

Flauber - physics

Alexander Lenz

Siegen University

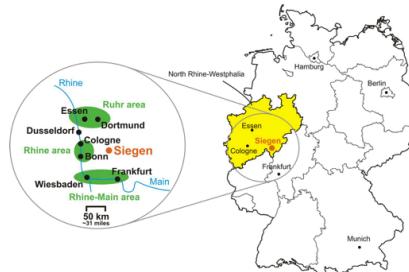
13./14./15. 2. 2023

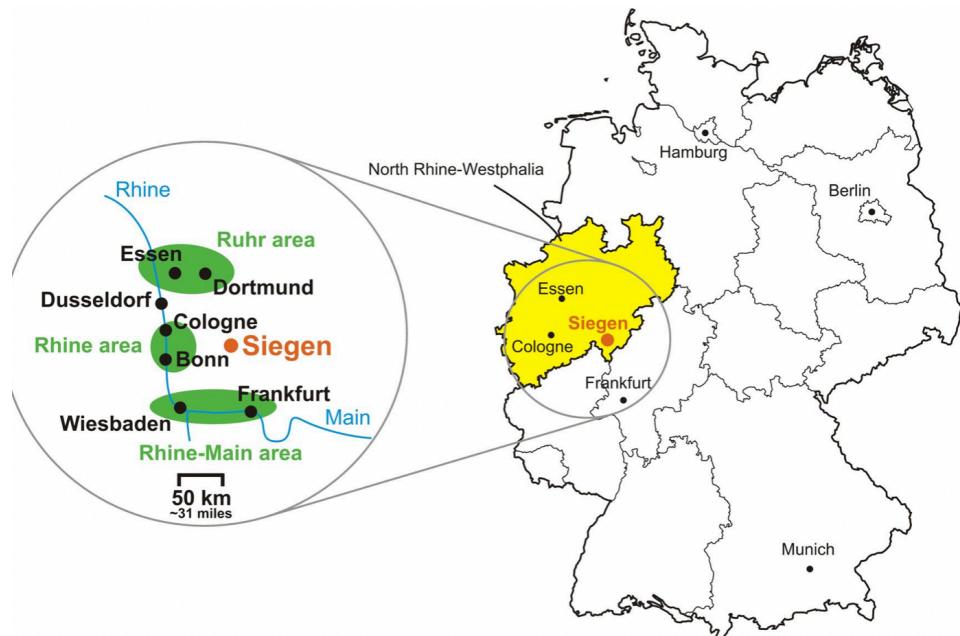
Siegen University

- 16 700 Students

Physics:

- 18 Professors
- approx. 65 post-docs
- Approx. 30 PhD students
- Each year around 30 first-year students





Lecture 1: SM, CKM, weak decay

Lecture 2: Theoretical Framework

Lecture 3: Mixing & CP

Lecture 1:

Unitarity triangle

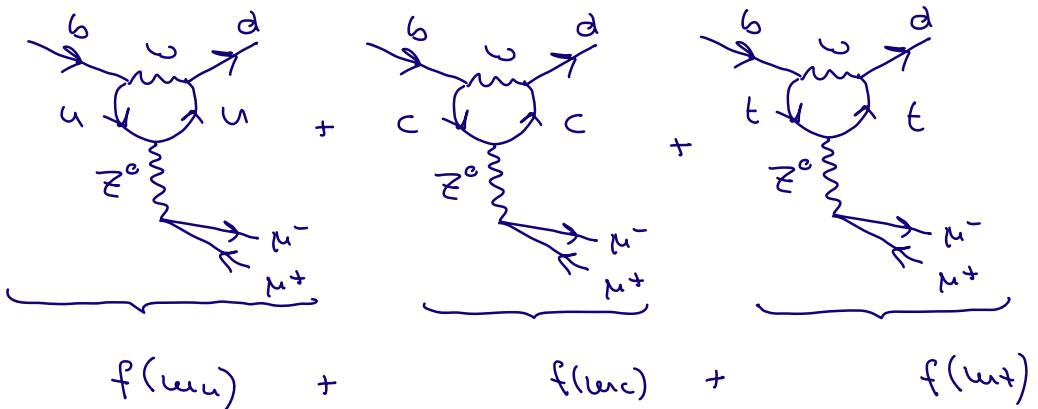
V_{CKM} is by construction unitary

$$V_{CKM}^\dagger V_{CKM} = \sum_{q_3=u,d,s} V_{q_3 d_1}^* V_{q_3 d_2} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

$$V_{CKM}^\dagger V_{CKM}^L = \sum_{q_3=d,s,b} V_{u_1 q_3}^* V_{u_2 q_3}^* = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\text{Bd: } 0 = V_{ub} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^*$$

* $b \rightarrow d \mu \bar{\mu}$ - Penguin



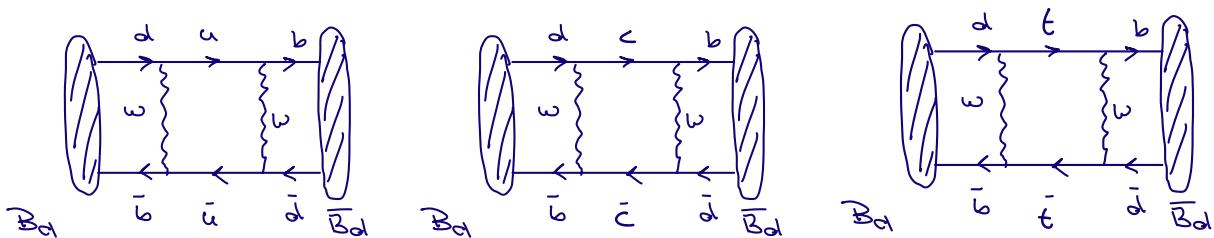
Ann 1: $u\bar{u} = c\bar{c} = t\bar{t}$

$\Rightarrow f(u\bar{u}) \{ V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} \} = 0$

GIM Mechanismus
Glashow
Illiopoulos

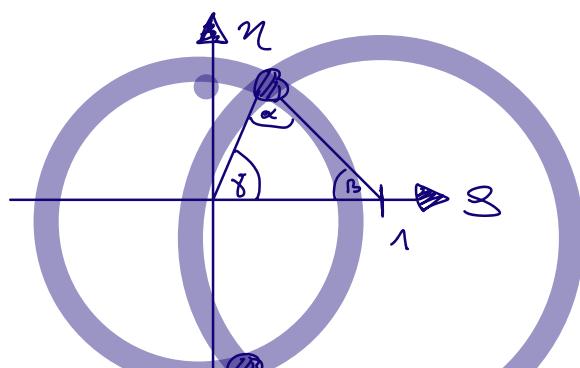
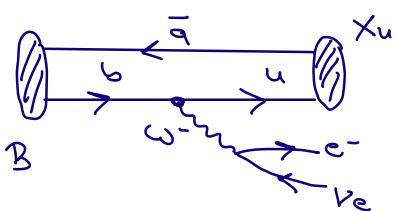
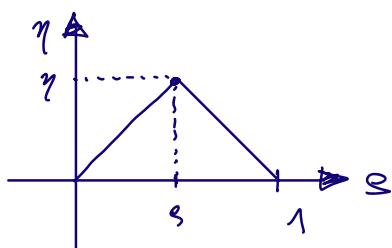
Ann 2: $f(c\bar{c}) = f_0 + \tilde{f}(c\bar{c})$
 $\tilde{f}(c\bar{c})$
 hebt sich weg

* B-Mischung: $B_d \leftrightarrow \bar{B}_d$

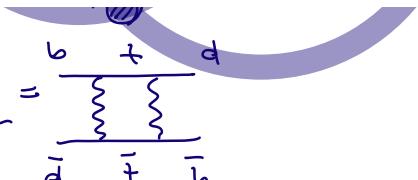


wieder $(V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td})^2$

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$$

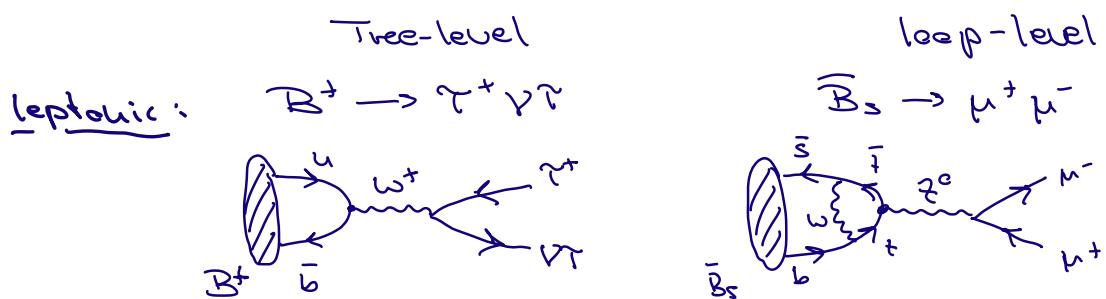


- $$\Delta m_d = \Gamma_{B_H} - \Gamma_{B_L} = \frac{b}{d} + \frac{d}{b}$$

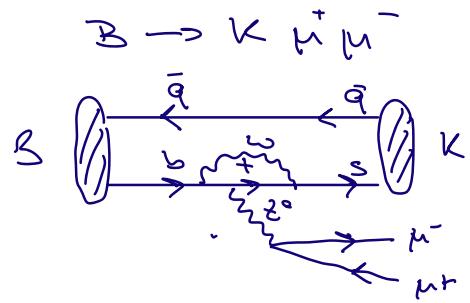
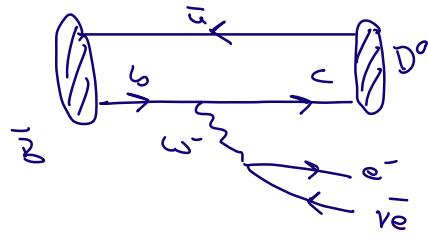


Exp:

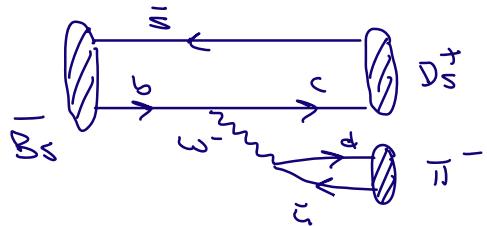
Classification of decays



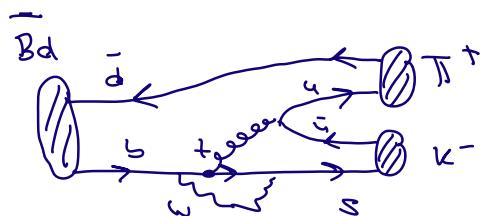
semi-leptonic: $\bar{B} \rightarrow D^0 e^- \bar{\nu}_e$



non-leptonic: $\bar{B}_s \rightarrow D_s^+ \pi^-$



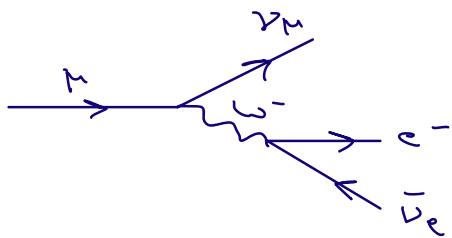
$\bar{B}_d \rightarrow \pi^+ K^-$



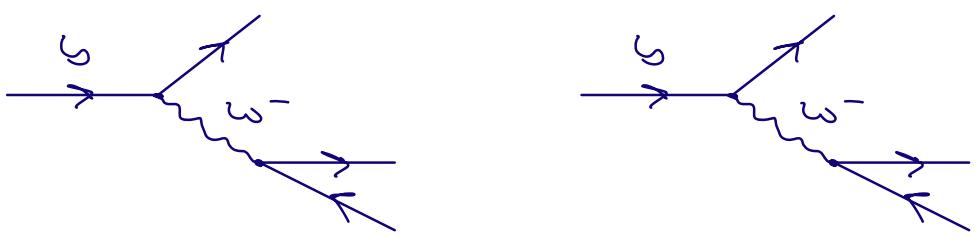
Lecture 2:

A. Buras

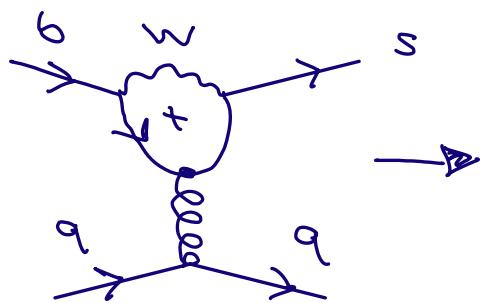
- hep-ph/9806471
Les Houches
- hep-ph/9512380
Review







Penguin operators



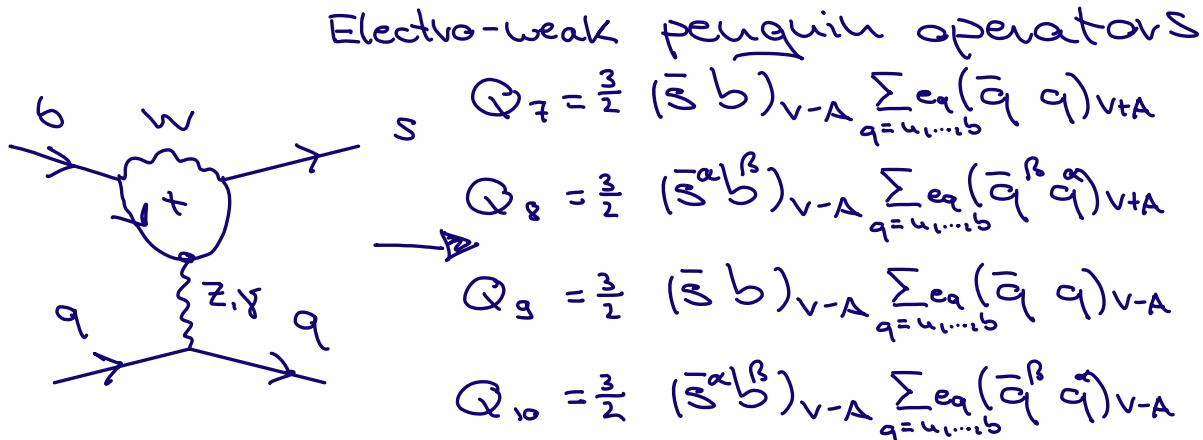
QCD penguin operators

$$Q_3 = (\bar{s} b)_{V-A} \sum_{q=u,d,s,b} (\bar{q} q)_{V-A}$$

$$Q_4 = (\bar{s}^\alpha b^\beta)_{V-A} \sum_{q=u,d,s,b} (\bar{q}^\beta q^\alpha)_{V-A}$$

$$Q_5 = (\bar{s} b)_{V-A} \sum_{q=u,d,s,b} (\bar{q} q)_{V+A}$$

$$Q_6 = (\bar{s}^\alpha b^\beta)_{V-A} \sum_{q=u,d,s,b} (\bar{q}^\beta q^\alpha)_{V+A}$$



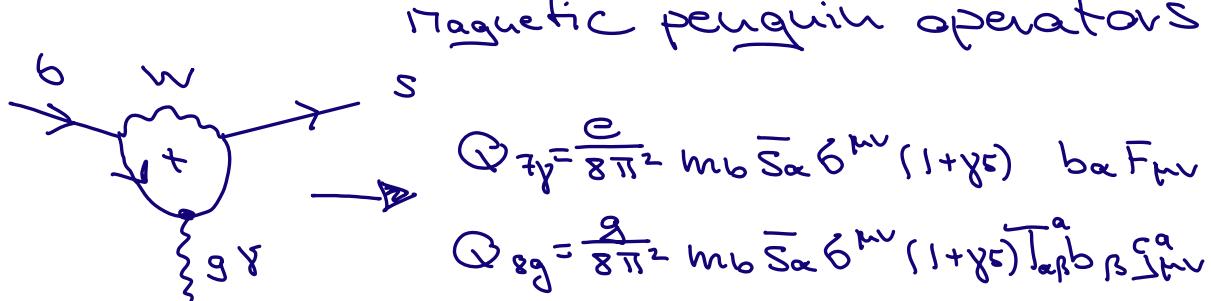
Electro-weak penguin operators

$$Q_7 = \frac{3}{2} (\bar{s} b)_{V-A} \sum_{q=u,d,s,b} e_q (\bar{q} q)_{V+A}$$

$$Q_8 = \frac{3}{2} (\bar{s}^\alpha b^\beta)_{V-A} \sum_{q=u,d,s,b} e_q (\bar{q}^\beta q^\alpha)_{V+A}$$

$$Q_9 = \frac{3}{2} (\bar{s} b)_{V-A} \sum_{q=u,d,s,b} e_q (\bar{q} q)_{V-A}$$

$$Q_{10} = \frac{3}{2} (\bar{s}^\alpha b^\beta)_{V-A} \sum_{q=u,d,s,b} e_q (\bar{q}^\beta q^\alpha)_{V-A}$$

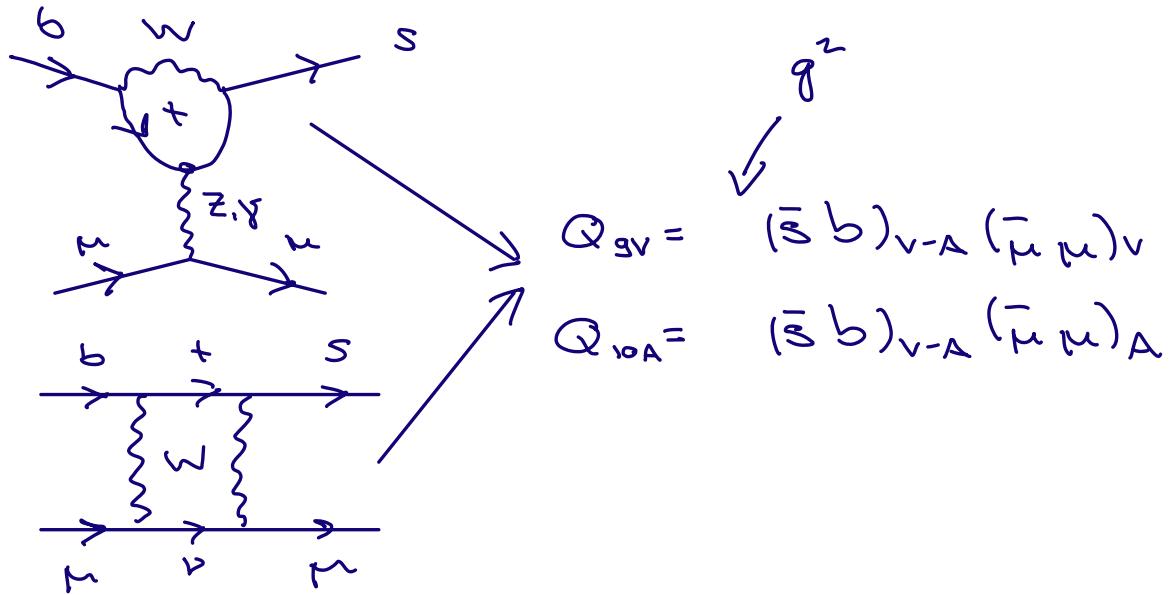


Magnetic penguin operators

$$Q_{7g} = \frac{e}{8\pi^2} m_b \bar{s}_\alpha \delta^{\mu\nu} (1 + \gamma_5) b_\alpha F_{\mu\nu}$$

$$Q_{8g} = \frac{g}{8\pi^2} m_b \bar{s}_\alpha \delta^{\mu\nu} (1 + \gamma_5) \tilde{T}_{\alpha\beta}^a b_\beta \tilde{c}_\nu^a$$

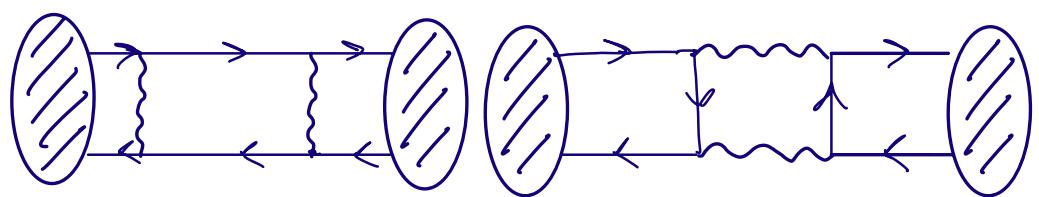
Semileptonic penguin operators

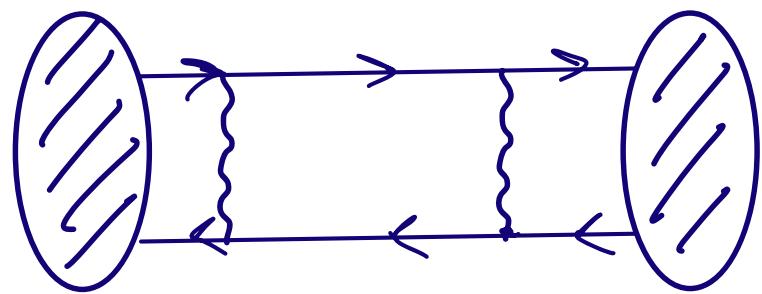


$$\mathcal{L}_{\text{left}} = \frac{G_F}{\sqrt{2}} \left[\sum_{q=u,c} V_q \left\{ C_1(\mu) Q_1^{q'q''} + C_2(\mu) Q_2^{q'q''} \right\} - V_P \sum_{j \geq 3} C_j(\mu) Q_j \right]$$

$q' = u, c$
 $q'' = d, s$

Lecture 3: Mixing & SP in the B-system





$$B_{S,H} = p B_S + q \frac{1}{B_S}$$
$$B_{S,L} = p B_S - q \frac{1}{B_S}$$

$$\Delta \pi_s = 2 |\pi_{12}^s| + O\left(\left|\frac{\pi_2}{\pi_{12}}\right|^2\right)$$

$$\Delta \tau_s = 2 |\tau_{12}^s| \cos(\phi_{12}^s) + O\left(\left|\frac{\pi_2}{\pi_{12}}\right|^2\right)$$

$$\phi_{12}^s = \arg\left(-\frac{\pi_{12}^s}{\tau_{12}^s}\right)$$

$$\Rightarrow g_+(t) = e^{-i\Gamma_{Bst}} e^{-\frac{\Gamma_{Bst}}{2}} \left[\begin{aligned} & \cosh \frac{\Delta\Gamma_{st}}{4} \cdot \cos \frac{\Delta\gamma_{st}}{2} \\ & - \sinh \frac{\Delta\Gamma_{st}}{4} \cdot \sin \frac{\Delta\gamma_{st}}{2} \end{aligned} \right]$$

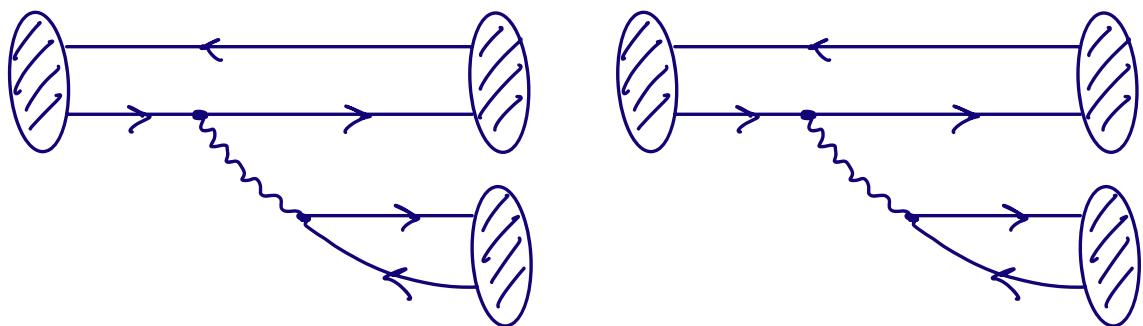
$$g_-(t) \equiv \quad \text{see e.g. 1511.09466}$$

$$\Gamma(\bar{B}_s(t) \rightarrow f) = n_f |A_f|^2 \frac{1 + |\lambda_f|^2}{2} e^{-\Gamma_{st}} \left\{ \begin{aligned} & \cosh \frac{\Delta\Gamma_{st}}{2} - \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2} \cos \Delta\gamma_{st} \\ & - \sinh \frac{\Delta\Gamma_{st}}{2} \frac{2\operatorname{Re}(\lambda_f)}{1 + |\lambda_f|^2} + \frac{2\operatorname{Im}(\lambda_f)}{1 + |\lambda_f|^2} \sin \Delta\gamma_{st} \end{aligned} \right\} (1 + \alpha_{fs}^s)$$

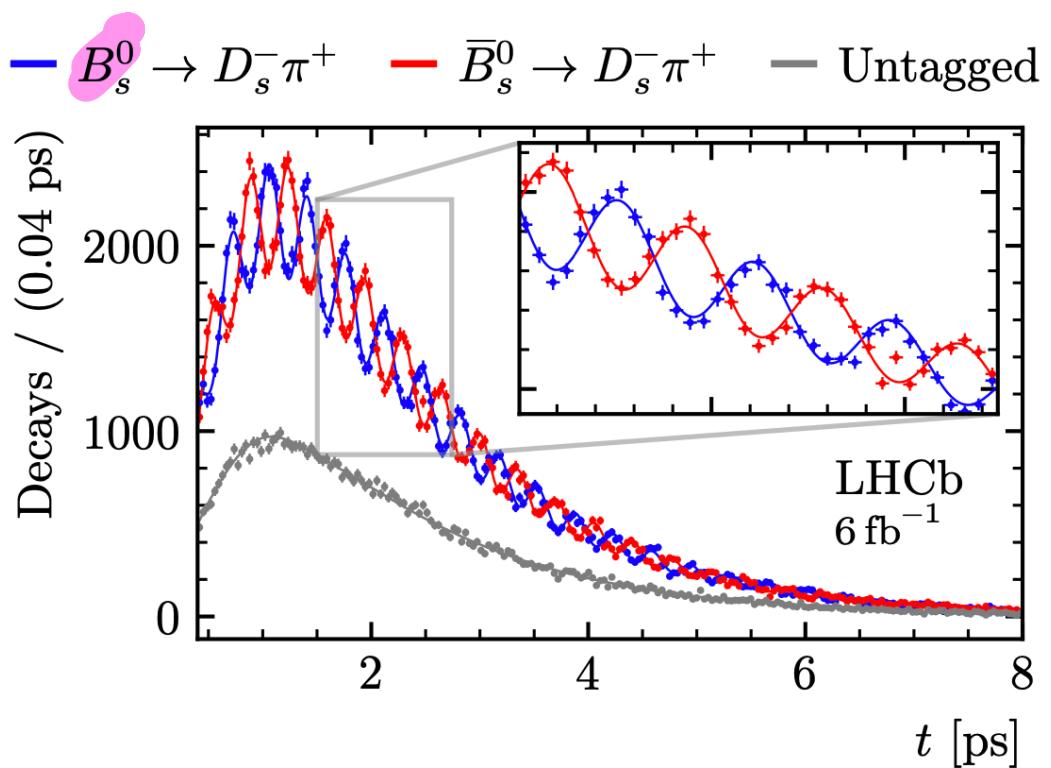
$$\mathcal{A}_f = \langle f | \gamma_{\text{eff}} | B_s \rangle$$

$$\bar{\mathcal{A}}_f = \langle f | \gamma_{\text{eff}} | \bar{B}_s \rangle \quad \bar{\mathcal{A}}_f = \langle f | \gamma_{\text{eff}} | \bar{B}_q \rangle$$

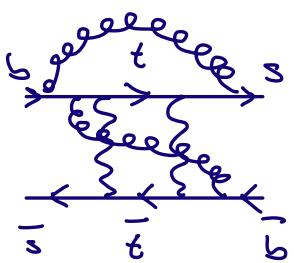
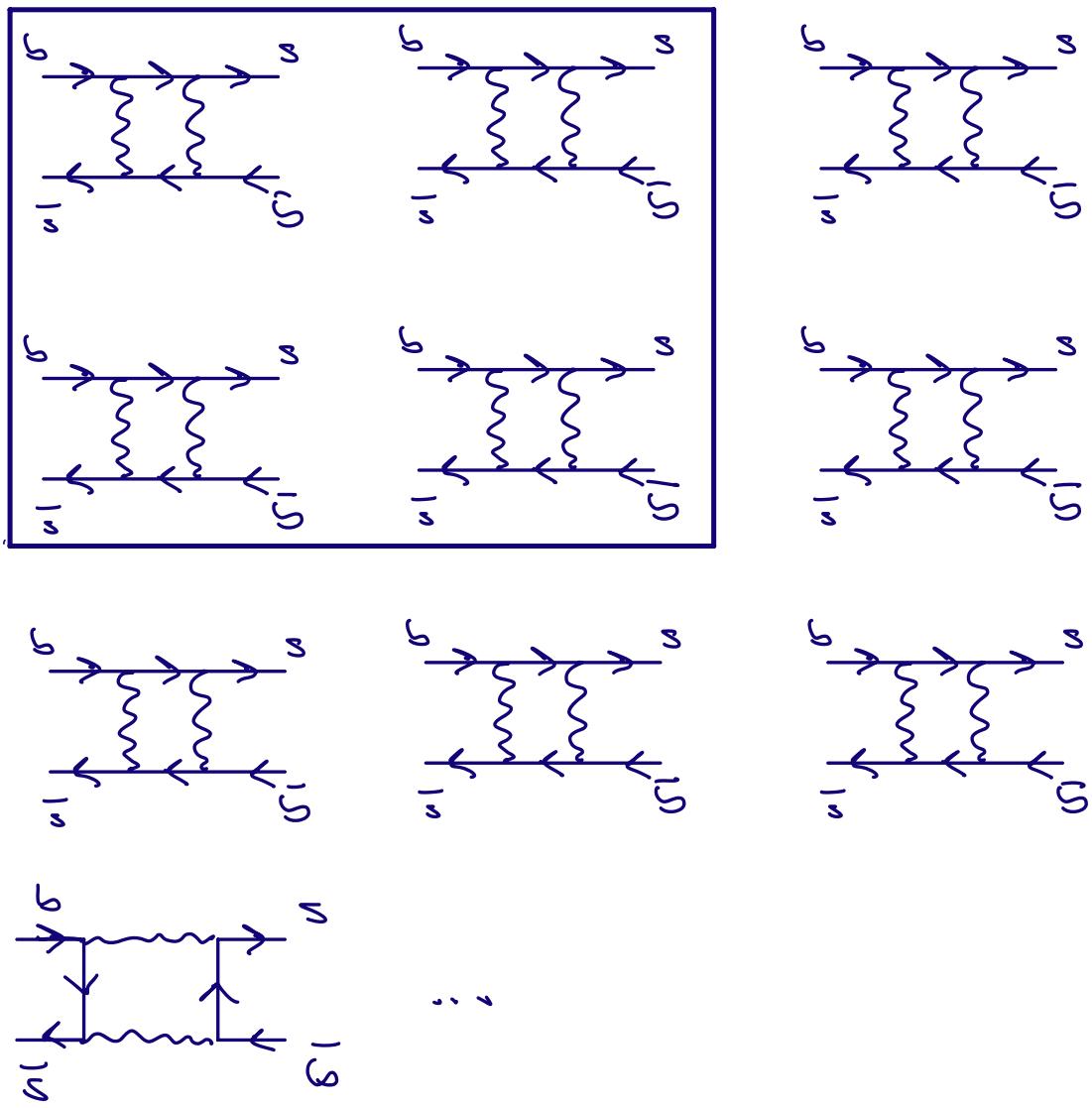
$$\lambda_f = \frac{g}{\rho} \frac{\bar{\mathcal{A}}_f}{\mathcal{A}_f}$$



How to measure $\Delta\tau$?



How to calculate $\Delta\tau$?



$$M_{12} = \lambda_u^2 F(u,u) + \lambda_u \lambda_c F(u,c) + \lambda_u \lambda_t F(u,t) \\ + \lambda_c \lambda_u F(c,u) + \lambda_c^2 F(c,c) + \lambda_c \lambda_t F(c,t) \\ + \lambda_t \lambda_u F(t,u) + \lambda_t \lambda_c F(t,c) + \lambda_t^2 F(t,t)$$

CKM:

$$\lambda_u = V_{us}^* V_{ub}$$

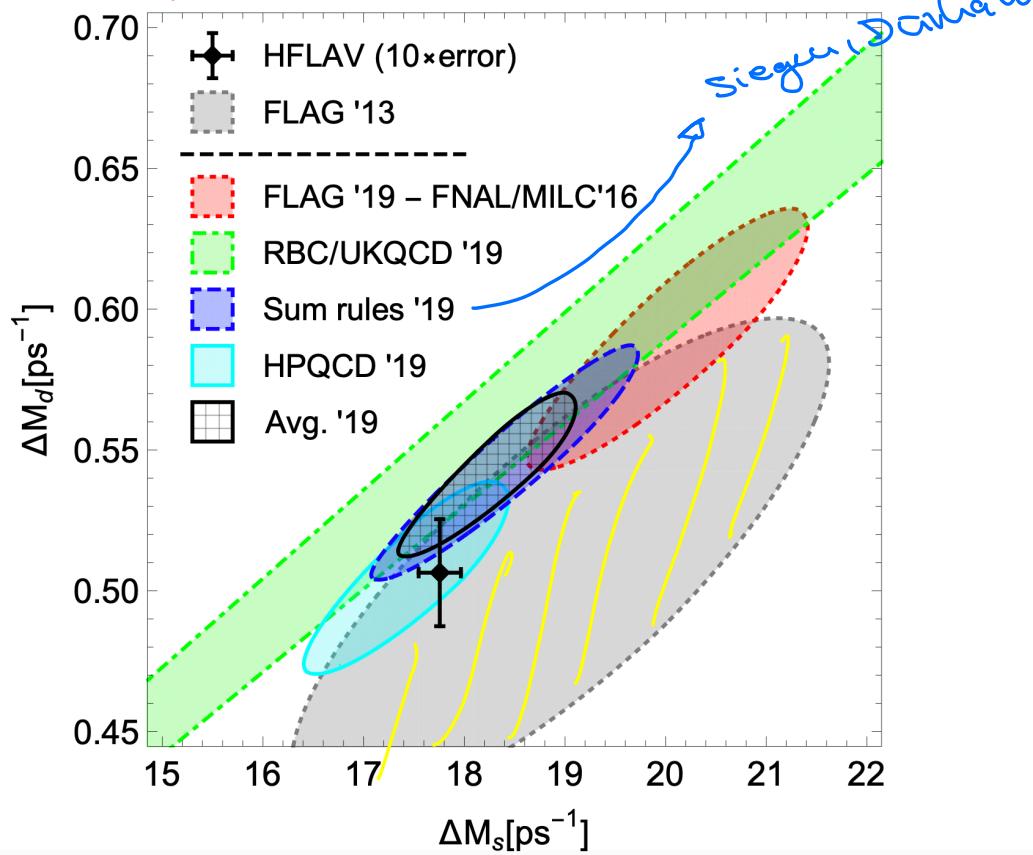
$$\lambda_c = V_{cs}^* V_{cb}$$

$$\lambda_t = V_{ts}^* V_{tb}$$

$$\Rightarrow M_{12} = \lambda_u^2 [F(c,c) - 2F(u,c) + F(u,u)]_1 \\ + 2\lambda_u \lambda_t [F(c,c) - F(u,c) + F(u,t) - F(c,t)]_2 \\ + \lambda_t^2 [F(c,c) - 2F(c,t) + F(t,t)]_3$$

$$S(x) = \frac{4x - (1+x^2+x^3)}{4(1-x)^2} - \frac{3 \times \ln x}{2(1-x)^2}$$

Average lattice & sum rules + 1909.11087



$$\Delta M_s^{\text{theory}} = (18.4 \pm 0.7) \text{ ps}^{-1}$$

1909.11087

High Energy Physics – Phenomenology

[Submitted on 18 Dec 2017 (v1), last revised 15 May 2018 (this version, v2)]

One constraint to kill them all?

Luca Di Luzio, Matthew Kirk, Alexander Lenz

Many new physics models that explain the intriguing anomalies in the b -quark flavour sector are severely constrained by B_s -mixing, for which the Standard Model prediction and experiment agreed well until recently. The most recent FLAG average of lattice results for the non-perturbative matrix elements points, however, in the direction of a small discrepancy in this observable. Using up-to-date inputs from standard sources such as PDG, FLAG and one of the two leading CKM fitting groups to determine ΔM_s^{SM} , we find a severe reduction of the allowed parameter space of Z' and leptoquark models explaining the B -anomalies. Remarkably, in the former case the upper bound on the Z' mass approaches dangerously close to the energy scales already probed by the LHC. We finally identify some model building directions in order to alleviate the tension with B_s -mixing.

Comments: 12 pages, 5 figures. To appear in PRD, matches the published version up to the title

Subjects: High Energy Physics – Phenomenology (hep-ph)

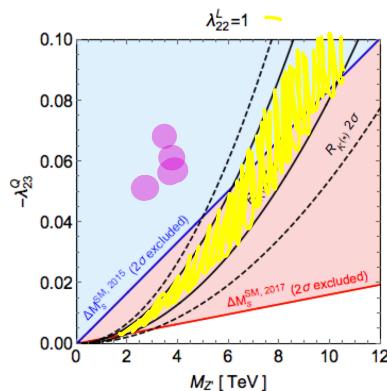
From: <prd@aps.org>
 Subject: DP11848 Di Luzio
 Date: 8 February 2018 at 17:08:14 GMT
 To: <luca.di-luzio@durham.ac.uk>
 Reply-To: <prd@aps.org>

Re: DP11848
 One constraint to kill them all?
 by Luca Di Luzio, Matthew Kirk, and Alexander Lenz

Dear Dr. Di Luzio,

Please suggest another title for the above paper that reflects more accurately the content of your manuscript and which facilitates information retrieval. We ask for a physically more informative title without reference to violence.

Yours sincerely,

Updated B_s -mixing constraints on new physics models for $b \rightarrow s\ell^+\ell^-$ anomalies

#5

Luca Di Luzio (Durham U, IPPP), Matthew Kirk (Durham U, IPPP), Alexander Lenz (Durham U, IPPP) (Dec 18, 2017)

Published in: *Phys.Rev.D* 97 (2018) 9, 095035 · e-Print: [1712.06572 \[hep-ph\]](https://arxiv.org/abs/1712.06572)[pdf](#) [DOI](#) [cite](#)

116 citations

