Experimental Physics at Lepton Colliders





Frank Simon @ Summer Student Lectures CERN - July 2023





Part II

Overview

A two-part story

- Part 1:
 - Scientific motivation
 - Future e⁺e⁻ colliders in broad strokes
 - Detectors at future e^+e^- and $\mu^+\mu^-$ colliders
- Part 2:
 - Higgs physics
 - Electroweak precision
 - Top quark physics
 - Into the unknown

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Disclaimer

- The point of the following discussions is not to compare projects in the sense of drawing performance projections shown here.
- but to illustrate certain features of measurements and facilities
- I am focussing on e⁺e⁻ colliders, only few remarks about $\mu^+\mu^-$ •

conclusions which one should be built - that is a multi-facetted question which extends beyond

• The numerical results may not always be perfectly up-to-date - again, the goal is not to compare,





Precision Higgs Measurements

Higgs Factories and beyond

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A rich field to explore



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A rich field to explore





250 GeV: Maximum of ZH production





A rich field to explore





250 GeV:

Maximum of ZH production

350 GeV:

WW fusion kicks in

(and top pair production)





A rich field to explore





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500 - 1000+ GeV:

ttH: direct access to top Yukawa coupling







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WW fusion kicks in (and top pair production)

500 - 1000+ GeV:

ttH: direct access to top Yukawa coupling

500 GeV; 1+ TeV: Higgs self-coupling







A rich field to explore



- 240 250 GeV: the minimum energy for a Higgs factory
- ~ 350 GeV: Additional production mode, also still access to ZH
- Higher energies: More processes
- 125 GeV, and extreme luminosity: A possibility to measure electron Yukawa coupling

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250 GeV:

Maximum of ZH production

350 GeV: WW fusion kicks in (and top pair production)

500 - 1000+ GeV:

ttH: direct access to top Yukawa coupling

500 GeV; 1+ TeV: Higgs self-coupling







Model Independence: The Pillar of Higgs Physics in e+e-

The ZH Higgsstrahlung process

- What model independence means: Measure the coupling of the Higgs Bosons to elementary particles free from model assumptions (e.g. how it decays)
 - Requires: The "tagging" of Higgs production without observing the particle directly
 - Not possible at hadron colliders













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Hadronic Recoils & Invisible Decays

Fully exploiting Higgsstrahlung



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Hadronic Recoils & Invisible Decays

Fully exploiting Higgsstrahlung



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Precision Measurements of Couplings

Exploring the Higgs Sector

• The main measurements to make:







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directly constrain the coupling of Higgs to Z in a model-independent way





Precision Measurements of Couplings

Exploring the Higgs Sector

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 σ x BR for specific Higgs decays - here the mass of 125 GeV is giving us many possibilities





HWW

~ 7.444



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Precision Measurements of Couplings

Exploring the Higgs Sector

• The main measurements to make:



 σ x BR for specific Higgs decays - here the mass of 125 GeV is giving us many possibilities



measure couplings to fermions and bosons using production and decay ~ 9.HW YHWW/W Experiments It Colliders - CERN Summer Student Lectures, July 2023



directly constrain the coupling of Higgs to Z in a model-independent way

- can be made model-independent in combination with the measurement of the HZ coupling in recoil





Unique Measurements at Lepton Colliders

Enabled by the clean environment

• H->bb: A difficult channel at LHC, a "simple" measurement in e+e-



of Higgs produced: ~4,000,000 significance: 5.4o

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 Low backgrounds, and highly capable detectors enable observations of final states that are hard or impossible at LHC

5.2σ







^{~400}



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Unique Méasurements at Lepton Colliders

Enabled by the clean environment



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- The Higgs coupling measurements at any present and future collider unfold their full potential in global fits of all observables - possibly beyond Higgs measurements alone
 - The evaluation of the potential of future colliders is based on such fits using projected precisions on various Higgs (and other) measurements as input







- of all observables possibly beyond Higgs measurements alone
 - various Higgs (and other) measurements as input





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total width as a free parameter: no constraints imposed on BSM decays

N.B.: Not fully model independent, does not account for certain possible BSM features of HV couplings







- of all observables possibly beyond Higgs measurements alone
 - various Higgs (and other) measurements as input

Typical fits used in this context:
$$C_{ZH} = g_{HZZ}^2$$
• "Model-independent" fit
minimize a χ^2 with
all measurements: $\chi^2 = \sum_i \frac{(C_i - 1)^2}{\Delta F_i^2}$ $C_{ZH,H \to b\bar{b}} = \frac{g_{HZZ}^2 g_{Hbb}^2}{\Gamma_H}$
 \dots ΔF_i : uncertainty of measurement
(σ or σxBR) ΔF_i

"Model-dependent κ" fit



• The Higgs coupling measurements at any present and future collider unfold their full potential in global fits

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the same as the MI fit, with the total width constrained to the sum of the SM decays

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$$\kappa_i^2 = \frac{\Gamma_i}{\Gamma_i|_{\rm SM}} \qquad \Gamma_{\rm H,md} = \sum_i \kappa_i^2 BR_i$$







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 $C_{\rm ZH} = g$ Typical fits used in this context:

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 ightarrow}$ "Model-independent" fit $C_{\mathrm{H} v_e \bar{v}_e,\mathrm{H}}$ minimize a χ^2 with all measurements: $\chi^2 = \sum_i \frac{(C_i - 1)^2}{\Delta F_i^2}$ ΔF_i : uno
- "Model-dependent κ" fit
- "Model-independent EFT" fit

A global fit of Higgs and other EW observables parametrizing deviations from the SM by various operators - allows for couplings not included in k fit, includes connections between W and Z couplings

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$$C_{ZH} = g_{HZZ}^{2}$$

$$C_{ZH,H\to b\bar{b}} = \frac{g_{HZZ}^{2}g_{Hbb}^{2}}{\Gamma_{H}}$$

$$C_{Hv_{e}\bar{v}_{e},H\to b\bar{b}} = \frac{g_{HWW}^{2}g_{Hbb}^{2}}{\Gamma_{H}}$$
....
$$\Delta F_{i}: \text{ uncertainty of measurement}$$
(σ or σ xBR)

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Model independent measurement at high precision

a few %:



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Model independent measurement at high precision

- e⁺e⁻ colliders provide the possibility for a model-independent measurement of the total width at the level of a few %:
- In the "model-independent fit" framework the total width is obtained from production and decay of the Higgs: $\sigma(\mathrm{ZH}) \times \mathrm{BR}(\mathrm{H} \to \mathrm{ZZ}) \propto \frac{g_{HZZ}^4}{\Gamma_{\mathrm{tot}}} \ \ \mathrm{and} \ \ \sigma(\mathrm{ZH}) \propto g_{HZZ}^2$



 \Rightarrow The low BR of H->ZZ and correspondingly large uncertainties make this determination relatively imprecise





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 \Rightarrow Profits substantially from higher energy, where WW fusion becomes relevant: $\sigma(\mathrm{H}\nu_e\nu_e) \times \mathrm{BR}(\mathrm{H} \to \mathrm{WW}^*) \propto \frac{g_{\mathrm{HWW}}^4}{\Gamma_{\mathrm{tot}}}$

$$\frac{\sigma(e^+e^- \to \mathrm{ZH}) \times \mathrm{BR}(\mathrm{H} \to b\bar{b})}{\sigma(e^+e^- \to \mathrm{H}\nu_e\nu_e) \times \mathrm{BR}(\mathrm{H} \to b\bar{b})} \propto \frac{g_{\mathrm{HZZ}}^2}{g_{\mathrm{HWW}}^2}$$



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Model independent measurement at high precision

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 $\frac{\sigma(e^+e^- \to \text{ZH}) \times \text{BR}(\text{H} \to b\bar{b})}{\sigma(e^+e^- \to \text{H}\nu_e\nu_e) \times \text{BR}(\text{H} \to b\bar{b})} \propto \frac{g_{\text{HZZ}}^2}{g_{\text{HWW}}^2}$ \Rightarrow Higher energies important for width measurements



 \Rightarrow The low BR of H->ZZ and correspondingly large uncertainties make this determination relatively imprecise



- \Rightarrow In EFT fits W and Z are connected, there the width can be well constrained also without WW fusion





Illustrating Interplay and Reach



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The Relevance of Higgs Coupling Measurements

One EFT Example for ILC

Integrated Luminosities [fb⁻¹]



- Precision measurements of couplings may show deviations from the Standard Model
 - "Fingerprinting" of deviation pattern reveals underlying mechanisms

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The Relevance of Higgs Coupling Measurements

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PMSSM2HDM2HDM2YComposite LHT-7 Radion Singlet

 Discrimination power between models illustrated with EFT fit of ILC projections

> arXiv:1708.08912 arXiv:1710.07621

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The Relevance of Higgs Coupling Meas One EFT Example for ILC

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PMSSM2HDM2HDM2YDM2YComposite LHT-7 Radion Singlet SM

- Discrimination power between models illustrated with EFT fit of ILC projections
 - higher energy may be decisive

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Accessing the Couplings to First Generation Leptons

Requiring extreme luminosities of circular colliders

- The only chance to access couplings to first generation: Study of s-channel Higgs production in e+ecollisions
 - Requires high luminosities and very small energy spread at 125.1 GeV



and 3 years may reach a result

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Requires special monochromatization to





Directly measuring the Coupling to the Top Quark

A higher-energy exclusive





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• Direct access to the top Yukawa coupling provided by ttH final state: requires energy \geq 500 GeV (ideal ~ 550 GeV - 1.5 TeV)







Directly measuring the Coupling to the Top Quark

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ILC: $\Delta g_{ttH}/g_{ttH} \sim 6.3\%$ with 4 ab⁻¹ @ 500 GeV would be ~ 3% @ 550 GeV (and ~ 13% @ 485 GeV: achieving design energy critical!)

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• Direct access to the top Yukawa coupling provided by ttH final state: requires energy \geq 500 GeV (ideal ~ 550 GeV - 1.5 TeV)



CLIC: Δg_{ttH}/g_{ttH} ~ 2.9% with 2.5 ab⁻¹ @ 1.4 TeV









• Two processes with sensitivity at e⁺e⁻ colliders:



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cross section depends nonlinearly on λ , measurements at different energies / of different processes lift









• Two processes with sensitivity at e⁺e⁻ colliders:



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cross section depends nonlinearly on λ , measurements at different energies / of different processes lift



Full potential unfolds in the multi-TeV region through growing σ of VBF process:

- 10% measurement feasible $\widehat{}$
- Significant observation also of ZHH channel in lower-energy running (up to $\sim 1.5 \text{ TeV}$)







Indirect Measurement of the Self Coupling

Accessible via particle loops

• The self-coupling also influences single Higgs production:



Interplay of different energies key. With optimised running, and increased Lint at 240 GeV and 365 GeV 20% may be doable.

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Model-dependent: assumptions required for interpretation!

Overall precision limited, ~ 33% at FCC-ee combined with HL-LHC (which provides ~ 50%)





Higgs Physics at Muon Colliders Brief overview

• In general the same processes as for e⁺e⁻, but with the backdrop of a much larger background, and reduced acceptance at small angles (which has an impact on WW fusion processes in particular). Here (much) higher energy can compensate!



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- $WW \rightarrow H$
- $ZZ \rightarrow H$
- $VV \rightarrow W^{\pm}H$
- VV→ZH
- ZΗ ----
 - VV→tīH
- ----- ttH
 - ZHH
 - $VV \rightarrow HH$





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Overall Precision Perspective

Including muon colliders

• An EFT fit, performed for Snowmass



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precision reach on effective Higgs couplings from SMEFT global fit MuC 3TeV 1 CLIC 380GeV₁ CLIC +1.5TeV_{2.5} CLIC +3TeV₅ MuC 10TeV ₁₀ MuC 125GeV_{0.02}+10TeV ₁₀ subscripts denote luminosity in ab⁻¹, Z & WW denote Z-pole & WW threshold δg_{H}^{cc} δg_H^{bb} $\delta g_{H}^{\mu\mu}$ δgH $\delta \Gamma_H$







Electroweak Precision

A Playground for Circular Colliders

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The FCC-ee Program at Z and WW

The ultimate electroweak program



- Building on the success of LEP & LEP II • High-statistics program at the Z - pole • W pair production - mass measurement and beyond

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with 2 IPs: 5x10¹² Zs (10⁵ x LEP) 10⁸ W pairs (2x10³ x LEP)

N.B.: Measurements also possible at linear colliders, but the statistics will be orders of magnitude smaller due to their lower luminosity at low energy.







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- Building on the success of LEP & LEP II • High-statistics program at the Z - pole • W pair production - mass measurement
- Improving electroweak precision observables, enter into global fits

with 2 IPs: 5x10¹² Zs (10⁵ x LEP) 10^8 W pairs (2x10³ x LEP)

Indirect searches for New Physics

N.B.: Measurements also possible at linear colliders, but the statistics will be orders of magnitude smaller due to their lower luminosity at low energy.







Electroweak Measurements

Cross sections and asymmetries



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Lineshapes and Thresholds

The things to explore



• Lineshapes, cross sections, asymmetries provide access to a wide range of electroweak precision measurements, putting the Standard Model to extremely stringent tests





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FCC-ee Electroweak Projections

Summary

Observable	Present	FCC-ee	FCC-ee	Comment and dominant exp. error
	value \pm error	Stat.	Syst.	
$m_{\rm Z}~({\rm keV})$	$91,186,700 \pm 2200$	4	100	From Z lineshape scan; beam energy calibration
$\Gamma_{\rm Z}~({\rm keV})$	$2,495,200 \pm 2300$	4	25	From Z lineshape scan; beam energy calibration
$R_{\ell}^{\rm Z}~(imes 10^3)$	$20,767\pm25$	0.06	0.2 - 1.0	Ratio of hadrons to leptons; acceptance for lept
$\alpha_S(m_{ m Z}^2)~(imes 10^4)$	$1,196\pm30$	0.1	0.4 - 1.6	From $R_{\ell}^{\rm Z}$ above
$R_b \; (\times 10^6)$	$216,290\pm 660$	0.3	< 60	Ratio of $b\overline{b}$ to hadrons; stat. extrapol. from SLI
$\sigma_{\rm had}^0$ (×10 ³) (nb)	$41,541\pm37$	0.1	4	Peak hadronic cross section; luminosity measure
$N_{\nu}~(imes 10^3)$	$2,996\pm7$	0.005	1	Z peak cross sections; luminosity measurement
$\sin^2 heta_{ m W}^{ m eff}$ (×10 ⁶)	$231,480\pm160$	1.4	1.4	From $A_{\rm FB}^{\mu\mu}$ at Z peak; beam energy calibration
$1/lpha_{ m QED}(m_{ m Z}^2)~(imes 10^3)$	$128,952\pm14$	3.8	1.2	From $A_{\rm FB}^{\overline{\mu}\overline{\mu}}$ off peak
$A_{\rm FB}^{b,0}~(imes 10^4)$	992 ± 16	0.02	1.3	b-quark asymmetry at Z pole; from jet charge
$A_{e}^{-}(\times 10^{4})$	$1,498\pm49$	0.07	0.2	from $A_{\rm FB}^{{\rm pol},\tau}$; systematics from non- τ background
$m_{ m W}~({ m MeV})$	$80,350\pm15$	0.25	0.3	From WW threshold scan; beam energy calibrat
$\Gamma_{\rm W} ~({\rm MeV})$	$2,085\pm42$	1.2	0.3	From WW threshold scan; beam energy calibrat
$N_{\nu} ~(\times 10^3)$	$2,920\pm50$	0.8	Small	Ratio of invis. to leptonic in radiative Z returns
$lpha_S(m_{ m W}^2)~(imes 10^4)$	$1,170\pm420$	3	Small	From R^W_ℓ

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Increasing interest

• An e⁺e⁻ collider running at the Z pole is also an excellent flavour factory! The 5 x 10^{12} Zs at FCC-ee will provide: 10^{12} bb events, 1.7 x 10^{11} $\tau^+\tau^-$ events



An excellent testing ground of universality, rare decays; precision measurements of masses and lifetimes





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High-statistics measurements to follow up on hints for Lepton Flavour non-universality seen in b->sll transitions



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Explore rare be decays with unprecedented precision. Study of CP violation, the CKM matrix, ...









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A precise study of the τ - extending beyond Belle II now beginning

Observable	Current precision	FCC-ee <mark>stat</mark> .	Possible syst.
m _τ [MeV]	1776.86 ± 0.12	0.004	0.1
τ _τ [fs]	290.3 ± 0.5 fs	0.001	0.04
B(τ→eνν) [%]	17.82 ± 0.05	0.0001	
B(τ→μνν) [%]	(τ→μνν) [%] 17.39 ± 0.05		0.003







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B(τ→μνν) [%]	17.39 ± 0.05	0.0001	

N.B.: Flavour physics introduces specific detector requirements such as PID, typically not front-andcenter in Higgs Factory detector designs

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The Top Quark

A new arena at 350 GeV and above

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Overview: Top Physics at e+e- Colliders

Understanding the Top, using the Top



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Overview: Top Physics at e⁺e⁻ Colliders

Understanding the Top, using the Top



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- Measuring the top quark mass (and other parameters) in theoretically welldefined frameworks
- Search for BSM decays in clean environment





Overview: Top Physics at e⁺e⁻ Colliders

Understanding the Top, using the Top



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- Measuring the top quark mass (and other parameters) in theoretically welldefined frameworks
- Search for BSM decays in clean environment
- Electroweak couplings of the top quark as a probe for New Physics





Overview: Top Physics at e⁺e⁻ Colliders

Understanding the Top, using the Top



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The Top Quark Mass (and other parameters) Possibilities & Precision

- The accelerator side: Requires sufficient collision energy for top pair production
 - So far thoroughly studied for ILC, CLIC, threshold studies common for CLIC, FCC-ee, ILC



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Extraction of the mass in theoretically well-defined mass definition (1S, PS): can directly be used in precision calculations, minimal conversion uncertainties to MSbar mass etc.

measurement of a "MC mass": Interpretation

uncertainties of several 100 MeV











Ultimate precision at the threshold





 Exploit precise theoretical calculations of cross section in the threshold region, in well-defined mass schemes (mt^{PS}, mt^{1S}...) -> Can be converted directly into MSbar mass.



Ultimate precision at the threshold



The threshold is sensitive to top quark properties

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• Exploit precise theoretical calculations of cross section in the threshold region, in well-defined mass schemes (m_t^{PS}, $m_t^{1S}...$) -> Can be converted directly into MSbar mass.



Ultimate precision at the threshold



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Ultimate precision at the threshold



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Electroweak Couplings of the Top Quark

Access via cross section and asymmetries



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• At Linear Colliders:

- Using different beam polarisations
- Measuring cross section, A_{FB}, and helicity angle (some studies)
- Particularly powerful with two (or more) energy points





Electroweak Couplings of the Top Quark

Access via cross section and asymmetries






Searching for New Physics

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Searching for Dark Matter

• A (very) wide range of possibilities - a few obvious examples: Search for Dark Matter



Experiments at Lepton Colliders - CERN Summer Student Lectures, July 2023





Frank Simon (<u>frank.simon@kit.edu</u>)

Searching for Dark Matter

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Searching for Dark Matter

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Sensitivity depends on

- Energy reach -> Mass coverage \bullet
- Background levels: Sensitivity to small \bullet couplings







Dark Sector Searches - an FCC-ee example



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mass vs mixing² - unique phase space covered by FCC-ee



and a second						
	→					
_						
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ed vertex						
IRs						
liggs						
$\textcircled{\ } 2 \propto 1 \Theta l^2 = 10 l^2 + 10 l^2$						
ced vertex						
BRs						
Higgs						
) @ 2o:	$ \Theta ^2 = \Theta_g ^2 + \Theta_g ^2$					





Indirect and direct exploration of the highest energy scales



Corrections to SM suppressed by 1/(mass scale)² Sensitivity grows with s







Indirect and direct exploration of the highest energy scales



For many generic models & new interactions: Corrections to SM suppressed by 1/(mass scale)² Sensitivity grows with s









Indirect and direct exploration of the highest energy scales



For many generic models & new interactions: Corrections to SM suppressed by 1/(mass scale)² Sensitivity grows with s









Indirect and direct exploration of the highest energy scales









Institute for

Data Processing and Electronics

The Strength of CLIC and Muon Colliders

• Pushing limits on dark matter



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The Strength of CLIC and Muon Colliders



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The Strength of CLIC and Muon Colliders



Lepton colliders: Full collision energy available for new particles -> Sensitivity up to kinematic limit.

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Conclusions

Wrapping up

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Compelling Scientific Opportunities

- An e⁺e⁻ collider operating around 250 380 GeV will provide a model-independent, precise investigation of the Higgs sector, and studies of unprecedented precision of the top quark
- A revisit to the Z pole with much higher luminosity than LEP will enable to electroweak precision tests of the Standard Model at completely new levels. At the same time, this will also be a high-statistics flavour physics program.
- Scales in the TeV region and above can directly be probed by high-energy lepton colliders CLIC, a (multi-)TeV ILC, and a muon collider. This also includes the measurement of the self-coupling of the Higgs.







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CERN is currently studying the feasibility of the Future Circular Collider:

- An e⁺e⁻ machine running from the Z-pole up to 365 GeV precision Higgs, Top, Electroweak.
- Followed by a \sim 100 TeV hadron collider exploration of the highest energy scales, measurement of the self-coupling of the Higgs.
- **CLIC** is studied as "Option B" in case FCC cannot go forward.

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The Way Forward

Strategies and Timescales - taken from last year's Snowmass Meeting





 Indicative timelines as discussed July '22 in Snowmass @ Seattle

resource realism varies - most developed for CERN projects







The Way Forward

Strategies and Timescales - taken from last year's Snowmass Meeting







some of them!

happen.

This will be *your* HEP facility!

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There are very exciting questions in high energy physics - a new e+e- collider may answer

Global large projects = long time scales - but contributions are needed now to make them





Extras Lecture 2

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FCC-ee Time Line to Physics

Making a new facility happen

FC	C-ee p	ohysics	srun
	2047 -	- 2047	
Start accelerator commissioning	2046 -	- 2046	Start detector commissioning
Start accelerator commissioning	2045 -	- 2043	
	2044 -	- 2044	2
	2042 -	- 2042	
End of HL-LHC operation	2041 -	- 2041	Start detector installation
Start accelerator installation	2040 -	_ 2040)
	2039 –	– 2039	3
	2038 –	– 2038	3
	2037 –	– 2037	7
Start accelerator component production	2036 –	- 2036	Start detector component production
Technical design & prototyping completed	2035 -	- 2035	Four detector TDRs completed
	2034 –	– 2034	+
	2033 –	- 2033	3
Ground-breaking and start civil engineering	2032 -	- 2032	2
Start engineering design	2031 -	- 2031	Detector CDRs (>4) submitted to FC ³
Completion of ULLUC: more ATS personnel available	2029 -	- 2030	Completion of ULLUC ungrade, more detector experts available
ECC Approval B&D start prototyping	2028 -	2028	EC ³ formation call for CDPs, callaboration forming
FCC Approval, K&D, start prototyping	2027 –	- 2027	FC Tormation, can for CDRS, conaboration forming
	2026 –	- 2026	European Strategy Update
FCC Feasibility Study Report	2025 -	- 2025	Detector EoI submission by the community
FCC-ee Accelerator	Key	dates	FCC-ee Detectors
ts at Lepton Colliders - CERN Summer Student Lect	ures July 2	2023	Frank Simon (frank simon@kit edu)

