Ising-like models on Euclidean black holes arXiv:2306.08547

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• The Ising model is a lattice of spins on a plane. Temperature is an external parameter. The effect of temperature on the order of spins is known.



• The transition between perfect order to disorder occurs at the critical temperature, $T_c = 2.27$.

- "General relativity is a theory of space, time and gravitation formulated by Einstein in 1915." [1]
- The gravitational field of a black hole (aka Lorentzian Schwarzschild spacetime) is a solution to Einstein's equations. The gravitational field of a black hole (with mass *M* and surface area *a*) satisfies [2]

$$dM = \left(\frac{1}{8\pi M}\right) \left(\frac{da}{4}\right).$$

• This resembles the first law of thermodynamics

$$dE = T dS$$

where E is energy, T is temperature and S is entropy.

- This correspondence motivated Gibbons and Hawking to treat black hole spacetime as a statistical mechanics system. They succeeded by using Euclidean Schwarzschild geometry [3].
- We asked the question: "Is there a way to construct an Ising-like model on this thermal background? If there is, then how do the spins behave as *M* is varied?"

- 1. Ising model
- 2. Euclidean black holes
- 3. Ising-like models on Euclidean black holes

Ising model

• Lattice of spins on a 2d plane in contact with a heat bath.

$$\begin{array}{c|c} & & \\ \uparrow & \downarrow & \uparrow \\ \mathfrak{f}_{m+1,n-1} & \mathfrak{f}_{m+1,n} & \mathfrak{f}_{m+1,n+1} \\ \downarrow & \uparrow & \uparrow \\ \mathfrak{f}_{m,n-1} & \mathfrak{f}_{m,n} & \mathfrak{f}_{m,n+1} \\ \mathfrak{f}_{m+1,n-1} & \mathfrak{f}_{m+1,n} & \mathfrak{f}_{m+n+1} \end{array}$$

Hamiltonian

$$H = -\sum_{m,n=1}^{N} s_{m,n} (s_{m+1,n} + s_{m,n+1}).$$

• A configuration of spins x has weight

$$p_x = rac{\exp\left(-rac{H_x}{T}
ight)}{Z}, \quad Z(T) = \sum_x \exp\left(-rac{H_x}{T}
ight).$$

Alignment

$$A=\frac{1}{N^2}\left|\sum_{m,n=1}^N s_{m,n}\right|.$$

• Statistical average of alignment

$$< A > = \sum_{x} A_{x} p_{x}.$$

• Phase transition of the Ising model occurs at $T_c = 2.27$. This is an exact result due to Onsager [4].



Used the Monte Carlo (MC) algorithm to calculate < A > for temperatures selected randomly from $T \in (0,5]$.



Summary: The 2d Ising model exhibits a phase transition. We reproduced the known result using numerical methods.

Euclidean black holes

• Spacetime event $\equiv (t, r, \theta, \phi)$.



• Distances in spacetime:

$$ds^2 = g_{ab}(t, \vec{x}) dx^a dx^b.$$

 $g_{ab}(t, \vec{x})$ is the metric.

- Einstein's theory of general relativity (GR) describes the interactions of matter and spacetime.
- Tenets of the theory:
 - Gravity is the curvature of spacetime.
 - John Wheeler: "Spacetime tells matter how to move; matter tells spacetime how to curve."



Schwarzschild spacetime:

$$ds^{2} = -\left(1 - 2Mr^{-1}\right)dt^{2} + \left(1 - 2Mr^{-1}\right)^{-1}dr^{2} + r^{2}d\Omega_{(2)}^{2}.$$

Here *M* is the mass of the black hole, $d\Omega_{(2)}^2$ is the metric describing a two-sphere, $t \in \mathbb{R}$, $r \in \mathbb{R}^+$; the spacetime has coordinate and physical singularities at r = 2M (the horizon) and r = 0 respectively.

• The spacetime has temperature

$$T=(8\pi M)^{-1}.$$

- Motivated by the thermodynamic properties of the black hole spacetime, Gibbons and Hawking tried to treat it as a statistical mechanics system. They succeeded by using Euclidean Schwarzschild background [3].
- To derive the Euclidean Schwarzschild space transform to Euclidean time $\tau = it$ and a new radial coordinate $\rho(r, M)$ and enforce Euclidean time to have periodicity $8\pi M$ [1]. The $\tau \rho$ part of the metric:

$$ds^2 = \rho^2 \left(\frac{d\tau}{4M}\right)^2 + \left[\frac{r(\rho)}{2M}\right]^4 d\rho^2, \quad \tau = [0, 8\pi M), \quad \rho = [0, 4M).$$

Summary: Euclidean black hole geometry is used to understand black hole thermodynamics; it can be derived from Lorentzian Schwarzschild spacetime through the application of a procedure which requires Euclidean time to have periodicity $8\pi M$ [1].

Ising-like models on Euclidean black holes

• Action of a scalar field on Euclidean background:

$$S = \frac{1}{2} \int d^4x \sqrt{g} \left(g^{ab} \partial_a \Phi \partial_b \Phi \right).$$

Choose Euclidean Scwarzschild background and $\Phi \equiv \Phi(\tau, \rho)$. *M* determines coordinate ranges and is present in the integrand.

 Discretize background into a N_τ × N_ρ lattice with lattice spacing ε (first arrow) and restrict the value of the field to ±1 (second arrow).



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 The action is the Ising-like model on Euclidean Schwarzschild background:

$$S = -4\pi \sum_{m,n=1}^{N_{\tau},N_{\rho}} \left\{ \Phi_{m,n} \left[\frac{16M^3}{\rho_n \left(1 - \left(\frac{\rho_n}{4M}\right)^2 \right)^4} \Phi_{m+1,n} + (\rho_n M) \Phi_{m,n+1} \right] \right\}.$$

- Features of this inhomogeneous Ising model:
 - Each interaction is a function of space.
 - M
- Influences the vertical and horizontal interactions differently,
- Determines lattice size.
- The spins are at temperature $T = (8\pi M)^{-1}$. But there is no heat bath.

The background has temperature.

• For fixed *M* a lattice configuration - *x* - has weight

$$P(x) = \frac{\exp(-S(x))}{Z}, \quad Z = \sum_{x} \exp(-S(x))$$

 Used Monte Carlo simulations (with ε = 0.01) to study < A > by varying M. Numerical results indicate a second order phase transition.



Summary: We constructed an Ising-like model on Euclidean black hole background, investigated the thermodynamic properties of spins by varying M and found evidence of a phase transition at sub-Planckian mass.

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