

Top-Bottom Interference Contribution to Fully-Inclusive Higgs Production

Marco Niggetiedt

with M. Czakon, F. Eschment, R. Poncelet and T. Schellenberger

based on

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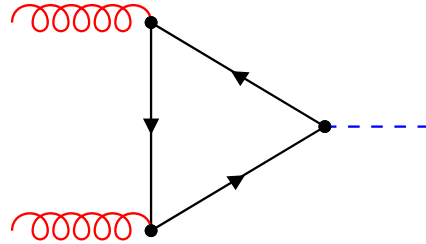
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Boston
June 4th 2024

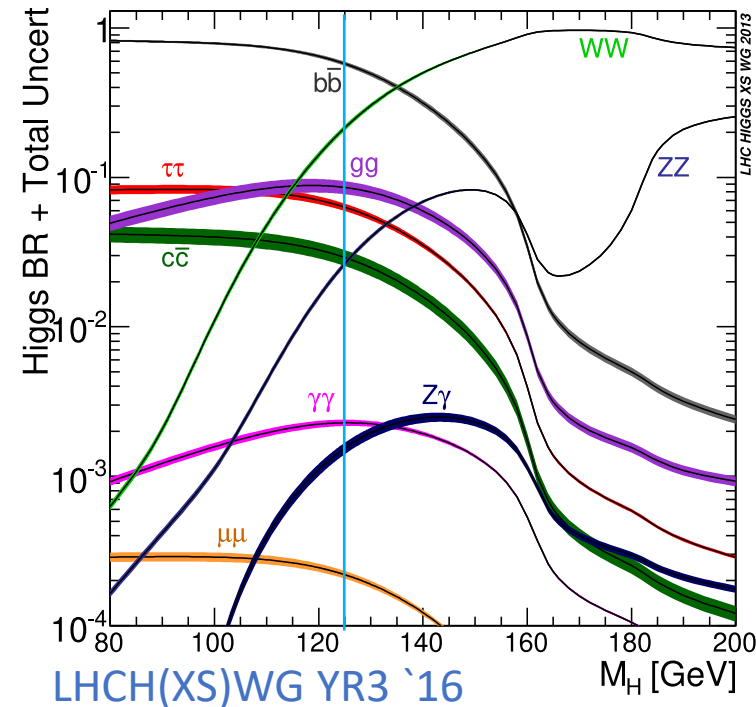
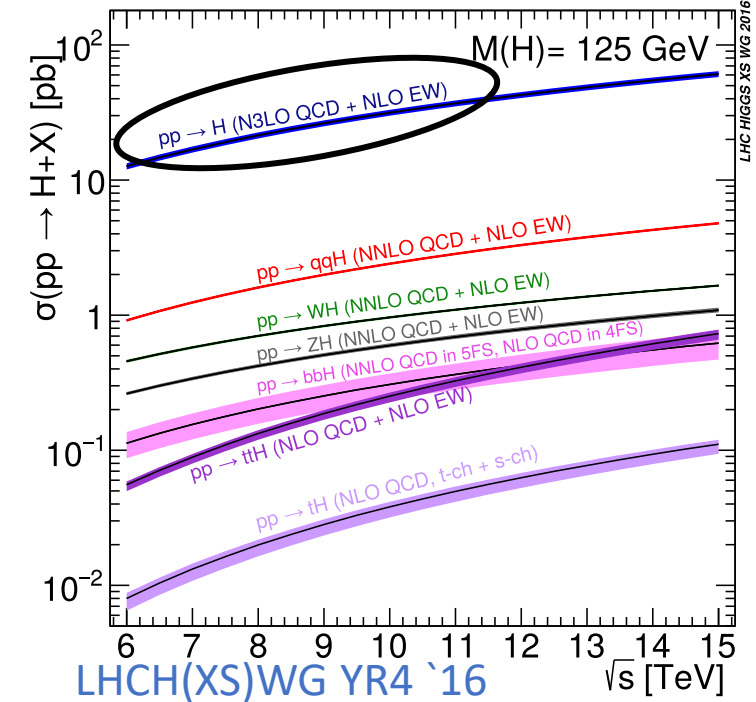
Gluon fusion

- Gluon fusion is the predominant Higgs boson production mode at the LHC
 - loop induced process



- Higgs boson plays unique role in the SM:
 - Only scalar particle
 - Only particle with Yukawa interactions to fermions
- Predictions for gluon fusion cross section directly impact extraction of Higgs couplings from experimental measurements
- Reducing theory uncertainty is crucial for facilitating high precision measurements of Higgs couplings at the LHC
- High luminosity LHC projections anticipate uncertainty $\mathcal{O}(2\%)$ and theory uncertainty to be halved

WG2 Report `19



Inclusive gluon fusion cross section (YR4 `16)

- Inclusive gluon fusion cross section according to [LHCH\(XS\)WG YR4 `16](#) at the LHC at 13 TeV:

48.58 pb =	16.00 pb	(+32.9%)	(LO, rEFT)
	+ 20.84 pb	(+42.9%)	(NLO, rEFT)
	- 2.05 pb	(-4.2%)	((<i>t, b, c</i>), exact NLO)
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[LHCH\(XS\)WG YR4 `16](#)

[Anastasiou, Duhr, Dulat, et al. `16](#)

- Sources of uncertainty:

$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
+0.10 pb -1.15 pb	± 0.18 pb	± 0.56 pb	± 0.49 pb	± 0.40 pb	± 0.49 pb
+0.21% -2.37%	$\pm 0.37\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

← All sources $\mathcal{O}(1\%)$

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Mistlberger `18

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[McGowan, Cridge, Harland-Lang et al. `22](#)
[Falcioni, Herzog, Moch, et al. `23/`24](#)
[NNPDF Collaboration `24](#)

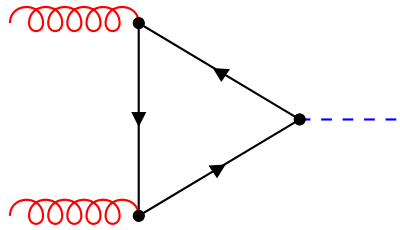
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Order by order in perturbation theory

- LO contribution exactly known for almost 50 years

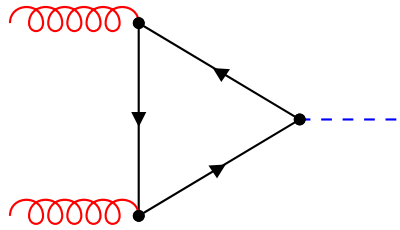


Georgi, Glashow, Machacek, et al. '78

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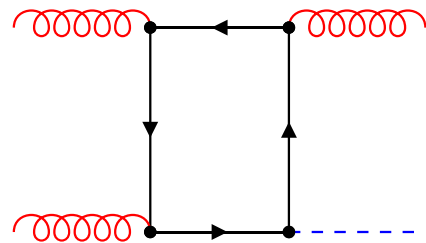
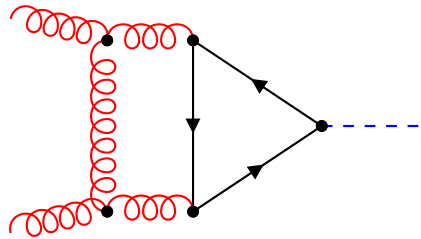
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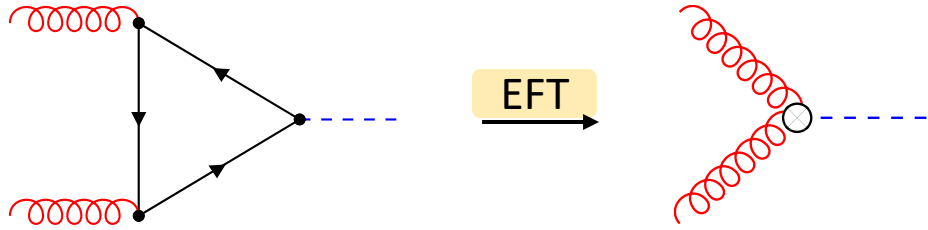
- NLO contribution exactly known for arbitrary quark masses running in the loop Graudenz, Spira, Zerwas '93



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Inclusive cross section in (r)EFT

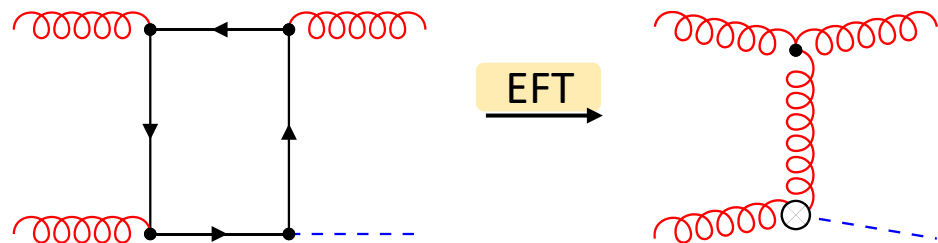
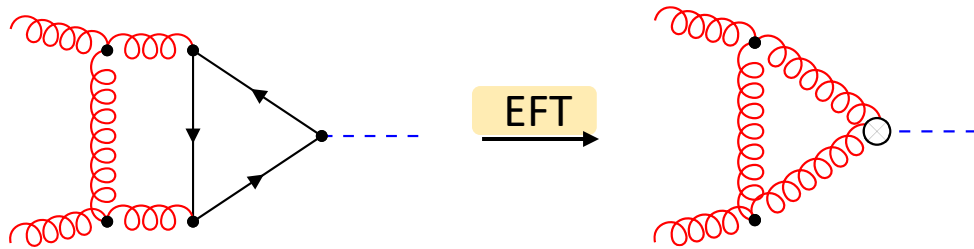
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Georgi, Glashow, Machacek, et al. '78

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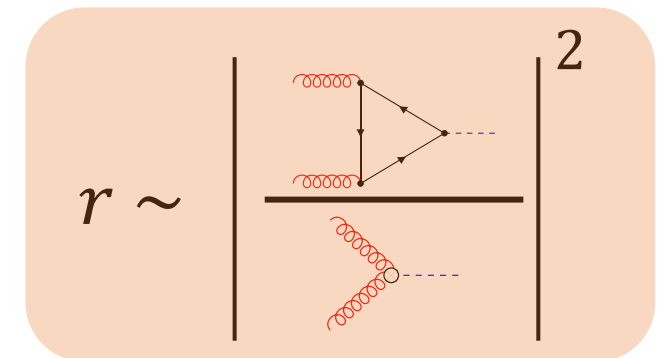
Dawson '91

Djouadi, Spira, Zerwas '91

$$\mathcal{L}_{\text{SM}} \xrightarrow{m_t \rightarrow \infty} \mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM},5} - \frac{1}{4} C H G_{\mu\nu}^a G_a^{\mu\nu}$$

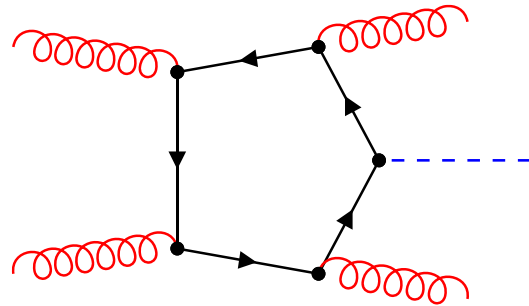
Chetyrkin, Kniehl, Steinhauser '98
Schröder, Steinhauser '06
Chetyrkin, Kühn, Sturm '06

$$\sigma_{\text{HEFT}}^{\text{HO}} = \left(\frac{\sigma^{\text{HO}}}{\sigma^{\text{LO}}} \right)_{M_t \rightarrow \infty} \sigma^{\text{LO}}$$



Computation

Ingredients for exact quark mass effects at NNLO

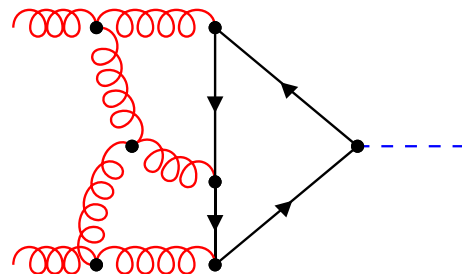
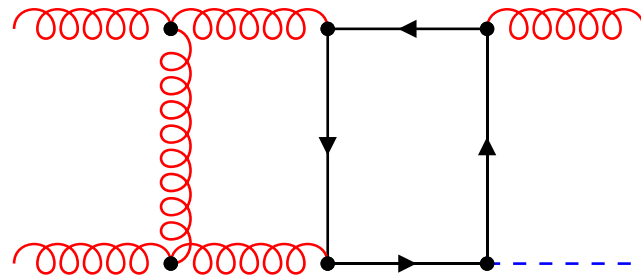


Analytically: [Del Duca, Kilgore, Oleari, et al. '01](#)

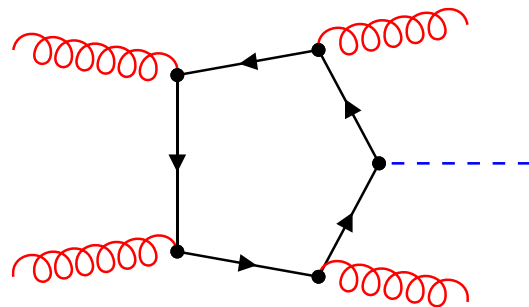
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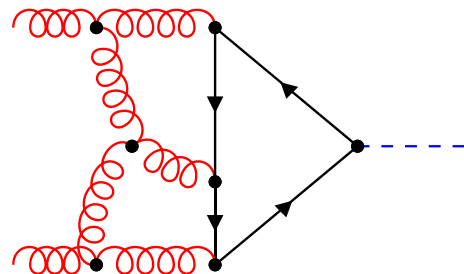
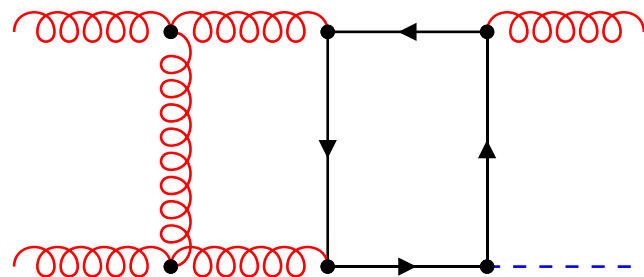
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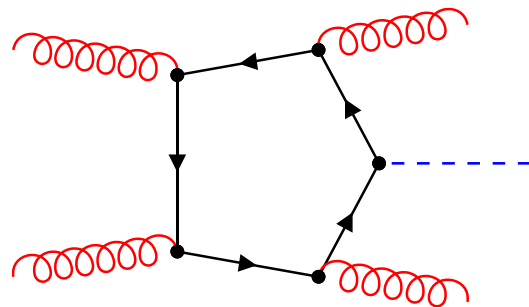
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Double-virtual corrections

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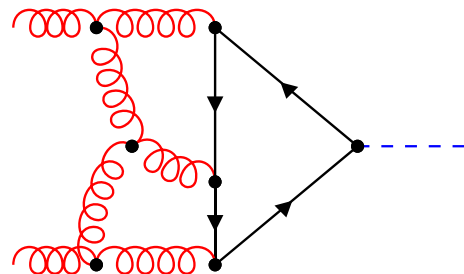
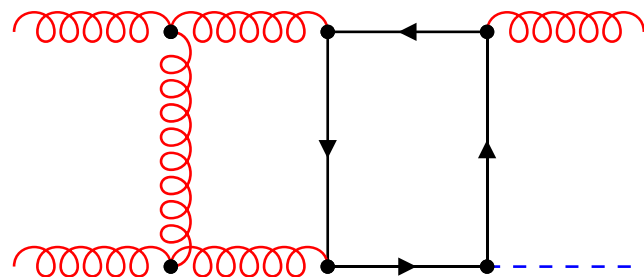


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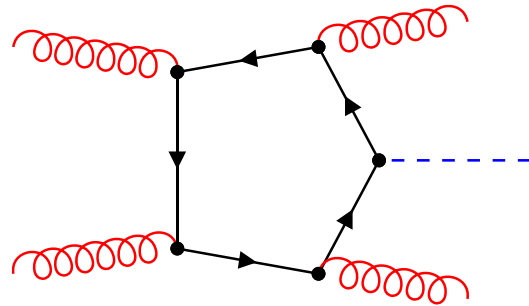


Double-virtual corrections

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[MN, Usovitsch '23](#)

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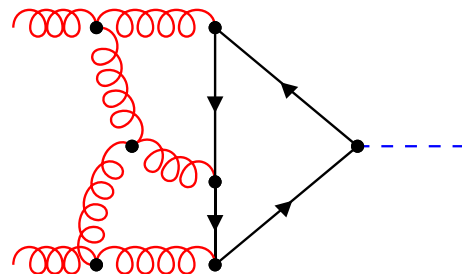
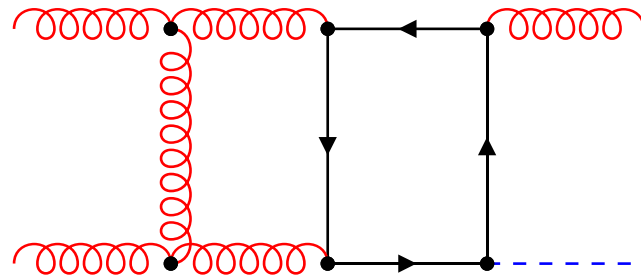


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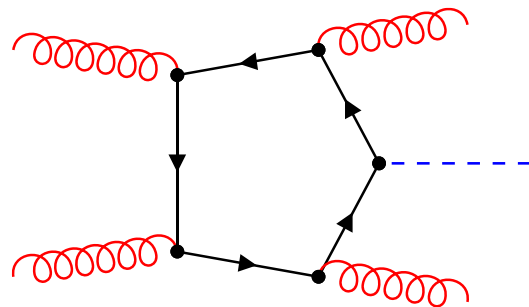
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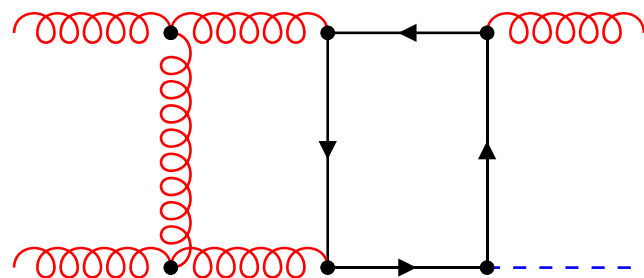


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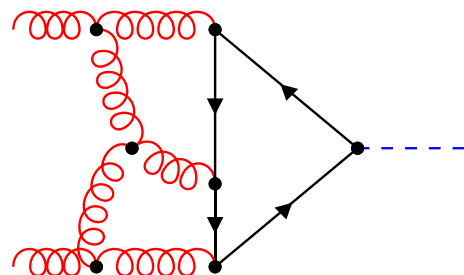
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Real-virtual corrections



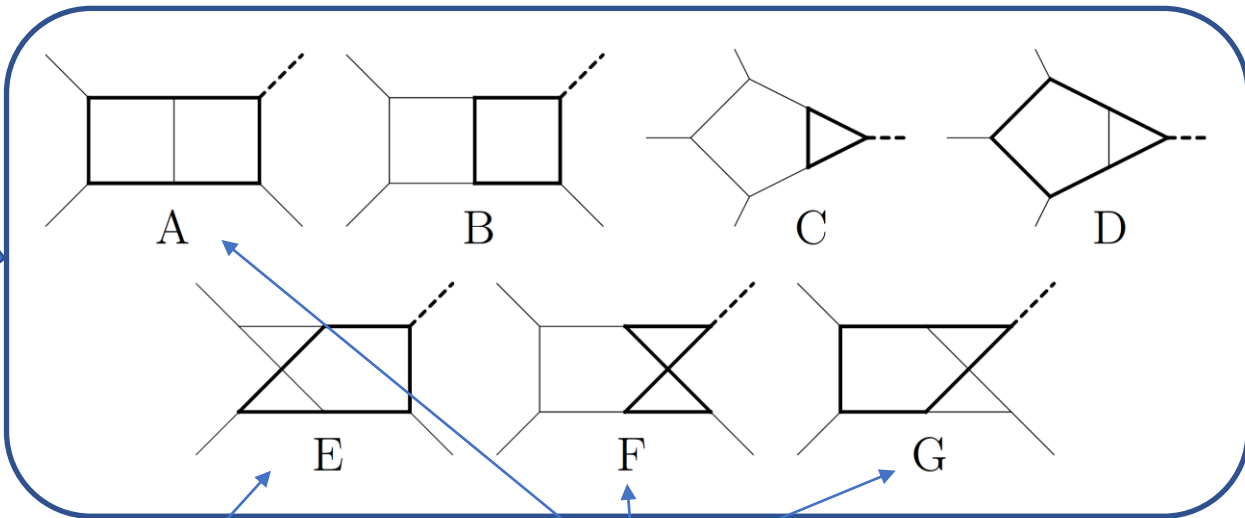
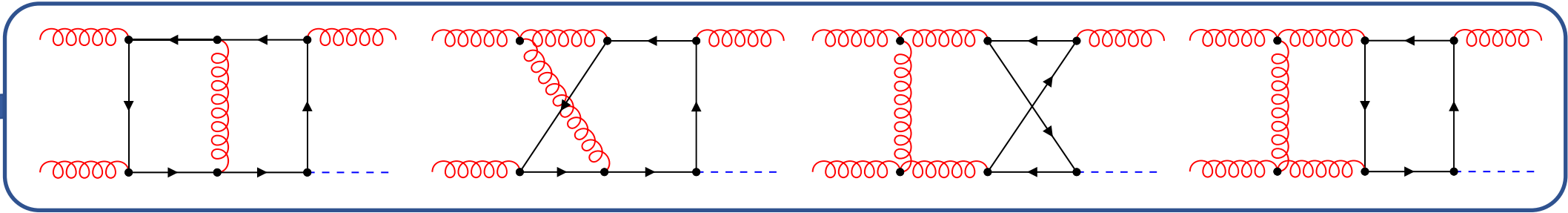
Double-virtual corrections

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Real-virtual corrections



vanishing color factor

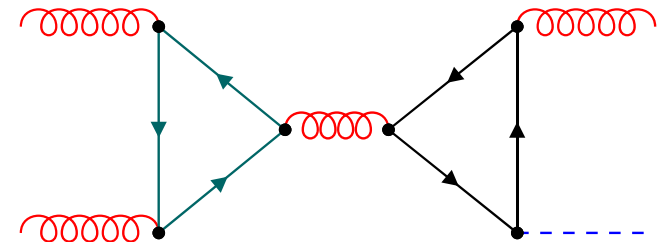
Elliptic sector

A,B,C,D: [Bonciani, Del Duca, Frellesvig, et al. '16](#)

F: [Bonciani, Del Duca, Frellesvig, et al. '19](#)

G: [Frellesvig, Hidding, Maestri, et al. '19](#)

Contributions with two closed fermion chains are always factorizable:



Parametrization

- Variables: \hat{s} , \hat{t} , \hat{u} , m_H^2 , m_q^2
- Introduce dimensionless variables and fix ratio m_q^2/m_H^2
 - z parametrizes **soft** limit
 - λ parametrizes **collinear** limit

$$\hat{t}/\hat{s} = z \lambda$$

$$\hat{u}/\hat{s} = z (1-\lambda)$$

$$z = 1 - m_H^2/\hat{s}$$

$$\lambda = \hat{t}/(\hat{t} + \hat{u})$$

$$z = 1 - m_H^2/\hat{s}$$

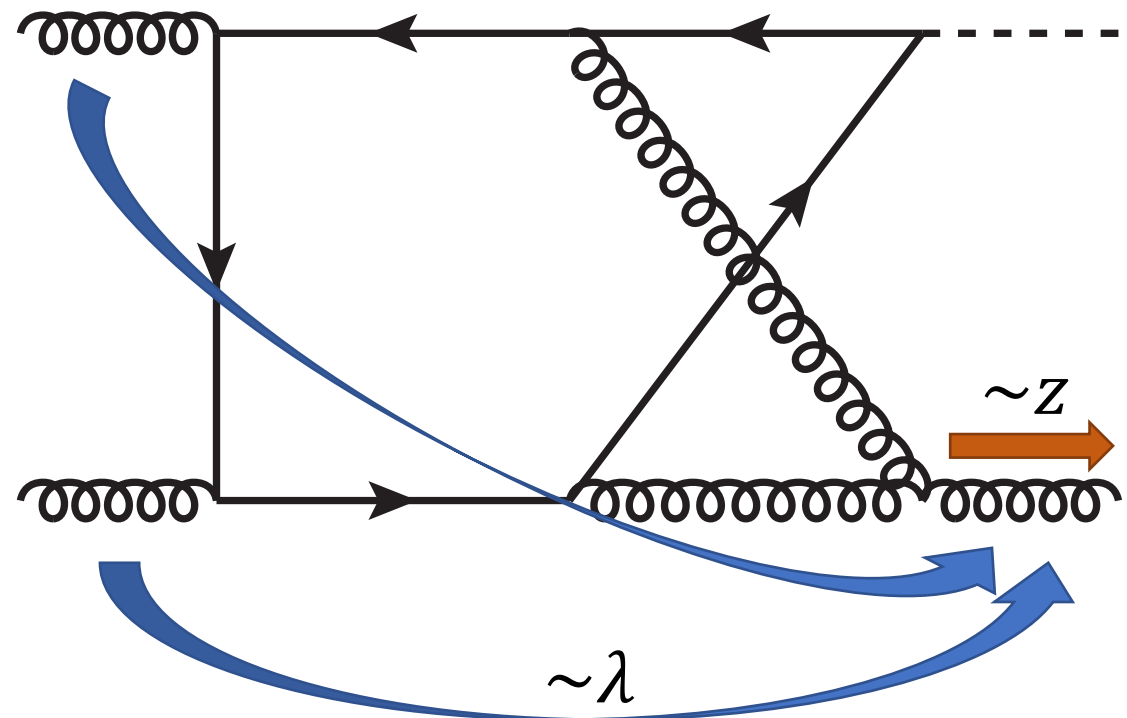
$$\lambda = \hat{t}/(\hat{t} + \hat{u})$$

$$m_t^2/m_H^2 = 23/12$$

$$m_b^2/m_H^2 = 1/684$$

Range of parameters:

- $\lambda \in (0,1)$
- $z \in (0,1)$

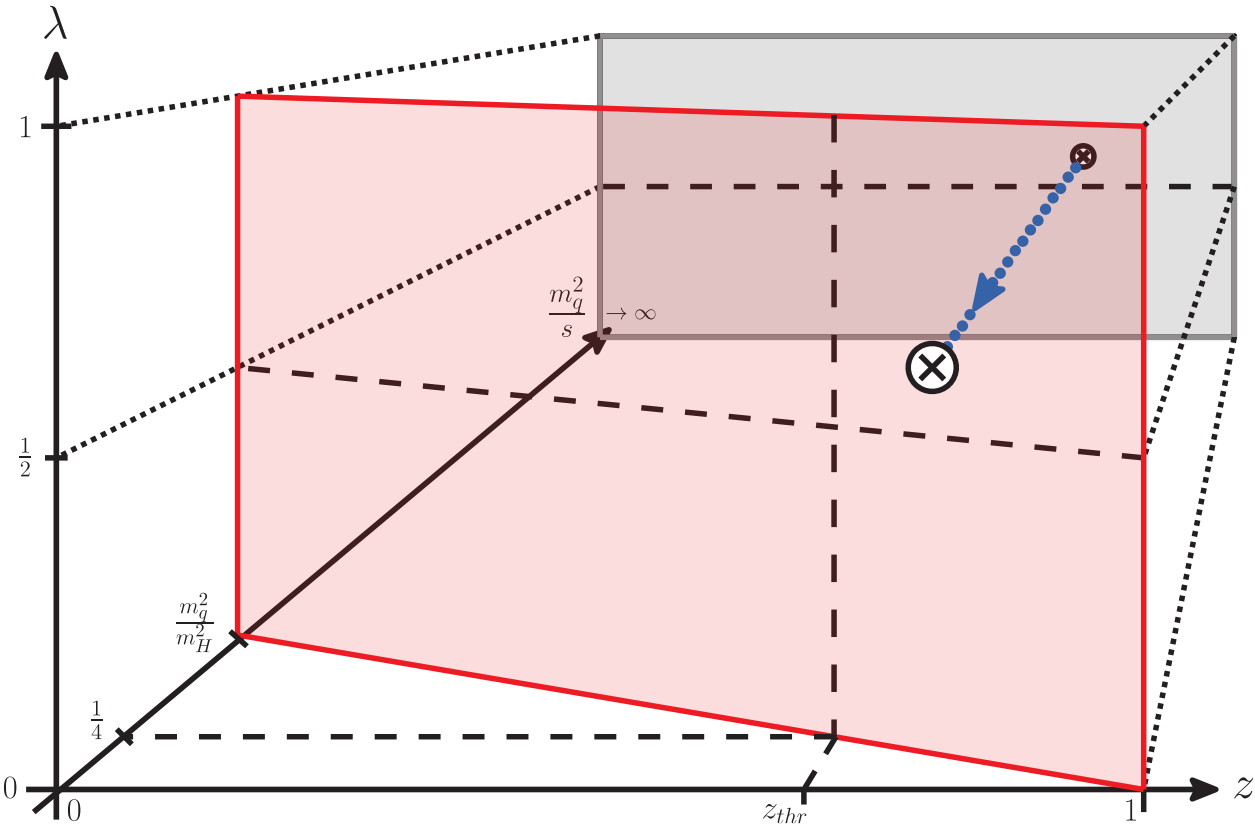


Numerical evolutions

$$z = 1 - m_H^2 / \hat{s}$$

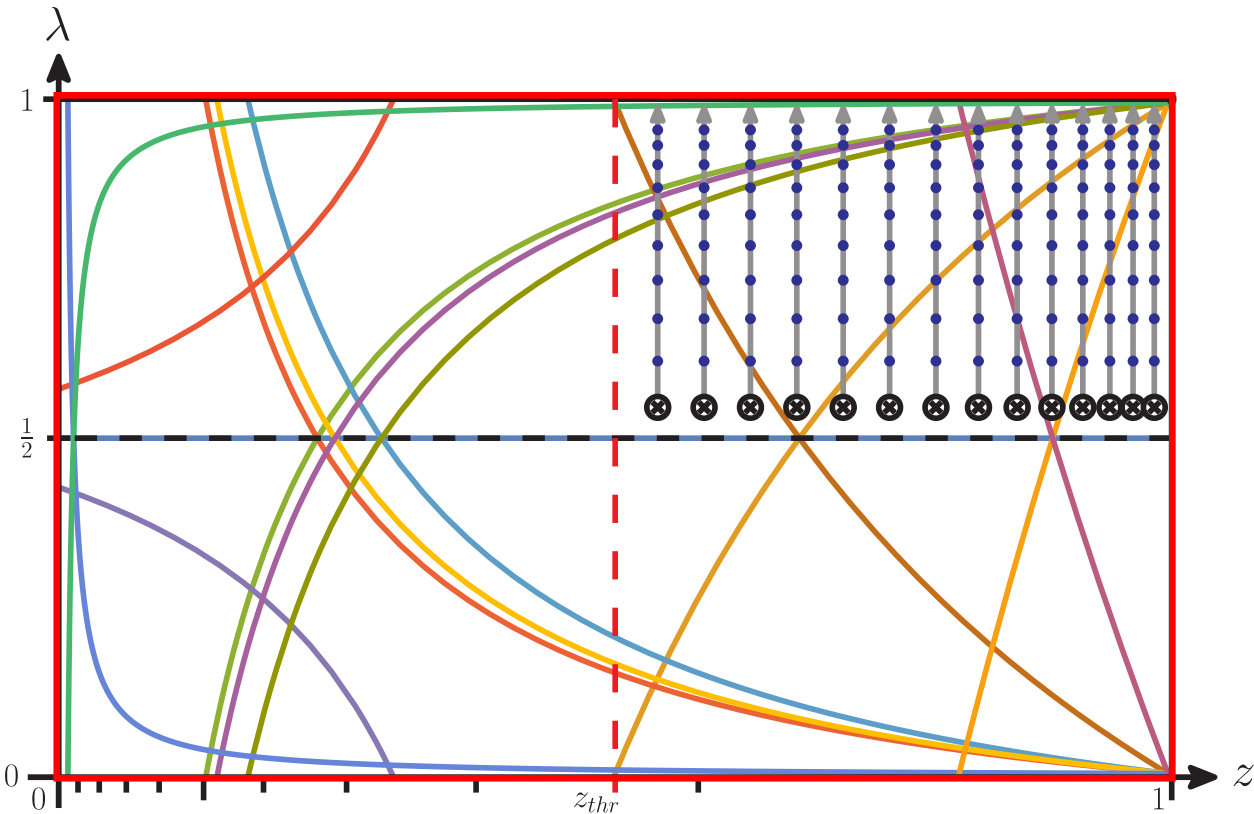
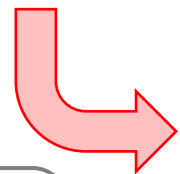
$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$m_t^2 / m_H^2 = 23/12$$



- Initial condition given in the limit $m_q^2 \rightarrow \infty$
- Transport boundary from $m_q^2 \rightarrow \infty$ to physical plane

- Collect numerical samples for MIs along straight contours in λ for different values of z



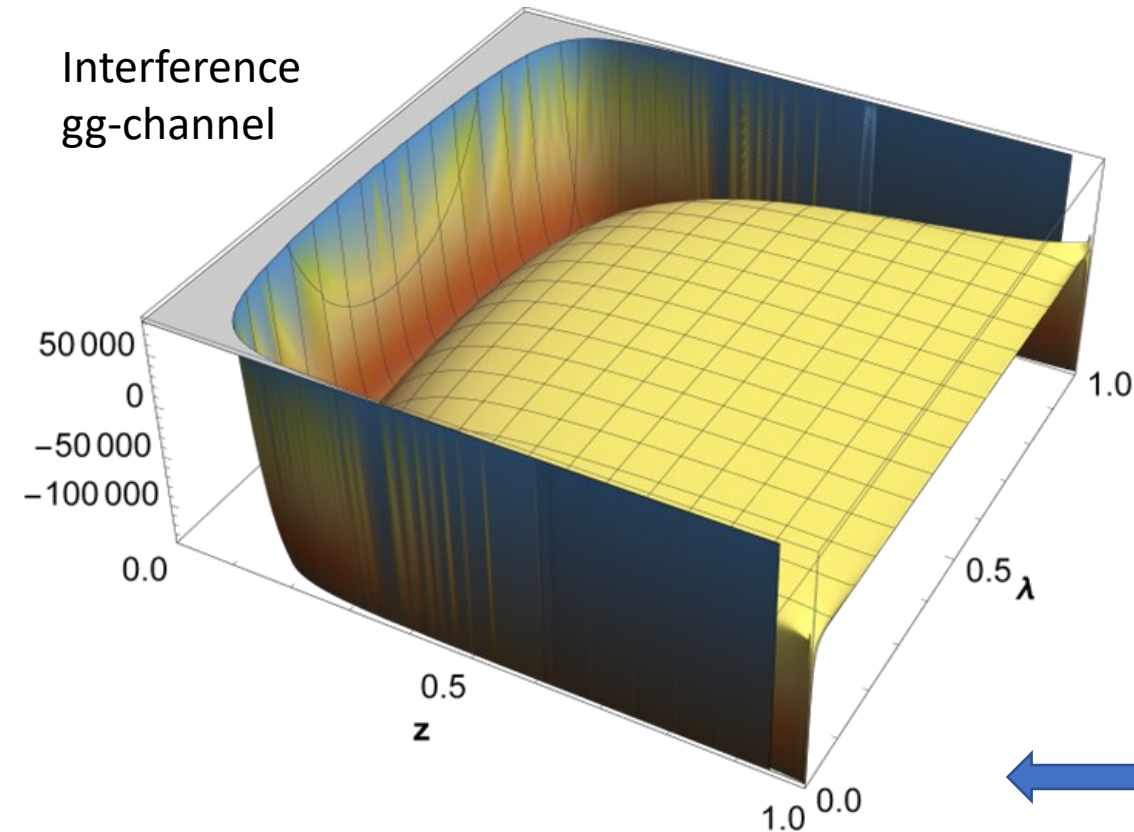
Construction of amplitudes

- Collected 2×10^6 numerical samples for MIs at m_t^2/m_H^2 by evaluation of the LME and numerical evolution above threshold
- Collected 1×10^6 numerical samples for MIs at m_b^2/m_H^2 via numerical evolution in the entire phase space

Insert into form factors and construct helicity amplitudes



Interference
gg-channel



$$\mathcal{M}_{+++} = \frac{1}{\sqrt{2}} \frac{s_{12}s_{23}}{\langle 12 \rangle \langle 23 \rangle \langle 31 \rangle} \left(F_1 + \frac{s_{13}}{s_{23}} F_2 + \frac{s_{13}}{s_{12}} F_3 + \frac{s_{13}}{2} F_4 \right)$$

$$\mathcal{M}_{---} = \mathcal{M}_{+++} \Big|_{\langle ij \rangle \leftrightarrow [ji]} = -\frac{1}{\sqrt{2}} \frac{s_{12}s_{23}}{[12][23][31]} \left(F_1 + \frac{s_{13}}{s_{23}} F_2 + \frac{s_{13}}{s_{12}} F_3 + \frac{s_{13}}{2} F_4 \right)$$

$$\mathcal{M}_{++-} = \frac{1}{\sqrt{2}} \frac{[12]^3}{[13][23]} \frac{s_{23}}{s_{12}} \left(F_1 + \frac{s_{13}}{2} F_4 \right)$$

$$\mathcal{M}_{--+} = \mathcal{M}_{++-} \Big|_{\langle ij \rangle \leftrightarrow [ji]}$$

$$\mathcal{M}_{+-+} = \frac{1}{\sqrt{2}} \frac{[13]^3}{[12][23]} \frac{s_{12}}{s_{13}} \left(F_2 + \frac{s_{23}}{2} F_4 \right)$$

$$\mathcal{M}_{-+-} = \mathcal{M}_{+-+} \Big|_{\langle ij \rangle \leftrightarrow [ji]}$$

$$\mathcal{M}_{-++} = \frac{1}{\sqrt{2}} \frac{[23]^3}{[12][13]} \frac{s_{13}}{s_{23}} \left(F_3 + \frac{s_{12}}{2} F_4 \right)$$

$$\mathcal{M}_{+--} = \mathcal{M}_{-++} \Big|_{\langle ij \rangle \leftrightarrow [ji]}$$

$$|\mathcal{M}|^2 = \sum_{h \in Hel.} |\mathcal{M}_h|^2$$

Subtraction of IR divergences for $gg \rightarrow gH$

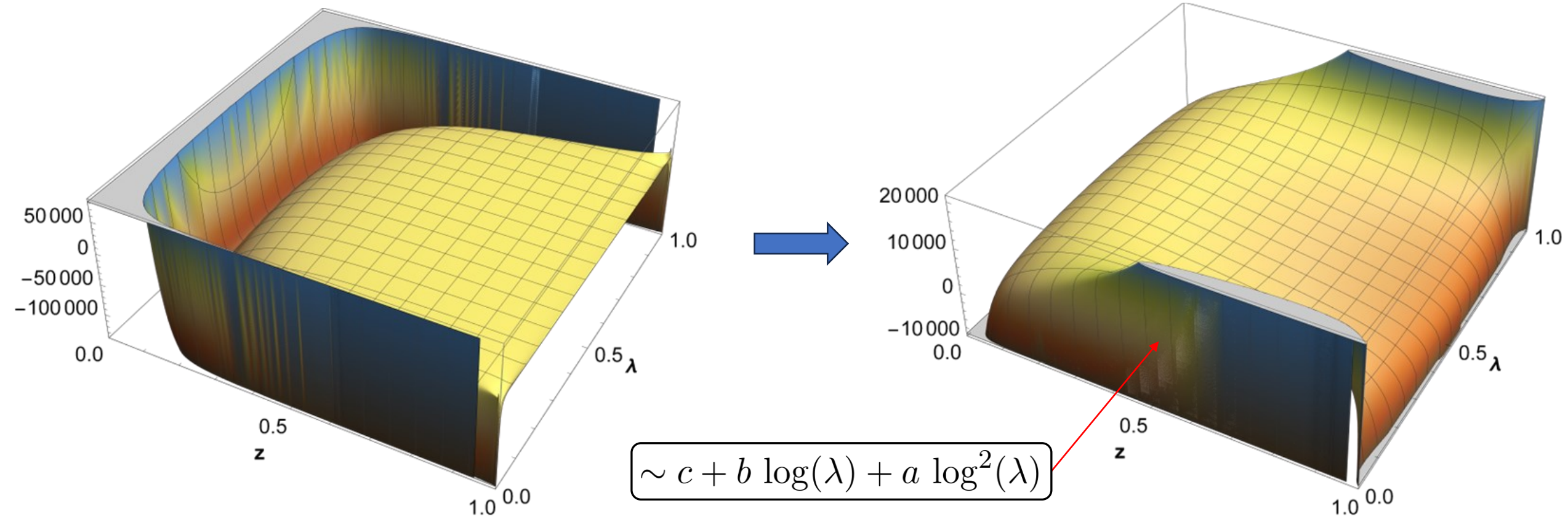
$$z = 1 - m_H^2 / \hat{s}$$

$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$\langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(2)} \rangle_{\text{regulated}} \equiv \langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(2)} \rangle - \left[\langle M_{\text{HEFT}}^{(1)} | M_{\text{HEFT}}^{(2)} \rangle + \frac{8\pi\alpha_s}{\hat{t}} \left\langle P_{gg}^{(0)} \left(\frac{\hat{s}}{\hat{s} + \hat{u}} \right) \right\rangle \langle F^{(1)} | (F_{\text{exact}}^{(2)} - F_{\text{HEFT}}^{(2)}) \rangle \right]$$

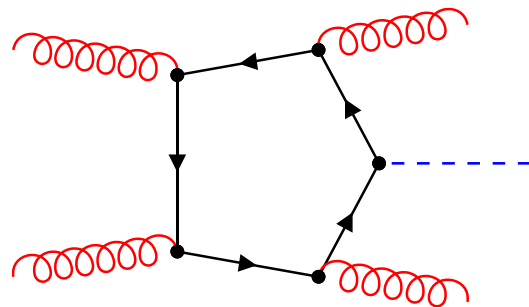
- From unregulated to regulated quantity at $\mu_R = m_H/2$:

$$r = \frac{\sigma_{t \times b}^{\text{LO}}}{\sigma_{\text{HTL}}^{\text{LO}}} \approx -0.129$$



$$\sim c + b \log(\lambda) + a \log^2(\lambda)$$

Ingredients for exact quark mass effects at NNLO

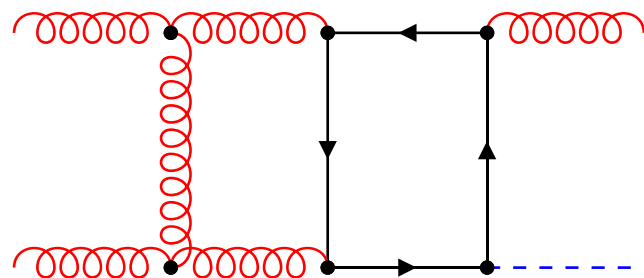


Analytically: [Del Duca, Kilgore, Oleari, et al. '01](#)

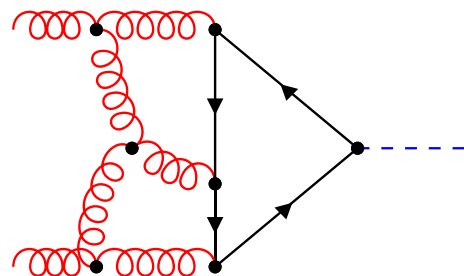
OpenLoops 2: [Buccioni, Lang, Lindert, et al. '19](#)

Analytically (more compact and

implemented in MCFM): [Budge, Campbell, De Laurentis, et al. '20](#)



Real-virtual corrections



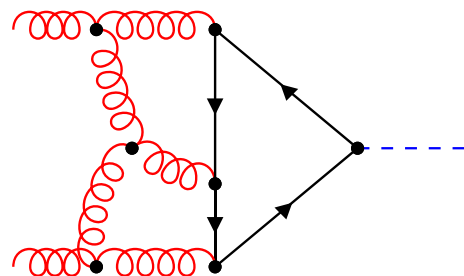
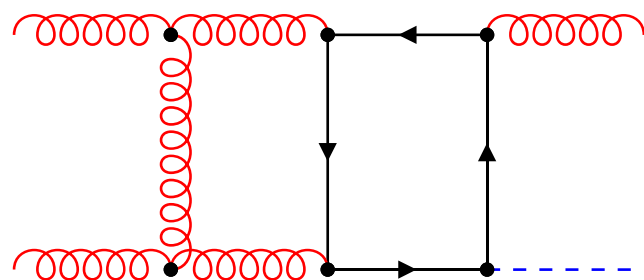
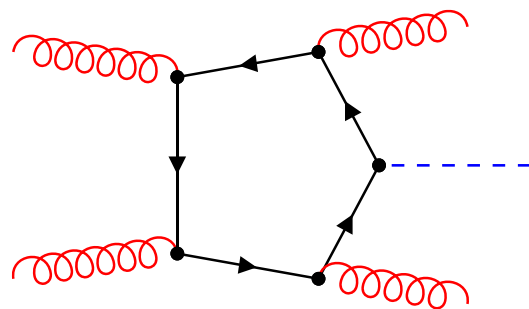
Double-virtual corrections

[Czakon, MN '20](#)

[MN, Usovitsch '23](#)



Ingredients for exact quark mass effects at NNLO



Analytically: Del Duca, Kilgore, Oleari, et al. '01

OpenLoops 2: Buccioni, Lang, Lindert, et al. '19

Analytically (more compact and

implemented in MCFM): Budge, Campbell, De Laurentis, et al. '20

Phase-space integration
with sector-improved residue
subtraction scheme
(**Stripper** Czakon '10)

Double-virtual corrections

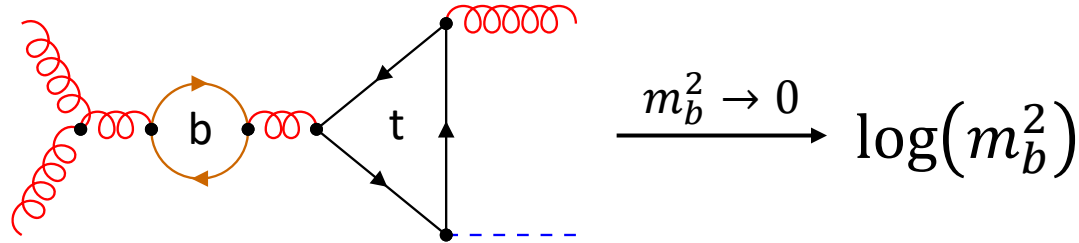
Czakon, MN '20

MN, Usovitsch '23

Results

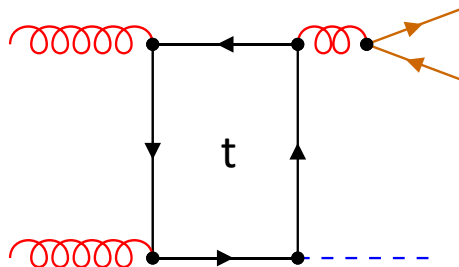
Note on the flavor scheme

- Subsets of diagrams in real-virtual and double virtual contribution give rise logarithmic mass divergences



4-flavor scheme

- Consistent treatment of massive t- and b-quarks
- Exclude b-quark from initial state
- Include massive b-quark splittings in final state



5-flavor scheme

- Treat b-quark as massless particle
- Massive b-quark only present in loops directly attached to the Higgs-boson
- Corresponds to theory with a replica b-quark carrying the mass of a heavy b-quark

Results

- Effects of interference of top- and bottom-quark amplitudes on Higgs production in gluon-fusion at the LHC

- PDF set: **NNPDF31_nnlo_as_0118** [NNPDF Collaboration `17](#)

- $\mu_R = \mu_F = m_H/2$ (central scale)

- $m_H = 125$ GeV $\Rightarrow m_t \approx 173.055$ GeV and $m_b \approx 4.779$ GeV (*both* in OS-scheme)

[Czakon, Eschment, MN, Poncelet, Schellenberger `23](#)

- HEFT values obtained with **SusHi** [Harlander, Liebler, Mantler `16](#)

Order	σ_{HEFT} [pb]	$(\sigma_t - \sigma_{\text{HEFT}})$ [pb]	$\sigma_{t \times b}$ [pb]	$\sigma_{t \times b}/\sigma_{\text{HEFT}}$ [%]
$\sqrt{s} = 8$ TeV				
$\mathcal{O}(\alpha_s^2)$	+7.39	–	–0.895	
LO	$7.39^{+1.98}_{-1.40}$	–	$-0.895^{+0.17}_{-0.24}$	–12
$\mathcal{O}(\alpha_s^3)$	+9.14	–0.0873	–0.268	
NLO	$16.53^{+3.63}_{-2.73}$	$-0.0873^{+0.030}_{-0.052}$	$-1.163^{+0.10}_{-0.08}$	$-7.0^{+1.0}_{-0.8}$
$\mathcal{O}(\alpha_s^4)$	+4.19	+0.0523(2)	+0.167(3)	
NNLO	$20.72^{+1.84}_{-2.06}$	$-0.0350(2)^{+0.048}_{-0.013}$	$-0.996(3)^{+0.12}_{-0.05}$	$-4.8^{+0.9}_{-0.8}$
$\sqrt{s} = 13$ TeV				
$\mathcal{O}(\alpha_s^2)$	+16.30	–	–1.975	
LO	$16.30^{+4.36}_{-3.10}$	–	$-1.98^{+0.38}_{-0.53}$	–12
$\mathcal{O}(\alpha_s^3)$	+21.14	–0.303	–0.446(1)	
NLO	$37.44^{+8.42}_{-6.29}$	$-0.303^{+0.10}_{-0.17}$	$-2.42^{+0.19}_{-0.12}$	$-6.5^{+0.9}_{-0.8}$
$\mathcal{O}(\alpha_s^4)$	+9.72	+0.147(1)	+0.434(8)	
NNLO	$47.16^{+4.21}_{-4.77}$	$-0.156(1)^{+0.13}_{-0.03}$	$-1.99(1)^{+0.30}_{-0.15}$	$-4.2^{+0.9}_{-0.8}$

Results

- Effects of interference of top- and bottom-quark amplitudes on Higgs production in gluon-fusion at the LHC
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Czakon, Eschment, MN, Poncelet,
Schellenberger `23

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➤ Interference effects much larger than pure top mass effect

[Czakon, Eschment, MN, Poncelet, Schellenberger `23](#)

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- Interference effects much larger than pure top mass effect
- Interference effect at NNLO cancels against NLO

[Czakon, Eschment, MN, Poncelet, Schellenberger `23](#)

Results

- Effects of interference of top- and bottom-quark amplitudes on Higgs production in gluon-fusion at the LHC
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- Interference effect at NNLO larger than NLO scale variation (similar in HEFT but less severe)

Czakon, Eschment, MN, Poncelet,
Schellenberger `23

Results

- Effects of interference of top- and bottom-quark amplitudes on Higgs production in gluon-fusion at the LHC
 - PDF set: **NNPDF31_nnlo_as_0118** [NNPDF Collaboration `17](#)
 - $\mu_R = \mu_F = m_H/2$ (central scale)
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[Czakon, Eschment, MN, Poncelet, Schellenberger `23](#)

Results

- Effects of interference of top- and bottom-quark amplitudes on Higgs production in gluon-fusion at the LHC
 - PDF set: **NNPDF31_nnlo_as_0118** [NNPDF Collaboration `17](#)
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- Interference NNLO scale variation increases compared to NLO
- Similar effects for different top quark mass ($m_t \approx 170.979$ GeV)

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Schellenberger `23

Results

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 - PDF set: **NNPDF31_nnlo_as_0118** [NNPDF Collaboration `17](#)
 - $\mu_R = \mu_F = m_H/2$ (central scale)
 - $m_H = 125$ GeV $\Rightarrow m_t \approx 173.055$ GeV and $m_b \approx 4.779$ GeV (OS-scheme, but Y_b in \overline{MS} with $\bar{m}_b(\bar{m}_b) \approx 4.18$ GeV)
 - HEFT values obtained with **SusHi** [Harlander, Liebler, Mantler `16](#)

Order	σ_{HEFT} [pb]	$(\sigma_t - \sigma_{\text{HEFT}})$ [pb]	$\sigma_{t \times b}$ [pb]	$\sigma_{t \times b}(Y_{b, \overline{MS}})$ [pb]
$\sqrt{s} = 13$ TeV				
$\mathcal{O}(\alpha_s^2)$	+16.30	–	–1.975	–1.223
LO	$16.30^{+4.36}_{-3.10}$	–	$-1.98^{+0.38}_{-0.53}$	$-1.22^{+0.29}_{-0.44}$
$\mathcal{O}(\alpha_s^3)$	+21.14	–0.303	–0.446(1)	–0.623(1)
NLO	$37.44^{+8.42}_{-6.29}$	$-0.303^{+0.10}_{-0.17}$	$-2.42^{+0.19}_{-0.12}$	$-1.85^{+0.26}_{-0.26}$
$\mathcal{O}(\alpha_s^4)$	+9.72	+0.147(1)	+0.434(8)	+0.019(5)
NNLO	$47.16^{+4.21}_{-4.77}$	$-0.156(1)^{+0.13}_{-0.03}$	$-1.99(1)^{+0.30}_{-0.15}$	$-1.83(1)^{+0.08}_{-0.03}$

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- Improved convergence in mixed renormalization scheme compared to OS-scheme

[Czakon, Eschment, MN, Poncelet, Schellenberger `23](#)

Results

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 - $\mu_R = \mu_F = m_H/2$ (central scale)
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$\sqrt{s} = 13$ TeV					
$\mathcal{O}(\alpha_s^2)$	+16.30	–	–1.975	–1.223	–1.118
LO	$16.30^{+4.36}_{-3.10}$	–	$-1.98^{+0.38}_{-0.53}$	$-1.22^{+0.29}_{-0.44}$	$-1.118^{+0.28}_{-0.43}$
$\mathcal{O}(\alpha_s^3)$	+21.14	–0.303	–0.446(1)	–0.623(1)	–0.647
NLO	$37.44^{+8.42}_{-6.29}$	$-0.303^{+0.10}_{-0.17}$	$-2.42^{+0.19}_{-0.12}$	$-1.85^{+0.26}_{-0.26}$	$-1.76^{+0.27}_{-0.28}$
$\mathcal{O}(\alpha_s^4)$	+9.72	+0.147(1)	+0.434(8)	+0.019(5)	+0.02(1)
NNLO	$47.16^{+4.21}_{-4.77}$	$-0.156(1)^{+0.13}_{-0.03}$	$-1.99(1)^{+0.30}_{-0.15}$	$-1.83(1)^{+0.08}_{-0.03}$	$-1.74(2)^{+0.13}_{-0.01}$

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- Improved convergence in mixed renormalization scheme compared to OS-scheme
- Similar pattern of corrections for m_b in \overline{MS} -scheme


preliminary

Summary and outlook


- The top-bottom interference contribution to the total Higgs production cross section was computed with *both* quarks renormalized in the OS-scheme
 - $\mathcal{O}(\alpha_s^4)$ correction at 8 TeV: +0.167 pb
 - $\mathcal{O}(\alpha_s^4)$ correction at 13 TeV: +0.434 pb

← $\mathcal{O}(1\%)$ effect
- NNLO correction at 13 TeV: $-1.99(1)_{-0.15}^{+0.30}$ pb compatible with previous estimate $-2.18_{-0.20}^{+0.20}$ pb [Anastasiou, Penin `20](#)
- Top-quark and interference contribution not sensitive to small variations of the top-quark mass
- Interference shows signs of poor perturbative convergence
 - Better convergence in \overline{MS} -scheme for the bottom-quark mass or Yukawa coupling only
- Cross checks: at the differential level
 - [Jones, Kerner, Luisoni `18](#)
 - [Caola, Lindert, Melnikov, et al. `18](#)
 - [Bonciani, Del Duca, Frellesvig, et al. `22](#)
- Next steps:
 - Complete calculation with quark masses renormalized in \overline{MS} -scheme
 - Consistent treatment of massive quarks in 4-flavor scheme
 - Top-charm interference contribution

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 - Complete calculation with quark masses renormalized in \overline{MS} -scheme 
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 - Top-charm interference contribution

preliminary

Order	$\sigma_{t \times b}$ [pb]	
	$\sqrt{s} = 13$ TeV	
	5FS	4FS
$\mathcal{O}(\alpha_s^2)$	-1.975	-1.971
LO	$-1.98_{-0.53}^{+0.37}$	$-1.97_{-0.56}^{+0.39}$
$\mathcal{O}(\alpha_s^3)$	-0.447(4)	-0.455(4)
NLO	$-2.42_{-0.12}^{+0.19}$	$-2.43_{-0.13}^{+0.21}$
$\mathcal{O}(\alpha_s^4)$	+0.434(8)	+0.389(11)
NNLO	$-1.99(1)_{-0.15}^{+0.30}$	$-2.04(1)_{-0.14}^{+0.29}$