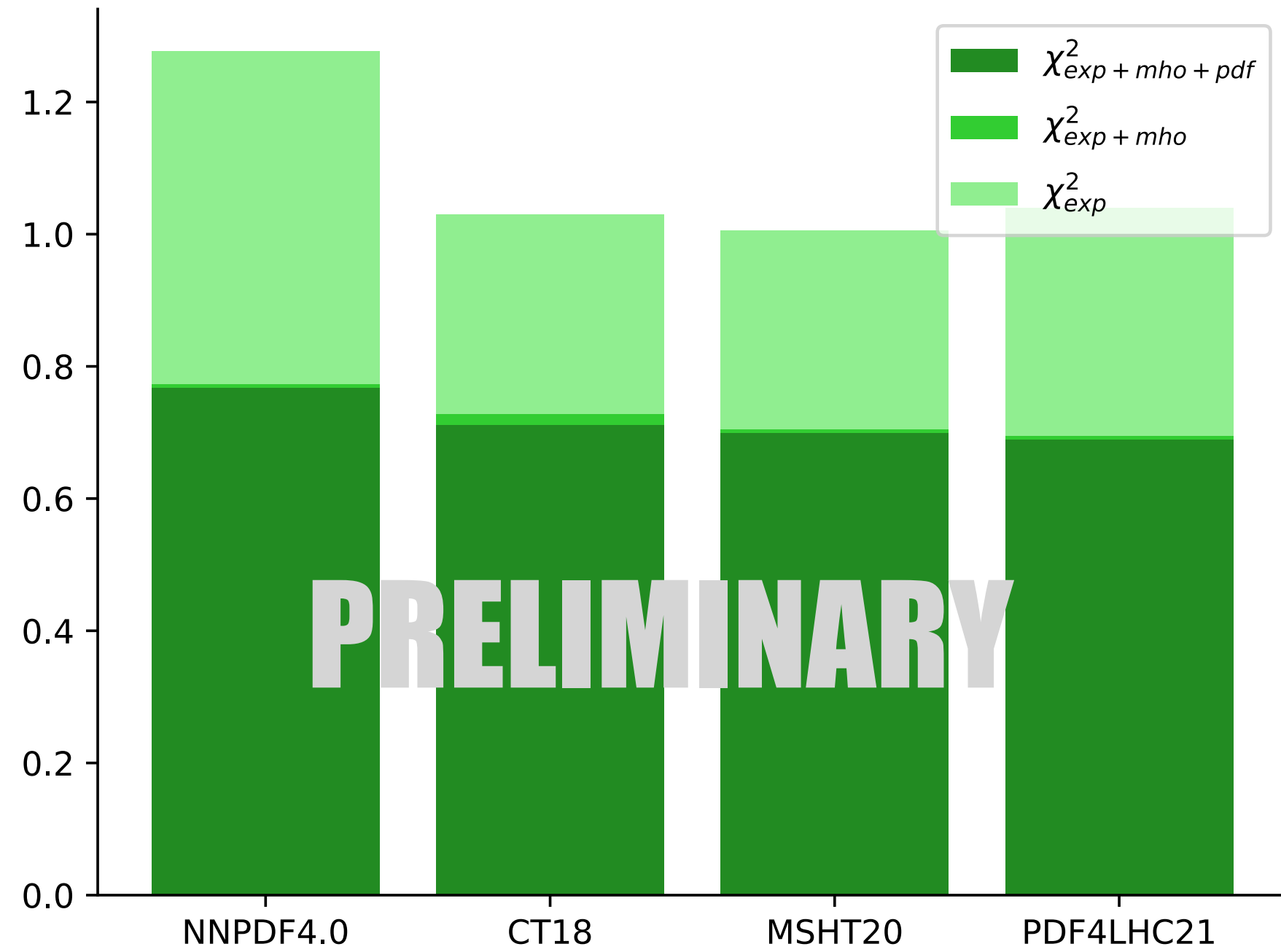
# THE EXPERIMENTAL DATA IN THE TEST SET

Sector	Exp.	$\sqrt{s}$ (TeV)	Channel	Observable	$\mathcal{L}$ (fb $^{-1}$ )	$N_{\text{dat}}$
$W, Z$	ATLAS	13	Z $p_T$ spectrum	$\frac{d\sigma}{dp_T^Z}$	36.1	10
	ATLAS	8	Z incl. prod.	$\frac{d\sigma}{d y_{ll} }$	20.2	7
	CMS	8	W incl. prod.	$\frac{d\sigma_{W^\pm}}{d\eta_l}$	35.9	36
	LHCb	13	Z incl. forward prod.	$\frac{d\sigma_Z}{dy_Z}$	5.1	17
top	ATLAS	13	hadronic	$(\frac{1}{\sigma}) \frac{d\sigma}{dm_{t\bar{t}}}, \frac{d\sigma}{d y_{t\bar{t}} }, \frac{d^2\sigma}{d y_{t\bar{t}} dm_{t\bar{t}}}$	36.1	9, 12, 11
	ATLAS	13	$\ell$ +jets	$(\frac{1}{\sigma}) \frac{d\sigma}{dm_{t\bar{t}}}, \frac{d\sigma}{dp_{T,t}}, \frac{d\sigma}{d y_t }, \frac{d^2\sigma}{d y_{t\bar{t}} }$	36.1	9, 8, 5, 7
	CMS	13	$\ell$ +jets	$(\frac{1}{\sigma}) \frac{d\sigma}{dm_{t\bar{t}}}, \frac{d\sigma}{dp_{T,t}}, \frac{d\sigma}{d y_{t\bar{t}} }, \frac{d\sigma}{d y_t }, \frac{d^2\sigma}{d y_{t\bar{t}} dm_{t\bar{t}}}$	137	15, 16, 10, 11, 35
jets	ATLAS	13	incl. jet R=0.4, 0.7	$\frac{d^2\sigma}{dp_{T,j}d y_j }$	3.2	177
	ATLAS	13	di-jets R=0.4	$\frac{d^2\sigma}{dm_{jj}d\Delta y}$	3.2	136
	CMS	13	incl. jets R=0.4, 0.7	$\frac{d^2\sigma}{dp_{T,j}d y_j }$	3.2	78
DIS jets	H1	0.319	incl. jet (low $q^2$ )	$\frac{d^2\sigma}{dq^2 dp_T}$	0.29	48
	H1	0.319	di-jets (low $q^2$ )	$\frac{d^2\sigma}{dq^2 d\langle p_T \rangle}$	0.29	48
	H1	0.319	incl. jet (high $q^2$ )	$\frac{d^2\sigma}{dq^2 dp_T}$	0.351	24
	H1	0.319	di-jets (high $q^2$ )	$\frac{d^2\sigma}{dq^2 d\langle p_T \rangle}$	0.351	24
	ZEUS	0.3	incl. jet	$\frac{d^2\sigma}{dE_T dq^2}$	0.038	30
	ZEUS	0.319	incl. jet	$\frac{d^2\sigma}{dE_T dq^2}$	0.082	30
	ZEUS	0.319	d-jets	$\frac{d^2\sigma}{dE_T dq^2}$	0.374	22

# THE TOP SECTOR

CMS top pair production at 13 TeV, l+jet channel (137fb<sup>-1</sup>)

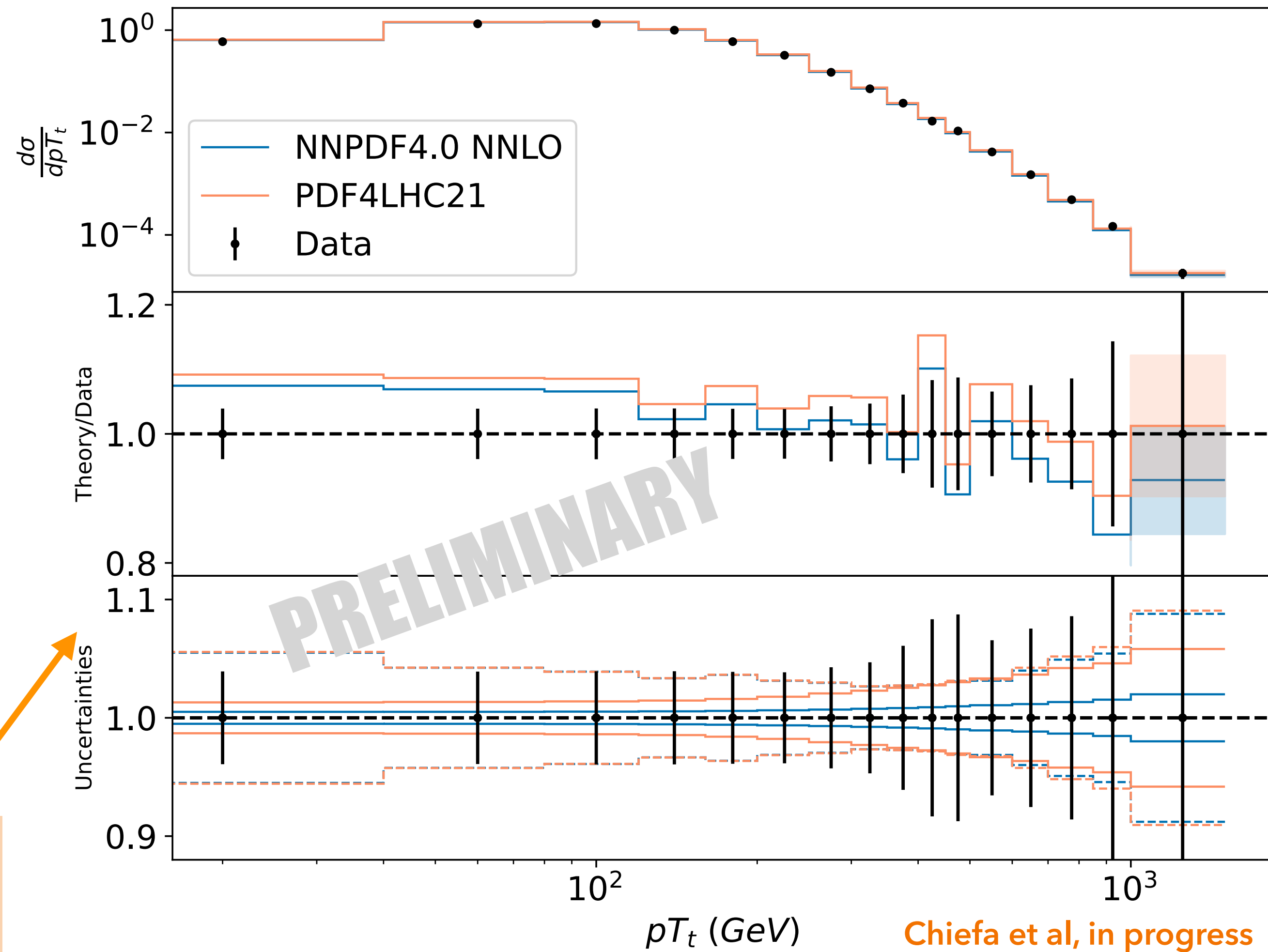
CMS collaboration [arXiv:2108.02803]



PDF set	$\frac{d\sigma}{dpT_t}$
ABMP16	0.986 (1.253)
CT18	0.712 (0.729)
CT18A	0.678 (0.686)
MSHT20	0.699 (0.705)
NNPDF3.1	0.782 (0.794)
NNPDF4.0	0.768 (0.773)
PDF4LHC15	0.709 (0.722)
PDF4LHC21	0.689 (0.695)

Absolute comparison, normalized & size of PDF (solid) and dashed (MHO)  $\Delta$  compared to data  $\Delta$ . Shaded band includes MHO (9 pts variation) and PDF uncertainties added in quadrature.

CMS 13 TeV top quark pair l+j channel:  $\frac{d\sigma}{dpT_t}$



Chiefa et al, in progress

Predictions: MATRIX

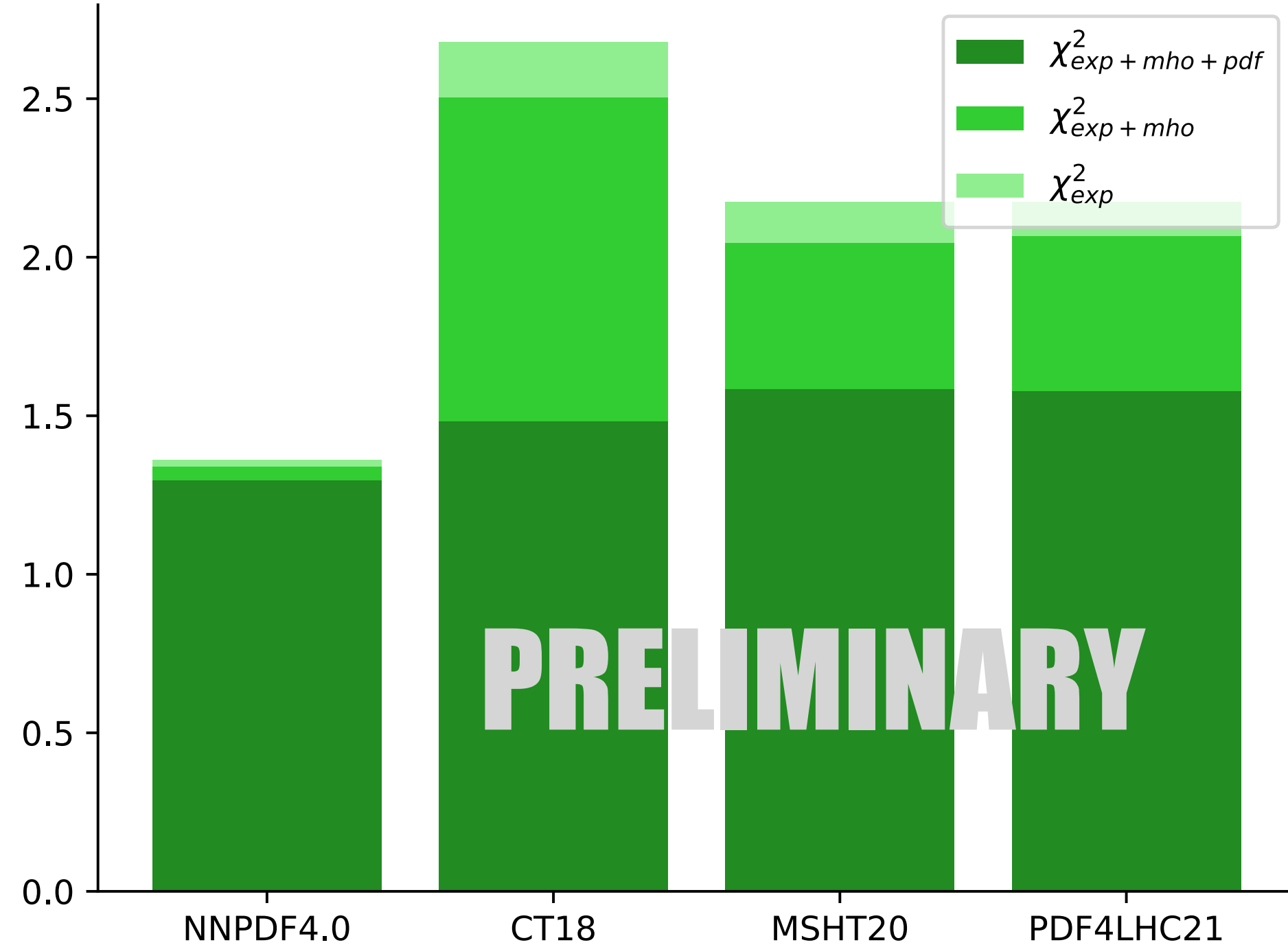
Hepdata: 10.17182/hepdata.102956



# THE TOP SECTOR

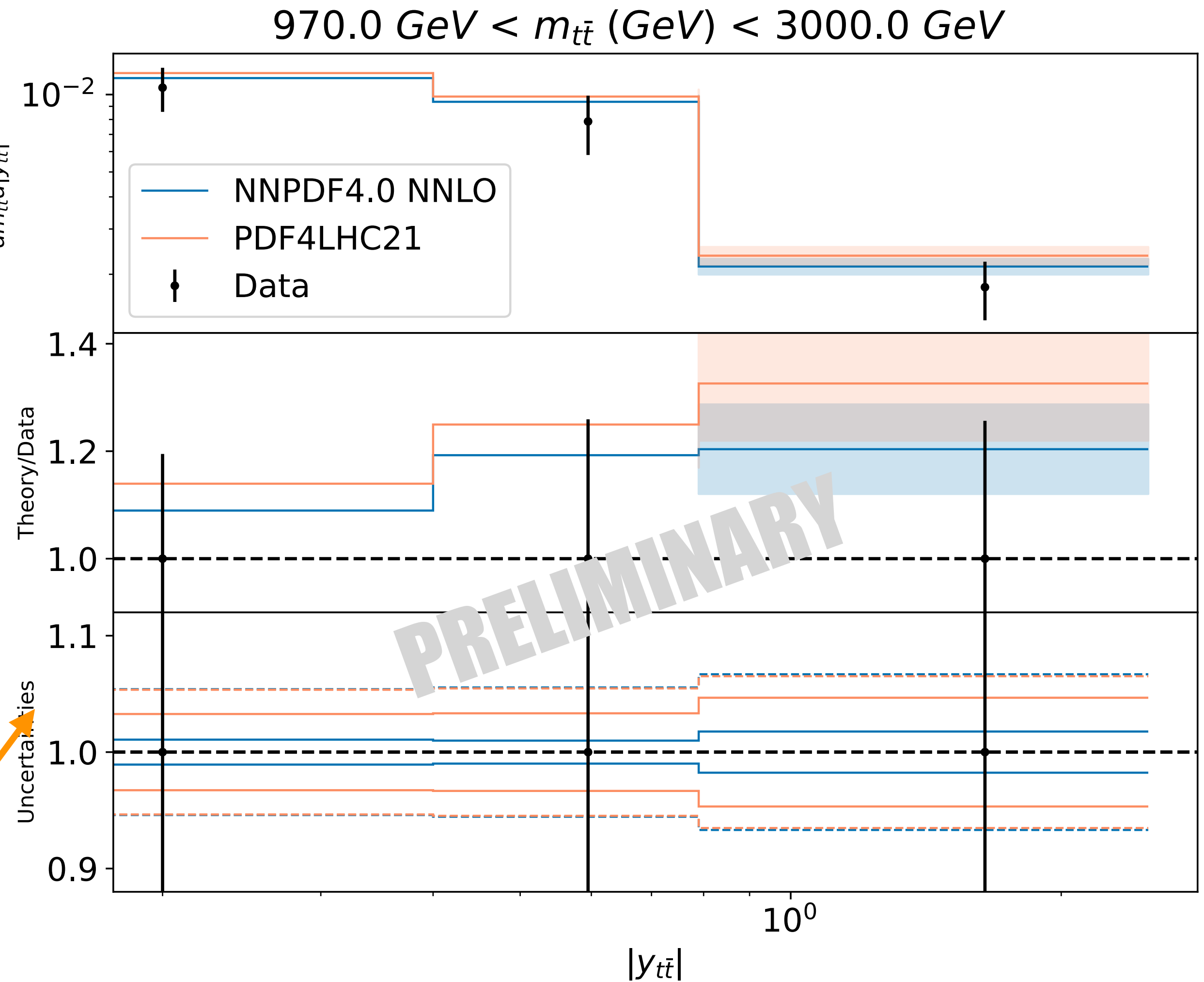
ATLAS top pair production at 13 TeV, l+jet channel, (36.1fb<sup>-1</sup>)

ATLAS collaboration [arXiv:1908.07305]



PDF set	$\frac{d^2\sigma}{dm_{t\bar{t}}d y_{t\bar{t}} }$
ABMP16	0.832 (0.835)
CT18	1.483 (2.504)
CT18A	1.489 (2.347)
MSHT20	1.585 (2.046)
NNPDF3.1	1.200 (1.244)
NNPDF4.0	1.297 (1.338)
PDF4LHC15	1.298 (2.088)
PDF4LHC21	1.577 (2.068)

Absolute comparison, normalized & size of PDF (solid) and dashed (MHO)  $\Delta$  compared to data  $\Delta$ . Shaded band includes MHO (9 pts variation) and PDF uncertainties added in quadrature.



Predictions: MATRIX

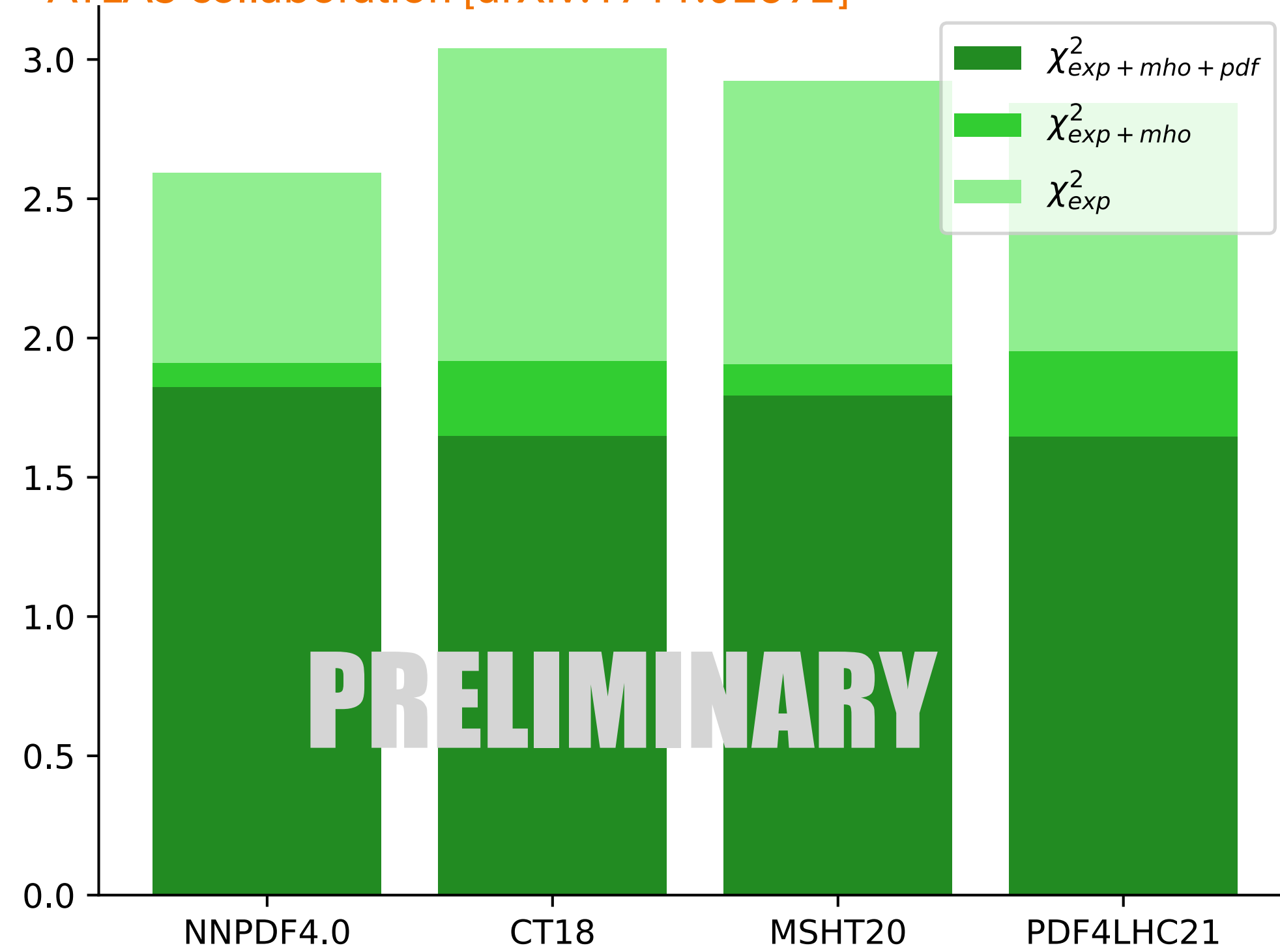
Hepdata: 10.17182/hepdata.95758

Chiefa et al, in progress

# THE JET SECTOR

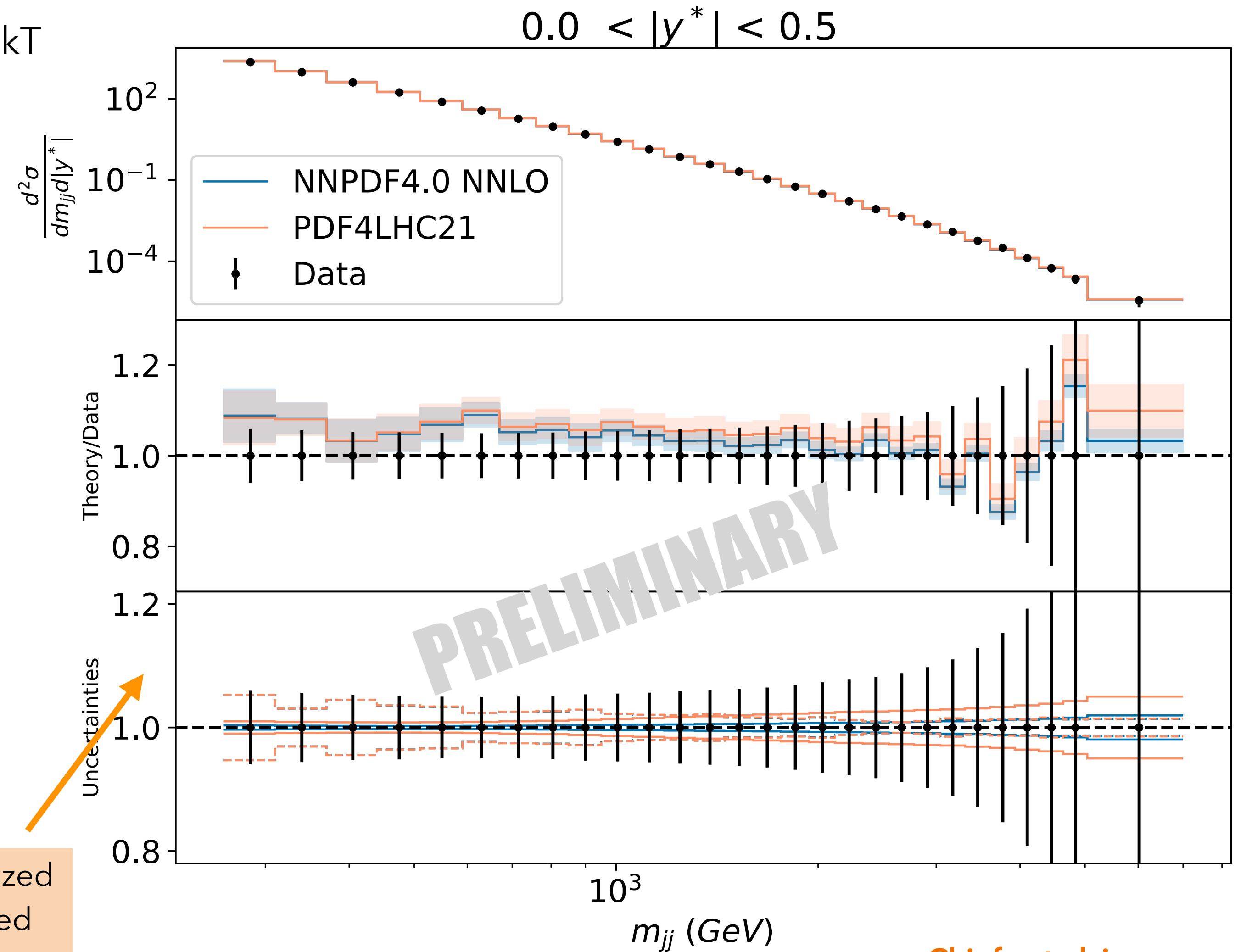
ATLAS di-jet cross-sections at 13 TeV (3.2fb<sup>-1</sup>), R=0.4, anti-kT

ATLAS collaboration [arXiv:1711.02692]



PDF set	
ABMP16	1.72 (2.12)
CT18	1.65 (1.92)
CT18A	1.67 (1.91)
CT18Z	1.61 ( <b>1.88</b> )
MSHT20	1.79 (1.91)
‘ NNPDF3.1	1.78 (2.03)
NNPDF4.0	1.82 (1.91)
PDF4LHC15	<b>1.59</b> (1.91)
PDF4LHC21	1.65 (1.95)

Absolute comparison, normalized & size of PDF (solid) and dashed (MHO)  $\Delta$  compared to data  $\Delta$ . Shaded band includes MHO (9 pts variation) and PDF uncertainties added in quadrature.



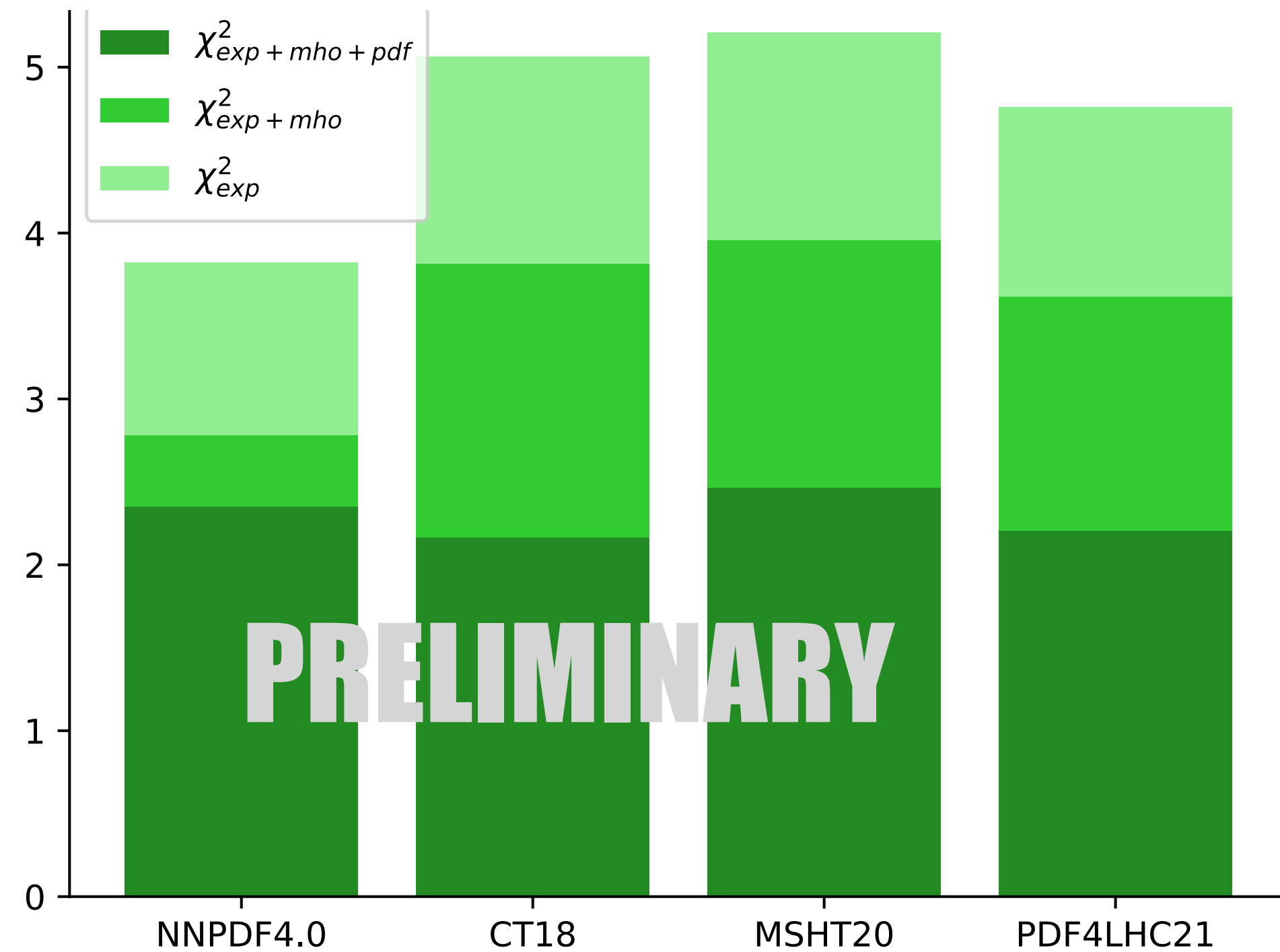
Chiefa et al, in progress

Leading color NNLO correction  
 Predictions: NNLOJET (plougshare)  
 Hepdata: 10.17182/hepdata.79952

# THE JET SECTOR

CMS inclusive jet cross-sections at 13 TeV (33.5fb<sup>-1</sup>) anti-kT, R=0.7

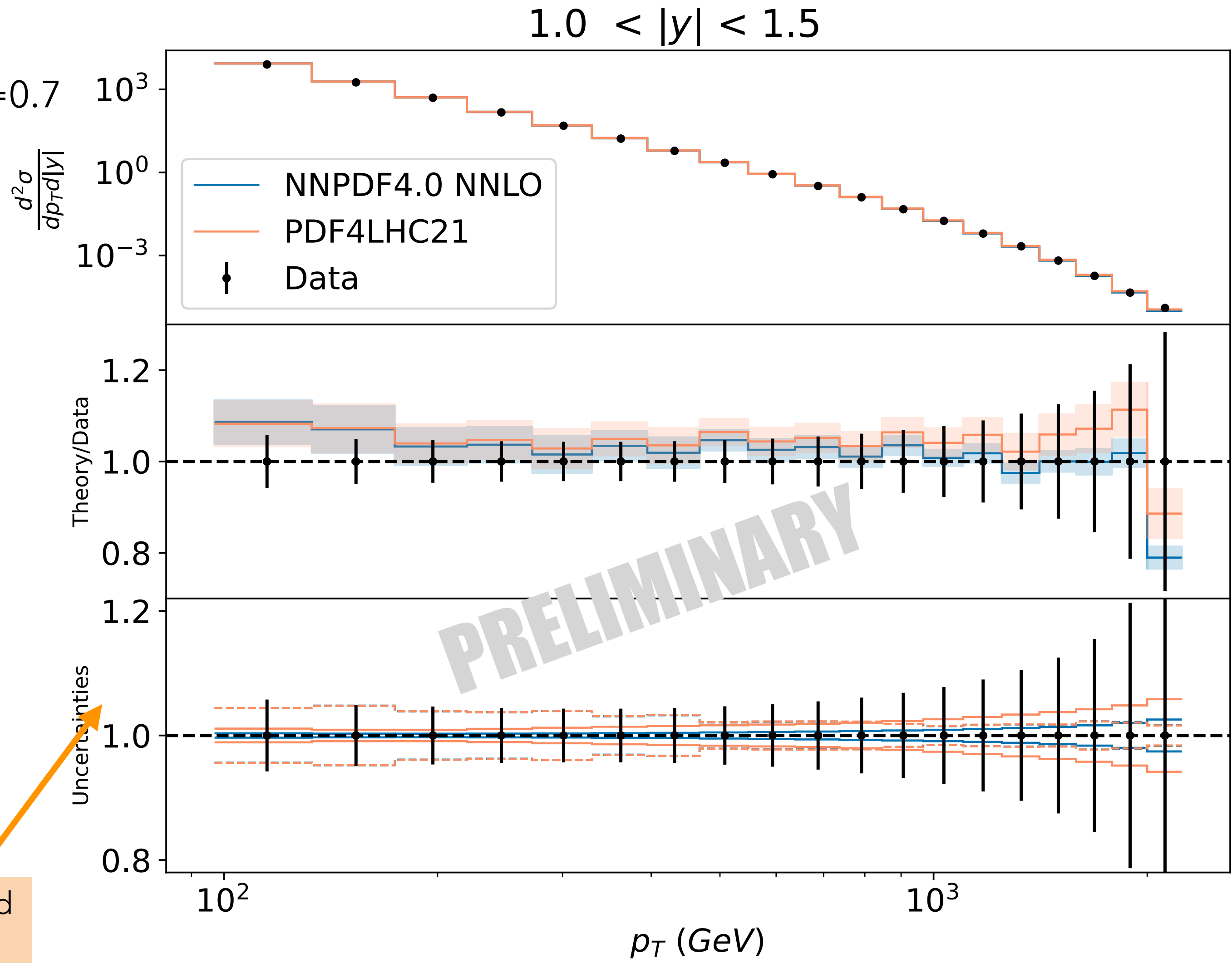
CMS collaboration [arXiv:2111.10431]



PRELIMINARY

PDF set	
ABMP16	2.77 (3.71)
CT18	2.16 (3.81)
CT18A	2.16 (3.77)
CT18Z	<b>2.06</b> (3.34)
MSHT20	2.47 (3.96)
‘ NNPDF3.1	2.52 (3.88)
NNPDF4.0	2.35 ( <b>2.78</b> )
PDF4LHC15	2.15 (3.80)
PDF4LHC21	2.21 (3.62)

Absolute comparison, normalized & size of PDF (solid) and dashed (MHO)  $\Delta$  compared to data  $\Delta$ . Shaded band includes MHO (9 pts variation) and PDF uncertainties added in quadrature.



1.0 < |y| < 1.5

— NNPDF4.0 NNLO  
— PDF4LHC21  
• Data

$$\frac{d^2\sigma}{dp_T dy}$$

Theory/Data

Uncertainties

10<sup>2</sup>

10<sup>3</sup>

p<sub>T</sub> (GeV)

Chiefa et al, in progress

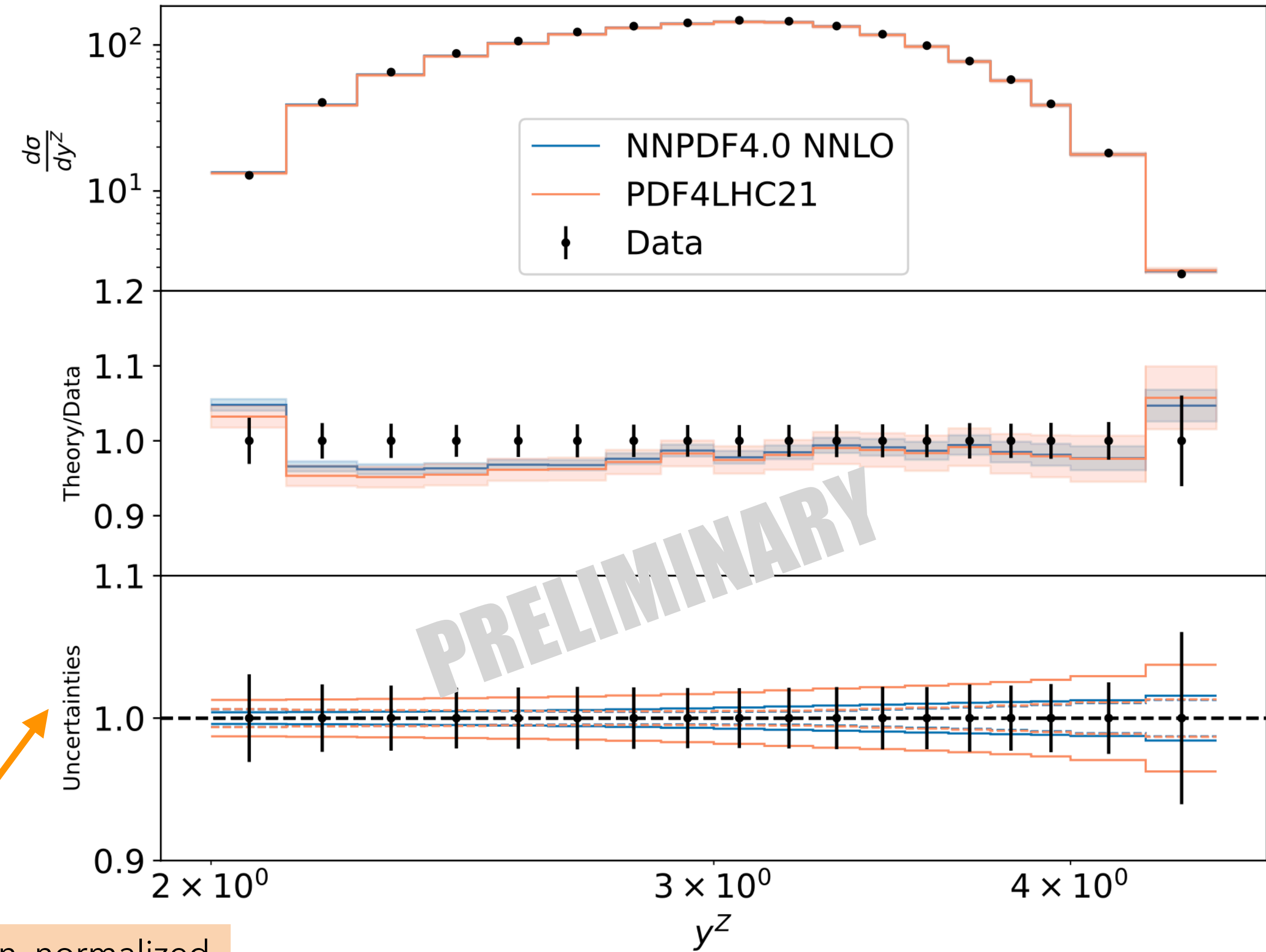
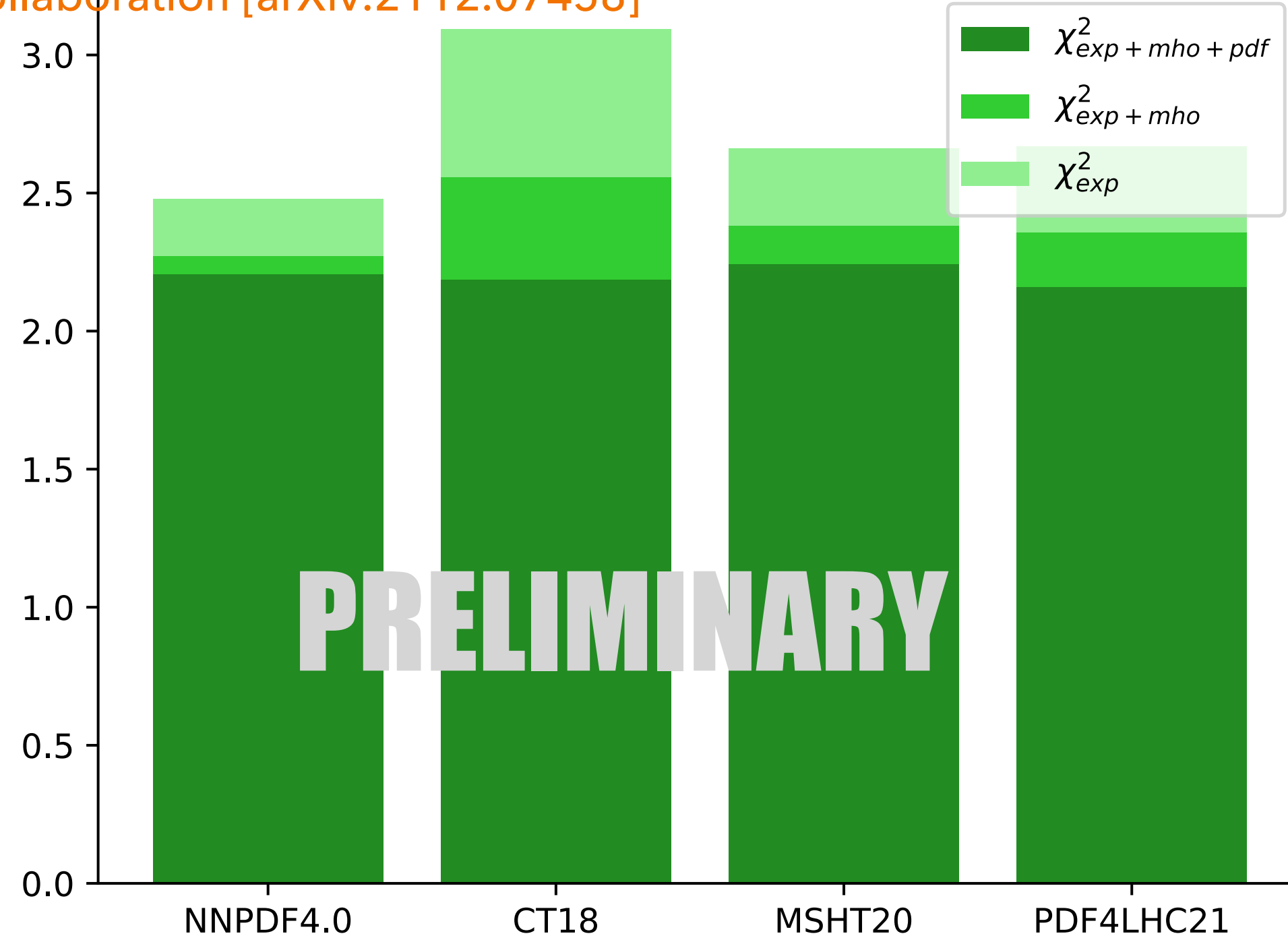
Leading color NNLO correction  
Predictions: NNLOJET (plougshare)  
Hepdata: 10.17182/hepdata.115022.v2

# THE DRELL-YAN SECTOR

LHCb Z forward production at 13 TeV (5.1fb<sup>-1</sup>)

LHCb collaboration [arXiv:2112.07458]

LHCb Z → μμ



PDF set	LHCb $\frac{d\sigma}{dy^Z}$ 13 TeV
ABMP16	2.22 ( <b>2.26</b> )
CT18	2.19 (2.56)
CT18A	2.25 (2.49)
MSHT20	2.24 (2.38)
NNPDF3.1	2.49 (2.73)
NNPDF4.0	2.21 (2.27)
PDF4LHC15	2.17 (2.59)
PDF4LHC21	<b>2.16</b> (2.36)

Absolute comparison, normalized & size of PDF (solid) and dashed (MHO)  $\blacktriangle$  compared to data  $\blacktriangle$ . Shaded band includes MHO (9 pts variation) and PDF uncertainties added in quadrature.

Chiefa et al, in progress

Predictions: MATRIX  
 Hepdata: 10.17182/hepdata.102956

- In an era of precision at LHC, need precise and accurate PDFs
- Tools like closure tests are crucial to test methodology robustness, effects of possible experimental inconsistencies and even possible effects of new physics in the high energy tails
- Progress possible thanks to automated tools for fast computation of observables and sharing (NNLOJET, MATRIX, PineAPPL, PloughShare, APPLgrid, FastNLO) and public PDF tools (XFitter, NNPDF)
- Comparing performance of broad range of PDF sets on unseen data we see that, while NNPDF4.0 PDF uncertainties are smaller, the data-theory description is similar to that of other PDF set
- Predictions for NNPDF4.0 seem to fall close the data, hinting that the improvement in precision (smaller uncertainties than NNPDF3.1) came with an improvement in accuracy
- Oscillations far from the data can have big MHO uncertainties, luckily the era of N3LO PDFs is starting now!



**EXTRA MATERIAL**

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# NNPDF4.0: THE ROLE OF LHC DATA

Data set	NNPDF3.1	NNPDF4.0	ABMP16	CT18	MSHT20
CMS $W$ asym. 7 TeV ( $\mathcal{L} = 36 \text{ pb}^{-1}$ )	✗	✗	✗	✗	✓
CMS $Z$ 7 TeV ( $\mathcal{L} = 36 \text{ pb}^{-1}$ )	✗	✗	✗	✗	✓
CMS $W$ electron asymmetry 7 TeV	✓	✓	✗	✓	✓
CMS $W$ muon asymmetry 7 TeV	✓	✓	✓	✓	✗
CMS Drell-Yan 2D 7 TeV	✓	✓	✗	(✓)	✓
CMS Drell-Yan 2D 8 TeV	(✓)	✗	✗	✗	✗
CMS $W$ rapidity 8 TeV	✓	✓	✓	✓	✓
CMS $W, Z$ $p_T$ 8 TeV ( $\mathcal{L} = 18.4 \text{ fb}^{-1}$ )	✗	✗	✗	(✓)	✗
CMS $Z$ $p_T$ 8 TeV	✓	✓	✗	(✓)	✗
CMS $W + c$ 7 TeV	✓	✓	✗	(✓)	✓
CMS $W + c$ 13 TeV	✗	✓	✗	✗	(✓)
CMS single-inclusive jets 2.76 TeV	✓	✗	✗	✗	✓
CMS single-inclusive jets 7 TeV	✓	(✓)	✗	✓	✓
CMS dijets 7 TeV	✗	✓	✗	✗	✗
CMS single-inclusive jets 8 TeV	✗	✓	✗	✓	✓
CMS 3D dijets 8 TeV	✗	(✓)	✗	✗	✗
CMS $\sigma_{tt}^{\text{tot}}$ 5 TeV	✗	✓	✗	✗	✗
CMS $\sigma_{tt}^{\text{tot}}$ 7, 8 TeV	✓	✓	✗	✗	✗
CMS $\sigma_{tt}^{\text{tot}}$ 8 TeV	✗	✗	✗	✗	✓
CMS $\sigma_{tt}^{\text{tot}}$ 5, 7, 8, 13 TeV	✗	✗	✓	✗	✗
CMS $\sigma_{tt}^{\text{tot}}$ 13 TeV	✓	✓	✓	✗	✗
CMS $t\bar{t}$ lepton+jets 8 TeV	✓	✓	✗	✗	✓
CMS $t\bar{t}$ 2D dilepton 8 TeV	✗	✓	✗	✓	✓
CMS $t\bar{t}$ lepton+jet 13 TeV	✗	✓	✗	✗	✗
CMS $t\bar{t}$ dilepton 13 TeV	✗	✓	✗	✗	✗
CMS single top $\sigma_t + \sigma_{\bar{t}}$ 7 TeV	✗	✓	✓	✗	✗
CMS single top $R_t$ 8, 13 TeV	✗	✓	✓	✗	✗
CMS single top 13 TeV	✗	✗	✗	✗	(✓)

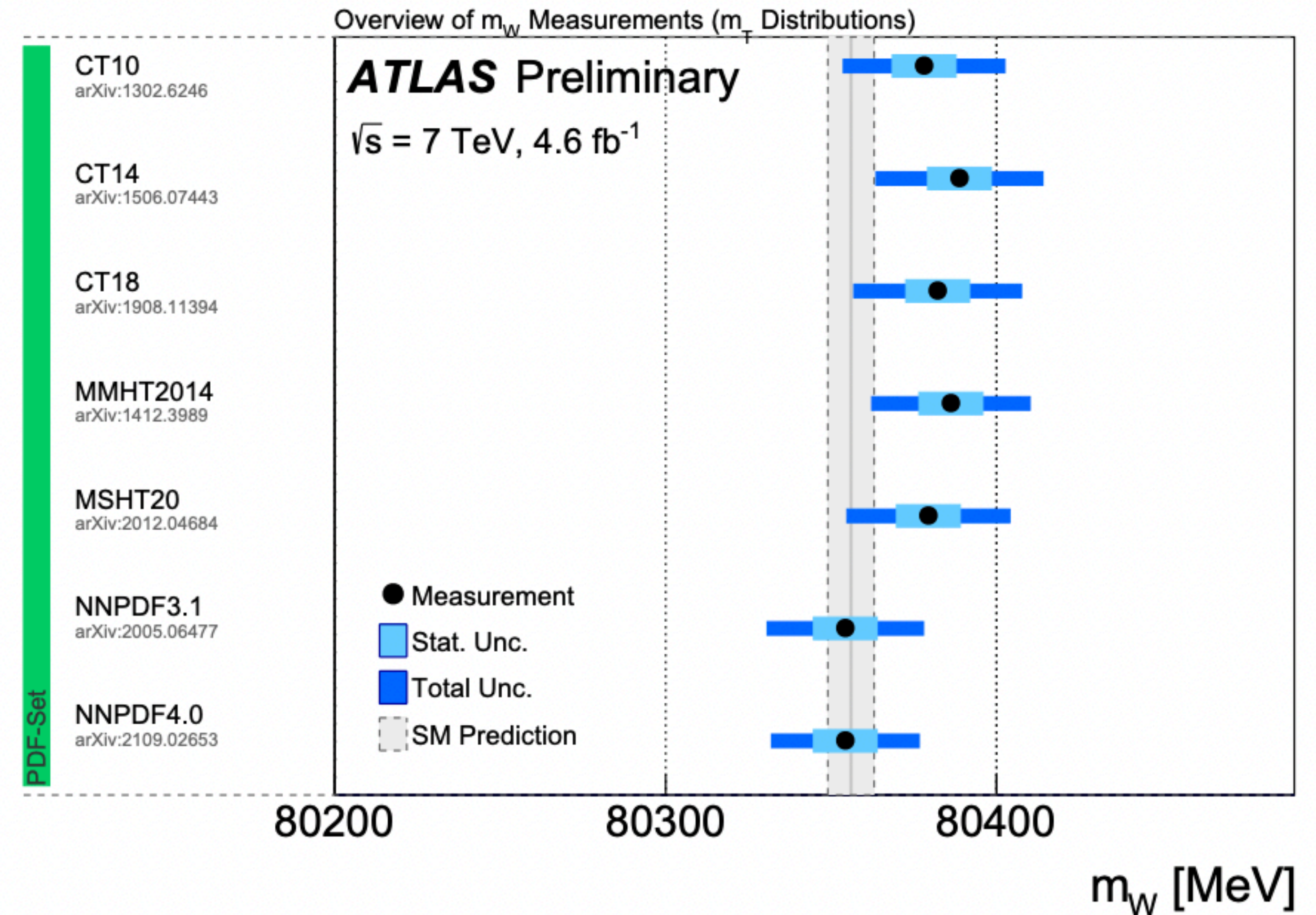
Data set	NNPDF3.1	NNPDF4.0	ABMP16	CT18	MSHT20
ATLAS $W, Z$ 7 TeV ( $\mathcal{L} = 35 \text{ pb}^{-1}$ )	✓	✓	✓	✓	✓
ATLAS $W, Z$ 7 TeV ( $\mathcal{L} = 4.6 \text{ fb}^{-1}$ )	✓	✓	✗	(✓)	✓
ATLAS low-mass DY 7 TeV	✓	✓	✗	(✓)	✗
ATLAS high-mass DY 7 TeV	✓	✓	✗	(✓)	✓
ATLAS $W$ 8 TeV	✗	(✓)	✗	✗	✓
ATLAS DY 2D 8 TeV	✗	✓	✗	✗	✓
ATLAS high-mass DY 2D 8 TeV	✗	✓	✗	(✓)	✓
ATLAS $\sigma_{W,Z}$ 13 TeV	✗	✓	✓	✗	✗
ATLAS $W$ +jet 8 TeV	✗	✓	✗	✗	✓
ATLAS $Z$ $p_T$ 7 TeV	(✓)	✗	✗	(✓)	✗
ATLAS $Z$ $p_T$ 8 TeV	✓	✓	✗	✓	✓
ATLAS $W + c$ 7 TeV	✗	✓	✗	(✓)	✗
ATLAS $\sigma_{tt}^{\text{tot}}$ 7, 8 TeV	✓	✓	✓	✗	✗
ATLAS $\sigma_{tt}^{\text{tot}}$ 7, 8 TeV	✗	✗	✓	✗	✗
ATLAS $\sigma_{tt}^{\text{tot}}$ 13 TeV ( $\mathcal{L} = 3.2 \text{ fb}^{-1}$ )	✓	✗	✓	✗	✗
ATLAS $\sigma_{tt}^{\text{tot}}$ 13 TeV ( $\mathcal{L} = 139 \text{ fb}^{-1}$ )	✗	✓	✗	✗	✗
ATLAS $\sigma_{tt}^{\text{tot}}$ and $Z$ ratios	✗	✗	✗	✗	(✓)
ATLAS $t\bar{t}$ lepton+jets 8 TeV	✓	✓	✗	✓	✓
ATLAS $t\bar{t}$ dilepton 8 TeV	✗	✓	✗	✗	✓
ATLAS single-inclusive jets 7 TeV, R=0.6	✓	(✓)	✗	✓	✓
ATLAS single-inclusive jets 8 TeV, R=0.6	✗	✓	✗	✗	✗
ATLAS dijets 7 TeV, R=0.6	✗	✓	✗	✗	✗
ATLAS direct photon production 8 TeV	✗	(✓)	✗	✗	✗
ATLAS direct photon production 13 TeV	✗	✓	✗	✗	✗
ATLAS single top $R_t$ 7, 8, 13 TeV	✗	✓	✓	✗	✗
ATLAS single top diff. 7 TeV	✗	✓	✗	✗	✗
ATLAS single top diff. 8 TeV	✗	✓	✗	✗	✗

- NNPDF4.0 based on larger set of data from LHC Run I and Run II
- Open-source public code [\[arXiv:2109.02671\]](https://arxiv.org/abs/2109.02671)
- O(4500) data points and first time inclusion of di-jets preferred over inclusive jets based on perturbative stability [\[Khalek et al, 2005.11327\]](https://arxiv.org/abs/2005.11327)

# PHENOMENOLOGICAL IMPLICATIONS

ATLAS-CONF-2023-004

PDF-Set	$p_T^\ell$ [MeV]	$m_T$ [MeV]	combined [MeV]
CT10	$80355.6^{+15.8}_{-15.7}$	$80378.1^{+24.4}_{-24.8}$	$80355.8^{+15.7}_{-15.7}$
CT14	$80358.0^{+16.3}_{-16.3}$	$80388.8^{+25.2}_{-25.5}$	$80358.4^{+16.3}_{-16.3}$
CT18	$80360.1^{+16.3}_{-16.3}$	$80382.2^{+25.3}_{-25.3}$	$80360.4^{+16.3}_{-16.3}$
MMHT2014	$80360.3^{+15.9}_{-15.9}$	$80386.2^{+23.9}_{-24.4}$	$80361.0^{+15.9}_{-15.9}$
MSHT20	$80358.9^{+13.0}_{-16.3}$	$80379.4^{+24.6}_{-25.1}$	$80356.3^{+14.6}_{-14.6}$
NNPDF3.1	$80344.7^{+15.6}_{-15.5}$	$80354.3^{+23.6}_{-23.7}$	$80345.0^{+15.5}_{-15.5}$
NNPDF4.0	$80342.2^{+15.3}_{-15.3}$	$80354.3^{+22.3}_{-22.4}$	$80342.9^{+15.3}_{-15.3}$



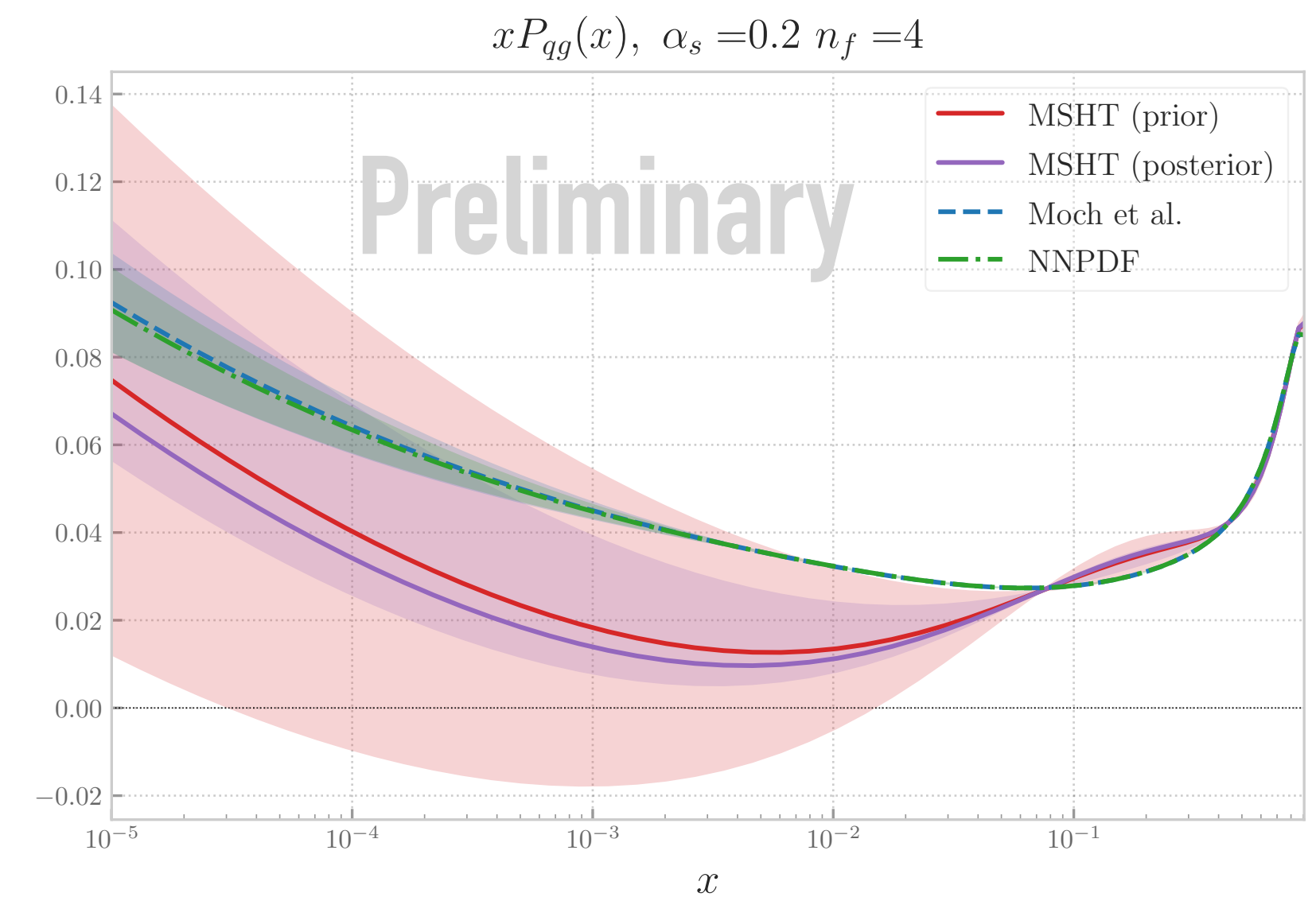
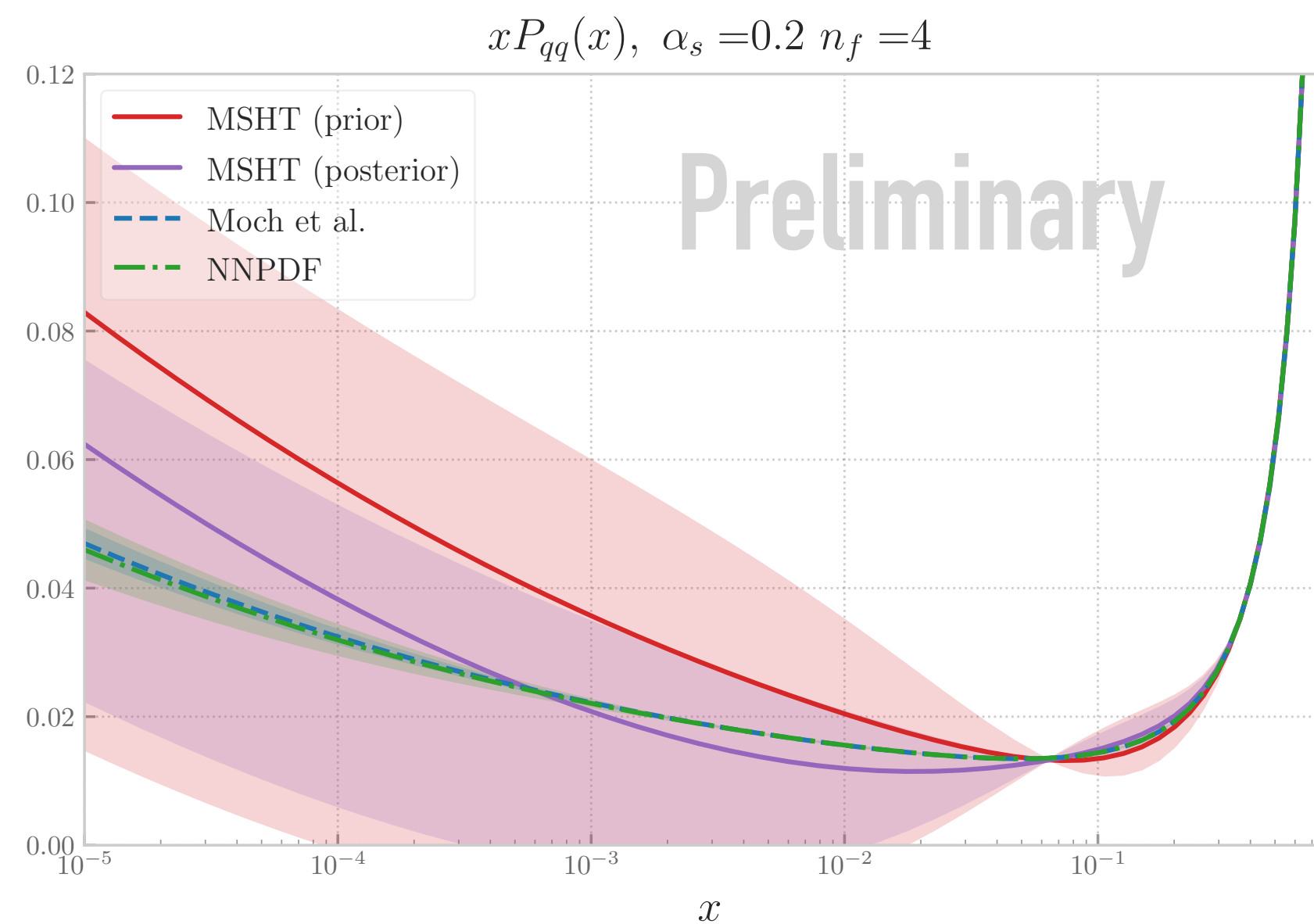
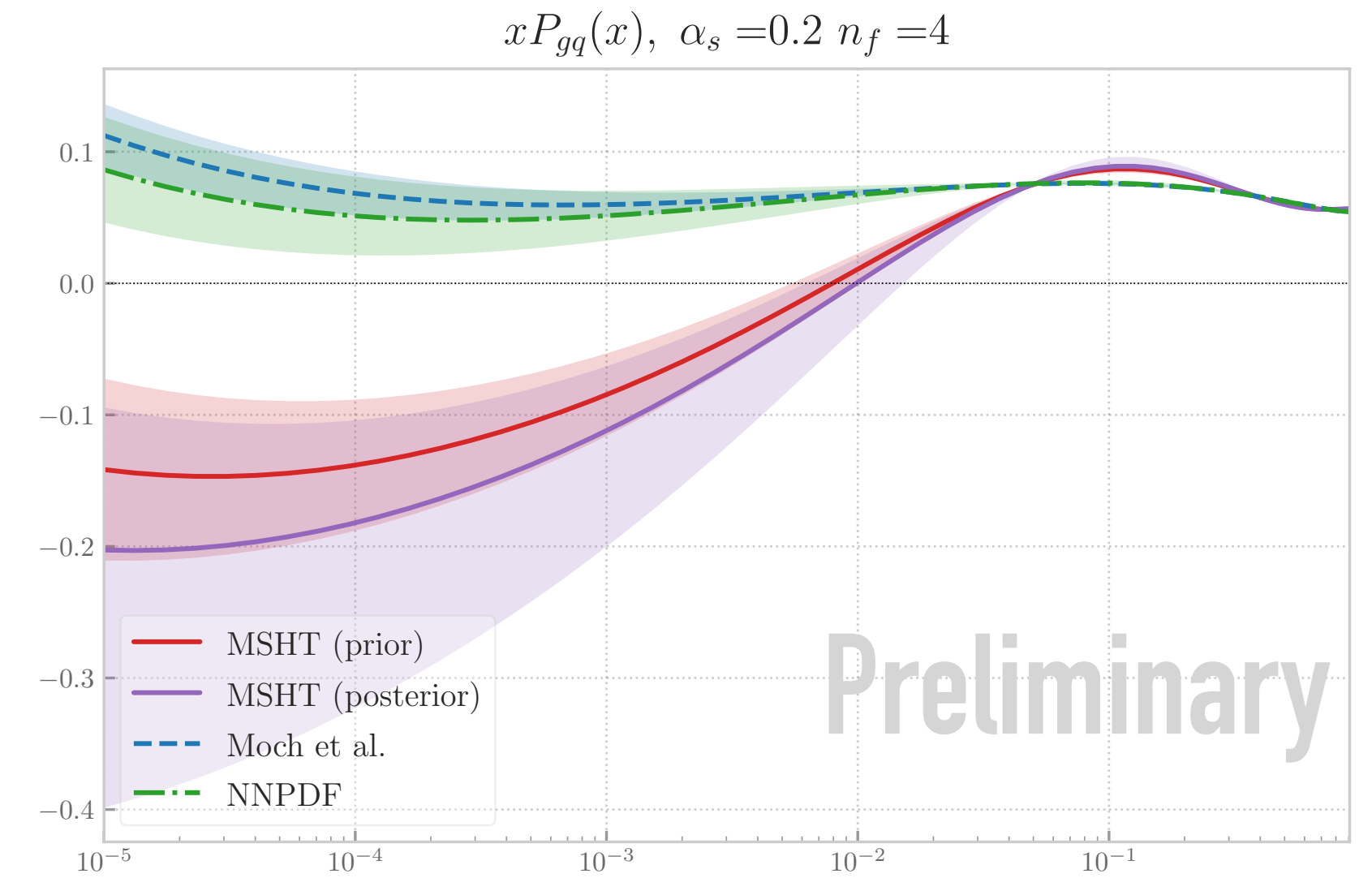
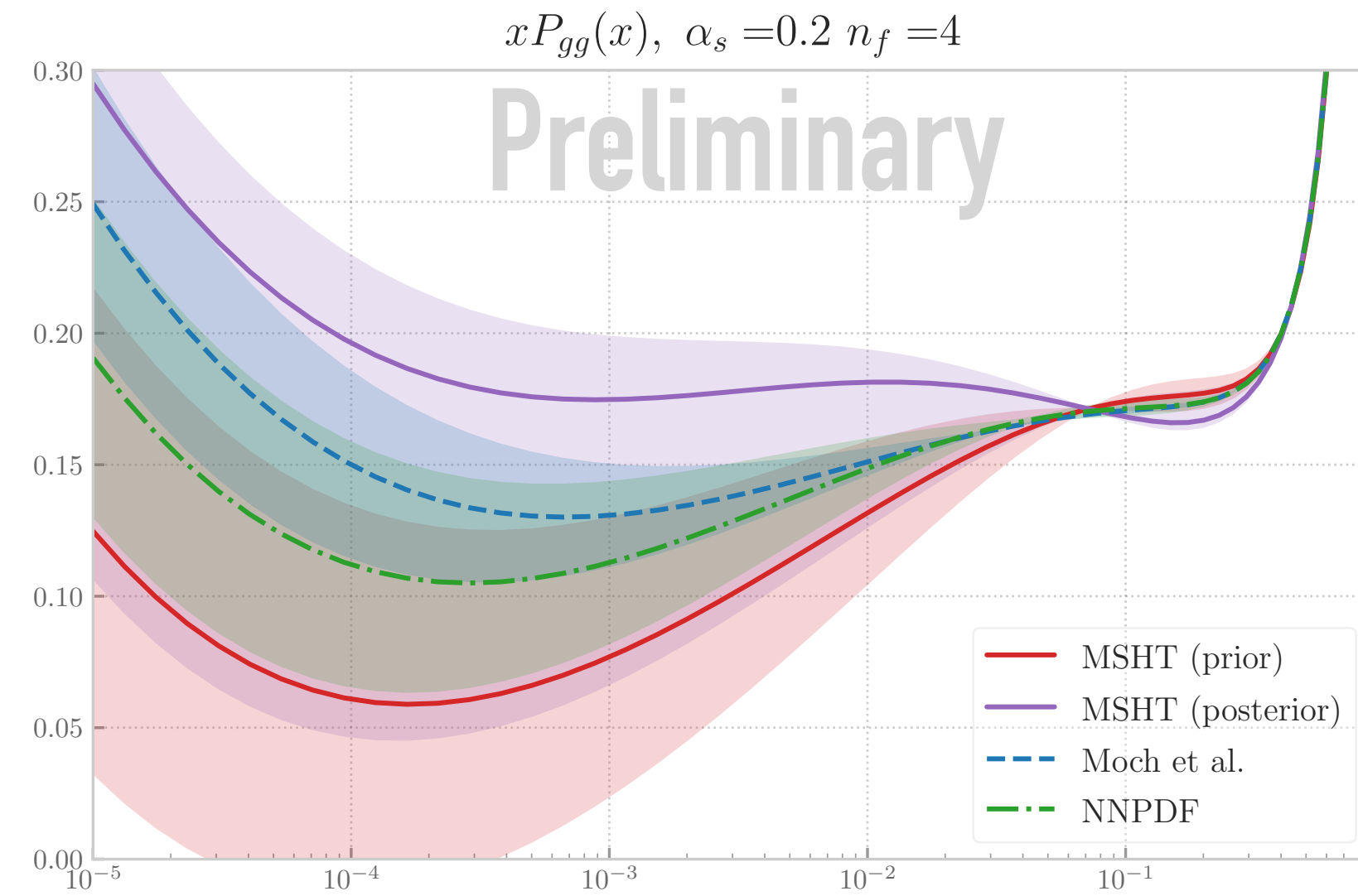
- Different PDF uncertainties yield slightly different total uncertainties.
- Differences of  $\sim 18 \text{ MeV}$  between central values obtained with NNPDF4.0 and CT18 comparable (larger) than measurement uncertainty
- Important to confirm results, understand reason behind the differences & stress-test methodologies



# #1 - APPROXIMATE N3LO PDFS

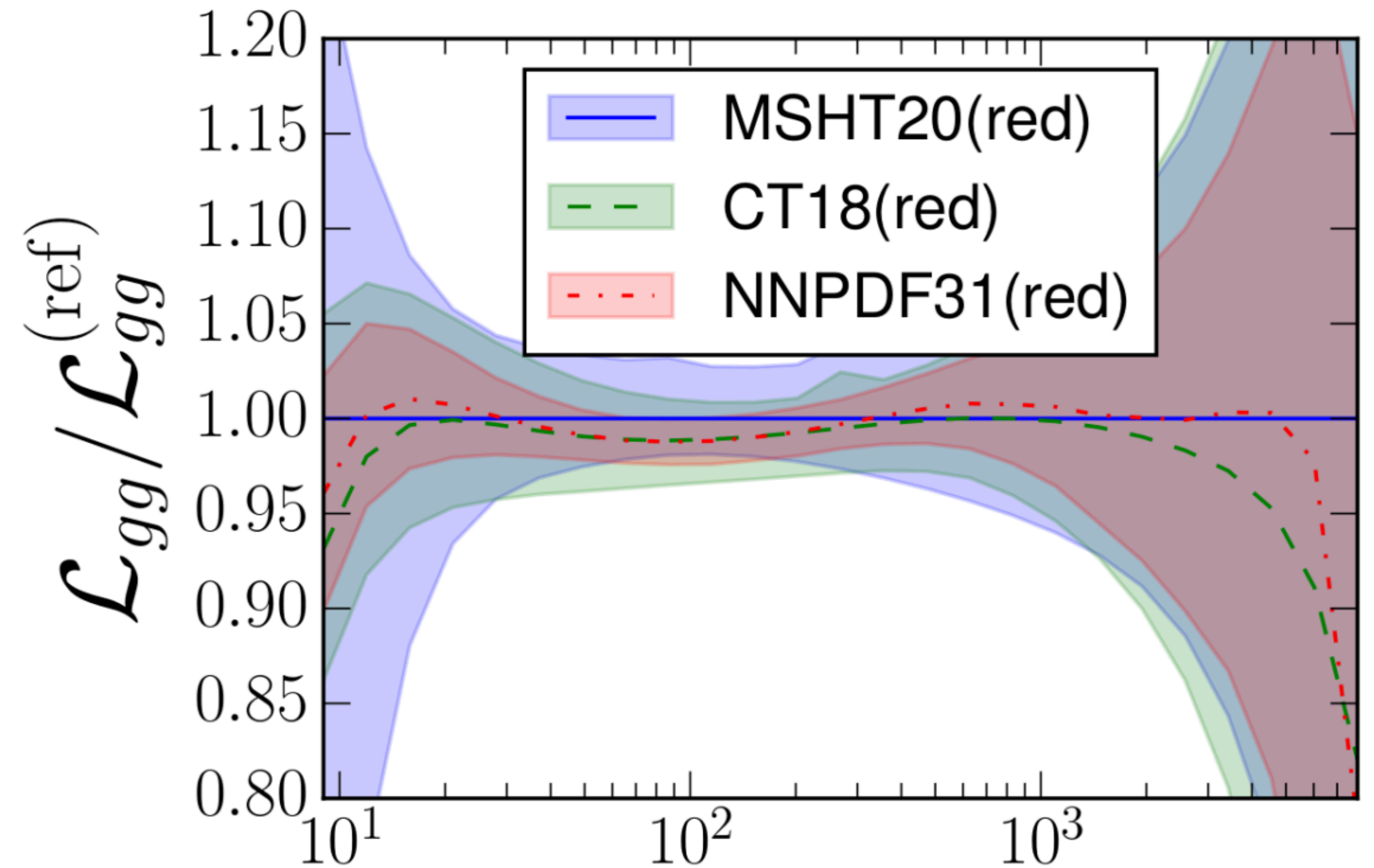
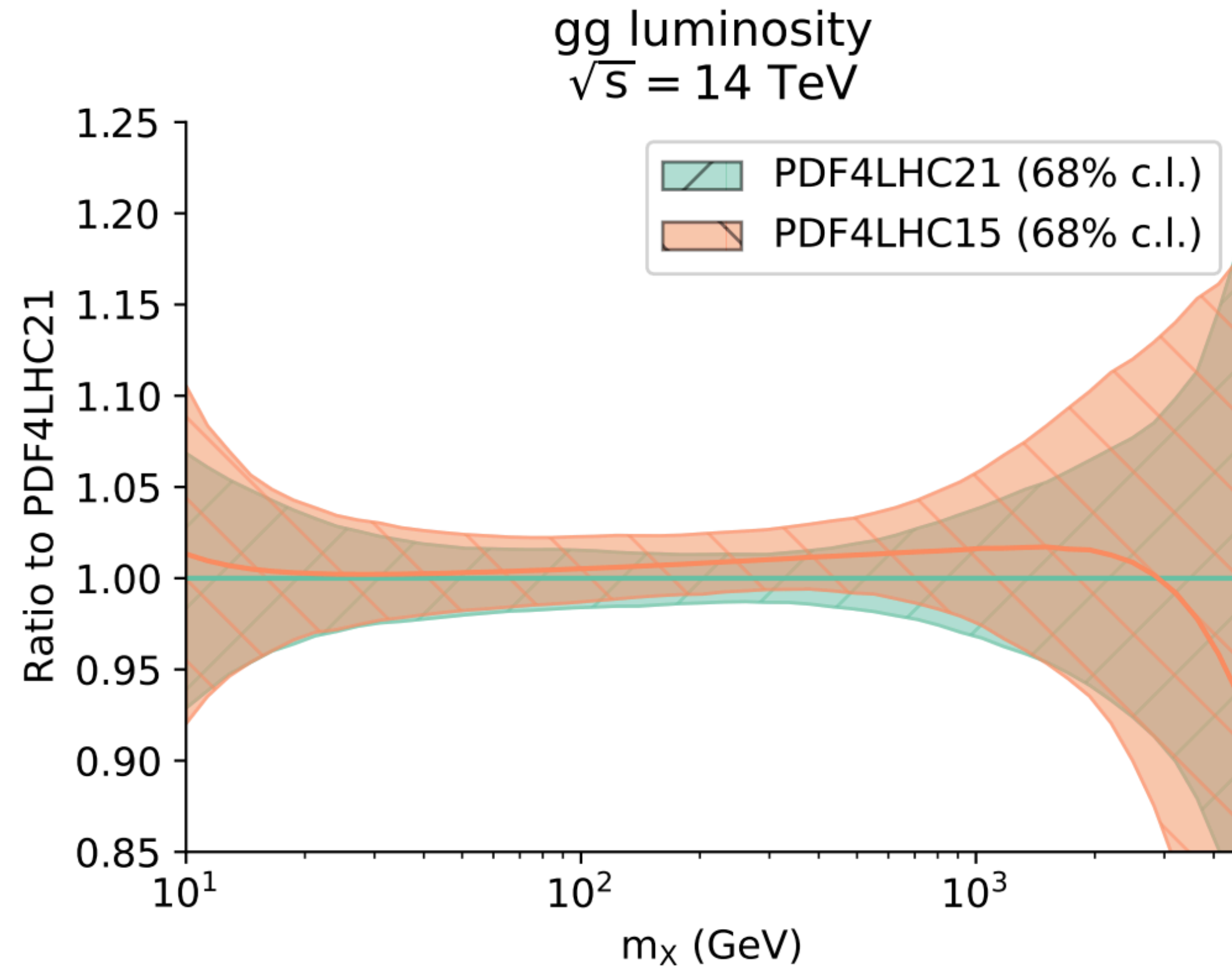
**MSHT prior** = pre-fit variation  
**MSHT posterior** = MSHTaN3LO  
**NNPDF** = NNPDF40aN3LO  
**Moch et al** = theory papers

- Different known ingredients included and different ways of parametrising sub-leading contributions produces different candidates (IHOU)
- Benchmarking exercise started before Summer in Les Houches to check the impact of the aN3LO splitting functions
- Write-up in preparation for Les Houches proceedings



# THE PDF4LHC21 COMBINATION

PDF4LHC WG, arXiv: 2203.05506



- In recent updates from global PDF fitting collaborations (NNPDF3.1, CT18, MSHT20, ABMP16) effect of LHC data driving PDF uncertainties down → PDF4LHC21 combination of NNPDF31', CT18' and MSHT20' smaller uncertainties than PDF4LHC15
- The increased precision of the data and their strong correlation demands methodological improvements (e. g. more flexible parametrisation that shifts PDFs outside nominal error band - MSHT) and shows some tension among (CT18 vs CT18Z)
- Benchmark among NNPDF3.1, MSHT20 and CT18: overall agreement, which improves once common dataset is used, differences in uncertainties with  $\Delta\text{CT18} \approx \Delta\text{MSHT20} \approx \Delta\text{NNPDF31}$
- NNPDF4.0 [[arXiv:2109.02653](https://arxiv.org/abs/2109.02653)], MSHT20aN3LO [[arXiv:2207.04739](https://arxiv.org/abs/2207.04739)] state-of-the-art post PDF4LHC21 sets

# #1 – APPROXIMATE N<sup>3</sup>LO PDFS

- ➔ Several ingredients required to perform N<sup>3</sup>LO PDF fits, many available some missing [[S. Moch's talk](#)]
- ➔ Only publicly available aN<sup>3</sup>LO PDF determination is from the MSHT collaboration [[McGowan et al, arXiv:2207.04739](#)]
- ➔ NNPDF has presented preliminary aN<sup>3</sup>LO results and paper is to appear soon

## MSHT

- All available theory input at the time of publication included (impact of new ingredients being explored)
- Incomplete N<sup>3</sup>LO terms added as variation in the prior and estimated by fitting nuisance parameters to the data (hence **posterior determined by fitting data**)
- No MHOU associated with NNLO contributions

## NNPDF

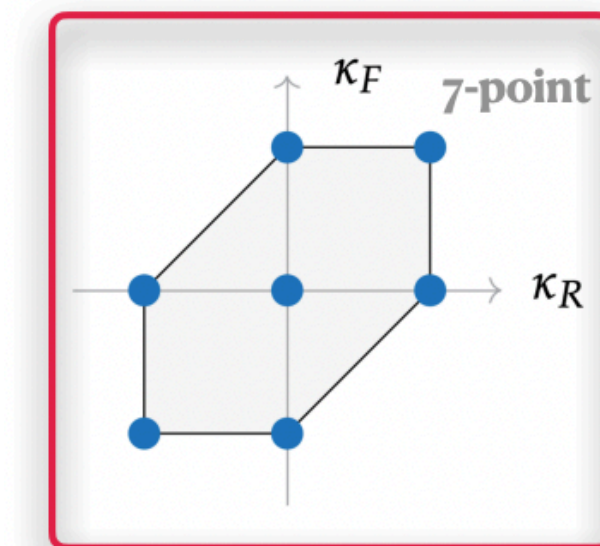
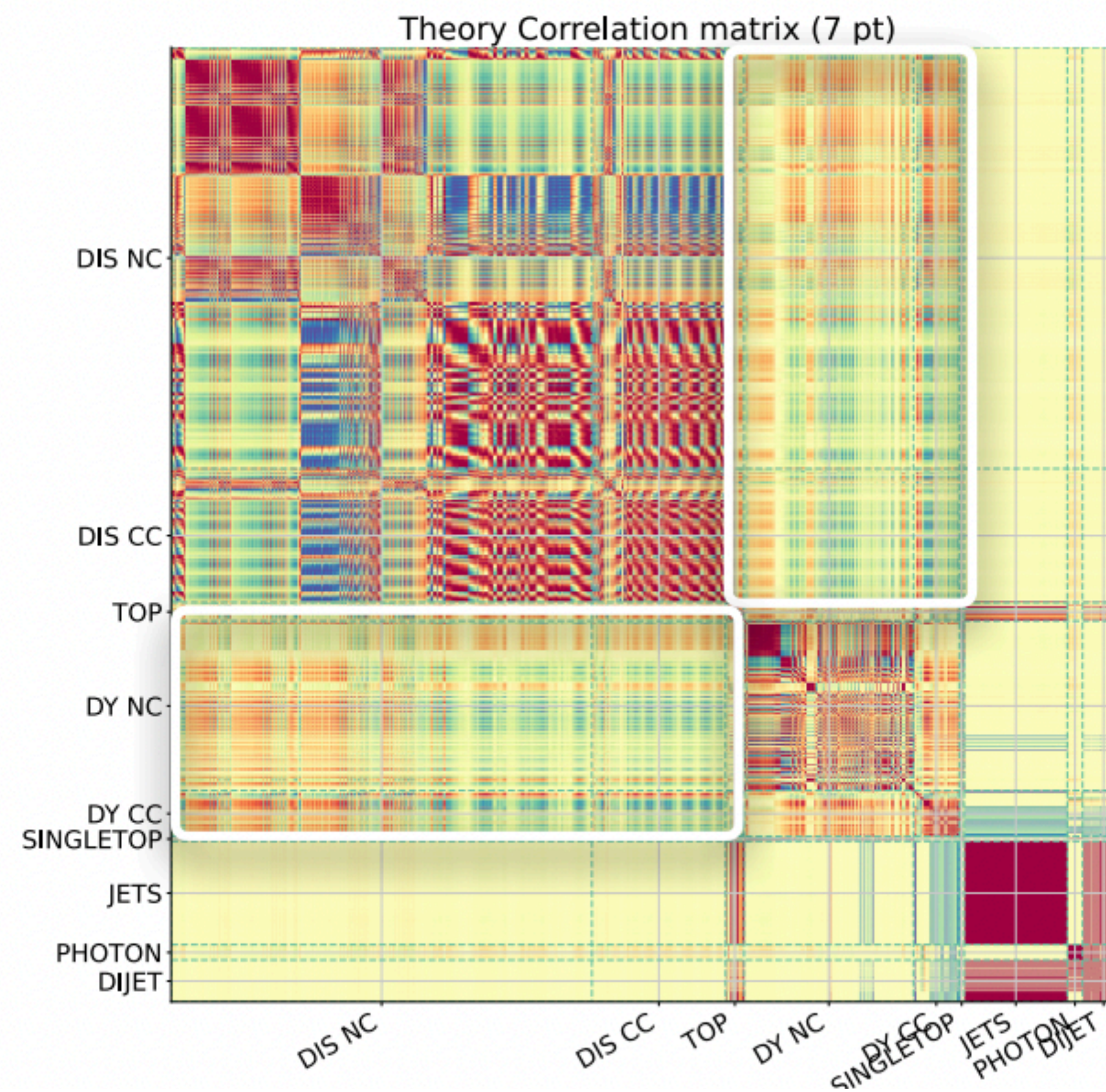
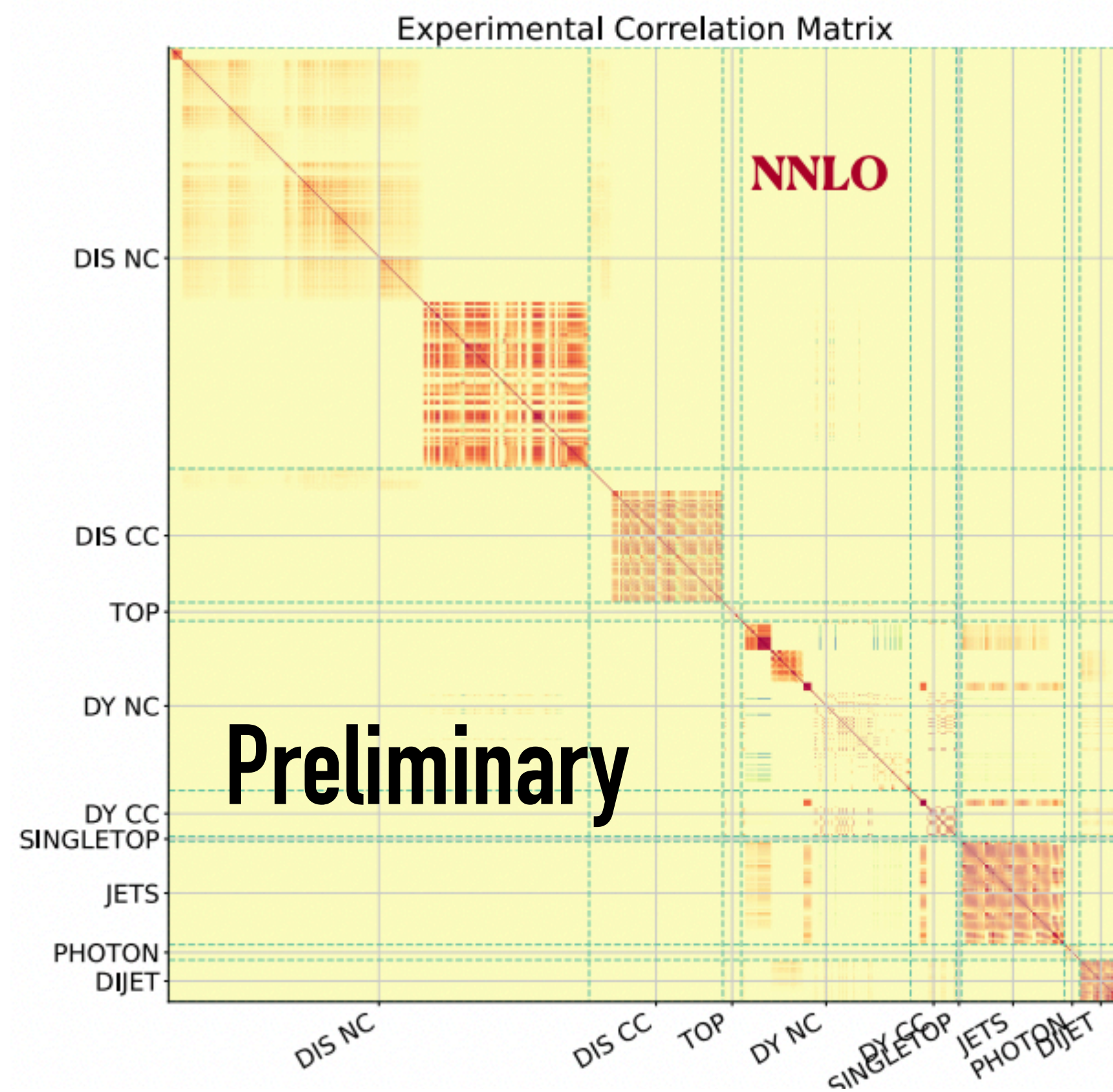
- More theory inputs published in between and included, in particular 6(1) extra momentum for  $P_{qg}$ ,  $P_{qq}$  ( $P_{gq}$ ,  $P_{gg}$ ), some terms in the large  $n_f$  limit, sub-leading small- $x$  and large- $x$  terms.
- **Only theory inputs** and their variations added to an additional theory covariance matrix associated with incomplete missing higher orders (IHOU)
- MHOU associated with NNLO included via theory covariance matrix

# #2 - MISSING HIGHER ORDER UNCERTAINTIES IN PDF FITS

- ➔ PDF uncertainties only include experimental component: how to include uncertainty associated to truncation of perturbative expansion in theoretical predictions used in PDF fits?
- ➔ One option is to construct a theory covariance matrix from scale-varied cross sections and combine it with the experimental covariance matrix

$$\chi^2 = \sum_{m,n=1}^N (d_m - t_m)(\text{cov}_{\text{exp}} + \text{cov}_{\text{th}})_{mn}^{-1} (d_n - t_n)$$

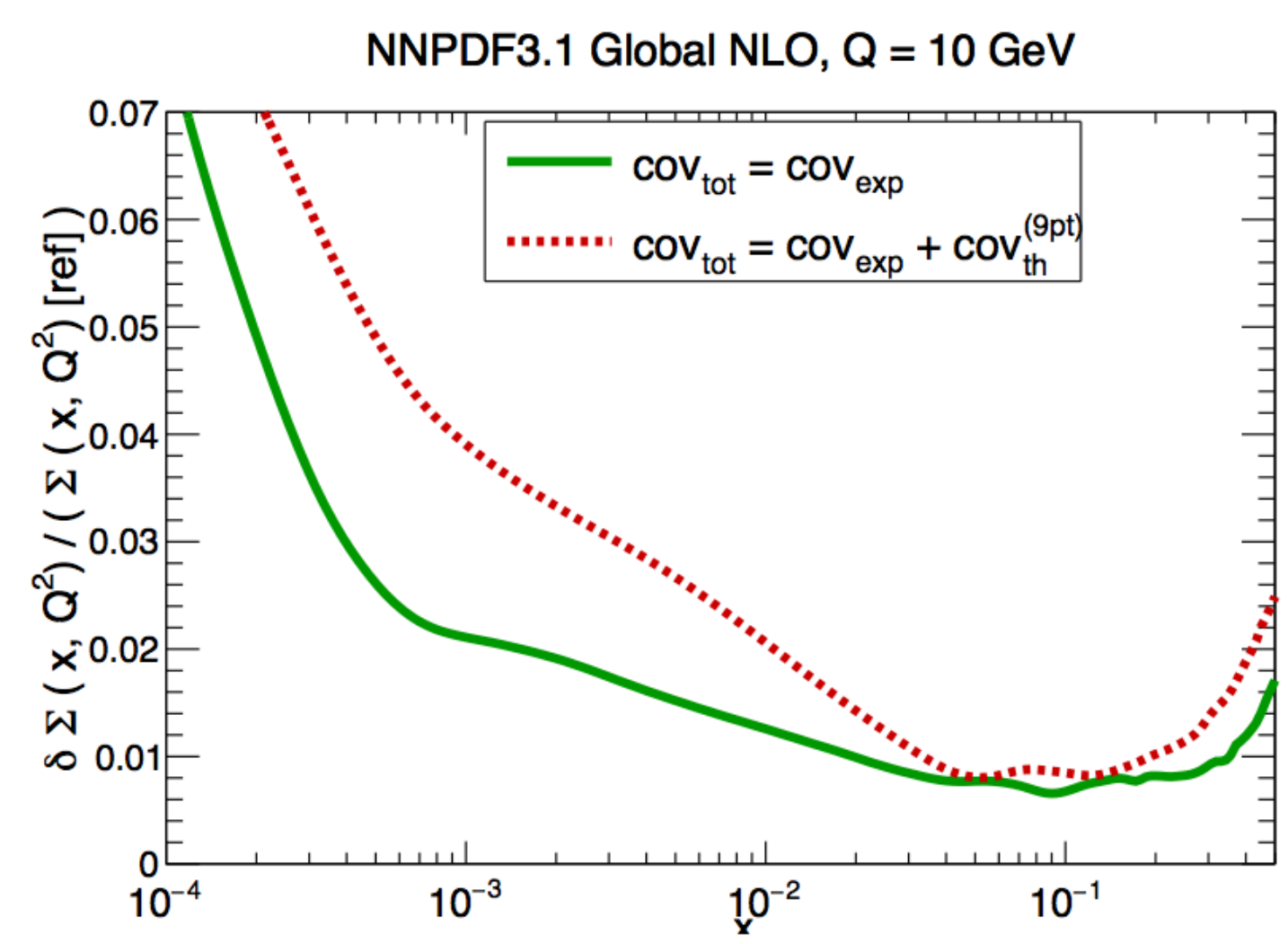
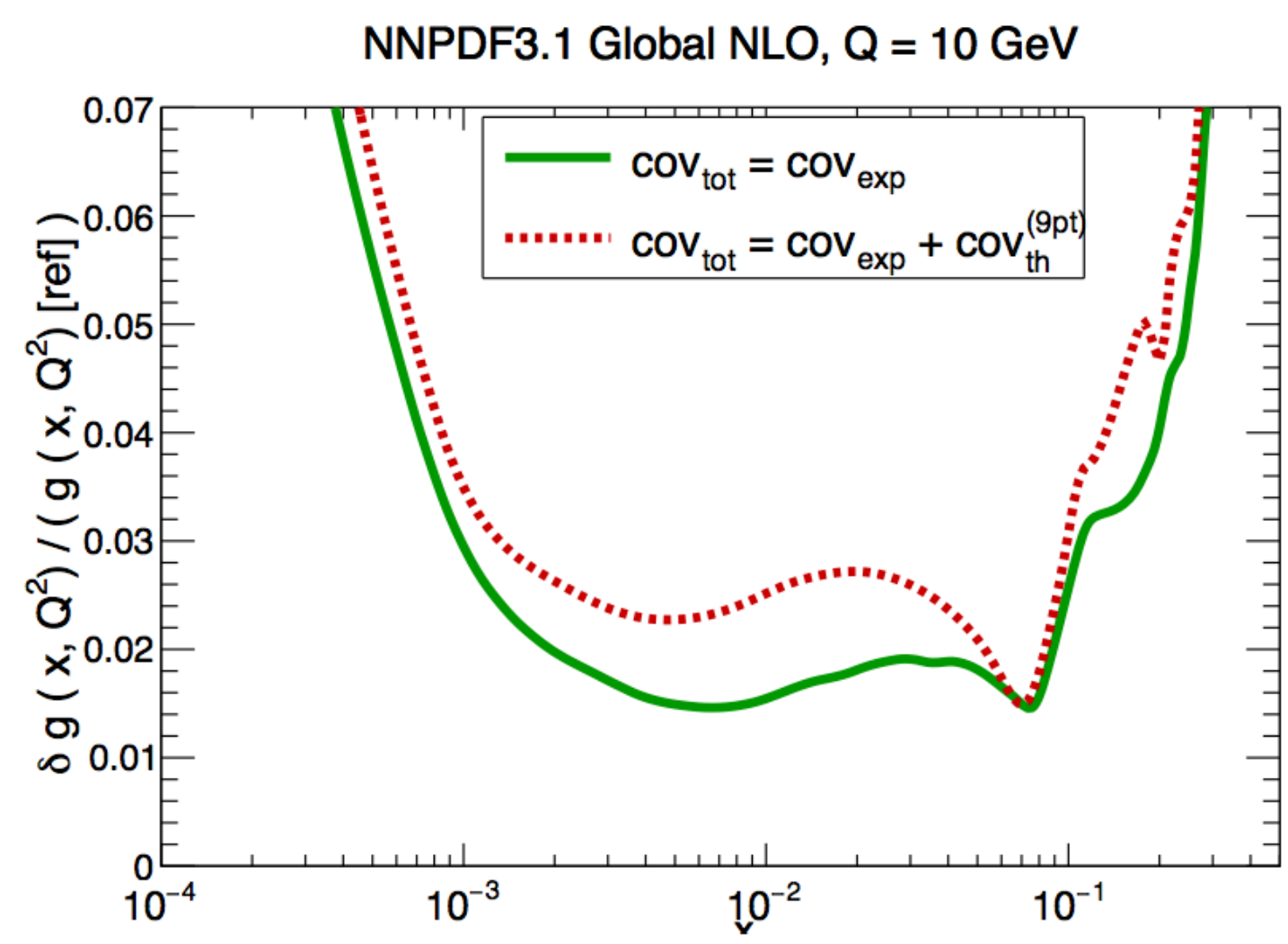
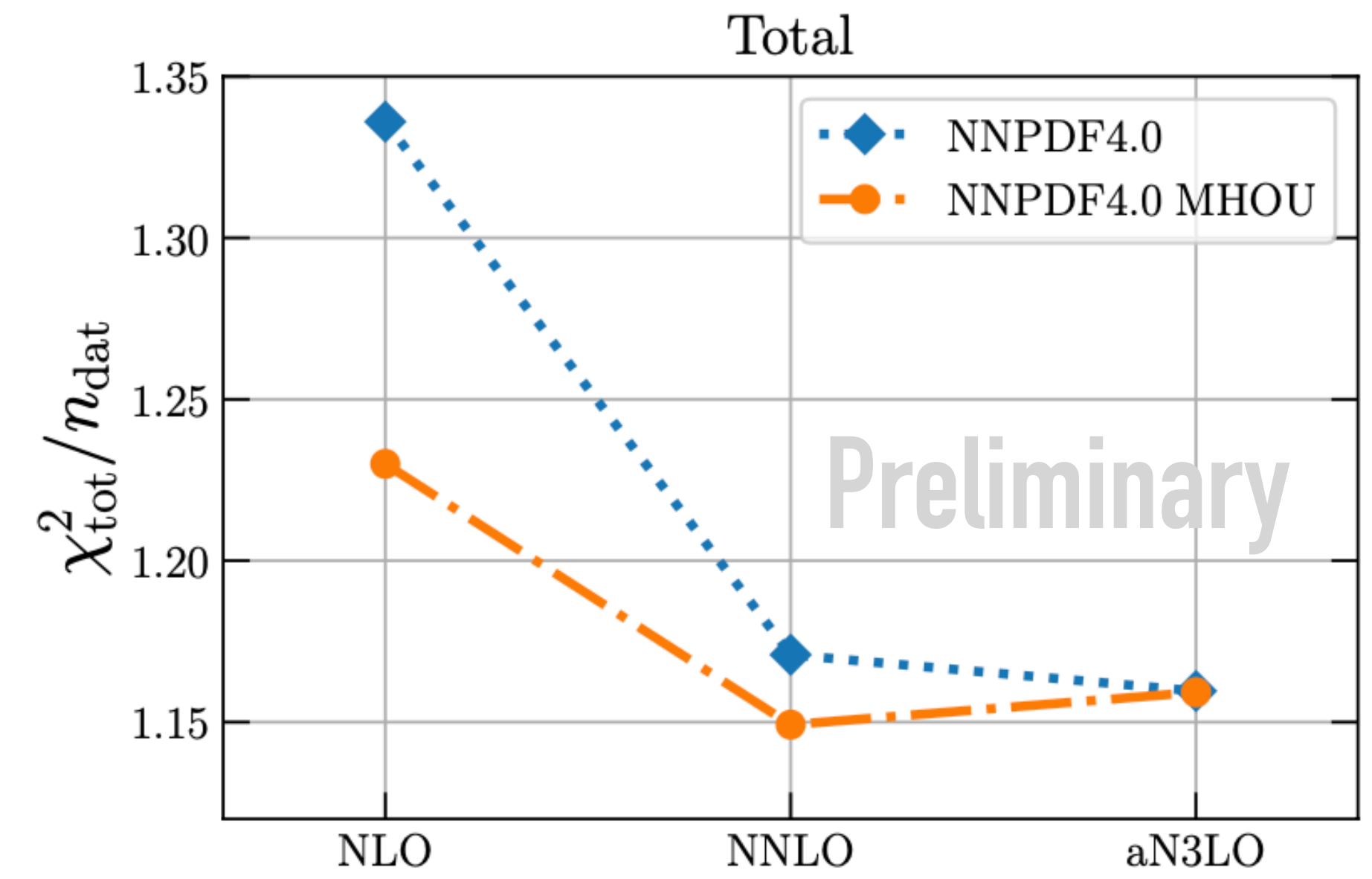
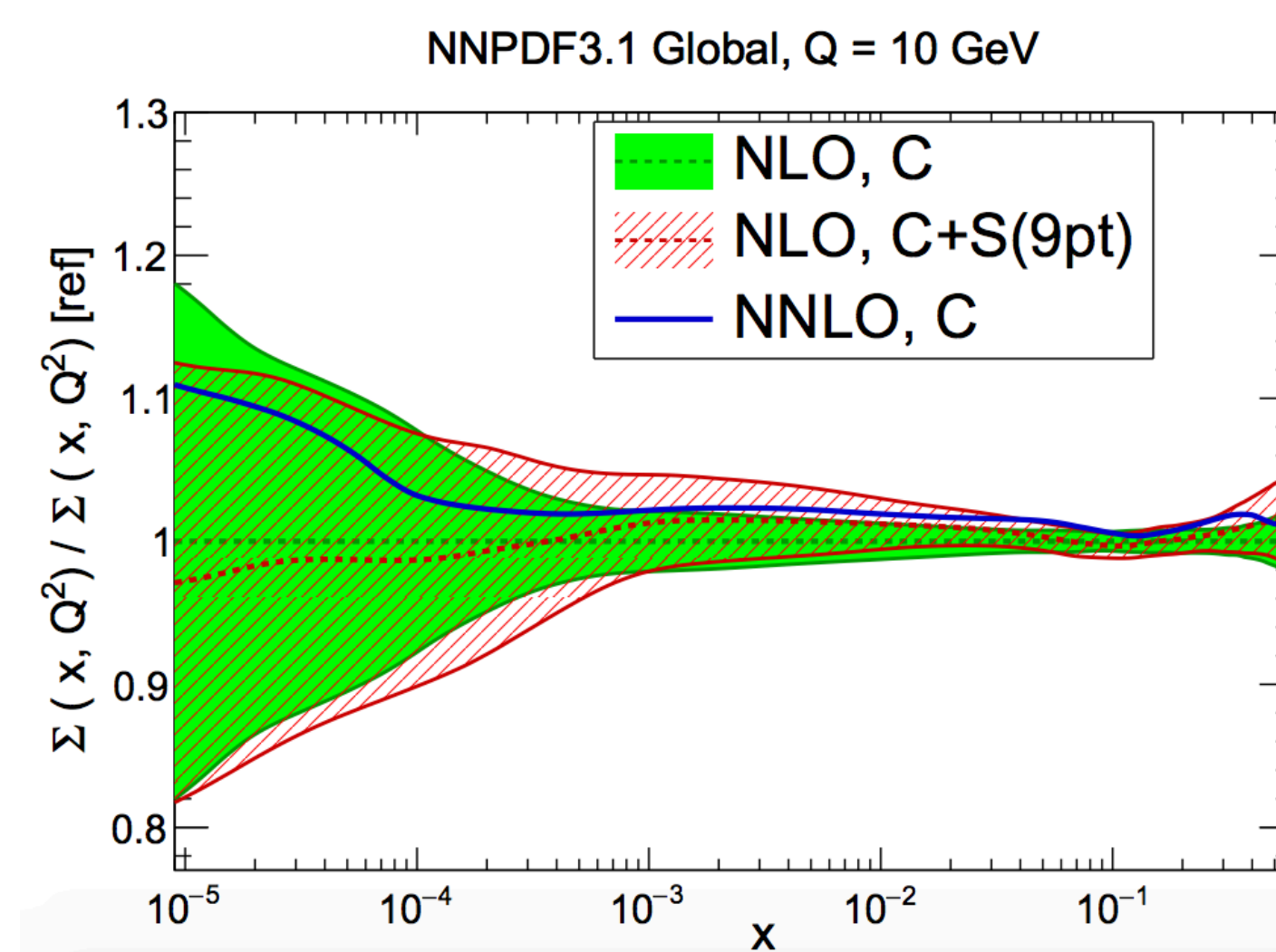
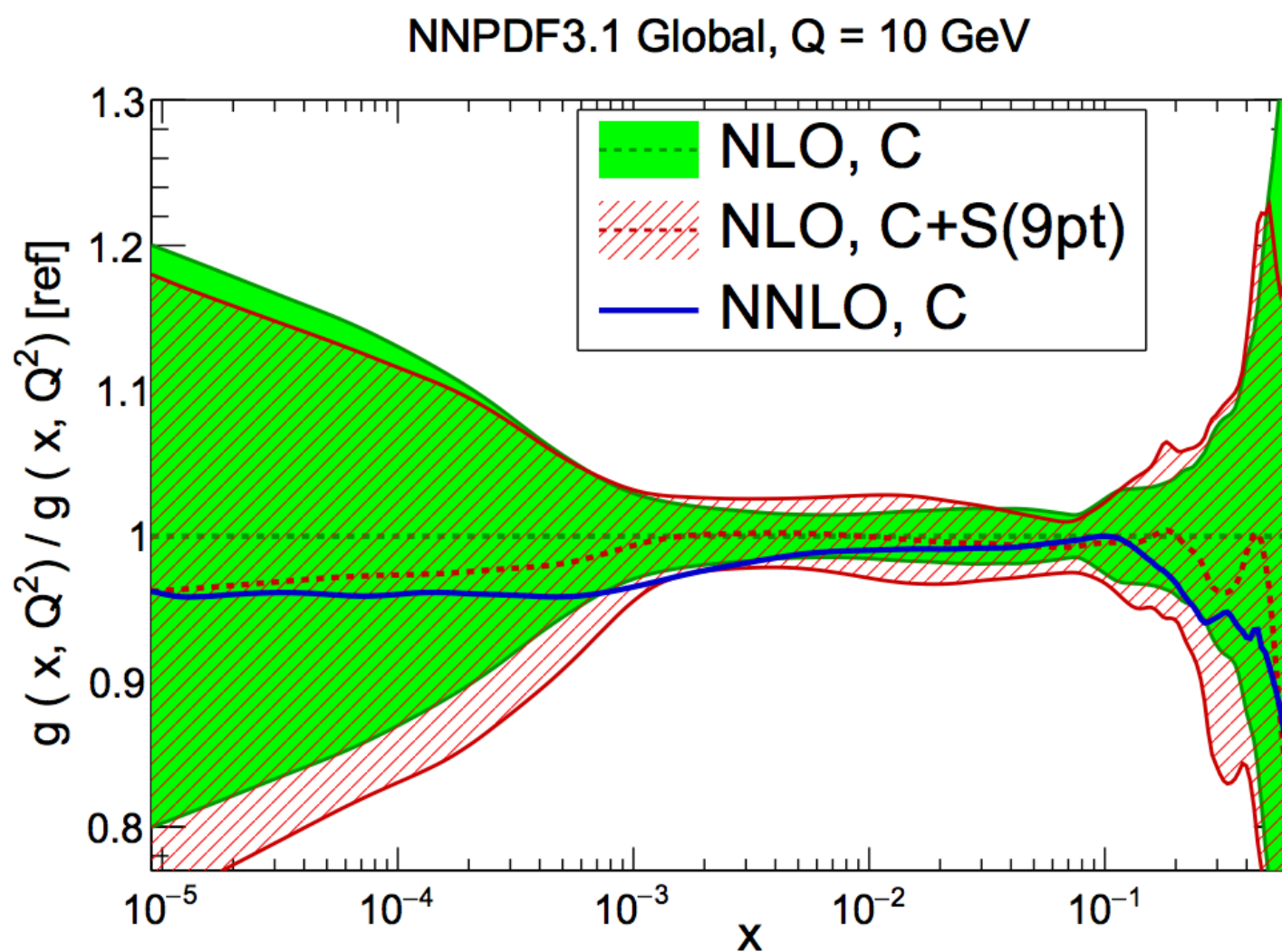
- Assumptions: experimental and theoretical errors independent and Gaussian



7-pts variation

[NNPDF: 1906.10698]  
NNLO paper in preparation

# #2 – MISSING HIGHER ORDER UNCERTAINTIES IN PDF FITS



➔ NNPDF40NNLO with MHOU and NNPDF40aN3LO will be published soon

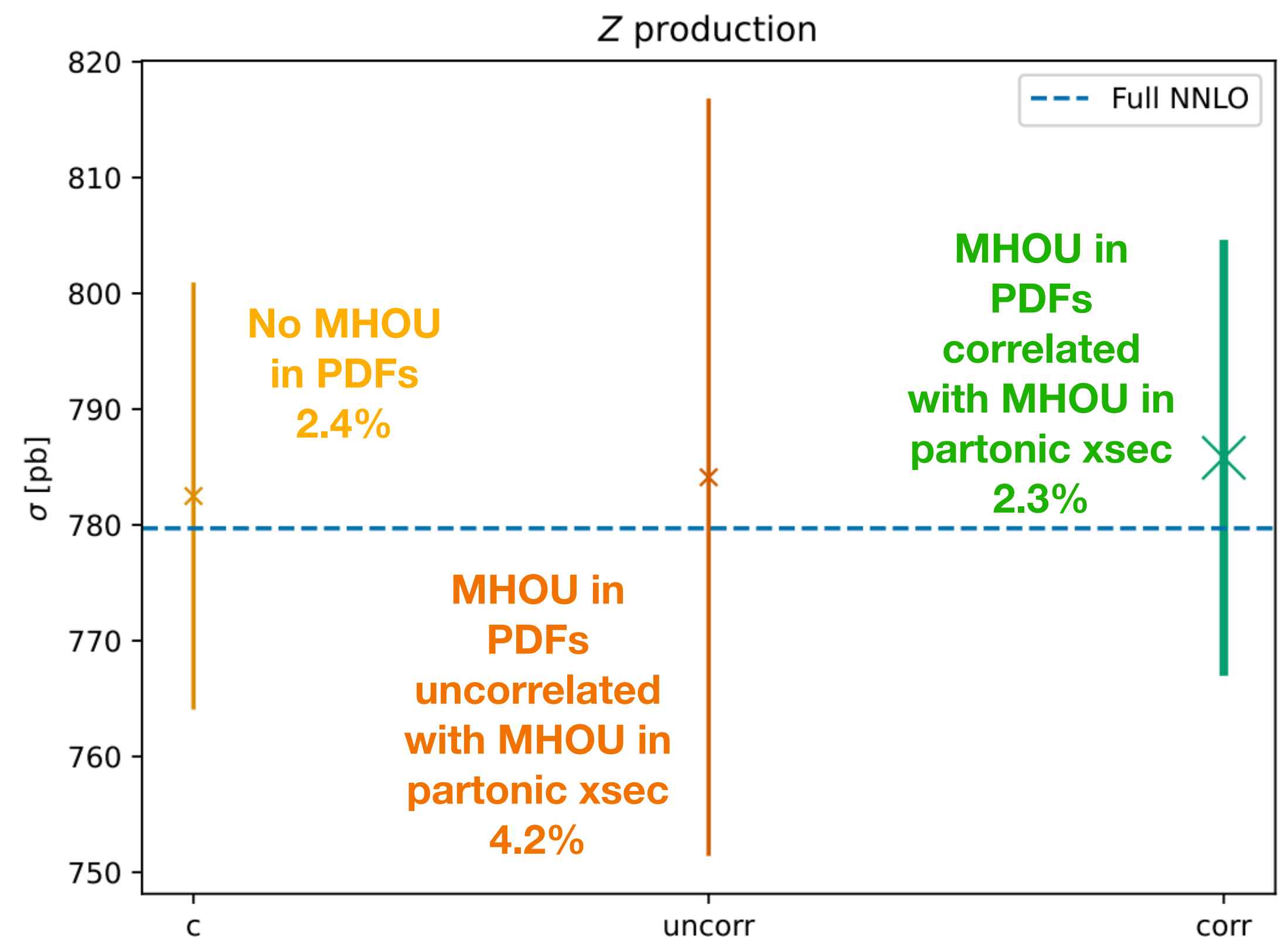
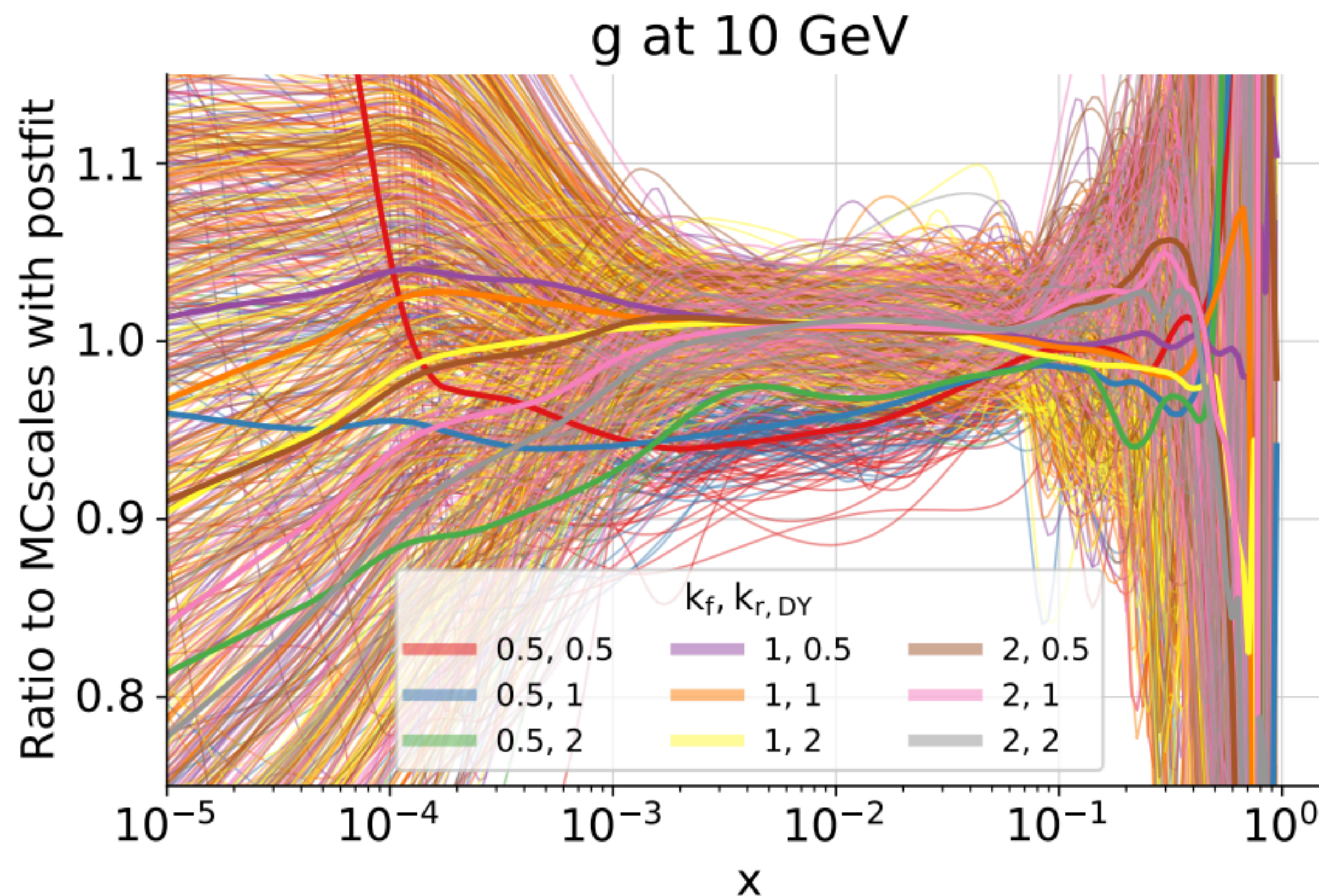
$$Cov_{ij} = Cov_{ij,EXP} + Cov_{ij,MHOU} + Cov_{ij,IHOU}$$

➔ Missing correlation between scale variation in PDF fits and hard cross sections [Harland-Lang, Thorne Eur.Phys.J.C 79 (2019) 3, 225] [Ball, Pearson arXiv:2105.05114]

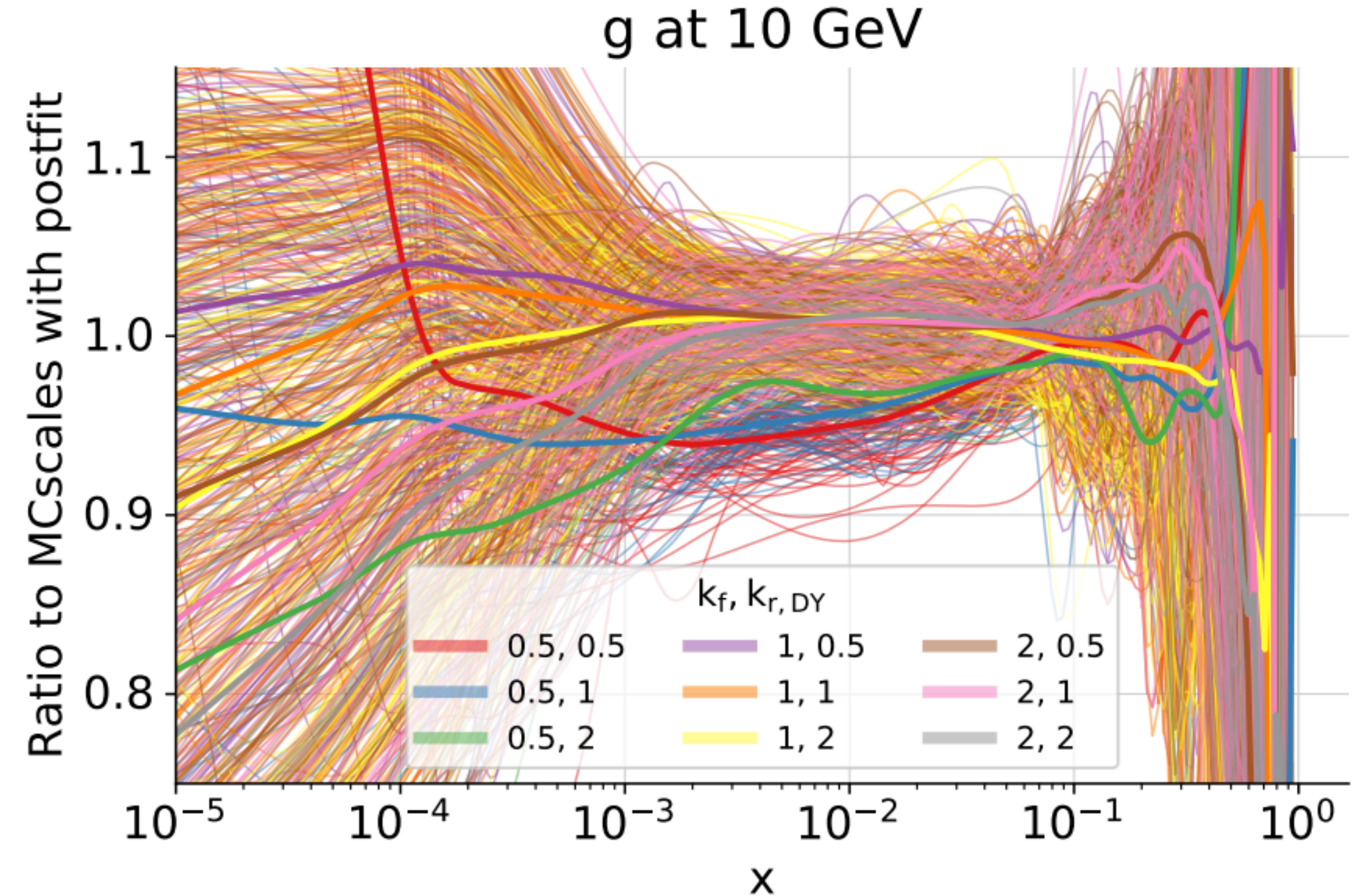
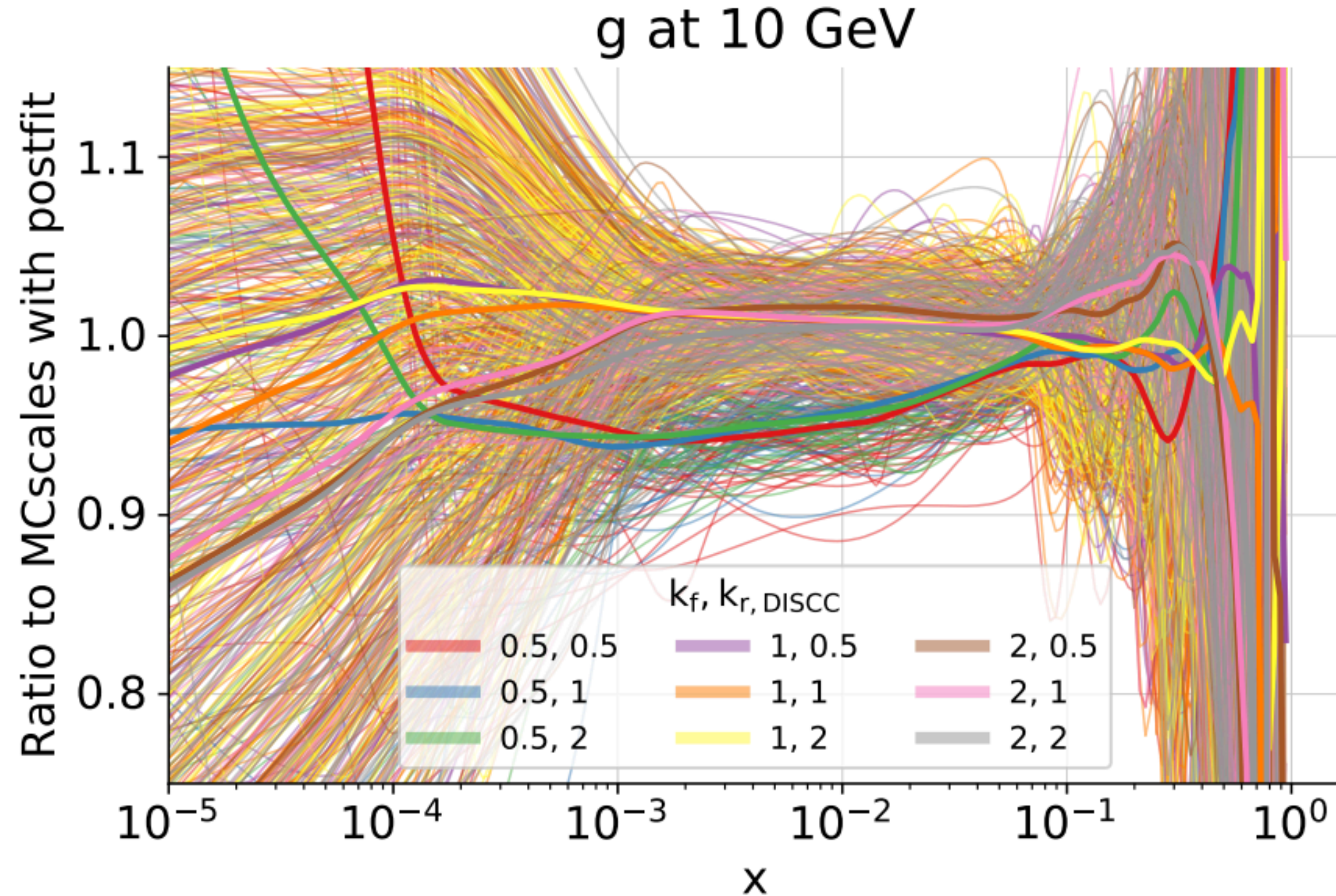
# #2 - MISSING HIGHER ORDER UNCERTAINTIES IN PDF FITS

- ➔ Alternatively, MCscales [Kassabov et al, 2207.07616]: renormalisation and factorisation scales treated as free parameters of the fixed-order theory, that induce an uncertainty on the theory predictions included in a PDF fit & need to be propagated.
- ➔ Joint sampling of experimental uncertainty (propagated to PDF uncertainty by MC sampling) by specifying a suitable prior probability distribution of all possible scale choices & a-posteriori criterion based on agreement with the data.
- ➔ Can compute full PDF+SCALE uncertainty in cross sections by matching the scales in the hard cross section computation with the scales in the MCscale PDF set: correlation fully taken into account

Kassabov et al: 2207.07616]



# THE MCSCALES APPROACH



- ✓ Main idea of MCscales: the renormalisation and factorisation scales are free parameters of the fixed-order theory, that induce an uncertainty on the theory predictions included in a PDF fit & need to be propagated
- ✓ Joint sampling of experimental uncertainty (propagated to PDF uncertainty by MC sampling) by specifying a suitable prior probability distribution of all possible scale choices & a-posteriori criterion based on agreement with the data.

$$P(k_f = \xi_f, k_{r_1} = \xi_1, \dots, k_{r_{N_p}} = \xi_{N_p}) = P(\omega)$$

$$\text{with } \omega \in \Omega = \{(\xi_f, \xi_1, \dots, \xi_{N_p}) \mid \forall \xi_f, \xi_1, \dots, \xi_{N_p} \in \Xi\}$$

$3^{1+N_p}$  elements, with  $N_p = 5$ ,  $p = \text{DIS NC, DIS CC, DY, JET, TOP}$

Choose prior = choose  $P(\omega)$

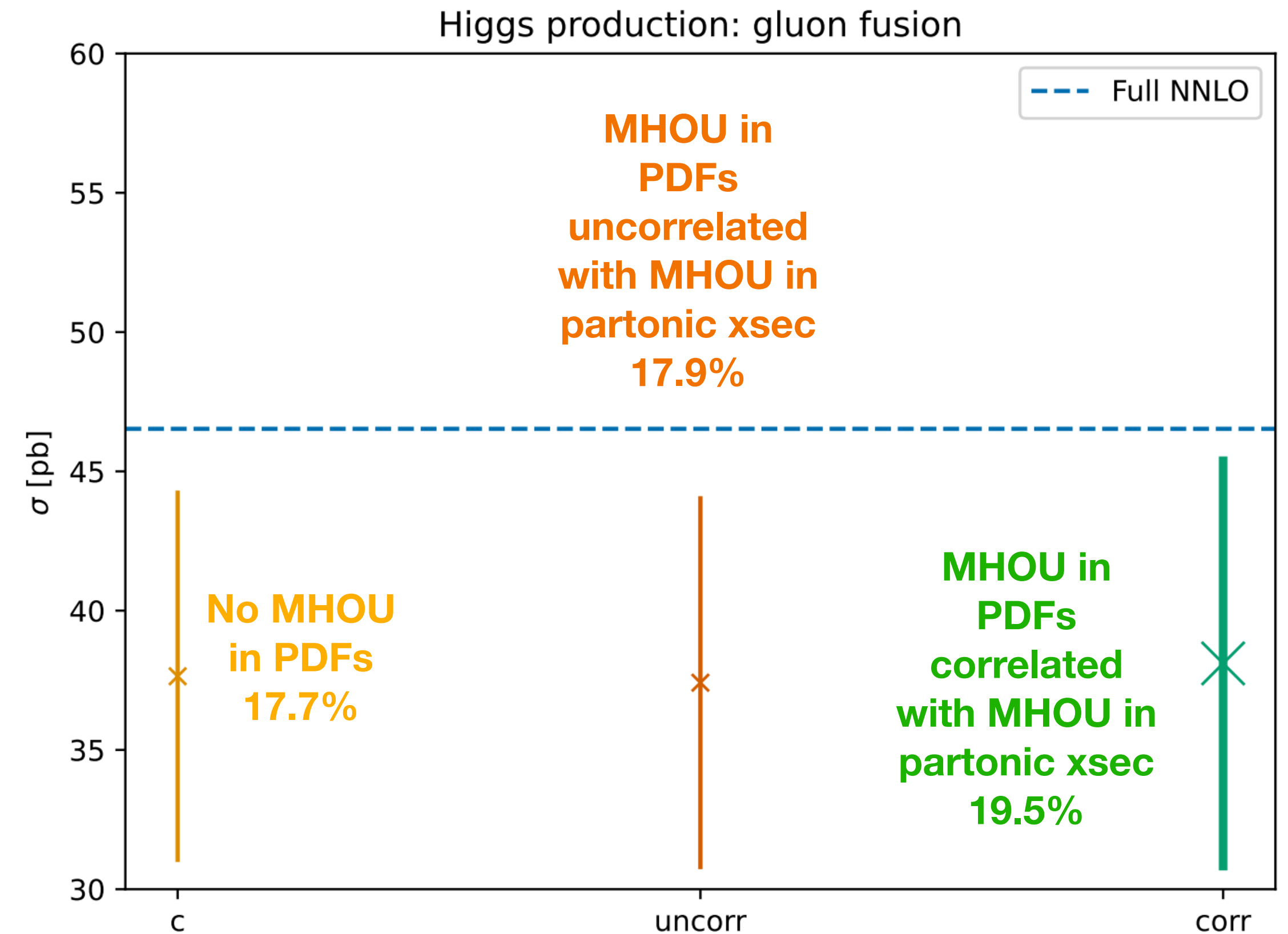
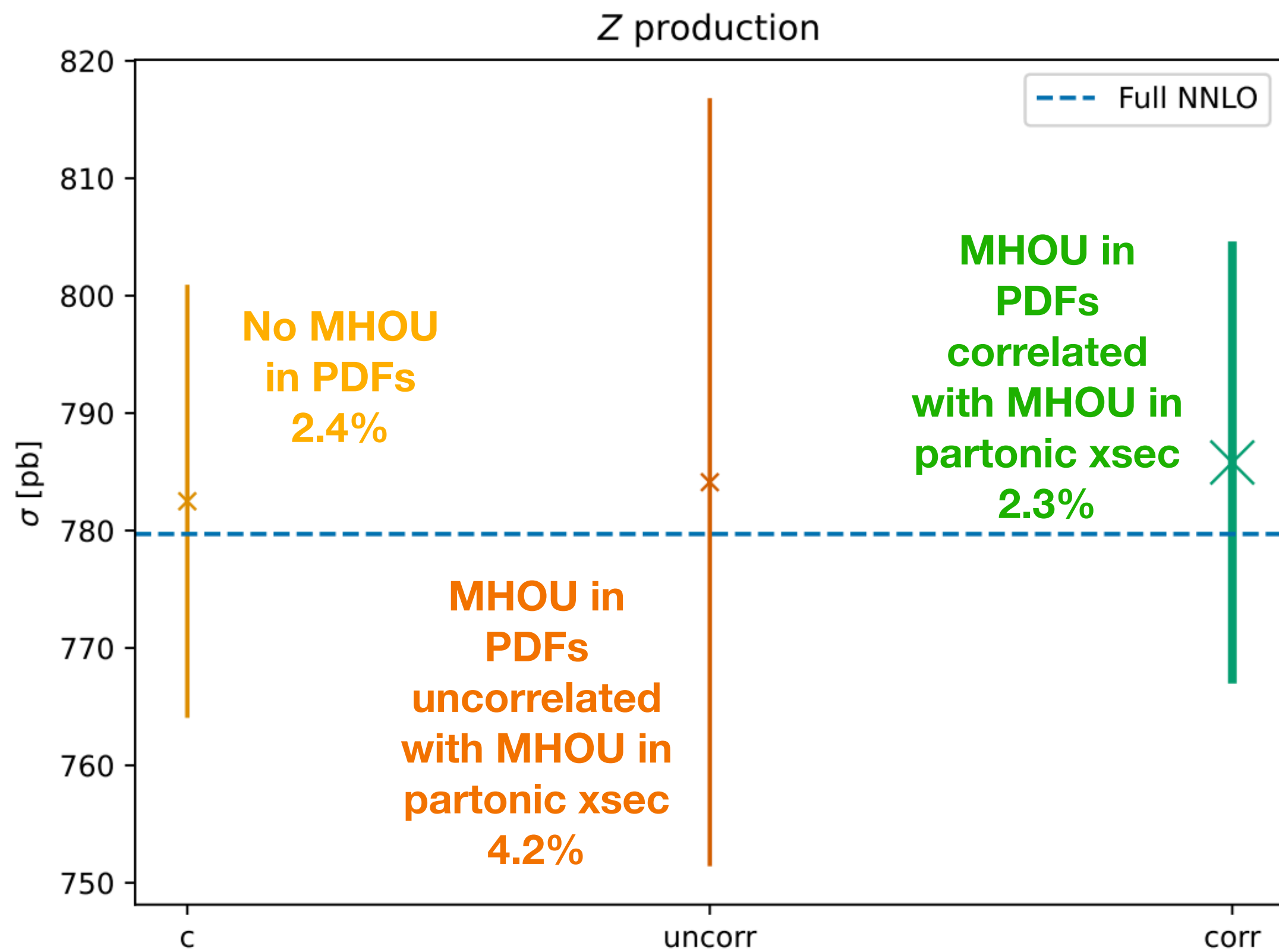
Posterior

$$\chi_n^2 > \langle \chi^2 \rangle_n | \omega^{(n)} = \{1, \dots, 1\} + 4 \text{std}(\chi^2)_n | \omega^{(n)} = \{1, \dots, 1\}$$

# THE MCSCALES APPROACH

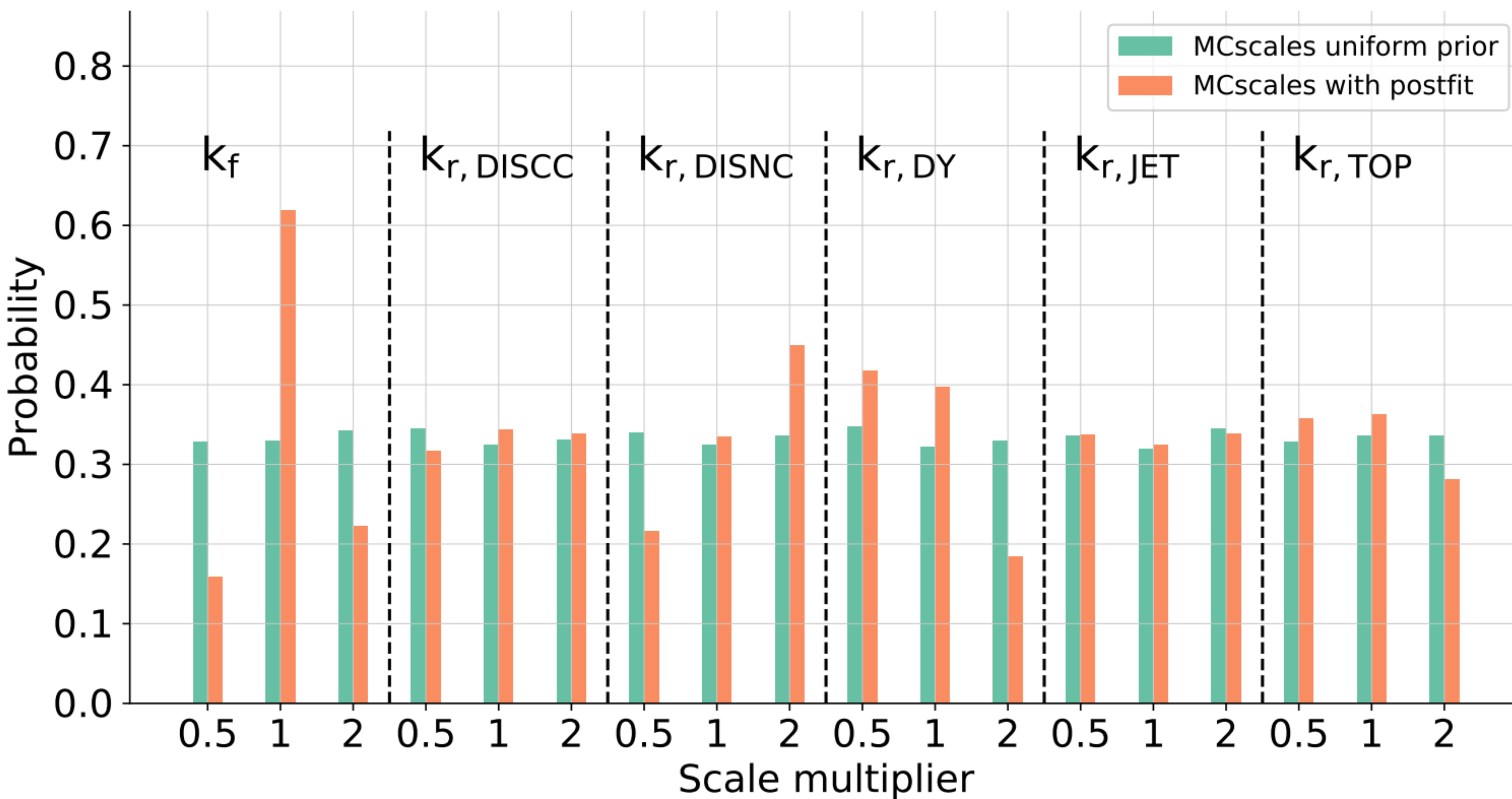
✓ Can compute full PDF+SCALE uncertainty in cross sections at NLO by matching the scales in the hard cross section computation with the scales in the MCscale PDF set: correlation fully taken into account

$$\left\{ \sigma_n = \hat{\sigma}_p(k_f^{(n)}, k_{r_p}^{(n)}) \otimes f_n(k_f^{(n)}, k_{r_p}^{(n)}) \quad \forall n = 1, \dots, N \right\}$$





# THE MCSCALES APPROACH

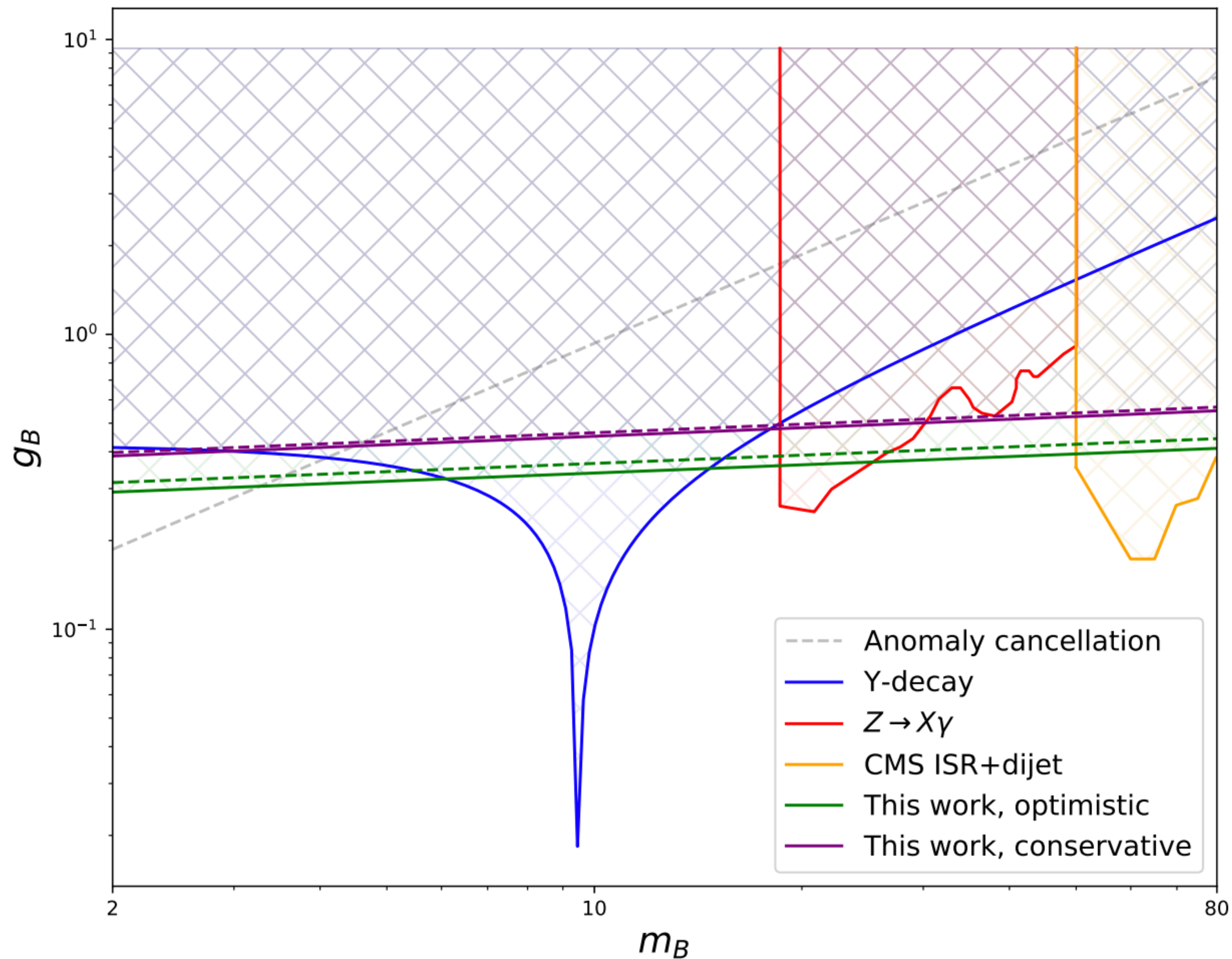


- ✓ Can look at the distribution of each of the scales over replicas.
- ✓ Flat distribution for the MCscales uniform prior.
- ✓ After applying postfit observe preference for central factorisation scale.
- ✓ Each process affected in a different way.

Scale multipliers	Process	Preferred values
$(k_f, k_r)$	DIS CC	(1, 1)
	DIS NC	(1, 2)
	DY	(1, 1)
	Jets	$(1, \frac{1}{2})$
	Top	(1, 1)

# DARK PHOTON

M. McCullough, J. Moore, MU, arXiv:2203.12628



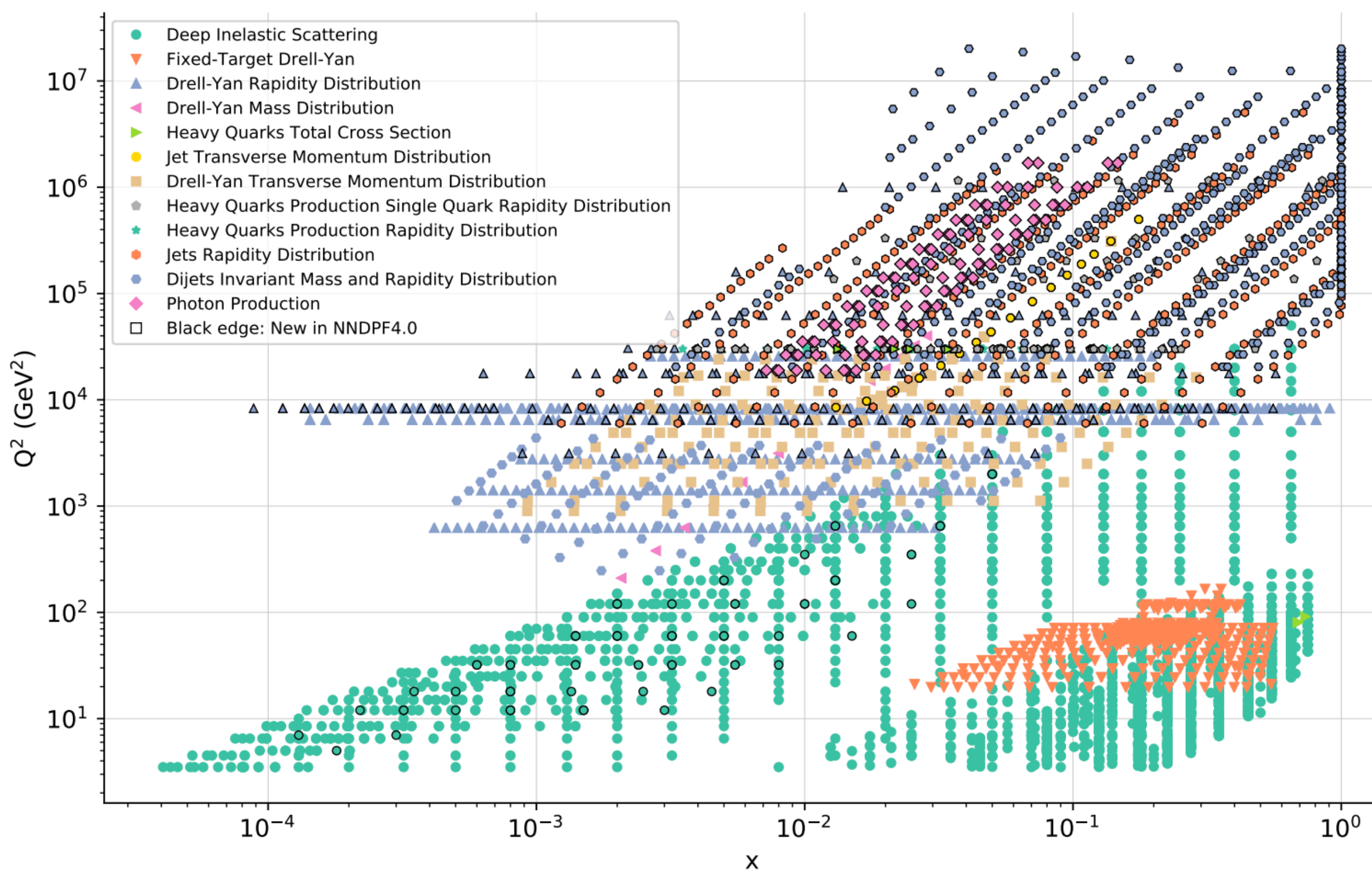
- If there was a lepto-phobic dark photon weakly coupled to quarks, it would appear among the partons of the proton.

$$\mathcal{L}_{\text{int}} = \frac{1}{3} g_B \bar{q} \not{B} q$$

- The presence of the dark Parton would modify the evolution of standard quarks and gluon.
- Precise LHC data can indirectly constrain parameter space of the dark photon in a competitive way compared to direct searches

# #3 – PDF AND NEW PHYSICS INTERPLAY

$$\chi^2 = \frac{1}{N_{\text{dat}}} \sum_{i=1}^{N_{\text{dat}}} (T_i(\{\theta\}, \{c\}) - D_i) \text{cov}_{ij}^{-1} (T_j(\{\theta\}, \{c\}) - D_j)$$



$$T_i(\{\theta\}, \{c\}) = \text{PDFs}(\{\theta\}, \{c\}) \otimes \hat{\sigma}_i(\{c\})$$



(B)SM parameters:  $\alpha_s(M_Z), M_W, \theta_W, \text{SMEFT WCs} \dots$

Parameters determining PDFs at initial scale

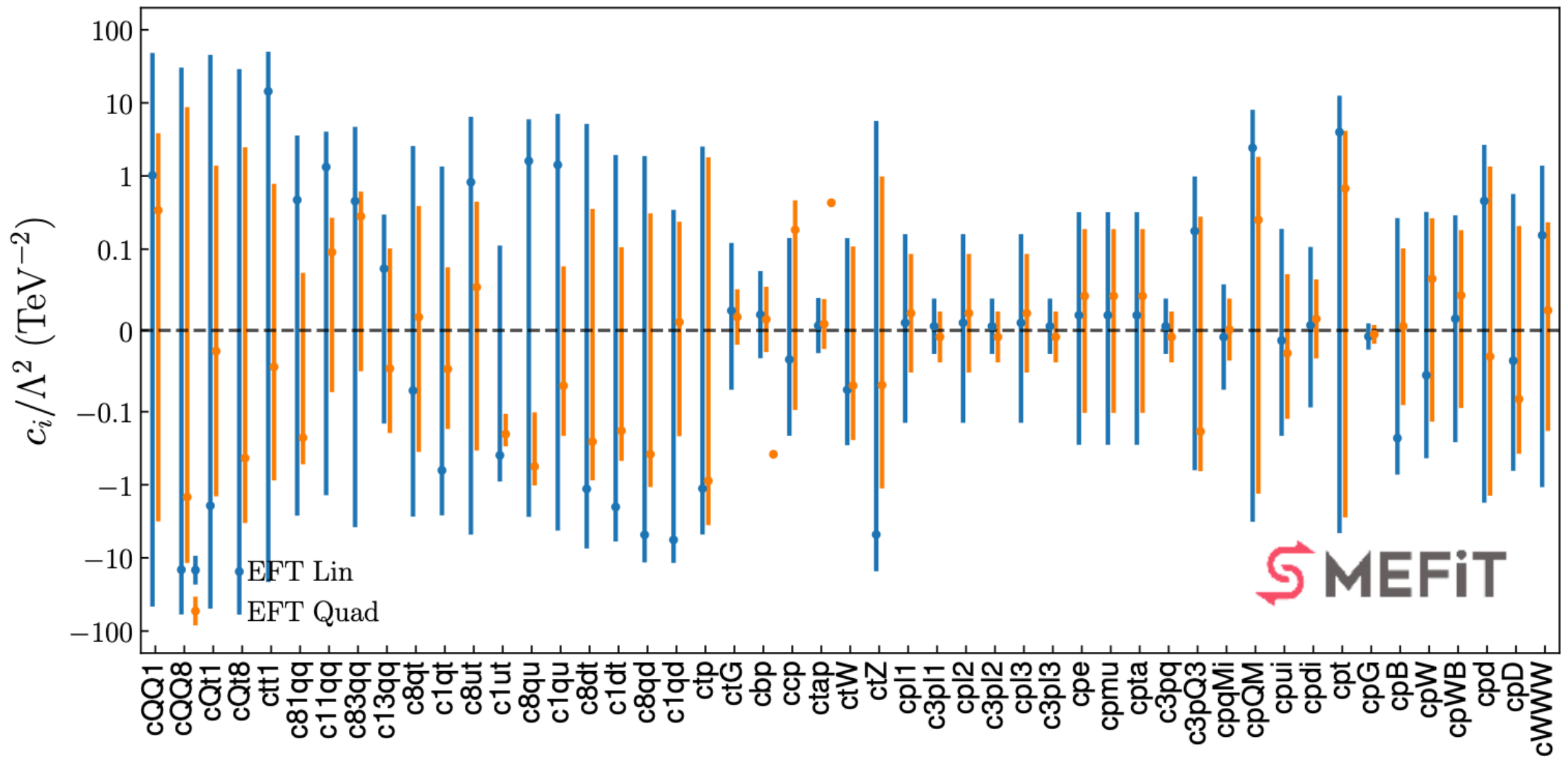
✓ In a PDF fit typically

$$T_i(\{\theta\}) = \text{PDFs}(\{\theta\}, \{c = \bar{c}\}) \otimes \hat{\sigma}_i(\{c = \bar{c}\})$$

# #3 – PDF AND NEW PHYSICS INTERPLAY

$$\chi^2 = \frac{1}{N_{\text{dat}}} \sum_{i=1}^{N_{\text{dat}}} (T_i(\{\theta\}, \{c\}) - D_i) \text{cov}_{ij}^{-1} (T_j(\{\theta\}, \{c\}) - D_j)$$

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$



$$T_i(\{\theta\}, \{c\}) = \text{PDFs}(\{\theta\}, \{c\}) \otimes \hat{\sigma}_i(\{c\})$$



(B)SM parameters:  $\alpha_s(M_Z), M_W, \theta_W, \text{SMEFT WCs} \dots$

Parameters determining PDFs at initial scale

✓ In a PDF fit typically

$$T_i(\{\theta\}) = \text{PDFs}(\{\theta\}, \{c = \bar{c}\}) \otimes \hat{\sigma}_i(\{c = \bar{c}\})$$

✓ In a fit of (B)SM parameters

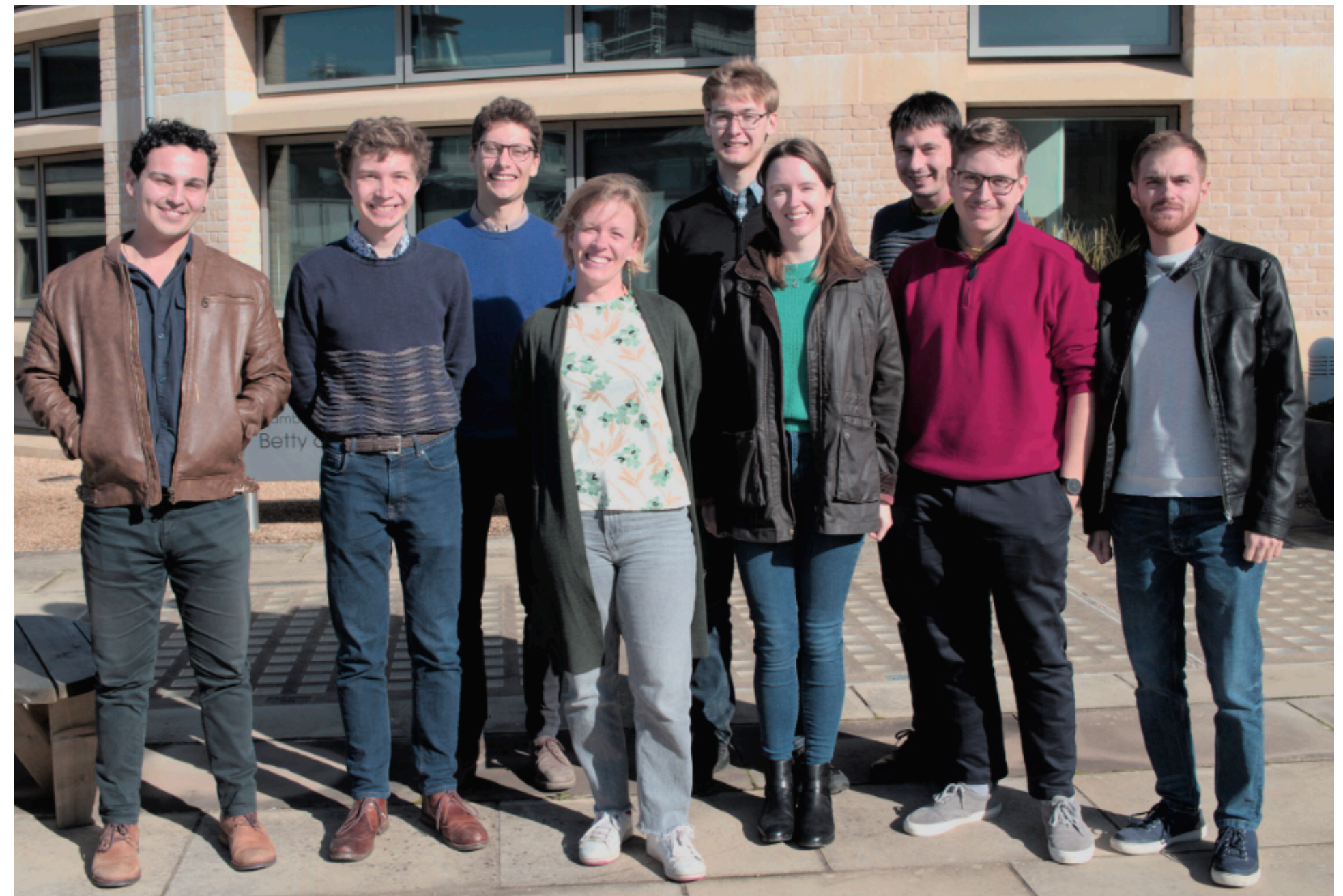
$$T_i(\{c\}) = \text{PDFs}(\{\bar{\theta}\}, \{\bar{c}\}) \otimes \hat{\sigma}_i(\{c\})$$

# #3 – PDF AND NEW PHYSICS INTERPLAY

- From the point of view of PDF fits:
  - ➔ How to make sure that new physics effects are not inadvertently fitted away in a PDF fit?
- From the point of view of SMEFT fits:
  - ➔ Should I make sure I am using a clean set of PDFs in a SMEFT analysis? How to define it? Is it enough?
  - ➔ How would the bounds change if I was consistently using PDFs that include in the fit the same operators that I am fitting?



MU + Z. Kassabov, M. Madigan,  
L. Mantani, J. Moore, M.  
Morales Alvarado, E. Hammou,  
M. Costantini



# #3 – PDF AND NEW PHYSICS INTERPLAY

- From the point of view of PDF fits:

➔ How to make sure that new physics effects are not inadvertently fitted away in a PDF fit?

- ◆ With current experiments, at the HL-LHC found a model involving a new heavy  $W'$  that would affect high energy Drell-Yan tails. If there was a new heavy  $W'$  PDF fits would have same quality but PDFs (hence theory predictions at LHC) would be significantly different. Large- $x$  PDFs would absorb New Physics in this case!
- ◆ Need orthogonal input from low-energy experiments!

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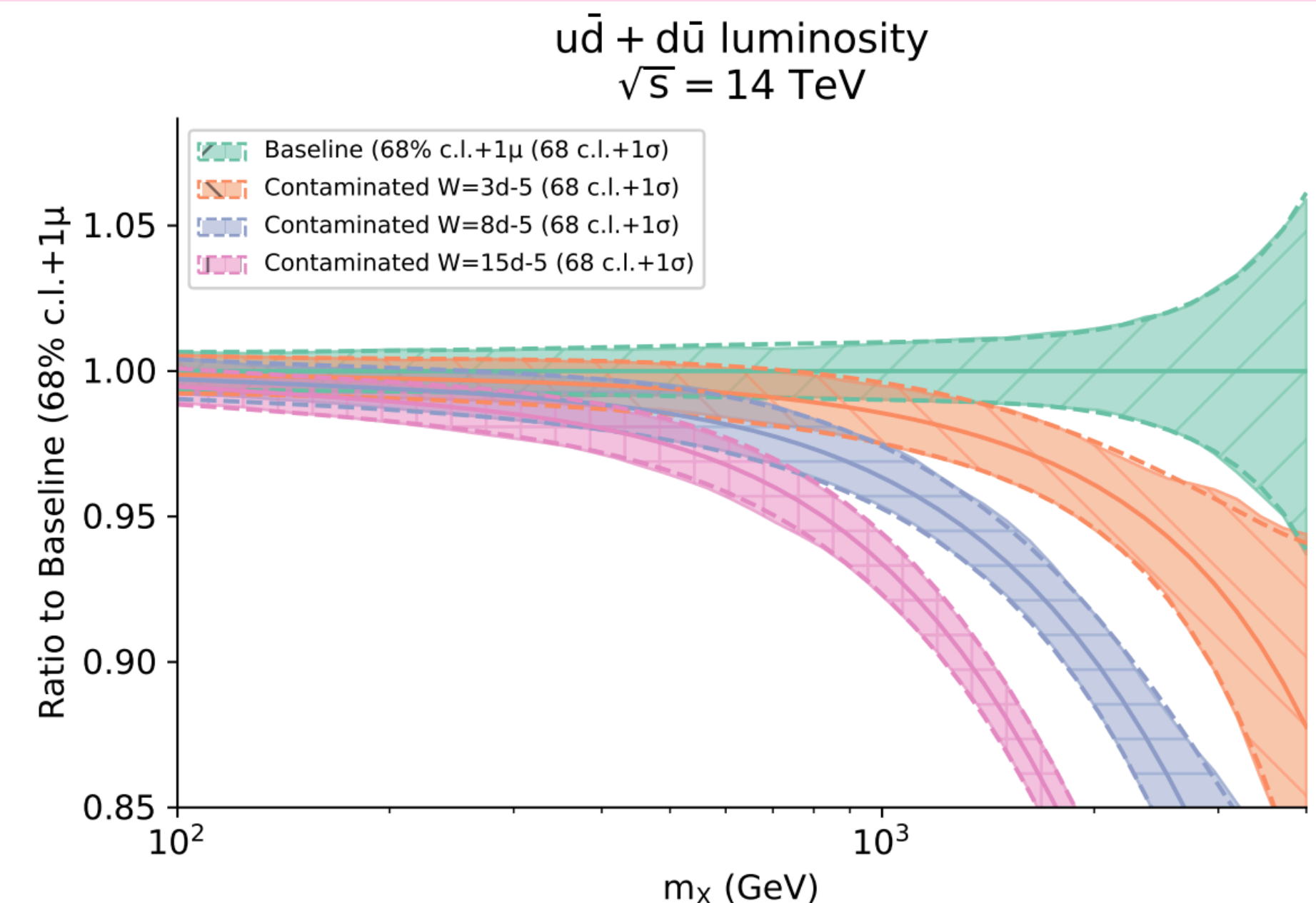
RECEIVED: August 21, 2023  
ACCEPTED: October 28, 2023  
PUBLISHED: November 16, 2023

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## Hide and seek: how PDFs can conceal new physics

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Elie Hammou,<sup>a</sup> Zahari Kassabov,<sup>a</sup> Maeve Madigan,<sup>b</sup> Michelangelo L. Mangano,<sup>c</sup>  
Luca Mantani,<sup>a</sup> James Moore,<sup>a</sup> Manuel Morales Alvarado<sup>a</sup> and Maria Ubiali<sup>a</sup>



# #3 – PDF AND NEW PHYSICS INTERPLAY

- From the point of view of PDF fits:
  - ➔ How to make sure that new physics effects are not inadvertently fitted away in a PDF fit?
- From the point of view of SMEFT fits:
  - ➔ Should I make sure I am using a clean set of PDFs in a SMEFT analysis? How to define it? Is it enough?
  - ➔ How would the bounds change if I was consistently using PDFs that include in the fit the same operators that I am fitting?

$$d\sigma^{pp \rightarrow ab} = \sum_{i,j} \boxed{f_i \otimes f_j \otimes d\hat{\sigma}^{ij \rightarrow ab}} + \dots$$

$f(\{\theta_k\})$        $\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$

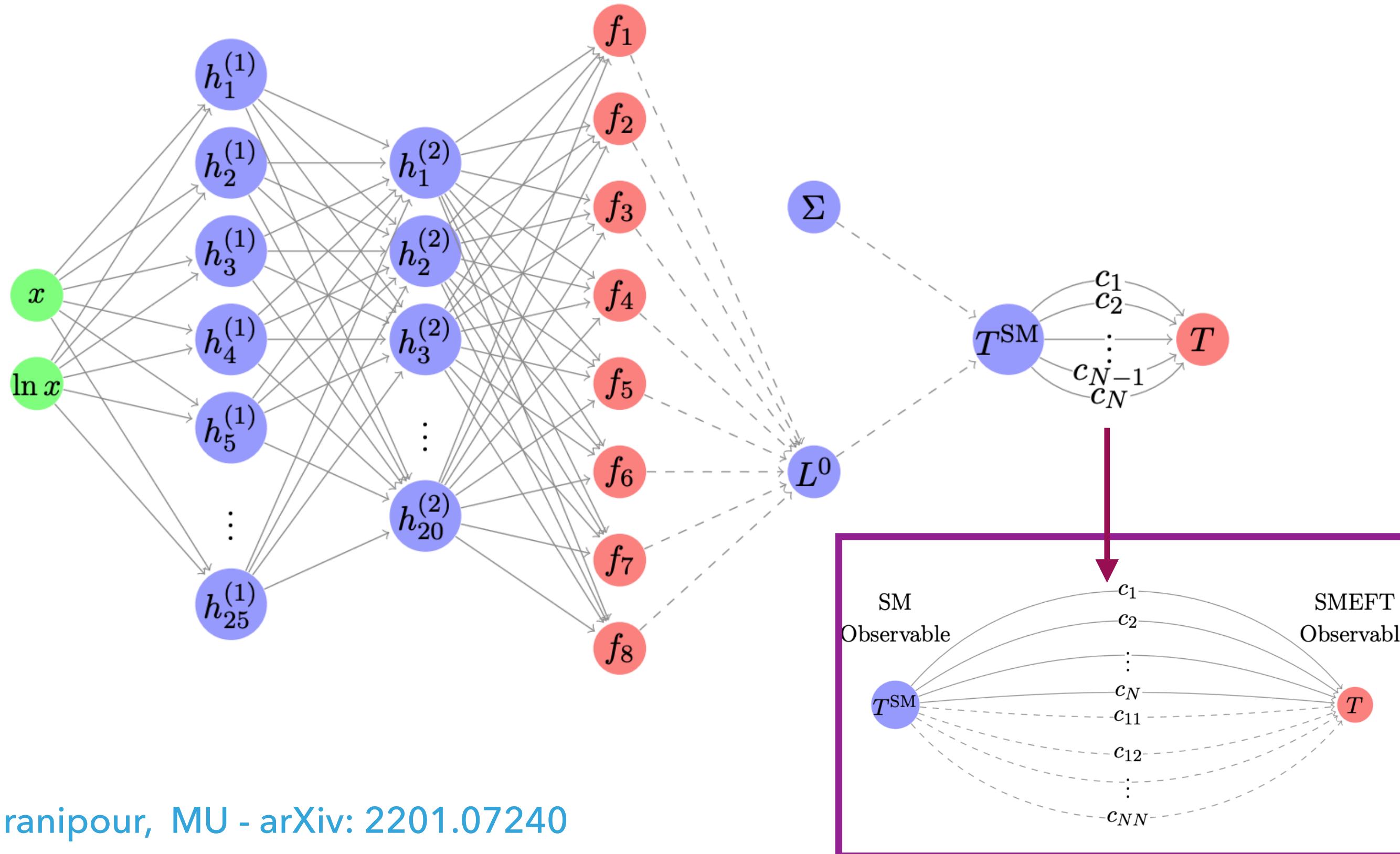
**Simultaneous fits of  
SMEFT and PDFs**

$T(\{\theta_k\}, \{c_i\})$

# #3 – PDF AND NEW PHYSICS INTERPLAY

- How to perform simultaneous fits?
- SimuNET [soon available open-source] yields simultaneous fit of PDFs and SMEFT coefficients, it does not have limit in number of parameters that can be fitted alongside PDFs at the initial scale!

Input layer      Hidden layer 1      Hidden layer 2      PDF flavours      Convolution step      SM Observable      SMEFT Observable



Linear dim-6 operator

$$T(\hat{\theta}) = \Sigma(\{c_n\}) \cdot L^0(\theta) = T^{\text{SM}}(\theta) \cdot \left( 1 + \sum_{n=1}^N c_n R_{\text{SMEFT}}^{(n)} \right)$$

$$T^{\text{SM}}(\theta) = \Sigma^{\text{SM}} \cdot L^0(\theta)$$

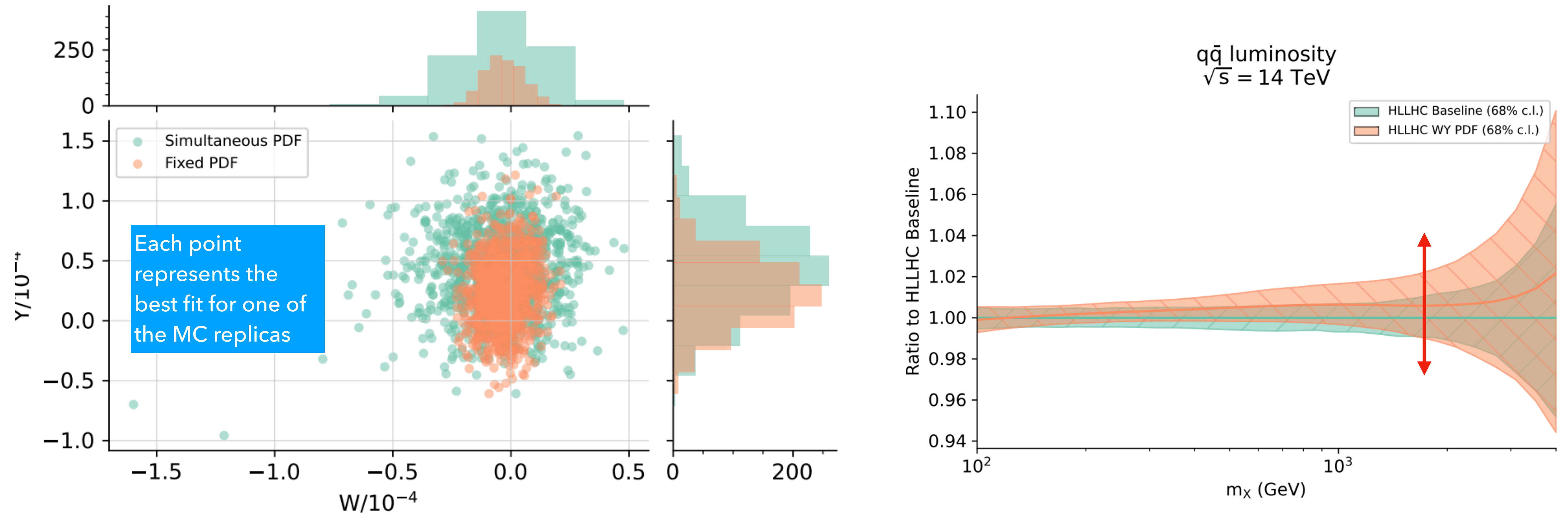
Quadratic dim-6 operator

$$T(\hat{\theta}) = T^{\text{SM}}(\theta) \cdot \left( 1 + \sum_{n=1}^N c_n R_{\text{SMEFT}}^{(n)} + \sum_{1 \leq n \leq m \leq N} c_{nm} R_{\text{SMEFT}}^{(n,m)} \right)$$

$c_n c_m$



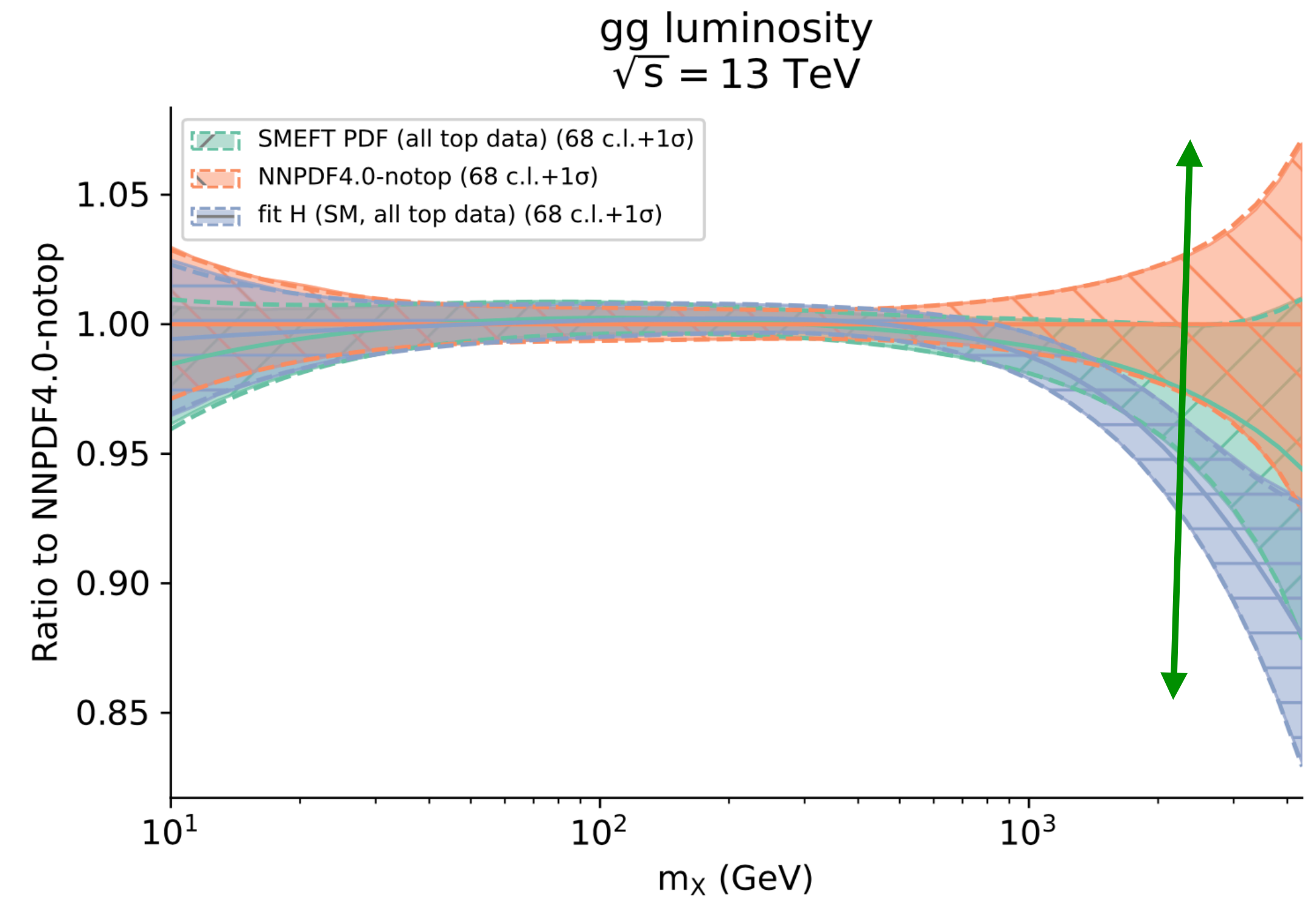
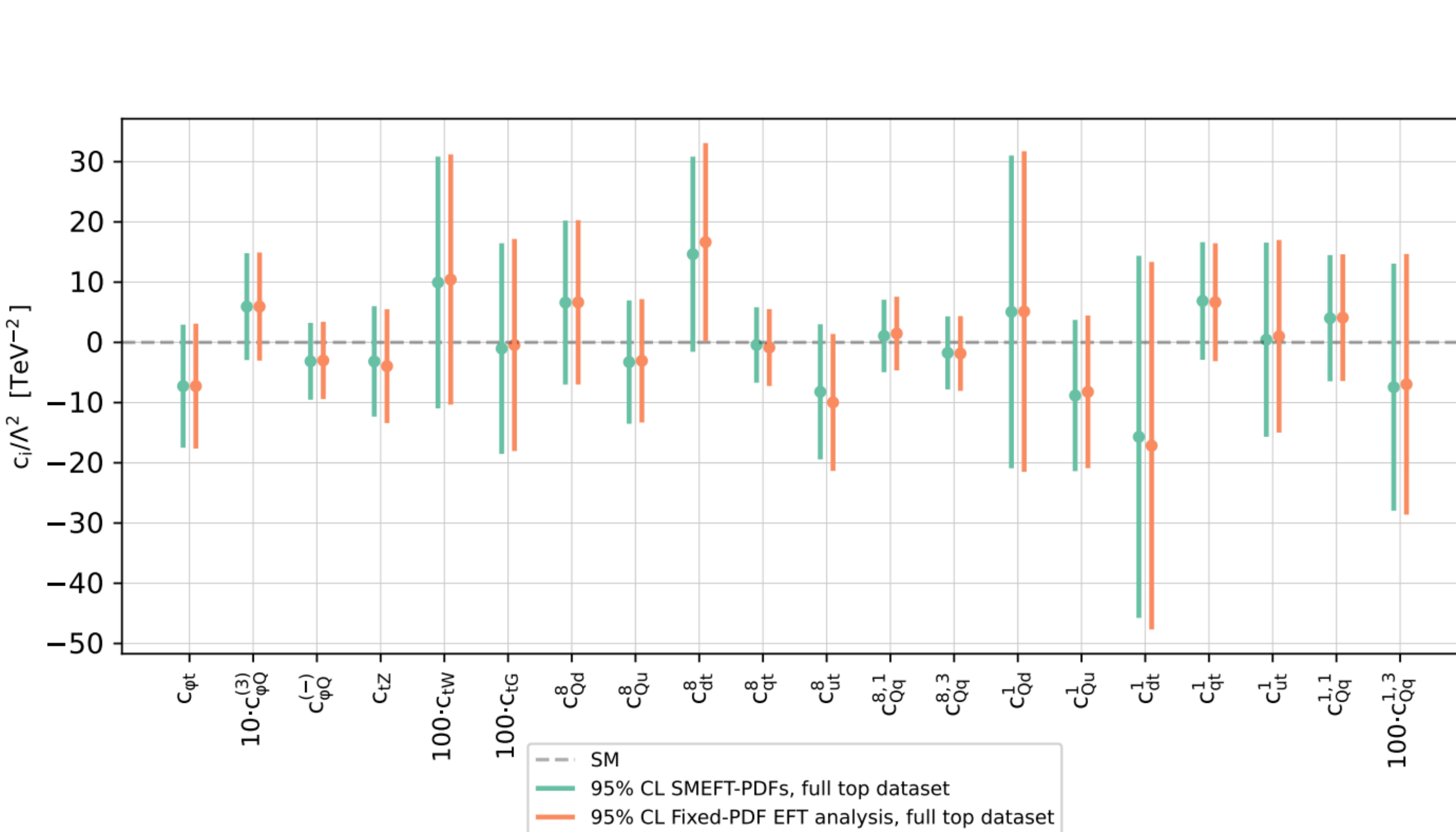
# #3 – PDF AND NEW PHYSICS INTERPLAY



S. Iranipour, MU - arXiv: 2201.07240

- ✓ Simultaneous analysis of PDFs and **Drell-Yan sector** Wilson coefficient of DIS + DY (including HL-LHC projections) using simuNET method shows that at HL-LHC the effect of interplay becomes important as WCs bounds broaden and PDF uncertainties change significantly once SMEFT effects allowed in theory predictions entering PDF fit

# #3 – PDF AND NEW PHYSICS INTERPLAY



Z. Kassabov, M. Madigan, L. Mantani, J. Moore, M. Morales, J. Rojo, MU - arXiv: 2303.06159

✓ Simultaneous analysis of PDFs and more than 20 operators in the **top sector** using simuNET method shows that WCs are stable, while PDF uncertainties broaden and PDF fitted simultaneously alongside WCs sit nicely between PDF-only fit without top data and PDF fit including all top data