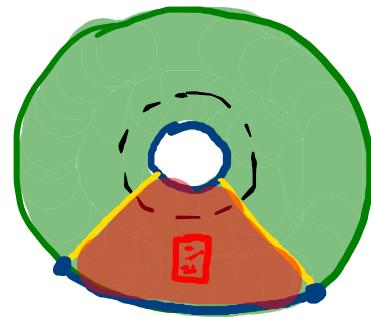
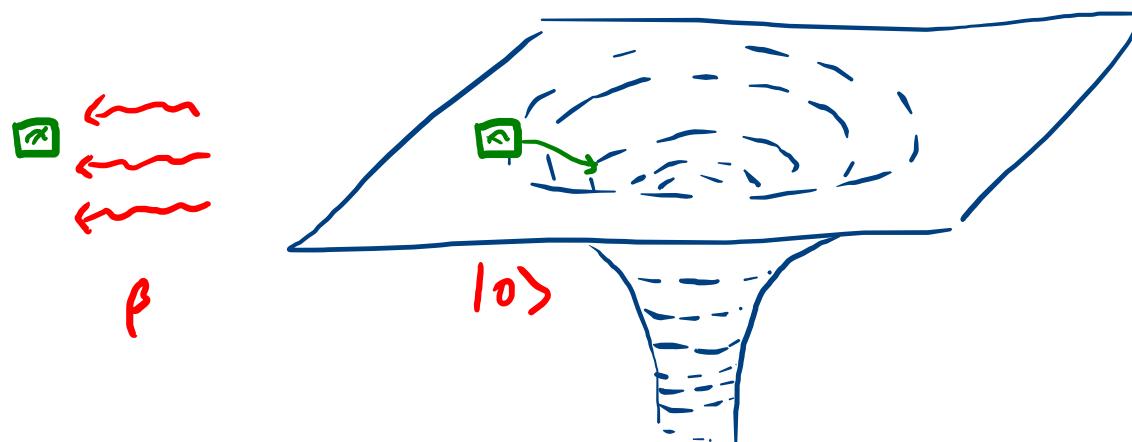


# HOW TO RECOVER YOUR HOMEWORK FROM A BLACK HOLE



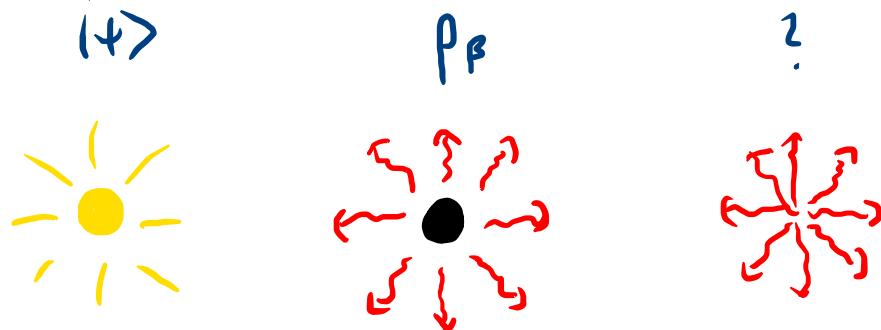
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UBC → Xanadu  
Theory Canada, 2023

- Alice takes a course on QFT in curved spacetime. Her midterm goes over the standard (Bogoliubov) derivation of Hawking radiation.



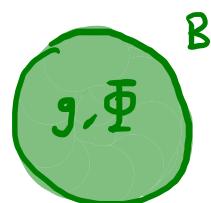
- An infalling observer  sees vacuum  $|0\rangle$ .
- A stationary observer  sees a thermal spectrum at  $\beta \propto M$ .

- This suggests information is turned into thermal noise.



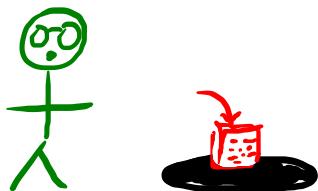
- Alice's midterm Euclidean hack ignored metric fluctuations. But she learns a for the calculation which sums over metrics:

$$\mathcal{Z} = \int_{\partial M = \beta} Dg D\Phi e^{-I(g, \Phi)}$$



- The entropy  $S = (1 - \beta \rho_\beta) \ln \mathcal{Z} \neq 0$ , so even with fluctuations the state is mixed. Semiclassical fluctuations don't help!

- Alice takes a shortcut through a condensed matter lab to give her midterm to the TA. She passes an **artificial black hole**.

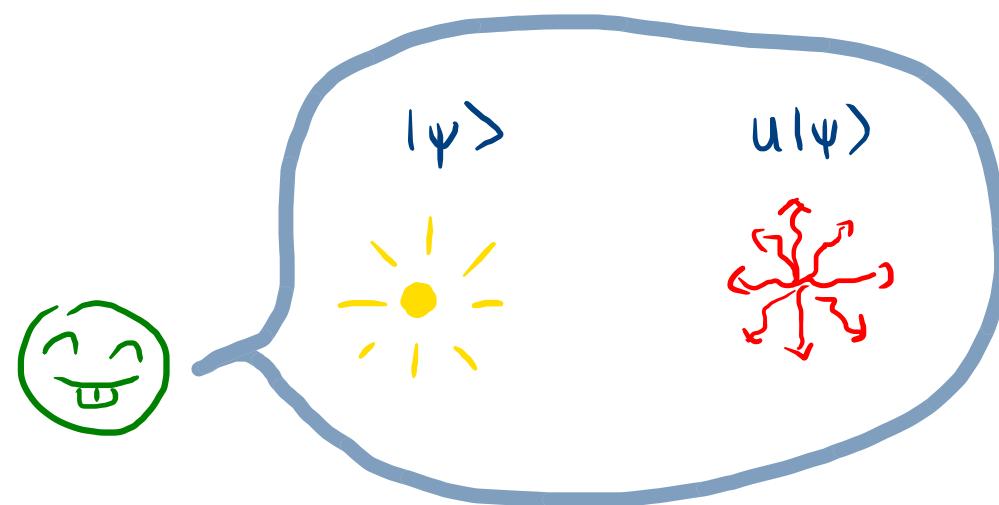


- Distracted, she accidentally drops her midterm in! Her course has become unexpectedly (and self-referentially) relevant.

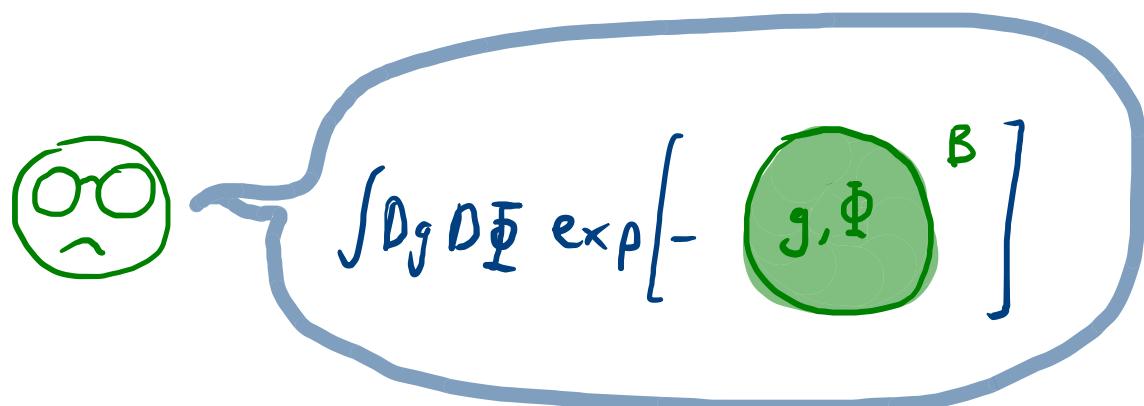


- From Hawking's calculation, Alice thinks it is lost forever. **Bolt disagrees!**

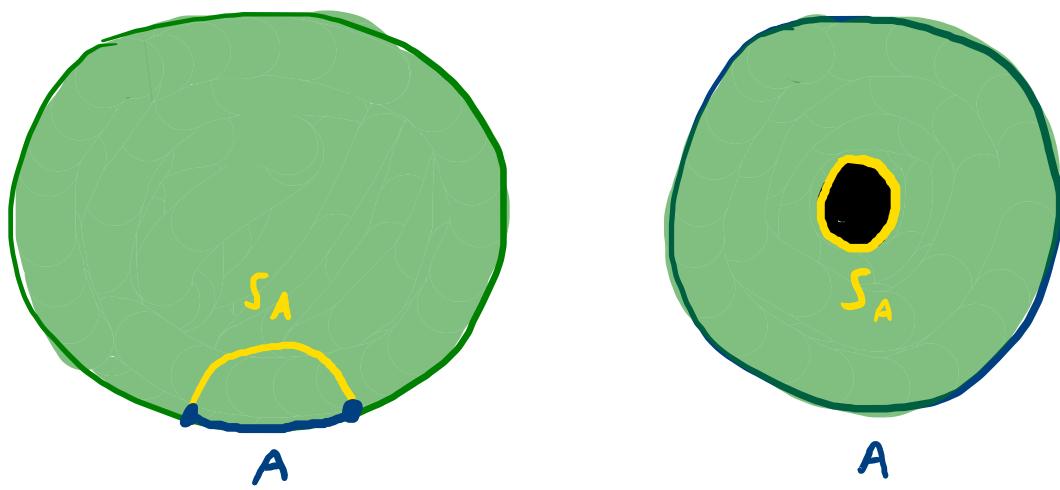
- "It can't be gone," he argues. "Nature is unitary. It only looks thermal."



- Alice wants to understand what goes wrong in the entropy calculation.

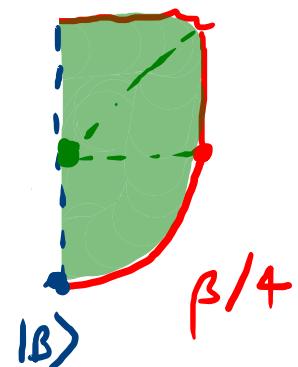
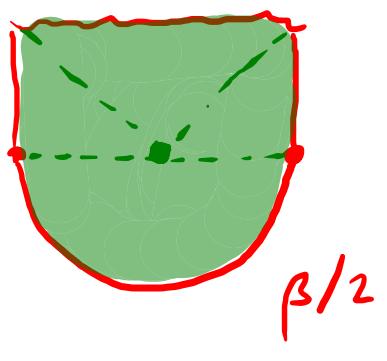


- "The partition function gives **coarse-grained entropy**. We need to compute **fine-grained quantum entropy**," Bob argues. "But how?" Alice asks.
- "Well, in AdS/CFT, **fine-grained entropy** of a boundary region is computed by the area of a **homologous minimal surface**."



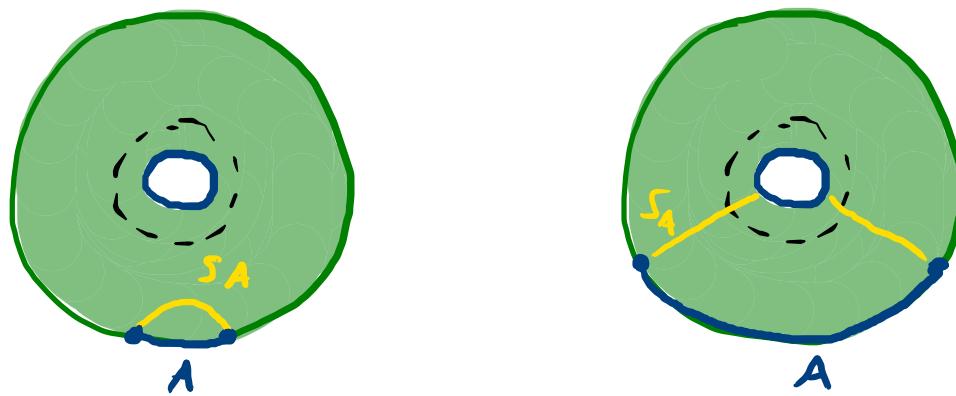
- Alice immediately notices a problem. "In Schwarzschild, **the horizon** is the minimal surface homologous to the whole boundary."

- Bolt pauses, then suggests; "The Schwarzschild geometry must be the result of averaging over geometries. We need pure state geometries!"
- In the Euclidean formalism, black holes are periodic in imaginary time.



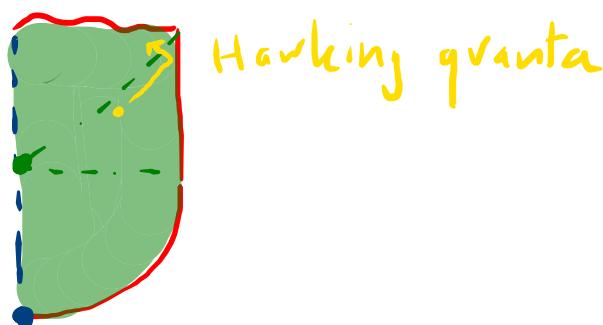
- Averaging over the half-disk prepares a Schwarzschild black hole. We can fold in half and impose boundary conditions.

- The fold terminates spacetime in an end-of-the-world (Eow) brane. These offer an alternative place for minimal surfaces to end!

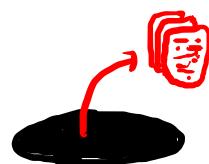


- Since the brane doesn't contribute to the minimal surface, the empty surface is now homologous to the whole boundary. The state is pure!
- A neat thing about AdS/CFT is that you can explicitly check this weird rule agrees with a microscopic calculation of entropy.

- To recap, we have a family of black hole pure states with branes, a way to compute fine-grained entropy of subregions, & consistency checks.

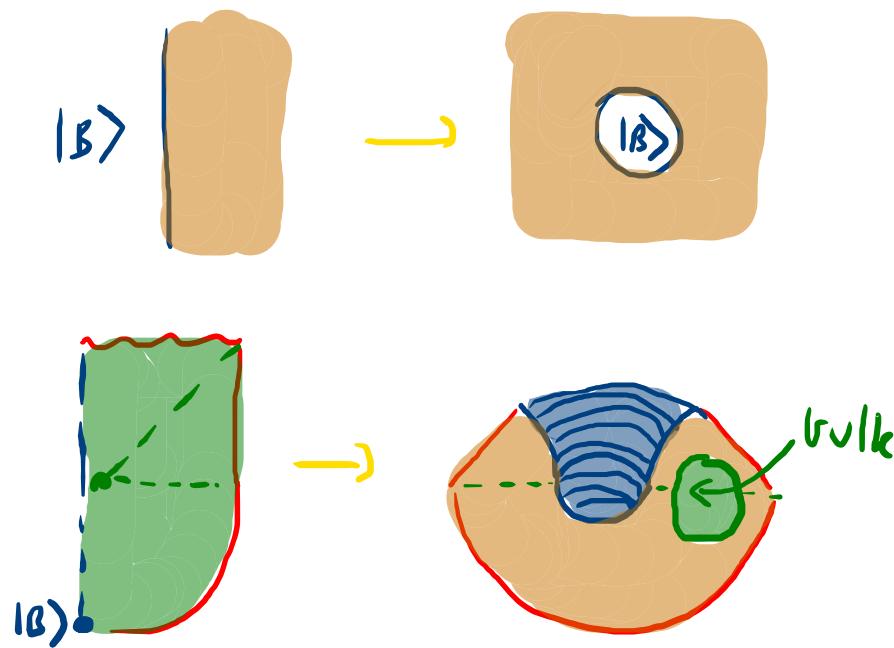


- This is all well and good, but as Alice points out to Bob:  
"The Hawking radiation bounces back, so there's no evaporation."



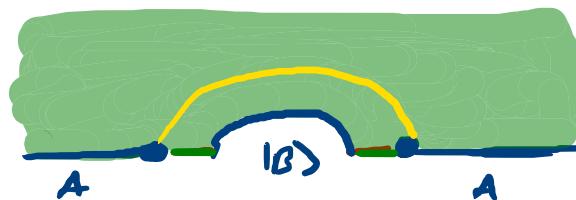
- Alice still needs her midterm!

- However, by a sneaky change of coordinates, Alice and Bob see how to make the spacetime dynamic.
- On the boundary the fold is a line. We turn it into a circle.

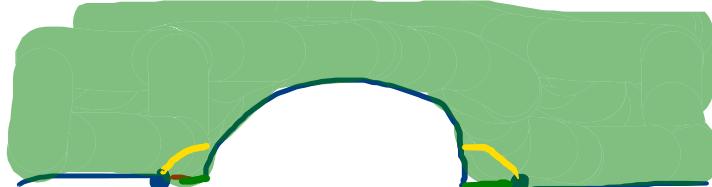


- This turns a static bulk black hole with spherical brane into a dynamic boundary black hole with an accelerating Rindler brane.

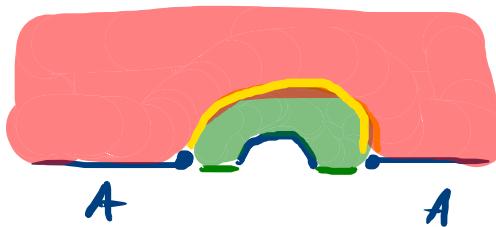
- The time-dependent behaviour of subregion entropy is of most interest.
- At early times, the minimal surface for a neighbourhood of infinity (analogous to an asymptotic particle detector) is connected.



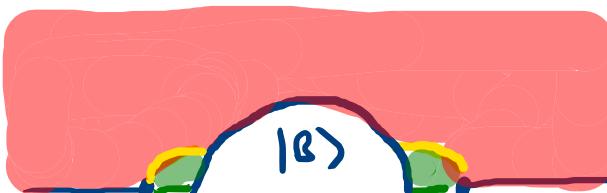
- At late times, it transitions to a disconnected surface. This looks like information escaping from the black hole!



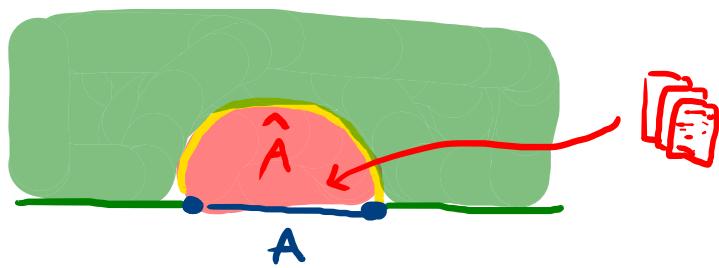
- In fact, this transition is exactly what's needed to show information is conserved, from the viewpoint of this lower-dimensional black hole.
- The pure state information is in the boundary condition  $|B\rangle$ . The green region is unknown to an observer with access to A.



- After the transition, the red region pinches off, and part of thebrane becomes known. This is sufficient to determine  $|B\rangle$ !

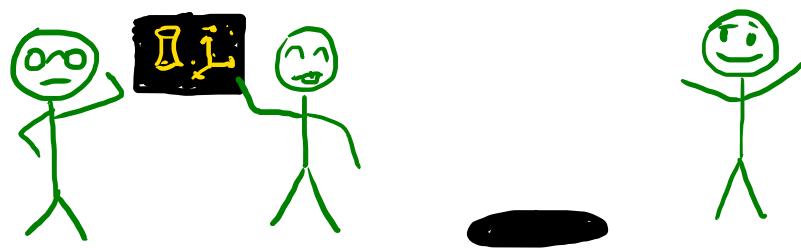


- But what does it mean to have "access" to a bulk region? And how does Alice finally get her midterm?
- The basic insight is that a boundary subregion  $A$  contains the same information as the bulk region  $\hat{A}$  between  $A$  and its minimal surface.



- The region  $\hat{A}$  is called the entanglement wedge, and in principle we can reconstruct midterms in  $\hat{A}$  from knowledge of  $A$ .

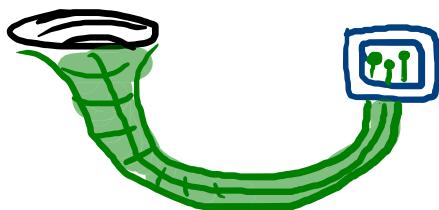
- At this point in their deliberations, Charlotte joins Alice and Bob. She is the **architect** of the artificial black hole.



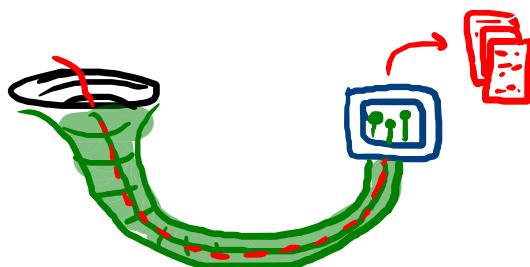
- "What are you kvetching about?" They explain their plan to use **entanglement wedge reconstruction** to recover the midterm.



- Charlotte sighs. "First, the black hole isn't in a boundary state. It's entangled with my quantum computer. It's a wormhole!"



- "Usually, that wormhole is non-traversable but luckily I was testing a double-trace deformation library which made it briefly traversable."



- "Your midtem popped up in my quantum inbox!"

- Alice breathes a sigh of relief. Her CPA is out of danger! Still, she regrets not being able to use her newfound knowledge.



- "I wouldn't worry too much," says Charlotte. "Sounds like you have the material for a pretty cool thesis!"



THANKS!