

# Electroweak corrections to double Higgs production at the LHC

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Based on: arxiv: 2311.16963

In cooperation with:

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# Introduction



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- Introduction
- Calculation strategy
- Results
- Summary

# Introduction to Higgs Boson



## Standard Model of Elementary Particles

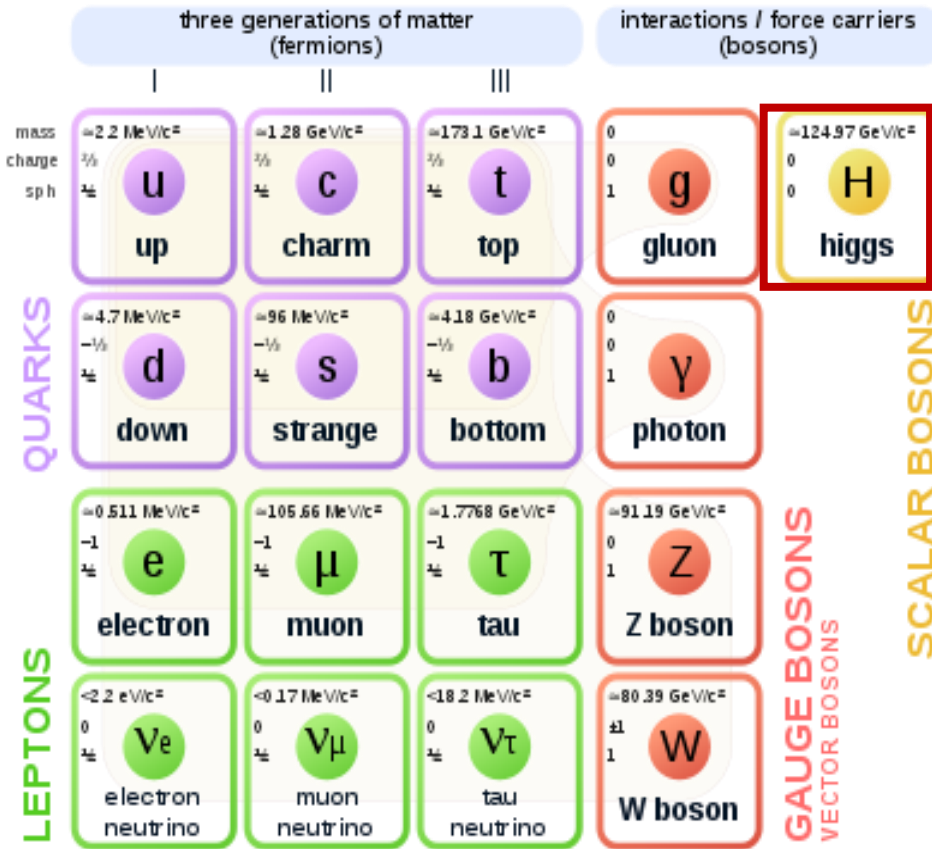
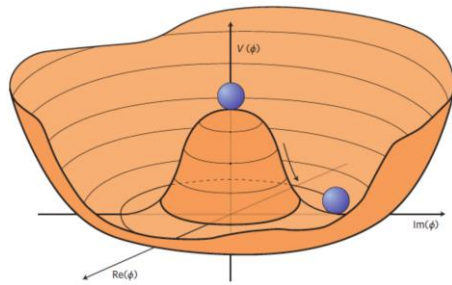


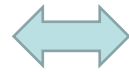
Figure taken from Wikipedia

- Discovery of Higgs boson(2012,LHC): last fundamental particle in SM.
- Experiments at the **ALTA**S and **CMS**: agrees with result SM predicted.
- Problems not solved**: electroweak symmetry breaking, Higgs coupling to SM particles/DM, hierarchy problem... Require new physics beyond SM.
- One promising way probing new physics: precision measurements of the properties of H (for e.g. **Higgs self coupling**).

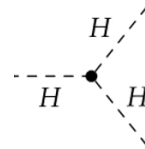
# Higgs self coupling



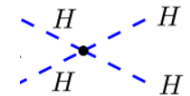
Plot taken from Ellis: [1312.5672](#)



$$V(H) = \frac{m_H^2}{2} H^2 + \lambda_{HHH}^{SM} v H^3 + \frac{1}{4} \lambda_{HHHH}^{SM} H^4$$



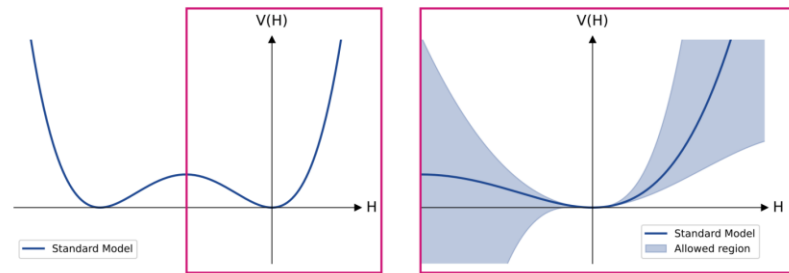
Feasible  
at HL-LHC



20??



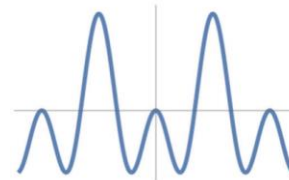
What the SM predicts vs what we know experimentally



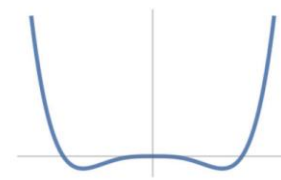
[using current ATLAS limits @ 95% CL]

Plot taken from Moser: [Higgs 2023](#)

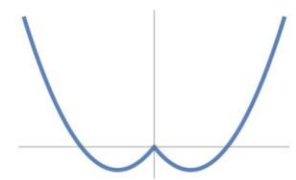
New physics



Nambu-Goldstone Higgs

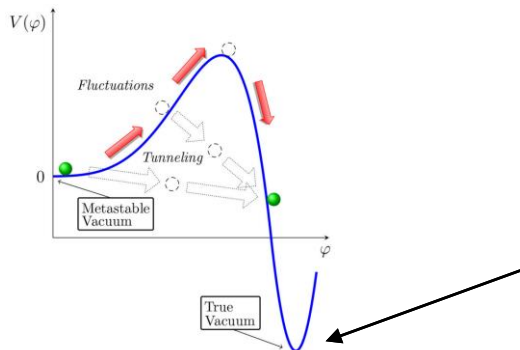


Coleman-Weinberg Higgs



Tadpole-Induced Higgs

Agrawal et al: [1907.02078](#)



What would  
Universe be  
like in such  
vacuum?



big consequences  
for the Universe

Markkanen et al: [1809.06923](#)

# Measurements of Higgs boson coupling

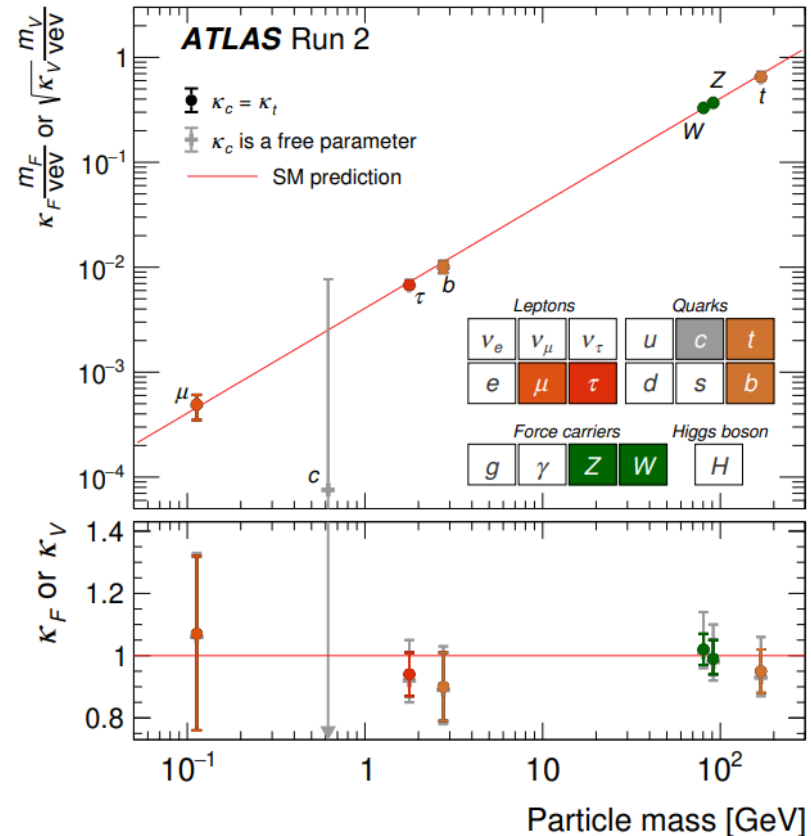


😊  $\mathcal{G}_{Hf\bar{f}}, \mathcal{G}_{HVV}$

- can be measured with high precision.

😞  $\lambda_{HHH}, \lambda_{HHHH}$

- require multi-Higgs production, small cross sections.
- Mixed with complicated background.



ATLAS:2207.00092

Run 2 $\delta_\mu^{\text{tot}}$ [%]	HL-LHC $\delta_\mu^{\text{tot}}$ ( $\delta_\mu^{\text{th}}$ ) [%]
$-1.0 < \lambda/\lambda_{\text{SM}} < 6.6$	$0.5 < \lambda/\lambda_{\text{SM}} < 1.5$

# Status of QCD corrections



- NLO QCD

- NLO QCD with full top-quark mass dependence, [Borowka et al:1604.06447](#)
- NLO QCD matched to parton shower, [Heinrich et al:1703.09252](#)
- NLO QCD with soft-gluon resummation, [Ferrera et al: 1609.01691](#)

- NNLO QCD

- NNLO QCD in heavy-top limit (HTL) approximation, [Florian et al:1305.5206](#)
- NNLO in HTL+ NLO with full top-quark mass dependence, [Florian et al:2106.14050](#)
- NNLO QCD in HTL matched to parton shower, [Alioli et al: 2212.10489](#)

- NNNLO QCD

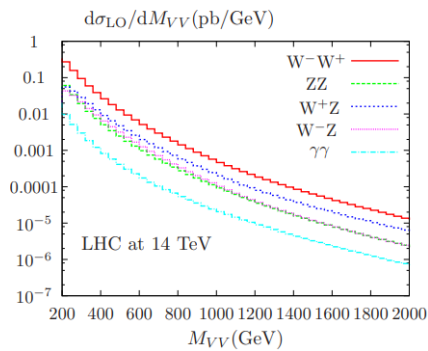
- NNNLO QCD in HTL, [Chen et al:1909.06808](#)
- NNNLO in HTL include the top-quark mass effects, [Chen et al:1912.13001](#)
- NNNLO in HTL + NLO with full top-quark mass dependence + soft-gluon resummation, [Ajith et al:2209.03914](#)

Process	Theory	$\sigma_{\text{th}}$ [pb]	$\delta_{\text{th}}$ [%]	$\delta_{\text{PDF}}$ [%]	$\delta_{\alpha_s}$ [%]
$ggF$ $HH$	$N^3\text{LO}_{\text{HTL}}$	0.03105	+2.2 -5.0	±2.1	±2.1
	$\text{NLO}_{\text{QCD}}$				

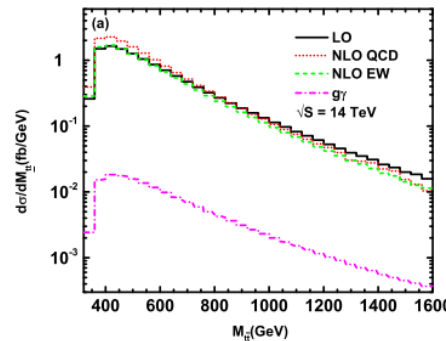
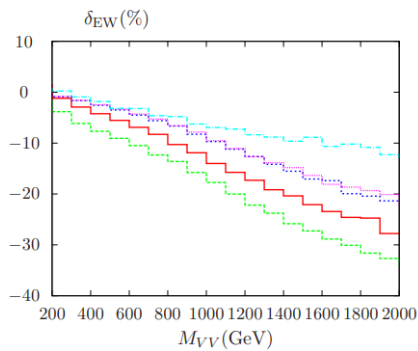


- Unknown size of EW corrections
  - Biggest uncertainties from theoretical side
- NLO EW corrections are notably significant at high energy region

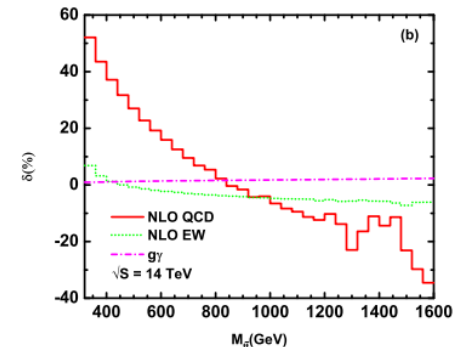
➤ Sudakov enhancement:  $\alpha \sim 0.7\% \rightarrow \frac{\alpha}{4\pi \sin^2\theta_W} \log^2\left(\frac{s}{m_Z^2}\right) \Big|_{s=2000^2} \sim 10\% \sim \alpha_s$



A Bierweiler et al:1305.5402

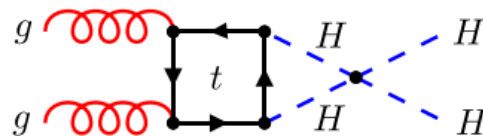


Zhang et al: 1407.1110



- Higgs quartic coupling only emerges at the NLO EW level

➤ Constrained on  $\lambda_{HHHH}^{SM}$  indirectly from NLO EW correction



# Status of NLO EW corrections



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- Partial results
  - Two-loop box diagrams, [Davies et al:2207.02587](#)
  - Top-quark Yukawa corrections, [Muhlleitner et al:2207.02524](#)
  - Higgs self-coupling corrections, [Borowka et al: 1811.12366](#)
  - HTL and Neglecting diagrams with massless fermion loops, [Davies et al: 2308.01355](#)
- Groups working on this topic:
  - See Hantian Zhang's talk at Higgs 2023: [HTL + partial results](#)
  - See Xiao Zhang's talk at Higgs 2023: [partial results](#)
  - See Thomas Stone's talk at Higgs 2023: [partial results](#)



# Calculation strategy



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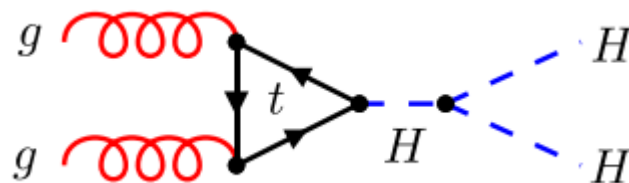


- Introduction
- **Calculation strategy**
- Results
- Summary

# EW corrections to double H production at the LHC

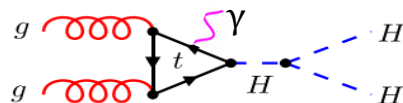
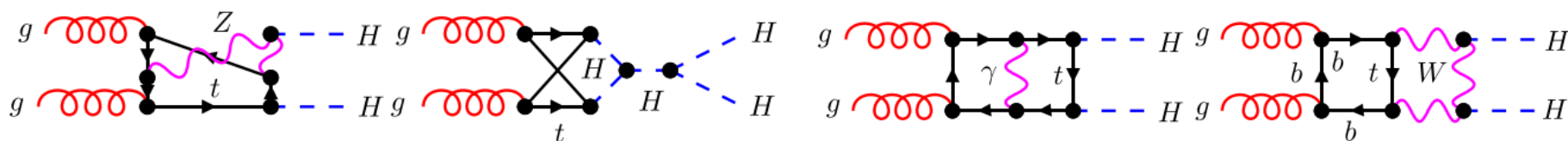


LO diagrams:



Typical Feynman diagrams at LO

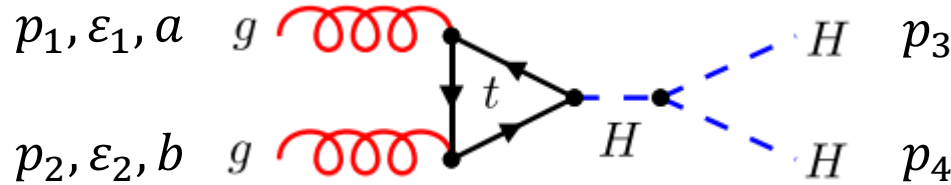
NLO diagrams:



Forbidden due to Furry Theorem

Typical Feynman diagrams at NLO EW

# Amplitudes of $gg \rightarrow HH$



## ● Amplitude Structure:

$$\mathcal{M}_{ab} = \delta_{ab} \epsilon_1^\mu \epsilon_2^\nu \mathcal{M}_{\mu\nu}$$

$$\mathcal{M}^{\mu\nu} = F_1 T_1^{\mu\nu} + F_2 T_2^{\mu\nu} + \Delta_0^{\mu\nu} + \Delta_5^{\mu\nu}$$

- General decomposition at any number of loop.
- $\Delta_0^{\mu\nu}$ : depends on  $p_1^\mu$  or  $p_2^\nu$ . No contribution at the matrix element level.
- $\Delta_5^{\mu\nu}$ : depends on Levi-Civita tensor. No contribution at the matrix element squared level at NLO EW.
- $F_1, F_2$ : Form factors.

# Calculation of form factors



- Form factors can be decomposed into:

$$F_{1,2}(x) = \sum_i d_i(x) FI_i(x)$$

$$\begin{aligned} x: \hat{s} &= (p_1 + p_2)^2, \\ \hat{t} &= (p_1 - p_3)^2. \end{aligned}$$

- Reduce  $FI_i(x)$  to master integrals (IBP):

$$\{FI_i(x)\} = \left\{ \sum_k c_{i,k}(x) I_k(x) \right\}$$

- $d_i(x)$  and  $c_{i,k}(x)$  are analytic.
- A huge number of  $I_k$  need to be calculated.
- The number of  $\{I_k\} < \{FI_i\}$ .
- The number of  $I_k$  is finite.
- We can construct the different equations for  $I_k$  and solve them. 12/22

# Different equations for $I_k$



Construct differential equations (DEs):  $\vec{I}(x) = \{I_1(x), I_2(x) \dots I_N(x)\}$

$$\frac{dI_m(x)}{dx} = \sum_n A_{m,n}(x) I'_n(x) \xrightarrow{\text{IBP}} \frac{d\vec{I}(x)}{dx} = A(x)\vec{I}(x)$$

- $\vec{I}(x)$  can be expanded as a **power expansion near  $x_0$** ,

- regular:  $S = \{0\}, k_0 = 0,$

- singular:  $S = \{-2\epsilon, 1 + \epsilon \dots\}, k_\mu \geq 0,$

$$I_i(x) = \sum_{\mu \in S} (x - x_0)^\mu \sum_{k=0}^{k_\mu} \log(x - x_0)^k \sum_{n=0}^m c_{i,\mu,k,n} (x - x_0)^n$$

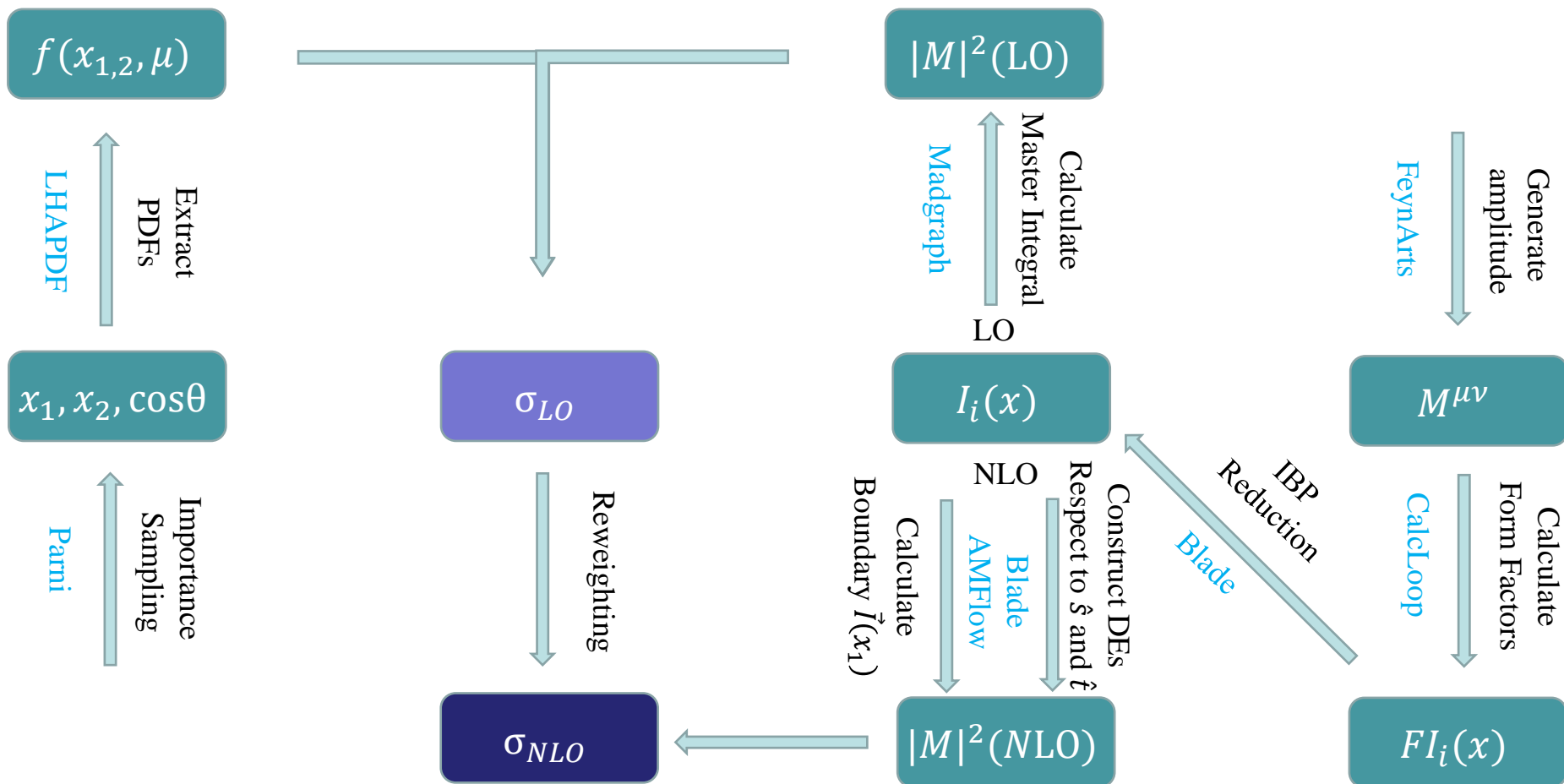
- $c_{i,\mu,k,n}$  can be determined once any boundary  $\vec{I}(x_1)$  are provided.

- $\vec{I}(x_1)$  can be determined by AMFlow

- Taking adequate **expansion order  $m$** , we can eventually achieve predictions with high precision.

- $\vec{I}(x)$  can be evaluated at any points of  $x$  efficiently.

# Calculation flowchart



# Calculation strategy



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- Introduction
- Calculation strategy
- **Results**
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# Input Parameters



$$m_t = 172.69 \text{ GeV}$$

PDG2022

$$\frac{m_H^2}{m_t^2} = \frac{12}{23}, \quad \frac{m_Z^2}{m_t^2} = \frac{23}{83}, \quad \frac{m_W^2}{m_t^2} = \frac{14}{65},$$

$$G_\mu = 1.166378 \times 10^{-5} \text{ GeV}^{-2}$$

$$\alpha = \frac{\sqrt{2}}{\pi} G_\mu m_W^2 \left( 1 - \frac{m_W^2}{m_Z^2} \right)$$

CKM=1

PDFs: NNPDF31\_nlo\_as\_0118

on-shell renormalization: masses and fields;  $G_\mu$ -scheme: Electromagnetic coupling

[Denner et al:1912.06823](#)

$$D=4-2\varepsilon, \quad \varepsilon = \pm 1/1000$$

$$\sigma(\varepsilon) = a_0 + a_1\varepsilon + a_3\varepsilon^2 + \dots$$

$$\sigma(0) \sim \frac{\sigma(+1/1000) + \sigma(-1/1000)}{2} = a_0 + a_3\varepsilon^2 + \dots$$



# Results: Total cross sections

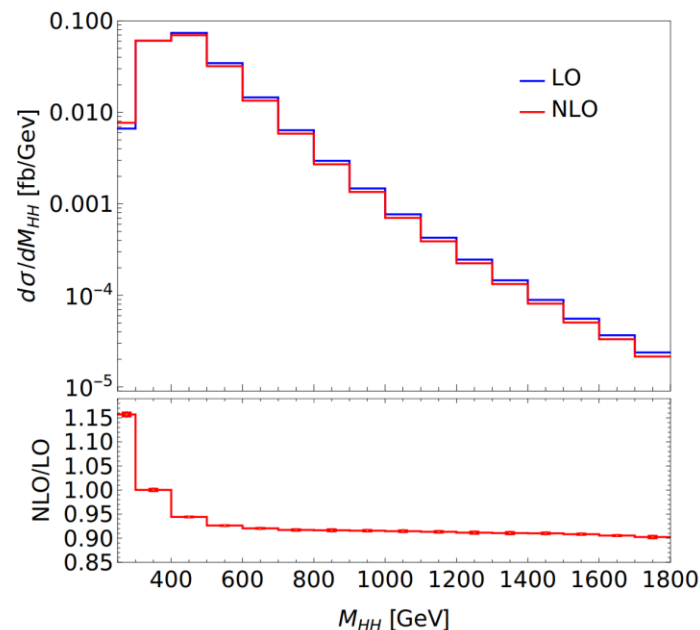


$\mu$	$M_{HH}/2$	$\sqrt{p_T^2 + m_H^2}$	$m_H$
LO	19.96(6)	21.11(7)	25.09(8)
NLO	19.12(6)	20.21(6)	23.94(8)
$\mathcal{K}$ -factor	0.958(1)	0.957(1)	0.954(1)

LO and NLO EW corrected integrated cross sections (in fb) 14 TeV LHC.

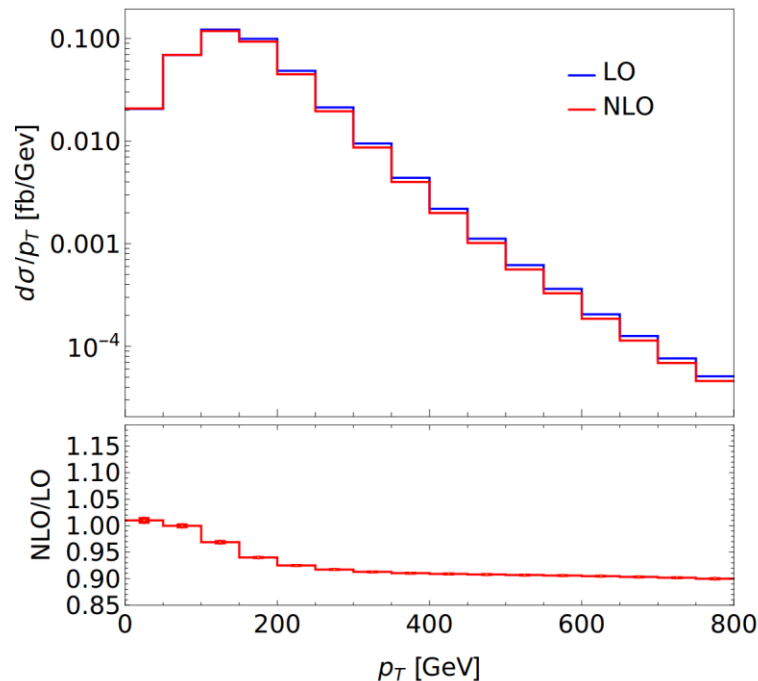
- Differences with varying scale choices are around 20%.
  - Huge scale uncertainties. Can be reduced by including QCD corrections.
- K-factor is insensitive to the scale choice.
  - EW corrections beyond NLO are on the order of a few thousandths.
- The statistical uncertainty for the K-factor is smaller than that of  $\sigma_{LO,NLO}$ .
  - K-factor can get a controllable error with far fewer events.

# Results: Differential cross sections



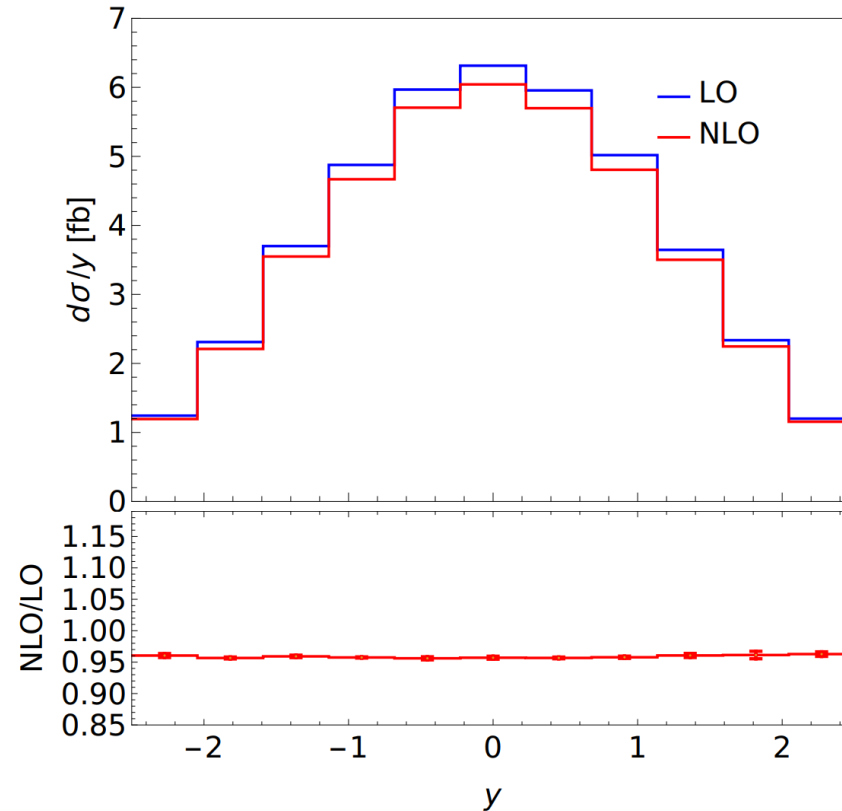
- Big positive corrections at the HH threshold.
  - Enhancement due to  $\sigma_{LO}(\sqrt{\hat{s}} = 2m_H) \sim 0$ .
- -10% correction at high energy region.
  - EW Sudakov effects.
- Tiny cross section at high energy region
  - Gluon PDFs are highly suppressed at high energy region.

# Results: Differential cross sections



- Positive corrections at the beginning of the spectrum.
  - The events in this region are mixed with high  $\sqrt{\hat{s}}$  and low  $\sqrt{\hat{s}}$ .
- -10% correction at high energy region.
  - EW Sudakov effects.

# Results: Differential cross sections

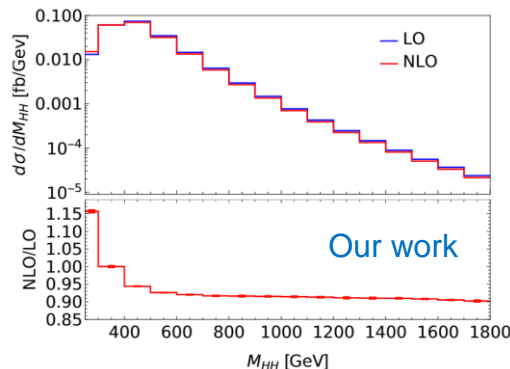
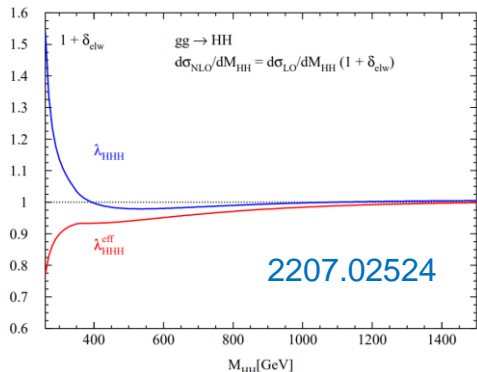


- Flat corrections at around -4%.
  - Similar to the total cross section

# Results: comparisons with other publications



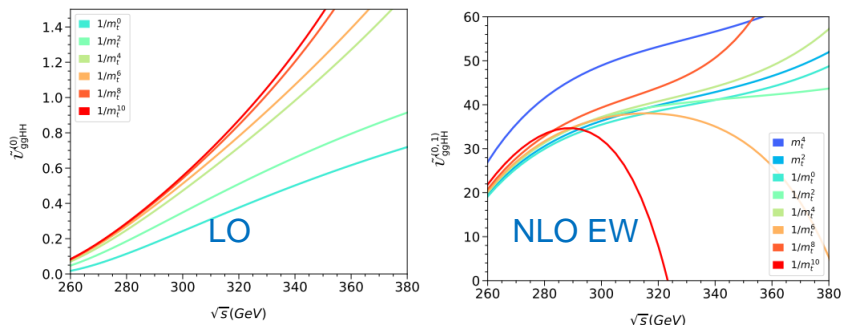
- Top-quark Yukawa corrections, [Muhlleitner et al:2207.02524](#)



- Similar Enhancement at Threshold
- Differences appear at for effective coupling

- HTL and neglecting diagrams with massless fermion loops, [Davies et al: 2308.01355](#)

2308.01355



- ~65% corrections at  $\sqrt{\hat{s}}=260$  GeV
- Our full results reveal the correction is 34% and 57% once neglect the diagram contains only mass less fermion

- Xiao Zhang's talk at Higgs 2023

- ~+1% corrections when only considering Top-Yukawa corrections and Higgs self coupling corrections.
- Our full results reveal the correction is ~-4%.

# Summary



- **Higgs self coupling** is important to identify the Higgs potential and to probe new physics.
- The study of  $\sigma(\text{HH})$  is the **best way** to extract the Higgs self coupling.
- Our **full calculation** includes all the diagrams and all the mass effects.
- **$\sim -4\%$**  EW corrections at total cross section level.
- For dimensionful observables, EW corrections reach up to  **$+15\%$**  at the beginning of the spectrum and  **$-10\%$**  in the tail.
- Our results suggest that the remained uncertainties from theoretical side is overall about **few percent** and it's **precise enough** for the measurements at the HL-HLC.

Thanks for your attention!

# Amplitudes of $gg \rightarrow HH$



$$\mathcal{M}^{\mu\nu} = F_1 T_1^{\mu\nu} + F_2 T_2^{\mu\nu}$$

- Tensor Deconposition:

- Is not unique, we adopt: [Plehn et al:9603205](#)

$$T_1^{\mu\nu} = g^{\mu\nu} - \frac{p_1^\nu p_2^\mu}{p_1 \cdot p_2},$$

$$T_2^{\mu\nu} = g^{\mu\nu} + \frac{1}{p_T^2 (p_1 \cdot p_2)} \left[ 2 (p_1 \cdot p_2) p_3^\nu p_3^\mu - 2 (p_1 \cdot p_3) p_3^\nu p_2^\mu - 2 (p_2 \cdot p_3) p_3^\mu p_1^\nu + m_H^2 p_1^\nu p_2^\mu \right],$$

- Projector

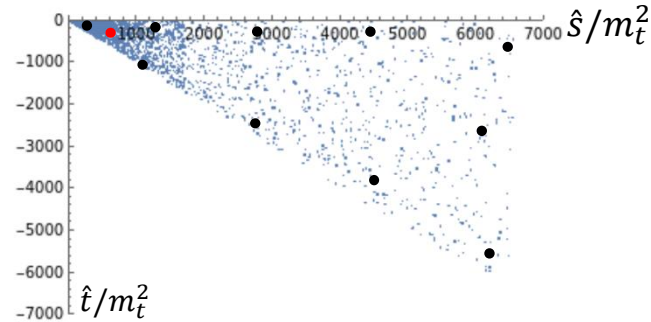
$$P_1^{\mu\nu} = +\frac{1}{4} \frac{D-2}{D-3} T_1^{\mu\nu} - \frac{1}{4} \frac{D-4}{D-3} T_2^{\mu\nu},$$

$$P_2^{\mu\nu} = -\frac{1}{4} \frac{D-4}{D-3} T_1^{\mu\nu} + \frac{1}{4} \frac{D-2}{D-3} T_2^{\mu\nu},$$

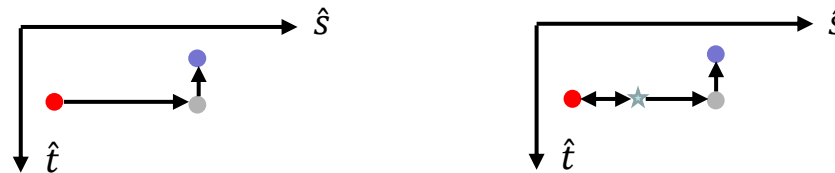
- Form Factor

$$F_1 = P_1^{\mu\nu} \mathcal{M}_{\mu\nu}, \quad F_2 = P_2^{\mu\nu} \mathcal{M}_{\mu\nu}.$$

# Evaluation of $I_k$



## ● The evaluation trajectory:



- Boundary point: ● evaluated by AMFlow
- Phase space point: ●
- Singular point: ★
- Auxiliary point: ●
- Checking point: ● evaluated by AMFlow