

Heavy Ion Physics

part II

Korinna Zapp

Lund University

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First lecture: summary

- ▶ heavy ion collisions: unique opportunity to study matter at extreme conditions (temperature and pressure)
- ▶ formation of new phase of QCD matter: quark-gluon plasma
→ deconfinement, chiral symmetry restoration
- ▶ QGP: only strongly coupled system of Standard Model microscopic degrees of freedom
- ▶ what we have learned so far
 - ▶ QGP created in heavy ion collisions at RHIC and LHC
 - ▶ “rapid hydrodynamisation”
 - ▶ QGP is a strongly coupled liquid

Summary: stages of a heavy ion collision

- ▶ **initial state:**
 - ▶ geometry (\rightarrow Glauber model)
 - ▶ gluon saturation (\rightarrow colour glass condensate)
- ▶ **pre-equilibrium dynamics:**
 - ▶ successful descriptions at weak (kinetic theory) and strong coupling (AdS/CFT)
 - ▶ leads to hydrodynamic regime on timescales $\mathcal{O}(1 \text{ fm}/c)$
- ▶ **hydrodynamic expansion:**
 - ▶ the QGP flows \rightarrow collective behaviour
 - ▶ least dissipative system known ($\eta/s \gtrsim 1/4\pi$)
- ▶ **hadronisation:** phenomenological modeling
- ▶ **hadronic re-scattering:** hydrodynamics or transport theory
more important at lower beam energies

Outline of today's lecture

Hard processes

- Factorisation and nuclear PDFs

- Quarkonia

- Electroweak bosons

- Jets

Small collision systems

Ultra-peripheral collisions

Summary

Factorisation of the cross section

- ▶ **factorisation** of the cross section:

$$\sigma(P_1, P_2) = \sum_{i,j} \int_0^1 dx_1 dx_2 f_i(x_1, Q^2) f_j(x_2, Q^2) \hat{\sigma}_{ij}(x_1 P_1, x_2 P_2, \alpha_s, Q^2)$$

- ▶ $\hat{\sigma}_{ij}$: partonic cross section
 - ▶ has perturbative expansion in α_s
 - known: NNLO for di-jets, NLO for up to ~ 5 jets**
 - ▶ short distance physics: insensitive to nature of incoming hadrons
 - i.e. no nuclear modifications**
- ▶ $f_i(x, Q^2)$: parton distribution function
 - nuclear pdf fits available**

Nuclear PDFs: EPPS16

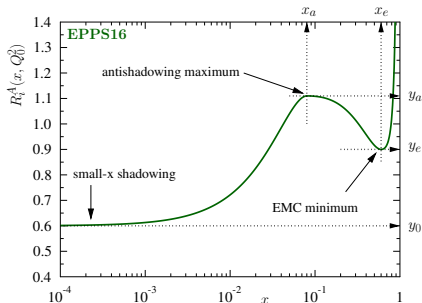
Eskola, Paakkinen, Paukkunen, Salgado, Eur. Phys. J. C 77 (2017) no.3, 163

- ▶ bound proton PDF defined as

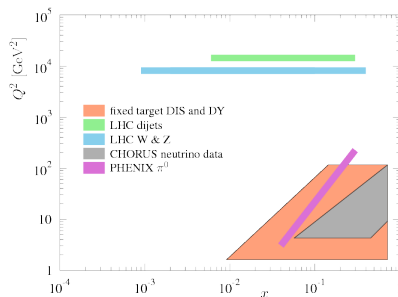
$$f_i^{p/A}(x, Q^2) = R_i^A(x, Q^2) f_i^p(x, Q^2)$$

- ▶ bound neutron PDF from isospin symmetry

- ▶ Q^2 dependence: DGLAP with 2-loop splitting functions
- ▶ parametrise R_i^A and fit for chosen free proton PDF (in this case CT14NLO)
- ▶ other nPDF sets: nCTEQ15, DSSZ



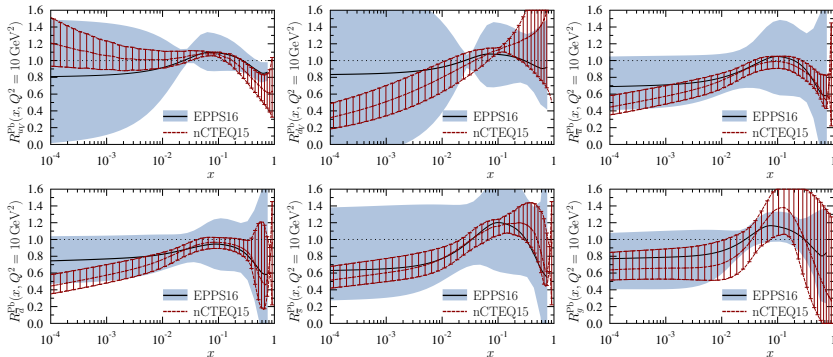
Nuclear PDFs



Escola, Paakinen, Paukkunen, Salgado, Eur. Phys. J. C 77 (2017) no.3, 163

- ▶ consistent fit to all available data
- ▶ constraining data sparse → sizeable uncertainties
- ▶ modifications less important at higher momentum scales

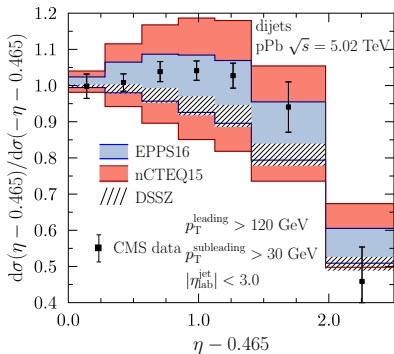
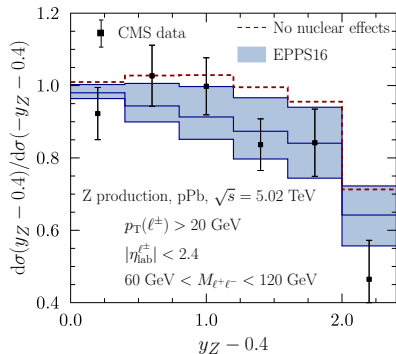
Nuclear PDFs



Eskola, Paakinen, Paukunen, Salgado, Eur. Phys. J. C 77 (2017) no.3, 163

Factorisation and nuclear PDFs

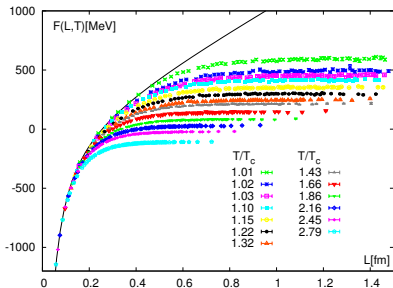
Nuclear PDFs



Eskola, Paakinen, Paukunen, Salgado, Eur. Phys. J. C 77 (2017) no.3, 163

Quarkonium suppression: idea

- ▶ in deconfined medium: colour screening
- ▶ attractive force between quark and anti-quark screened

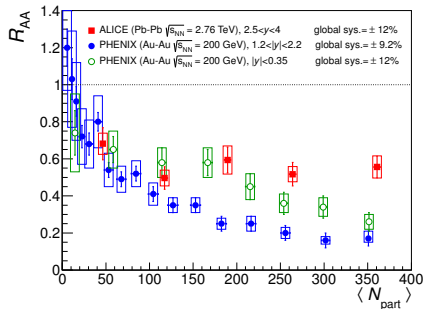
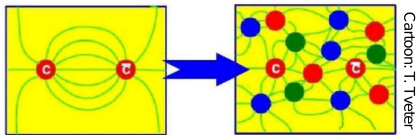


Ewerz, Kaczmarek, Samberg, JHEP 1803 (2018) 088 [arXiv:1605.07181]

- ▶ should lead to dissociation of bound states
- ▶ suitable states to observe this: charmonium & bottomonium

Matsui & Satz, 1986

J/ψ suppression

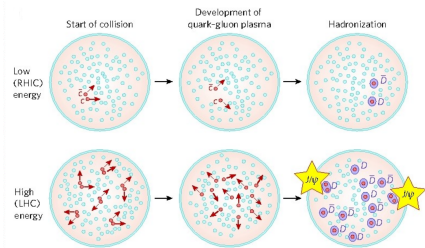


ALICE, Phys. Rev. Lett. 109 (2012) 072301 [arXiv:1202.1383]

- ✓ clear suppression of J/ψ
- ✗ weaker suppression at higher cms energy

J/ψ regeneration

- ▶ charm quarks produced early in a hard processes
 too heavy for thermal production
- ▶ quark move quasi-freely in QGP phase



picture: J. Stachel

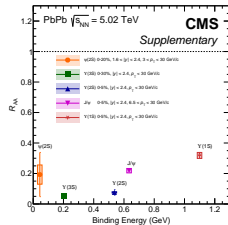
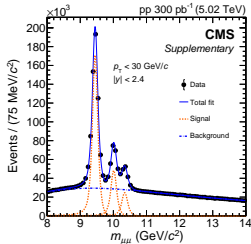
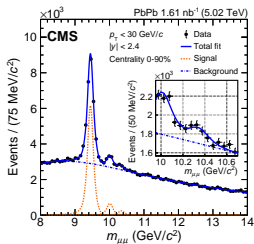
- ▶ at phase boundary J/ψ production by combinatorics
- ▶ J/ψ production rate depends sensitively on charm cross section
- ▶ effect increases rapidly with cms energy

Sequential suppression

state	J/ψ	ψ'	Υ	Υ'	Υ''
mass [GeV]	3.10	3.68	9.46	10.02	10.36
ΔE [GeV]	0.64	0.05	1.10	0.54	0.20
r [fm]	0.25	0.45	0.14	0.28	0.39

- ▶ expect dissociation at different temperatures
- ▶ “QGP thermometer”

Sequential suppression



CMS, arXiv:2303.17026

- ▶ there is an indication of sequential suppression
- ▶ the story is actually much more complicated

static potential over-simplified, quarkonia not produced at rest, ...

Nuclear modification factors

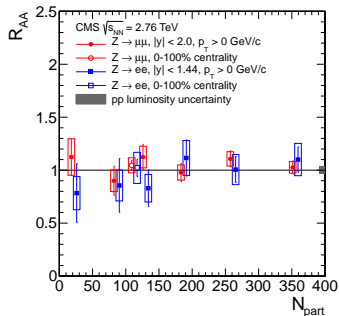
- ▶ standard procedure for quantifying QGP effects: divide by result from p+p
- ▶ but: have to scale p+p result by number of N-N collisions

hard processes scale like N_{coll}

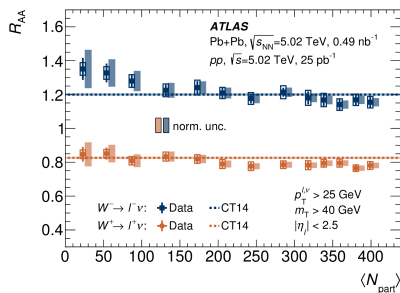
$$R_{AA}(\{\cdot\}) = \frac{\left. \frac{d\sigma_X}{d\{\cdot\}} \right|_{AA}}{\langle N_{\text{coll}} \rangle \left. \frac{d\sigma_X}{d\{\cdot\}} \right|_{pp}}$$

- ▶ deviations from unity signal effects beyond incoherent superposition of nucleon-nucleon collisions
- ▶ in case of electroweak processes: have to take different isospin composition into account

W & Z production



CMS, JHEP 1503 (2015) 022 [arXiv:1410.4825]

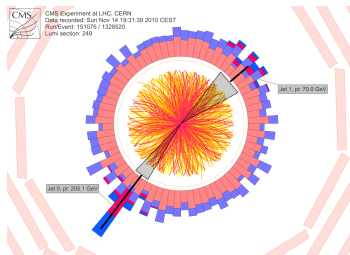
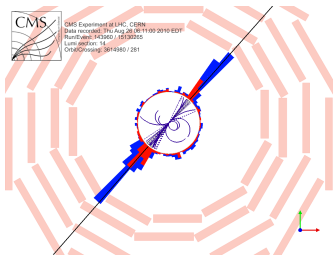


ATLAS arXiv:1907.10414

- ▶ no significant deviation from unity
- ▶ W and Z don't interact strongly → leave dense interaction region without re-scattering
- ▶ important confirmation of N_{coll} scaling and nPDFs

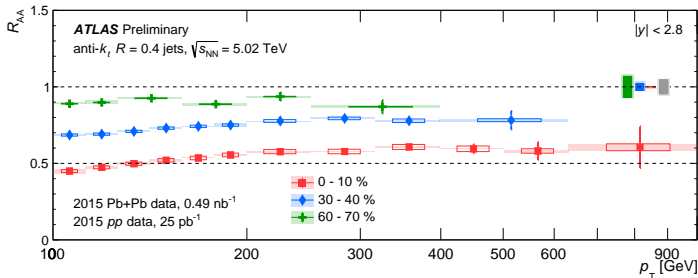
Second main discovery of heavy ion physics

Jet quenching: Hard jets are suppressed and their structure modified.



- ▶ proton-proton collisions: 2 jets with **balancing transverse momentum**
- ▶ in heavy ion collisions: significant **softening** of jets
- thermalisation of a far-from-equilibrium system
- **jet quenching informs us about equilibration in QCD**

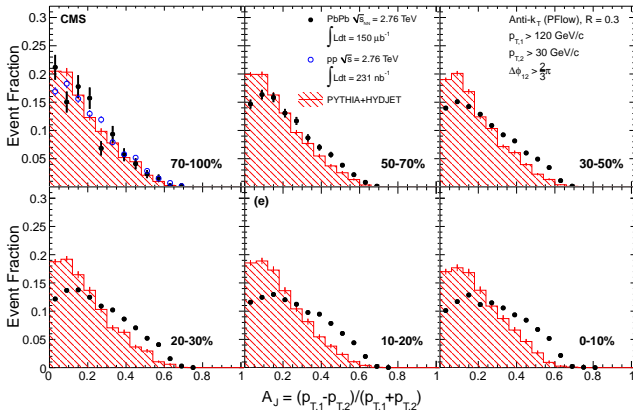
Suppression of single-inclusive jets



ATLAS-CONF-2017-009

- **suppression** of jets by factor 2 relative to expectation from p+p
 need to scale p+p reference by number of hard N+N collisions

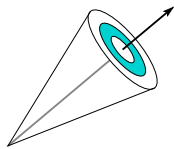
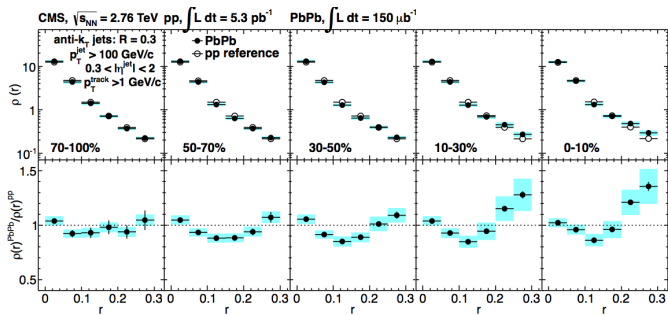
Di-jet momentum asymmetry



CMS, Phys. Lett. B 712 (2012) 176

► enhancement of asymmetric configurations

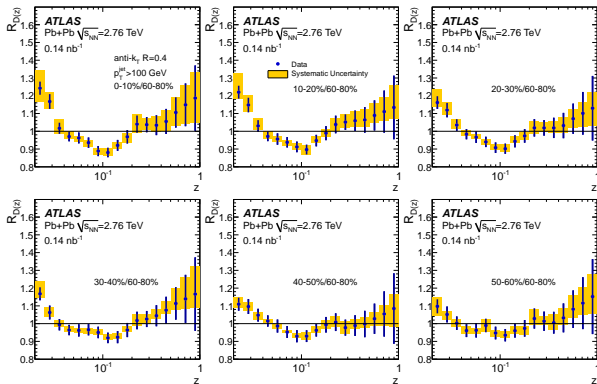
Intra-jet energy distribution: Jet profile



CMS, Phys. Lett. B 730 (2014) 243

- ▶ suppression of activity at intermediate r
- ▶ **increase** near the **edge** of the jet

Intra-jet energy distribution: fragmentation function



$$z = \frac{p_{\perp,h}}{p_{\perp,J}} \cos(\Delta R_{hJ})$$

ATLAS, Phys. Lett. B 739 (2014) 320

- ▶ distribution of hadrons inside jets
- ▶ **suppression** at **intermediate** & **enhancement** of **soft** momenta

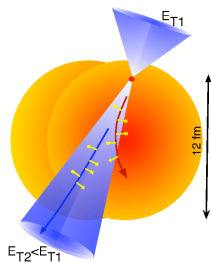
A first lesson

Summary of experimental results

- ▶ jet production suppressed by factor ~ 2 up to p_{\perp} 's of 1 TeV
- ▶ hard structures inside jets survive largely unmodified
- ▶ enhancement of soft activity far away from jet axis

A first lesson to be learned

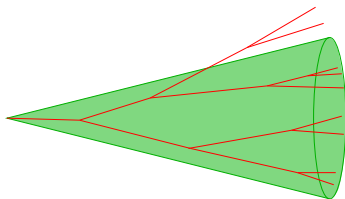
- ▶ energy loss from jets driven by peeling off of soft gluons
- ▶ rare, semi-hard emissions not an important effect
- ▶ both were equivalent in early calculations



Casalderrey-Solana, Milhano, Wiedemann, J. Phys. G 38 (2011) 035006 [arXiv:1012.0745]

What is unique about jet quenching

- ▶ “calibrated” probe: well understood in $p+p$
fixed order matrix elements + resummation (parton showers)



- ▶ jet quenching allows to observe *process of equilibration*
soft observables see *result of equilibration*
- ▶ jets give access to *scale dependence* of medium properties

What happens to jets in medium?

Scenario I: hard partons don't resolve quasi-particles

- ▶ interactions between jet & medium at large coupling
- ▶ AdS/CFT techniques

Scenario II: hard partons do resolve quasi-particles

- ▶ jet – medium interactions at weak(ish) coupling
- ▶ perturbative techniques
- ▶ thermalisation through elastic re-scattering (slow)
- ▶ parton energy loss through QCD bremsstrahlung
- ▶ destructive interference in multiple scattering LPM effect

relevant scale: momentum transfer q between hard parton and medium

Jet quenching at strong coupling

Advantages:

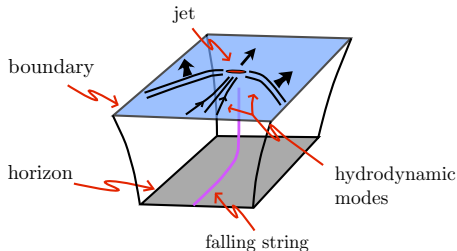
1. exact solution of gauge theory at strong coupling
2. no uncertainty in jet-medium interaction

Fundamental problems:

1. jets are a weak-coupling phenomenon
2. QCD is not $\mathcal{N} = 4$ SYM

“holographic jet”:

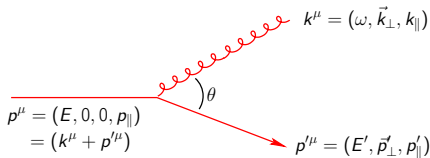
$q\bar{q}$ pair in QGP dual to classical string in black hole geometry



Chesler, Jensen, Karch, Yaffe, Phys. Rev. D 79 (2009) 125015

Chesler & Rajagopal, JHEP 1605 (2016) 098

How long does it take to radiate a gluon?



assume: massless quark

- ▶ virtual state: $p^2 = E^2 - \mathbf{p}^2 \neq 0 \rightarrow m_{\text{virt}}^2 = p^2$
- ▶ uncertainty principle: $1 = \Delta t \Delta E = \Delta t m_{\text{virt}}$
- ▶ gluon formation time: $t_{\text{form}} = \Delta t \times (\text{boost factor})$

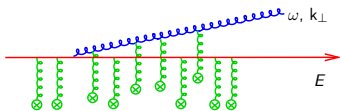
$$t_{\text{form}} = \frac{1}{m_{\text{virt}}} \frac{E}{m_{\text{virt}}} = \frac{E}{2p_\mu k^\mu} \simeq \frac{E}{\omega E \theta^2} \simeq \frac{\omega}{k_\perp^2}$$

- ▶ time for entire jet evolution: $\mathcal{O}((1 - 10) \text{ fm})$
- ▶ (transverse) size of medium: $\mathcal{O}((1 - 10) \text{ fm})$

Bremsstrahlung in medium: heuristic discussion

Baier, Dokshitzer, Mueller, Peigne, Schiff, Nucl. Phys. B 483 (1997) 291

ω, k_{\perp} Baier, Schiff, Zakharov, Ann. Rev. Nucl. Part. Sci. 50 (2000) 37



$$E \gg \omega \gg k_{\perp}, q_{\perp}$$

Brownian motion of the gluon: $\langle k_{\perp}^2 \rangle = \hat{q}L$

formation time of the radiated gluon:

$$t_f \simeq \frac{\omega}{k_{\perp}^2} \simeq \frac{\omega}{\hat{q}t_f} \Rightarrow t_f = \sqrt{\frac{\omega}{\hat{q}}} \quad \text{and} \quad N_{\text{coh}} = \frac{t_f}{\lambda}$$

gluon energy spectrum:

$$\frac{d^2I_{\text{coh}}}{d\omega dy} \simeq \frac{1}{N_{\text{coh}}} \frac{d^2I_{\text{incoh}}}{d\omega dy} \propto \frac{\alpha_s}{\omega\lambda} \lambda \sqrt{\frac{\hat{q}}{\omega}} = \alpha_s \sqrt{\hat{q}} \omega^{-3/2}$$

radiative energy loss:
$$\Delta E = \int_0^L dy \int_0^{\omega_c} d\omega \omega \frac{d^2I}{d\omega dy} \propto \alpha_s \hat{q} L^2$$

Beyond parametric estimates

Eikonal (straight line) propagation



- ▶ interactions result in colour phase rotation
- ▶ in-medium propagator: Wilson line

$$W(x_{f+}, x_{i+}; x_{\perp}) = \mathcal{P} \exp \left\{ ig \int_{x_{i+}}^{x_{f+}} d\xi A_{-}(\xi, x_{\perp}) \right\}$$

path ordering

medium colour field

Further developments

- ▶ non-eikonal propagation
- ▶ 2-gluon emission
- ▶ resummation of incoherent emissions
- ▶ ...

From partons to jets: antennae

- ▶ consider colour singlet dipole in medium
- ▶ in vacuum colour coherence leads to angular ordering of radiation
- ▶ re-scattering in medium introduces new scale: q_{\perp}
- ▶ medium modifications depend on whether dipole is resolved

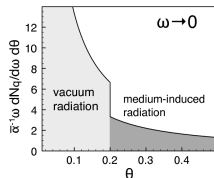


- ▶ in particular: unresolved structures remain unperturbed

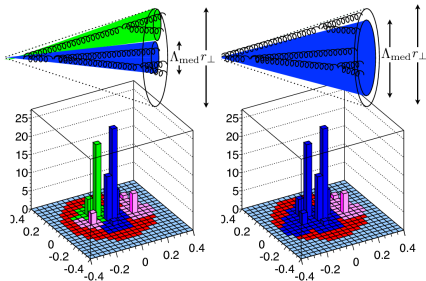
Casalderrey-Solana, Mehtar-Tani, Salgado, Tywoniuk, Phys. Lett. B 725 (2013) 357

From partons to jets: antennae

- ▶ decoherence time scale: $\tau_d \sim \left(\frac{1}{\hat{q}\theta^2_{q\bar{q}}} \right)^{1/3}$
- ▶ decoherence opens up phase space
- ▶ in soft limit: geometric separation



Mehtar-Tani, Salgado, Tywoniuk, JHEP 1204 (2012) 064 doi:10.1007/JHEP04(2012)064 [arXiv:1112.5031]

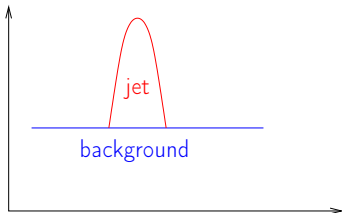


- ▶ can be generalised to jets

Casalderrey-Solana, Mehtar-Tani, Salgado, Tywoniuk,
Phys. Lett. B 725 (2013) 357 [arXiv:1210.7765]

Medium response

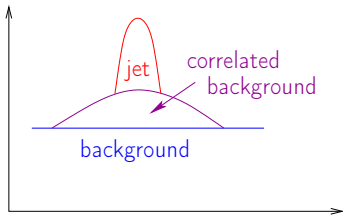
- ▶ medium response: medium's reaction to energy & momentum deposited by jet
- ▶ gives rise to additional soft activity
- ▶ momentum conservation → additional particles follow jet direction
- ▶ correlated background



- ▶ correlated background cannot and should not be subtracted
- ▶ rather, it is part of jet
- ▶ fluctuations matter

Medium response

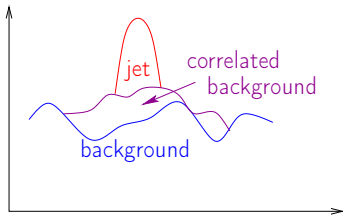
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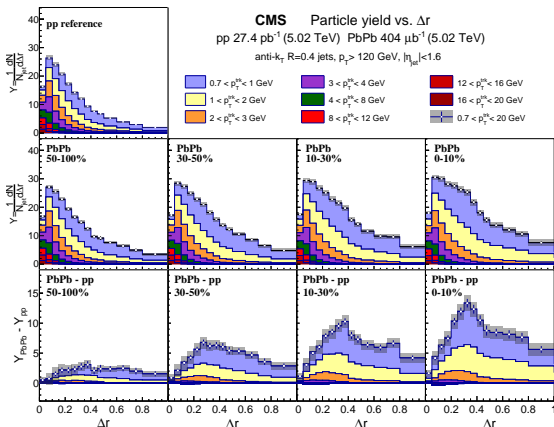
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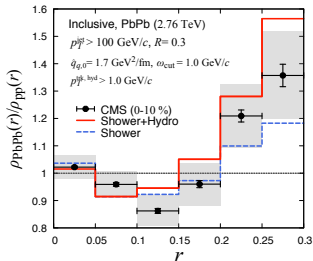
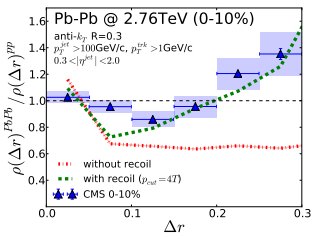
- ▶ correlated background cannot and should not be subtracted
- ▶ rather, it is part of jet
- ▶ fluctuations matter

Particle distribution outside jets



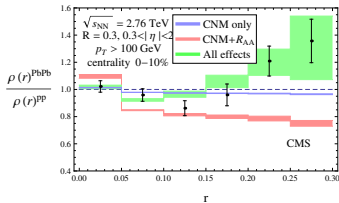
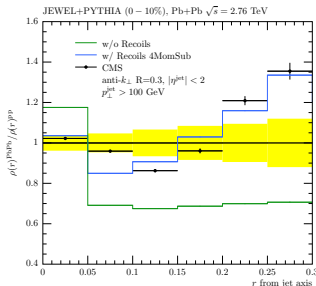
► enhancement of soft particles far away from jet

Some theory results



Tachibana, Chang, Qin,
 Phys. Rev. C 95 (2017) no.4, 044909

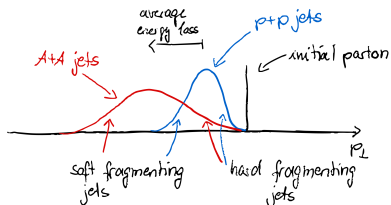
Park, Jeon, Gale, arXiv:1807.06650



Kunnawalkam Elayavalli, Zapp,
 JHEP 1707 (2017) 141

Chien, Vitev,
 JHEP 1605 (2016) 023

Selection bias



- ▶ jet spectrum steeply falling
- selecting on final jet p_{\perp} favors jets with small energy loss
- ▶ these are predominantly hard fragmenting, narrow jets
- ▶ bias stronger in A+A than in p+p
- ▶ complicates interpretation of jet shapes and sub-structure of quenched jets

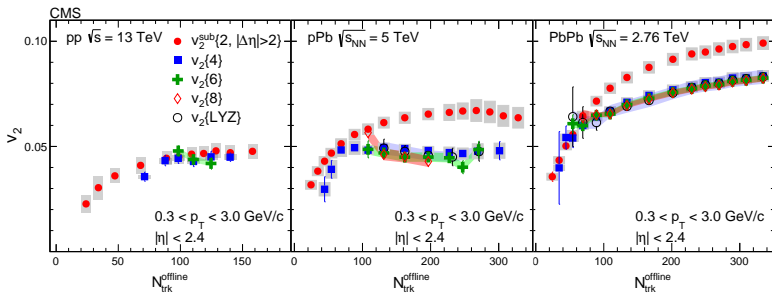
Third main discovery of heavy ion physics

- ▶ (partial) equilibration of final state in A+A
- ▶ expectation: p+p qualitatively different from A+A

p+p collisions too dilute to develop collectivity

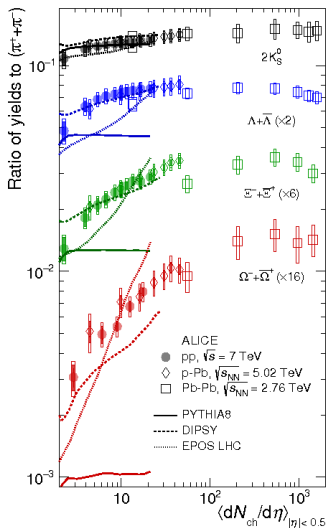
▶ observed:

Soft particle production in high-multiplicity p+p and p+A closely resembles A+A, but so far no jet quenching is observed.



CMS, Phys. Lett. B 765 (2017) 193

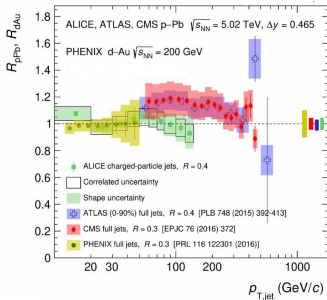
Strangeness enhancement



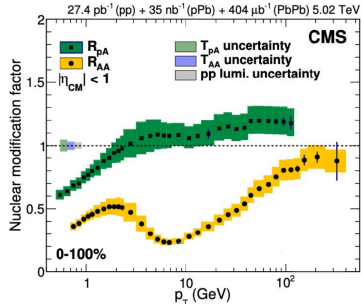
- ▶ strong increase of strange particle production with multiplicity
- ▶ traditional view: strangeness production in p+p suppressed due to s quark mass
- ▶ strangeness practically equilibrated in Pb+Pb collisions at LHC
- ▶ strangeness enhancement counted as signal for formation of hot system

ALICE, Nature Phys. 13 (2017) 535

Jet quenching



ALICE, arXiv:2307.10860



CMS, JHEP 04 (2017) 039

- ▶ no suppression of jets and leading hadrons in p+Pb
- ▶ large suppression in Pb+Pb collisions

Possible explanations for v_n 's in p+p

Final state interactions

azimuthal anisotropies due to response to initial geometry

- ▶ hydrodynamics

Weller, Romatschke, Phys. Lett. B 774 (2017) 351

- ▶ kinetic theory

Kurkela, Wiedemann, Wu, Phys. Lett. B 783 (2018) 274

Initial state correlations

initial state correlations imprinted on final state

- ▶ color-glass condensate

Schenke, Schlichting, Venugopalan, Phys. Lett. B 747 (2015) 76

Alternative explanations

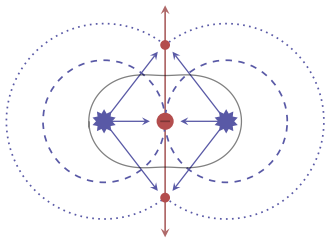
models inspired by p+p physics

- ▶ string shoving Bierlich, Gustafson, Lönnblad, Phys. Lett. B 779 (2018) 58 [arXiv:1710.09725]

at hadronisation densely packed strings repel each other

The escape mechanism

Kurkela, Wiedemann, Wu, Phys. Lett. B 783 (2018) 274

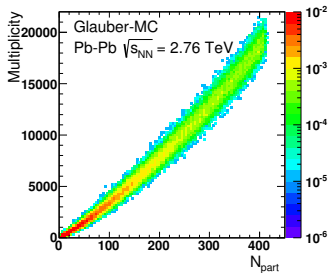
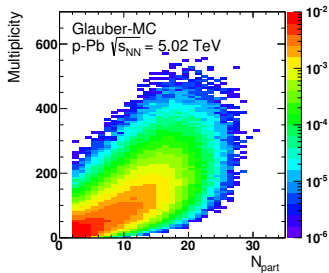


- ▶ initially isotropic particle distribution with spatial anisotropy
 - ▶ local anisotropy where scatterings happen
 - ▶ local isotropisation due to scattering induces global anisotropy
- ▶ requires only $\mathcal{O}(1)$ scattering per particle
 - ▶ also observed in transport code AMPT

He, Edmonds, Lin, Liu, Molnar, Wang, Phys. Lett. B 753 (2016) 506

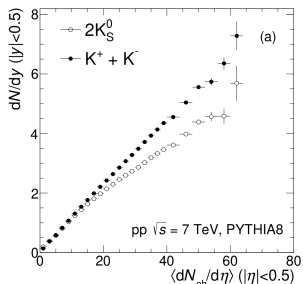
Important differences between small and large systems

- ▶ multiplicity generated in different ways:
 - ▶ Pb+Pb: requiring high multiplicity biases towards central events
 - ▶ p+p/Pb: requiring high multiplicity biases towards jetty events
- ▶ not clear what multi-particle correlations measure in jetty events
- ▶ no strong correlation between centrality and multiplicity



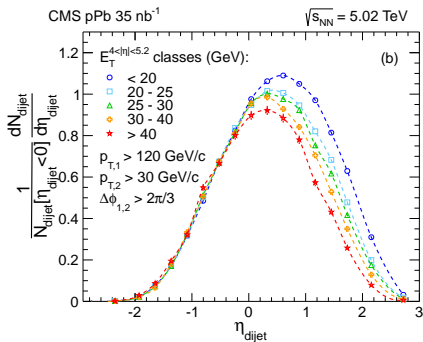
ALICE, Phys. Rev. C 91 (2015) no.6, 064905

Auto-correlations in small collision systems



O. Matonoha, PhD thesis

- ▶ autocorrelation with charged particle multiplicity at mid-rapidity
- ▶ improves when measuring N_{ch} at forward rapidity



CMS, Eur. Phys. J. C 74 (2014) no.7, 2951

- ▶ but: rapidity shift of di-jets due to energy conservation

Armesto, Gülhan, Milhano, Phys. Lett. B 747

(2015) 441

Outline

Hard processes

- Factorisation and nuclear PDFs

- Quarkonia

- Electroweak bosons

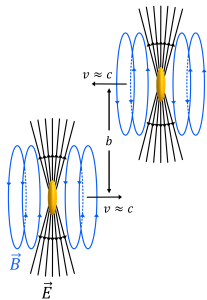
- Jets

Small collision systems

Ultra-peripheral collisions

Summary

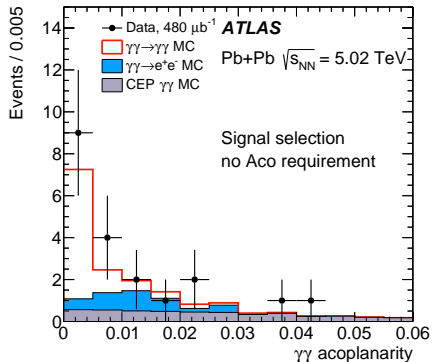
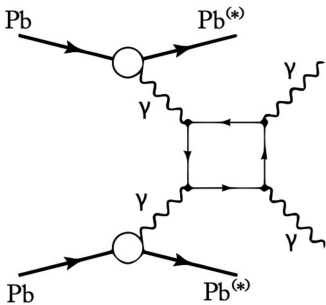
Ultra-peripheral collisions



- ▶ relativistic ions are sources of strong EM fields
- ▶ $E_{\text{max}} \approx 10^{16} - 10^{18} \text{ V/m}$
 $B_{\text{max}} \approx 10^{14} - 10^{16} \text{ T}$
- ▶ strongest EM fields in universe
- ▶ very short-lived
- ▶ two types of interactions: photon-photon scattering and photo-nuclear reactions

Light-by-light scattering

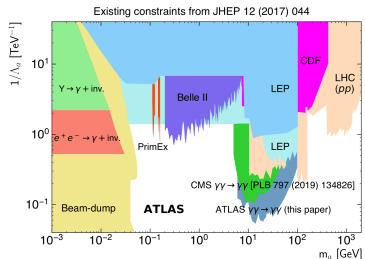
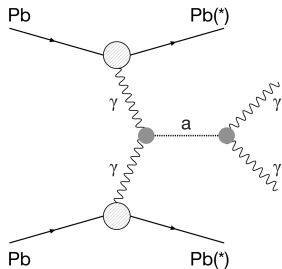
► first evidence for $\gamma\gamma \rightarrow \gamma\gamma$ scattering



ATLAS, Nature Phys. 13 (2017) no.9, 852 [arXiv:1702.01625]

Light-by-light scattering

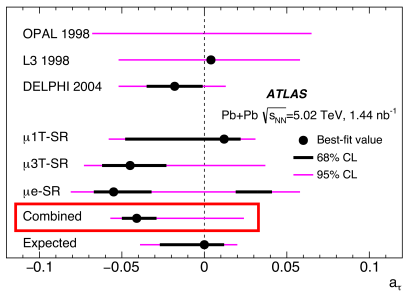
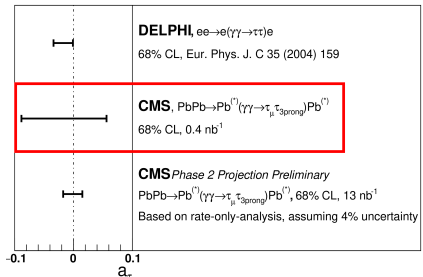
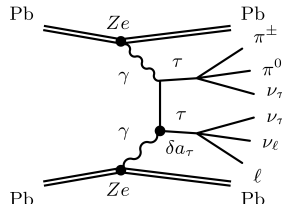
► sensitive to axion-like particles



ATLAS, JHEP 03 (2021), 243 [erratum: JHEP 11 (2021), 050]

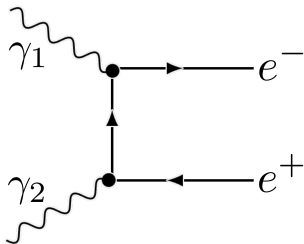
τ pair production

- ▶ sensitive to τ anomalous magnetic moment a_τ
 - ▶ BSM sensitivity $\delta a_I \propto m_I^2$
- 280 × larger than for muon

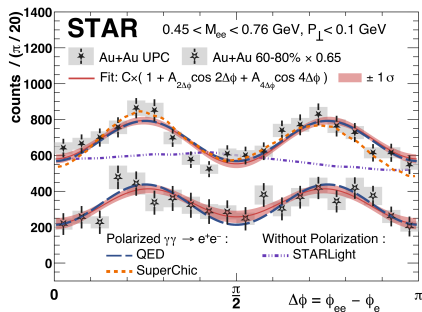


CMS, arXiv:2206.05192 & ATLAS, arXiv:2204.13478

Breit-Wheeler process and photon spin

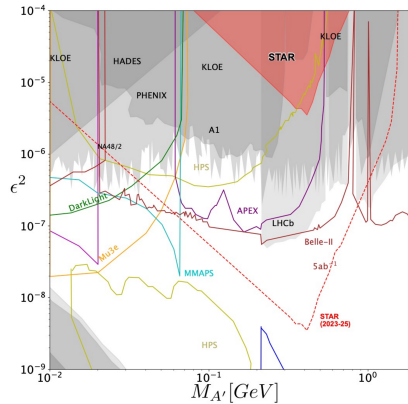
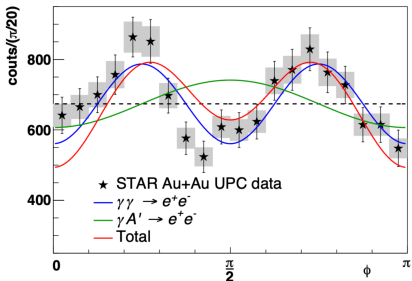
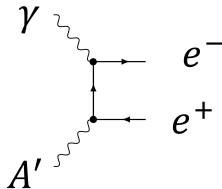


- ▶ real photons: only transverse polarisation
- ▶ imprinted on leptons as angular orbital momentum
- angular asymmetry



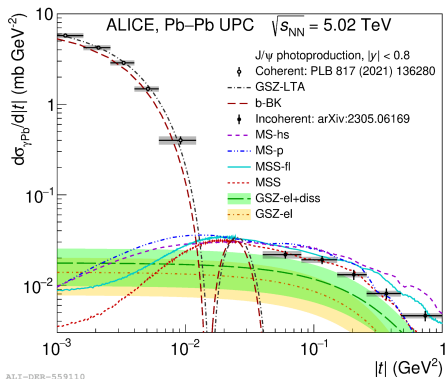
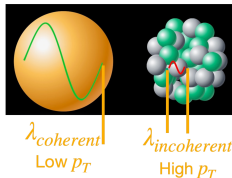
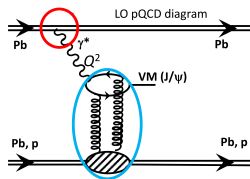
STAR, Phys. Rev. Lett. 127 (2021) no.5, 052302

Breit-Wheeler process and dark photons



Xu, Lewis, Wang, Brandenburg, Ruan, arXiv:2211.02132

Photonuclear J/ψ production



ALICE, Phys. Lett. B 817 (2021), 136280

ALICE, arXiv:2305.06169

- sensitive to nuclear and sub-nucleonic structure

Outline

Hard processes

Factorisation and nuclear PDFs

Quarkonia

Electroweak bosons

Jets

Small collision systems

Ultra-peripheral collisions

Summary

Wrap-up

- ▶ heavy ion collisions allow to study fundamental questions in strong interaction physics
- ▶ three main discoveries of heavy ion physics:
 1. hydrodynamic behaviour
 - ▶ QGP behaves as an almost perfect liquid
 - ▶ lowest η/s observed, close to theoretical lower bound of $1/4\pi$
 2. jet quenching
 - ▶ partial thermalisation of a far-from-equilibrium system
 - ▶ informs us about how the QGP came into being
 - ▶ gives access to scale dependence of QGP properties
 3. similarity between p+p/p+A and A+A collisions
 - ▶ exciting, but no firm conclusions yet
- ▶ many more exciting phenomena

some of which I discussed, others I had to leave out

Wrap-up

- ▶ heavy ion physics is new physics
- ▶ rich and diverse phenomenology
- ▶ have to develop new theoretical tools using
 - ▶ classical field theory
 - ▶ AdS/CFT
 - ▶ thermal field theory
 - ▶ kinetic theory/transport theory
 - ▶ relativistic hydrodynamics
 - ▶ effective field theories
 - ▶ zero-temperature field theory (in particular perturbation theory)
 - ▶ phenomenological models
- ▶ also new experimental techniques needed
 - e.g. variety of correlations, background subtraction for jets, ...

Wrap-up

- ▶ heavy ion physics is ~~new~~ unexplored physics
- ▶ rich and diverse phenomenology
- ▶ have to develop new theoretical tools using
 - ▶ classical field theory
 - ▶ AdS/CFT
 - ▶ thermal field theory
 - ▶ kinetic theory/transport theory
 - ▶ relativistic hydrodynamics
 - ▶ effective field theories
 - ▶ zero-temperature field theory (in particular perturbation theory)
 - ▶ phenomenological models
- ▶ also new experimental techniques needed
 - e.g. variety of correlations, background subtraction for jets, ...
- ▶ heavy ion physics is a broad, multifaceted and exciting field