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Light scalars at e^+e^- colliders

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zoom, 04/2023

1. Motivation

2. Interpretation

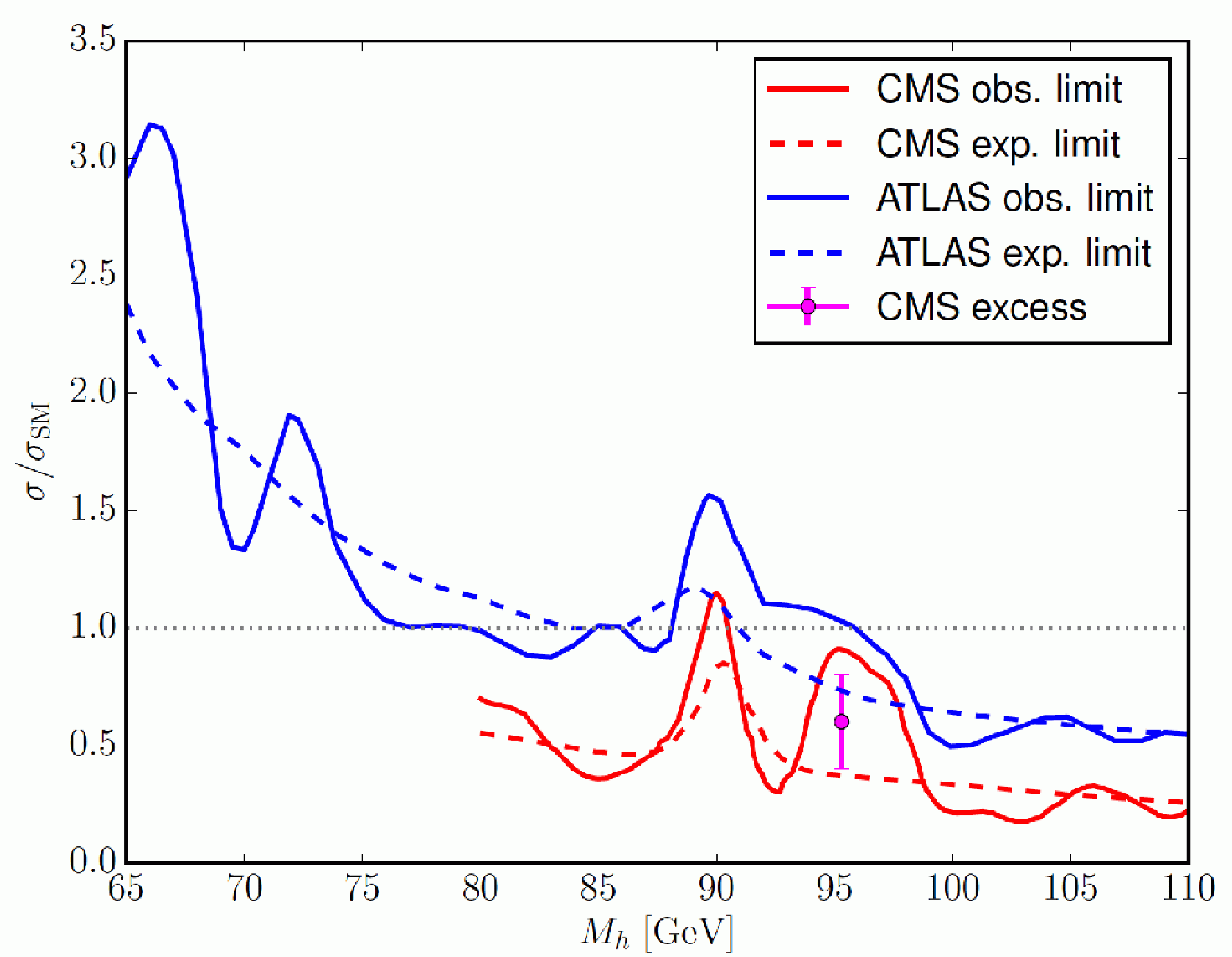
3. Physics opportunities at e^+e^- colliders

4. Conclusions

1. Motivation

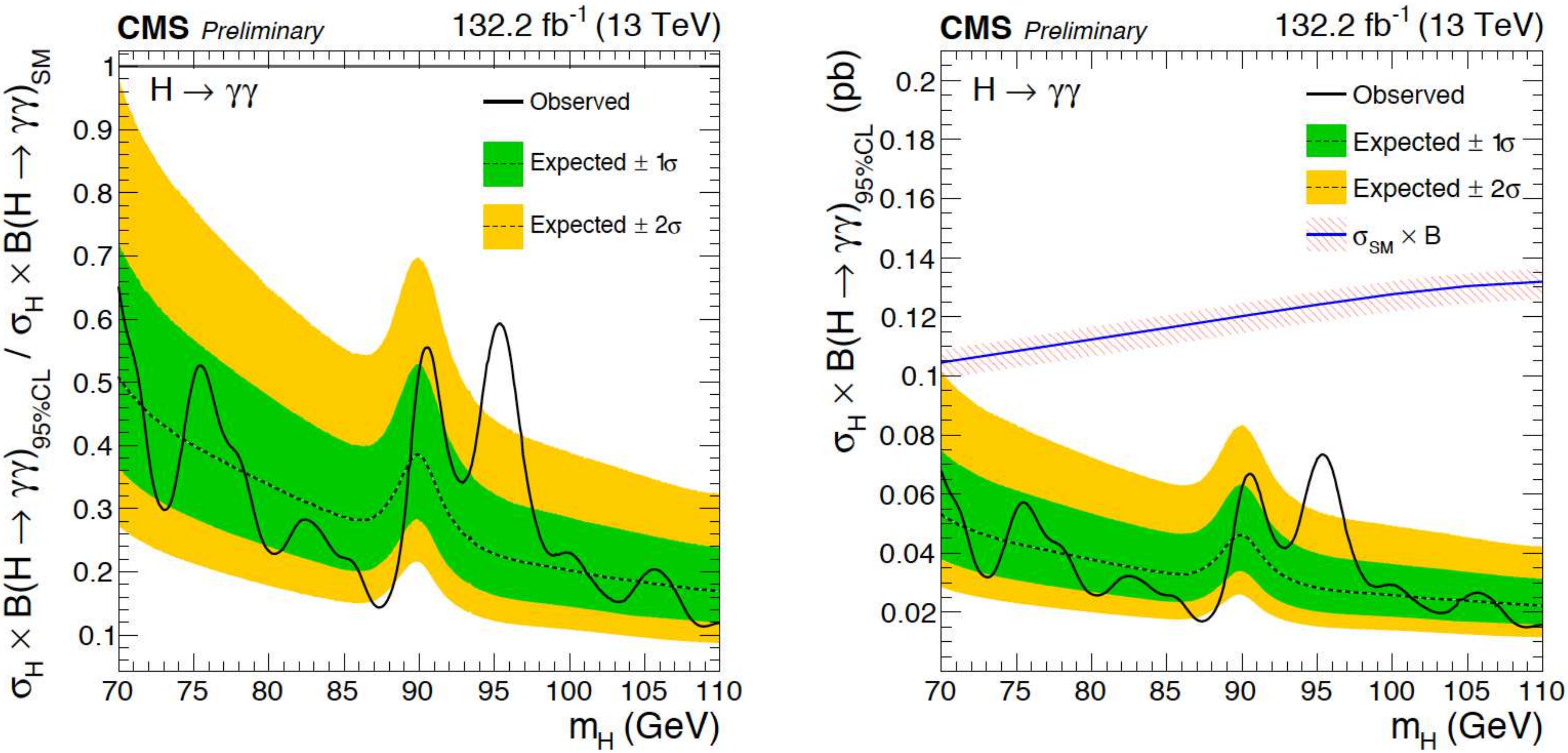
Case study: Search for $pp \to \phi \to \gamma\gamma$: excess at $m_\phi \sim 95$ GeV

[CMS '17, ATLAS '18, S.H., T. Stefaniak '18] $\mu_{\rm CMS} = 0.6 \pm 0.2$



$\Rightarrow$ if there is something, it would look exactly like this!

NEW: Full Run 2 results from CMS: [CMS '23]



$$\mu_{\gamma\gamma} = [\sigma(gg \to h_{95}) \times \mathrm{BR}(h_{95} \to \gamma\gamma)]_{\mathrm{exp/SM}} = 0.33^{+0.19}_{-0.12}$$

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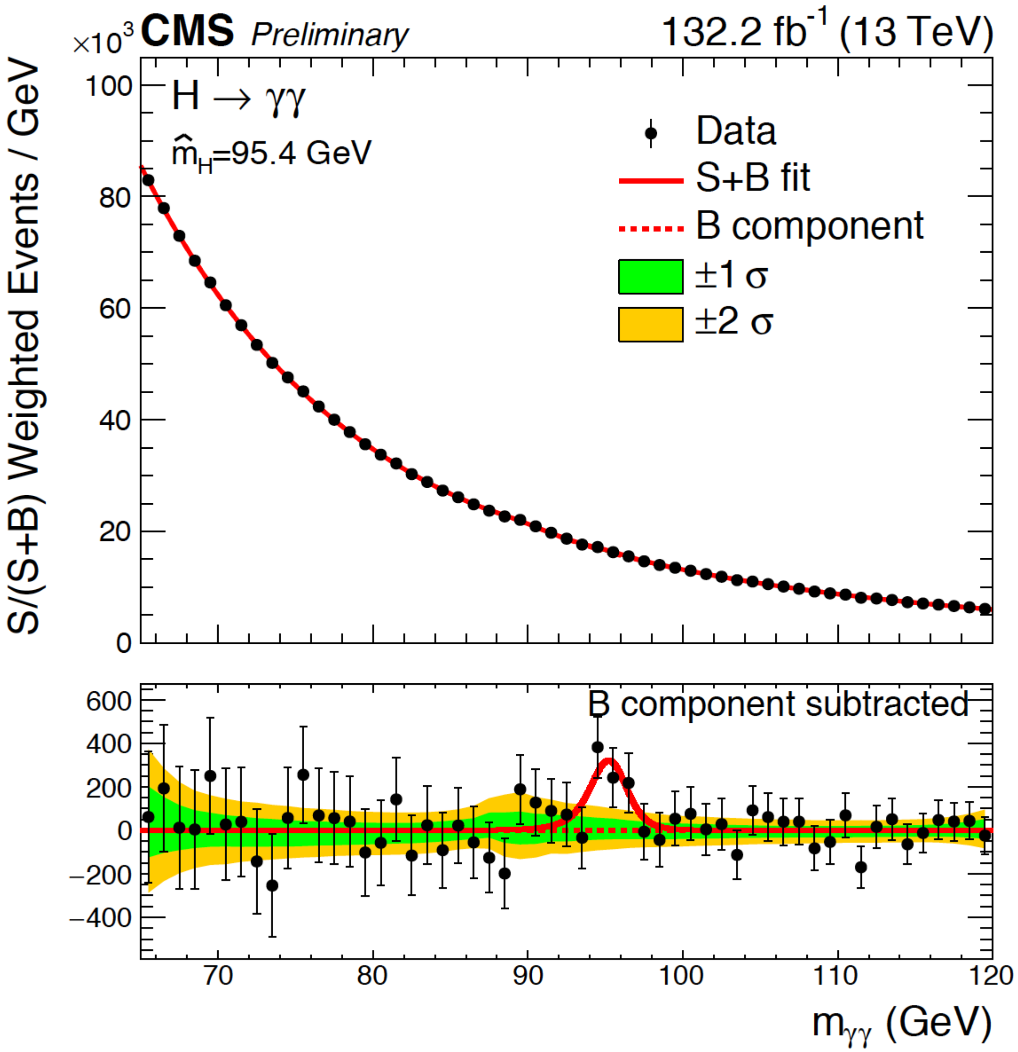
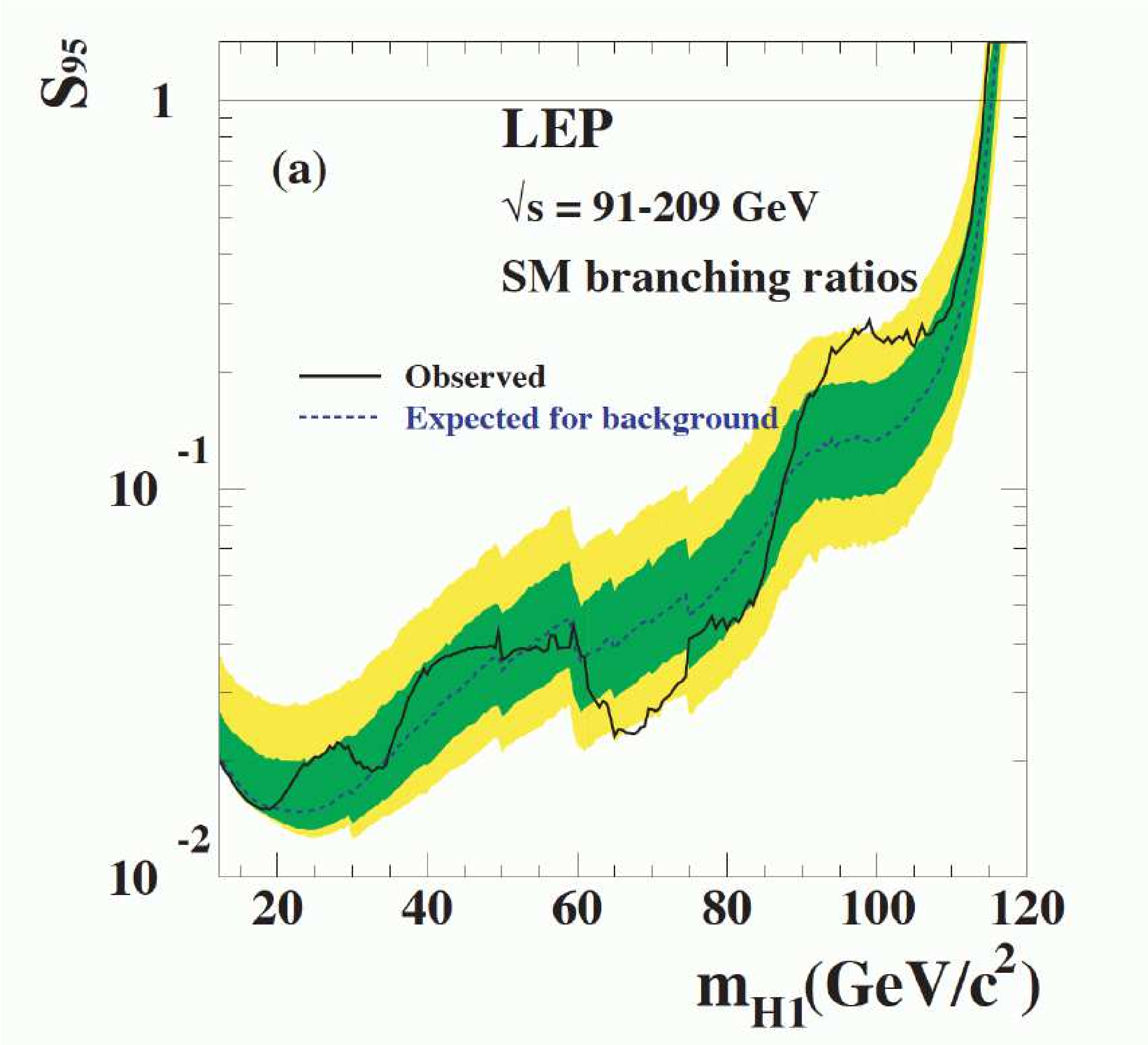


$$\mu_{\gamma\gamma} = [\sigma(gg \to h_{95}) \times \mathrm{BR}(h_{95} \to \gamma\gamma)]_{\mathrm{exp/SM}} = 0.33^{+0.19}_{-0.12}$$

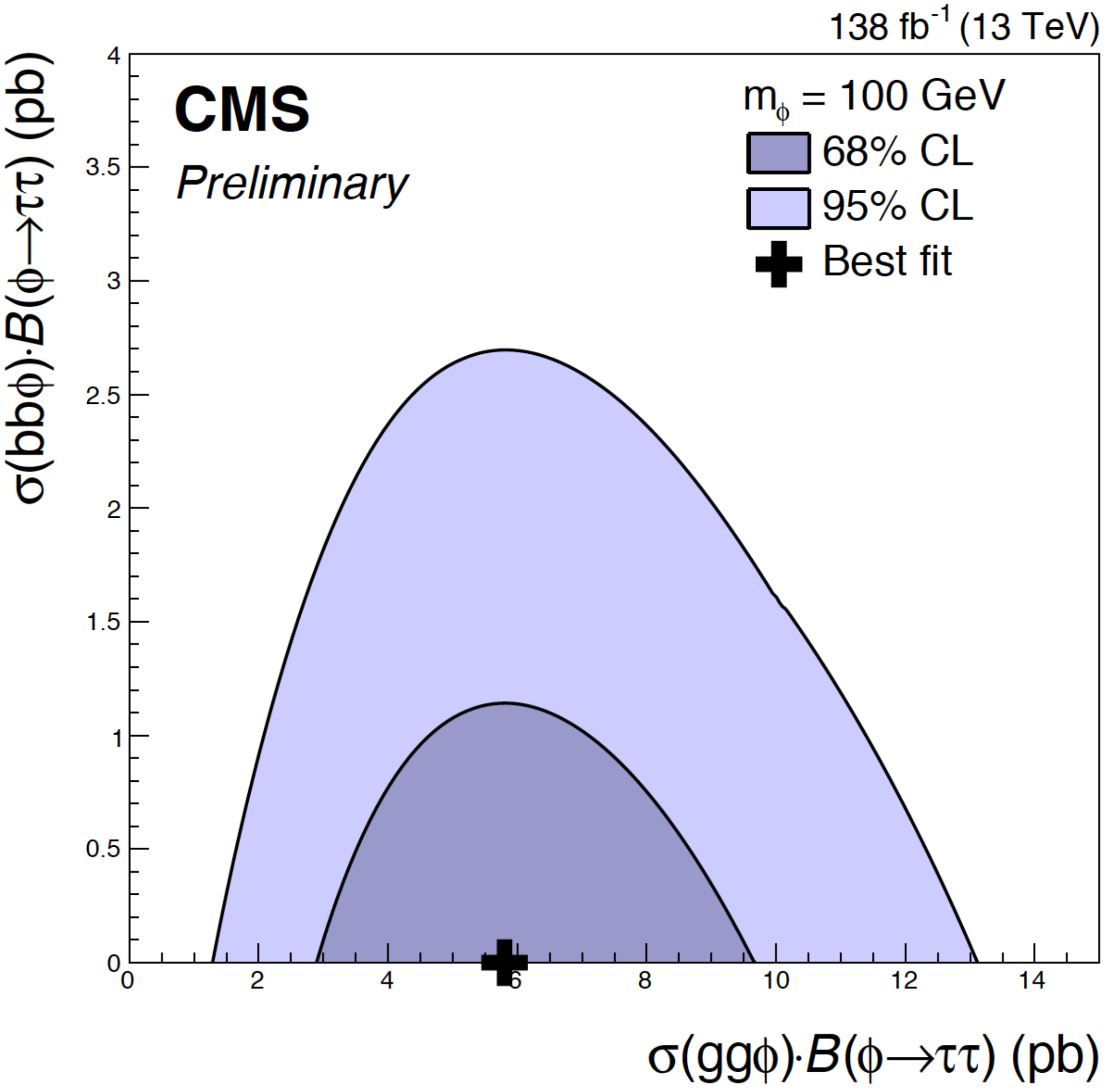
Remember the LEP excess?



$$\mu_{bb}(98\text{ GeV}) = \left[\sigma(e^+e^- \to Zh) \times \text{BR}(h \to b\bar{b})\right]_{\text{exp/SM}} = 0.117 \pm 0.057$$

The new $\tau^+\tau^-$ excess

[CMS '22]



$$\mu_{\tau\tau} = [\sigma(gg \to h_{95}) \times \mathrm{BR}(h_{95} \to \tau\tau)]_{\mathrm{exp/SM}} = 1.2 \pm 0.5$$

Now we have three excesses at ~ 95 GeV

$$\mu_{bb}^{\text{exp}} = 0.117 \pm 0.057, \quad \mu_{\gamma\gamma}^{\text{exp}} = 0.35 \pm 0.12, \quad \mu_{\tau\tau}^{\text{exp}} = 1.2 \pm 0.5$$

corresponding to

$$\mu_{bb}^{\text{exp}} \sim 2\,\sigma, \quad \mu_{\gamma\gamma}^{\text{exp}} \sim 3\,\sigma, \quad \mu_{\tau\tau}^{\text{exp}} \sim 2.4\,\sigma$$

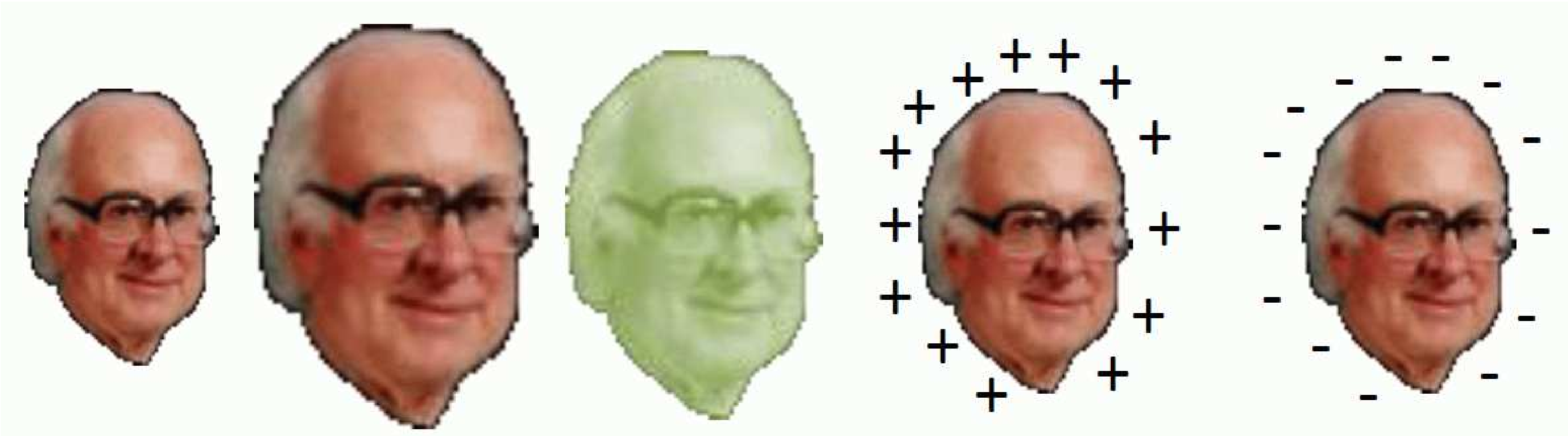
Three (effectively) independent channels
$\Rightarrow$ no LEE (as theorist I am allowed to add naively)

$$\Rightarrow \sim 4.3\,\sigma$$

$$\chi^2_{95} = \frac{(\mu_{bb}^{\text{theo}} - 0.117)^2}{(0.057)^2} + \frac{(\mu_{\gamma\gamma}^{\text{theo}} - 0.35)^2}{(0.12)^2} + \frac{(\mu_{\tau\tau}^{\text{theo}} - 1.2)^2}{(0.5)^2}$$

Can we fit all excesses together?

2. Interpretation



Toy example: **N2HDM** [**S2HDM** similar with complex singlet]

Fields:

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + \rho_1 + i\eta_1) \end{pmatrix} , \; \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \rho_2 + i\eta_2) \end{pmatrix} , \; \Phi_S = v_S + \rho_S$$

Potential:

$$\begin{aligned} V \;=\; & m_{11}^2|\Phi_1|^2 + m_{22}^2|\Phi_2|^2 - m_{12}^2(\Phi_1^\dagger\Phi_2 + h.c.) + \frac{\lambda_1}{2}(\Phi_1^\dagger\Phi_1)^2 + \frac{\lambda_2}{2}(\Phi_2^\dagger\Phi_2)^2 \\ & +\lambda_3(\Phi_1^\dagger\Phi_1)(\Phi_2^\dagger\Phi_2) + \lambda_4(\Phi_1^\dagger\Phi_2)(\Phi_2^\dagger\Phi_1) + \frac{\lambda_5}{2}[(\Phi_1^\dagger\Phi_2)^2 + h.c.] \\ & +\frac{1}{2}m_S^2\Phi_S^2 + \frac{\lambda_6}{8}\Phi_S^4 + \frac{\lambda_7}{2}(\Phi_1^\dagger\Phi_1)\Phi_S^2 + \frac{\lambda_8}{2}(\Phi_2^\dagger\Phi_2)\Phi_S^2 \end{aligned}$$

Z_2 symmetry: $\Phi_1 \to \Phi_1$, $\Phi_2 \to -\Phi_2$, $\Phi_S \to \Phi_S$

Z_2' symmetry: $\Phi_1 \to \Phi_1$, $\Phi_2 \to \Phi_2$, $\Phi_S \to -\Phi_S$ (broken by $v_S \Rightarrow$ no DM)

Physical states: h_1, h_2, h_3 ($\mathcal{CP}$-even), A ($\mathcal{CP}$-odd), $H^\pm$ (charged)

Extension of the Z_2 symmetry to fermions determines four types:

	u-type	d-type	leptons
type I	Φ_2	Φ_2	Φ_2
type II	Φ_2	Φ_1	Φ_1
type III (lepton-specific)	Φ_2	Φ_2	Φ_1
type IV (flipped)	Φ_2	Φ_1	Φ_2

$\Rightarrow$ exactly as in 2HDM

Three neutral $\mathcal{CP}$-even Higgses:

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \rho_1 \\ \rho_2 \\ \rho_S \end{pmatrix} , \quad R = \begin{pmatrix} c_{\alpha_1} c_{\alpha_2} & s_{\alpha_1} c_{\alpha_2} & s_{\alpha_2} \\ -(c_{\alpha_1} s_{\alpha_2} s_{\alpha_3} + s_{\alpha_1} c_{\alpha_3}) & c_{\alpha_1} c_{\alpha_3} - s_{\alpha_1} s_{\alpha_2} s_{\alpha_3} & c_{\alpha_2} s_{\alpha_3} \\ -c_{\alpha_1} s_{\alpha_2} c_{\alpha_3} + s_{\alpha_1} s_{\alpha_3} & -(c_{\alpha_1} s_{\alpha_3} + s_{\alpha_1} s_{\alpha_2} c_{\alpha_3}) & c_{\alpha_2} c_{\alpha_3} \end{pmatrix}$$

Coupling to massive gauge bosons: (identical for all four types)

	$c_{h_iVV} = c_\beta R_{i1} + s_\beta R_{i2}$
h_1	$c_{\alpha_2} c_{\beta-\alpha_1}$
h_2	$-c_{\beta-\alpha_1} s_{\alpha_2} s_{\alpha_3} + c_{\alpha_3} s_{\beta-\alpha_1}$
h_3	$-c_{\alpha_3} c_{\beta-\alpha_1} s_{\alpha_2} - s_{\alpha_3} s_{\beta-\alpha_1}$

Coupling to fermions: (same pattern as in 2HDM)

	u-type (c_{h_itt})	d-type (c_{h_ibb})	leptons ($c_{h_i\tau\tau}$)
type I	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$
type II	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i1}}{c_\beta}$
type III (lepton-specific)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$
type IV (flipped)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i2}}{s_\beta}$

"Physical" input parameters:

$$\alpha_{1,2,3}\,, \quad \tan\beta\,, \quad v\,, \quad v_S\,, \quad m_{h_{1,2,3}}\,, \quad m_A\,, \quad M_{H^\pm}\,, \quad m_{12}^2$$

Needed to fit the $\gamma\gamma$ and $b\bar{b}$ excesses: $m_{h_1} \sim 95$ GeV, $m_{h_2} \sim 125$ GeV

- $c^2_{h_1VV}$ strongly reduced for μ_{LEP}
- c_{h_1bb} reduced to enhance $\text{BR}(h_1 \to \gamma\gamma)$
- c_{h_1tt} not reduced for μ_{CMS}
- $c_{h_1\tau\tau}$ possibly reduced to enhance $\text{BR}(h_1 \to \gamma\gamma)$

	Decrease $c_{h_1b\bar{b}}$	No decrease $c_{h_1t\bar{t}}$	No enhancement $c_{h_1\tau\bar{\tau}}$
type I	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-(	$(\frac{R_{12}}{s_\beta})$:-)
type II	$(\frac{R_{11}}{c_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{11}}{c_\beta})$:-)
type III	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-(	$(\frac{R_{11}}{c_\beta})$:-(
type IV	$(\frac{R_{11}}{c_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-)	$(\frac{R_{12}}{s_\beta})$:-)

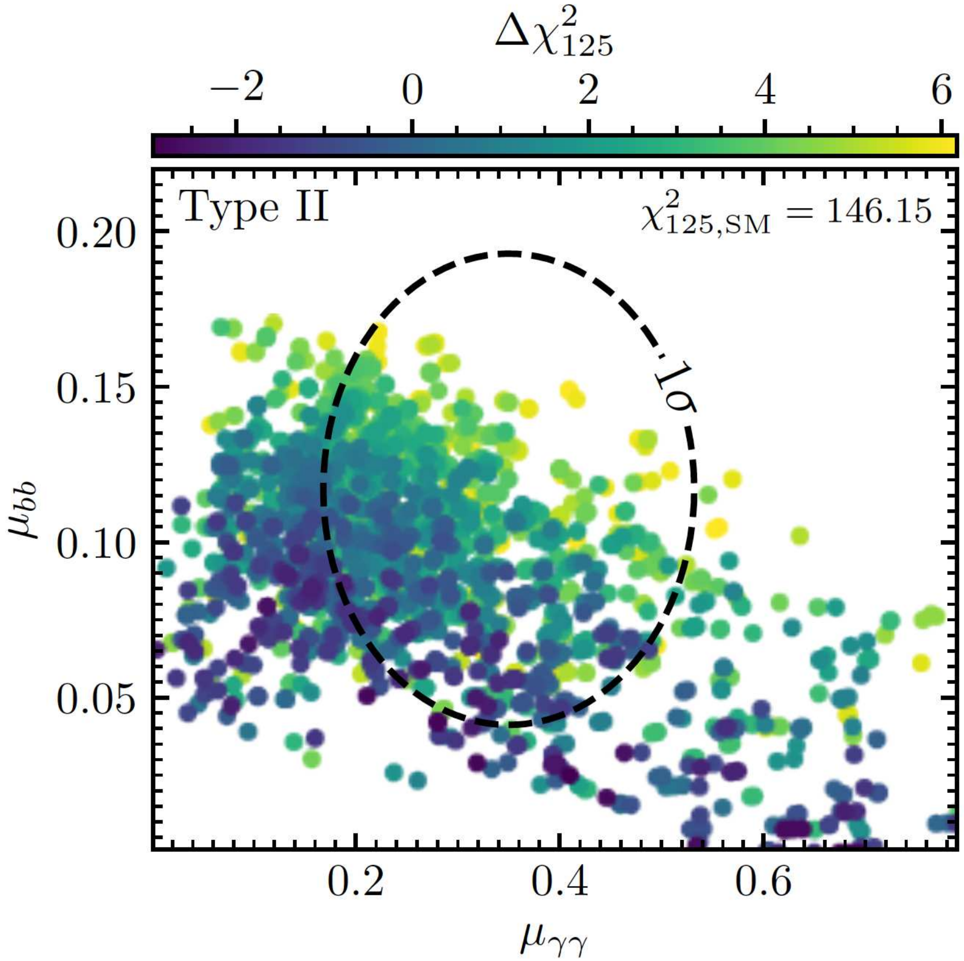
Type II and IV: c_{h_1bb} and c_{h_1tt} independent

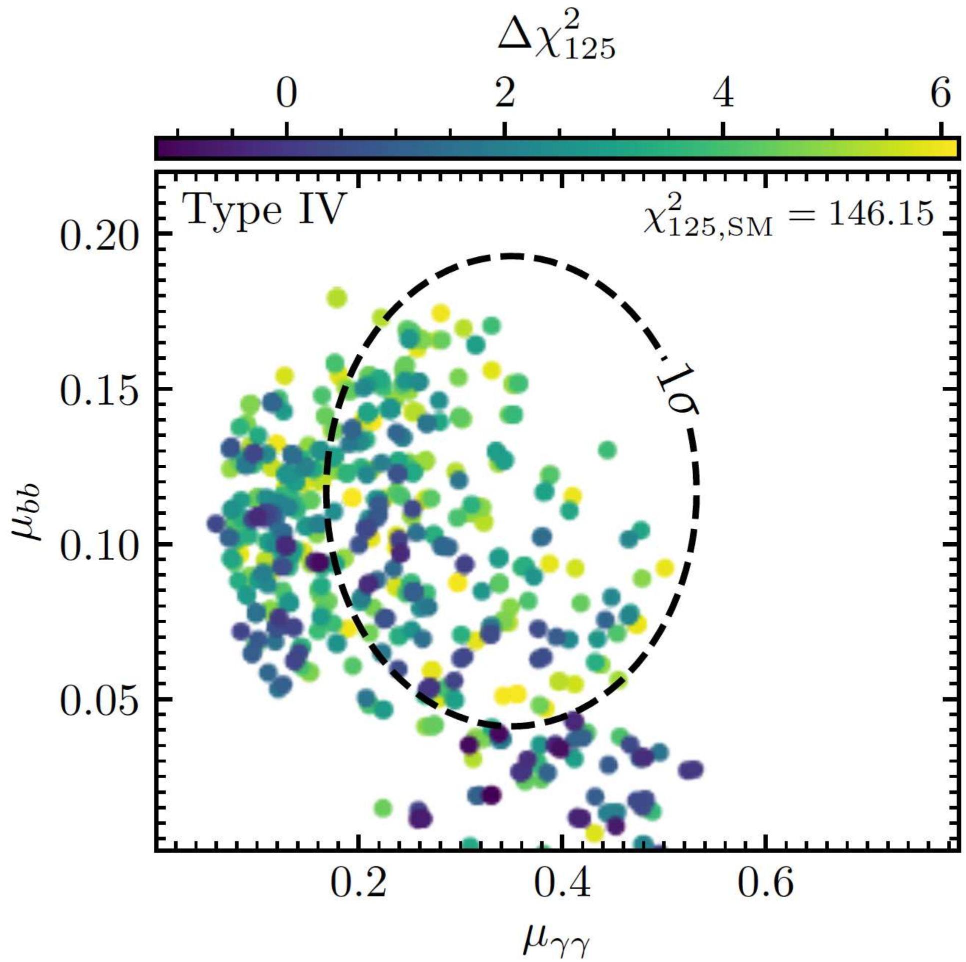
Type II vs. IV: $c_{h_1\tau\tau}$ can be suppressed or enhanced

$\Rightarrow$ possible explanations: $\gamma\gamma$, $b\bar{b}$: type II/IV, $\tau\tau$: type IV only

S2HDM type II vs. type IV

[*T. Biekötter, S.H., G. Weiglein '23*]

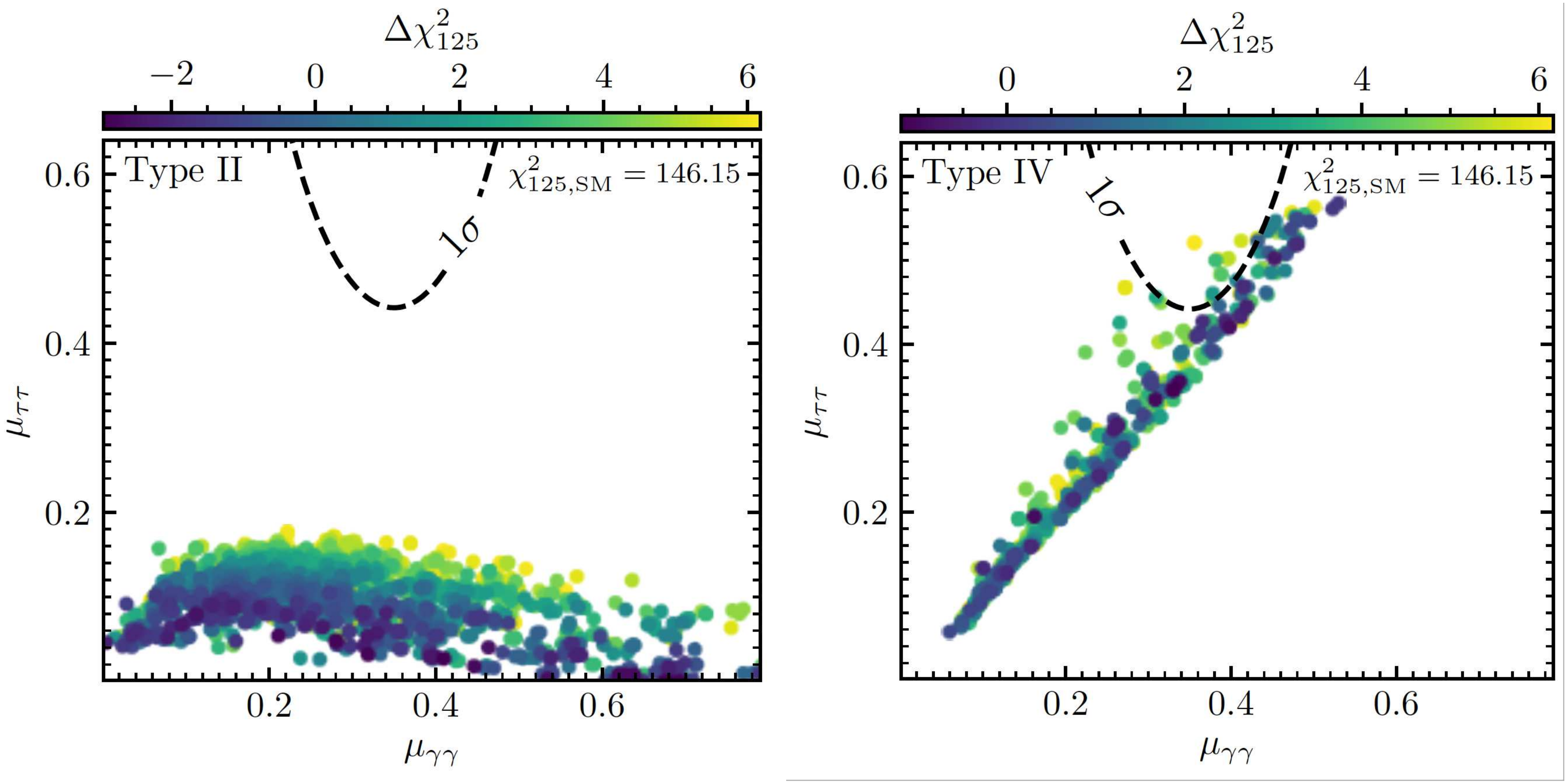




Color coding: χ^2_{125} from HiggsSignals

$\Rightarrow$ both type II and IV can fit the $\gamma\gamma$ and bb excesses

[T. Biekötter, S.H., G. Weiglein '23]



Color coding: χ^2_{125} from HiggsSignals

$\Rightarrow$ only type IV can fit marginally the $\gamma\gamma$ and $\tau\tau$ excesses

3. Physics opportunities at e^+e^- colliders

What can we learn from future measurements?

– LHC h_{125} coupling measurements

– HL-LHC h_{125} coupling measurements

– ILC h_{125} coupling measurements

– direct production of ϕ_{95} at the LHC

– direct production of ϕ_{95} at the HL-LHC

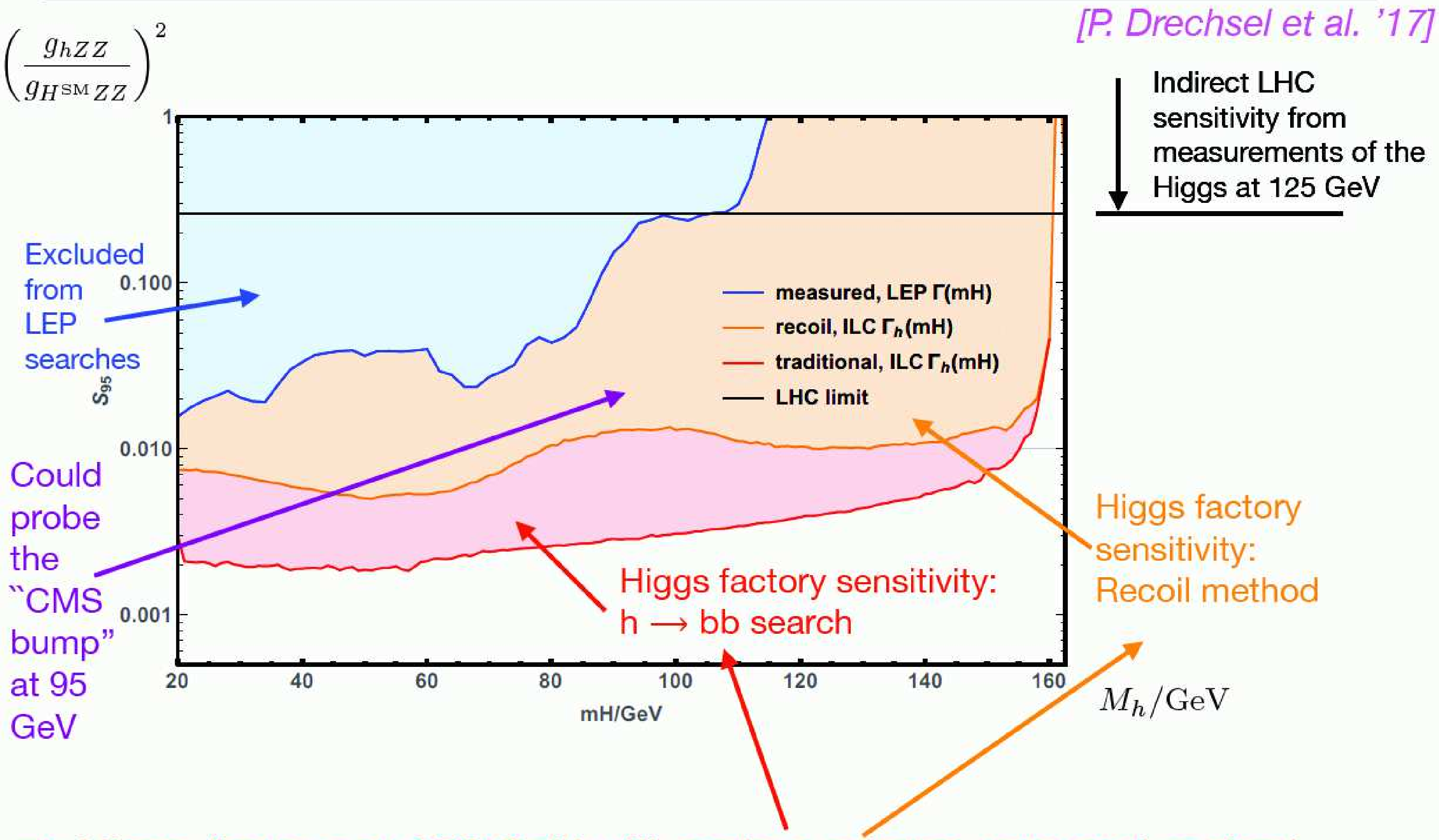
– direct production of ϕ_{95} at the ILC

– ILC ϕ_{95} coupling measurements

– production of other BSM Higgs bosons at the LHC/HL-LHC/ILC/...

ILC = ILC (or other e^+e^- collider)

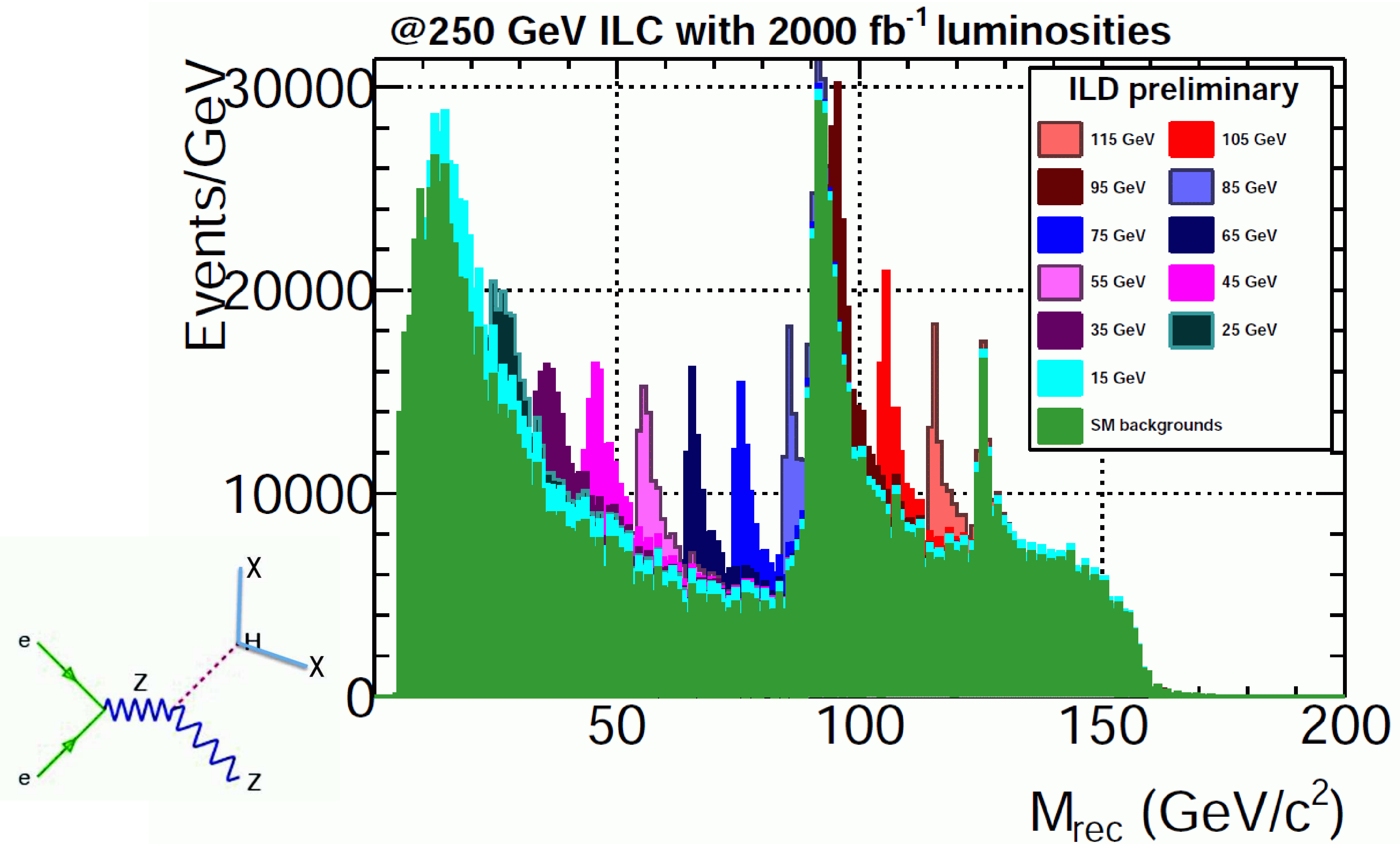
Example for discovery potential for new light states: Sensitivity at 250 GeV with 500 fb^{-1} to a new light Higgs



⇒ Higgs factory at 250 GeV will explore a large untested region!

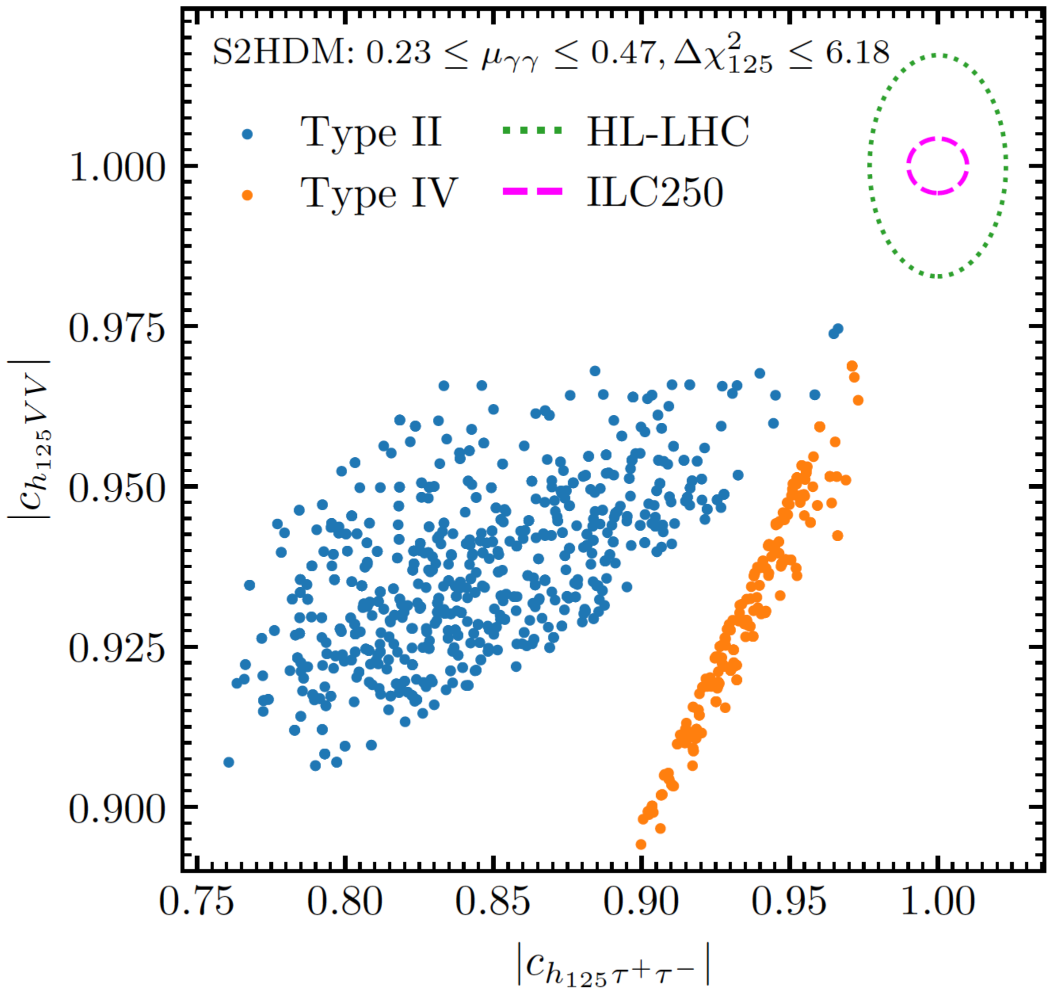
[*Taken from G. Weiglein '18*]

Search for invisibly decaying light Higgses via recoil method: [*Y. Wang et al. '18*]



h_{125} coupling measurements at the HL-LHC/ILC

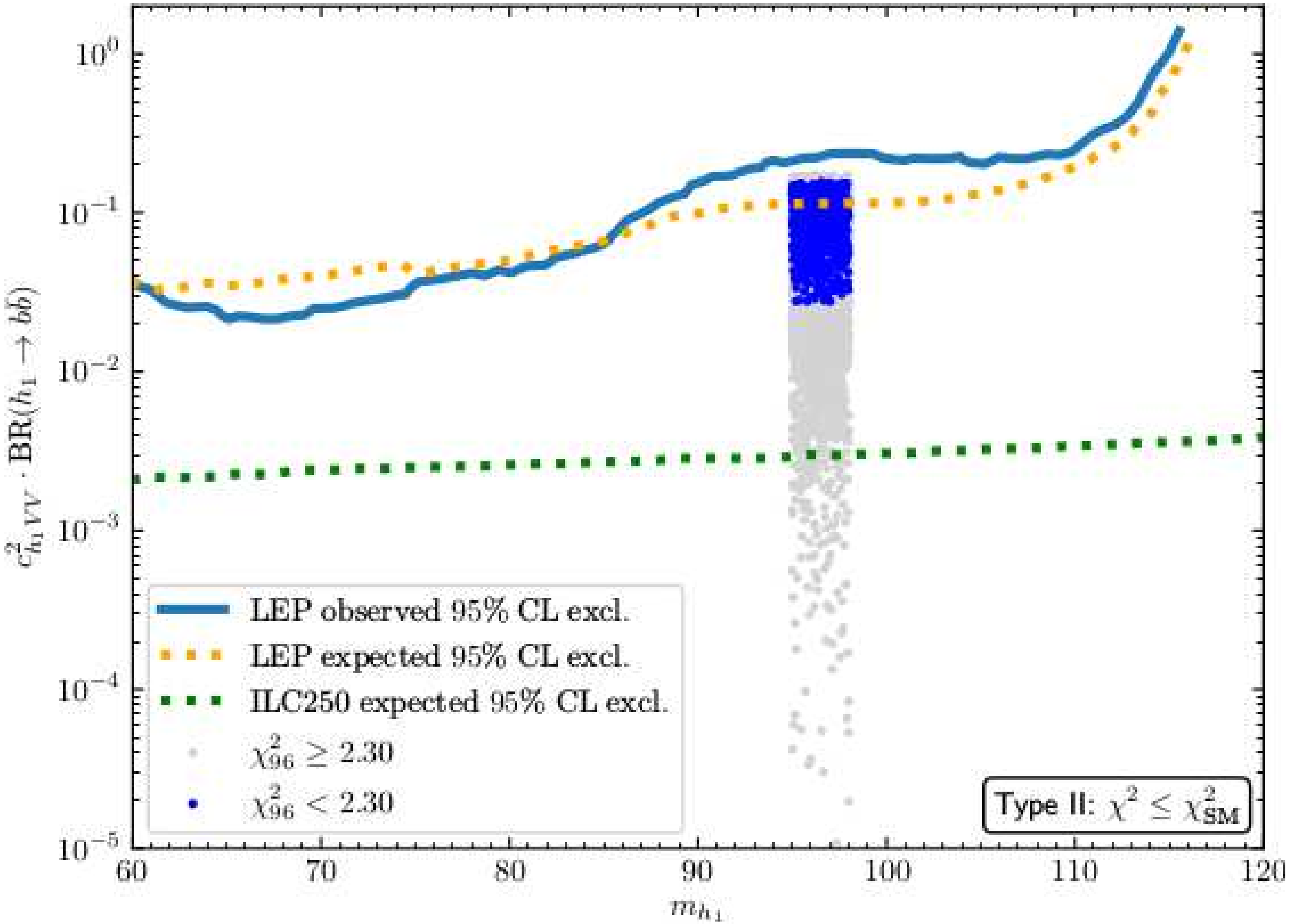
[*T. Biekötter, S.H., G. Weiglein '23*]



$\Rightarrow$ both types show some deviation from SM

Production of the light Higgs at the ILC:

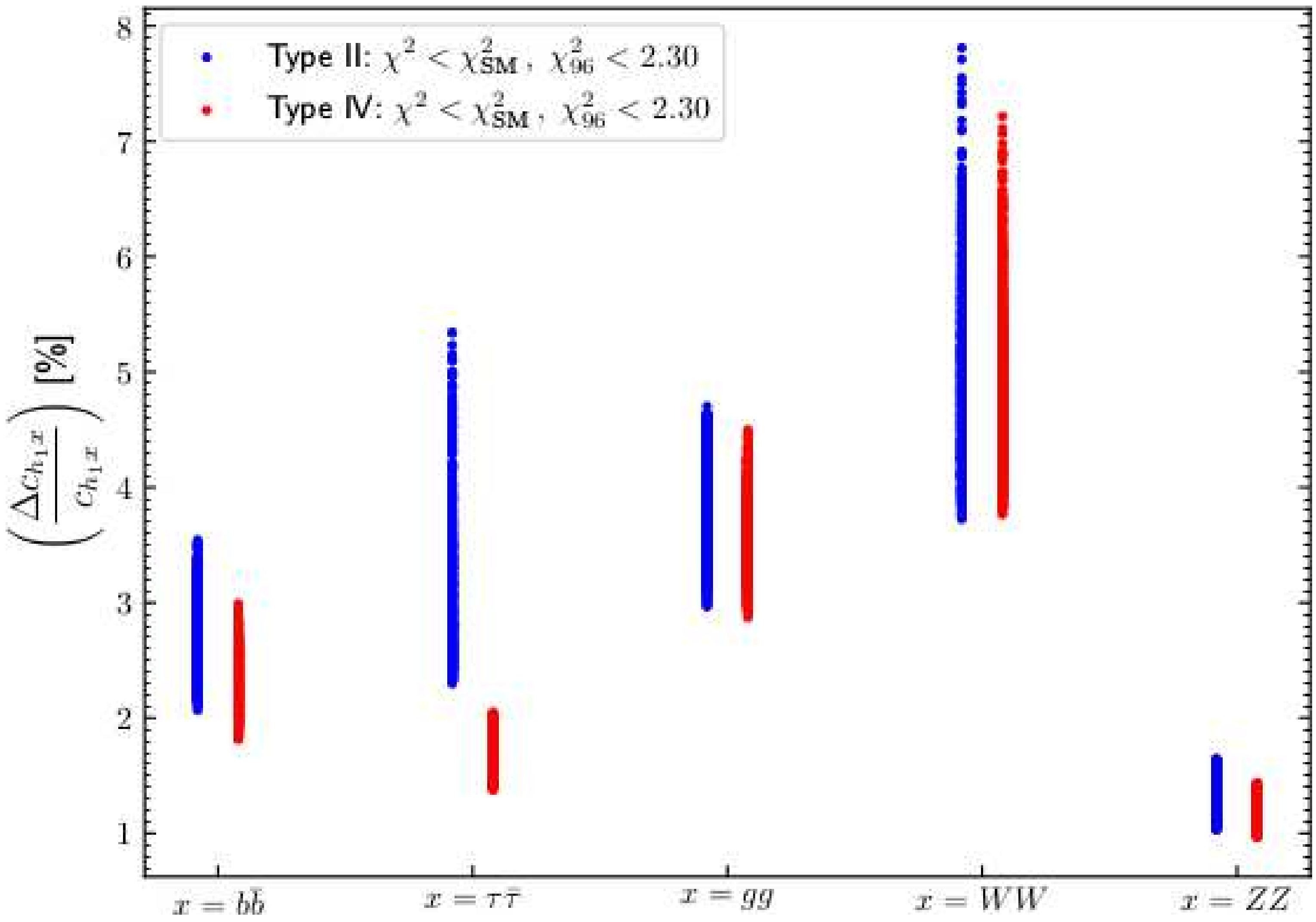
[*T. Biekötter, S.H., G. Weiglein – PRELIMINARY*]



⇒ new state easily in the reach of the ILC ⇒ coupling measurements

h_{95} coupling measurements at the HL-LHC/ILC

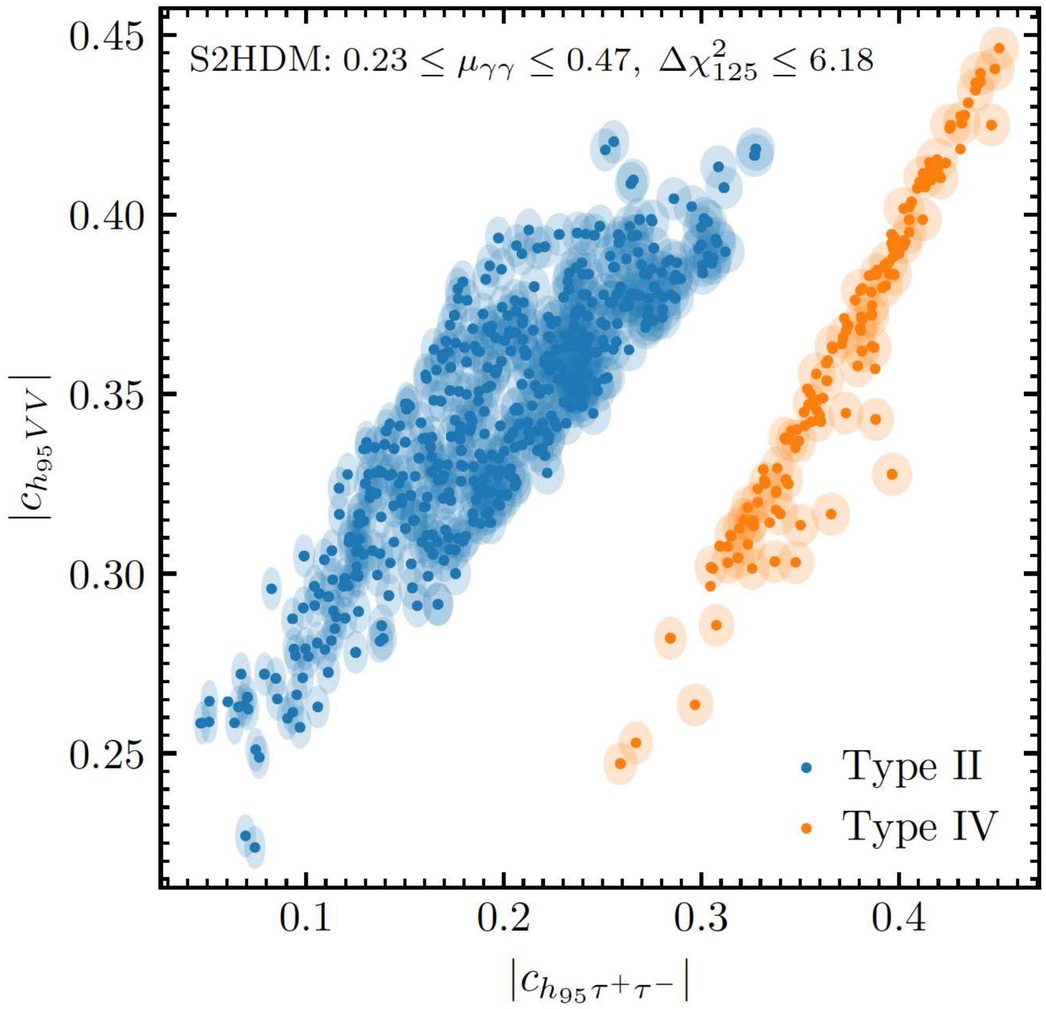
[*T. Biekötter, S.H., G. Weiglein – PRELIMINARY*]



⇒ clear difference in $g_{h_{95}\tau\tau}$ as expected

h_{95} coupling measurements at the HL-LHC/ILC

[*T. Biekötter, S.H., G. Weiglein '23*]



⇒ models clearly distinguishable!

4. Conclusinos

- Interesting excesses at ~ 95 GeV:
  CMS: $pp \to \phi \to \gamma\gamma$ ($3\,\sigma$ local) ATLAS: no sensitivity (yet)
  LEP: $e^+e^- \to Z\,\phi \to Z\,b\bar{b}$ ($2\,\sigma$ local)
  CMS: $pp \to \phi \to \tau\tau$ ($2.5\,\sigma$ local)

- $\Rightarrow$ N2HDM analysis (also S2HDM)
  $\Rightarrow$ possible explanations: $\gamma\gamma$, $b\bar{b}$: type II/IV, $\tau\tau$: type IV only

- ILC250: analysis of h_{125}:
  – precision measurements of couplings can distinguish N2HDM vs. SM
  – possible distinction between type II and IV

- ILC250: analysis of h_{95}:
  – h_{95} can be produced abundantly
  – precision in couplings: 1-8%: g_Z best from production
  – coupling measurements ($\tau\tau$, ZZ) clearly distinguishes type II and IV


Higgs Days at Santander 2023
Theory meets Experiment
4 - 8 September
Contact: Sven.Heinemeyer@cern.ch
Local: Alicia.Calderon@cern.ch
Gervasio.Gomez@cern.ch
http://hdays.csic.es


Further Questions?

SUSY realizations

What about SUSY??

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– NMSSM
– $\mu\nu$SSM
– . . .

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– NMSSM
– $\mu\nu$SSM
– . . .

Q: Can the models fit the excesses despite the additional SUSY constraints on the Higgs sector **???**

What about the NMSSM?

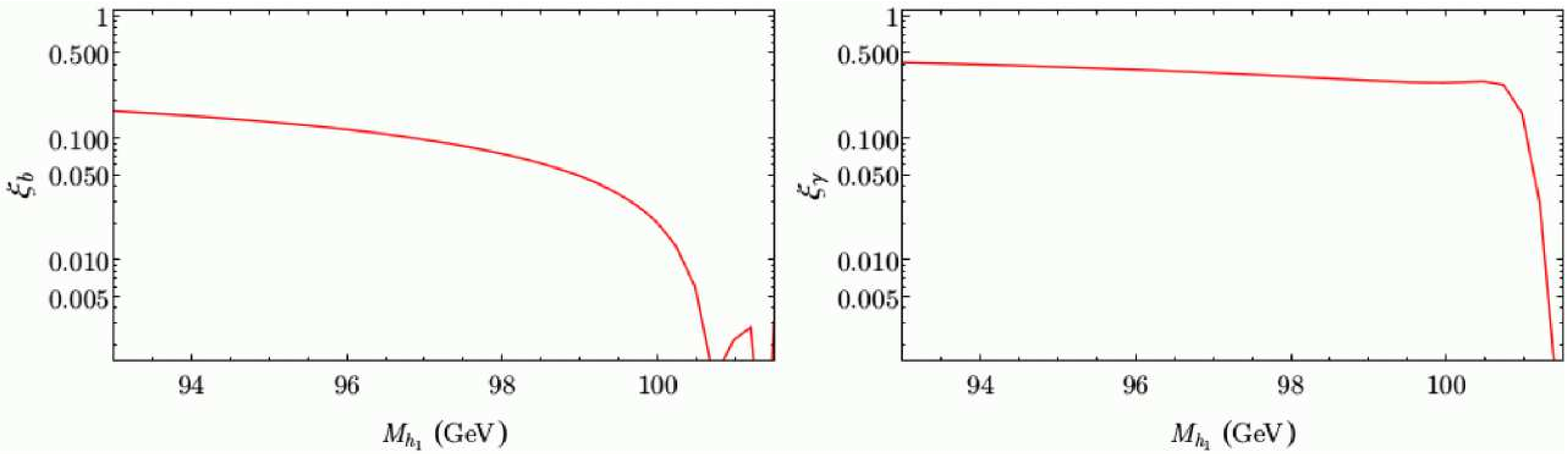
[F. Domingo, S.H., S. Passehr, G. Weiglein '18]

Parameters:

$\lambda = 0.6,\ \kappa = 0.035,\ \tan\beta = 2,\ \mu_{\text{eff}} = (397 + 15x)$ GeV, $M_{H^\pm} = 1$ TeV, $A_\kappa = -325$ GeV, $M_{\text{SUSY}} = 1$ TeV, $A_t = A_b = 0$

$$\xi_b \equiv \frac{\Gamma[h_1 \to ZZ] \cdot \text{BR}[h_1 \to b\bar{b}]}{\Gamma[H_{\text{SM}}(M_{h_1}) \to ZZ] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \to b\bar{b}]} \sim \frac{\sigma[e^+e^- \to Z(h_1 \to b\bar{b})]}{\sigma[e^+e^- \to Z(H_{\text{SM}}(M_{h_1}) \to b\bar{b})]}$$

$$\xi_\gamma \equiv \frac{\Gamma[h_1 \to gg] \cdot \text{BR}[h_1 \to \gamma\gamma]}{\Gamma[H_{\text{SM}}(M_{h_1}) \to gg] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \to \gamma\gamma]} \sim \frac{\sigma[gg \to h_1 \to \gamma\gamma]}{\sigma[gg \to H_{\text{SM}}(M_{h_1}) \to \gamma\gamma]}\ .$$



$\Rightarrow$ both excesses can be fitted simultaneously well with new $\mu_{\gamma\gamma}$!

What about the $\mu\nu$SSM?

$\mu\nu$SSM: [*D. Lopez-Fogliani, C. Muñoz '06*]

$\mu\nu$SSM: NMSSM + well motivated RPV (in simple terms)

$\Rightarrow$ EW scale seesaw to reproduce the neutrino data

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$\mu\nu$SSM: NMSSM + well motivated RPV (in simple terms)

$\Rightarrow$ EW scale seesaw to reproduce the neutrino data

Can the $\mu\nu$SSM explain the two excesses?

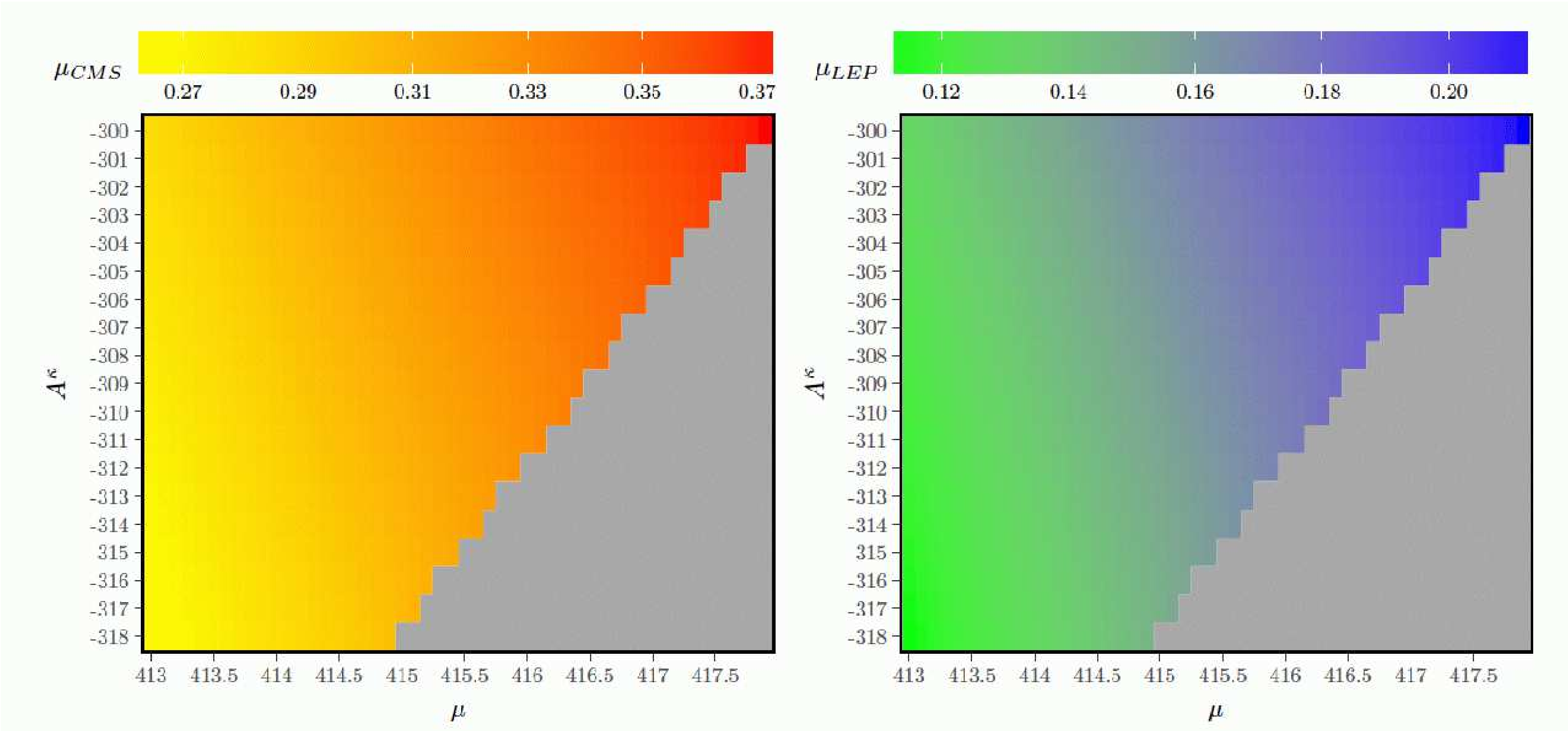
[*T. Biekötter, S.H., C. Muñoz '17*]

v_{iL}	Y_i^ν	A_i^ν	$\tan\beta$	μ	λ	A^λ	κ	A^κ	M_1
$\sqrt{2}\cdot 10^{-5}$	10^{-7}	-1000	2	$[413; 418]$	0.6	956.035	0.035	$[-300; -318]$	100

M_2	M_3	$m^2_{\tilde{Q}_{iL}}$	$m^2_{\tilde{u}_{iR}}$	$m^2_{\tilde{d}_{iR}}$	A_1^u	$A_{2,3}^{u,d}$	$(m^2_{\tilde{e}})_{ii}$	A_{33}^e	$A_{11,22}^e$
200	1500	800^2	800^2	800^2	0	0	800^2	0	0

Can the $\mu\nu$SSM explain the two excesses?

[*T. Biekötter, S.H., C. Muñoz '17*]

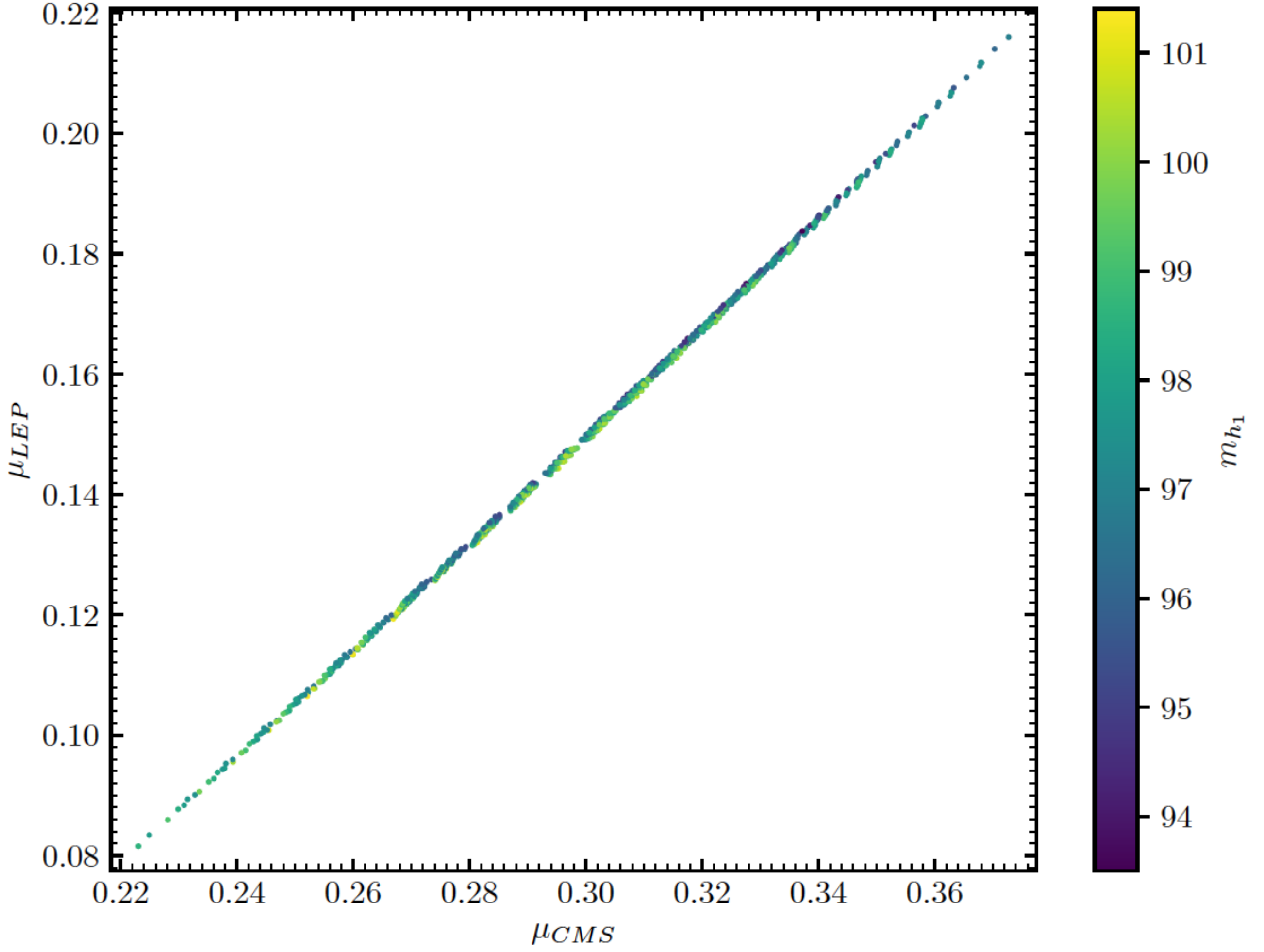


$\Rightarrow$ Yes! :-)

using the <u>new</u> $\mu_{\gamma\gamma}$!

Why does SUSY prefer the new $\mu_{\gamma\gamma}$?

[*T. Biekötter, S.H., C. Muñoz '19*]



⇒ SUSY enforces strong correlation!

⇒ LEP excess enforces $\mu_{\gamma\gamma} \lesssim 0.35$

How to evaluate the precision of ϕ_{95} coupling measurements?

Start with data of the SM Higgs:

SM Higgs BRs: [YR4 LHCHXSWG]

final state	$b\bar{b}$	gg	$\tau^+\tau^-$	WW^*	σ_{ZH}
BR	0.582	0.082	0.063	0.214	206 fb

SM Higgs coupling uncertainties:

ILC, $\mathcal{L}_{\text{int}} = 2\,\text{ab}^{-1}$ at $\sqrt{s} = 250$ GeV [T. Barklow et al. '17]

coupling	$b\bar{b}$	gg	$\tau^+\tau^-$	WW	ZZ
rel. unc. [%]	1.04	1.60	1.16	0.65	0.66

SM Higgs S/B: [S. Dawson et al. '13] [J. Tian, priv. commun.]

coupling	$H \to b\bar{b}$	$H \to gg$	$H \to \tau^+\tau^-$	$H \to WW$	σ_{ZH}
S/B	1/0.89	1/13	1/0.44	1/0.96	1/1.65

Some more basics:

$$f := S/B \equiv N_S/N_B$$

$$\frac{\Delta N_S}{N_S} = \frac{1}{\sqrt{N_S}} \sqrt{1 + 1/f}$$

Holds is background is known perfectly and the overall uncertainty is dominated by statistical precision

Uncertainty improves with $1/\sqrt{N_S}$ for $f = S/B \gg 1$

Cross section for ϕ_{95}:

$$\sigma(e^+e^- \to \phi Z) = \sigma_{\text{SM}}(e^+e^- \to Z\, H_{\text{SM}}^{\phi_{95}}) \times |c_{\phi VV}|^2$$

$$\sigma_{\text{SM}}(e^+e^- \to Z\, H_{\text{SM}}^{\phi_{95}}) = 0.332\,\text{pb}$$

$\Rightarrow \mathcal{O}\left(10^5\right)$ ϕ_{95}'s can be produced at $\sqrt{s} = 250$ GeV and $\mathcal{L}_{\text{int}} = 2\,\text{ab}^{-1}$

Evaluating uncertainties:

- Coupling is measured via decay
  A new Higgs boson ϕ couples with g_x to xx

$$\Gamma(\phi \to xx) \propto g_x^2$$

$$\mathrm{BR}(\phi \to xx) =: 1/p$$

$$\frac{\Delta N_S}{N_S} = 2\,\frac{\Delta g_x}{g_x}\left(1 - \frac{1}{p}\right)$$

- Coupling is measured via production: g_Z

$$\sigma(e^+e^- \to Z\phi) \propto g_Z^2$$

$$\frac{\Delta N_S}{N_S} = 2\,\frac{\Delta g_x}{g_x}$$

- Final assumption: $$\left(\frac{N_S}{N_B}\right)_H \Big/ \left(\frac{N_S}{N_B}\right)_\phi = f_H/f_\phi =: D$$

  with $D = 3$ as starting point

Evaluating uncertainties of ϕ_{95}:

- Coupling is measured via decay

$$\left(\frac{\Delta g_x}{g_x}\right)_\phi = \left(\frac{\Delta g_x}{g_x}\right)_H \times \frac{\left(\frac{\Delta N_S}{N_S}\right)_\phi}{\left(\frac{\Delta N_S}{N_S}\right)_H} \times \frac{\left(1-\frac{1}{p_H}\right)}{\left(1-\frac{1}{p_\phi}\right)}$$

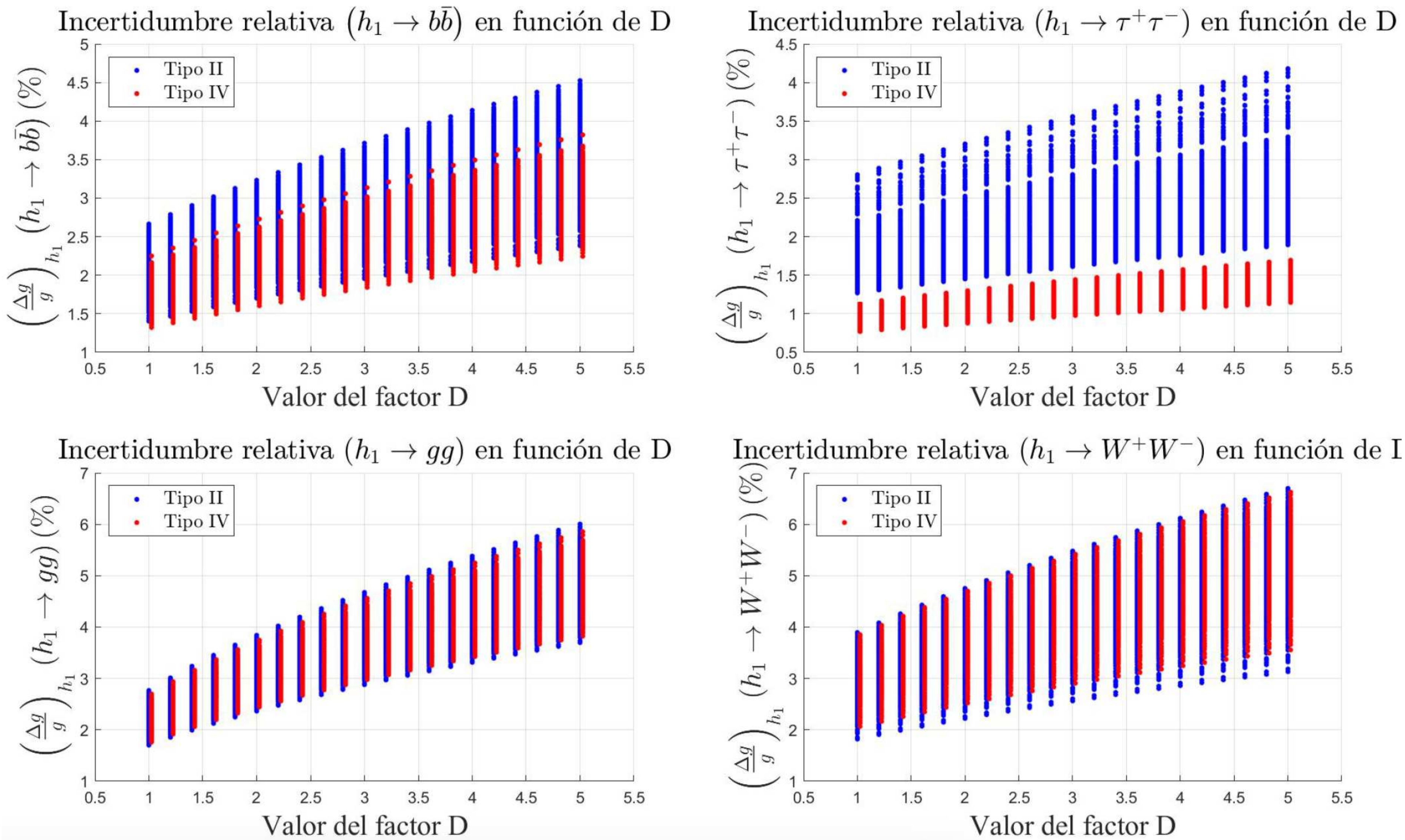
$$\to \sqrt{\frac{D+f_H}{1+f_H}} \times \sqrt{\frac{\sigma(e^+e^-\to ZH)}{\sigma(e^+e^-\to Z\phi)}} \times \sqrt{\frac{\mathrm{BR}(H\to xx)}{\mathrm{BR}(\phi\to xx)}} \times \frac{(1-\mathrm{BR}(H\to xx))}{(1-\mathrm{BR}(\phi\to xx))}$$

- Coupling is measured via production: g_Z (S/B does not change)

$$\left(\frac{\Delta g_Z}{g_Z}\right)_\phi = \left(\frac{\Delta g_Z}{g_Z}\right)_H \times \frac{\left(\frac{\Delta N_S}{N_S}\right)_\phi}{\left(\frac{\Delta N_S}{N_S}\right)_H}$$

$$\to \sqrt{\frac{\sigma(e^+e^-\to ZH)}{\sigma(e^+e^-\to Z\phi)}}$$

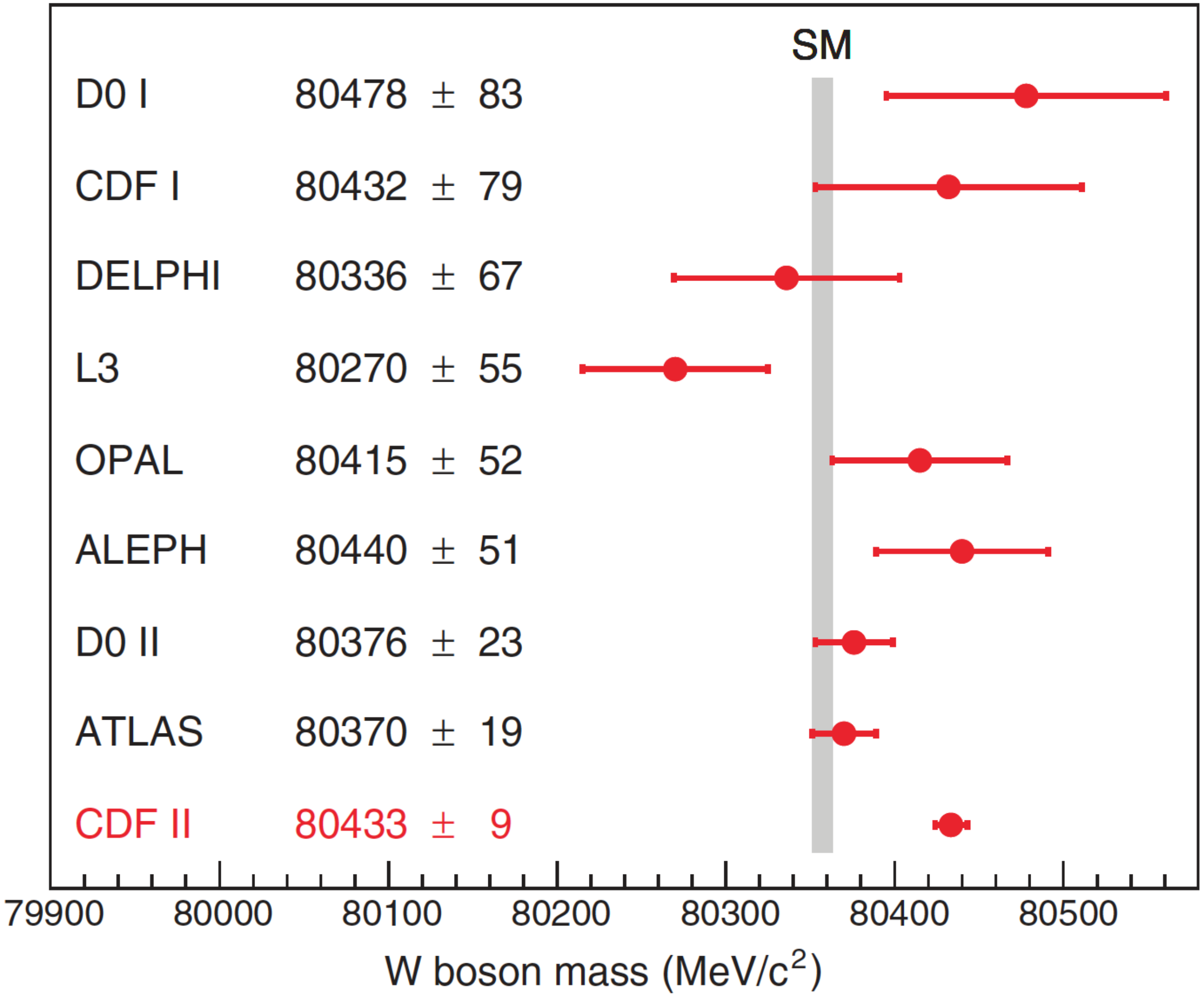
N2HDM: dependence on $D = f_H/f_\phi$: [S.H., P. Toledo '20]



$\Rightarrow$ non-negligible, but small $\Rightarrow$ "robust" result

The mass of the W boson: theory vs. experiment

[*CDF '22*]



⇒ large discrepancy with the SM prediction

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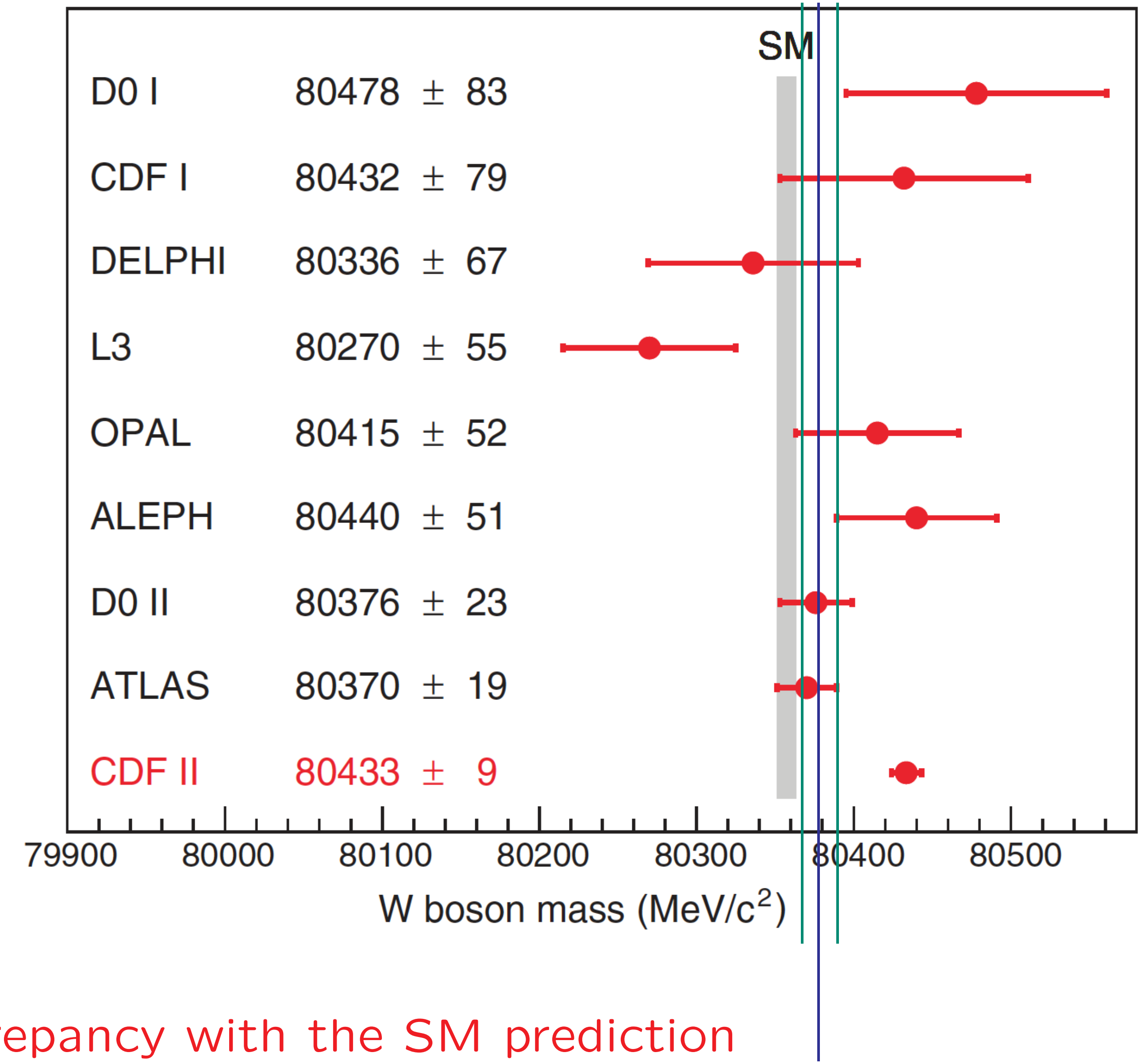


⇒ large discrepancy with the SM prediction

⇒ large discrepancy with other measurements: $M_W^{\text{PDG}} = 80379 \pm 12$ MeV

Approximation of M_W with S, T, U: [M. Peskin, T.Takeuchi '90]

→ capture the gauge boson self-energies

⇒ good approximation in multi-Higgs models

$$M_W^2 = M_W^2\big|_{\rm SM}\left(1 + \frac{s_w^2}{c_w^2 - s_w^2}\Delta r'\right) ,$$

$$\Delta r' = \frac{\alpha}{s_w^2}\left(-\frac{1}{2}S + c_w^2 T + \frac{c_w^2 - s_w^2}{4 s_w^2}U\right) .$$

Main contribution:

$$+\frac{\alpha c_{\rm W}^2}{s_{\rm W}^2}\frac{s_{\rm W}^2}{c_{\rm W}^2 - s_{\rm W}^2}T$$

$$=: +\frac{c_{\rm W}^2}{c_{\rm W}^2 - s_{\rm W}^2}(\alpha T)$$

$$=: +\frac{c_{\rm W}^2}{c_{\rm W}^2 - s_{\rm W}^2}\Delta\rho \qquad \alpha T \equiv \Delta\rho = \frac{\Sigma_Z(0)}{M_Z^2} - \frac{\Sigma_W(0)}{M_W^2}$$

Implications for BSM models

Contribution from 2HDM Higgs sector to $\Delta\rho$:

$$\Delta\rho^{(1)}_{\text{non-SM}} = \frac{\alpha}{16\pi^2 s_W^2 M_W^2} \left\{ \frac{m_A^2 m_H^2}{m_A^2 - m_H^2} \ln \frac{m_A^2}{m_H^2} - \frac{m_A^2 m_{H\pm}^2}{m_A^2 - m_{H\pm}^2} \ln \frac{m_A^2}{m_{H\pm}^2} - \frac{m_H^2 m_{H\pm}^2}{m_H^2 - m_{H\pm}^2} \ln \frac{m_H^2}{m_{H\pm}^2} + m_{H\pm}^2 \right\}$$

$\Rightarrow$ large $\Delta\rho$ needed to accomodate M_W^{CDF}

Before M_W^{CDF}:

$\Rightarrow$ small mass splittings between $m_{H\pm}$-m_H and $m_{H\pm}$-m_A

After M_W^{CDF}:

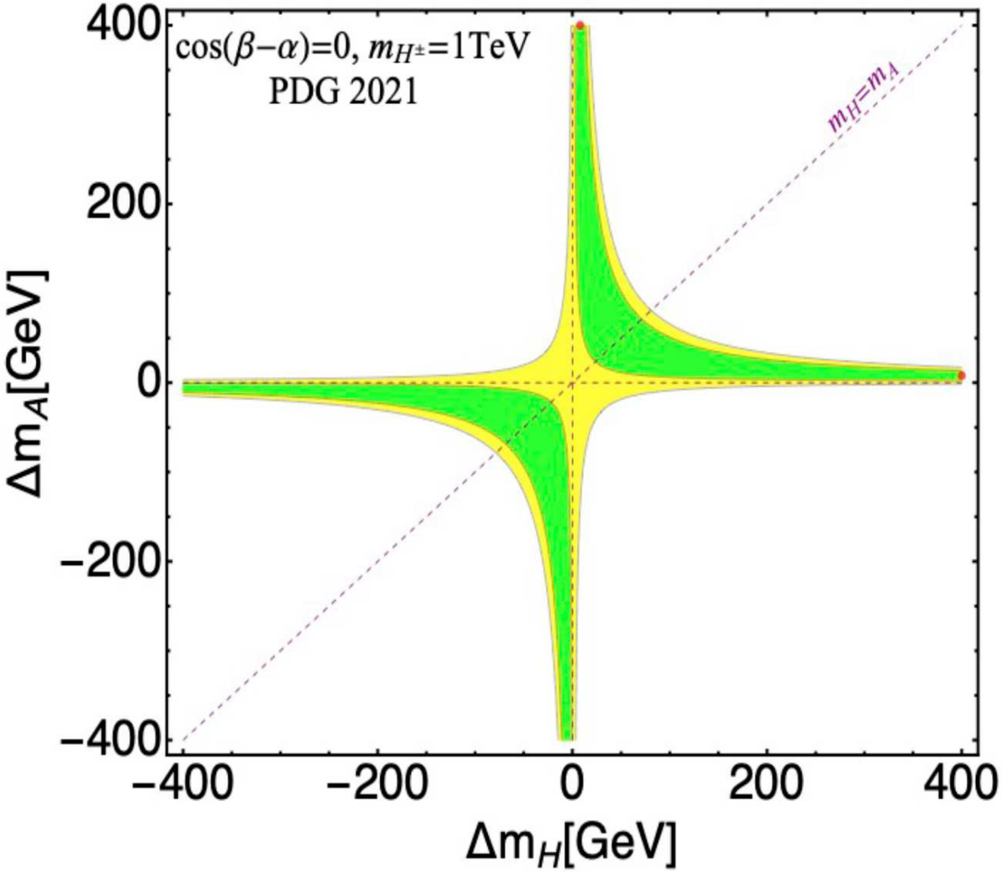
$\Rightarrow$ increased mass splittings to accomodate M_W^{CDF}

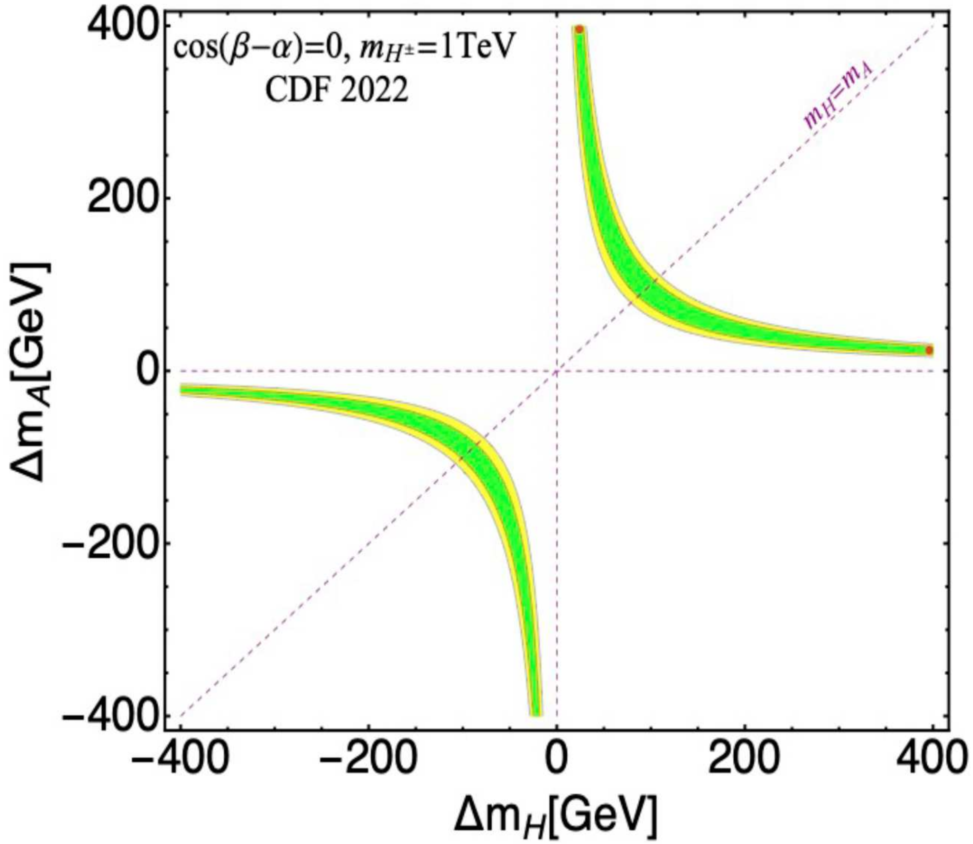
Example: $m_{H^\pm} = 1000$ GeV, $\cos(\beta - \alpha) = 0$

[C. Lu, L. Wu, Y. Wu, B. Zhu '22]

PDG 2021

M_W^{CDF}





⇒ nearly no overlap of the 2σ regions

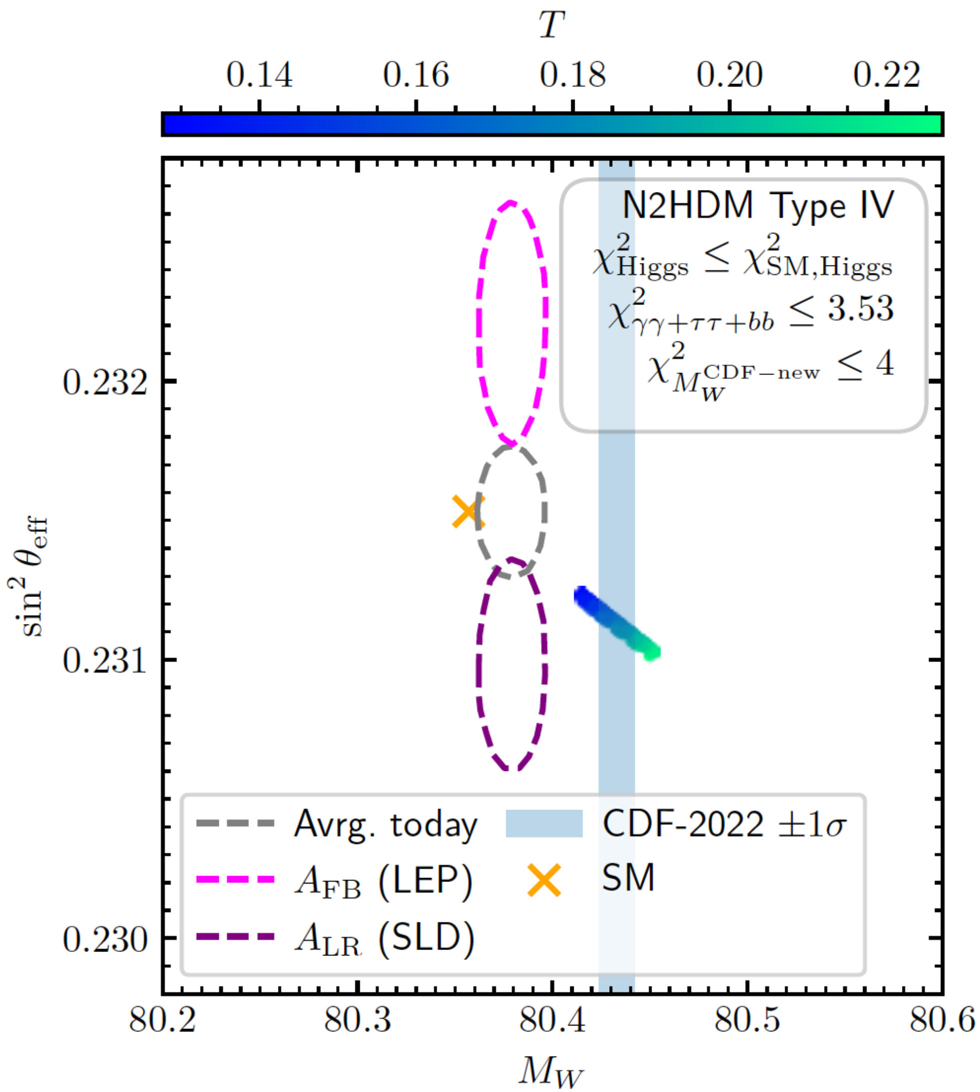
⇒ new CDF value requires relatively large BSM Higgs mass splitting

⇒ upper limit on heavy Higgs-boson masses from $M_W^{\text{CDF}} \oplus$ unitarity/stability!

Can we fit three 95 GeV excesses and M_W^{CDF}?

Can we fit three 95 GeV excesses and M_W^{CDF}? $\Rightarrow$ N2HDM type IV

[*T. Biekötter, S.H., G. Weiglein '22*]



Remember: $\Delta\rho$ goes up $\Rightarrow M_W$ goes up, $\sin^2\theta_{\text{eff}}$ goes down

$\Rightarrow$ agreement only with SLD value of $\sin^2\theta_{\text{eff}}$

Dedicated workshop at CERN:

MWDays23 workshop

17–20 Apr 2023
CERN
Europe/Zurich timezone

Confirmed speakers:

- M. Boonekamp (ATLAS, MW combination working group)
- S. Camarda (ATLAS)
- C. Hays / A. Kotwal (CDF)
- M. Ramon-Pernas (LHCb)
- G. Wilson (future e^+e^- colliders)
- L. Cieri (SM theory)
- S. Dittmaier (SM theory)
- W. Hollik (SM theory)
- A. Huss (SM theory)
- T. Neumann (SM theory)
- A. Vicini (SM theory)
- P. Nadolsky (PDF)
- M. Ubiali (PDF)
- L. Silvestrini (EW fit)
- K. Mimasu (SMEFT global fits)
- J. Erler (PDG)
- G. Arcadi (BSM)
- J. Braathen (BSM)
- C.-W Chiang (BSM)
- A. Crivellin (BSM)
- J. Evans (BSM)
- F. Sannino (BSM)

Organizing committee

- E. Bagnaschi (CERN/INFN Laboratori Nazionali di Frascati, LOC)
- P. Monni (CERN, LOC)
- P. Giardino (IGFAE)
- S. Heinemeyer (IFT Madrid)
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- G. Weiglein (DESY)