

LHC Run 3 and HL-LHC

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Lecture 2

Nikhef



Higgs couplings to 2nd generation: c quarks



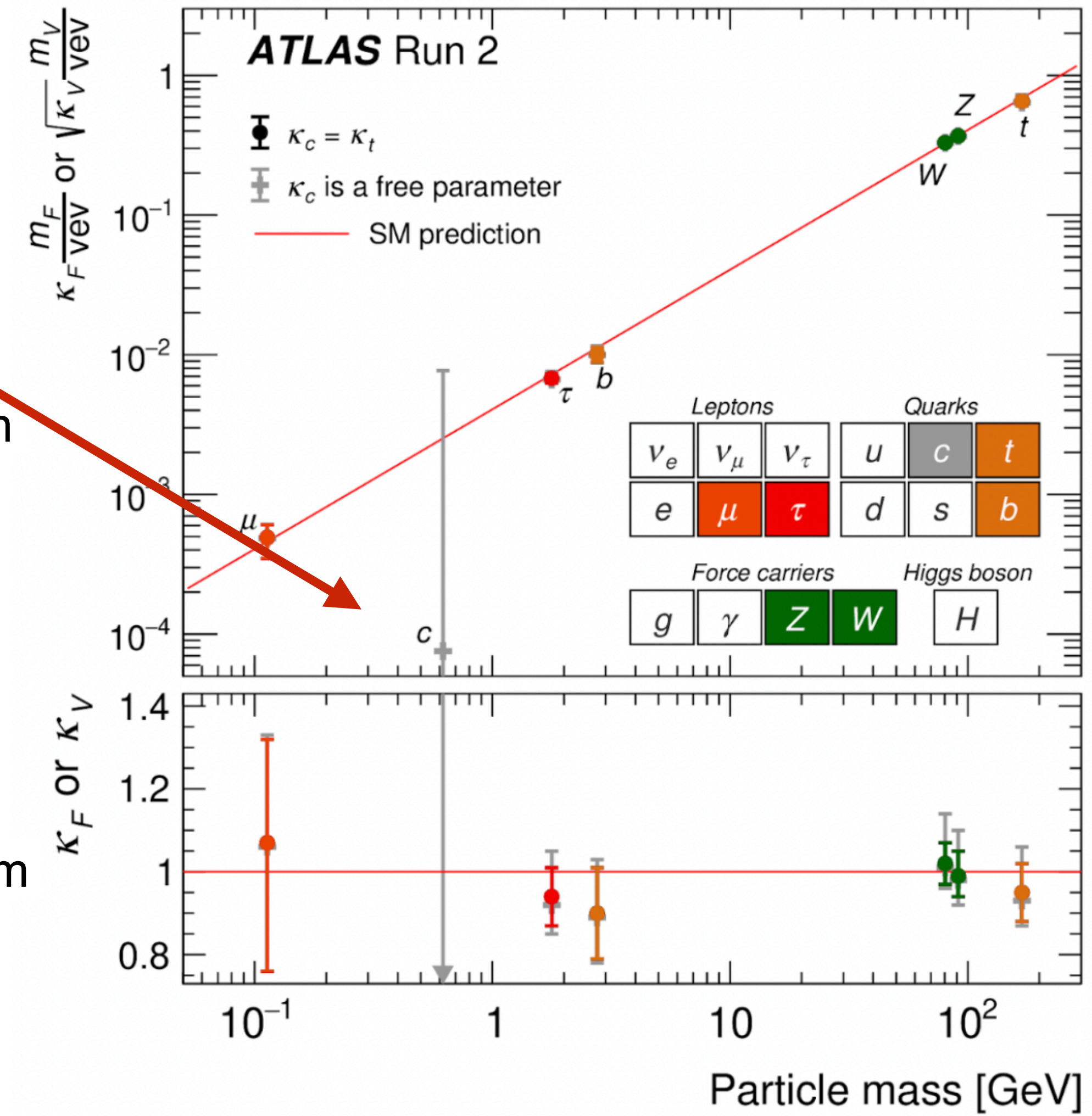
Run 3 will serve as a fundamental benchmark for studying the coupling to second-generation quarks, specifically the charm quark.

Decays of the Higgs boson into a pair of c ("charm") quarks are relatively common; however, the challenge lies in accurately identifying them based on their detector signature.

When high-energy quarks transform into collimated jets of bound states known as hadrons, those originating from b or c quarks travel a finite distance before decaying (D lifetime 10^{-15} s, B lifetime 10^{-12} s)

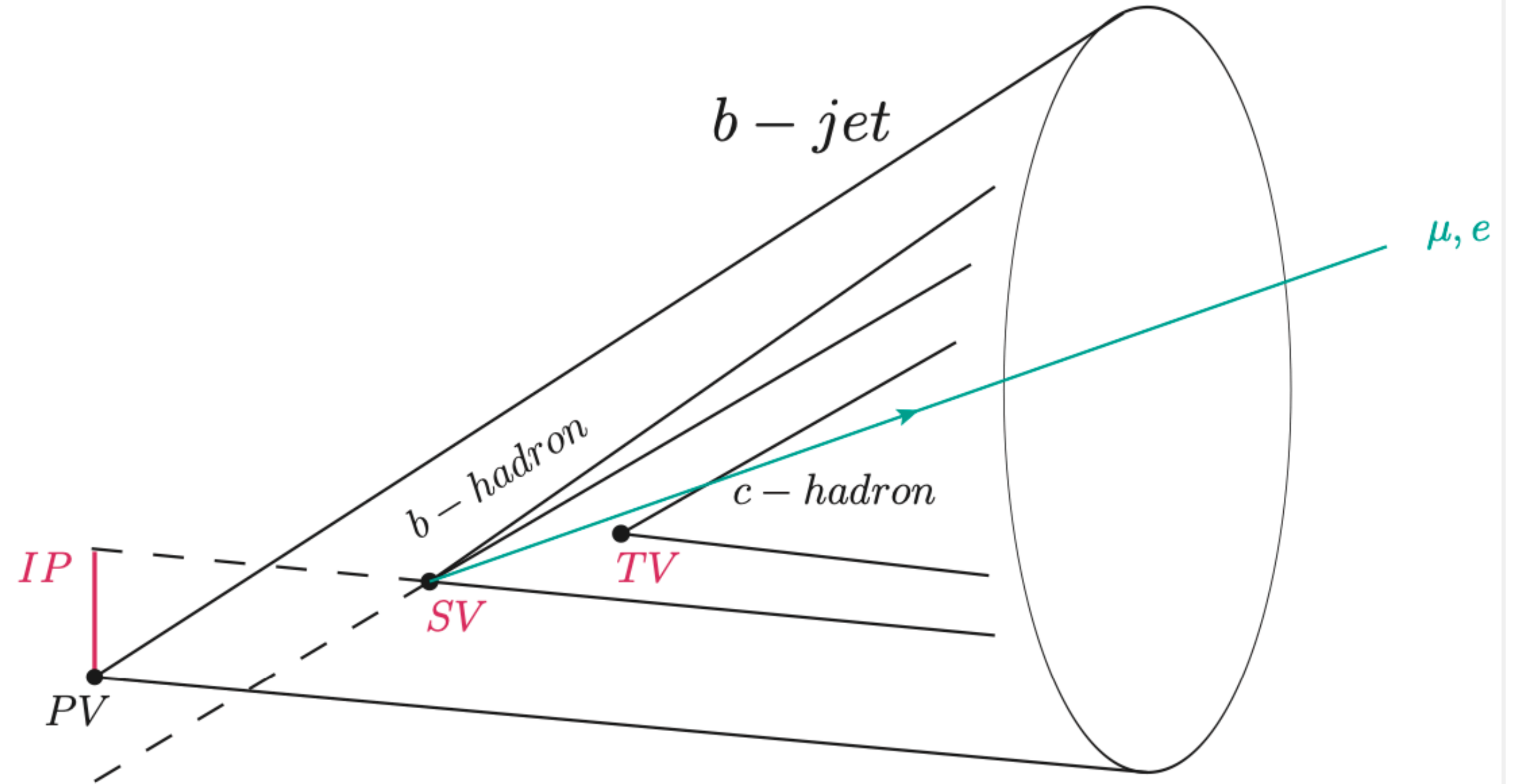
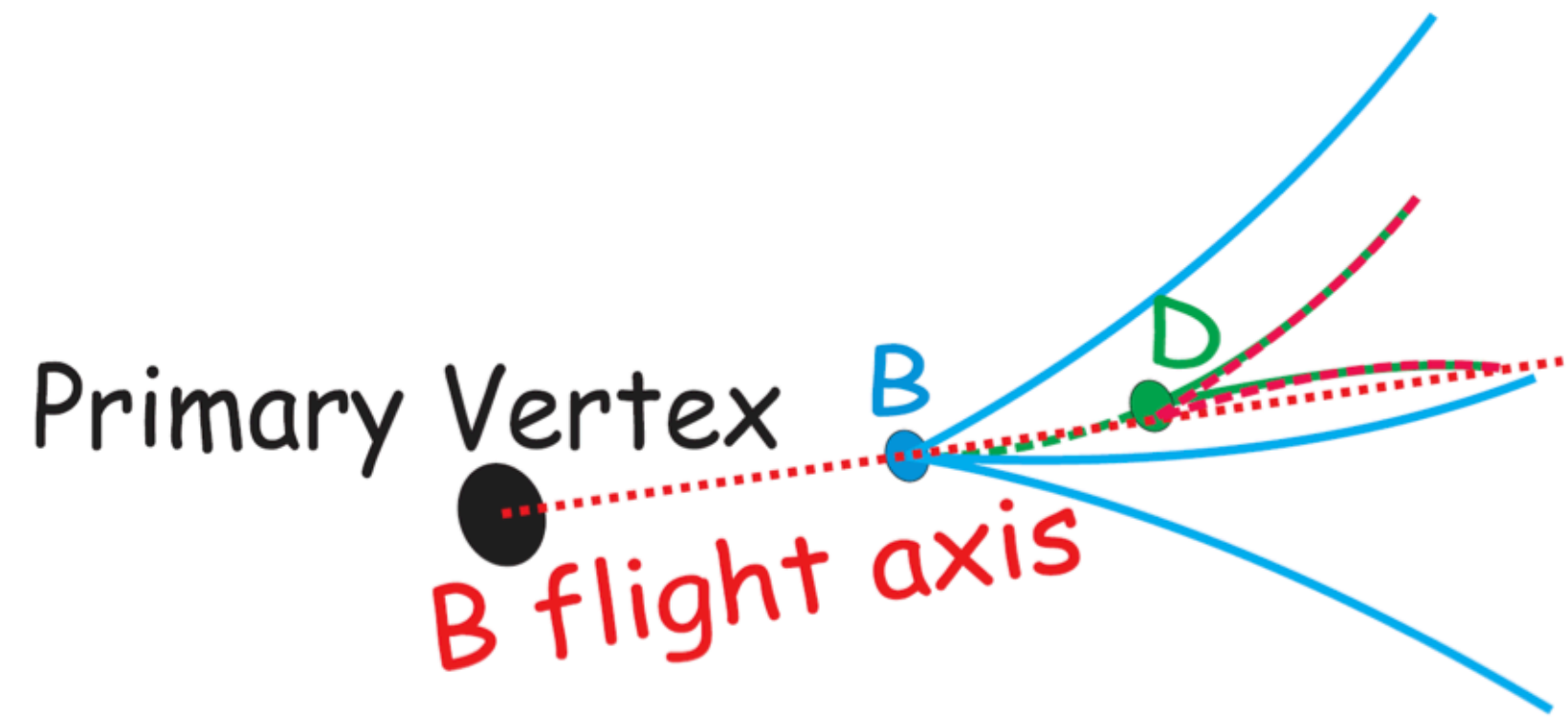
Techniques based on distance measurements have proven effective in identifying the long-lived and heavy b quarks of the third generation.

To address the more challenging scenario of the shorter-lived and lighter charm quarks, innovative analysis techniques and the utilization of boosted Higgs decays have brought the charm quark within reach for the High-Luminosity phase of the LHC. Run 3 will be instrumental in testing and establishing new analysis strategies to pave the way forward.



Flavor tagging

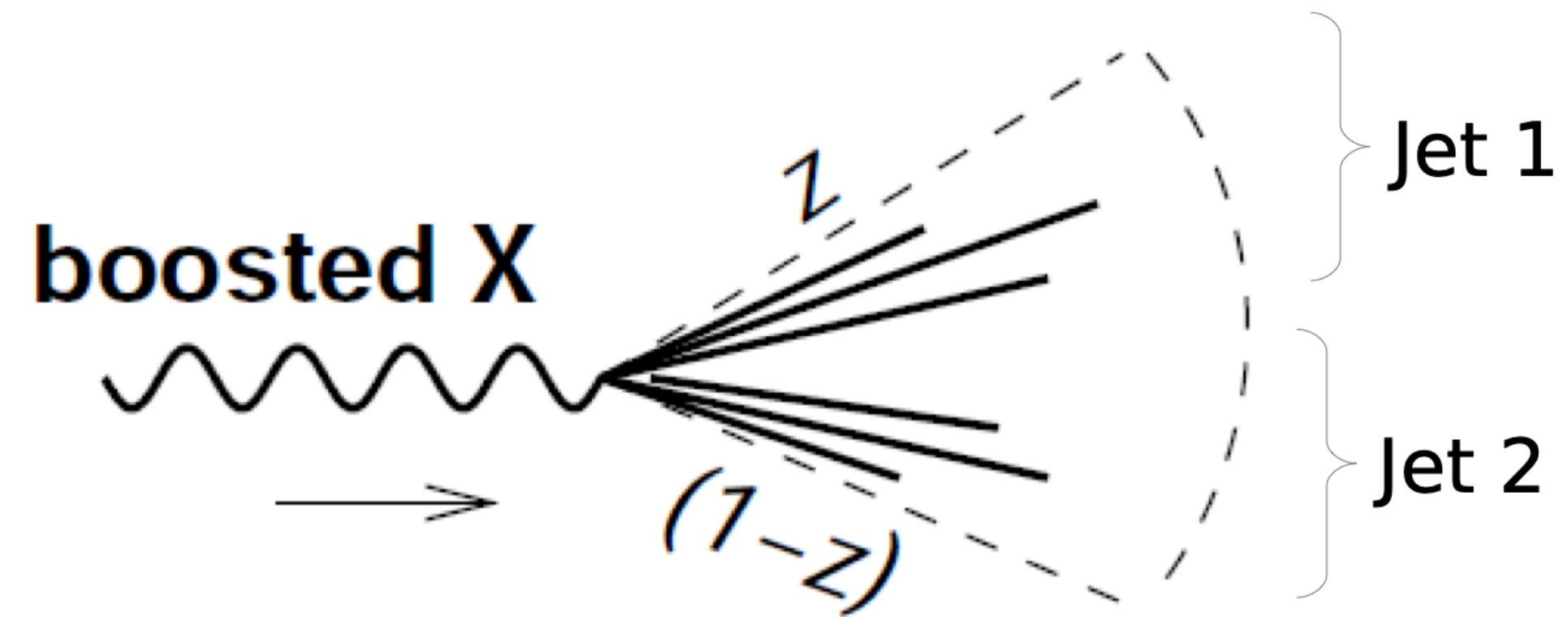
We tag b-hadrons and c-hadrons thanks to the fact that there is a secondary vertex



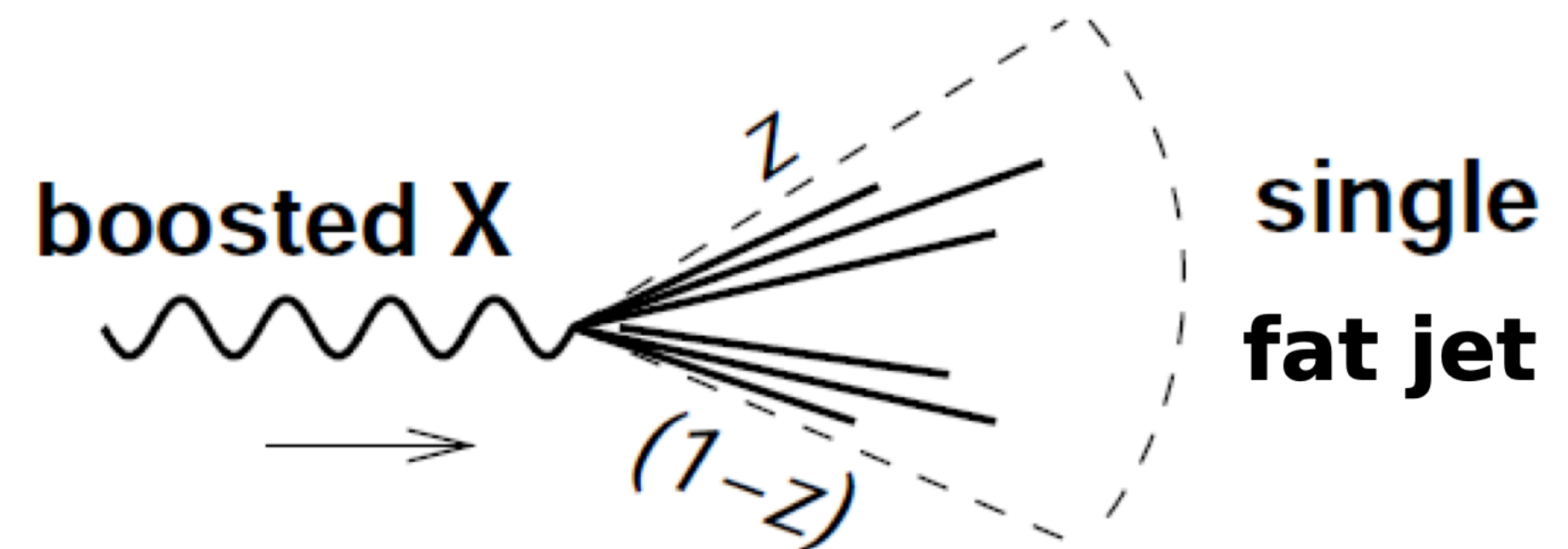
Boosted objects

At the LHC given the large center of mass energy and given that the SM particles have masses below 200 GeV, also the heaviest SM particles often acquire large momentum $\gg m \rightarrow$ production of “boosted objects”

Normally we reconstruct jets with $R=0.4$, if the object is boosted the jets in which it decays cannot be resolved in small r -jets



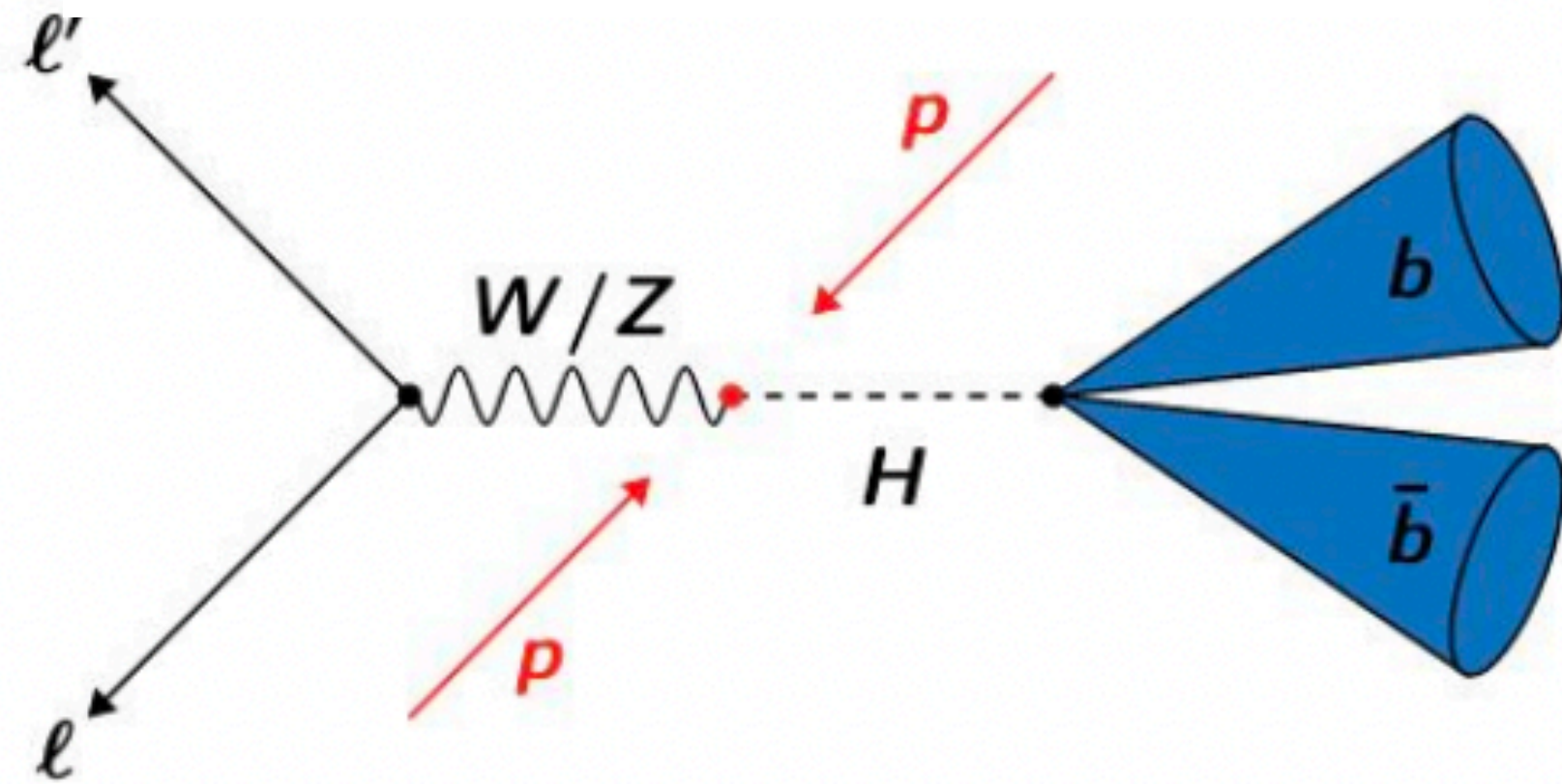
Recover sensitivity to boosted objects by developing boosted taggers, using larger R



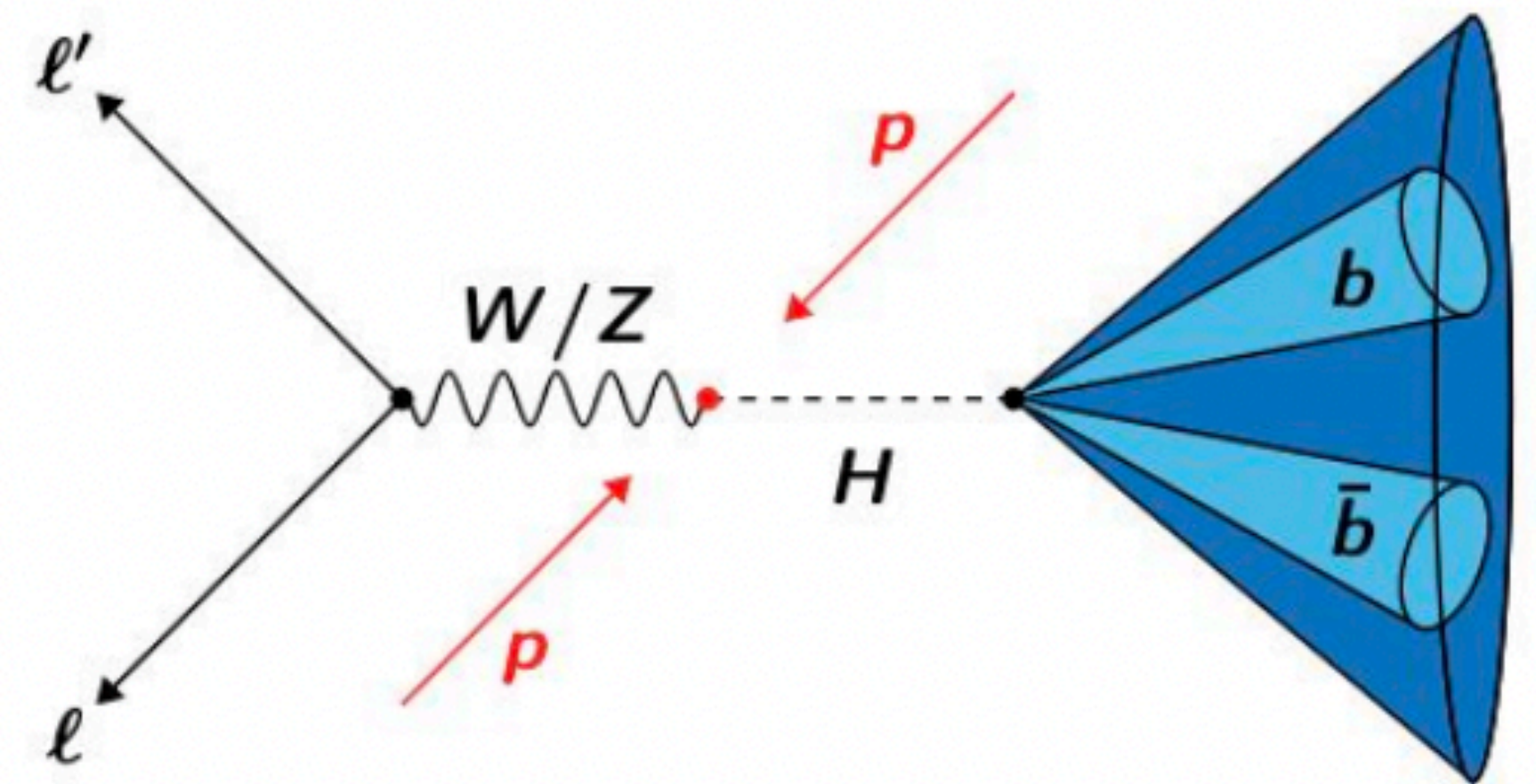
Boosted objects

At the LHC given the large center of mass energy and given that the SM particles have masses below 200 GeV, also the heaviest heaviest SM particles often acquire $p_T \gg m \rightarrow$ production of “boosted objects”

Resolved

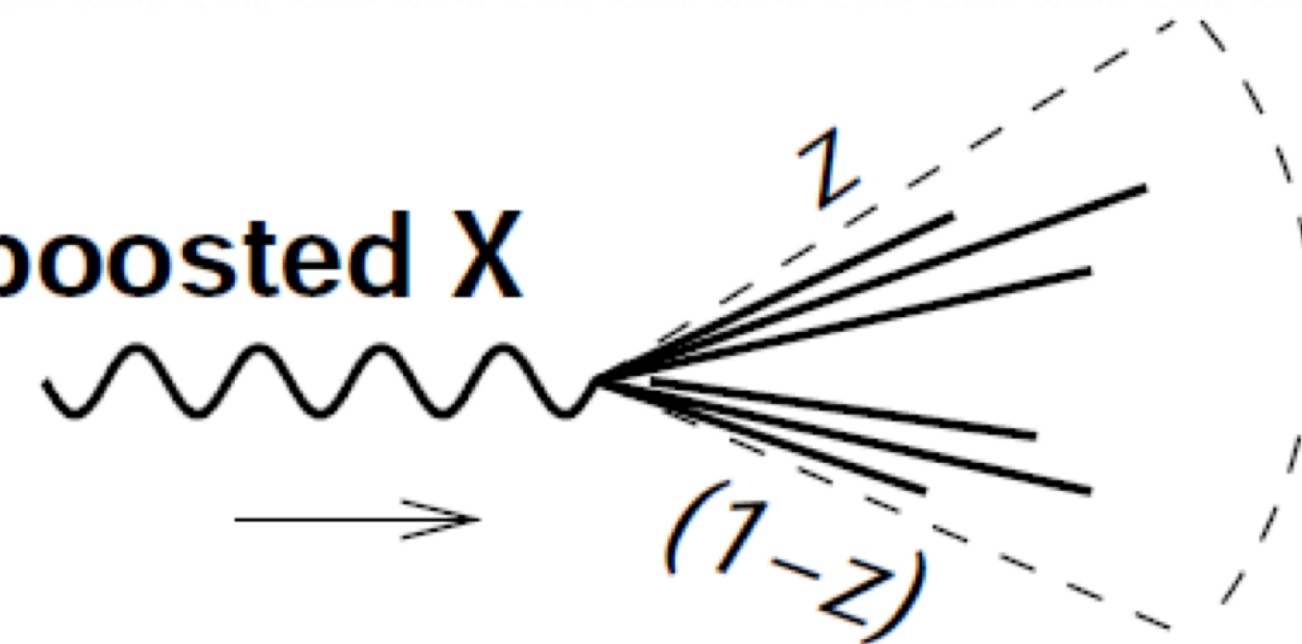


Boosted



Recover sensitivity to boosted objects by developing boosted taggers, using larger R

boosted X



single fat jet

Higgs boosted taggers for $H \rightarrow c\bar{c}$

arXiv:2205.05550



Latest CMS Run 2 results (dataset 20 times smaller than HL-LHC) has sensitivity of 3.4 times the SM coupling in VH (WH, ZH) production mode. When the V has a large p_T , the Higgs is boosted.

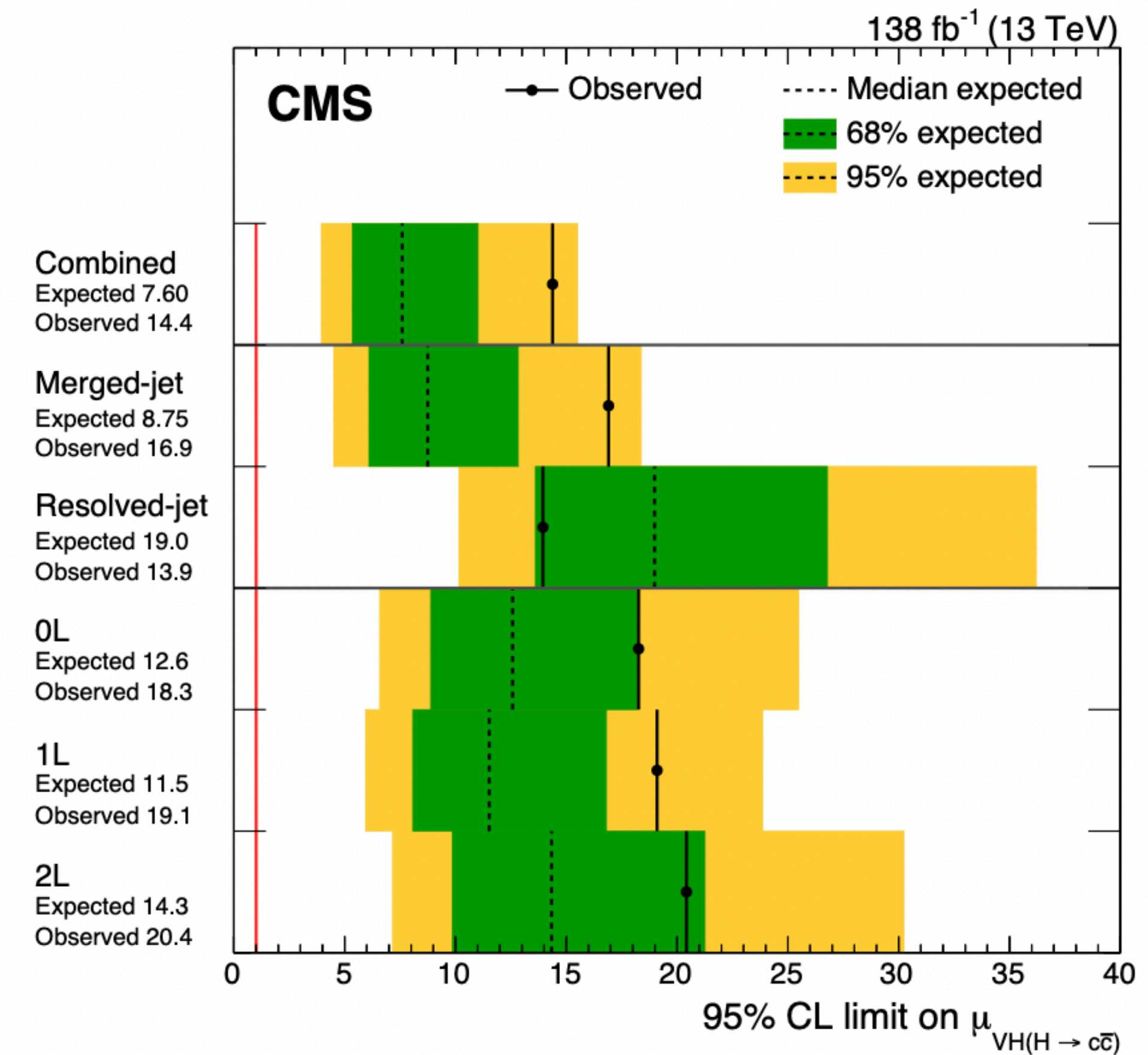
expected $|k_c| < 3.4$ observed $1.1 < |k_c| < 5.5$ @95% CL

thanks to exploitation of flavour tagging + reconstruction of the m_{Higgs} through boosted large R-jet using modern Machine learning techniques.

HL-LHC Lumi → Expected sensitivity ↓	VH(→ c \bar{c})	
	3 ab ⁻¹	2.5 ab ⁻¹
ATLAS	-	-
CMS	1.3	1.2
Combined	1.9	1.7

Adding inclusive Higgs and the VBF production modes +various improvements could lead to first direct evidence for the Yukawa coupling of the Higgs boson to charm at HL-LHC

It is therefore extremely important as an intermediate goal of Run 3 that progress is shown by all experiments in improving their sensitivity in this channel:



Graph nets



Graph nets

- Graph nets can be neural networks operating on graphs, but can be implemented with functions very different from neural networks. [arXiv:1806.01261v3](https://arxiv.org/abs/1806.01261v3)
Networks acting on a “graph” rather than a vector of inputs, with output being a graph: Lot of activity on this in the past years in industry

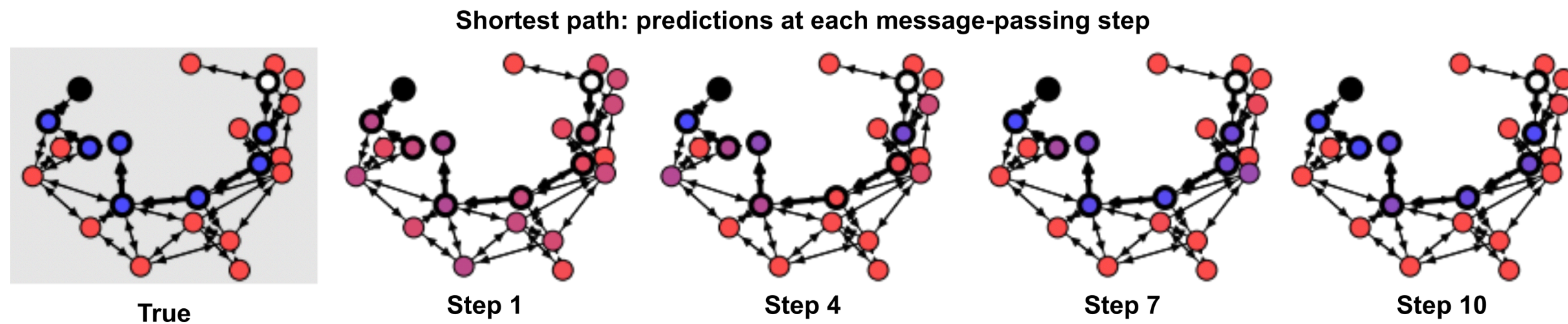
Here one can find open-source software library for building graph nets, with demonstrations on how to use them:

https://github.com/deepmind/graph_nets

- Quite some possible applications: they have been used already for a variety of cases:
 - to learn the dynamics of physical systems (Battaglia et al., 2016; Chang et al., 2017; Watters et al., 2017; van Steenkiste et al., 2018; Sanchez-Gonzalez et al., 2018)
 - to predict the chemical properties of molecules (Duvenaud et al., 2015; Gilmer et al., 2017)
 - to predict traffic on roads (Li et al., 2017; Cui et al., 2018)
 - to classify and segment images and videos (Wang et al., 2018c; Hu et al., 2017)
 - to perform semi-supervised text classification (Kipf and Welling, 2017)
 - in machine translation (Vaswani et al., 2017; Shaw et al., 2018; Gulcehre et al., 2018)...

Graph nets: demo

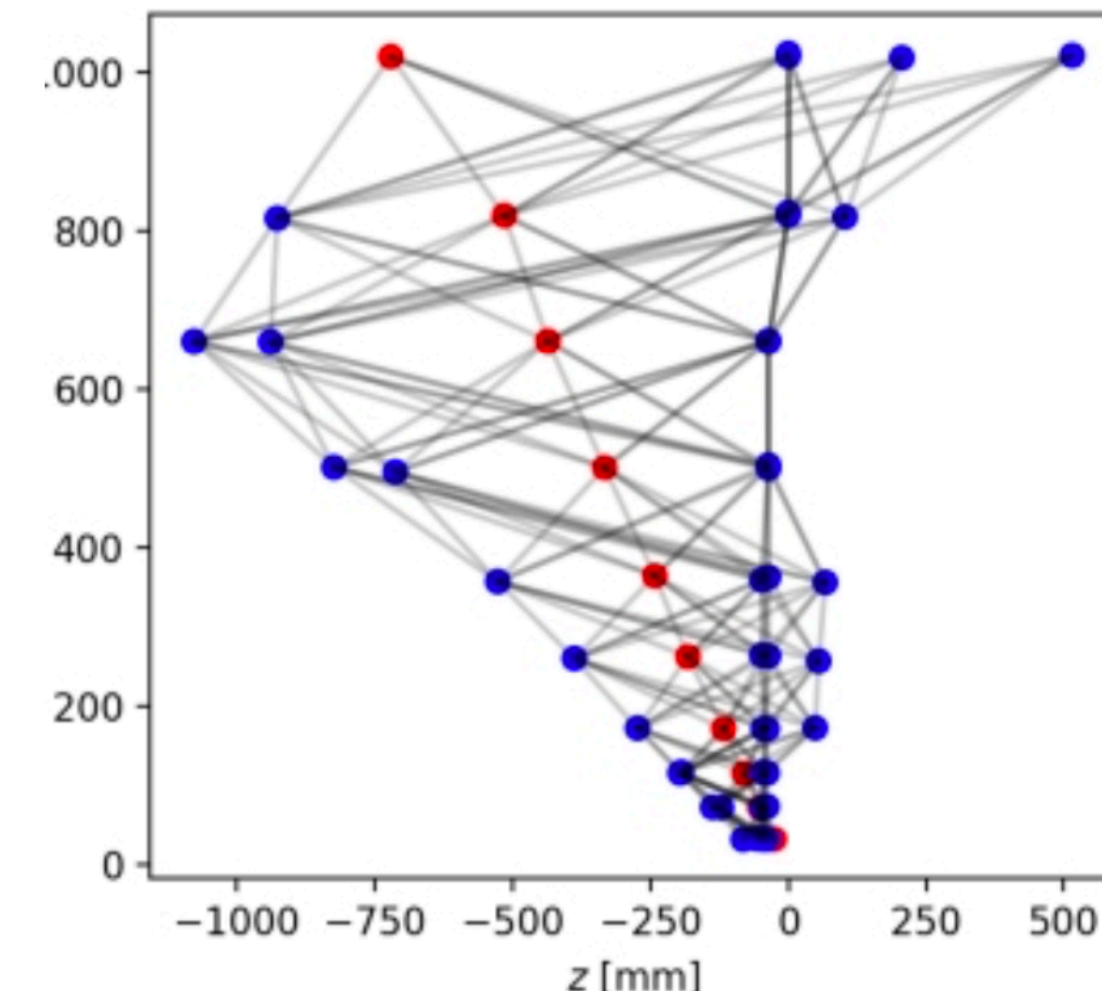
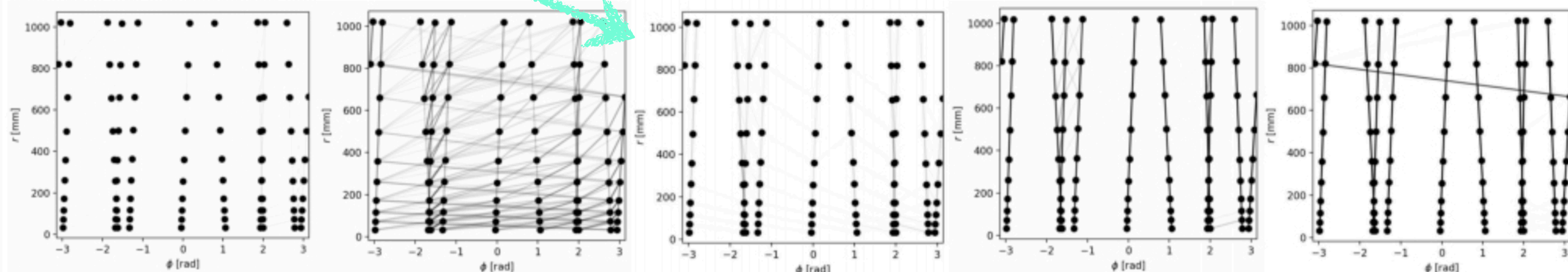
- Find the shortest path in a graph: [demo: tinyurl.com/gn-shortest-path-demo](http://demo.tinyurl.com/gn-shortest-path-demo)
This demo creates random graphs, and trains a GN to label the nodes and edges on the shortest path between any two nodes. Over a sequence of message-passing steps (as depicted by each step's plot), the model refines its prediction of the shortest path.



Where could we apply graph-nets?

- A great improvement could be achieved by applying graph-nets to tracking
- Tracking is a very time consuming reco task at LHC (most consuming?)
- When applying graph-nets to track building one could for example use them to pair hits

Successive iterations on an event

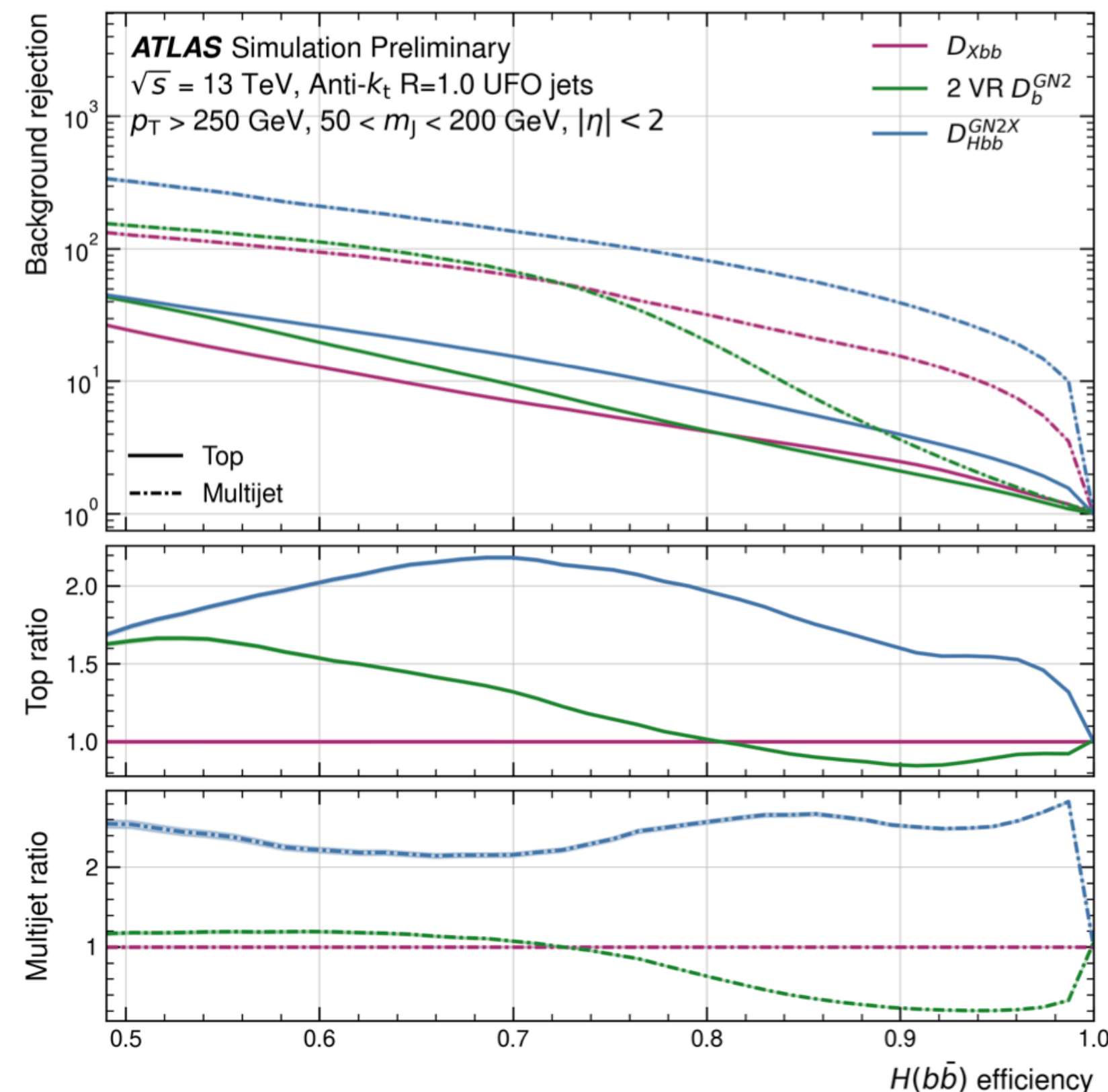
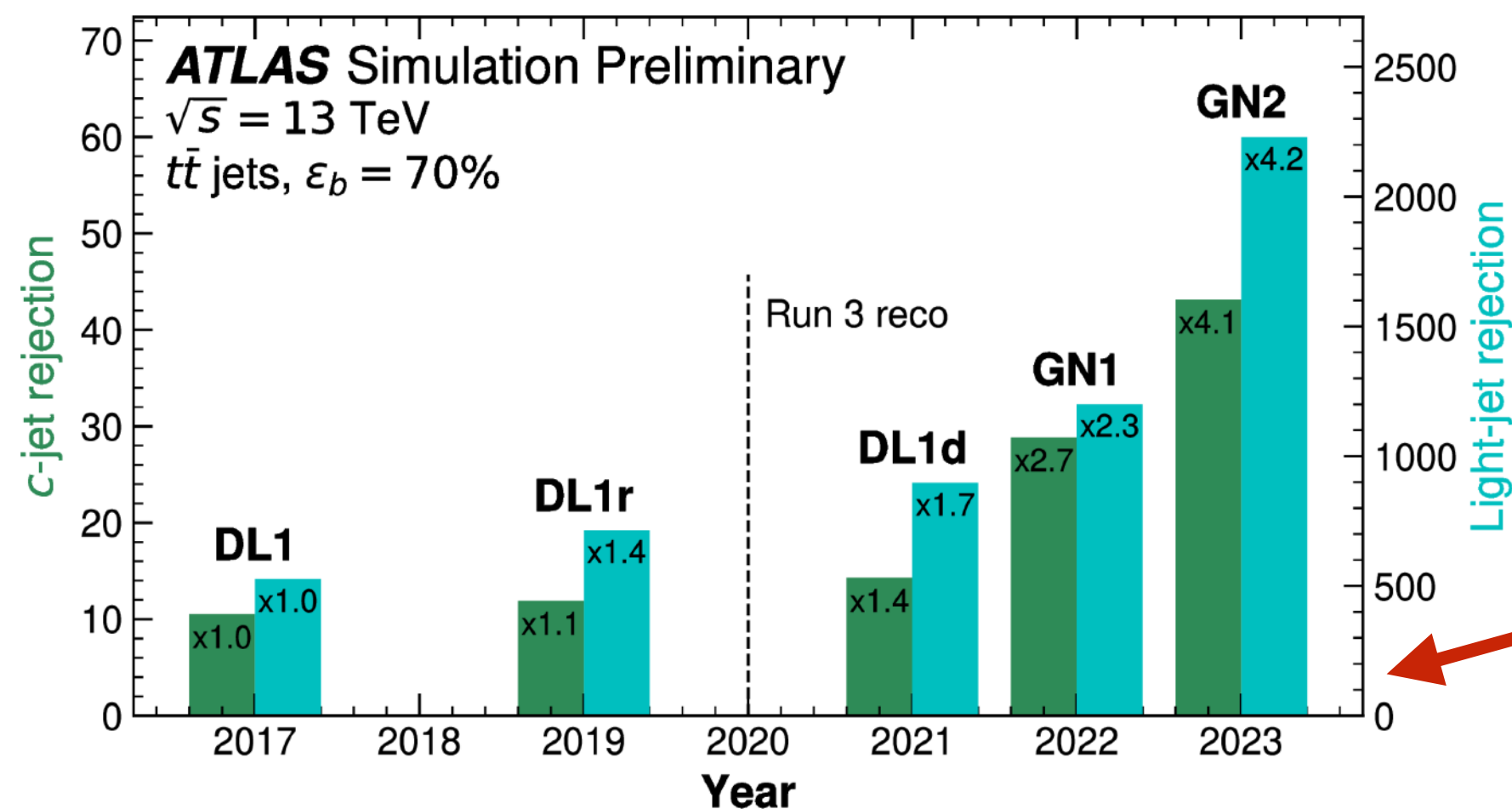


Flavor tagging in continuous evolution

Boosted H->bb/cc tagging ATL-PHYS-PUB-2023-021

- Boosted b-tagging: new algorithm, GN2X for large-radius jets: tagging boosted H(bb) jets and H(cc) jets.

small R-jet tagging Jet Flavour Tagging With GN1 and DL1d

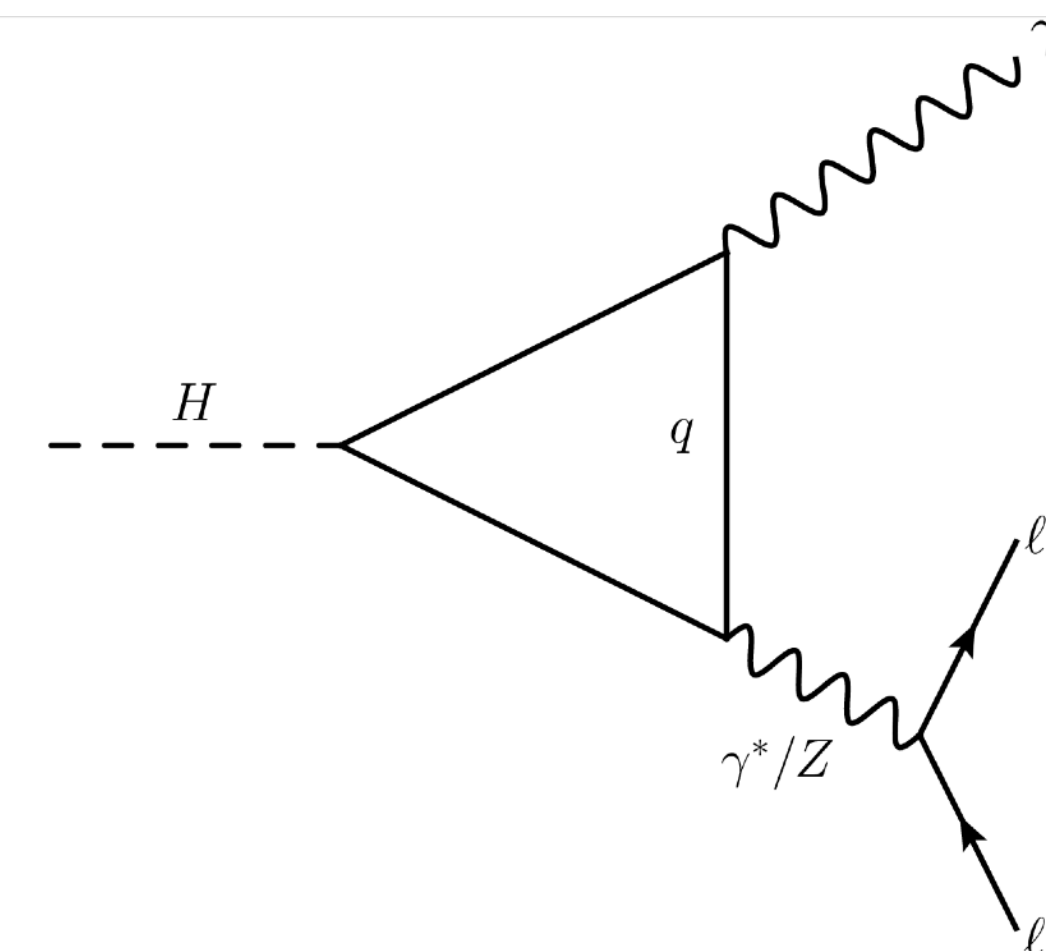


- GN2X benefits from advances in flavour tagging of small-radius jets with Graph Neural Networks (GNNs)

Rare processes: back on the envelope calculation based on SM expectations

Higgs

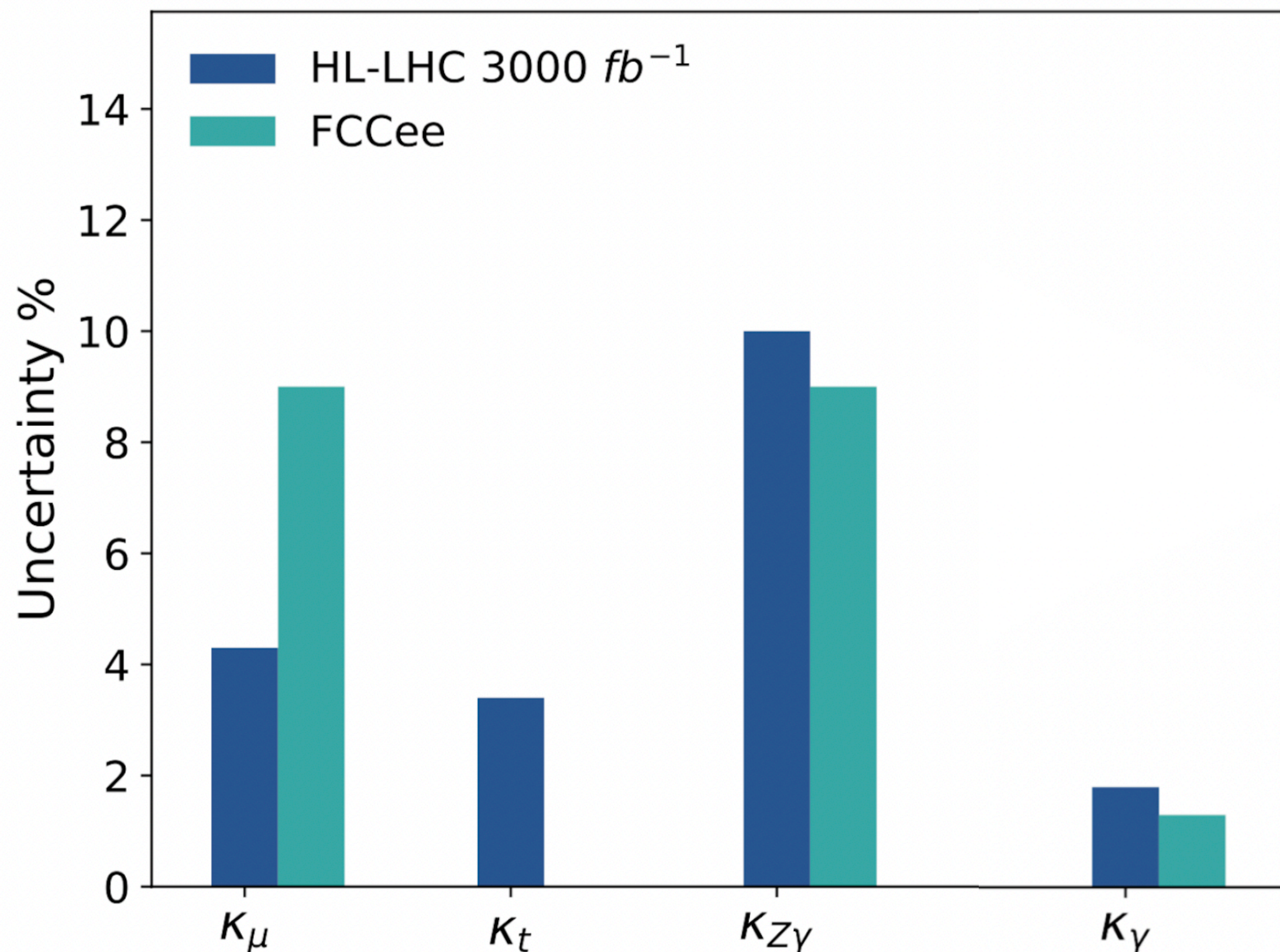
run3 Lumi→	H→μμ		H→yy*		H→Zy		HH	
	250 fb	200 fb	250 fb	200 fb	250 fb	200 fb	250 fb	200 fb
ATLAS	2.8	2.6	3.5	3.3	2.0	1.9	1.2	1.1
CMS	4.2	4.0	-	-	2.0	1.9	1.3	1.2
Combined	5.0	4.8	5.0	4.7	2.8	2.6	1.8	1.6



Can reach combined observation in Run 3

γ* is a virtual particle with (non zero) mass, decays instantly to two leptons $m_{\ell\ell} < 30$ GeV (typically < 1 GeV)
 p_T , small leptons separation \sim cm
 (challenge for electrons)

What will the HL_LHC bring?



The ultimate precision on Higgs couplings reachable at HL-LHC and FCCee.

Improvements in experimental techniques and theoretical calculations will be needed to reach as close as possible to a $O(1\%)$ precision for all these observables.

Higgs factories cannot probe κ_t in a model independent way, and can only reach a $O(10\%)$ accuracy on κ_μ , κ_t , $\kappa_{Z\gamma}$, through loop effects in other decays, assuming no competing new physics contributions.



Higgs mass: great example of gain that reconstruction improvements can lead to

[arXiv:2308.07216](https://arxiv.org/abs/2308.07216)
[arXiv:2308.04775](https://arxiv.org/abs/2308.04775)

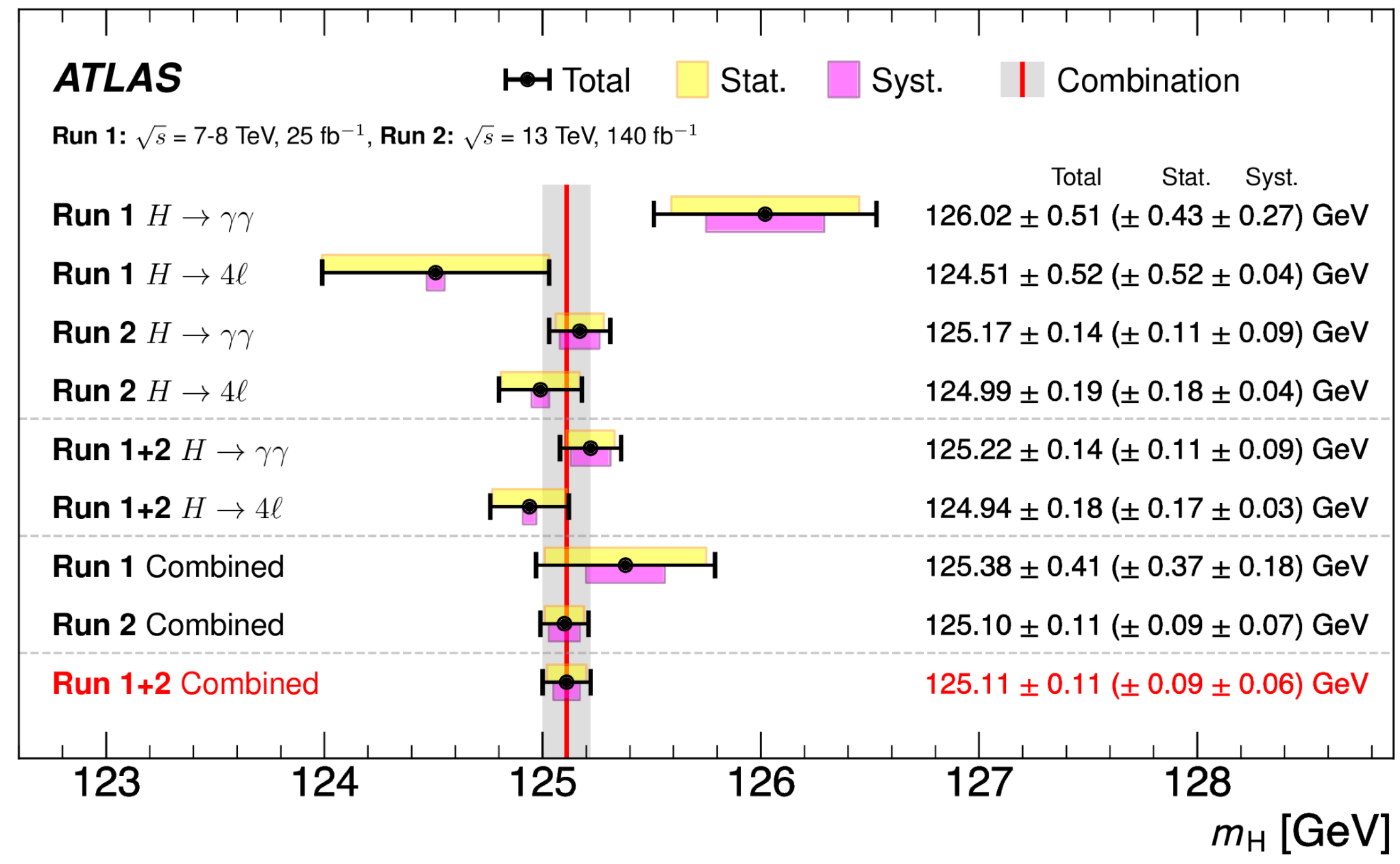
Combination of $H \rightarrow ZZ$ and $H \rightarrow \gamma\gamma$ provides:
most precise m_H measurement at 0.09%

$m_H = 125.11 \pm 0.11 \text{ GeV}$

Profits of various performance improvements:

- ~4x improvements in photon energy calibration!
 - due to 30% improvement in systematics: EM calorimeter layer calibration, measure of E lost around e/ γ clusters.
 - Residual electron E scale non-linearities used for first time to constrain systematic uncertainties \rightarrow further x2 improvement

**\rightarrow Reduces $H \rightarrow \gamma\gamma$ systematics by factor 4:
320 MeV \rightarrow 80 MeV**





Higgs mass great example of gain that reconstruction improvements can lead to

ATLAS measures

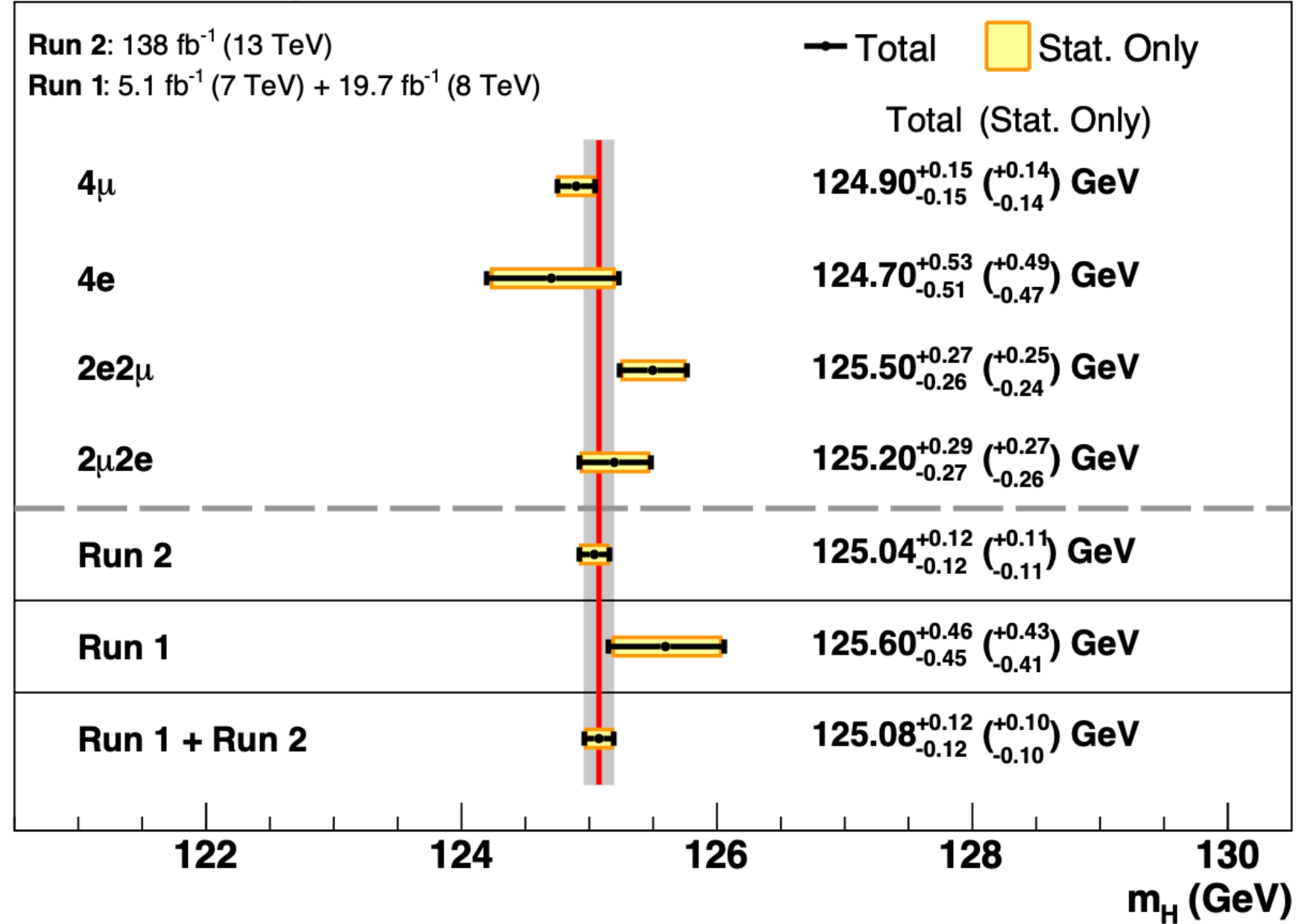
$$m_H = 125.11 \pm 0.11 \text{ GeV}$$

CMS measures in H->ZZ channel

$$m_H = 125.04 \pm 0.12 \text{ (stat.)} \pm 0.05 \text{ (syst.) GeV}$$

Great agreement among the 2 experiments!

CMS Preliminary



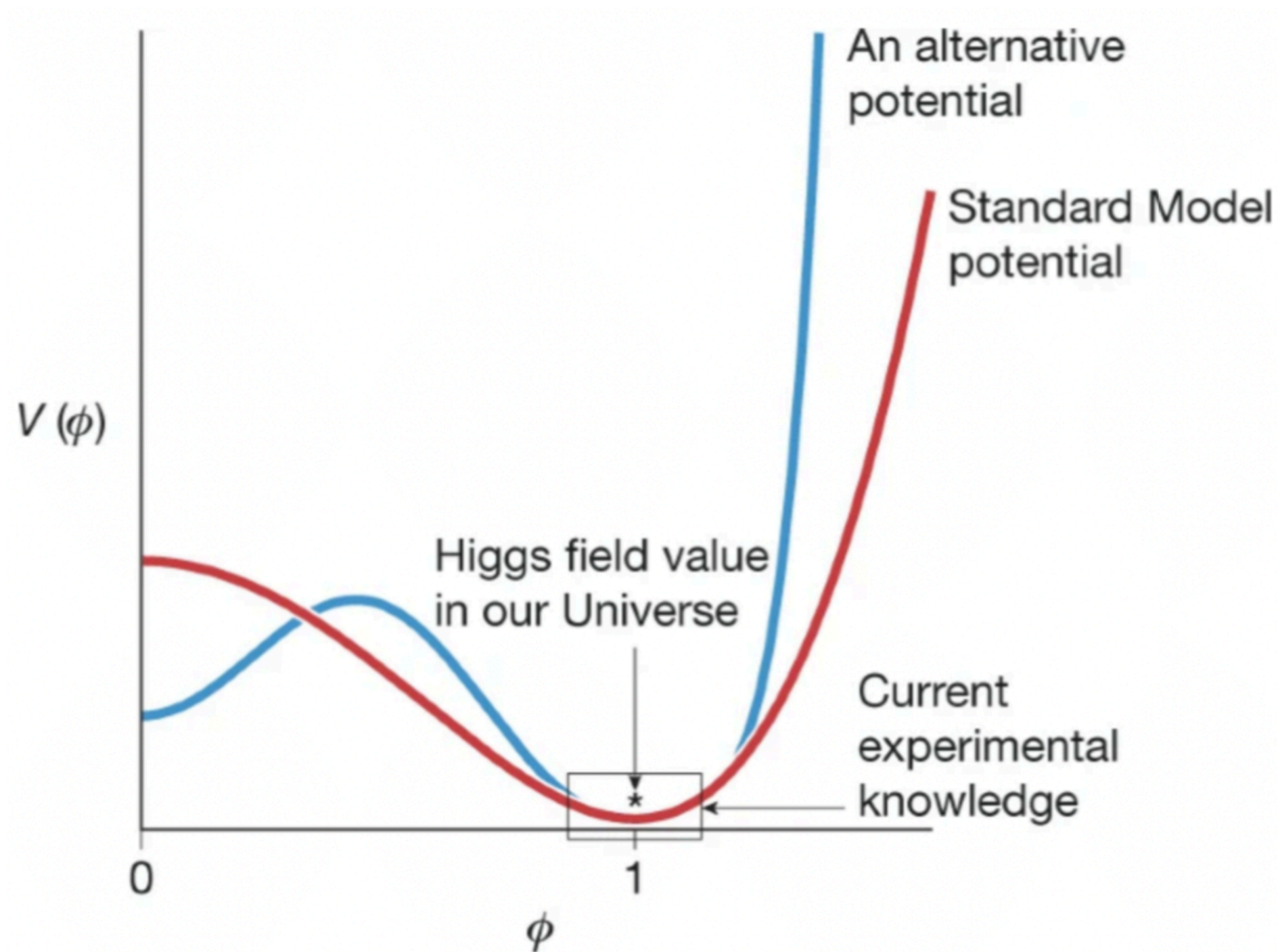
→ CMS $H \rightarrow ZZ$ most precise single measurement



Higgs potential and self couplings

Understanding the shape of the Higgs potential is fundamental

[G. Salam, Nature 607, 41–47 (2022)]



mass term, indicating a physical particle, the Higgs boson

$$V(\phi) = -\frac{\mu^4}{4\lambda} - \mu^2 H^2 + \lambda v H^3 + \dots$$

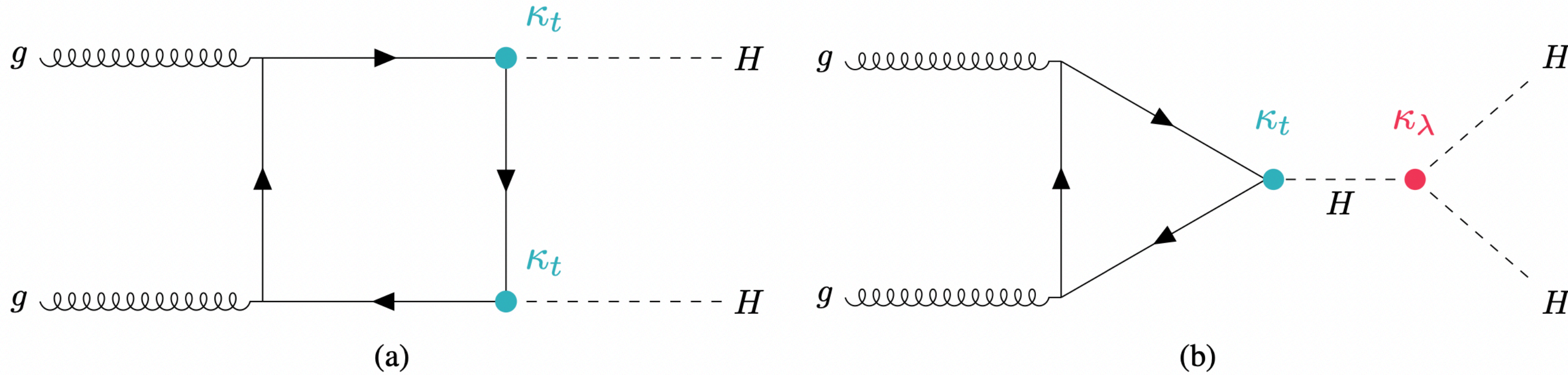
Higgs self-interaction term
 direct probing of Higgs self-interaction
 and the shape of Higgs potential

deviations from the SM would indicate new physics



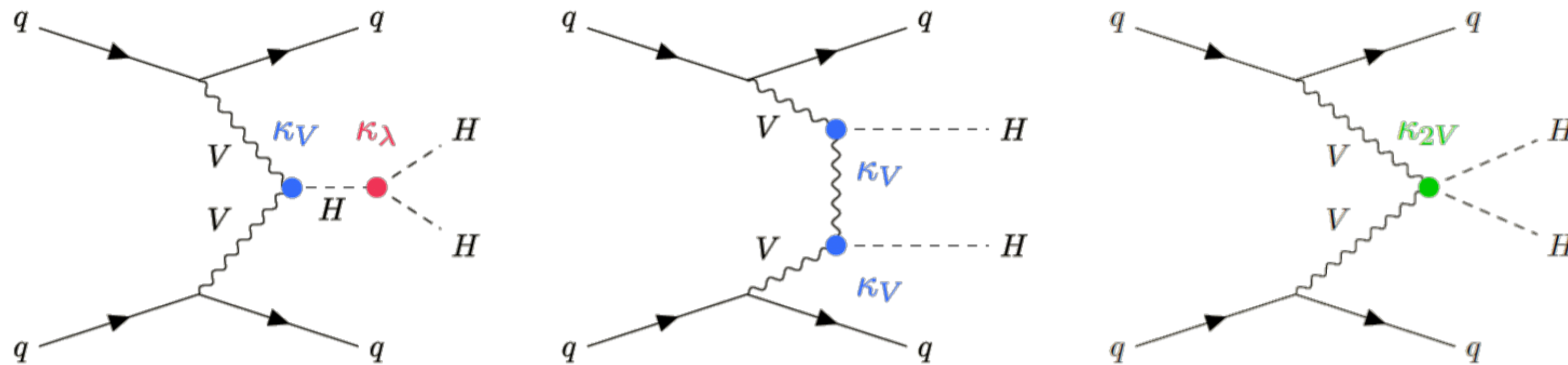
di-Higgs production at LHC

dominant production mode ggF 31 fb[13TeV] with 2 diagrams that have destructive interference



κ_λ = ratio of the Higgs boson self-coupling to its SM value

other dominant modes



Associated productions, HHV, HHtt have much smaller production cross-sections



di-Higgs decay modes

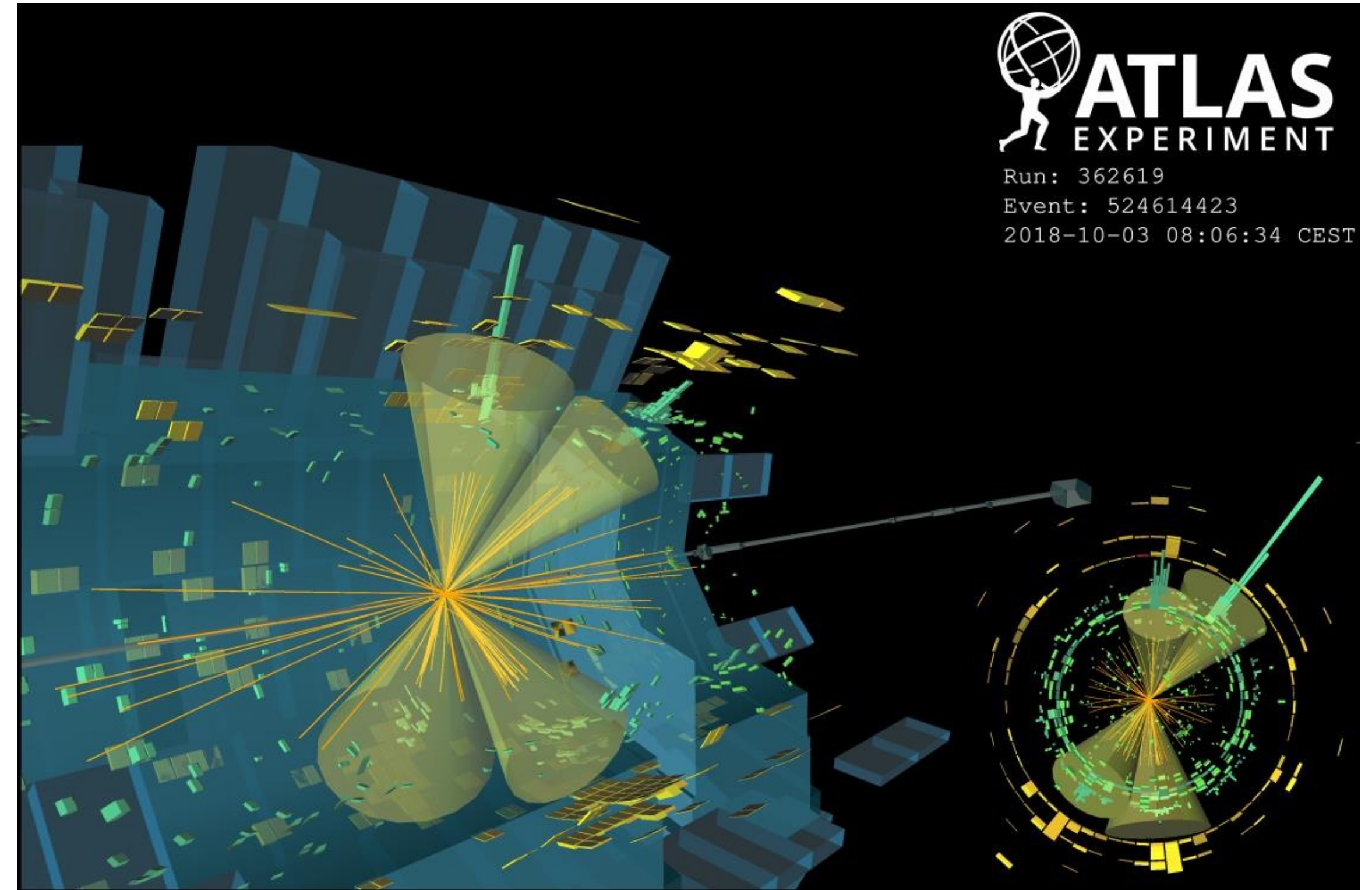
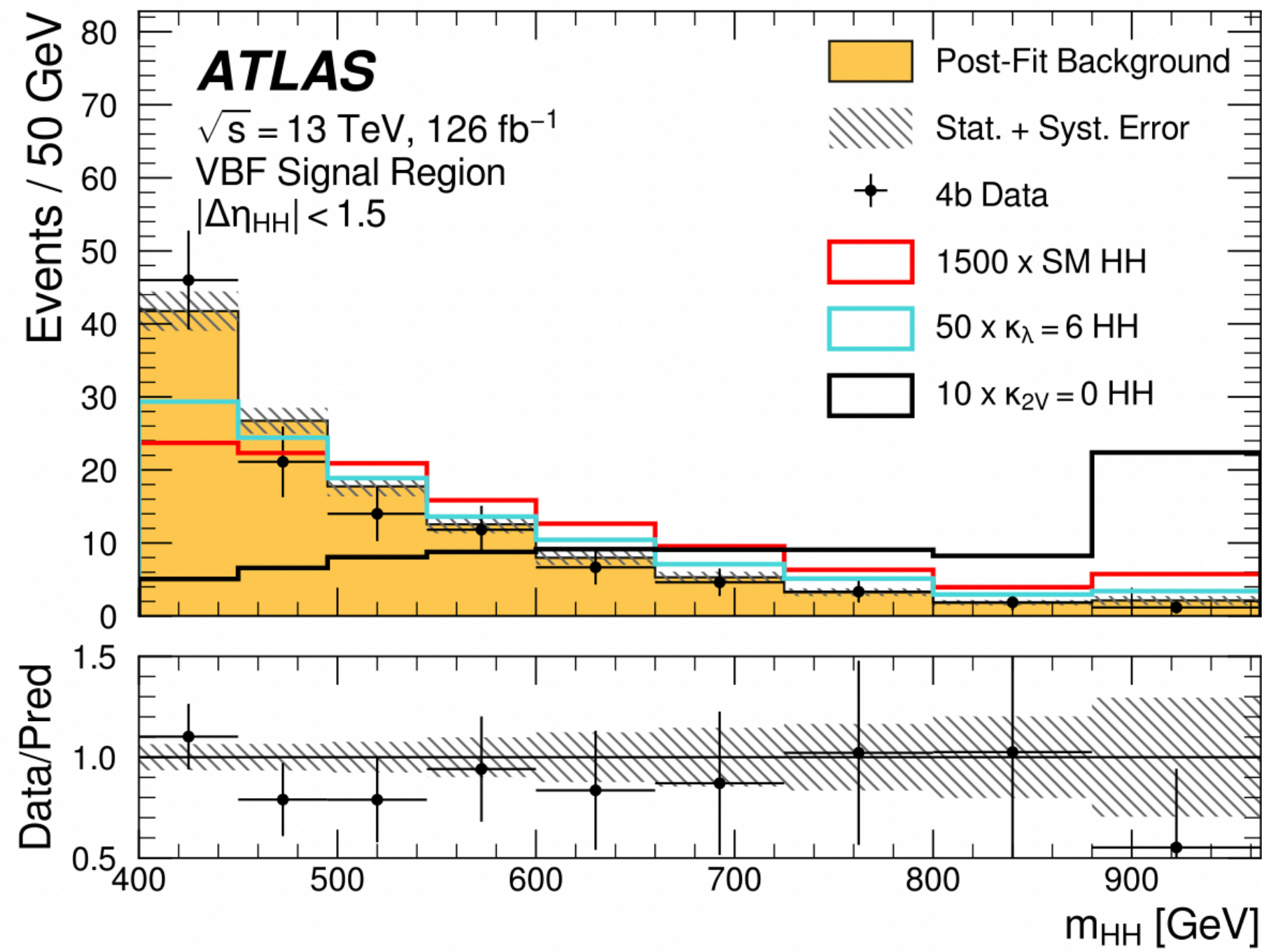
by Katharine Leney

	bb	WW	$\tau\tau$	ZZ	$\gamma\gamma$
bb	34%				
WW	25%	4.6%			
$\tau\tau$	7.3%	2.7%	0.39%		
ZZ	3.1%	1.1%	0.33%	0.069%	
$\gamma\gamma$	0.26%	0.10%	0.028%	0.012%	0.0005%

The golden channels



HH \rightarrow *bbbb*

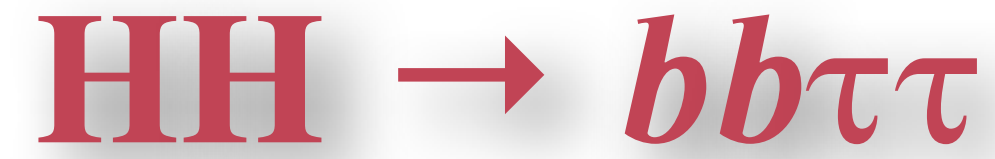


Highest BR, sensitive to high p_T of the Higgs

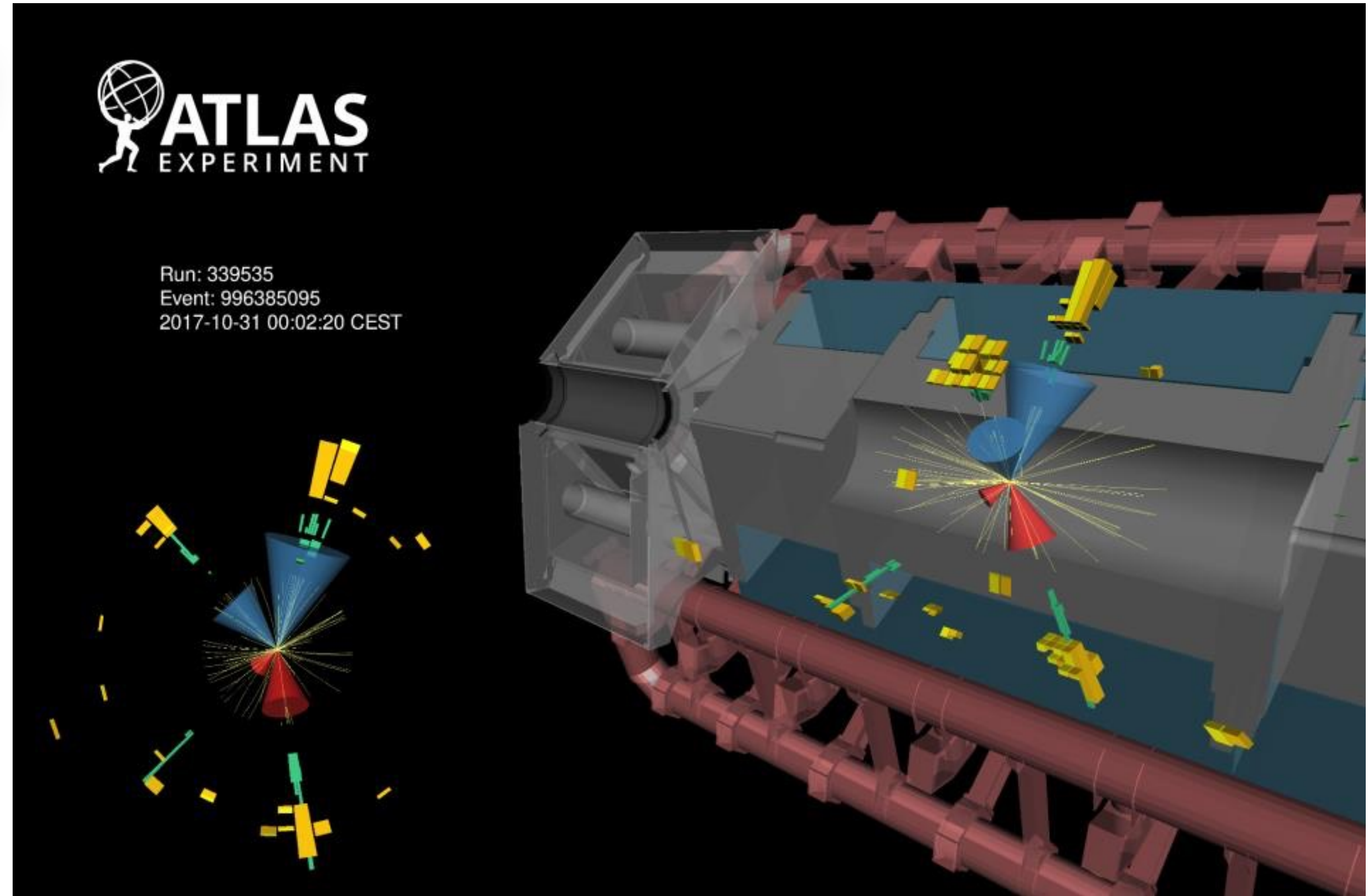
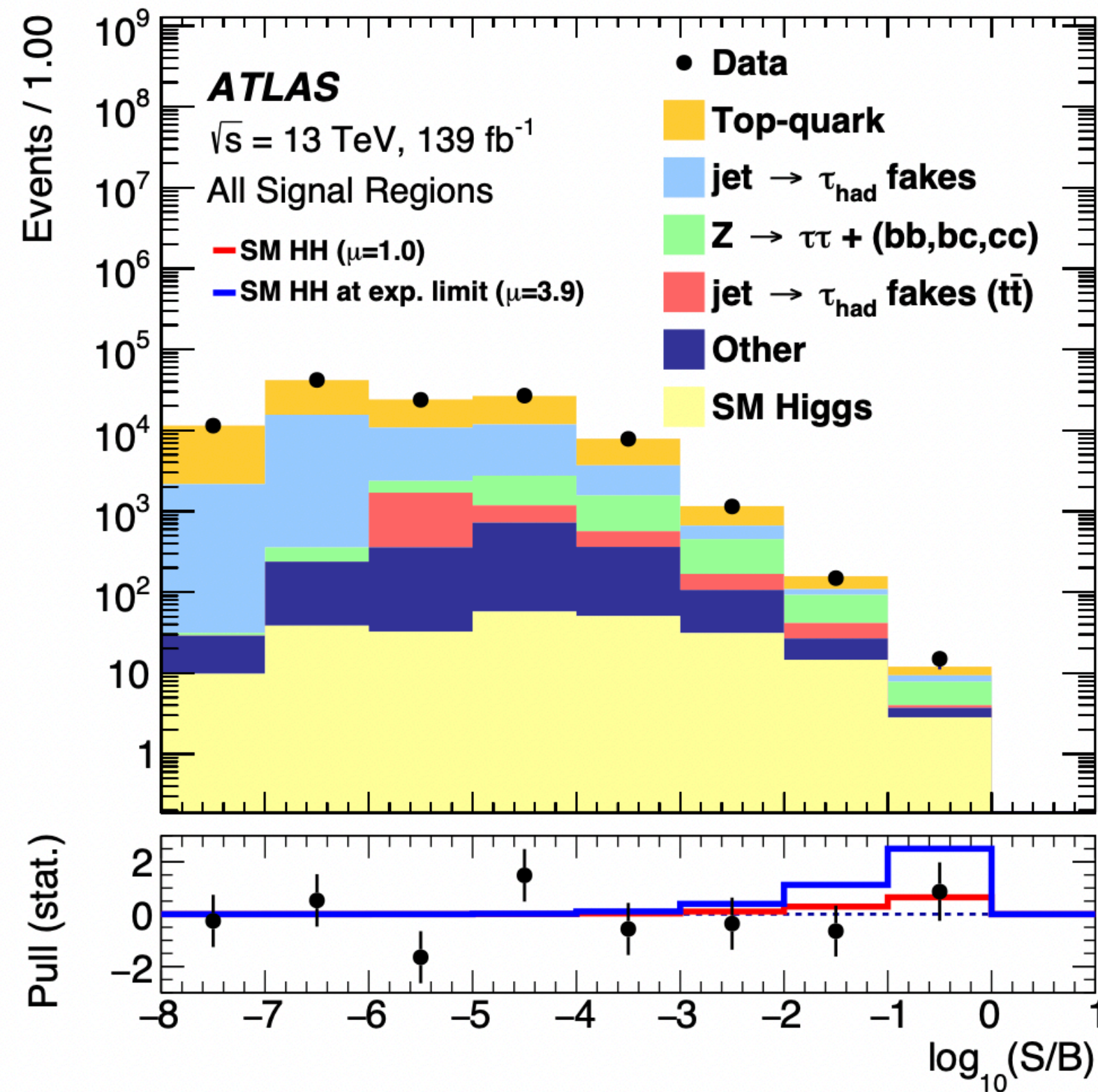
The sensitivity of the analyses is improved relative to previous iterations by using more sophisticated background modeling techniques, event categorization and improved jet reconstruction and flavor identification algorithms, in addition to the increased integrated luminosity of the analyzed data.

	Observed Limit	-2σ	-1σ	Expected Limit	$+1\sigma$	$+2\sigma$
μ_{ggF}	5.5	4.4	5.9	8.2	12.4	19.6
μ_{VBF}	130	70	100	130	190	280
$\mu_{ggF+VBF}$	5.4	4.3	5.8	8.1	12.2	19.1

both in $T_{hadThad}$ and $T_{lepThad}$ channel
and in $ggf+VBF$ production



10



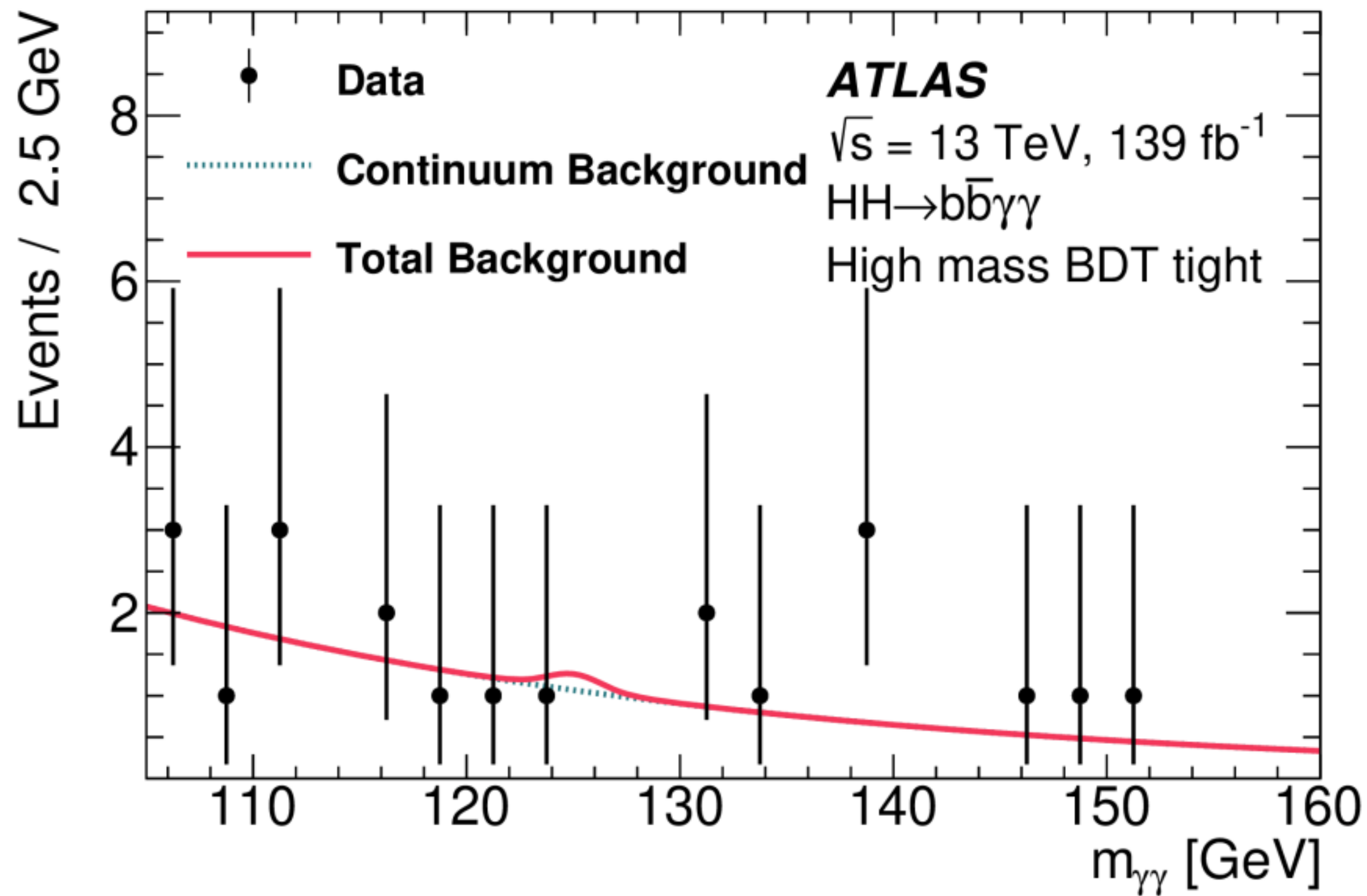
Factor 4 improvement wrt to previous version of analysis.
Half of this improvement is due to the larger dataset,
while most of the remaining sensitivity gain is due to significant improvements in the $\tau_{had-vis}$ and b -jet reconstruction and identification.

HH

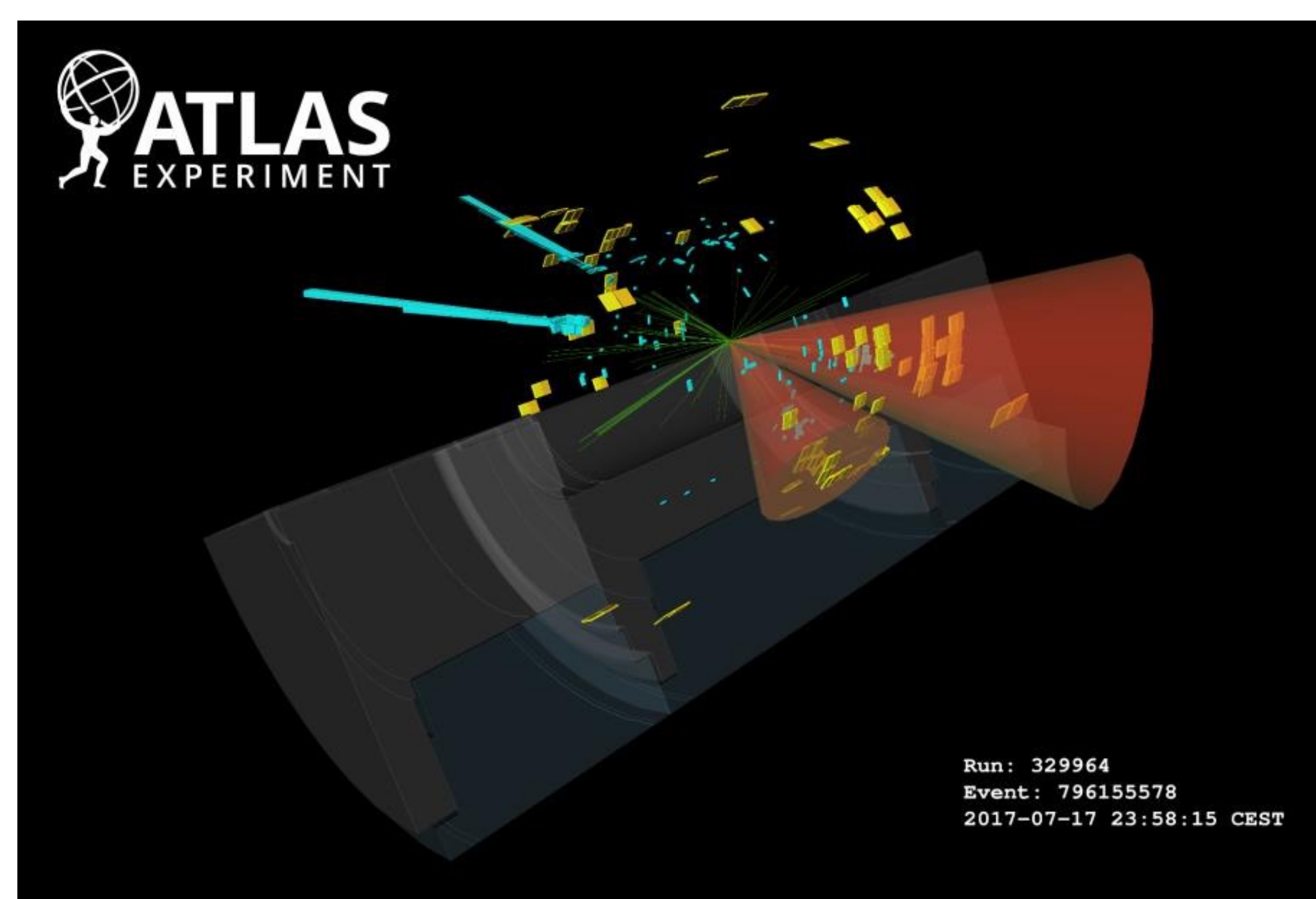
HH

HH \rightarrow $bb\gamma\gamma$

Phys. Rev. D 106 (2022) 052001



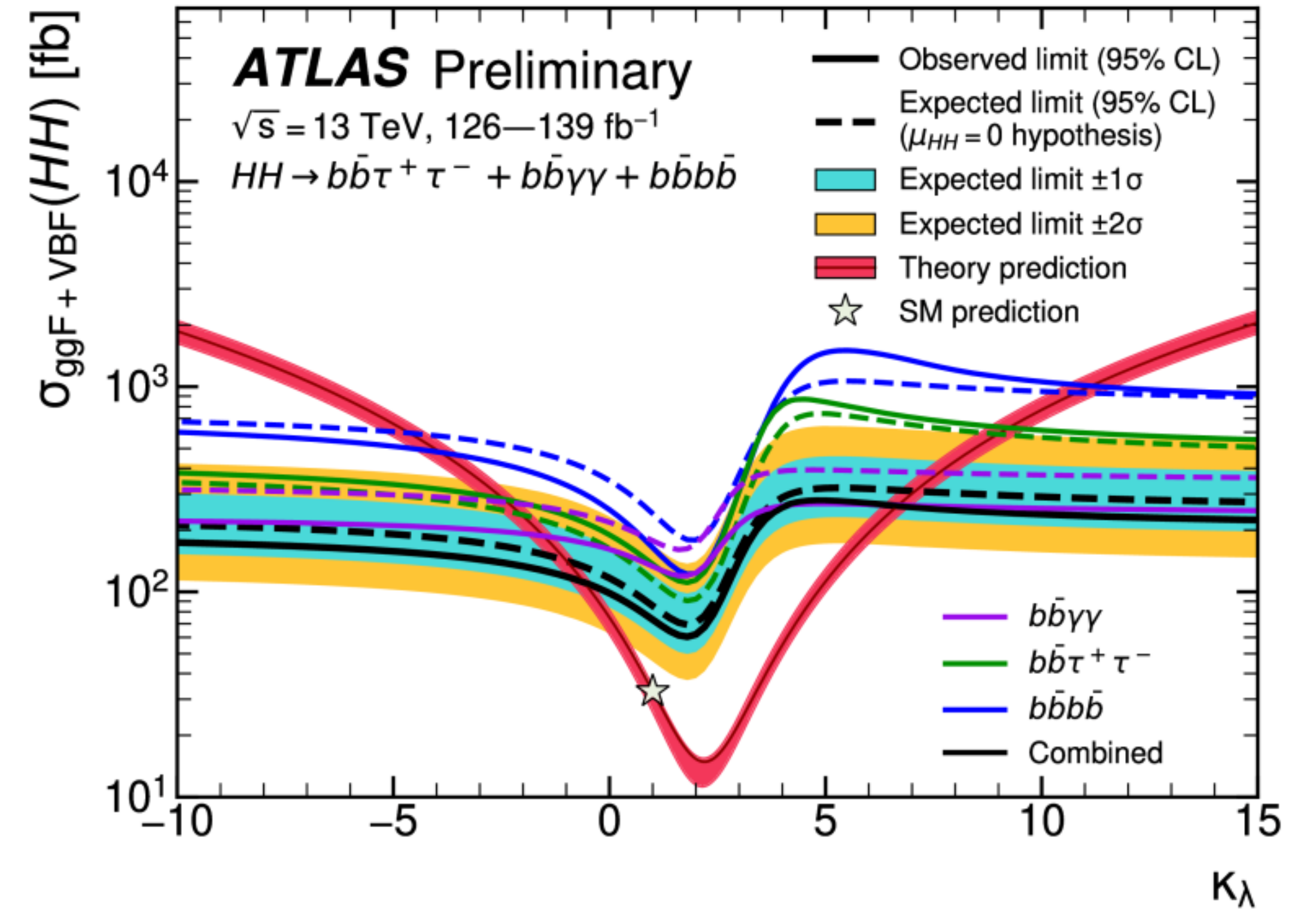
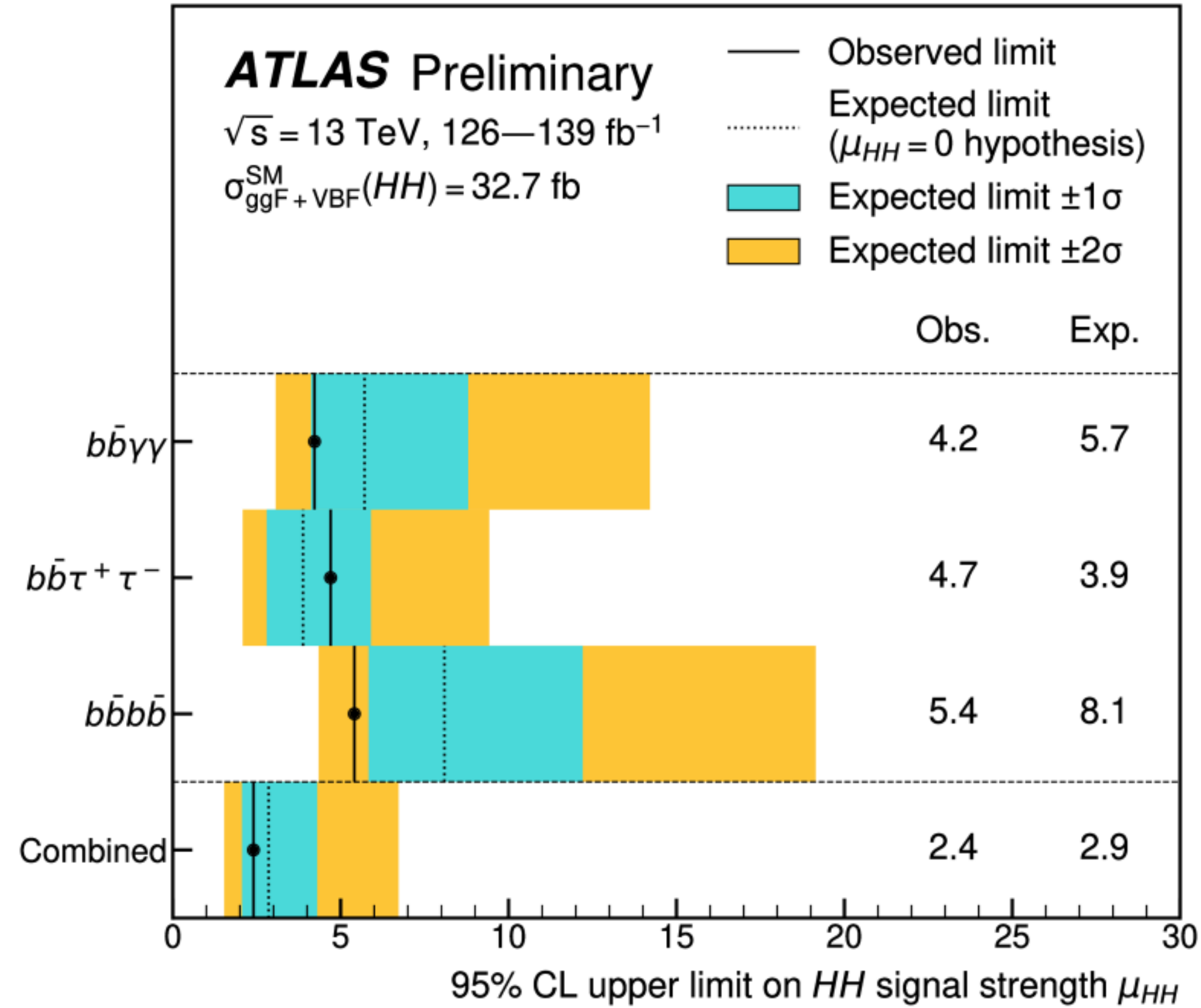
Limit on σ 4.2 (4.7) times the SM prediction @ 95% CL



four times larger dataset, incorporates a categorization based on $m_{bb\bar{\gamma}\gamma}$ and multivariate event selections, and expands analyzed mass range of the resonance search to lower values.

The results improve upon the previous ATLAS limits on the $HH \rightarrow b\bar{b}\gamma\gamma$ production cross section by up to a factor of 5 (half due to improved analysis)

HH combination

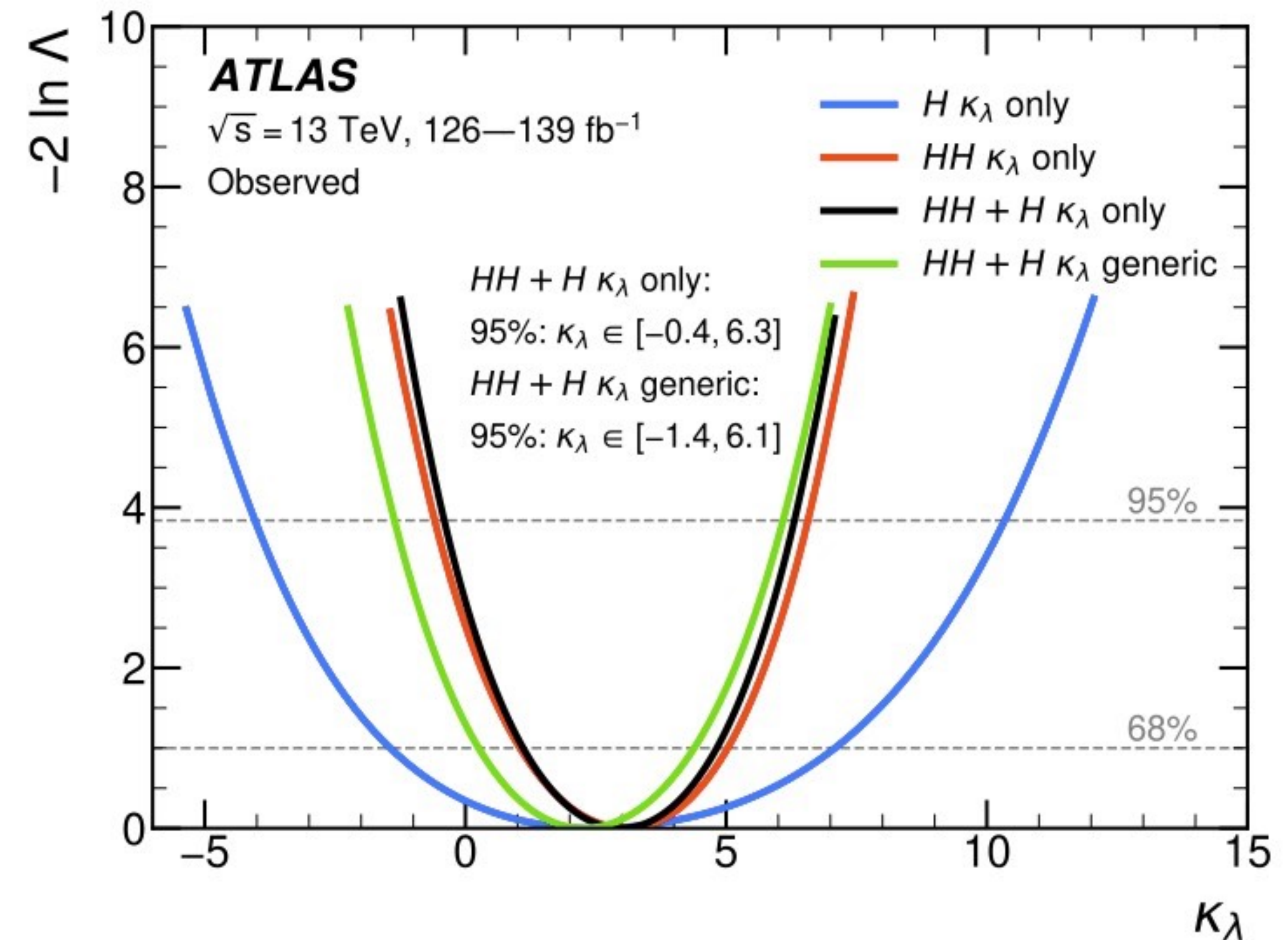
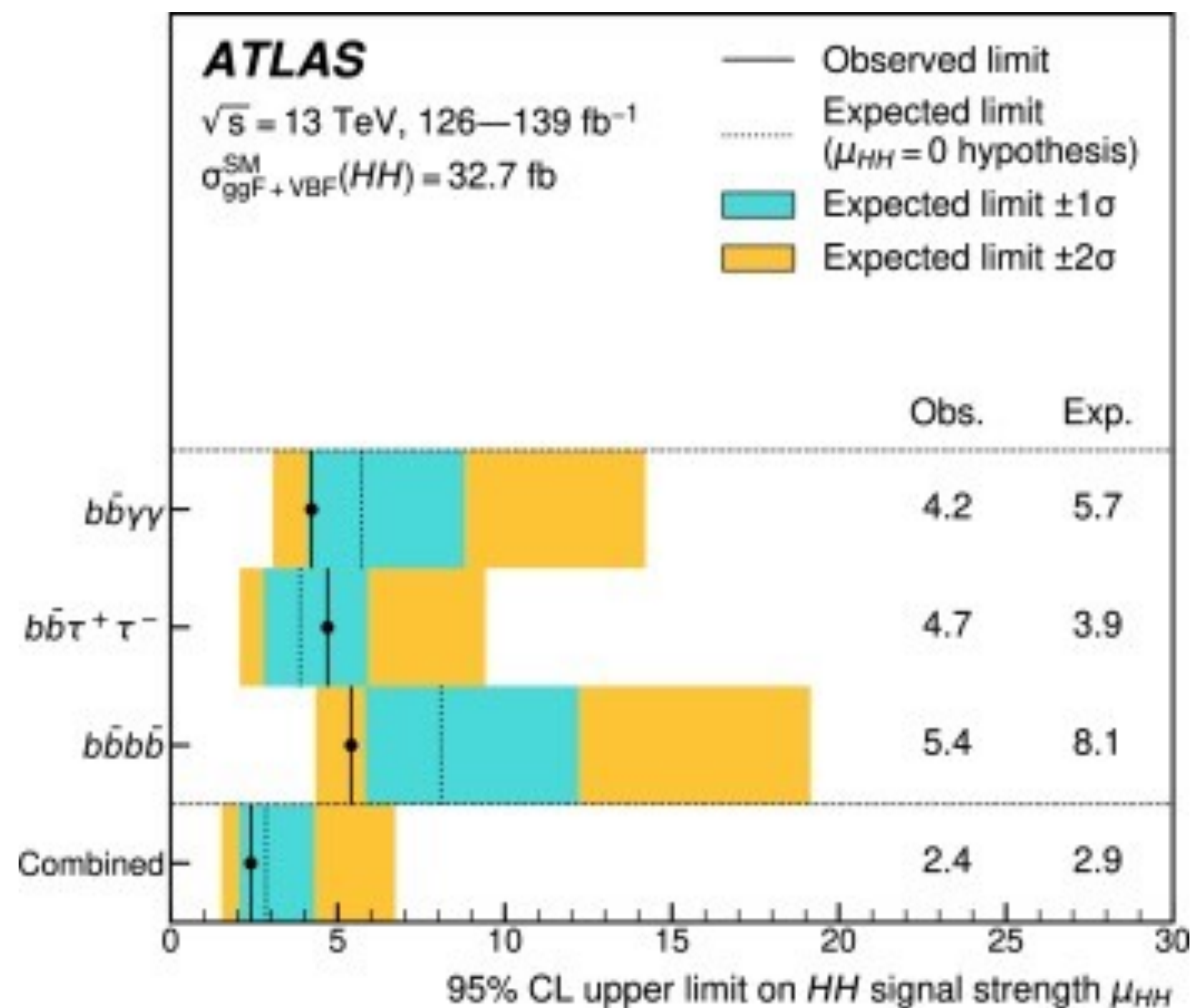


Limit on σ 2.4 (2.9) times the SM prediction at 95% CL

Combination of di-Higgs searches

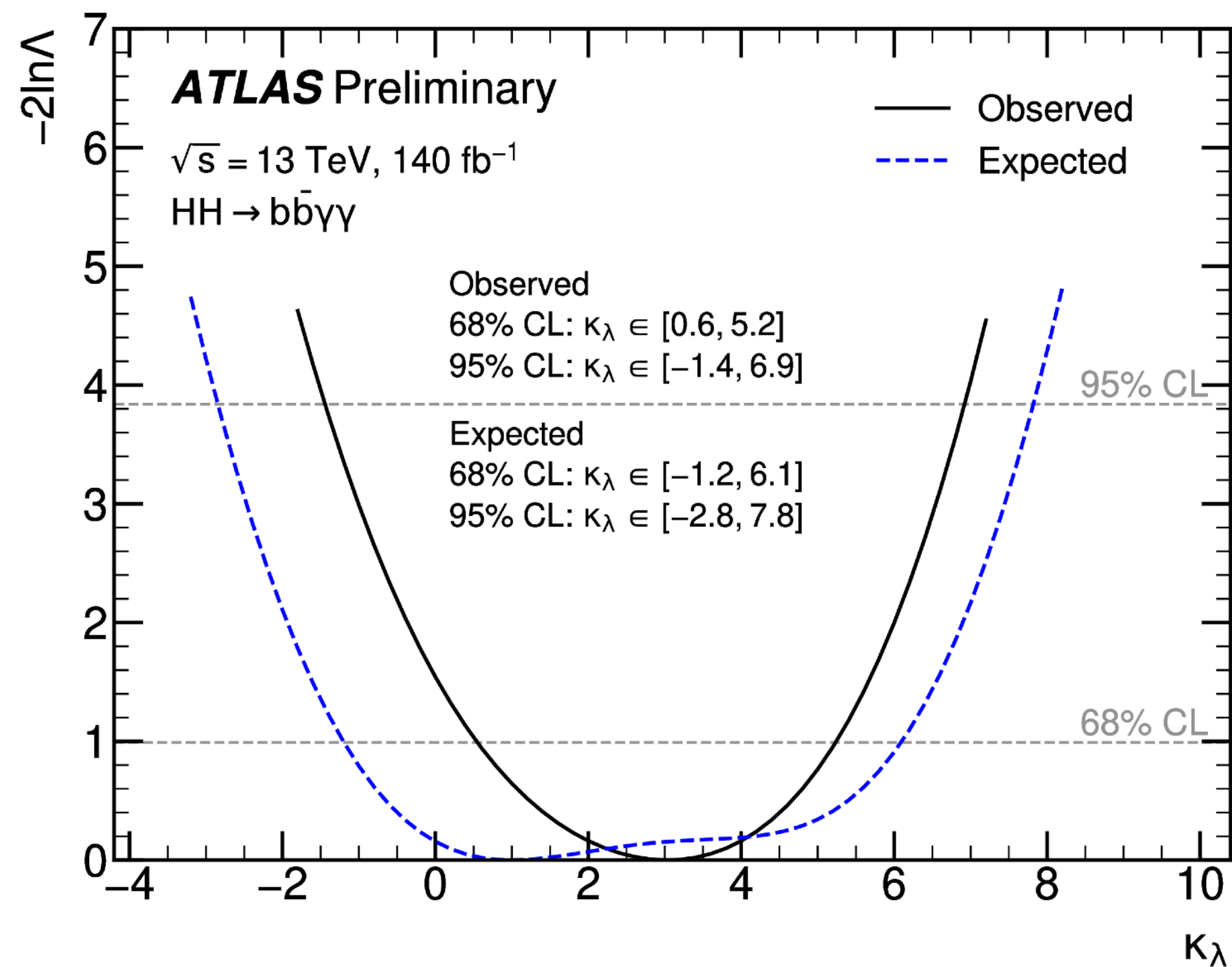
Phys. Lett. B 843 (2023) 137745

- $HH \rightarrow 4b, b\bar{b}\gamma\gamma, b\bar{b}\tau^+\tau^-$ have been combined with single Higgs results
- μ_{HH} : $2.4 \times \text{SM}$ ($2.9 \times \text{SM}$ exp.) at 95% CL
- $-0.4 < \kappa_\lambda < 6.3$ @95%CL (HH+H combination)

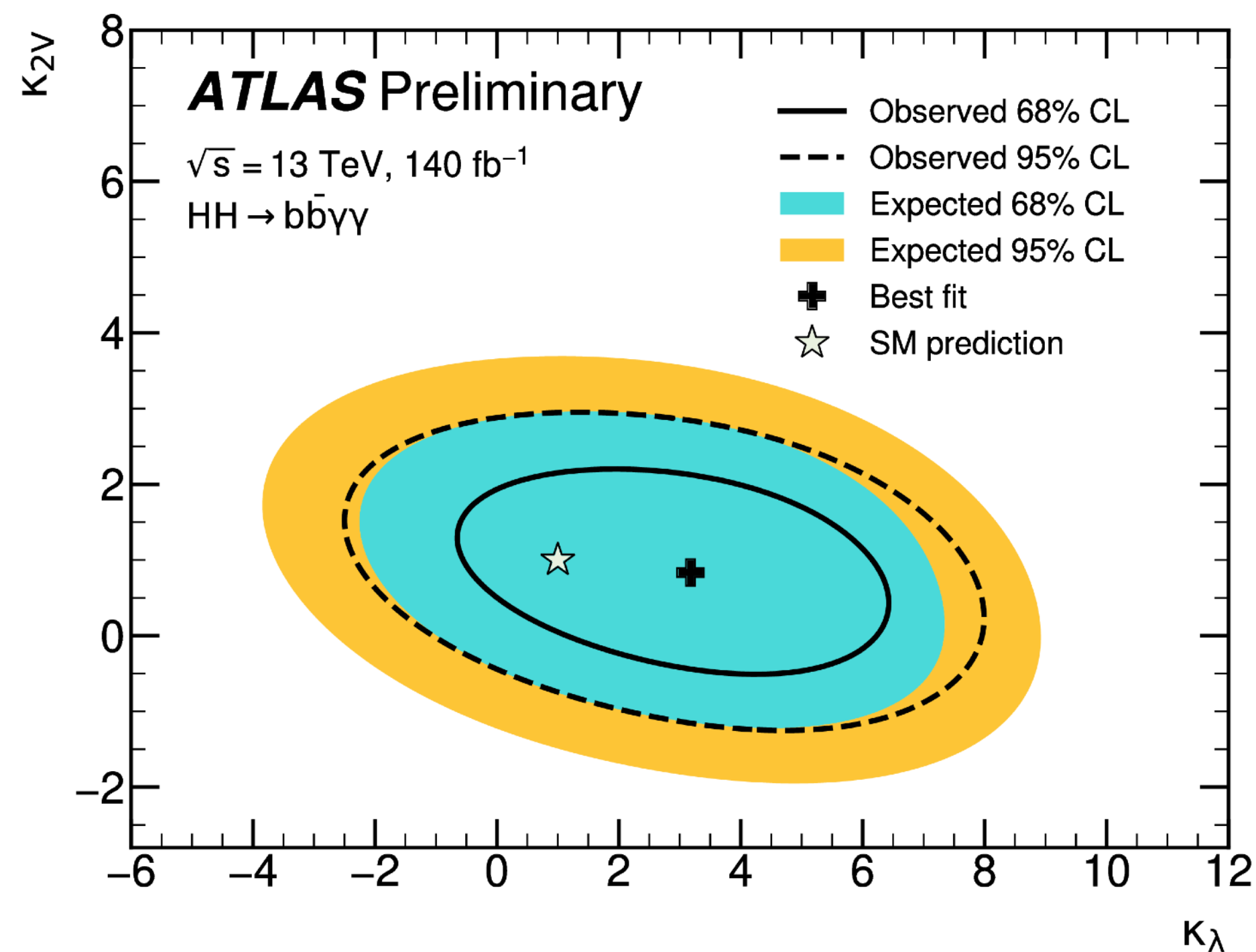
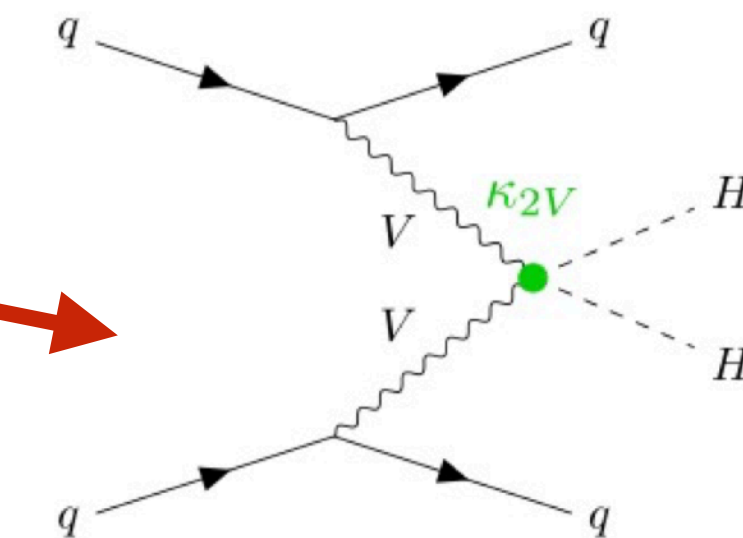


Most stringent limits to date

Further improved HH \rightarrow $\gamma\gamma bb$ search



- BDT used in 7 categories
- sensitive to H self-coupling λ and κ_{2V}

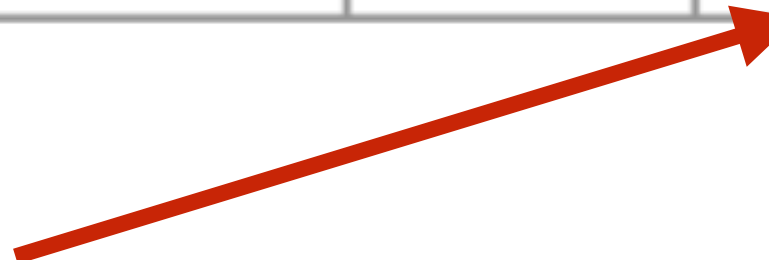
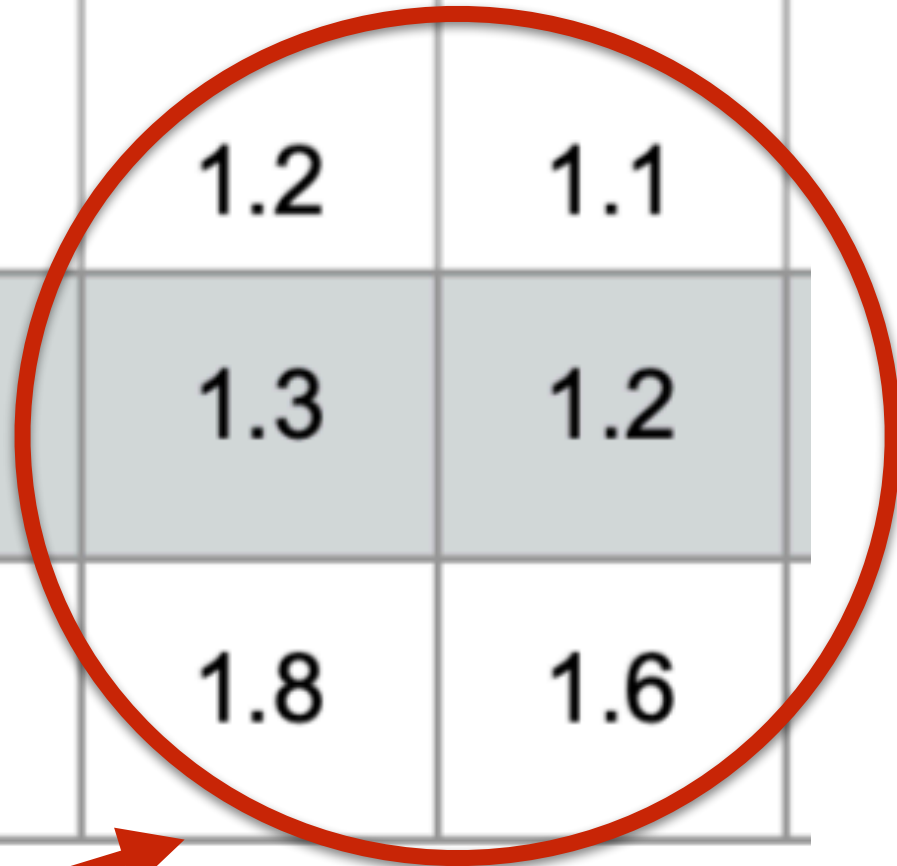


- Expected improvement μ_{HH} (12%), κ_λ (6%), κ_{2V} (17%) (mostly owing to event categorization)
- $\mu_{HH} < 4$ @95%CL
- not including most recent improvements in b-tagging!

Rare processes: back on the envelope calculation based on SM expectations



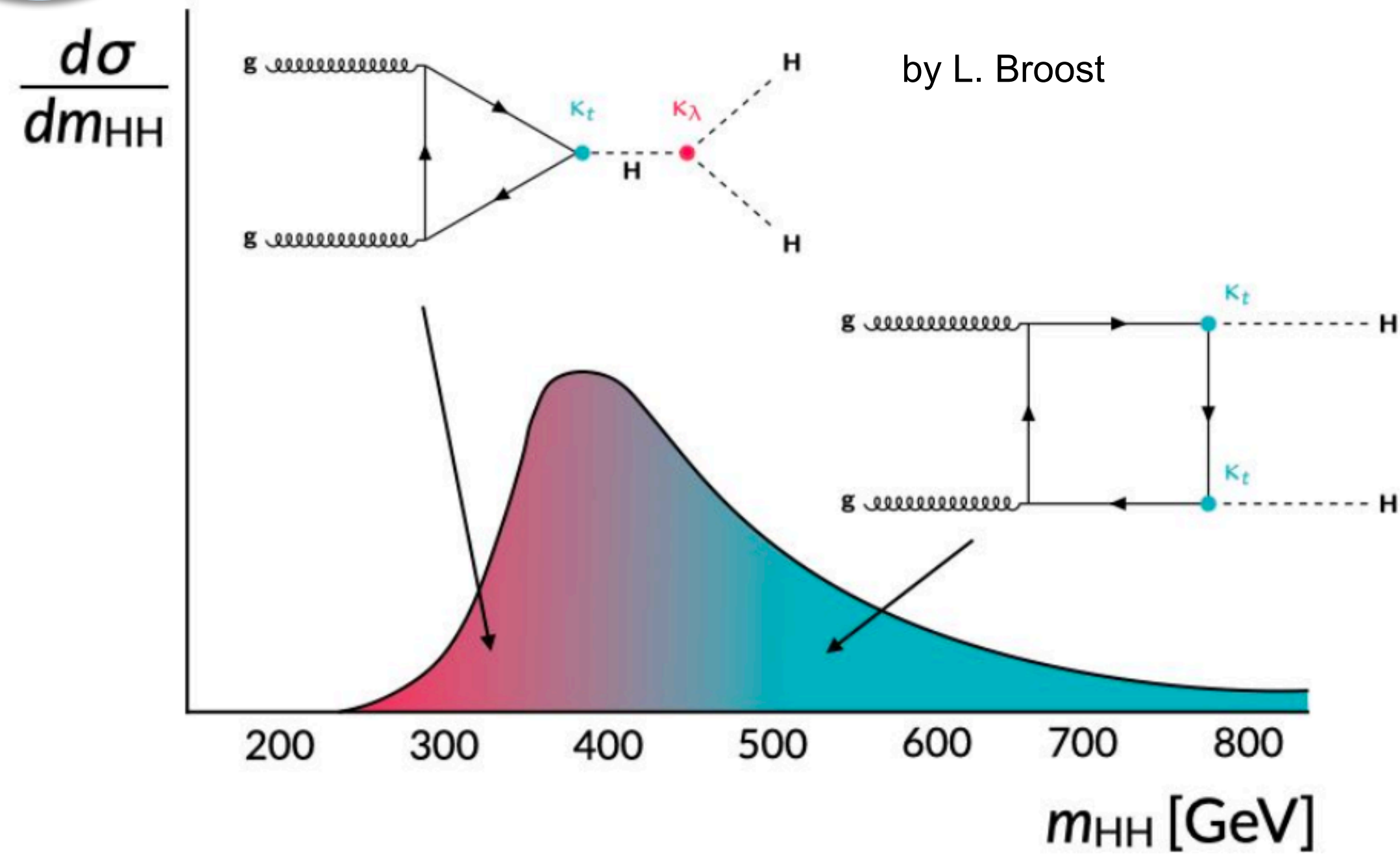
run3 Lumi→	H→μμ		H→yy*		H→Zγ		HH	
	250 fb.	200 fb.	250 fb.	200 fb.	250 fb.	200 fb.	250 fb.	200 fb.
ATLAS	2.8	2.6	3.5	3.3	2.0	1.9	1.2	1.1
CMS	4.2	4.0	-	-	2.0	1.9	1.3	1.2
Combined	5.0	4.8	5.0	4.7	2.8	2.6	1.8	1.6



Doesn't include new b-taggers etc etc.

HH

Prospects



In addition to new flavor taggers, new channels etc,
the low m_{HH} regions drives sensitivity
therefore lowering thresholds including trigger is
fundamental, especially for the future

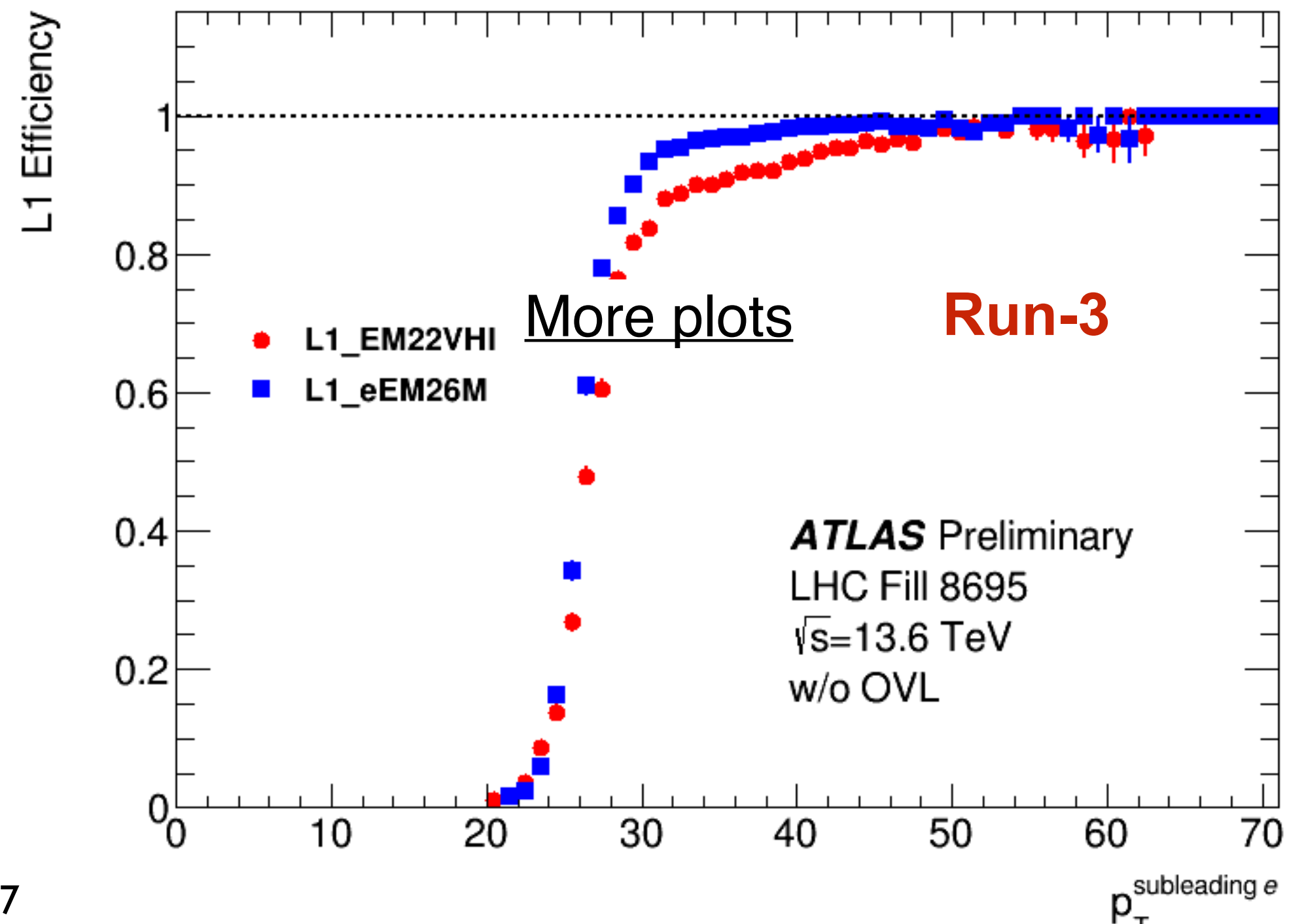
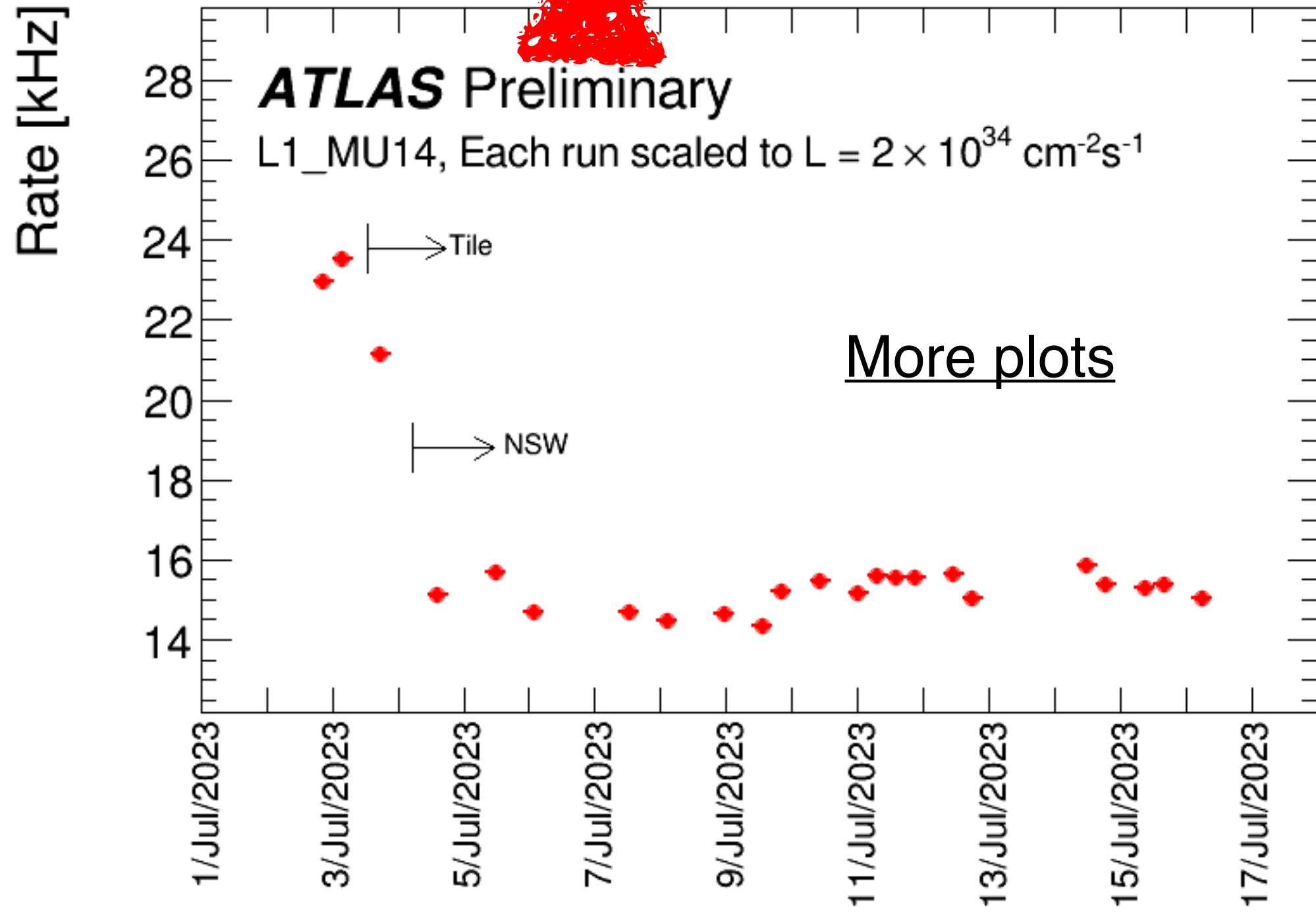
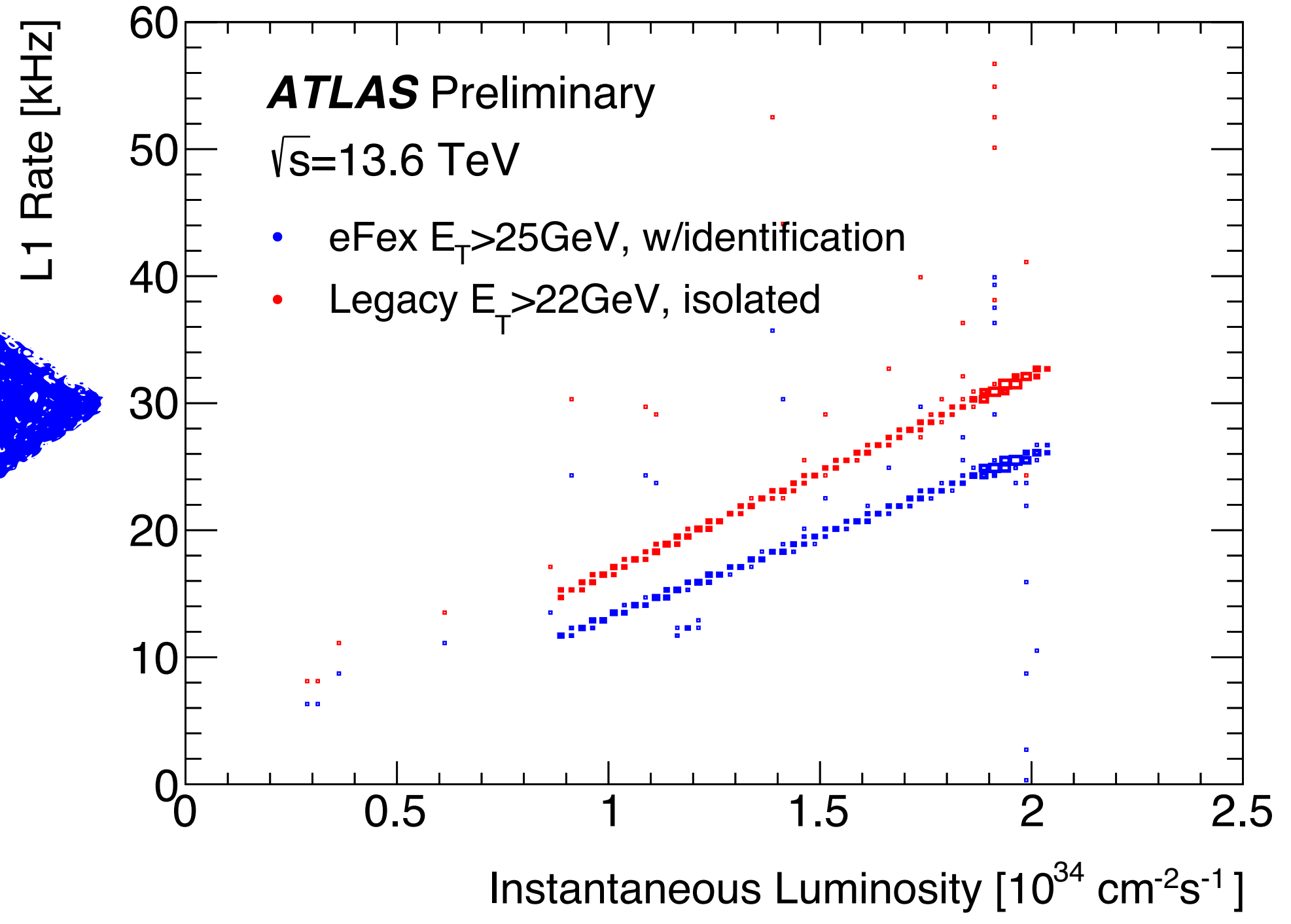
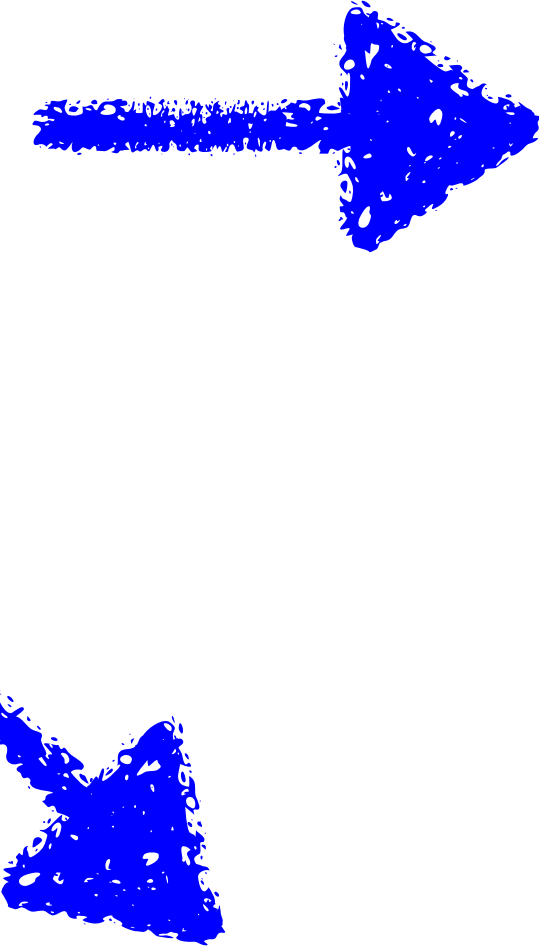
As shown at the beginning of these lectures
both experiments are improving their trigger
capabilities in Run3 but even more at the
HL_LHC



Run3 triggers, reduce rates and increase efficiencies

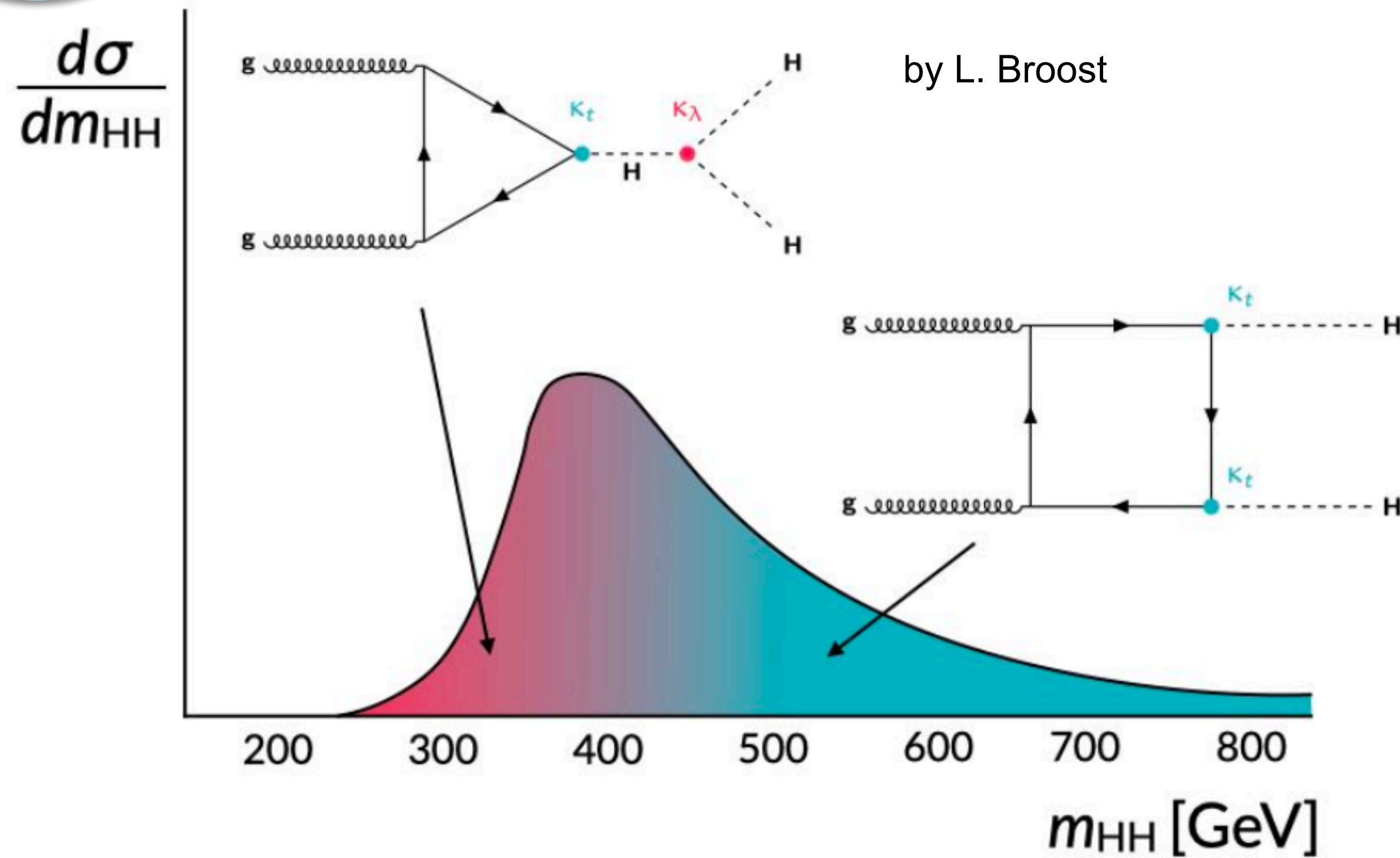
Level 1 Calo single electron rates are decreased and trigger efficiencies increased

NSW and Tile calorimeter coincidences decrease significantly the muon rate



HH

Prospects



In addition to new flavor taggers, new channels etc,
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As shown at the beginning of these lectures
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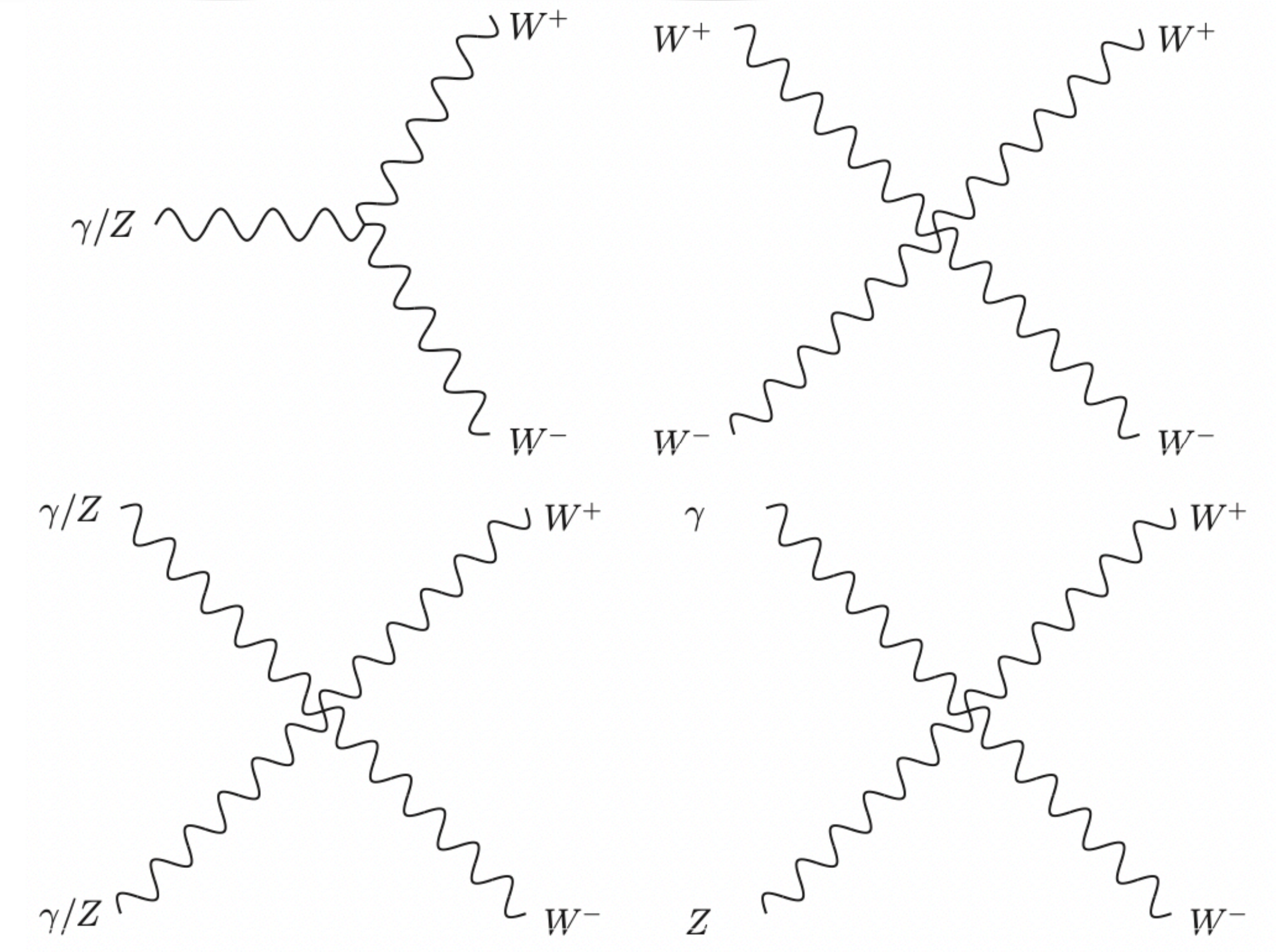
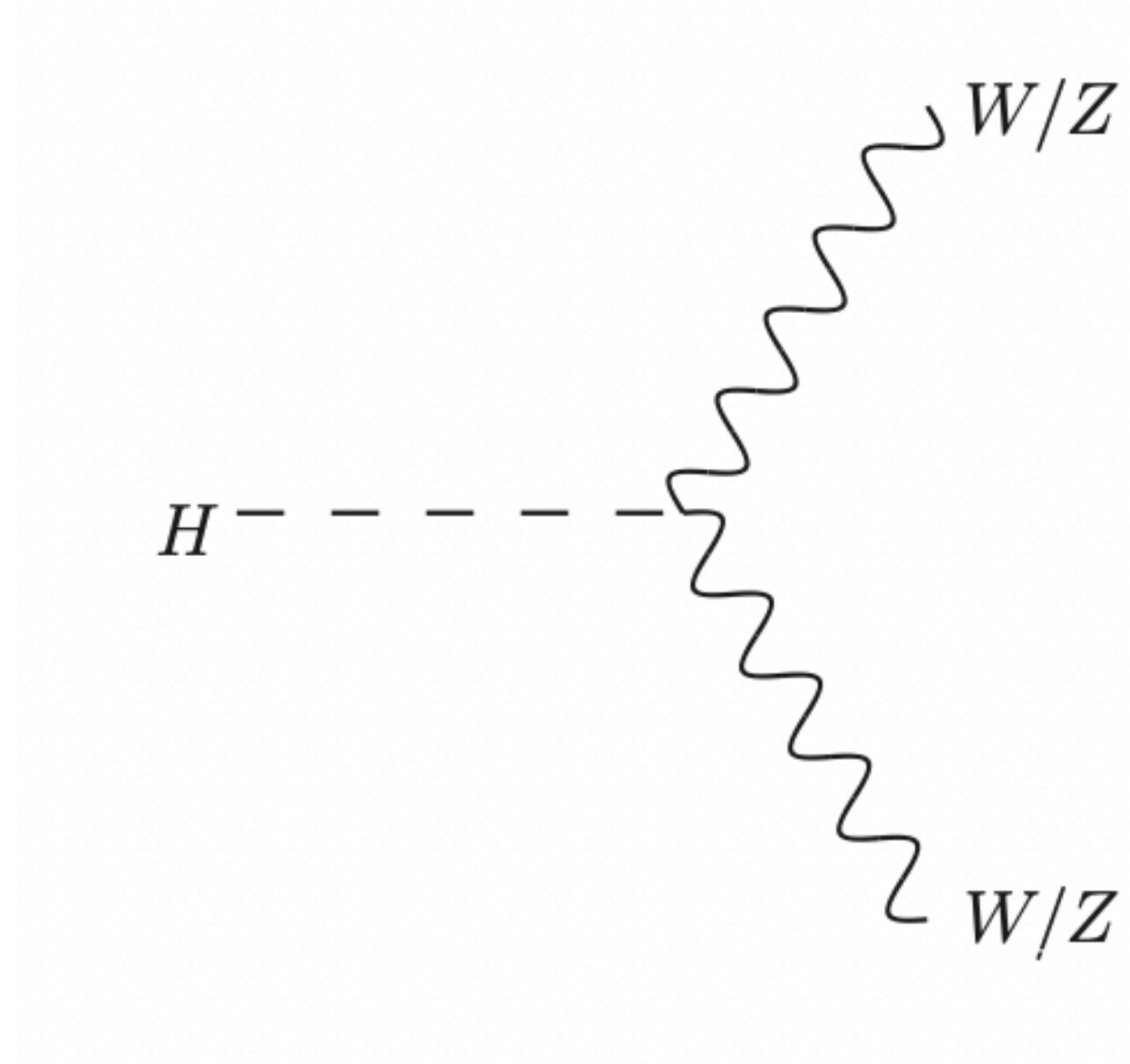
HL-LHC Lumi → Expected sensitivity ↓	HH		VH(→ cc)		VBS long. polarised	
	3 ab ⁻¹	2.5 ab ⁻¹	3 ab ⁻¹	2.5 ab ⁻¹	3 ab ⁻¹	2.5 ab ⁻¹
ATLAS	3.4	3.2	-	-	3.0	2.7
CMS	3.7	3.4	1.3	1.2	2.7	2.5
Combined	5.0	4.7	1.9	1.7	4.0	3.7

Back of the envelope
calculation (no official source)

Testing the Electroweak symmetry breaking via VectorBoson Scattering: Another approach

Electroweak symmetry breaking

Self interactions of the Gauge bosons are predicted by the SM precisely

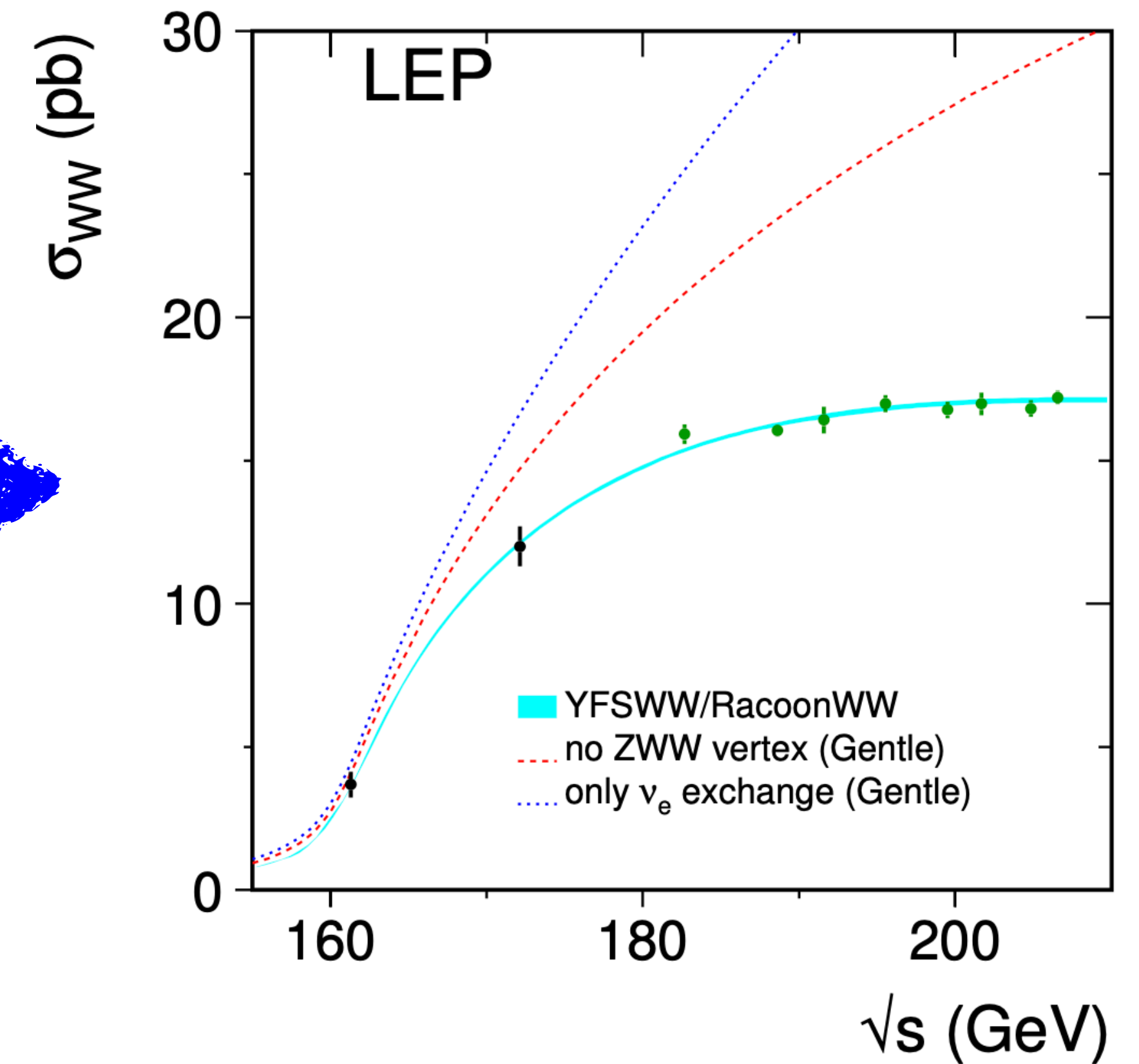
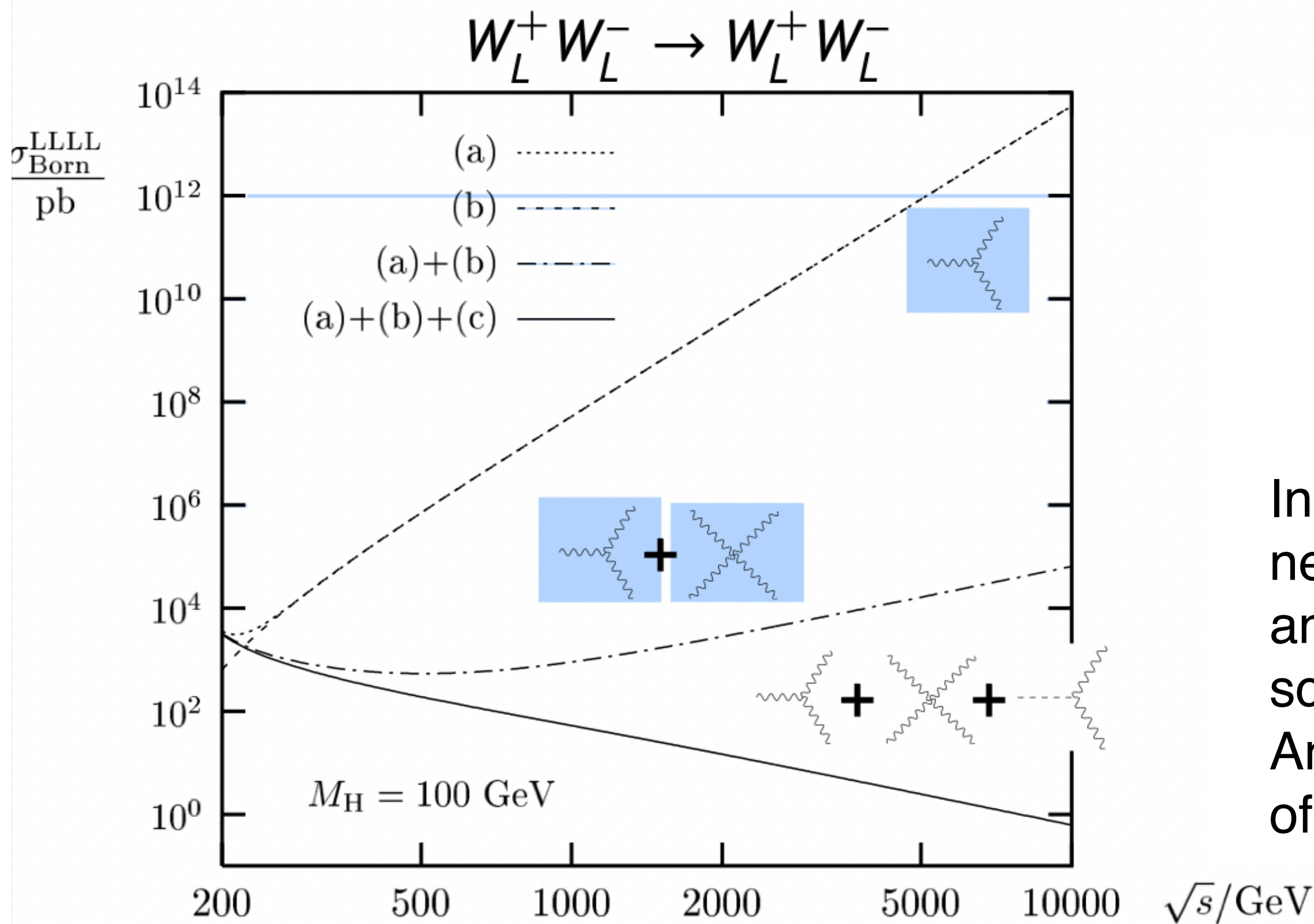
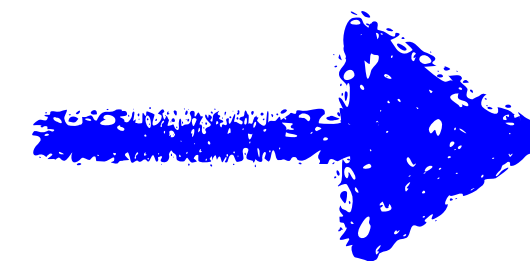


They interact even with the Higgs boson.

Electroweak symmetry breaking

Gauge-boson self interactions play a crucial role for the renormalisability of the electroweak theory

Large cancellations of divergences arising in individual diagrams are exact if couplings take the values of the SM



In vector boson scattering, the presence of the Higgs boson is needed to exactly cancel out the otherwise diverging scattering amplitudes at high energies and prevent unitarity violation at the TeV scale.

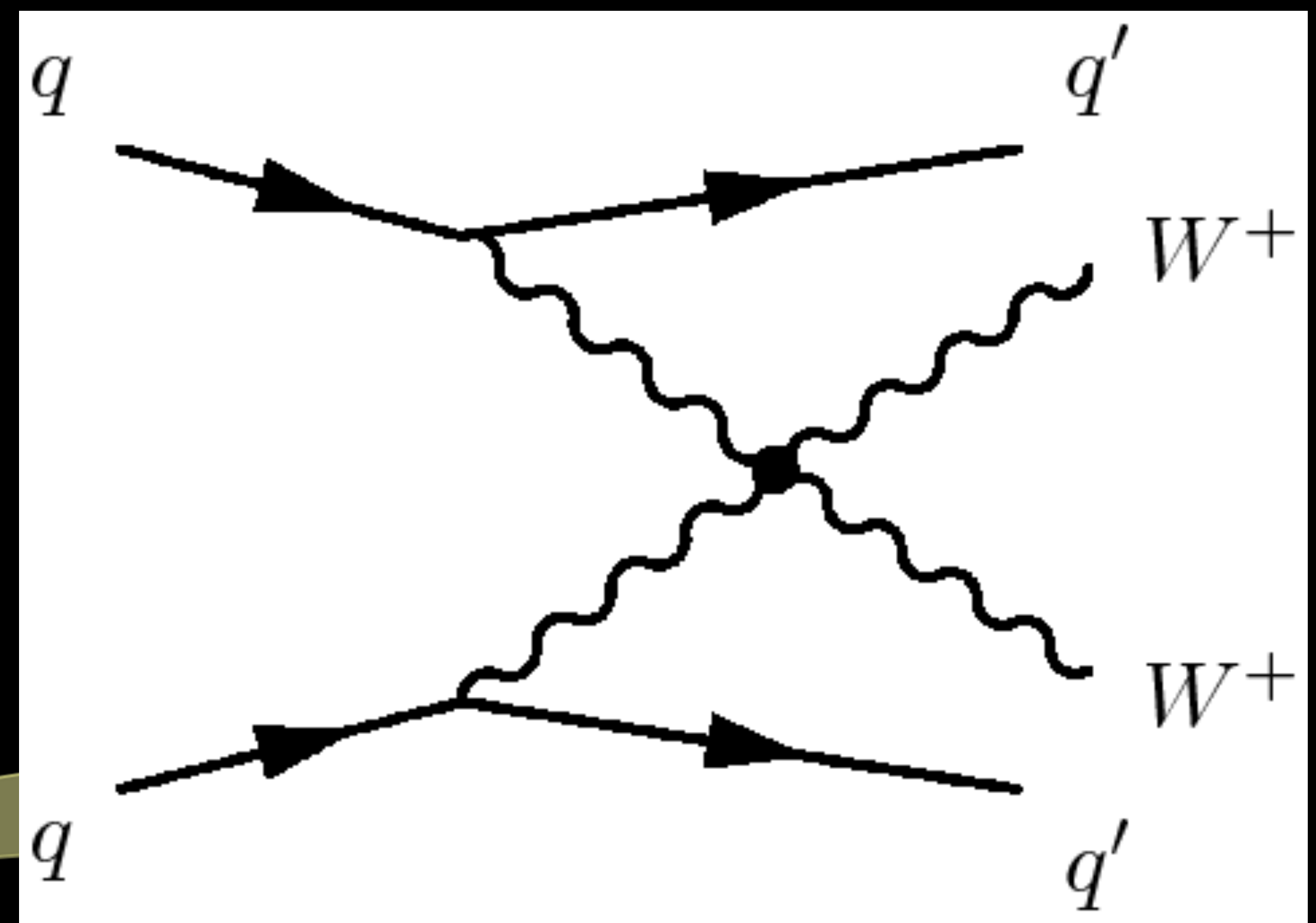
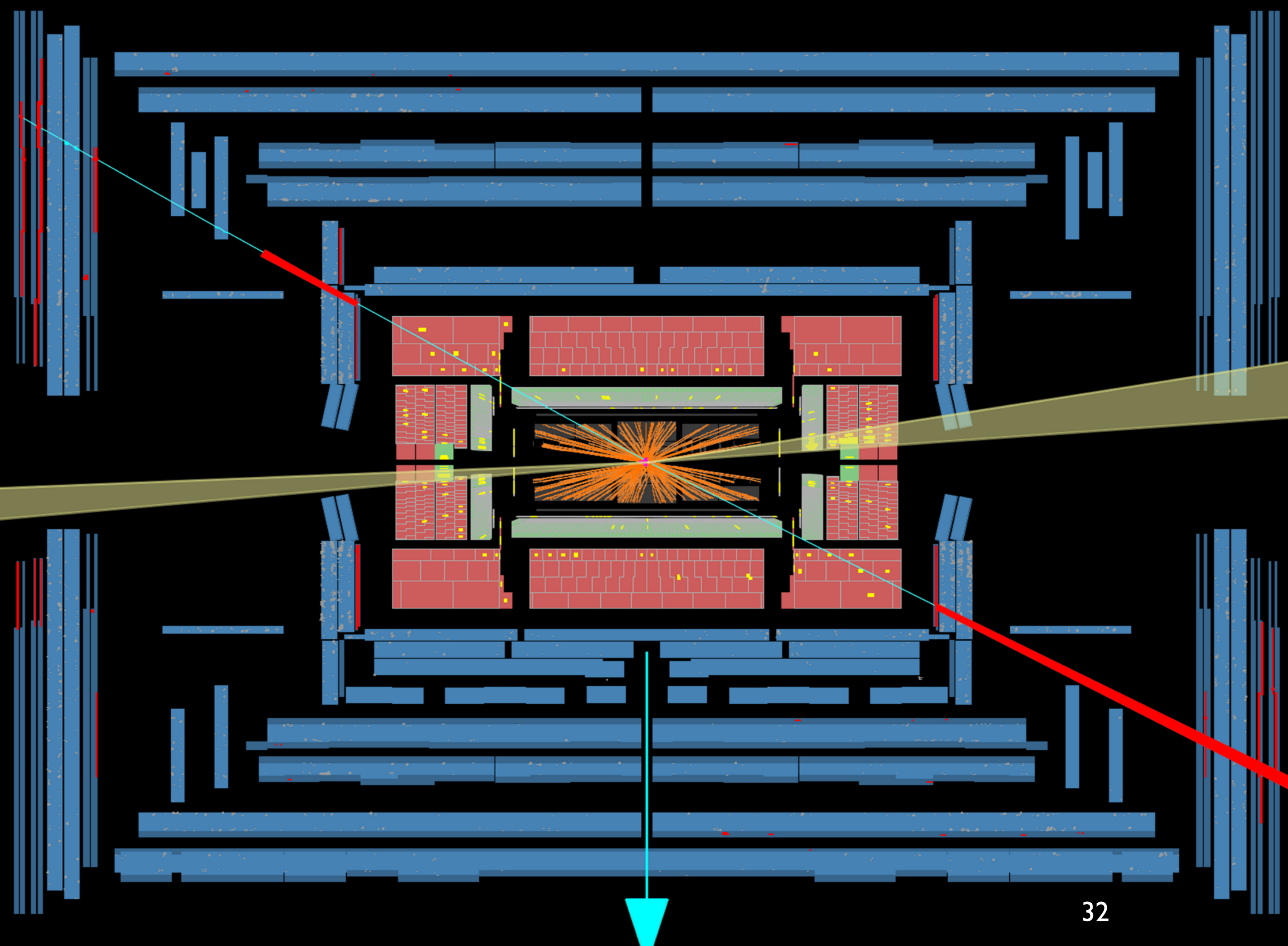
Any significant deviation from the predicted high-energy behaviour of vector boson scattering would point to new phenomena.

Vector boson scattering: Probing EW symmetry

$\mu^+\mu^+jj$ Candidate Event

$m_{jj} = 2800$ GeV

$|\Delta y_{jj}| = 6.3$



ATLAS EXPERIMENT

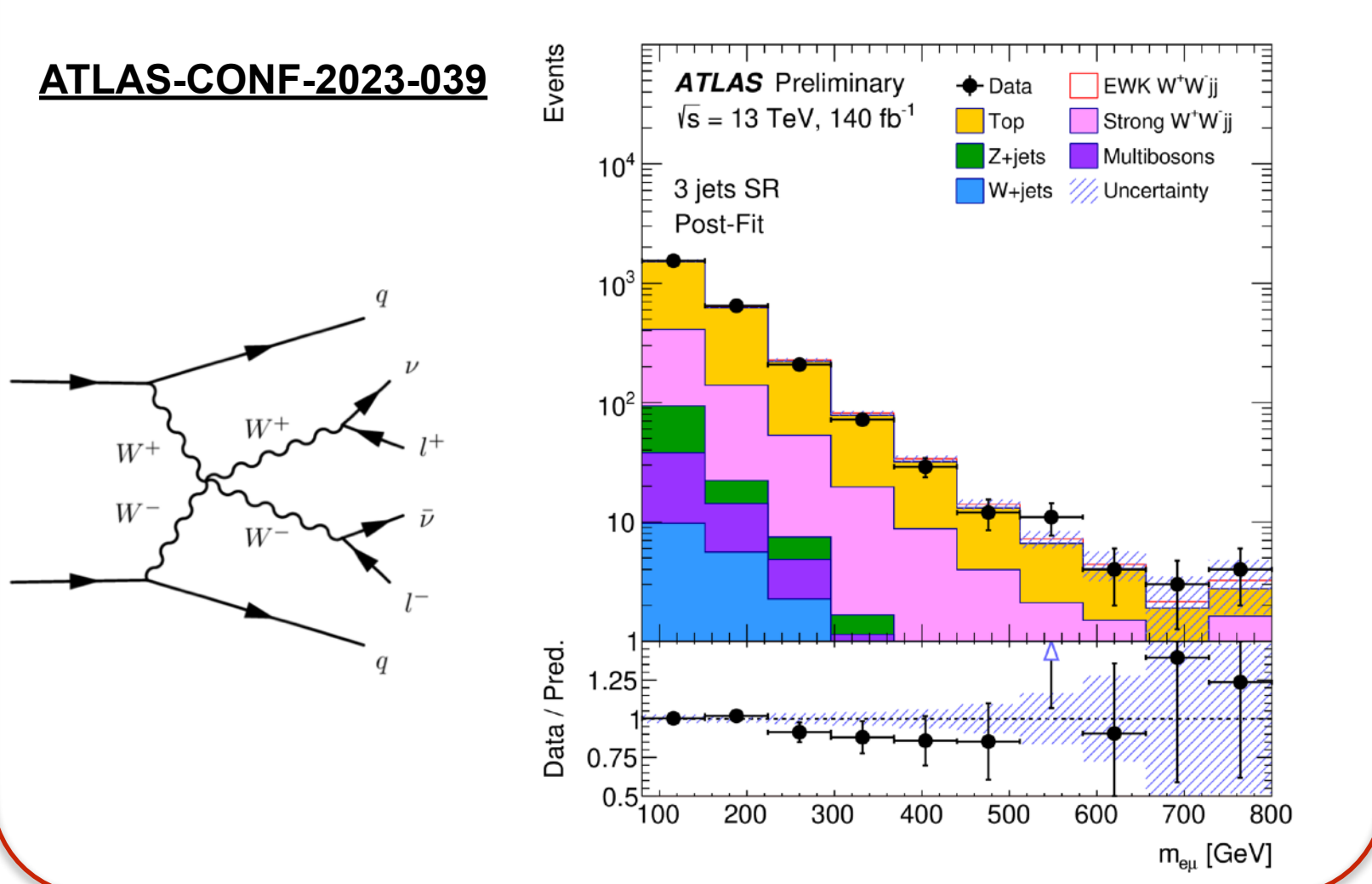
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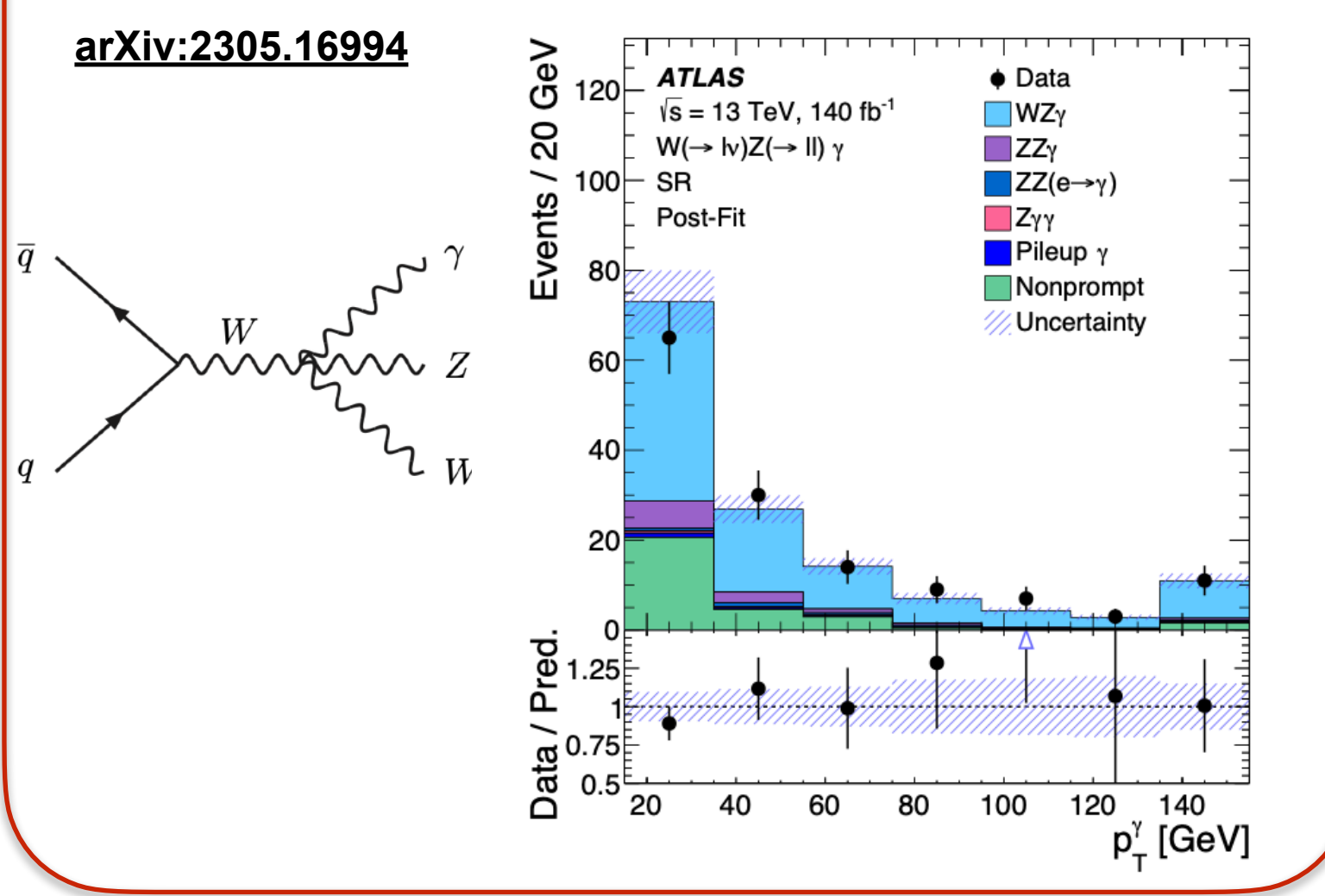
VBS+triboson to study EWK symmetry breaking

ATLAS has a broad research programme to study VBS, recent key results below

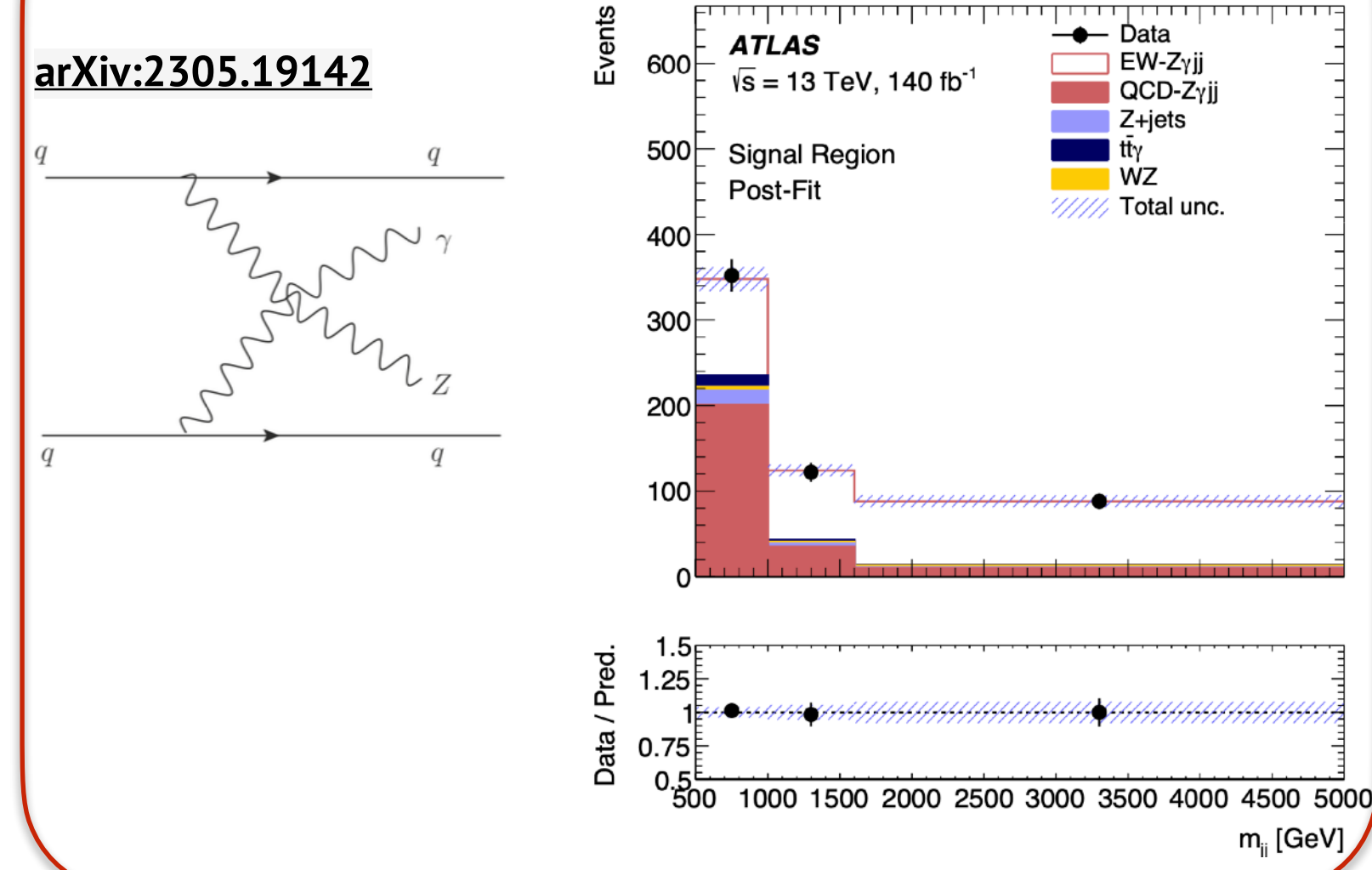
W+W-jj ATLAS observation at 7.1σ (6.2σ exp)



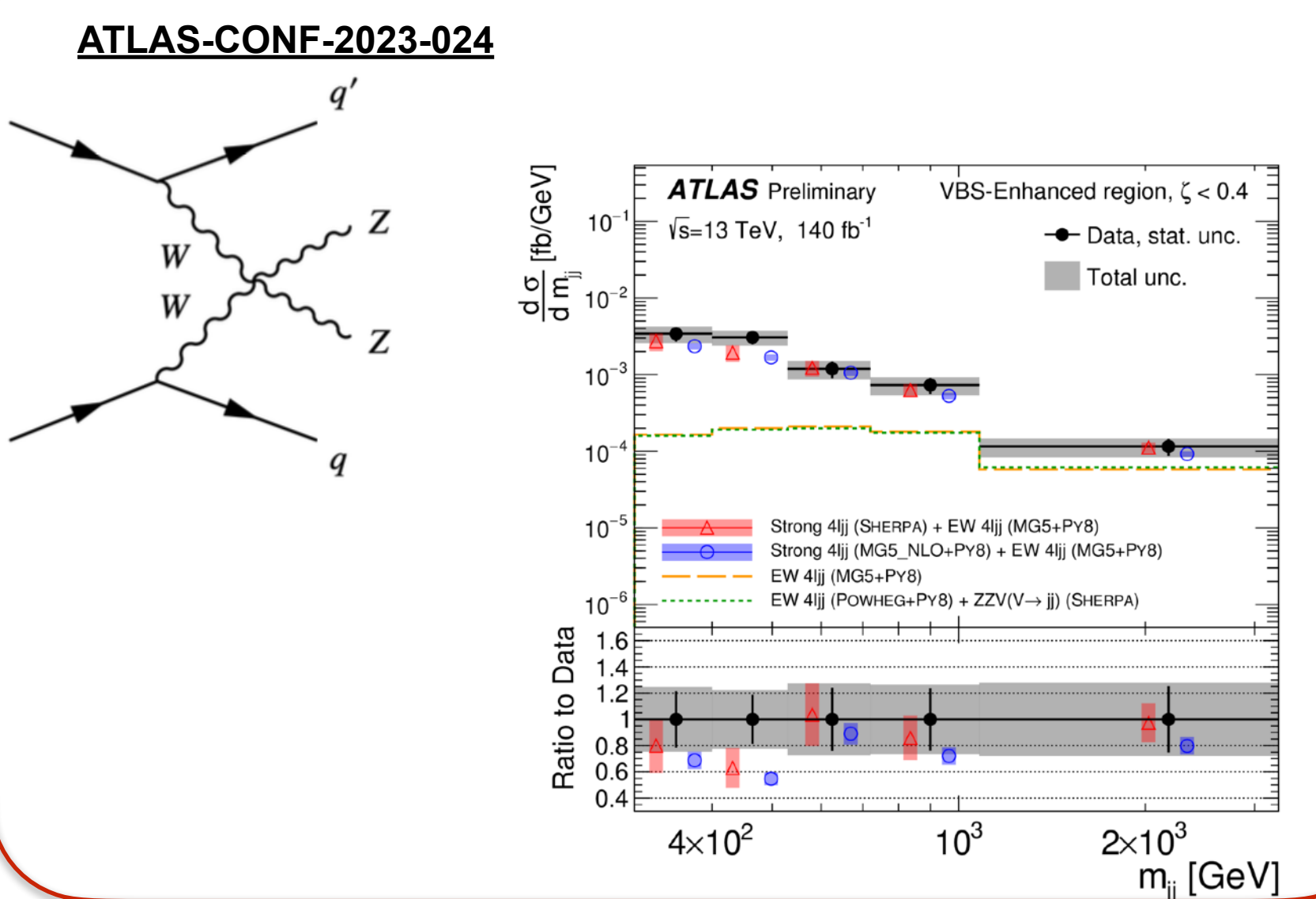
WZγ observation 6.3σ (5.0σ exp)



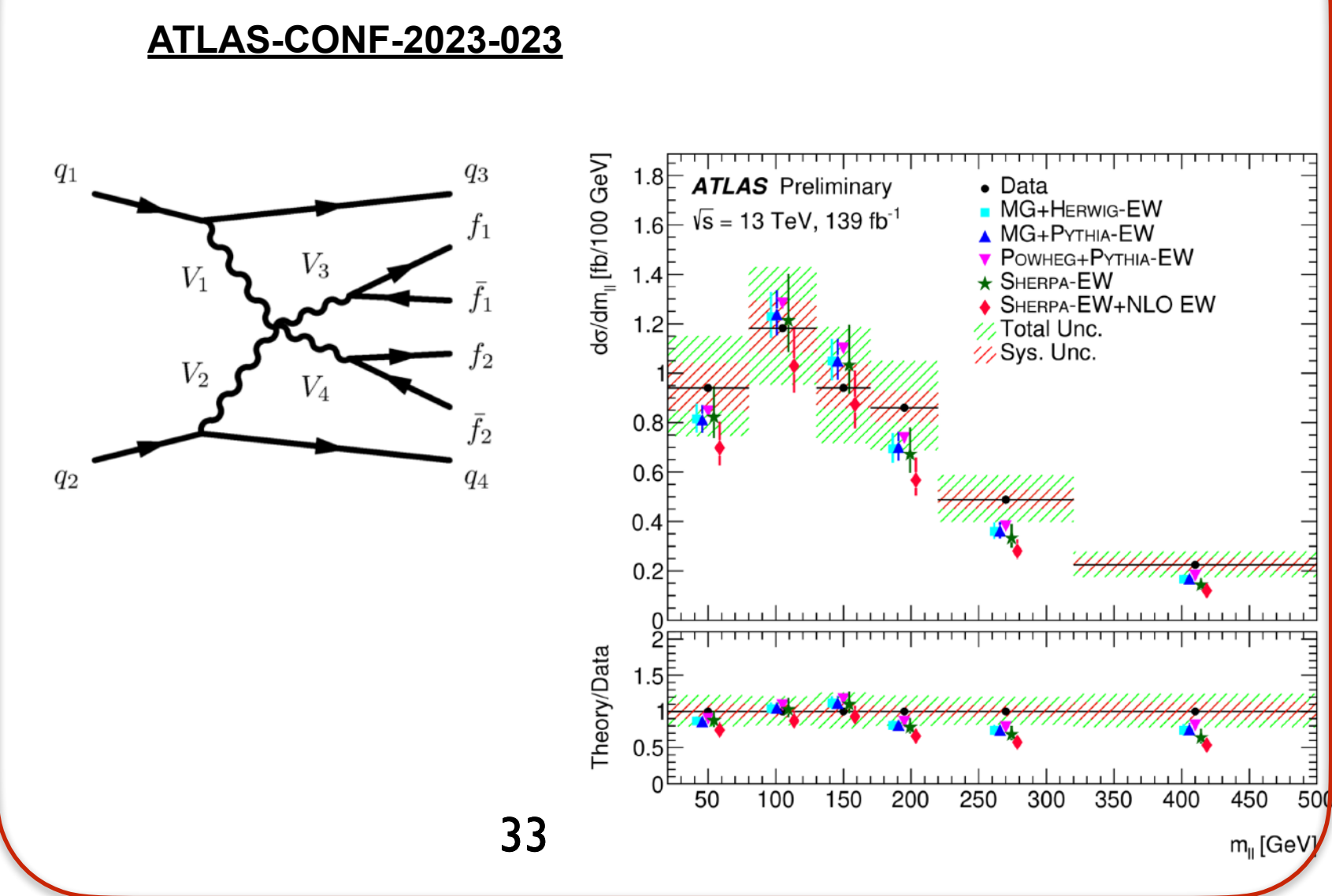
ATLAS observation of EW Zyjj



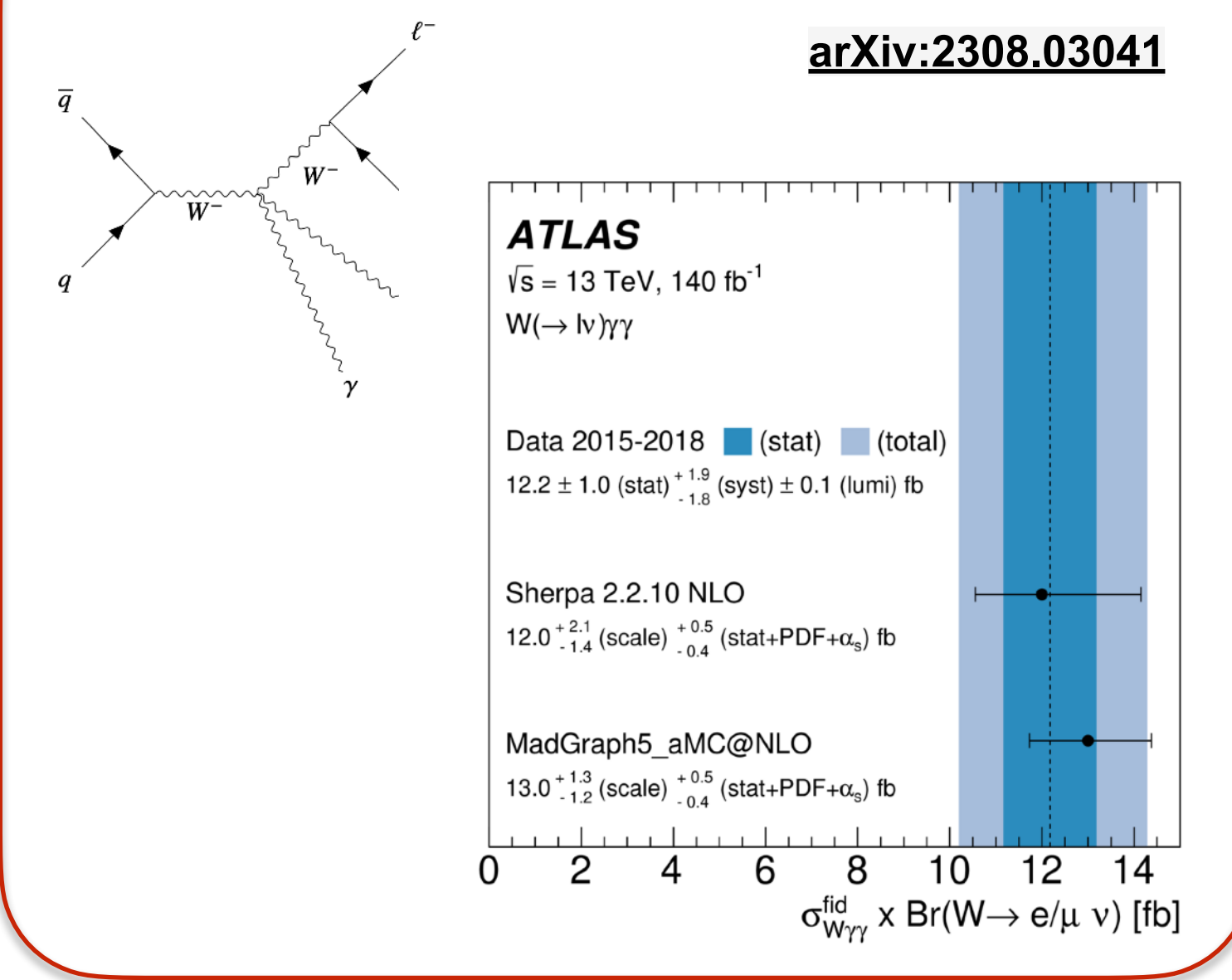
ZZjj differential distributions



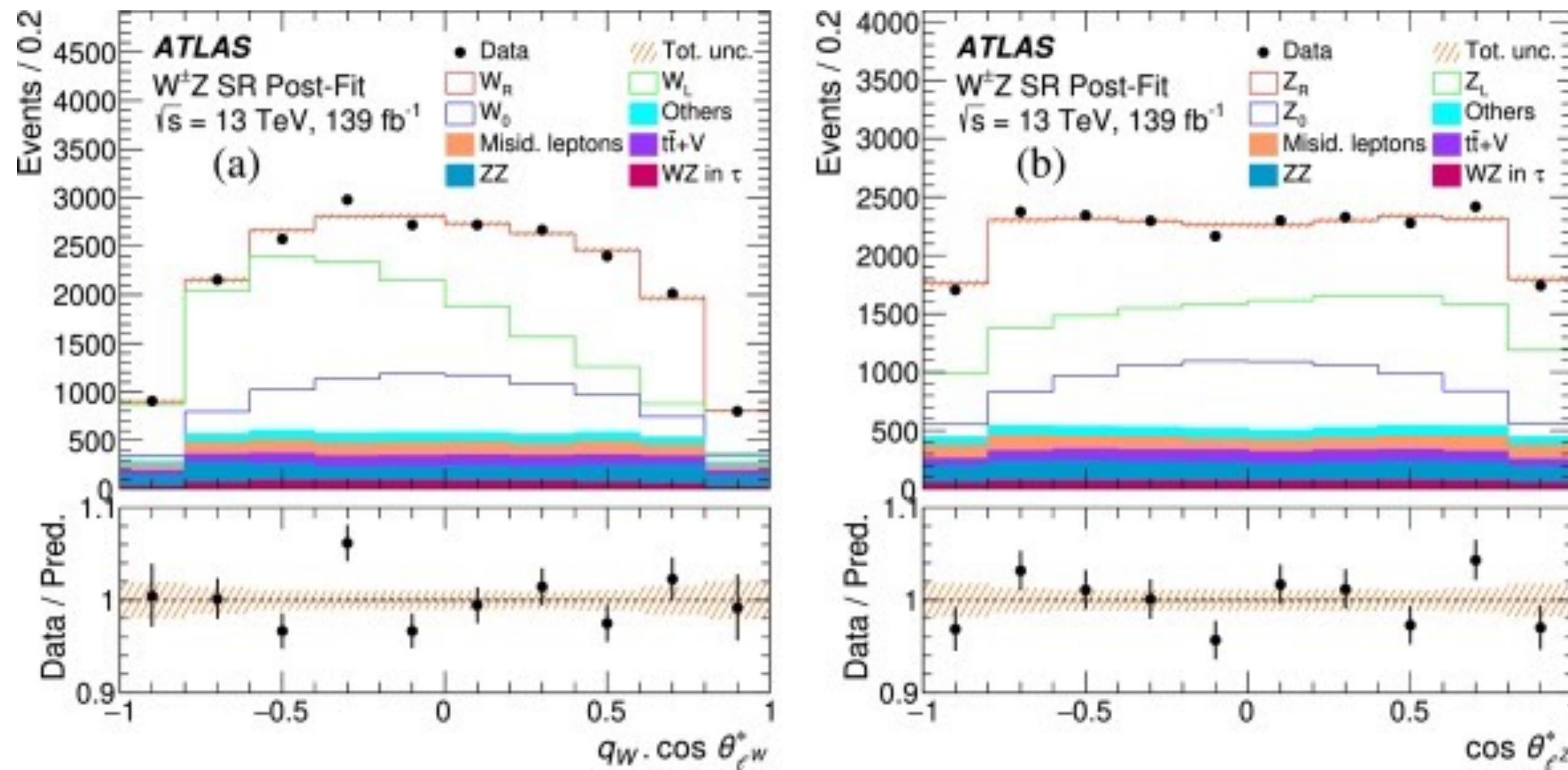
differential VBS Same Sign WWjj



Triboson Wyy observation 5.6σ (5.6σ exp)



W±Z and WW polarization



- longitudinally polarised state of weak bosons is a consequence of the non-vanishing mass of the bosons generated by the electroweak symmetry breaking mechanism.
- these measurements give insights to the way the symmetry is spontaneously broken

ATLAS observed for the first time the production of di-boson polarisation in the $W^\pm Z$ final state.

The most sensitive channel to probe for anomalies is the scattering of two longitudinally W bosons.

While the cross section of same-sign WW production was observed for the first time using Run 2 data,

it is one of the goals of the Run 3 program to measure the polarization

HL-LHC Lumi→ Expected sensitivity↓	VBS long. polarised	
	3 ab ⁻¹	2.5 ab ⁻¹
ATLAS	3.0	2.7
CMS	2.7	2.5
Combined	4.0	3.7

Probing new physics with precision measurements


- The absence of definitive signals indicating physics beyond the SM at the LHC suggests the possibility of a scale separation between the SM and any potential new physics at higher energies. This motivates the utilization of the **Standard Model Effective Field Theory** (SMEFT) as a valuable tool for indirectly searching for new physics in LHC data
- Effective Field theories introduce new-physics states at a high mass scale Λ , significantly larger than the electroweak scale. By expanding in terms of E/Λ , where E represents the typical energy exchanged in a process, the theory provides predictions for experimental observables. This expansion is achieved through a series of operators, which are constructed as gauge-invariant combinations of SM fields with energy dimensions greater than four.

Testing the deviations from the SM via precision measurements

EFT/BSM study

- EFT and BSM interpretation using 10 year's Higgs anniversary [Nature 607 \(2022\) 52](#) publication. In addition differential x-sec combined results for $H \rightarrow \gamma \gamma$ ([JHEP 08 \(2022\) 027](#)) and ZZ decays ([Eur. Phys. J. C 80 \(2020\) 942](#)) are used as well.

Effective Field Theory parametrizes Beyond Standard Model (BSM) effects at high energies ($\Lambda \gg v$, above electroweak scale) at low energies, $E \ll \Lambda$, in terms of higher-dimensional operators in an effective Lagrangian:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i^{N_{d=6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_j^{N_{d=8}} \frac{b_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots,$$


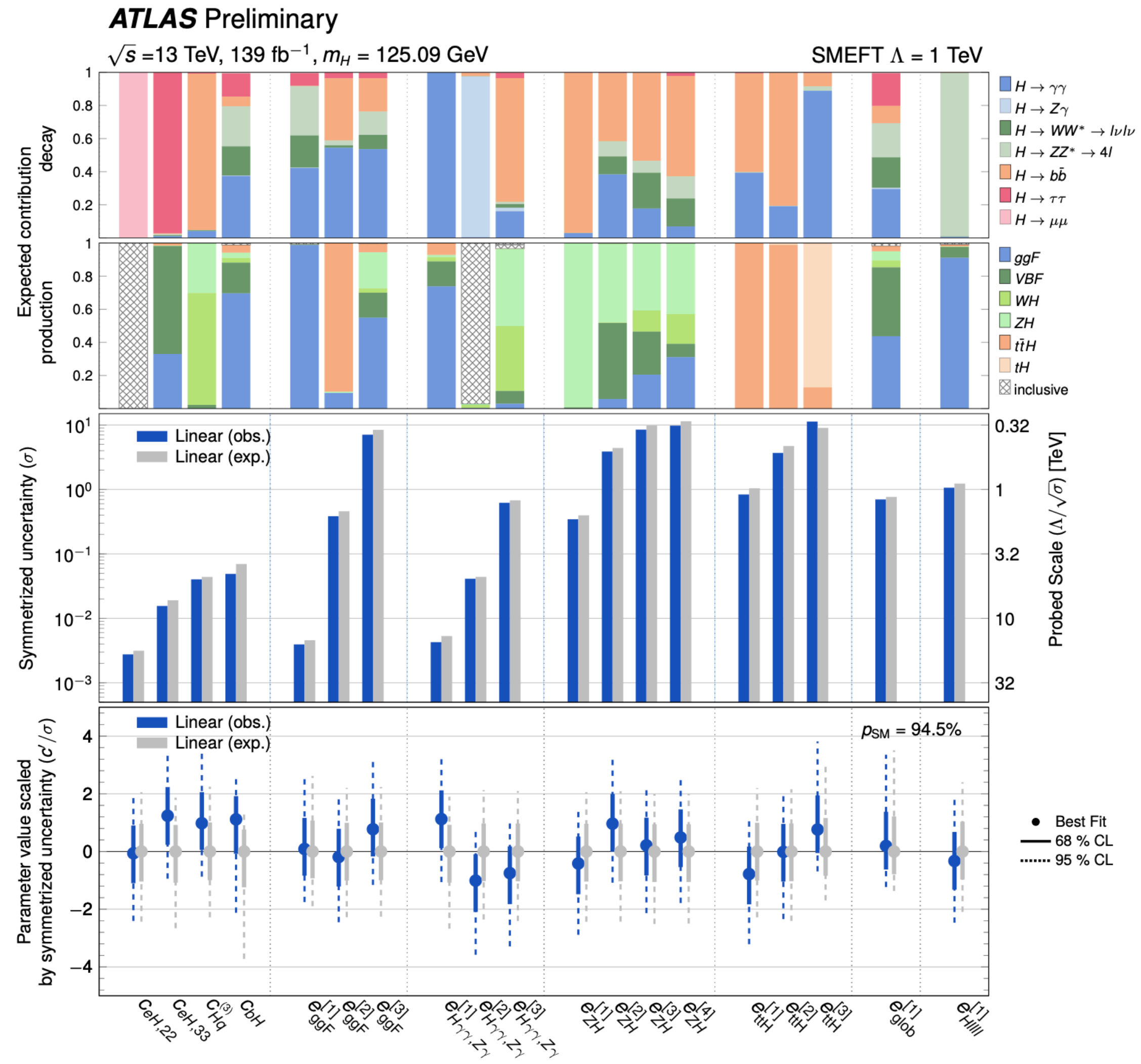
dimensionless Wilson coefficients of higher dimension operators



Comprehensive Higgs EFT/BSM study

- EFT and BSM interpretation using 10 year's Higgs anniversary [Nature 607 \(2022\) 52](#) publication. In addition uses differential x-sec combined results for $H \rightarrow \gamma\gamma$ ([JHEP 08 \(2022\) 027](#)) and ZZ decays ([Eur. Phys. J. C 80 \(2020\) 942](#)).

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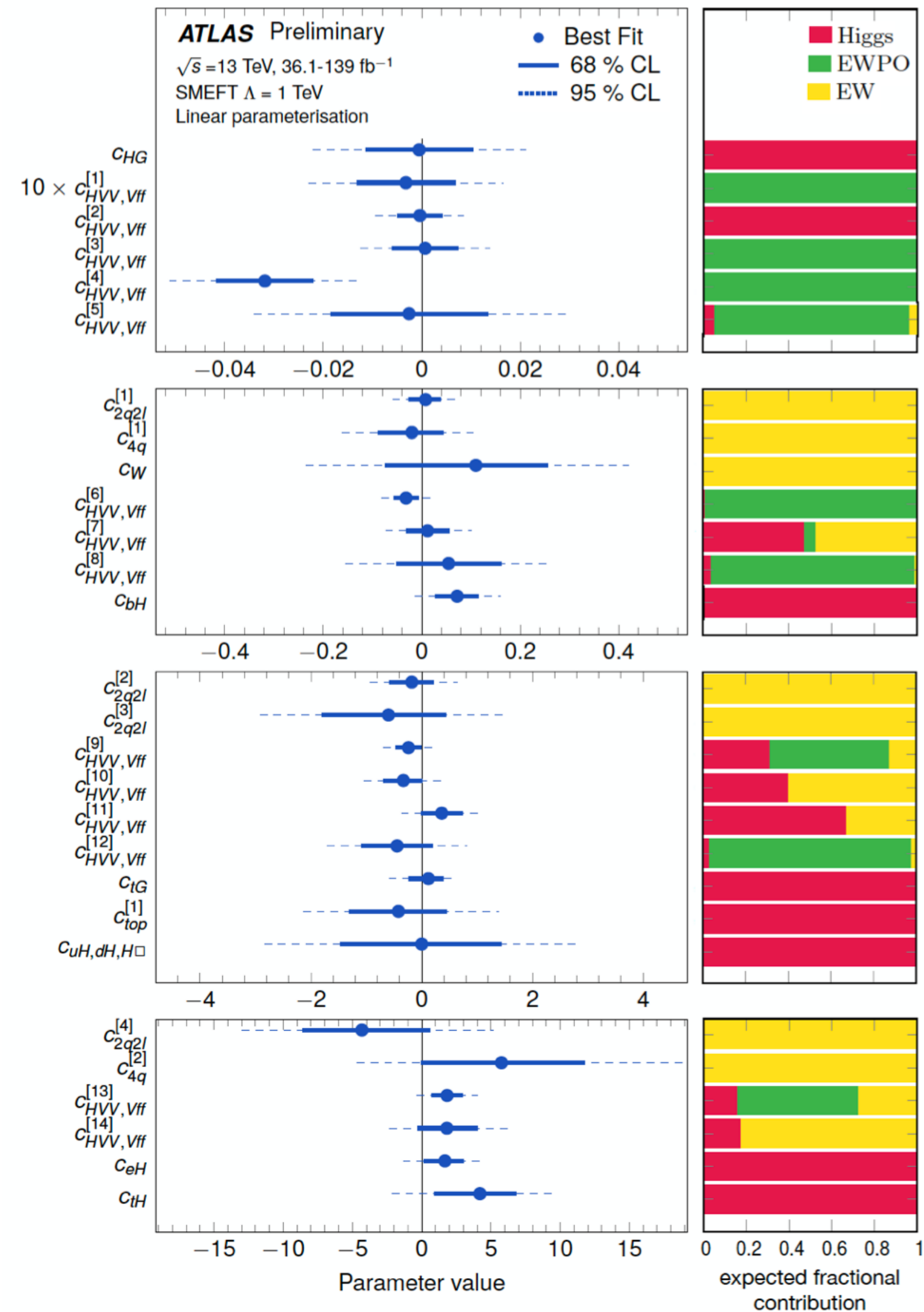


EFT combinations of SM and Higgs channels

Decay channel	Target Production Modes	\mathcal{L} [fb ⁻¹]	Ref.
$H \rightarrow \gamma\gamma$	ggF, VBF, WH , ZH , $t\bar{t}H$, tH	139	[10]
$H \rightarrow ZZ^*$	ggF, VBF, WH , ZH , $t\bar{t}H(4\ell)$	139	[11]
$H \rightarrow WW^*$	ggF, VBF	139	[12]
$H \rightarrow \tau\tau$	ggF, VBF, WH , ZH , $t\bar{t}H(\tau_{\text{had}}\tau_{\text{had}})$	139	[13]
	WH, ZH	139	[14–16]
$H \rightarrow b\bar{b}$	VBF	126	[17]
	$t\bar{t}H$	139	[18]

We can add also SM measurements to constrain the same operators
 Planning to add also top quark

Process	Important phase space requirements	Observable	\mathcal{L} [fb ⁻¹]	Ref.
$pp \rightarrow e^\pm \nu \mu^\mp \nu$	$m_{\ell\ell} > 55 \text{ GeV}$, $p_{\text{T}}^{\text{jet}} < 35 \text{ GeV}$	$p_{\text{T}}^{\text{lead. lep.}}$	36	[19]
$pp \rightarrow \ell^\pm \nu \ell^+ \ell^-$	$m_{\ell\ell} \in (81, 101) \text{ GeV}$	m_{T}^{WZ}	36	[20]
$pp \rightarrow \ell^+ \ell^- \ell^+ \ell^-$	$m_{4\ell} > 180 \text{ GeV}$	m_{Z2}	139	[21]
$pp \rightarrow \ell^+ \ell^- jj$	$m_{jj} > 1000 \text{ GeV}$, $m_{\ell\ell} \in (81, 101) \text{ GeV}$	$\Delta\phi_{jj}$	139	[22]



The same wilson coefficients are constrained by different processes!

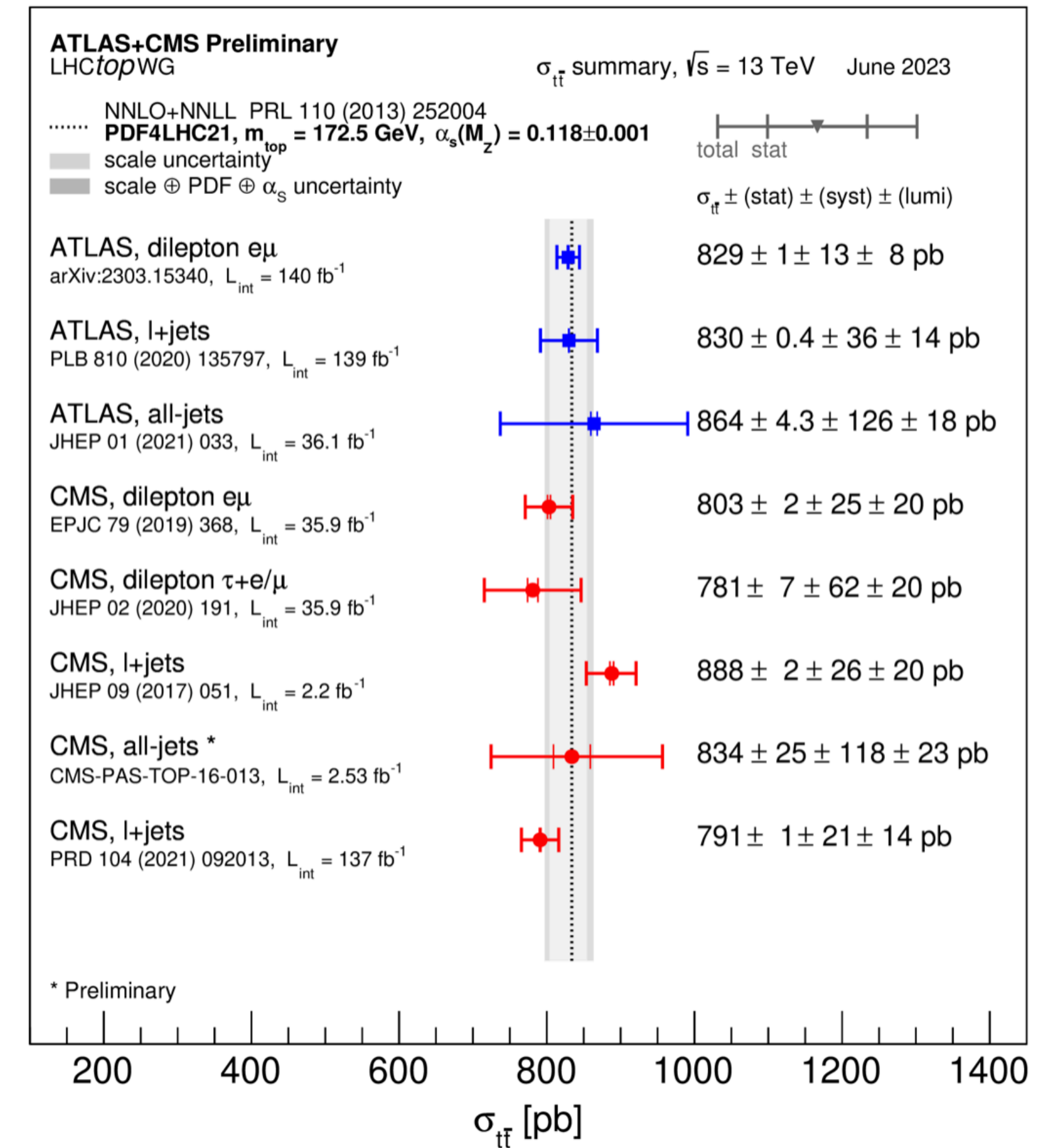
The statistical uncertainties are very large

Run 3 will be pivotal for such EFT interpretations and top+SM+higgs results will be combined,

Some more SM: top quark and more

Top quark pair production cross-section

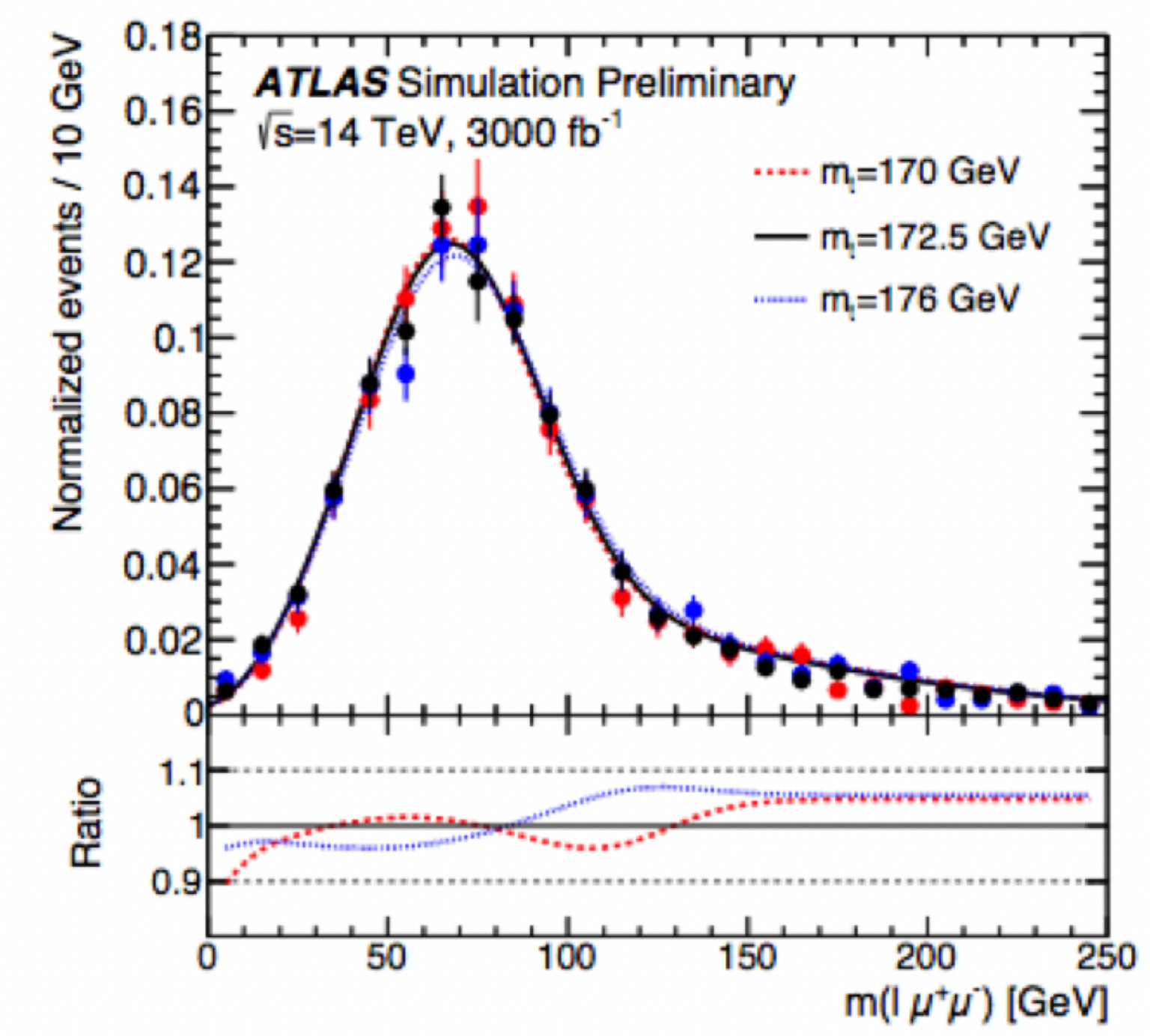
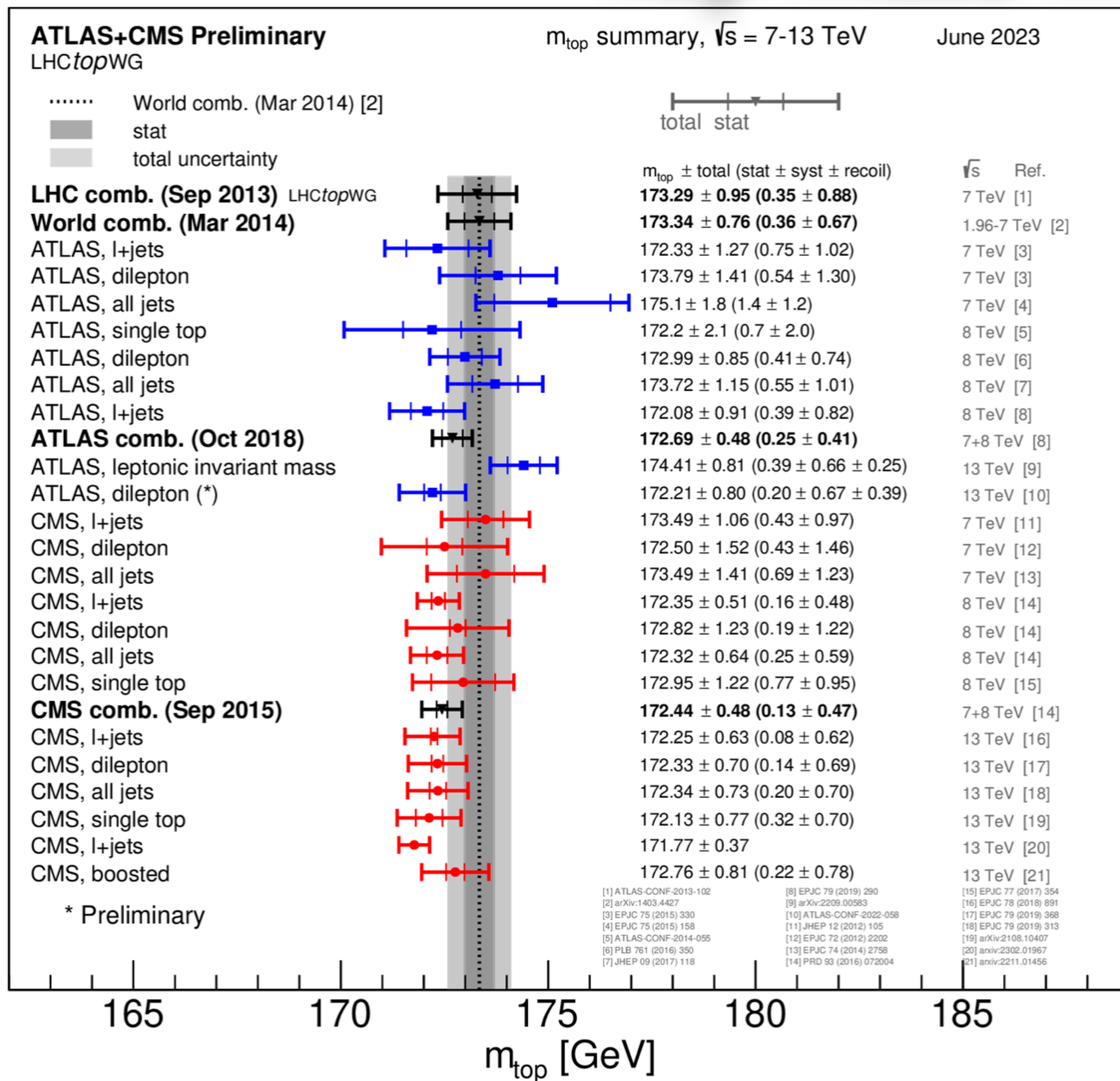
Channel	\sqrt{s} [TeV]	$\int \mathcal{L} dt$ [fb^{-1}]	$\sigma_{t\bar{t}}$ [pb]	Reference
Dilepton, ℓ +jets	5	0.257	67.5 ± 2.7	[7]
$e\mu$	7	4.6	183 ± 7	[8]
ℓ +jets	7	4.7	169 ± 7	[3]
$e\mu$	8	20.2	242 ± 9	[8]
ℓ +jets	8	20.2	248 ± 14	[9]
$e\mu$	13	140	829 ± 15	[4]
ℓ +jets	13	139	830 ± 39	[10]
all-jets	13	36.1	864 ± 127	[11]
$e\mu$	13.6	11	859 ± 29	[2]



For top physics the theory modelling uncertainties are more important than more data in Run 3!

Generally for the top cross-section, top mass and top+X processes the top modelling uncertainties are playing a very important role, even for analyses that are statistically dominated.

Top mass measurements



The analysis will profit from high luminosity at HL-LHC and larger acceptance of the inner detector. It will be limited by jet energy uncertainty scale on m_{top} but tt modelling will also be relevant.

EW fit constrains m_W , m_{top} and m_H

The global electroweak fit enabled prediction of m_{top} and m_H before their discoveries:

- Measure different observables
- Calculate relations between observables

$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2} \right) = \frac{\pi\alpha}{\sqrt{2}G_\mu} (1 + \Delta r)$$

The W boson mass in the SM is related with the Z-boson mass, m_Z , the fine structure constant, α , and the Fermi constant, G_μ

Δr includes the quantum corrections to m_W , which depend m_{top} quadratically and m_H , logarithmically.

One can indirectly constrain these parameters with great precision.

By the end of the LHC, we might have results in indirect precisions of $\Delta m_W \approx 4$ MeV, $\Delta m_{Top} \approx 1.3$ GeV, $\Delta m_H \approx 13$ GeV

The EW fits generically impose stringent constraints on any theory of electroweak symmetry breaking

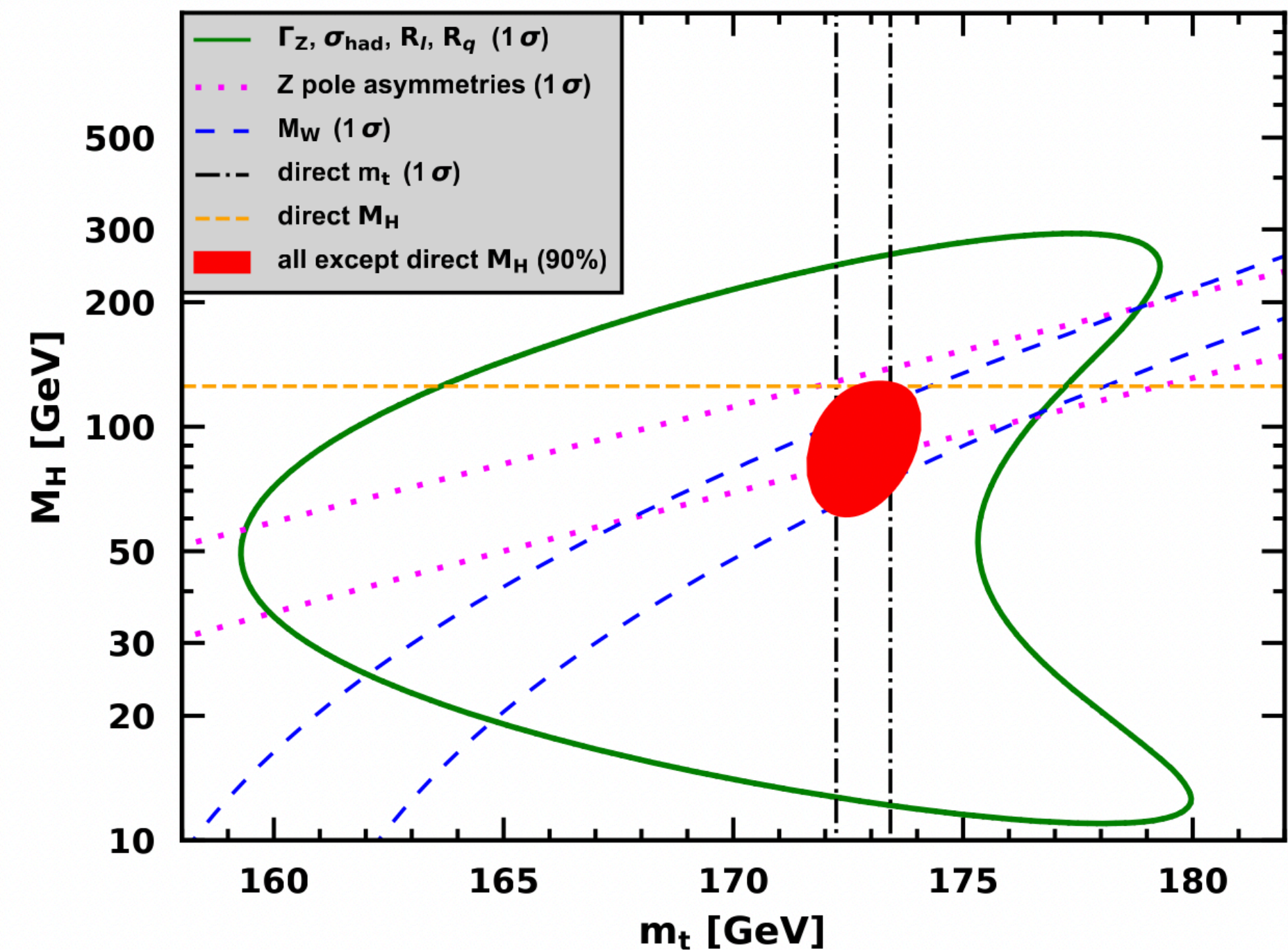
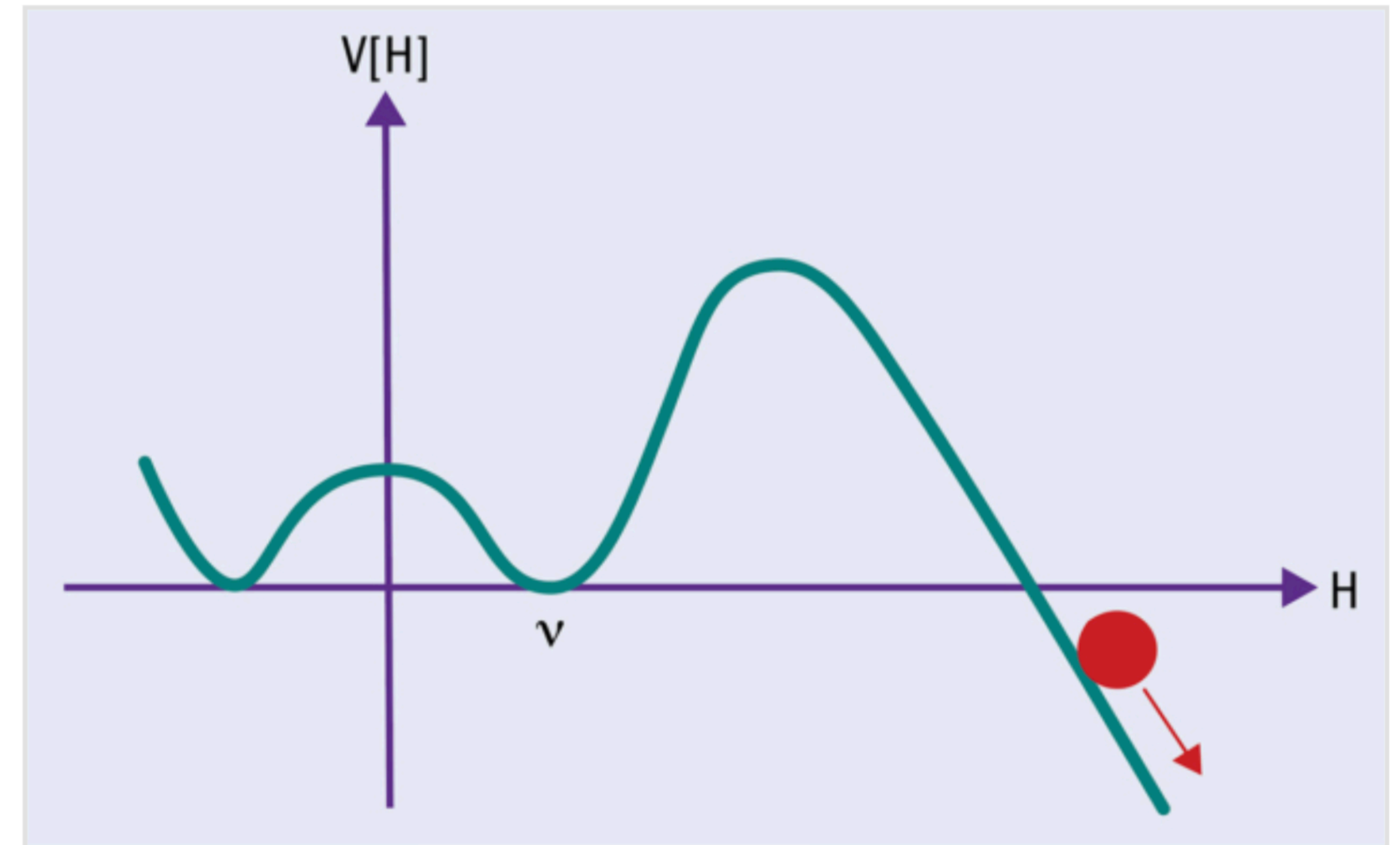
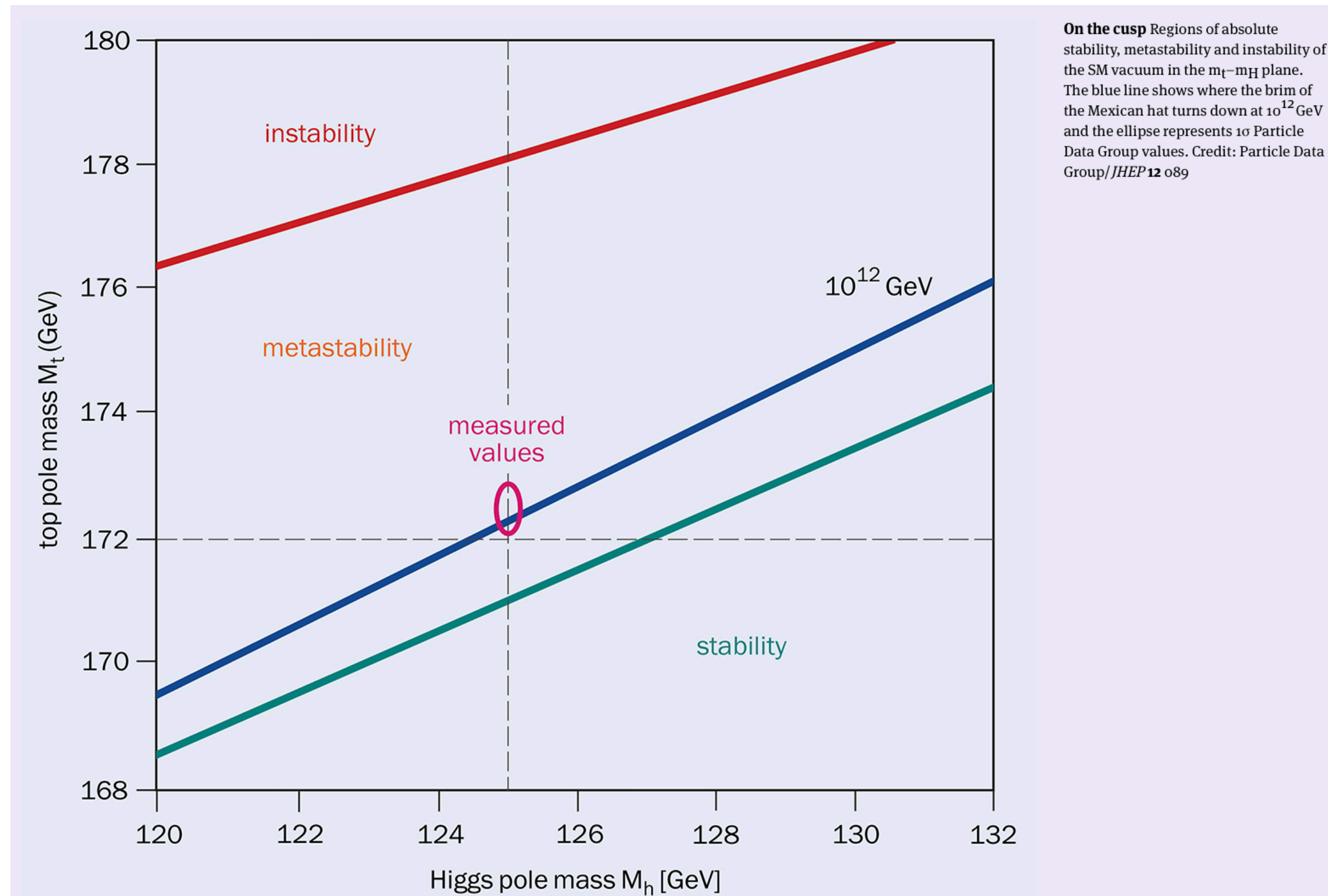


Figure 10.4: Fit result and one-standard-deviation (39.35% for the closed contours and 68% for the others) uncertainties in M_H as a function of m_t for various inputs, and the 90% CL region ($\Delta\chi^2 = 4.605$) allowed by all data. $\alpha_s(M_Z) = 0.1185$ is assumed except for the fits including the Z lineshape. The width of the horizontal dashed band is not visible on the scale of the plot.

m_{top} and m_H

Quantum effects can change the shape of the Mexican hat Higgs potential. The Higgs field has self-interactions that make the hat turn upwards, additional quantum effects can turn it downwards, due to interactions with the fundamental particles to which the Higgs gives mass. The top mass is the heaviest and therefore the most important.

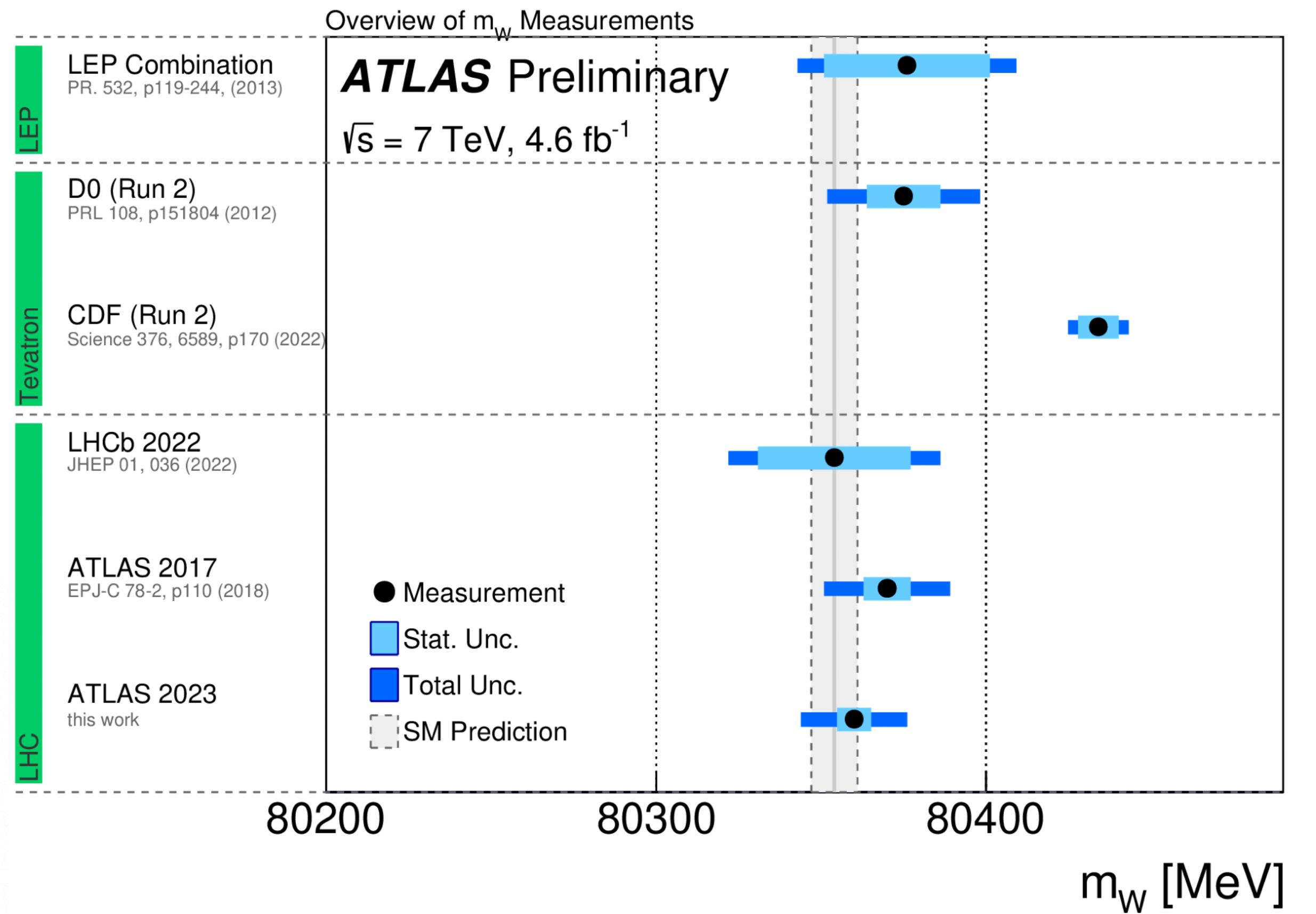
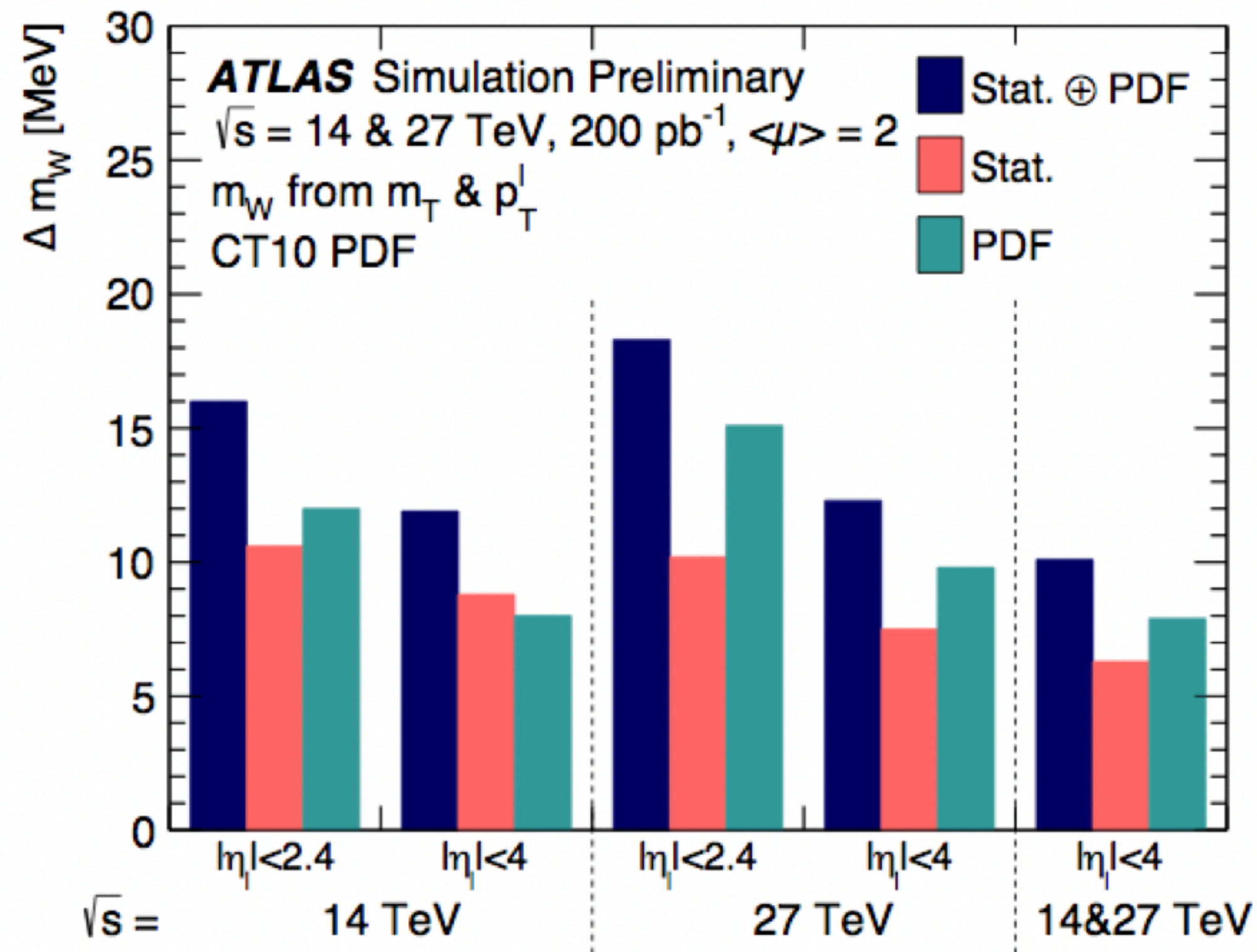


The present measurements indicate that the current minimum of the Higgs potential is not the lowest and that universe could be metastable and that it could end up in the different minimum. New physics could stabilize the vacuum.

W mass at HL_LHC

at HL-LHC a total of 2×10^6 will be produced at the HL-LHC in 1 week

Understanding of PDFs will be crucial



- Re-analysis: new log-likelihood fit constrains systematic uncertainties with data, more modern PDFs!

- **Reduces systematic uncertainties from 18 to 15 MeV**

$$m_W = 80360 \pm 5 \text{ (stat.)} \pm 15 \text{ (syst.)} = 80360 \pm 16 \text{ MeV}$$

Compatible with SM

Errors at 10 MeV or lower will be achieved

Direct searches for New Physics

Searches: we are exploring in all directions

dark sector

Dark Matter

Invisible decays

**heavy
neutrinos**

Extra dimensions

compositeness

resonances

**contact
interactions**

Long Lived Particles

**Highly ionizing
particles**

Leptoquarks

vector like quarks

Axions

**WIMPS
weakly interacting
massive particles**

Supersymmetry

Strategy for Run 3 searches

- covering wider phase space
- going more model independent
- explore wider range of signatures
- exploit the new triggering features of the new detector
- exploit better reconstruction performance in particular flavor tagging large r-jets
- exploit better tracking capabilities: few examples.

Let's start with the analyses that
will profit of new Triggers

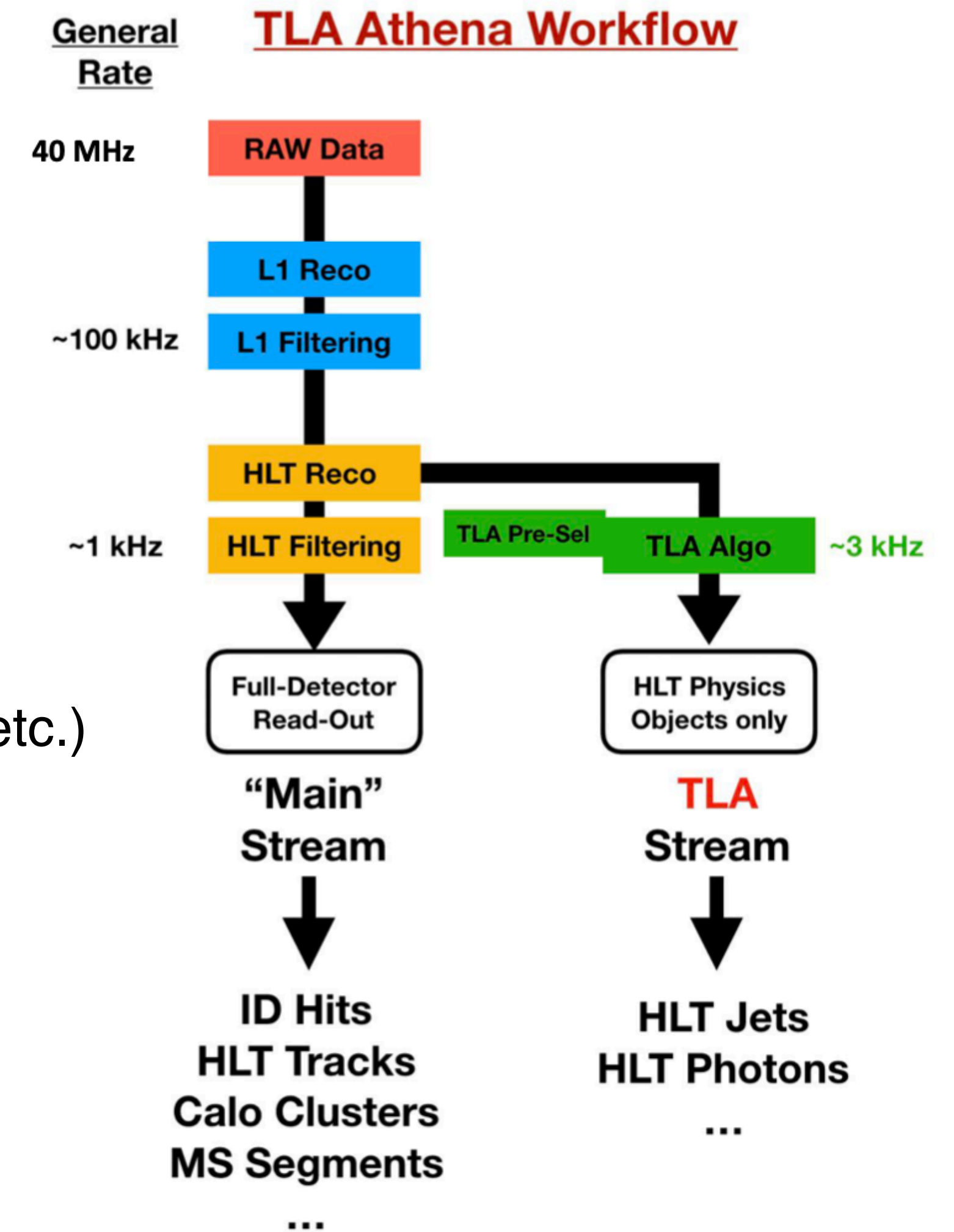
Trigger Level analysis ATLAS

TLA for Inclusive Searches: TLA idea:

- Events only seen by the trigger contain compelling physics:
- Discarded due to trigger thresholds

○ But already reconstructed to perform the trigger decision
 Recover these events → Store trigger reconstruction outcome

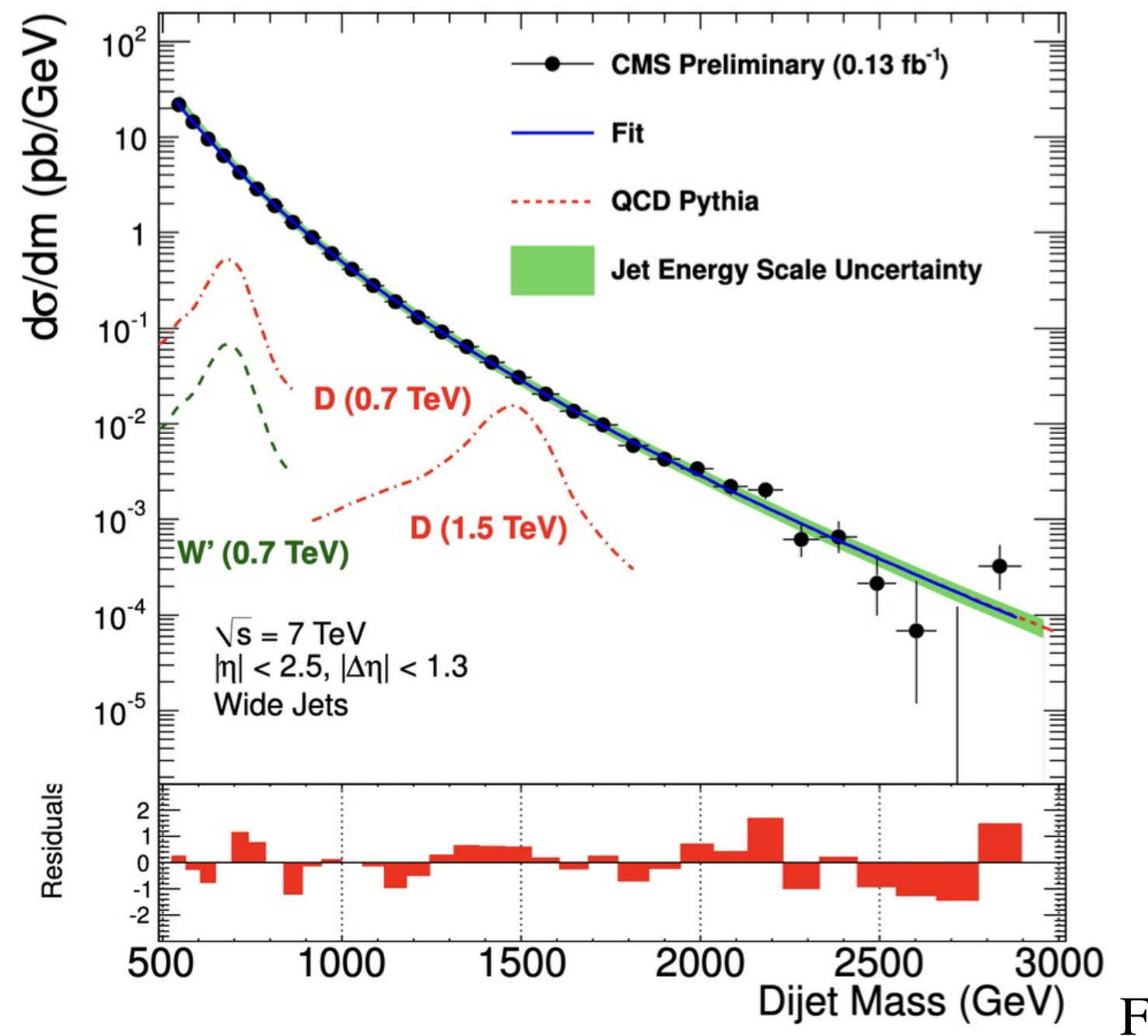
- Run 2/3 baseline TLA:
- SAVE ONLY RESULT OF HLT RECONSTRUCTION (HLT jets, photons, etc.)
- No RAW data stored to output



Bandwidth	=	Rate	x	Event Size	
ATLAS Physics Stream	~ 1.7kHz	x	1.5MB/s	→	3 GB/s
TLA Stream	~ 6kHz	x	4.5kB/s	→	27 MB/s

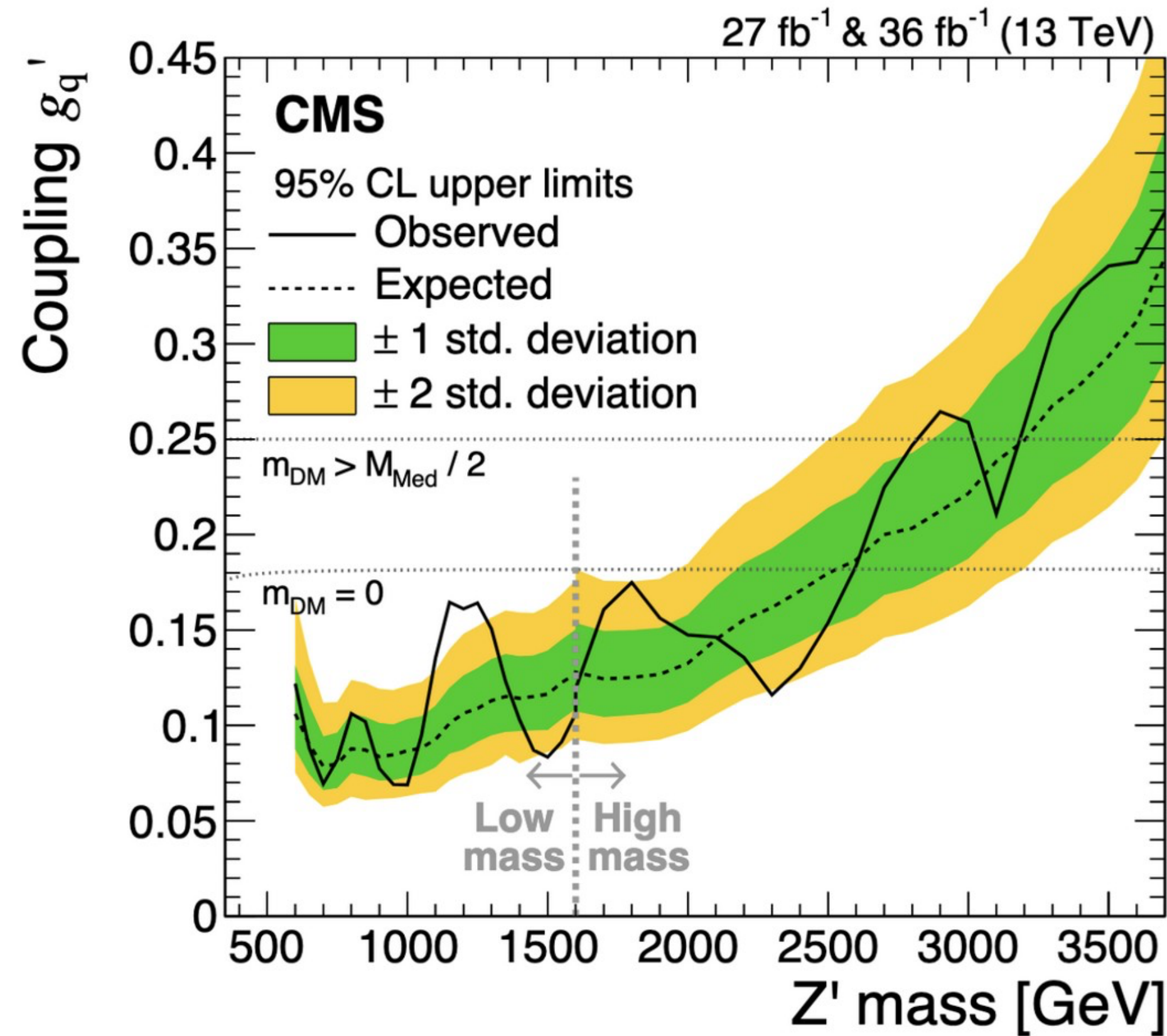
Data-scouting in CMS (same concept)

First employed for di-jet searches by CMS in LHC Run 1



First employed for Dijet searches by CMS in LHC Run 1

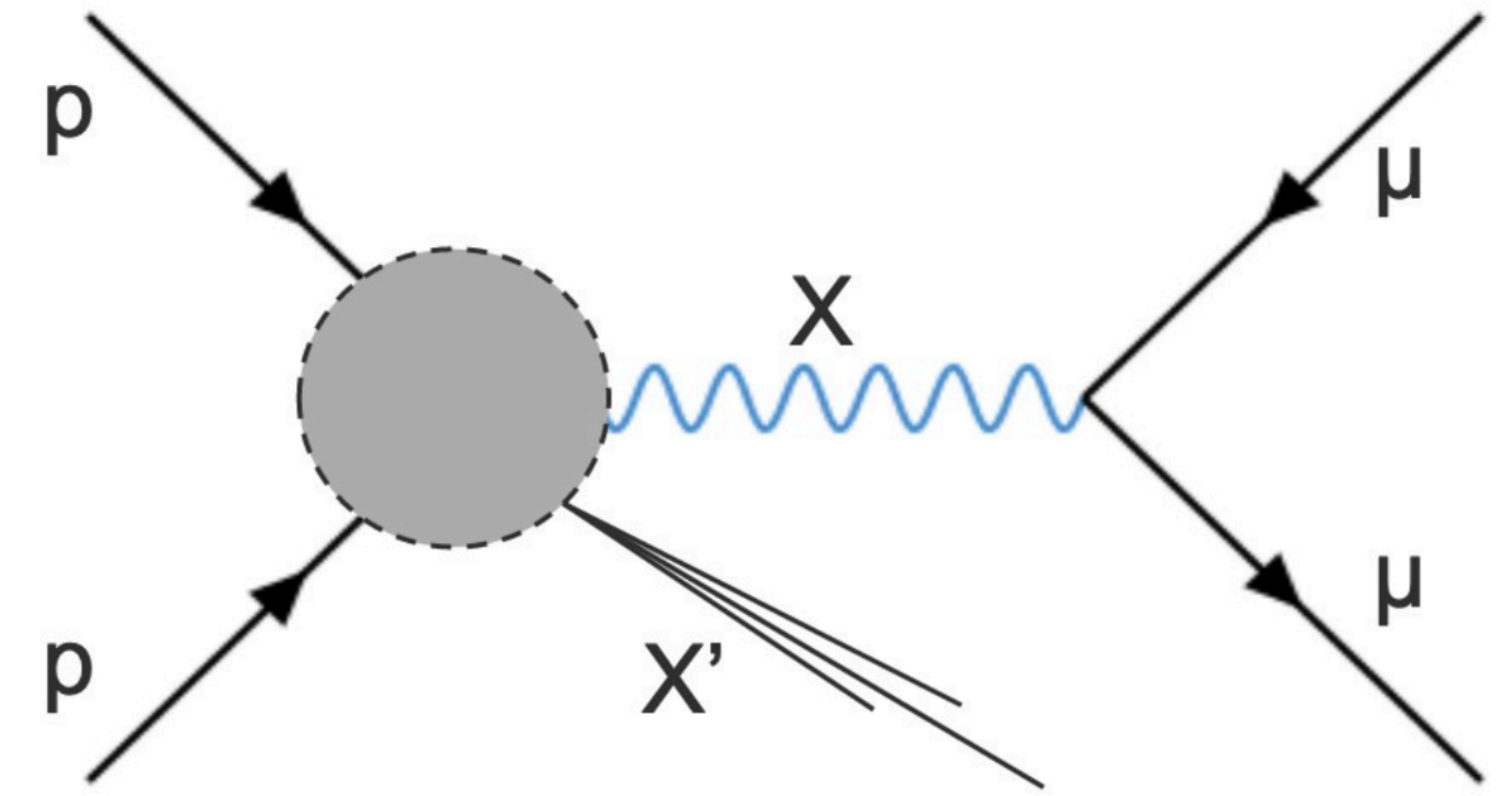
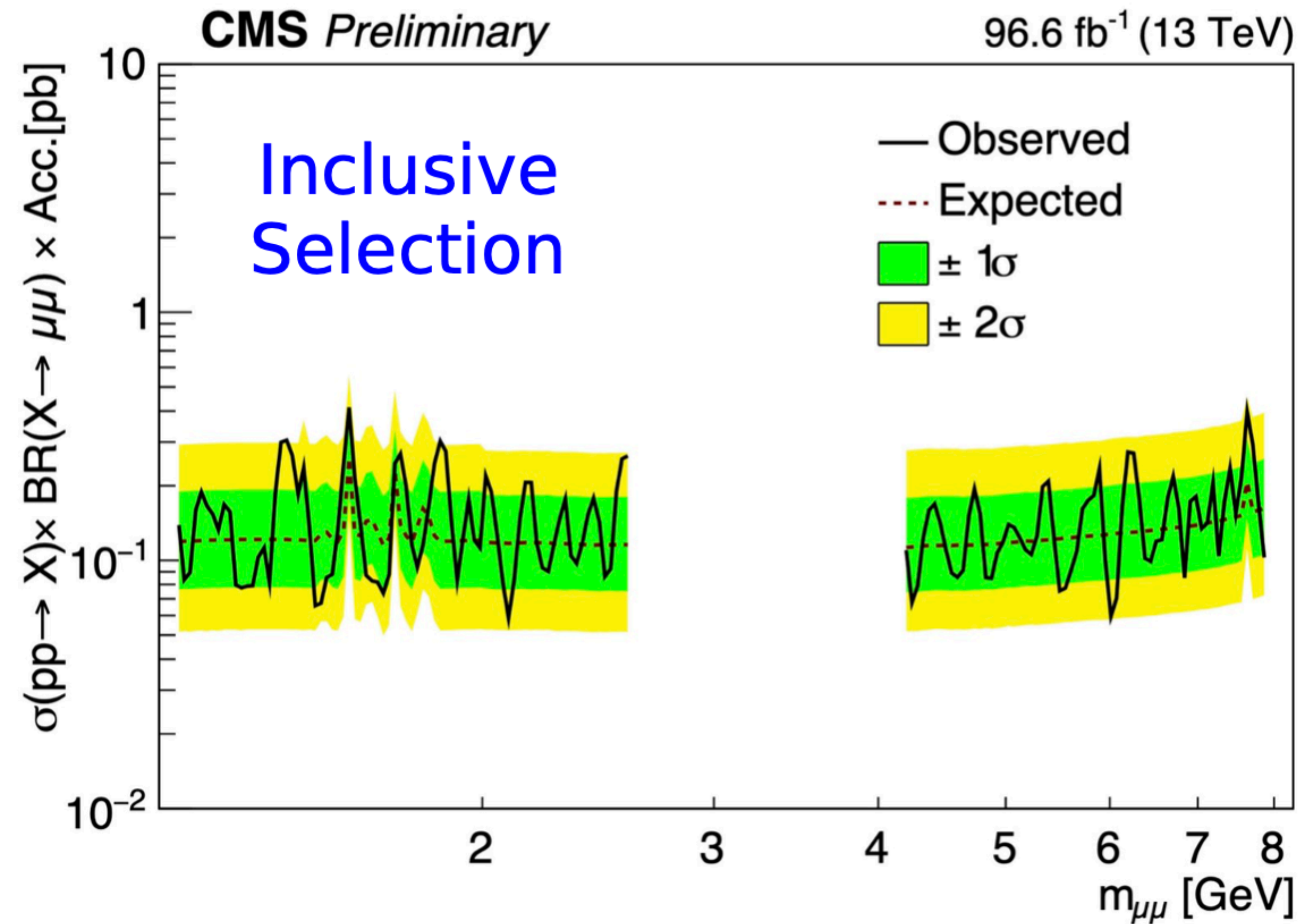
Then for many hadronic searches



Finally CMS has it Fully commissioned for multi-muon final states

Also used for the search for unknown resonances

Bump hunt on the dimuon mass

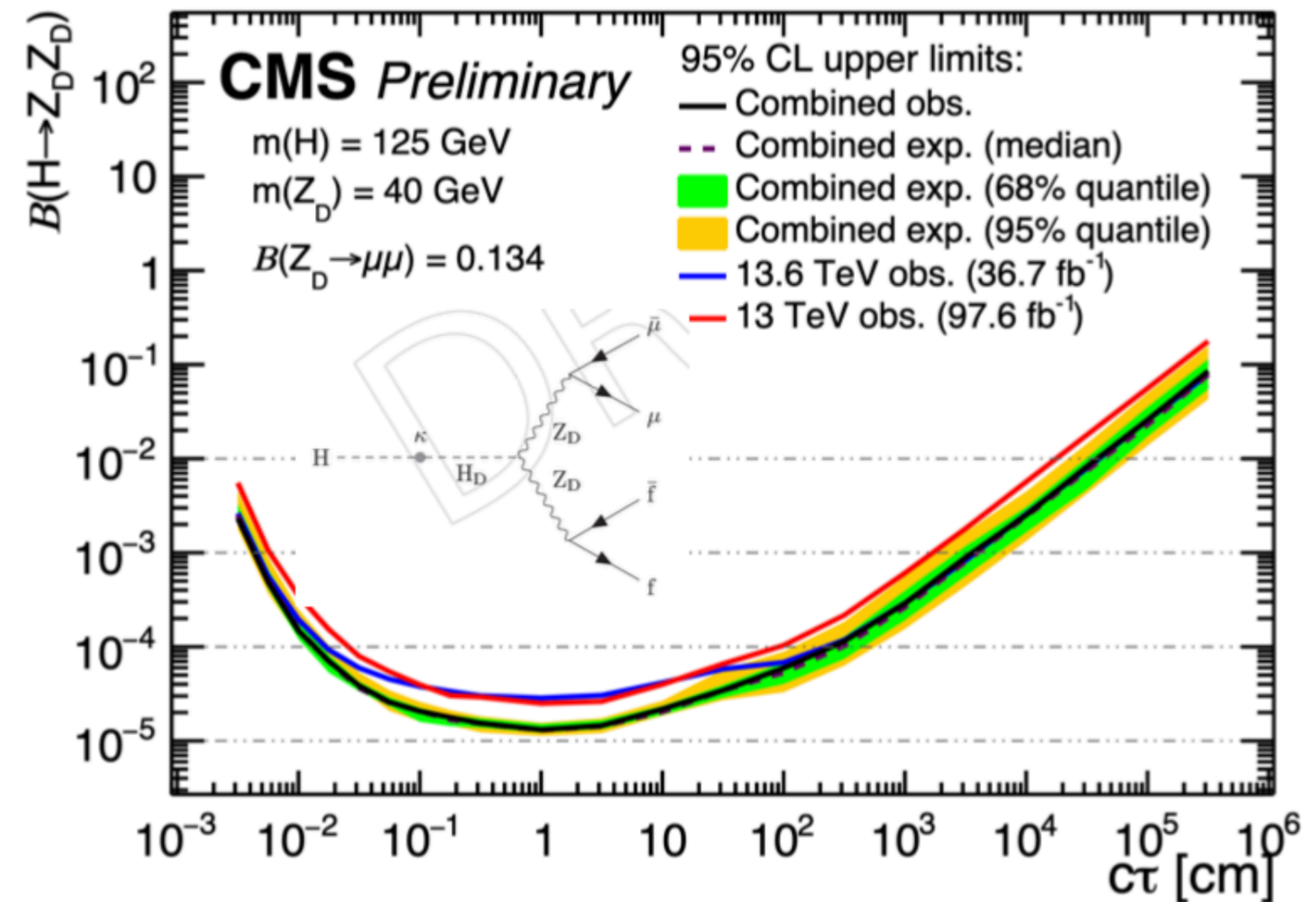
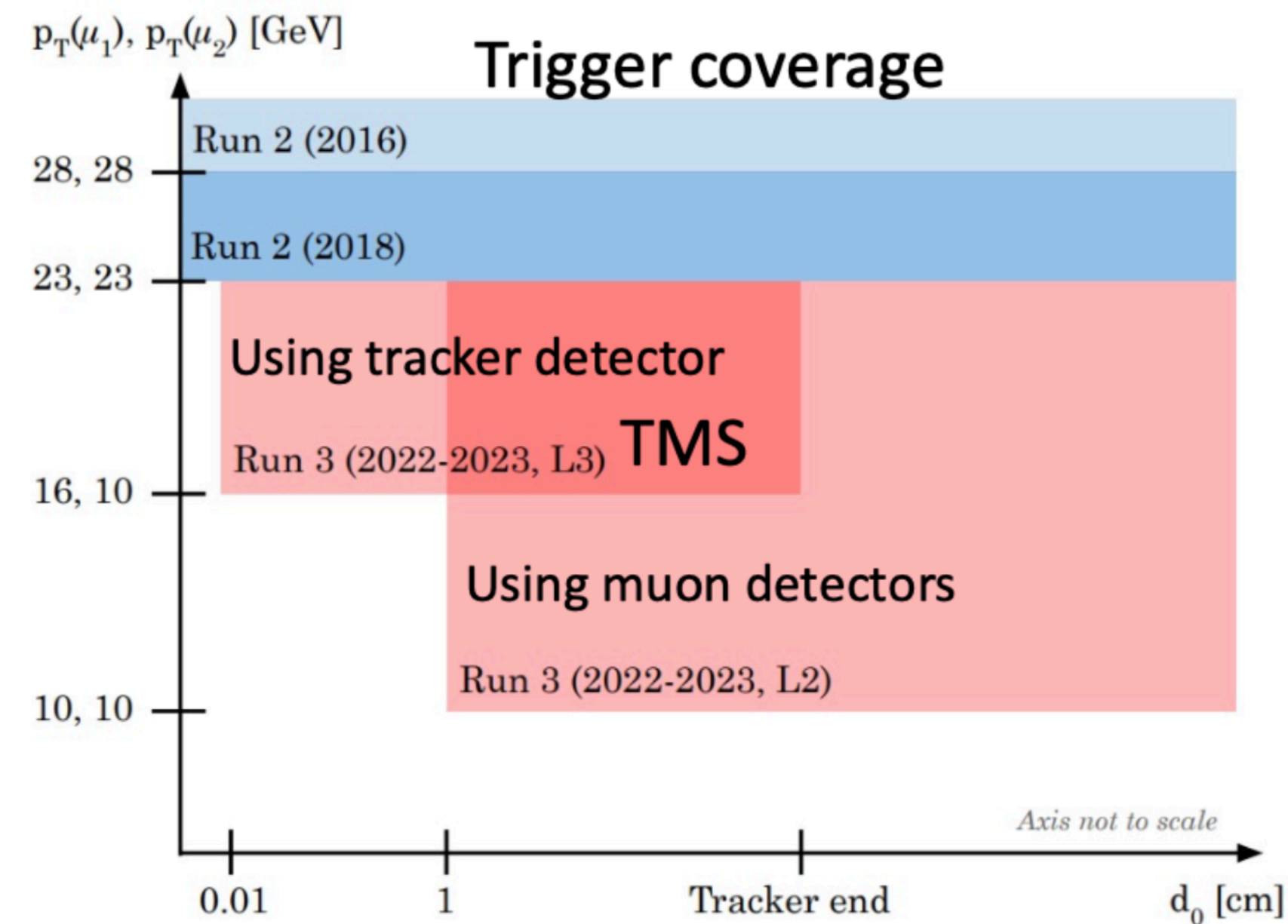


Reaches very low di muon masses!

Long Lived particles @Run 3

First search for new physics at Run 3, looking for long lived particles decaying into muon pairs: selects muon originating from a common secondary vertex spatially separated from the primary interaction point from few hundred μm to several meters.

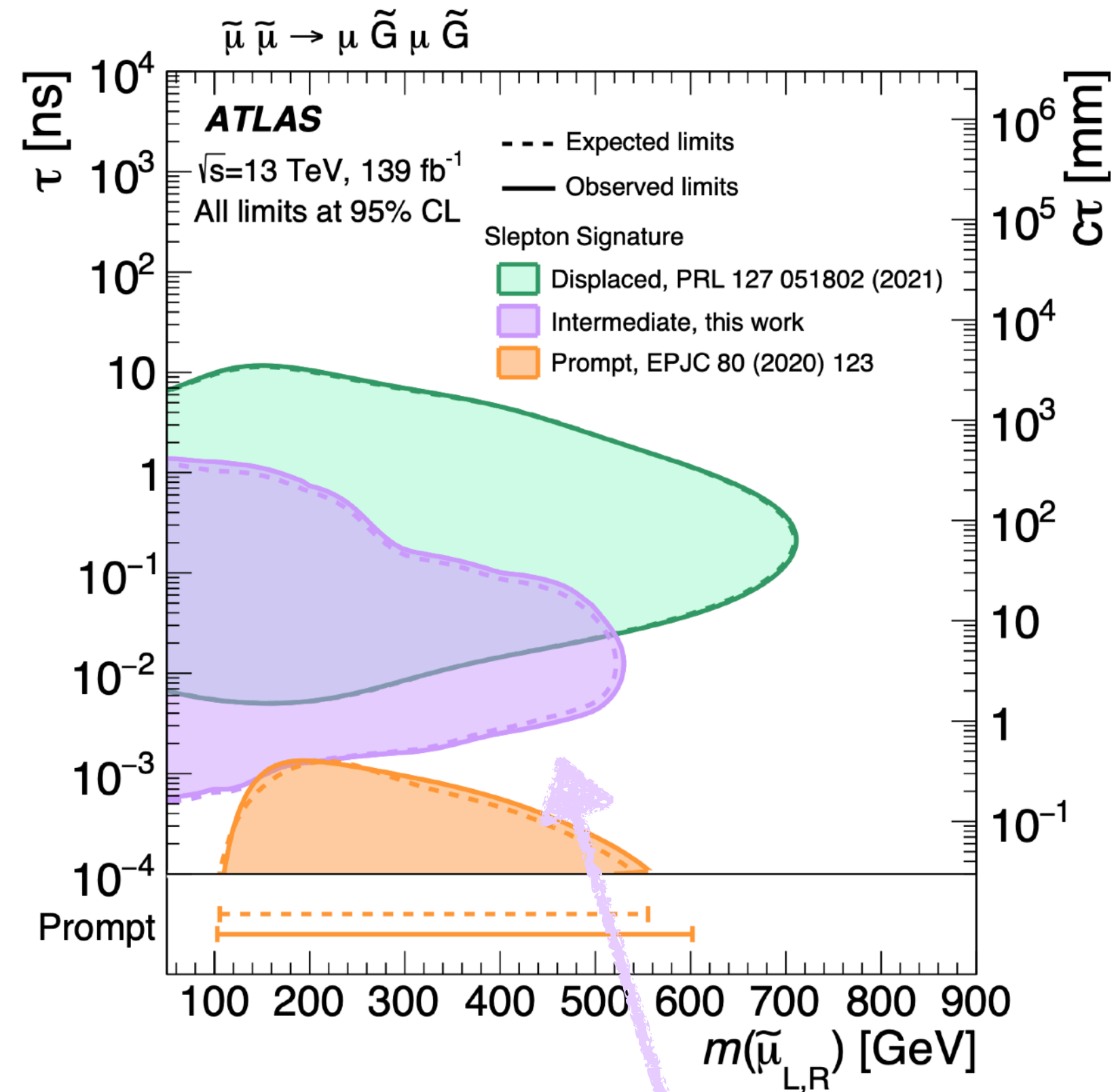
Substantial improvement of efficiency due to improved triggers for displaced muons (and also analysis techniques)



Displaced particles

<https://arxiv.org/abs/2305.14931>

Search for muons with small displacements

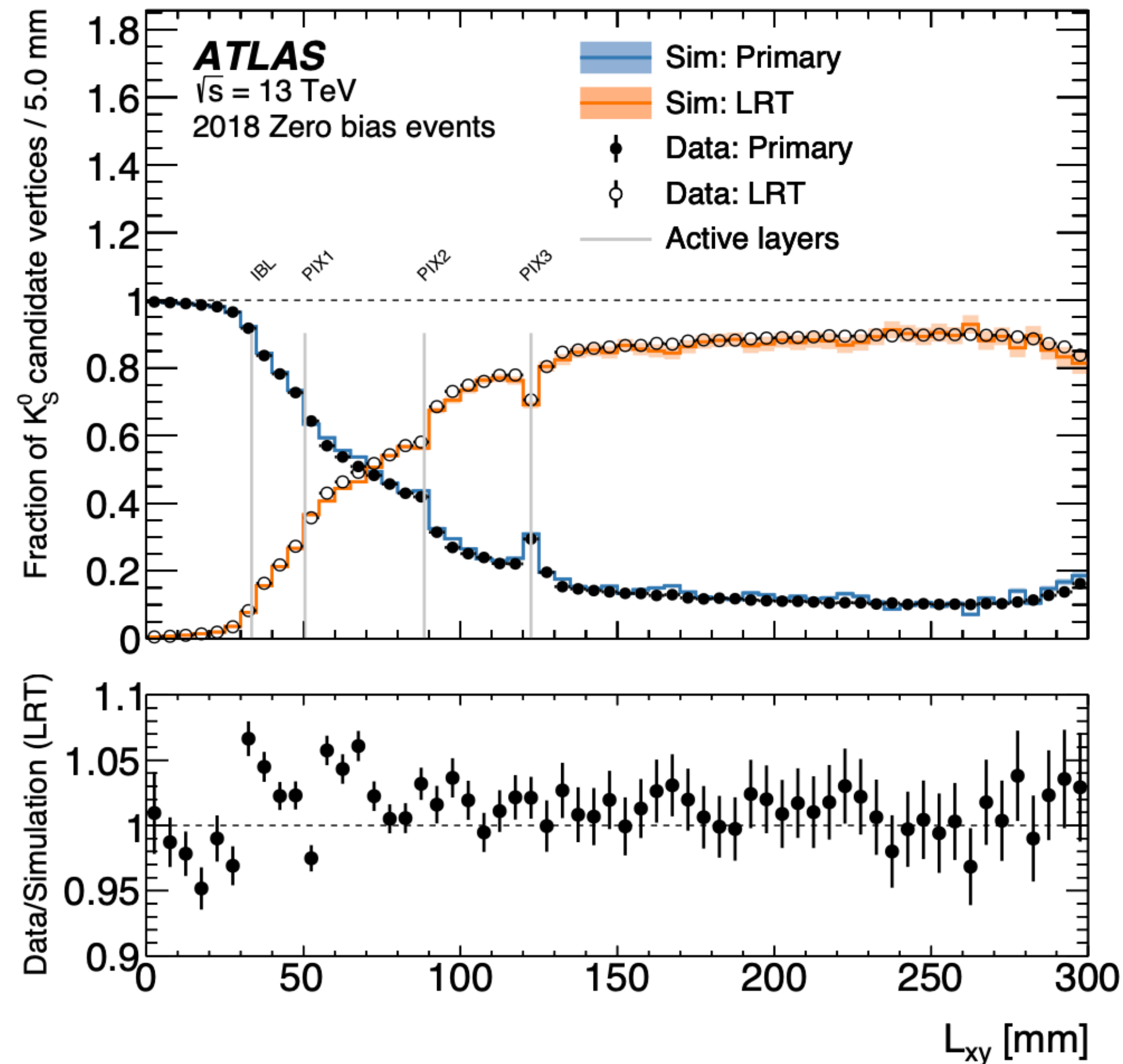


Covers phase space
between prompt & displaced
muons

ATLAS Tracking improved performance! [arXiv:2308.09471](https://arxiv.org/abs/2308.09471)

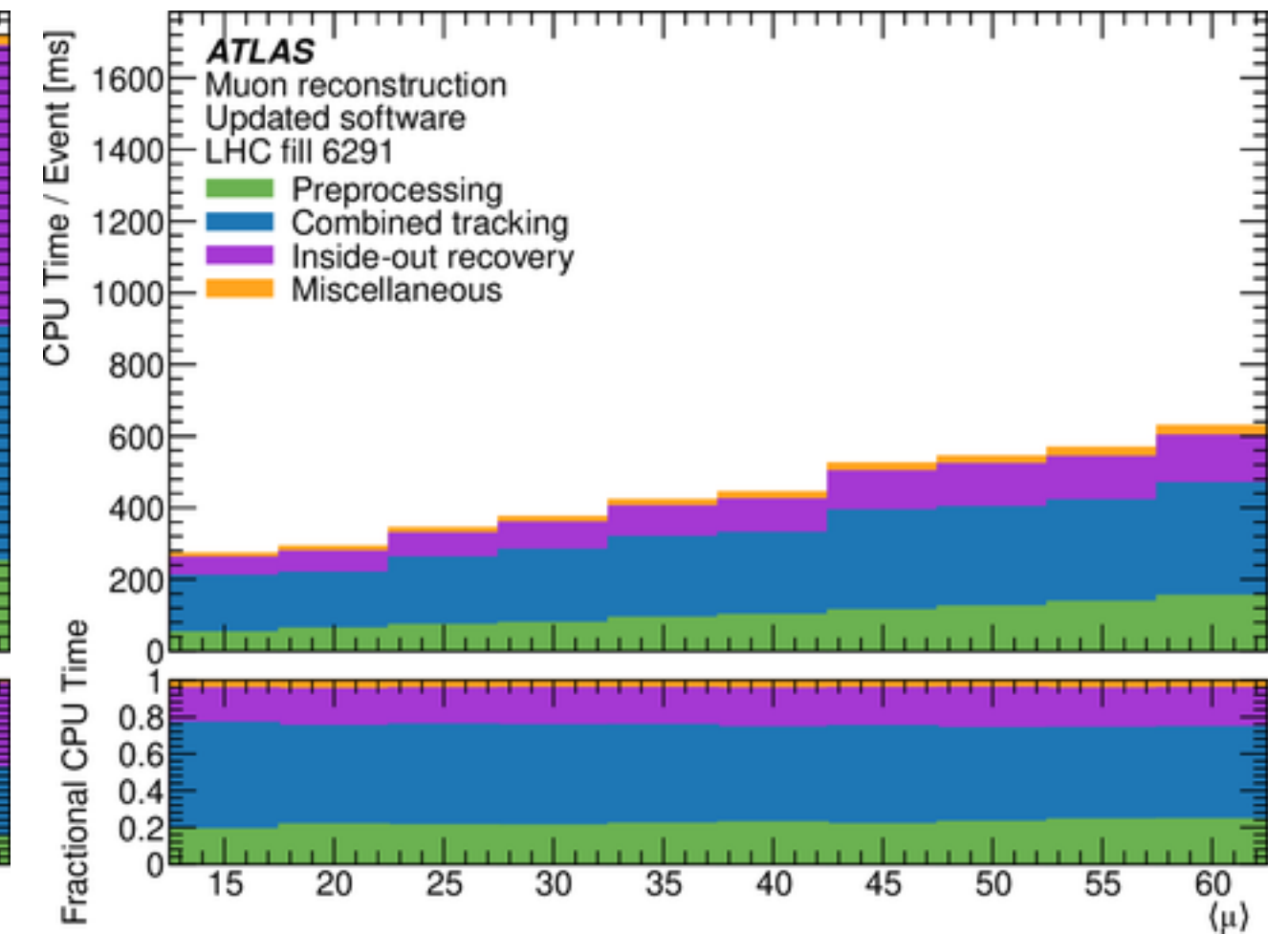
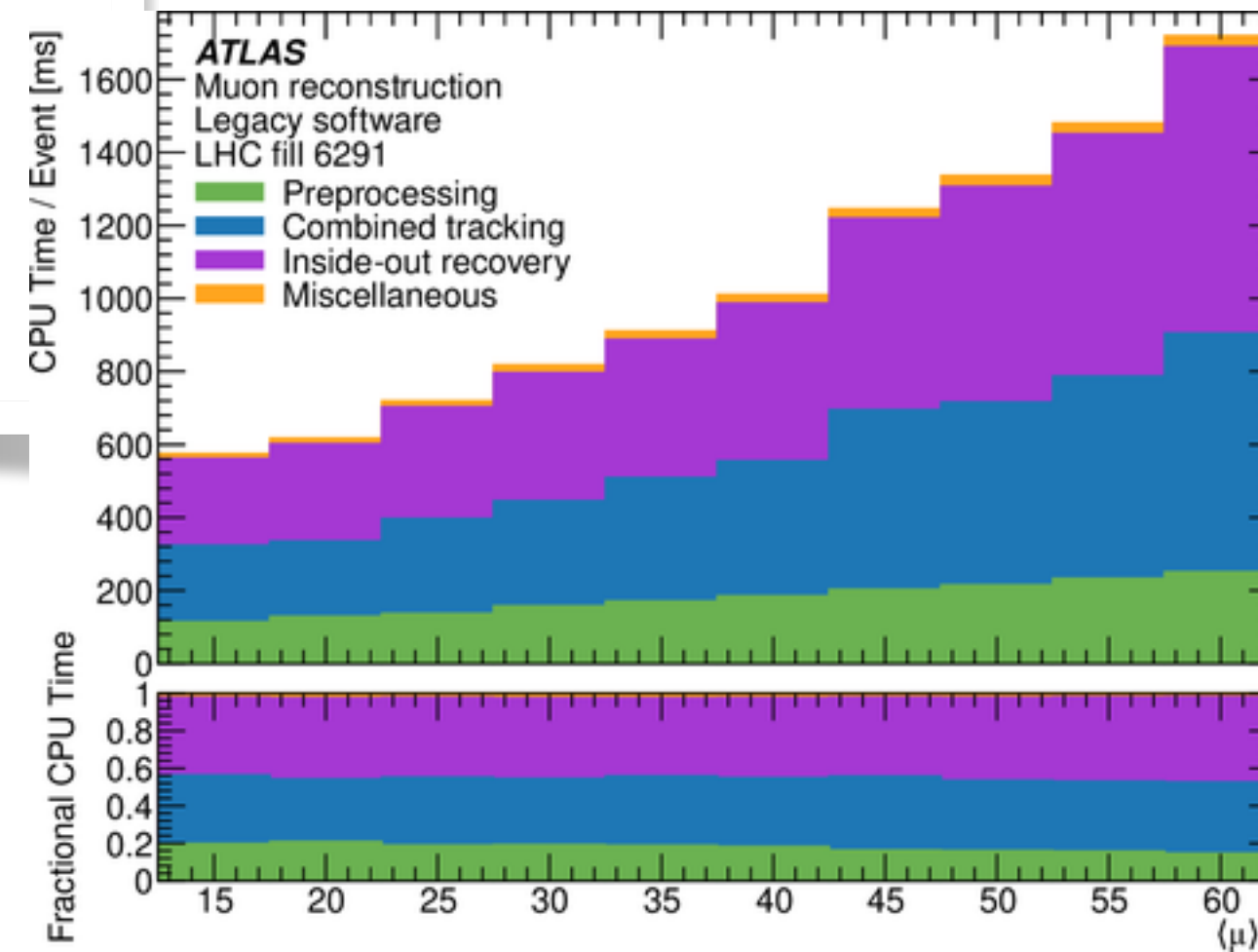
Tracking: For 60 pp collisions ($\langle \mu \rangle$) per bunch crossing:

- track reconstruction nearly 3 times as fast
- no significant reduction in reco efficiency
- large reduction in combinatorial fake tracks rate.



Legacy software

Updated software



[arXiv:2304.12867](https://arxiv.org/abs/2304.12867)

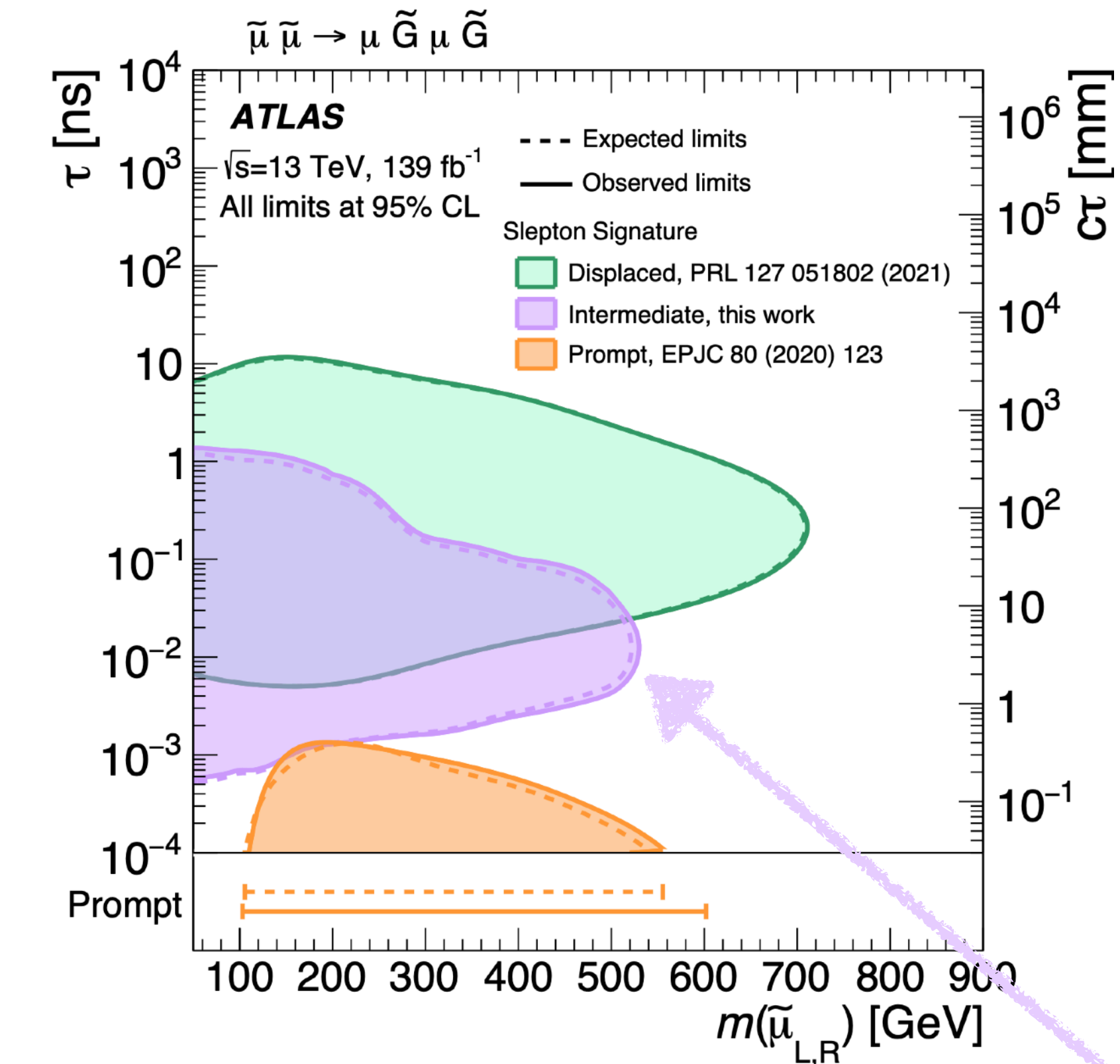
Improved Large Radius Tracking (LRT) deployed
LRT available in standard reconstruction

improves long lived particles searches!
10x(50x) improvement in CPU usage(disk usage)
(also present at HighLevelTrigger HLT for Run 3).

Run3/HL_LHC Long Lived particles ATLAS

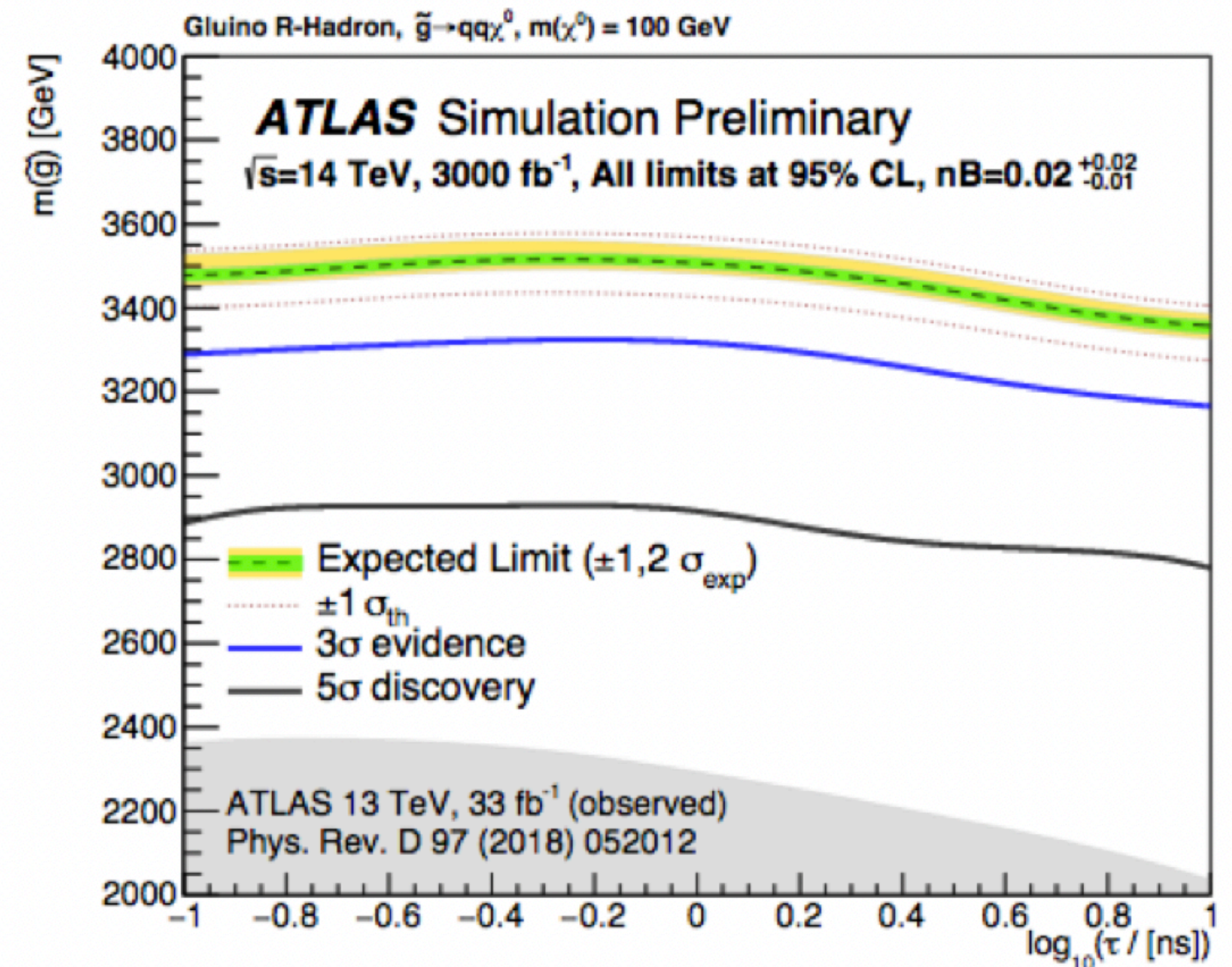
<https://arxiv.org/abs/2305.14931>

Search for muons with small displacements



Goal for Run 3: exploring uncovered phase spaces!
Covers phase space between prompt & displaced muons

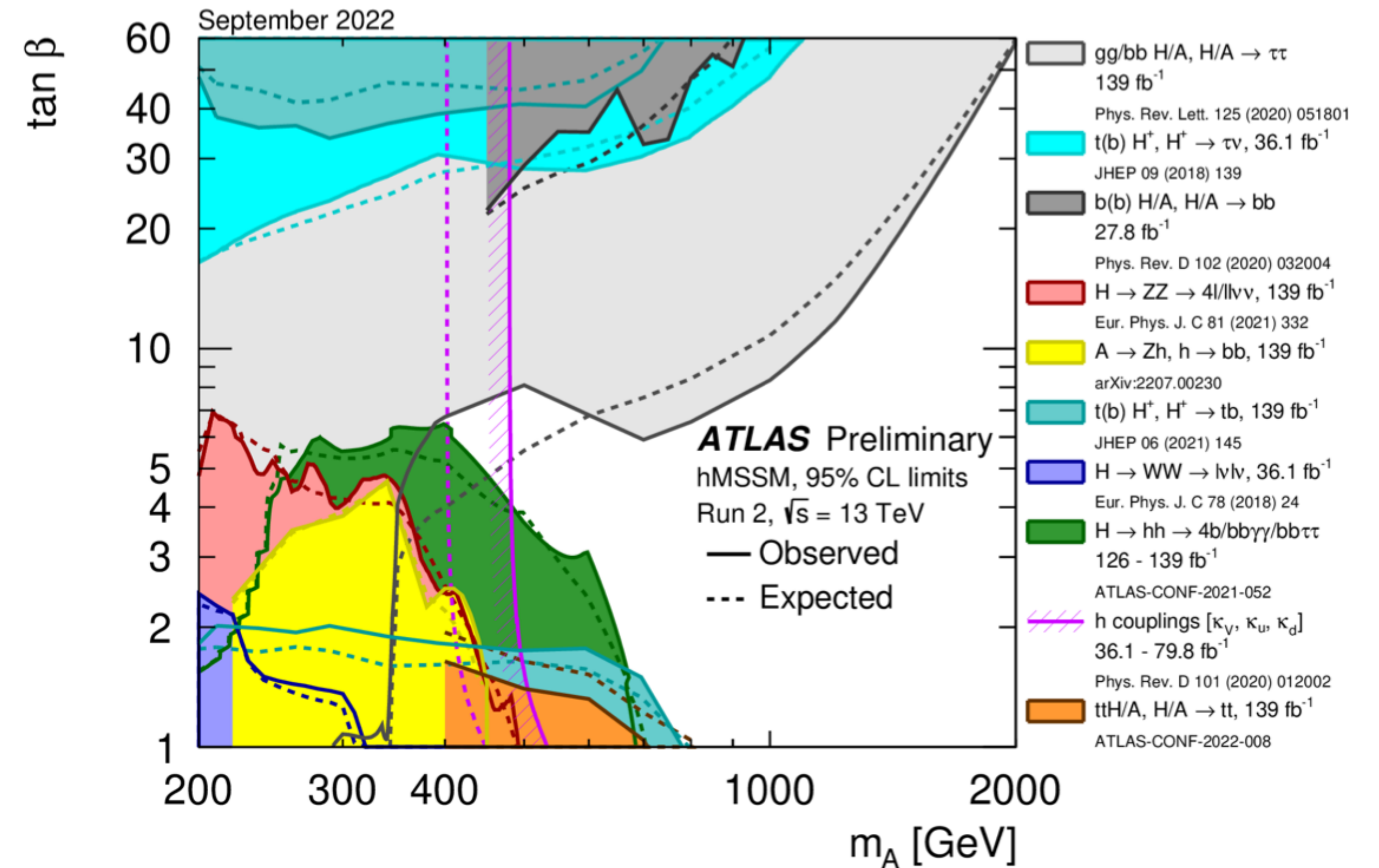
HL-LHC: long-lived particles: projected sensitivity for the mass of a long lived gluino which hadronizes after production into an R-hadron, and then decays through a virtual squark into a pair of SM quarks and a neutralino. In the analysis high Emiss and one displaced T vertex are required.



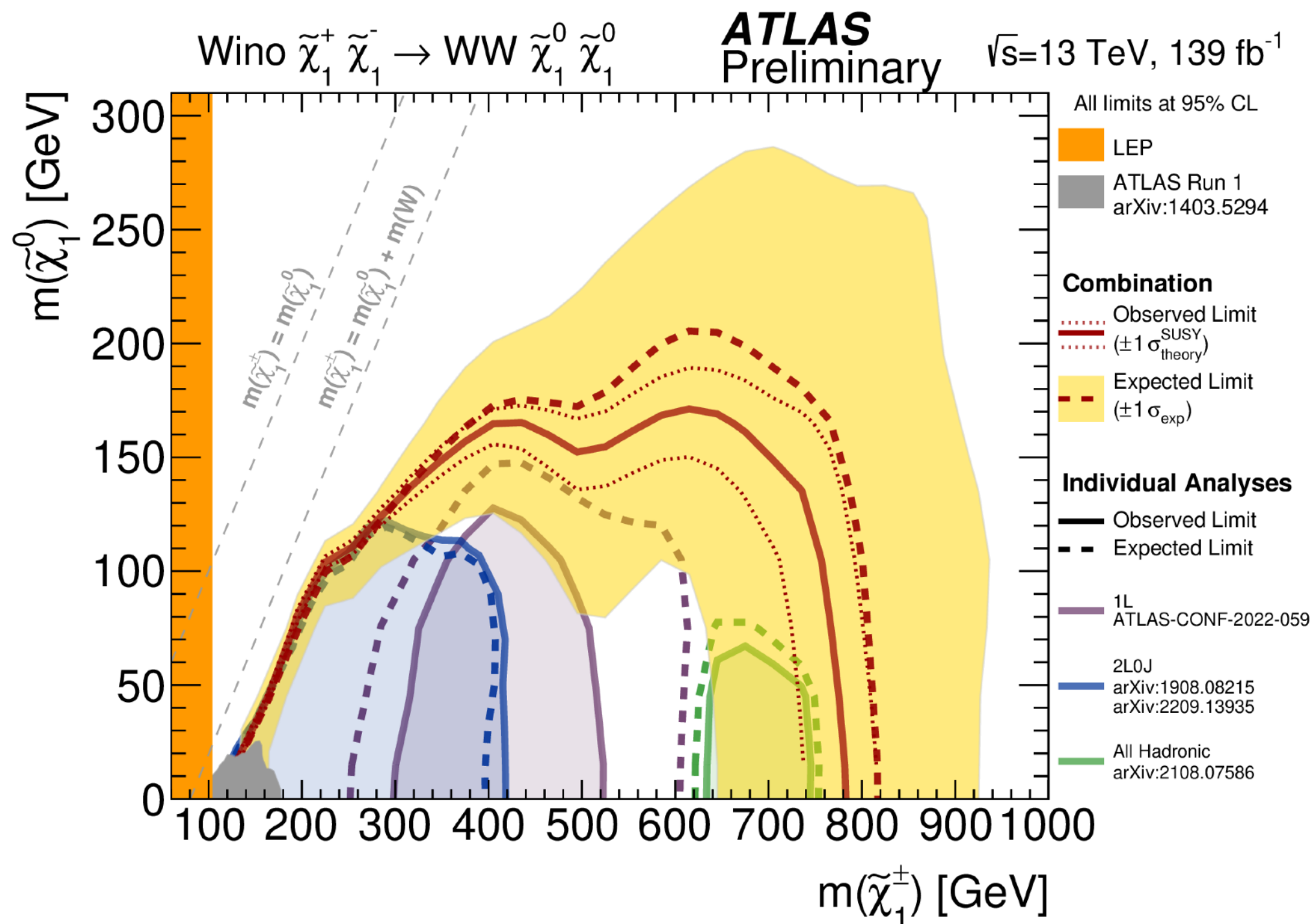
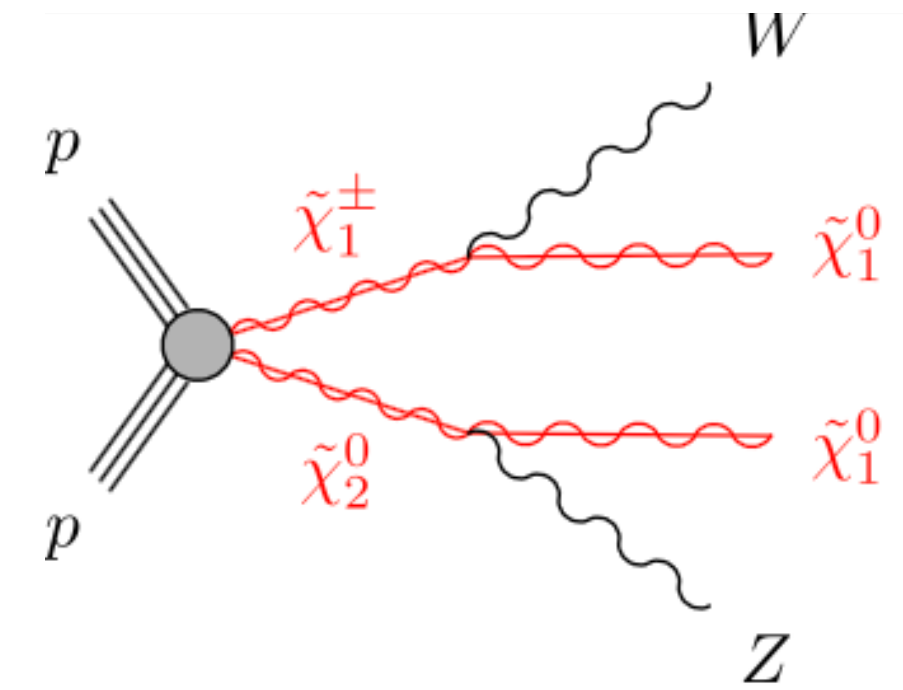
Extended Higgs sector and Supersymmetry

Strategy for Run 3 extended Higgs searches

- For probing the Electroweak sector, continued searches of extended Higgs sectors are of particular interest.
- An extended Higgs sector is, for example, needed to lead to a first-order phase transition in the early universe.
- Using Run 2 data, the searches for an extended Electroweak sector have been vast, covering searches for new diboson resonances, exotic Higgs decays and direct and indirect searches for additional Higgs bosons.
- see hMSSM model exclusion: wide range and complementary of different final states and search modes, but also unexplored gaps.
- The added data of Run 3 will benefit us greatly in closing these gaps.



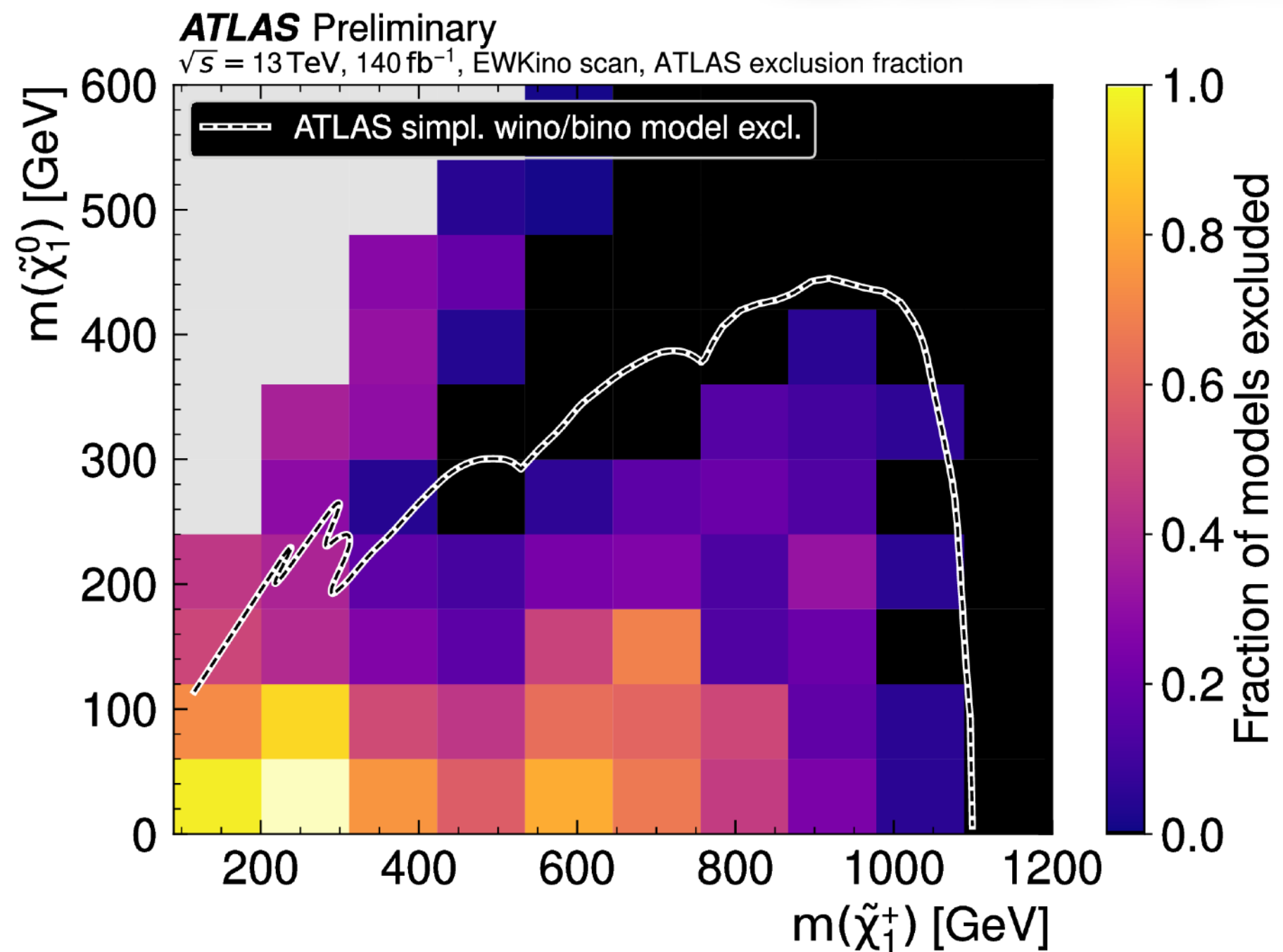
- Normally in SUSY we use simplified models. Here we present
- **Statistical** Combination of multiple SUSY EWK analyses, improving exclusion limits and exclusion depth
- chargino, neutralino production decaying via W,Z



The combined result fills the gap between the individual analyses

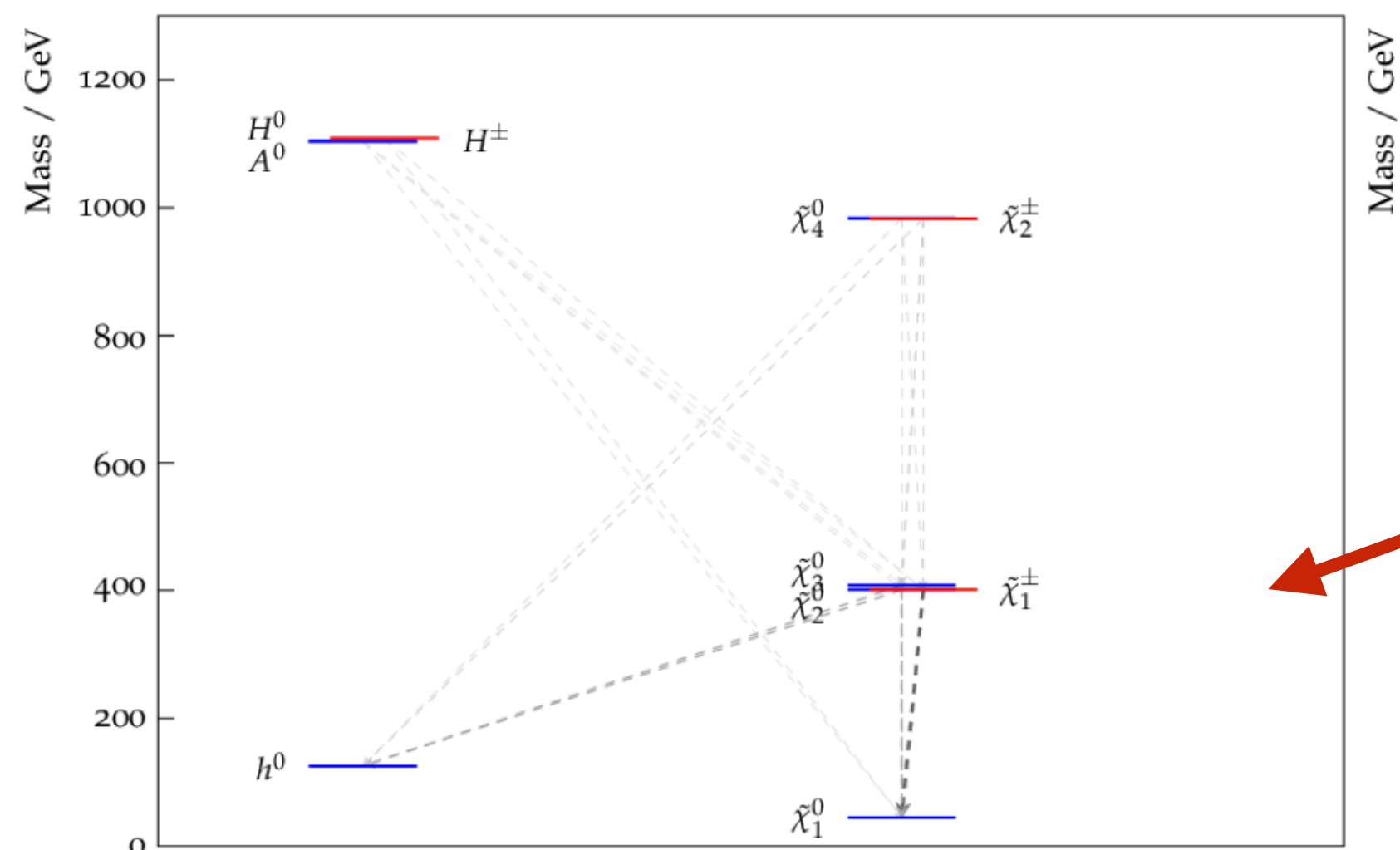
Simplified models come with shortcomings

It is mandatory to make sure that we are searching in the correct phase space



- scan exploring Phenomenological Minimal Supersymmetry (pMSSM) a UV complete Model (normally simplified models are used)
- imposes LHC + external constraints (LEP, flavor, precision EWK, Dark Matter)

- Almost full exclusion of low-mass χ^0_1 in regions where a low-mass neutralino would not oversaturate the dark matter relic abundance



(a) Z/h funnel region

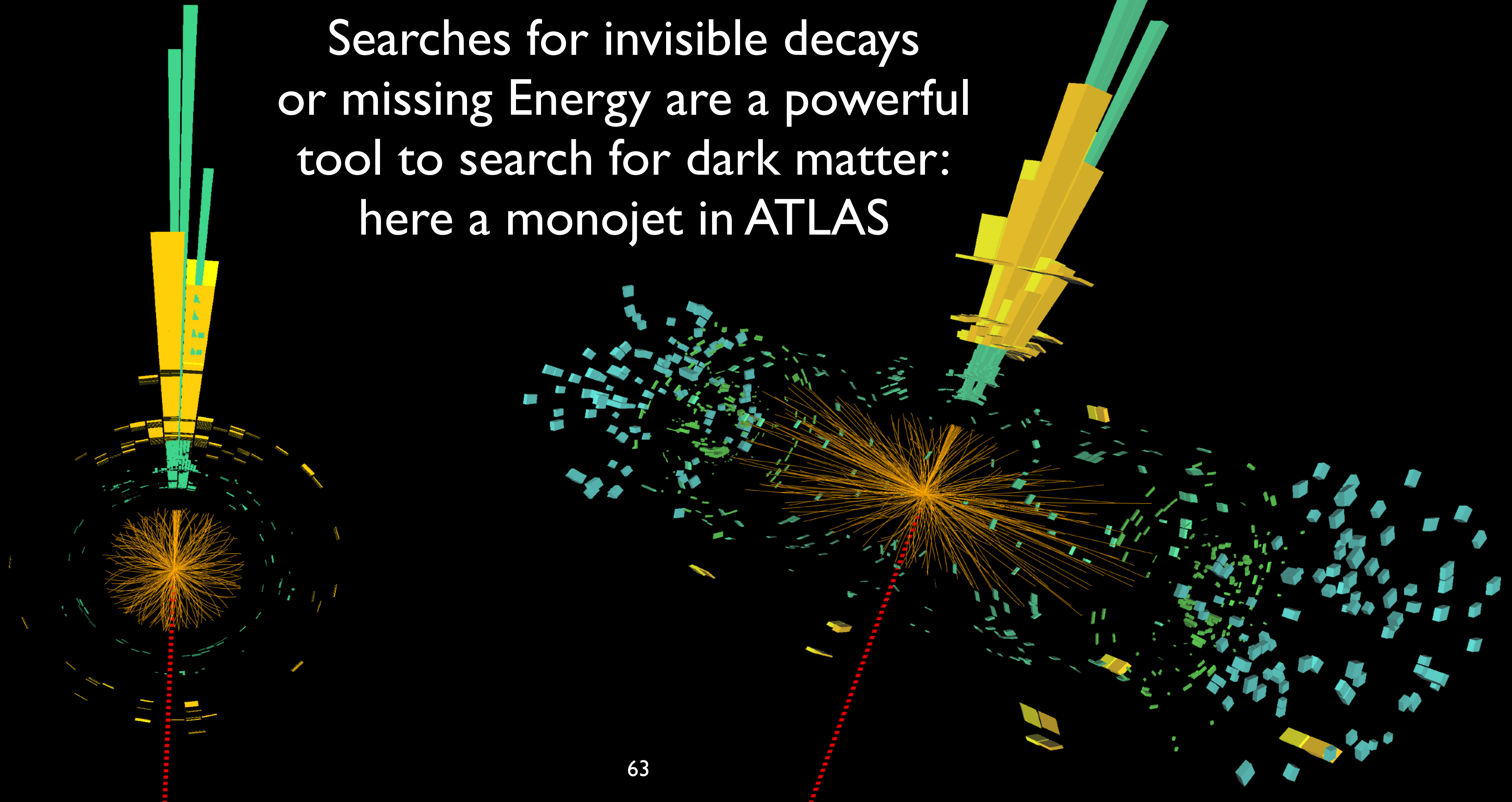
- Example spectra for surviving supersymmetry models that are not excluded despite having a mass-spectrum within published ATLAS simplified model contours.

Dark matter and invisible decays

Run: 337215
Event: 2546139368
2017-10-05 10:36:30 CEST

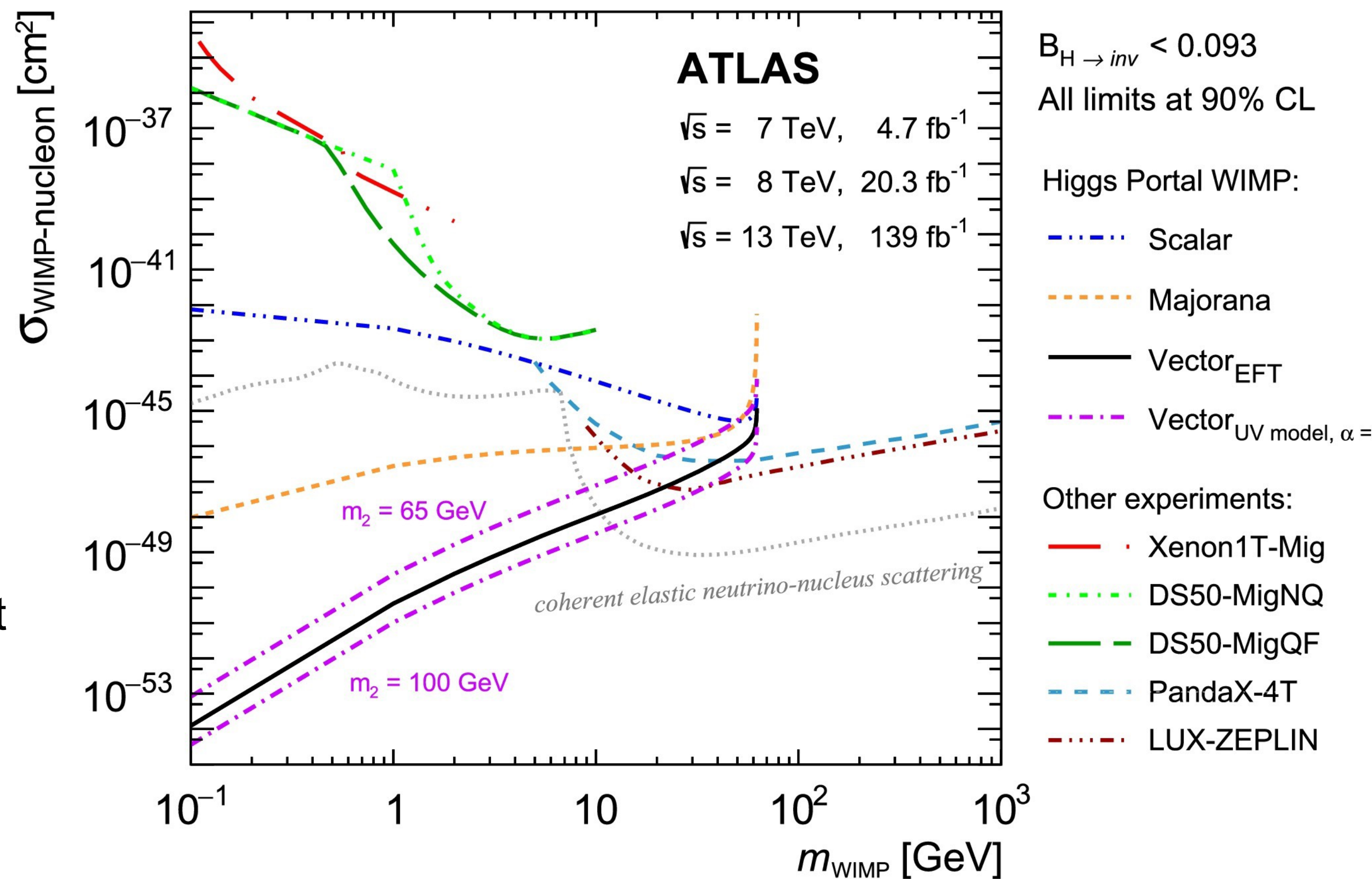
$E_T^{\text{miss}} = 1.9 \text{ TeV}$
jet $p_T = 1.9 \text{ TeV}$

Searches for invisible decays
or missing Energy are a powerful
tool to search for dark matter:
here a monojet in ATLAS



- SM particles get mass through the Higgs. Dark matter could behave the same way and be produced in Higgs decays
- SM Higgs invisible decays are $< 0.1\%$
- **The analysis: $B(H \rightarrow \text{inv})_{\text{obs}} < 10.7\%$ @95%CL**
 $B(H \rightarrow \text{inv})_{\text{exp}} < (7.7\%)$ @95%CL best to date
- These results are also interpreted in the context of models where the SM Higgs boson acts as a portal to dark matter
- exclusion regions extend to very low DM mass-
> very important to improve

90% CL on the spin-independent WIMP-nucleon scattering cross-section

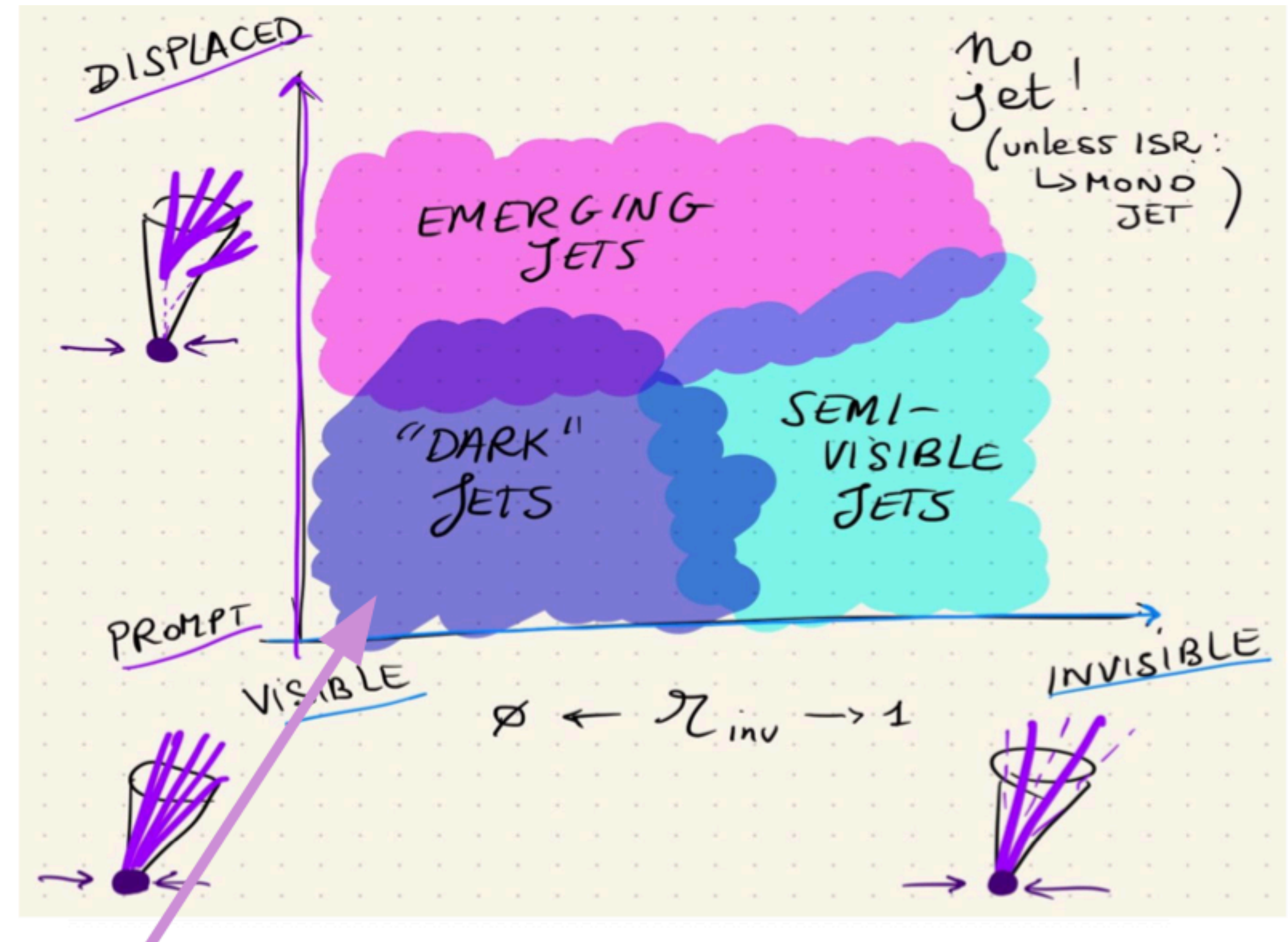


The regions above the limit contours are excluded

Dark Matter searches not only WIMPs: QCD like dark sector

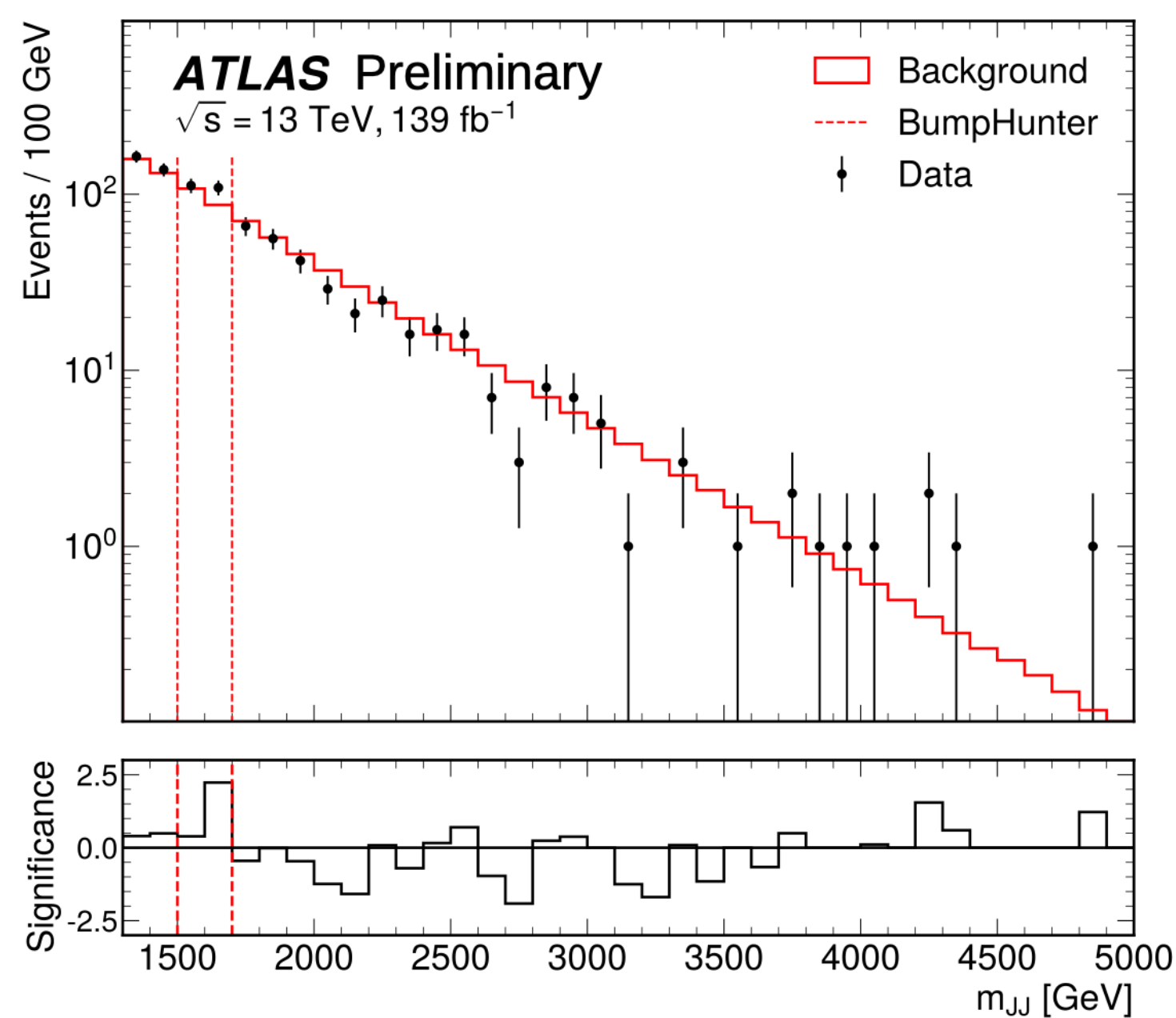
QCD-like dark sector producing dark showers
Dark hadrons can decay completely or partially in a QCD-like fashion:

- Semi-Visible jets
- Emerging jets
- dark jets from Stable dark hadrons with unusual large R-jet dijet signatures (higher charged-particle multiplicity)



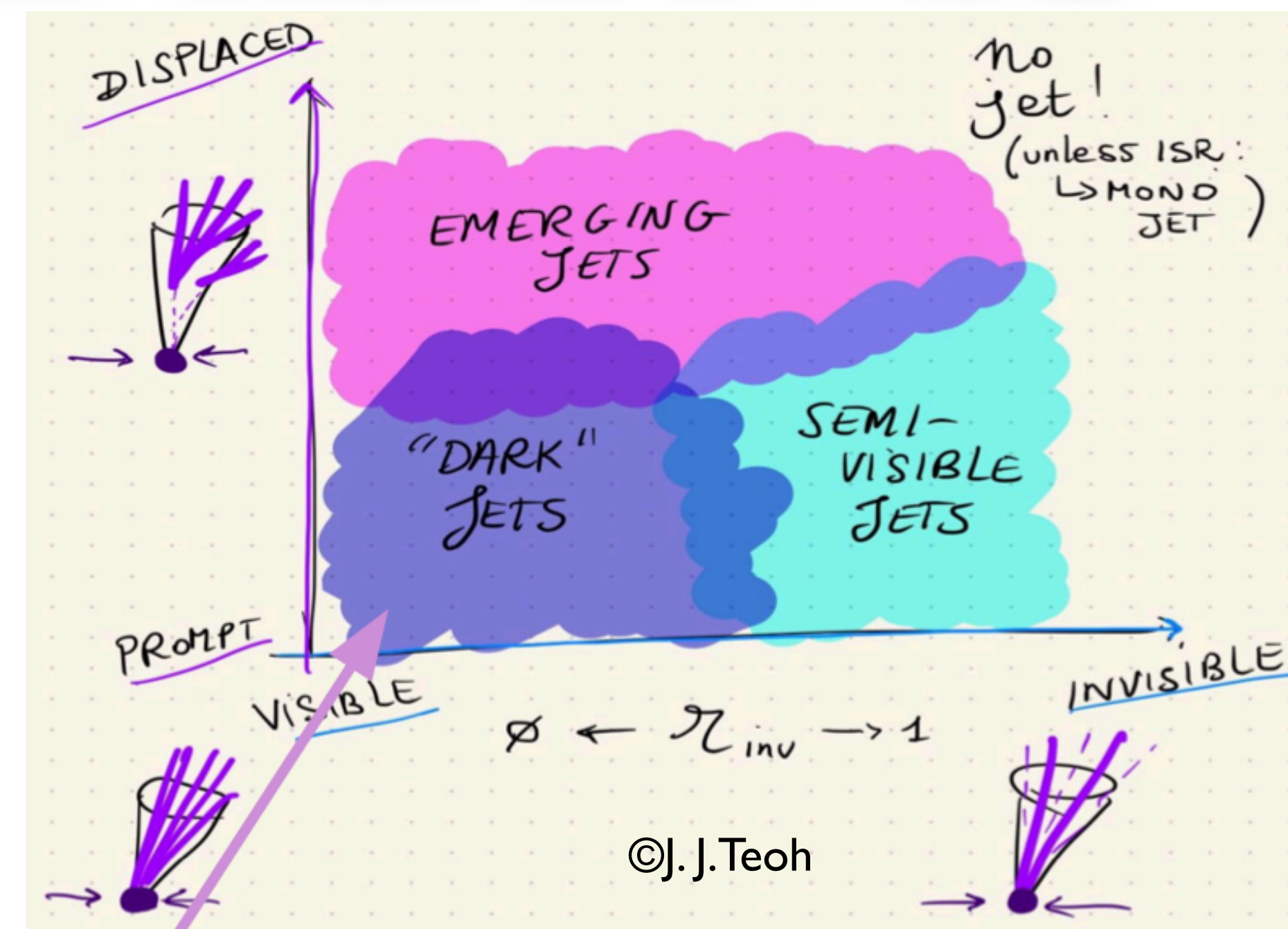
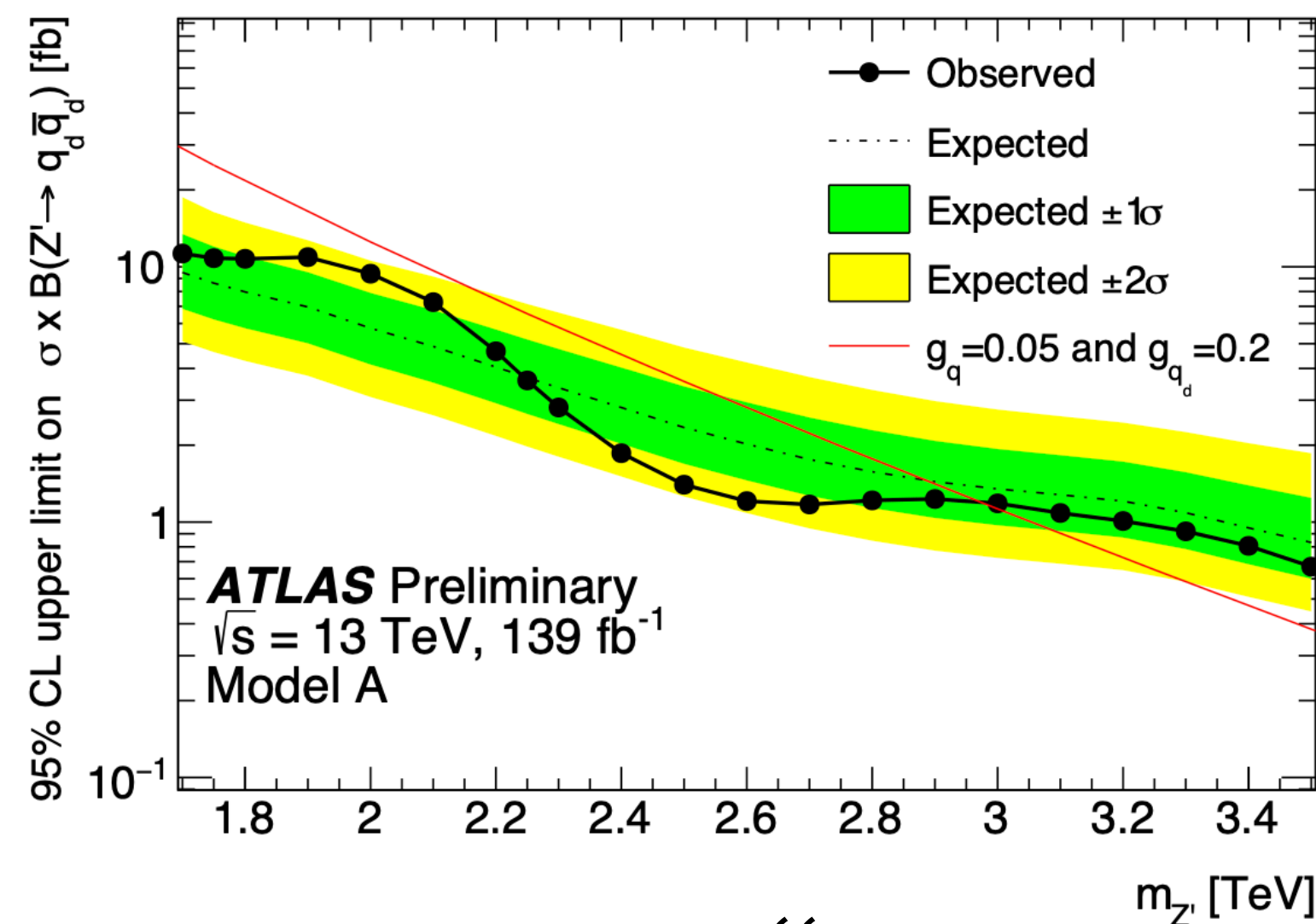
Dark Matter: QCD-like dark sector linked to SM via Z'

- Stable dark hadrons with unusual dijet signatures (higher charged-particle multiplicity)
- Search for dark jets bump in the mass spectrum of two large-R jets.

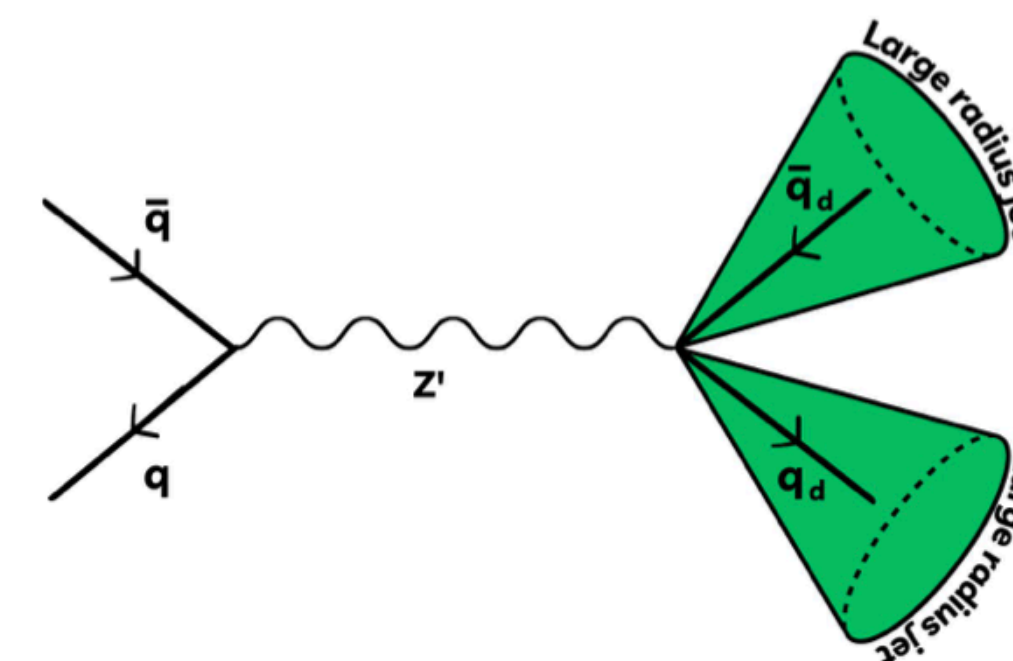


no excess

ATLAS-CONF-2023-047



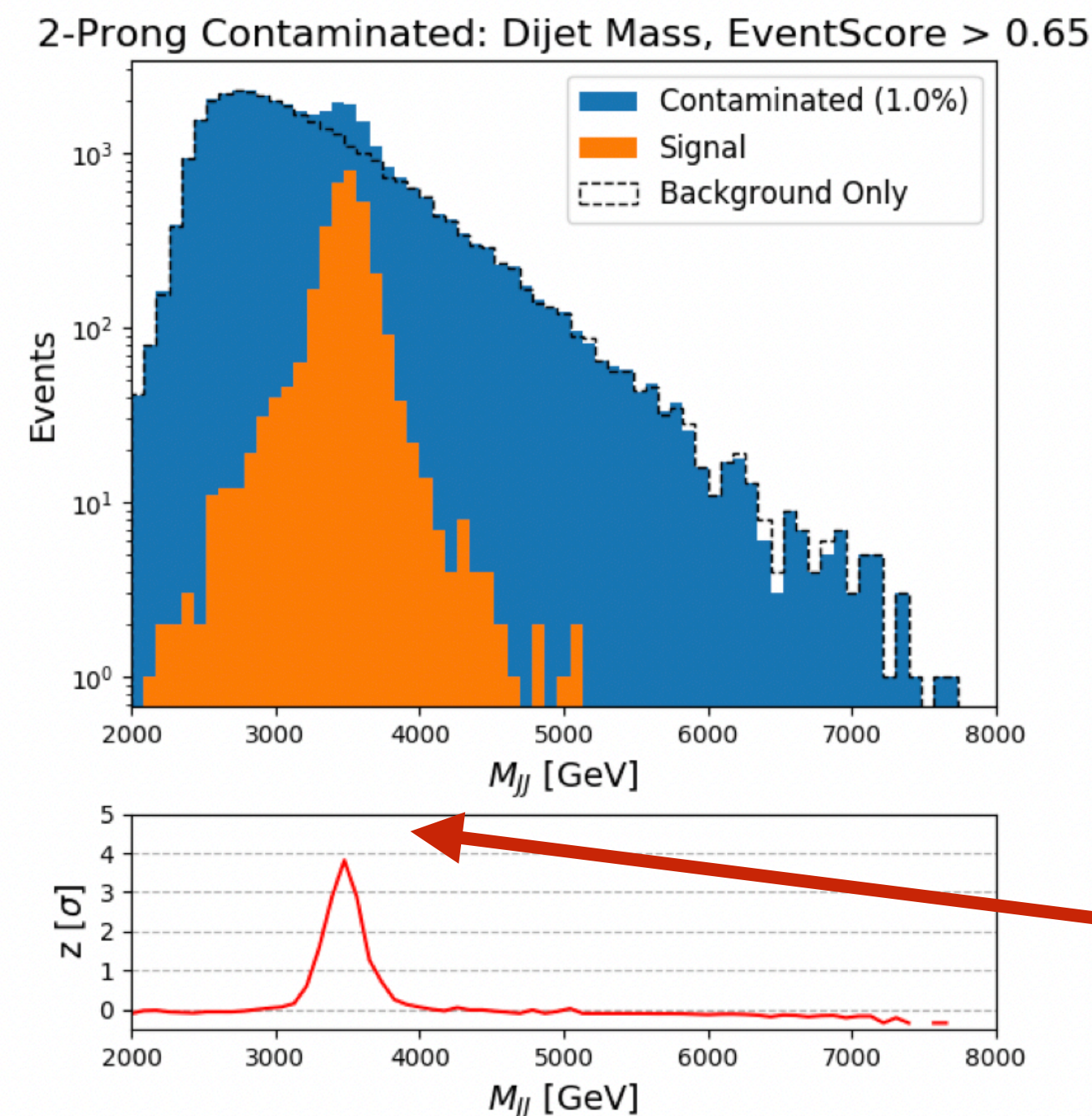
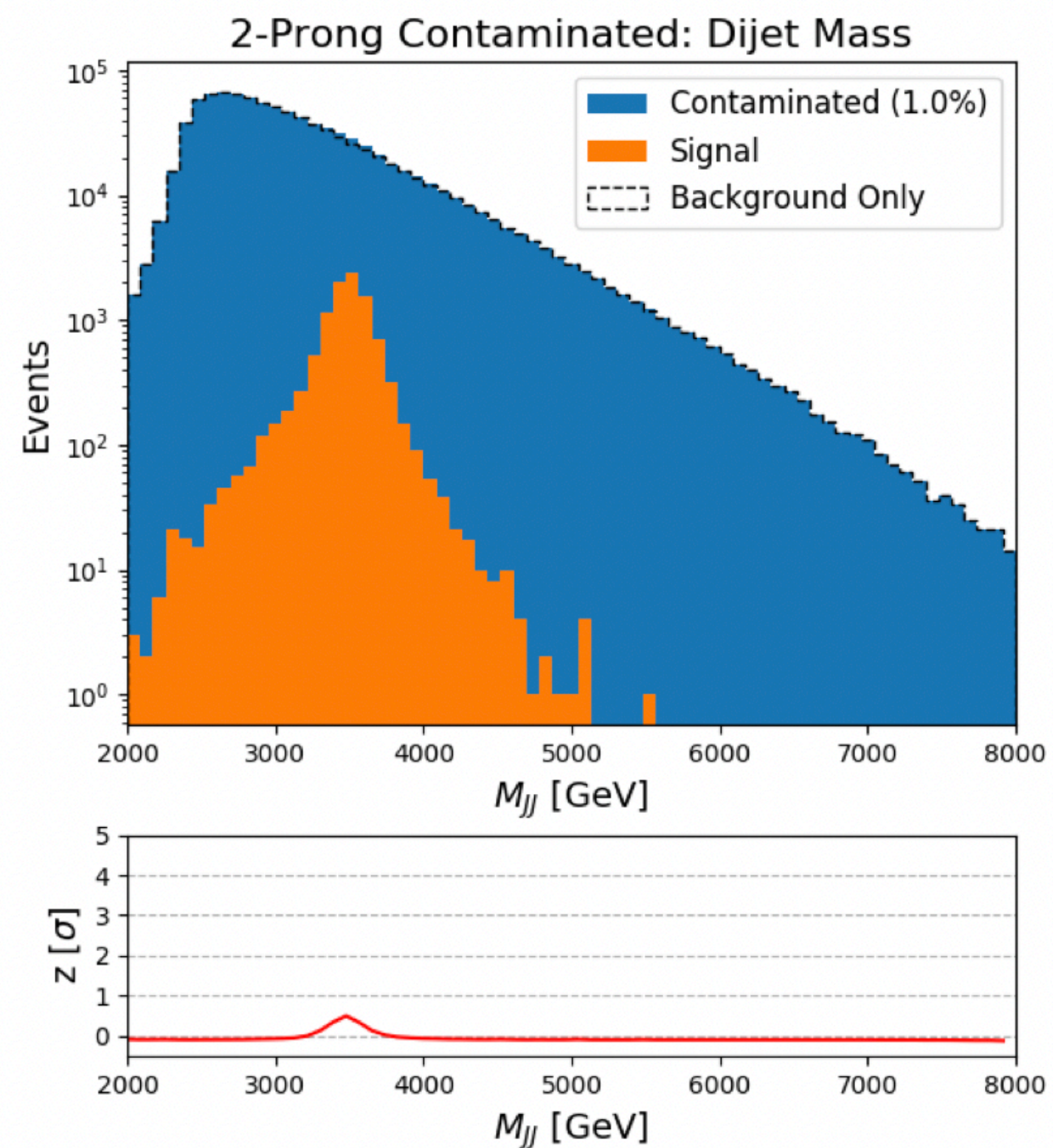
uncovered phase space!



Going model independent

Anomaly detection

- Detecting anomalies using unsupervised Machine learning!
- use Model-independent discovery region introduced with novel, data-driven anomaly score (AS). For example searching for boosted hadronically decaying objects by treating them as anomalous elements of a contaminated dataset.
- the AS for example in this analysis [link](#): is determined from fully unsupervised variational recurrent neural network (VRNN) trained over jets modeled as sequence of constituent four-vectors.

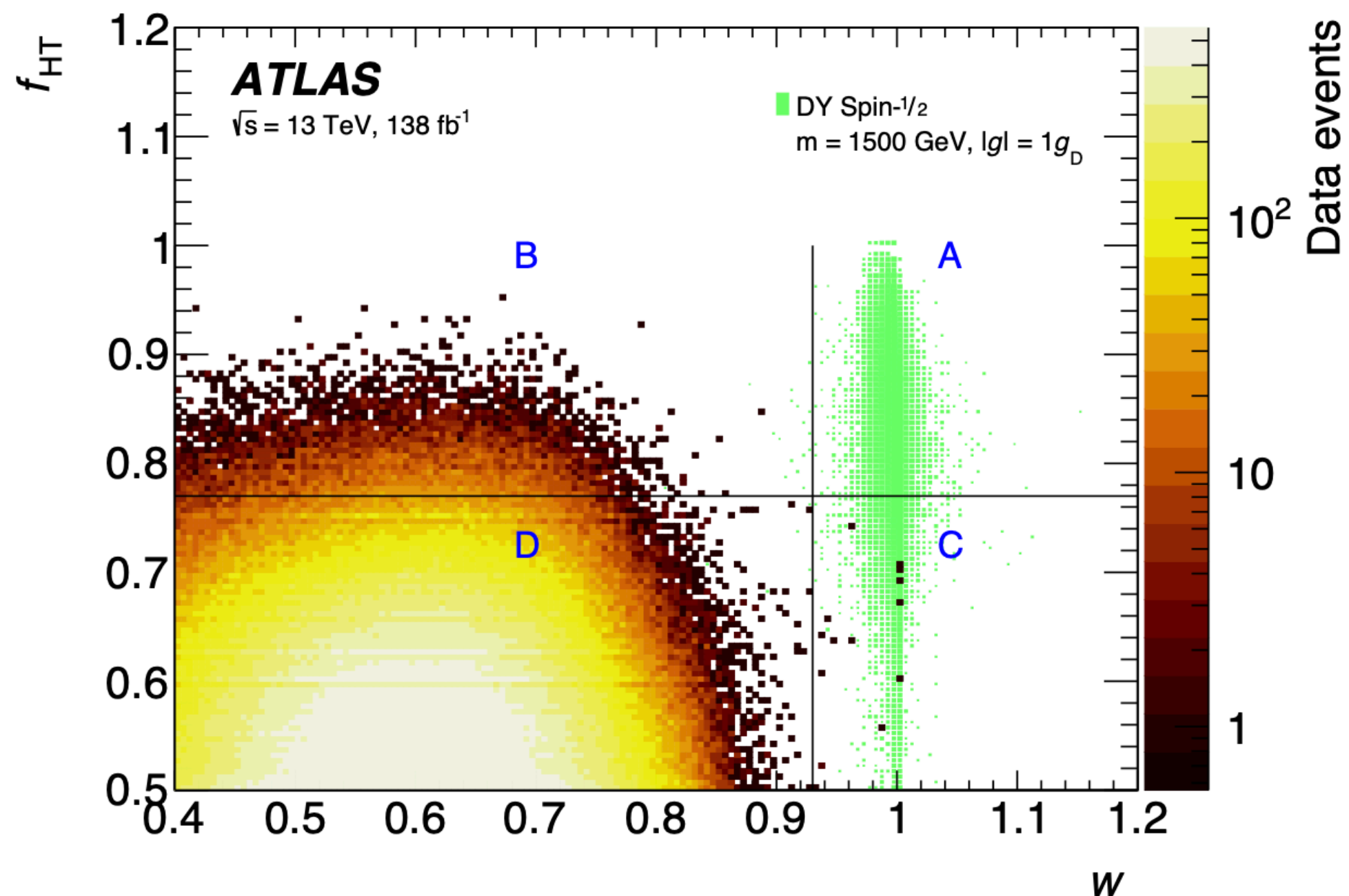
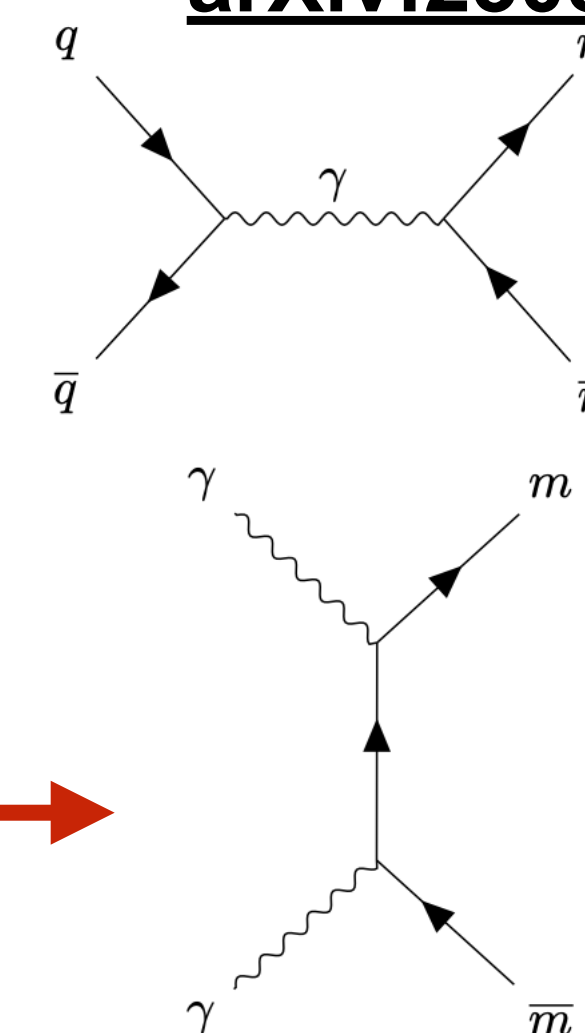


- Variational Autoencoders (VAEs) are built on the idea of standard AEs, with the extension that they are designed to perform Bayesian inference. This assumes that observed data x is generated by some hidden random variable z whose posterior distribution $p(z|x)$ is intractable. The goal of a VAE is to learn an approximate posterior distribution, $q(z|x)$, through training.
- After cutting on anomaly score

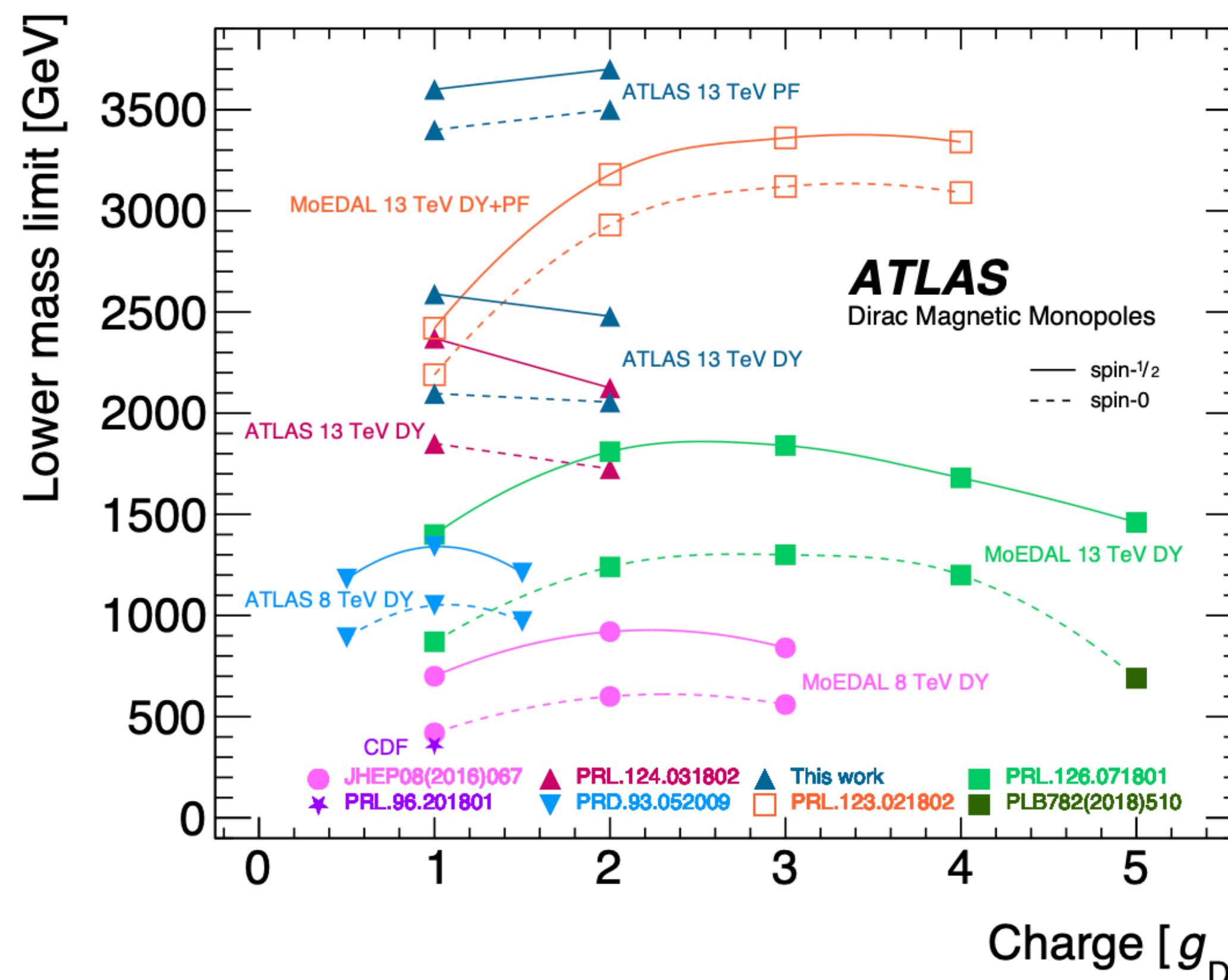
Highly ionizing particles?

Search for highly ionising particles

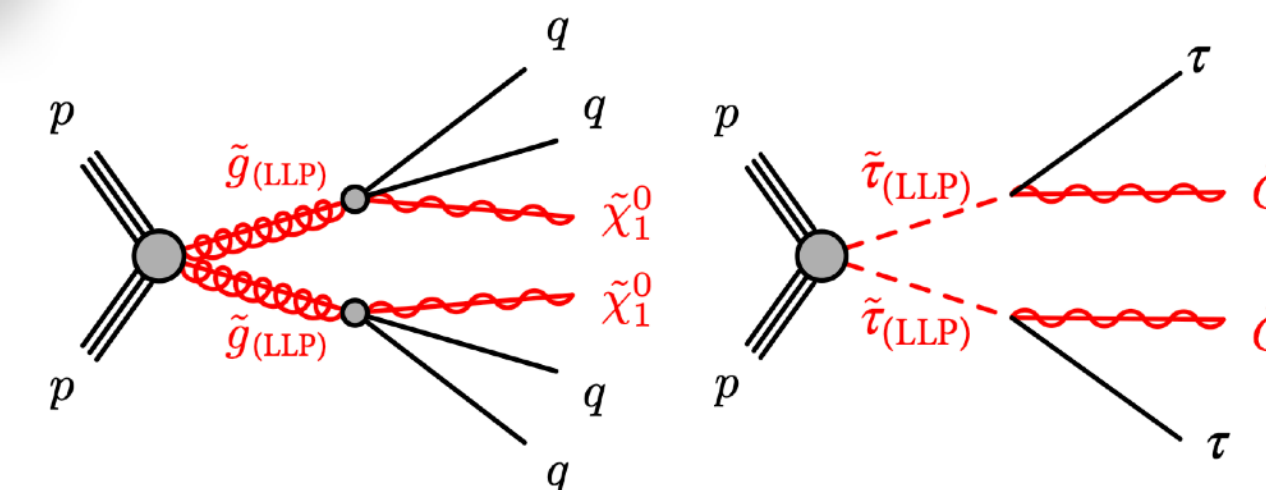
- Search for magnetic monopoles and stable particles with high electric charges
- **improves by factor 3 the previous x-section limits** by ATLAS 36fb⁻¹
- first ATLAS limits on photon-fusion pair production mechanism.



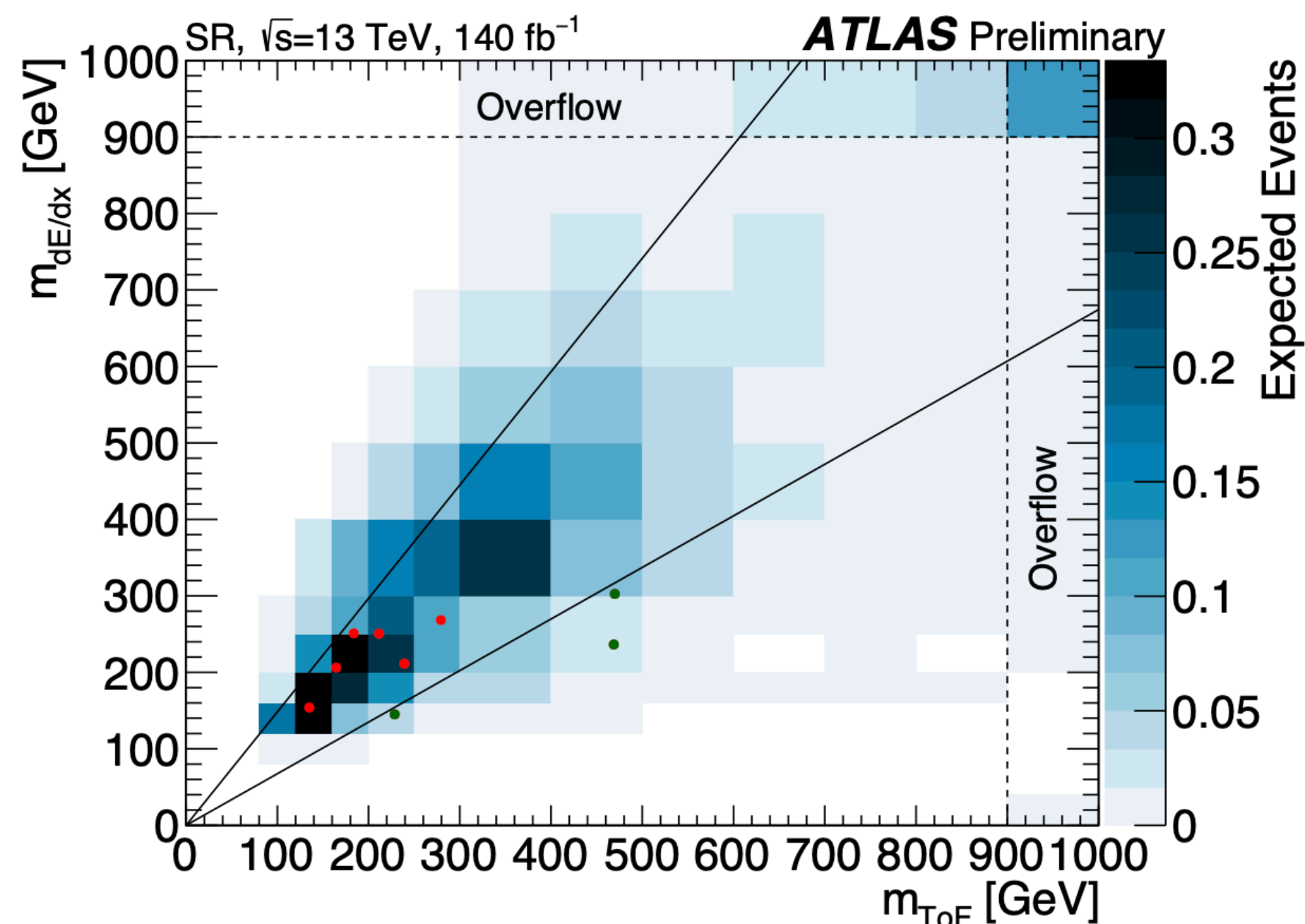
- HIPs produce TRT tracks with δ -rays \Rightarrow many high TRT hits (HT)
- too massive to produce shower in EM calo \Rightarrow low lateral dispersion (w)



Highly ionizing particles (large Energy deposition in pixel detectors)

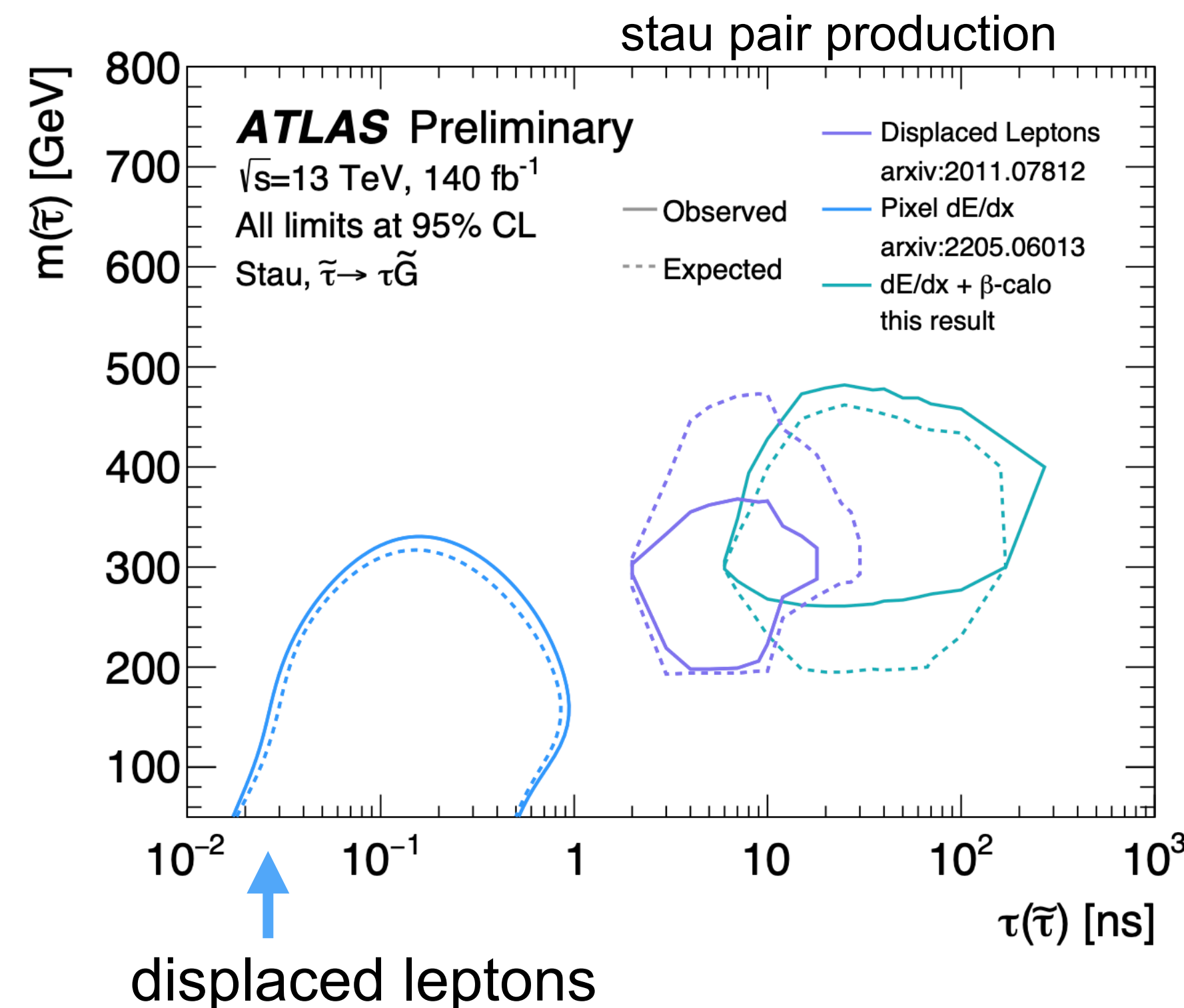


- **complements previous analysis** [arXiv:2022.06013](https://arxiv.org/abs/2022.06013) using large E depositions in pixel detector by including calorimeter ToF info
- targeting massive slow particles with high charge



- previous analysis $\rightarrow 3.3 \sigma$ excess $m \sim 1.4$ TeV not confirmed to be due to slow high-mass particles by ToF

- **no significant excess: 6 data vs 3.7 bkg events**

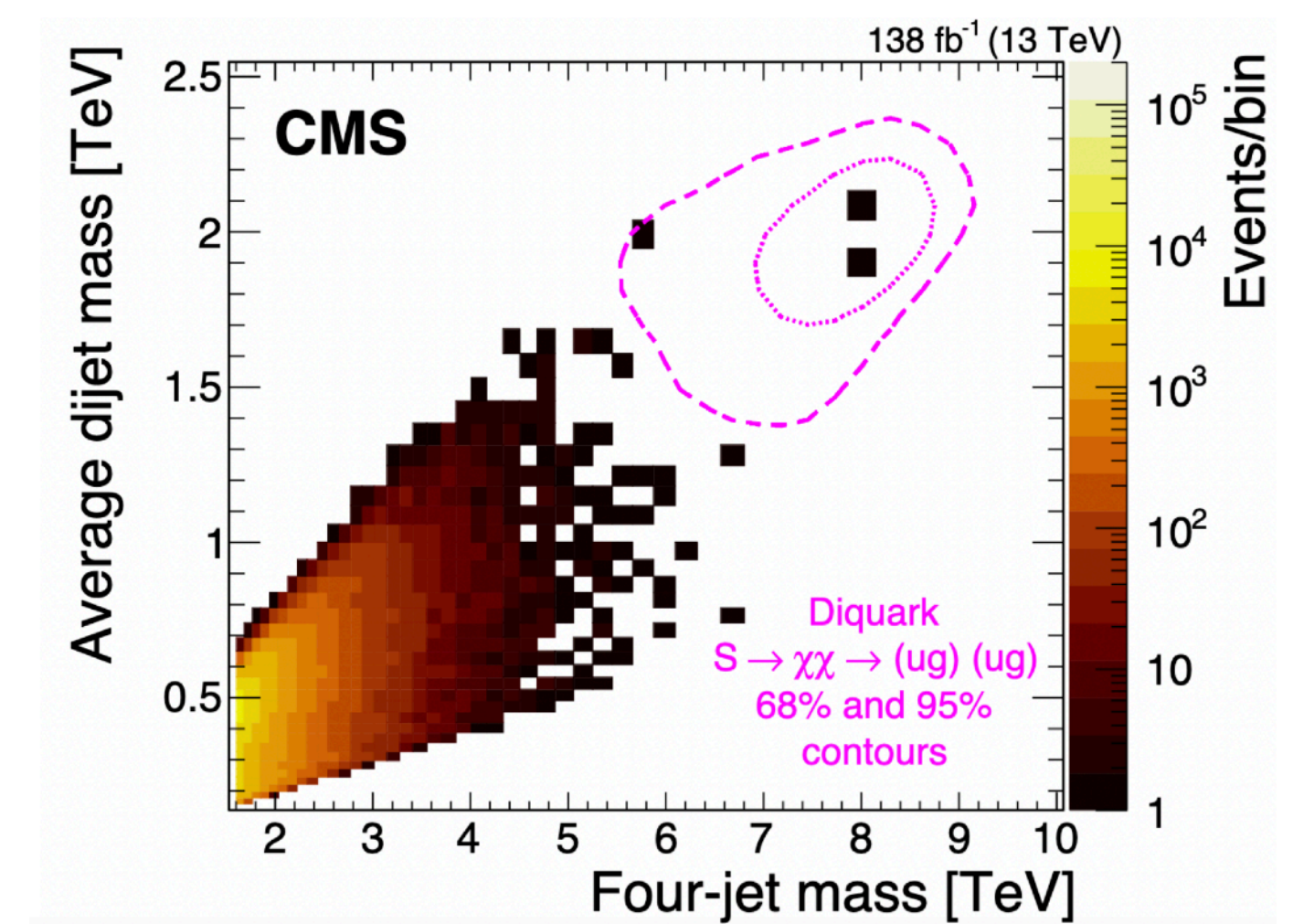
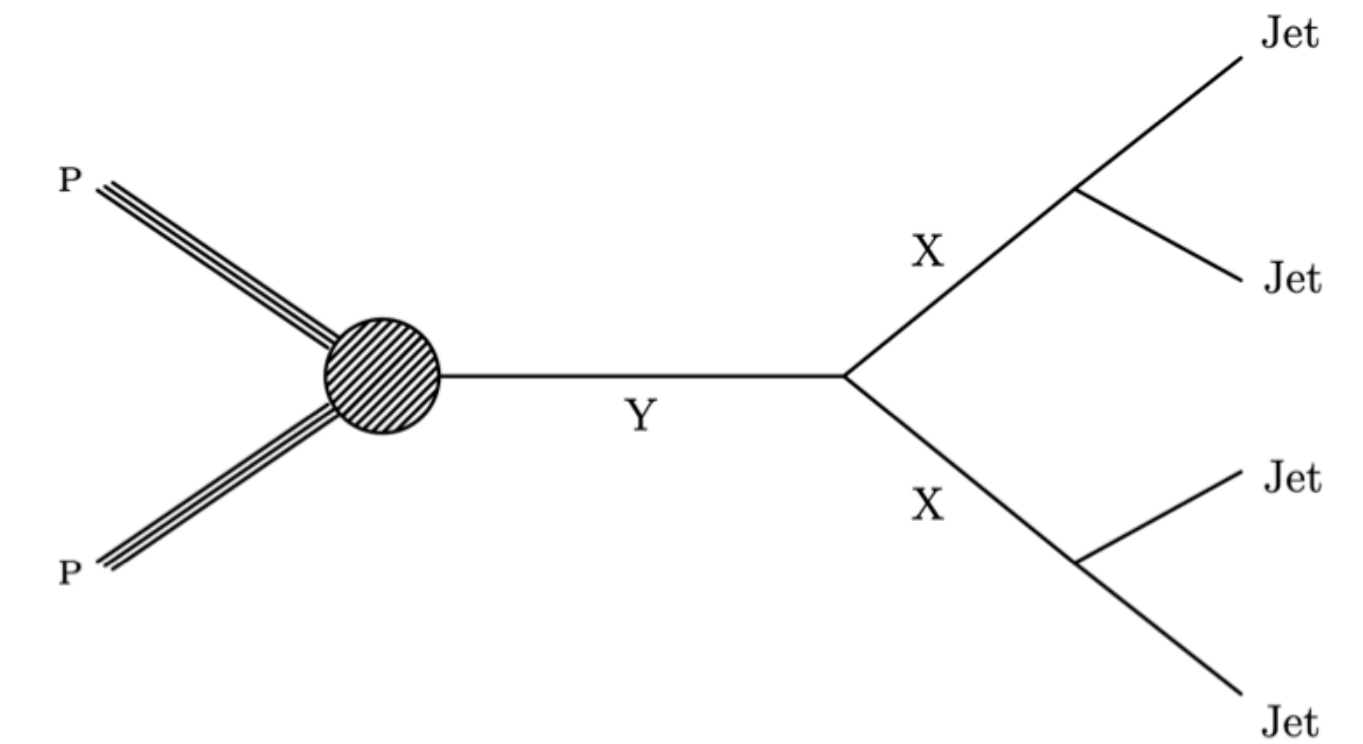
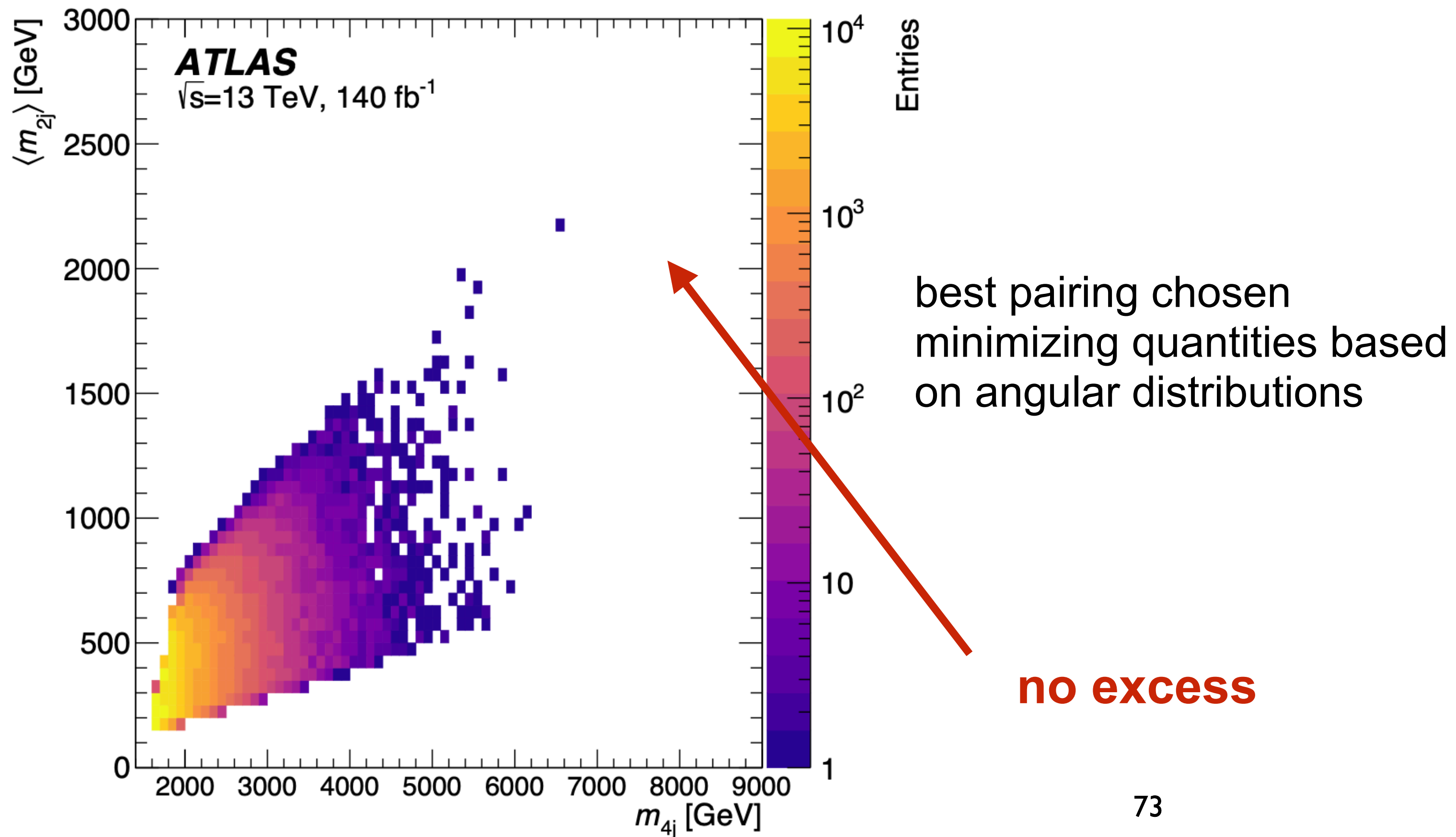


Resonances

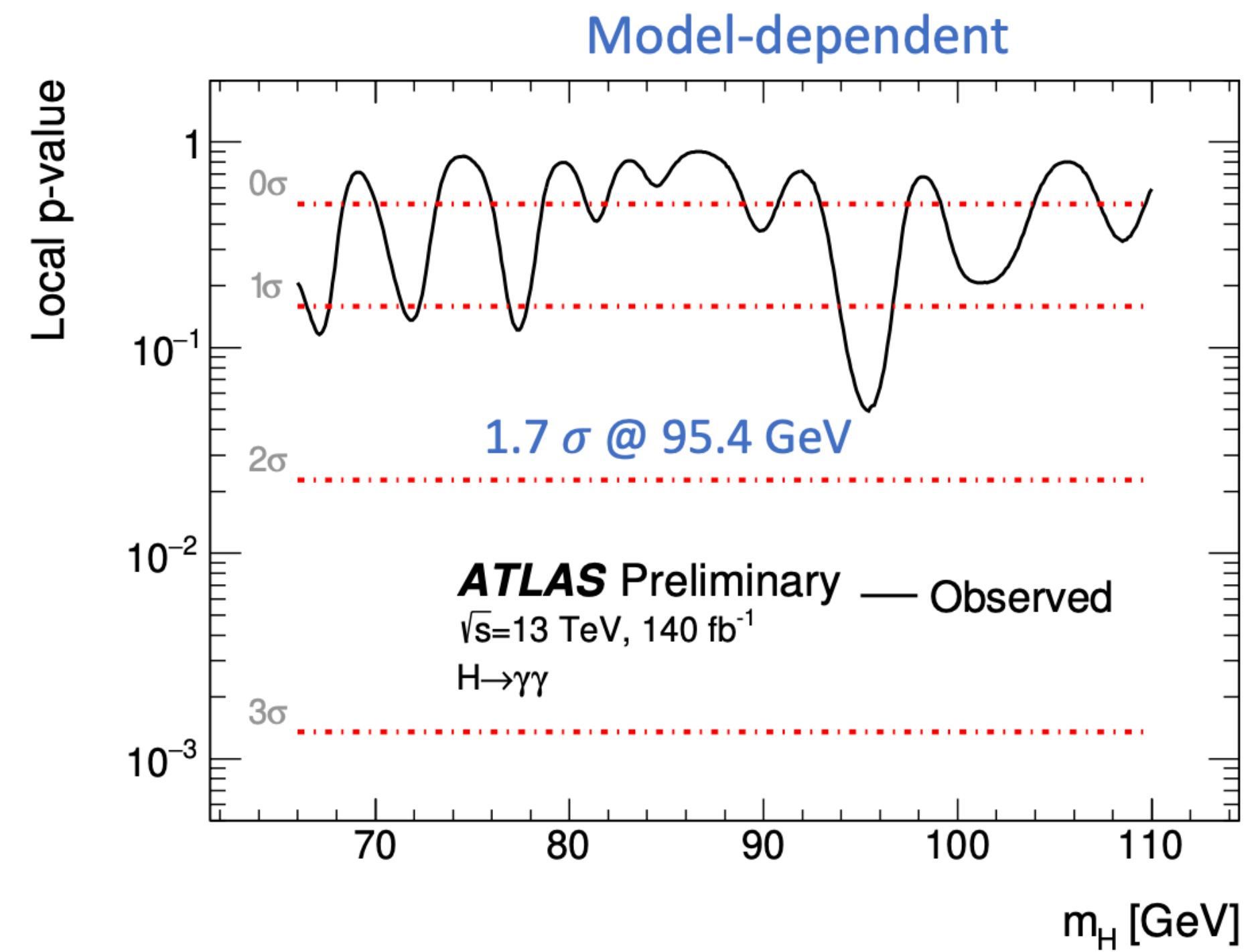
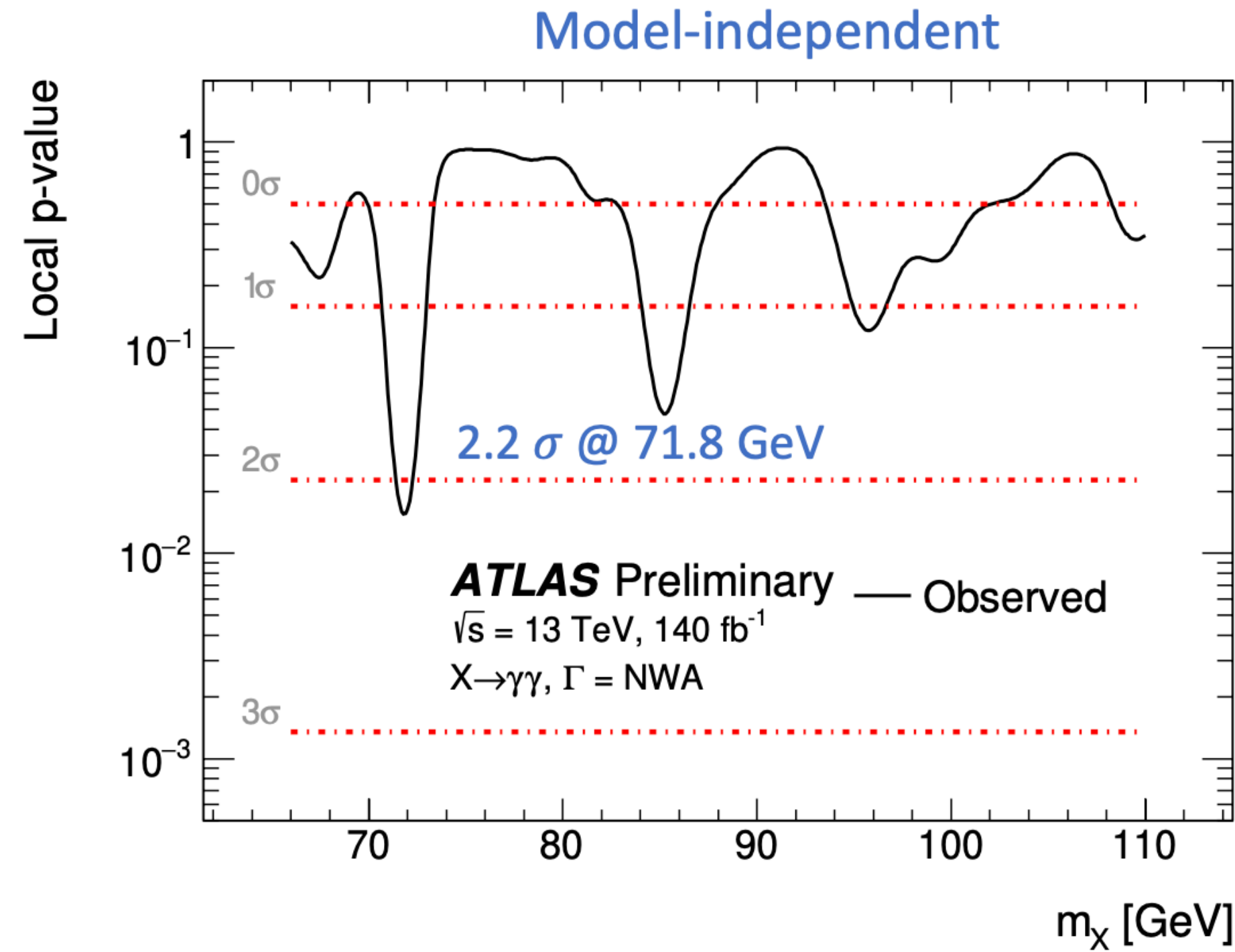
Resonances: Tetrajets $Y \rightarrow XX \rightarrow jjjj$

[arXiv:2307.14944](https://arxiv.org/abs/2307.14944)

- Search for generic massive resonance Y decaying to intermediate resonances X
- Bumphunter search in m_{4j} & di-jet average inv. mass $\langle m_{2j} \rangle$
- Follow up on 3.6σ CMS excess paired dijets [[arXiv: 2206.09997](https://arxiv.org/abs/2206.09997)]



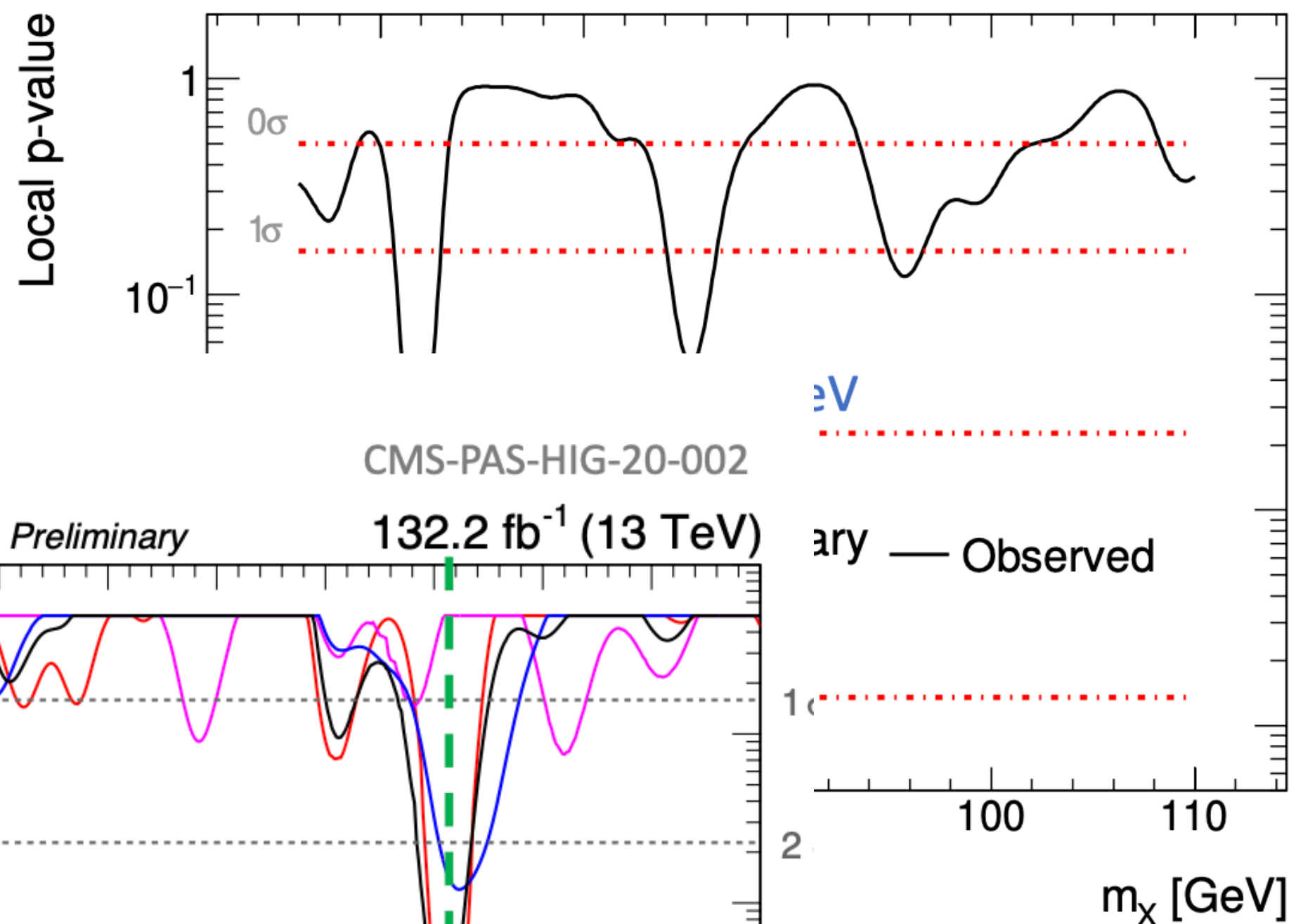
Low mass $h \rightarrow \gamma\gamma$



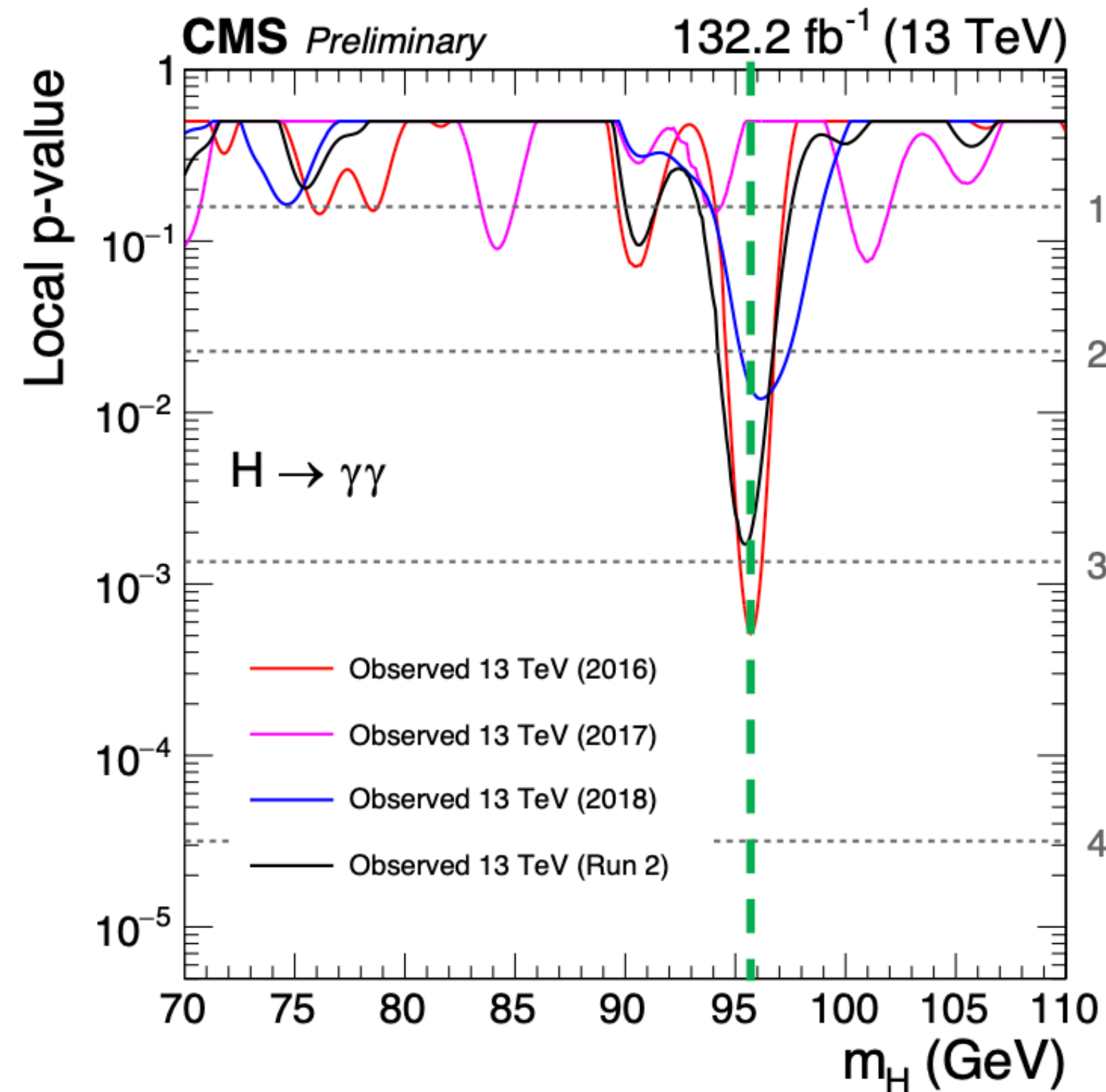
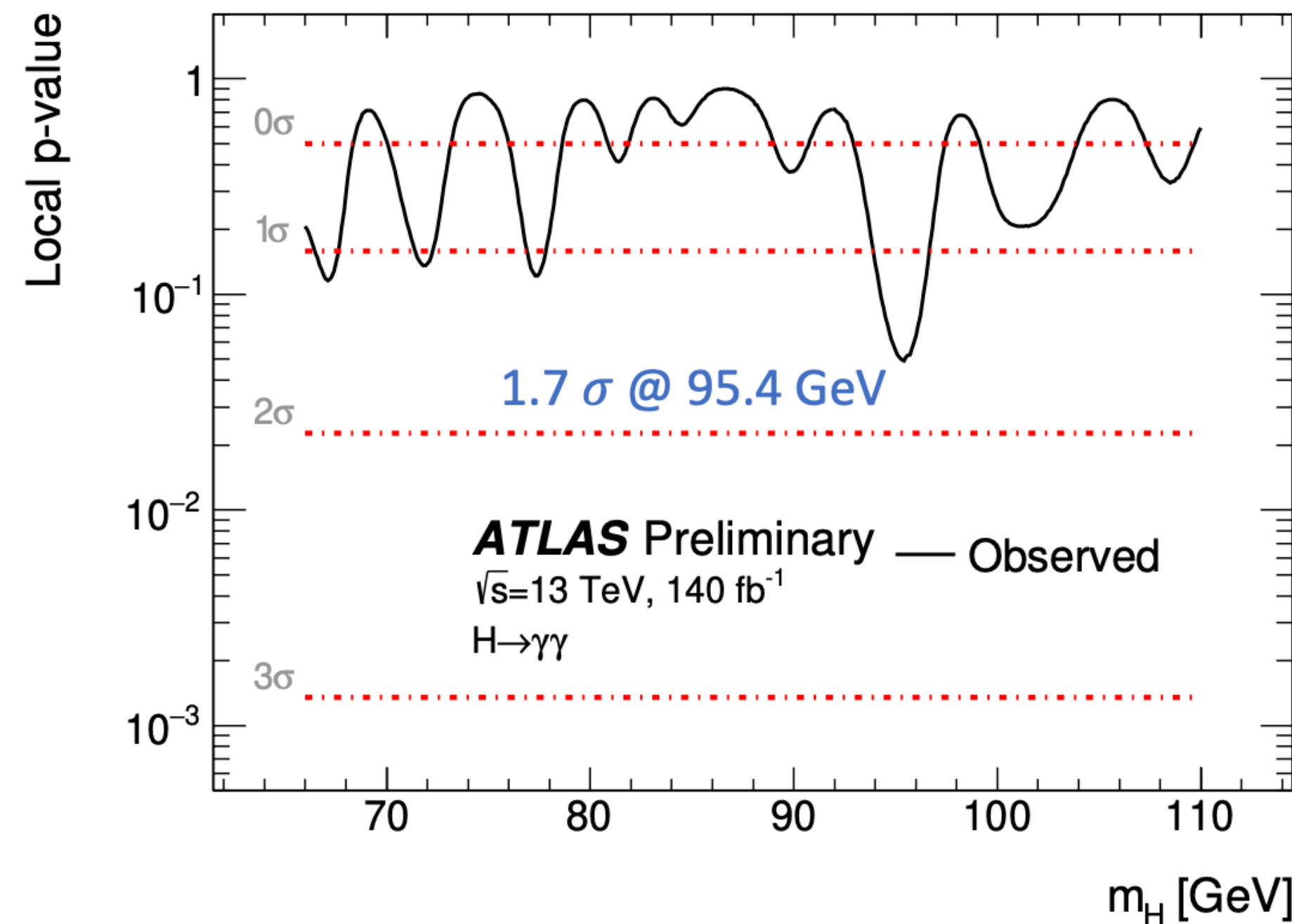
Low mass Higgs resonance
could come from Axion like
particles in SUSY, or from
2HDMS

Low mass $h \rightarrow \gamma\gamma$

Model-independent



Model-dependent



best pairing chosen
minimizing quantities based
on angular distributions

Conclusions

- There are many exciting searches and results behind the corner
- The efforts in improving the trigger, the analysis techniques are paying off.
- This shows that much more can be achieved, especially by exploiting even more the data
- as a tool to control or systematic errors
- Theory and experiment also have to go hand in hand
- We have great chances at Run 3 and the HL-LHC to find the clue that will lighten up the path for the search of new physics



Back-up

Several recording strategies to circumvent limitations in practice

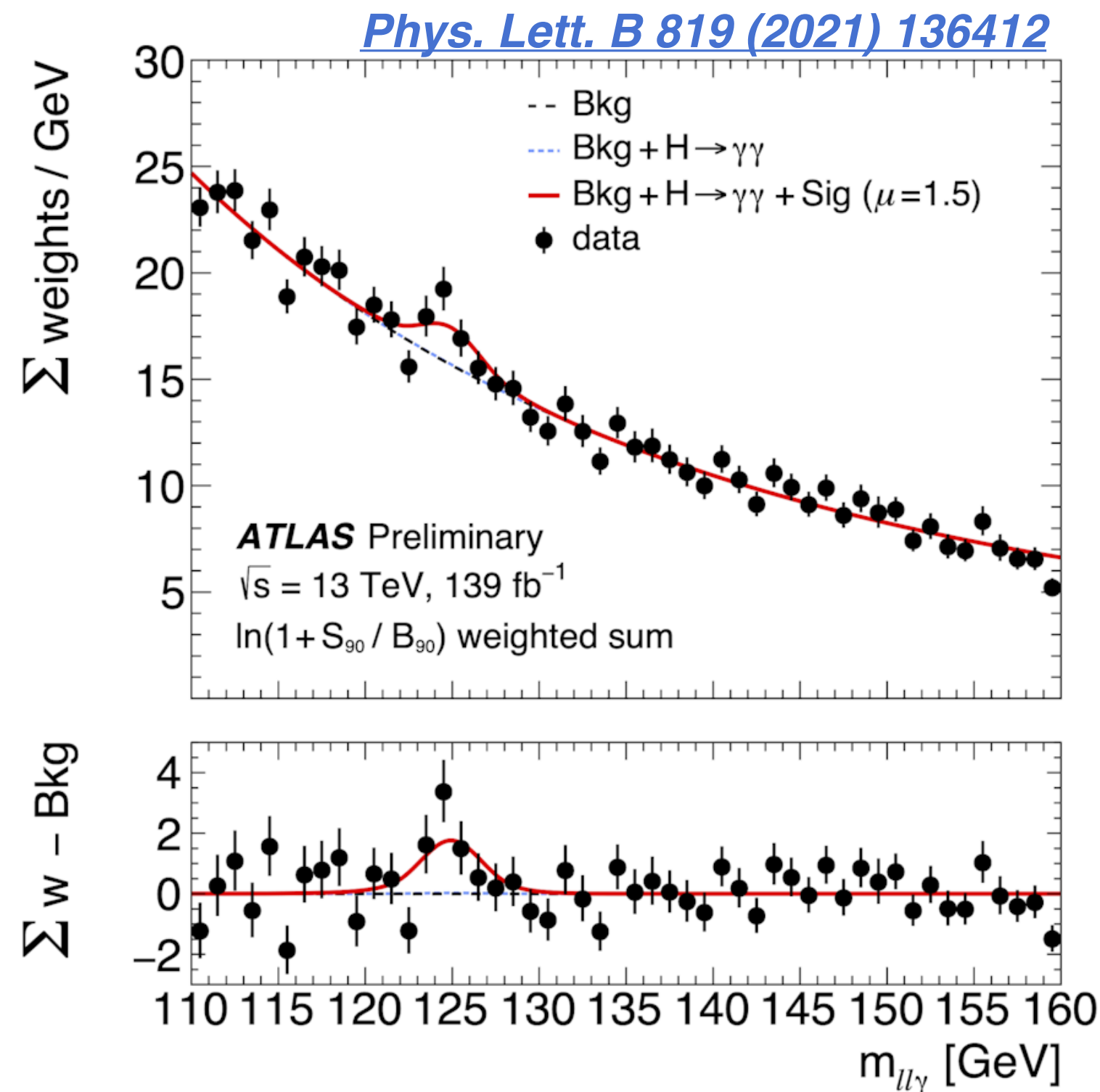
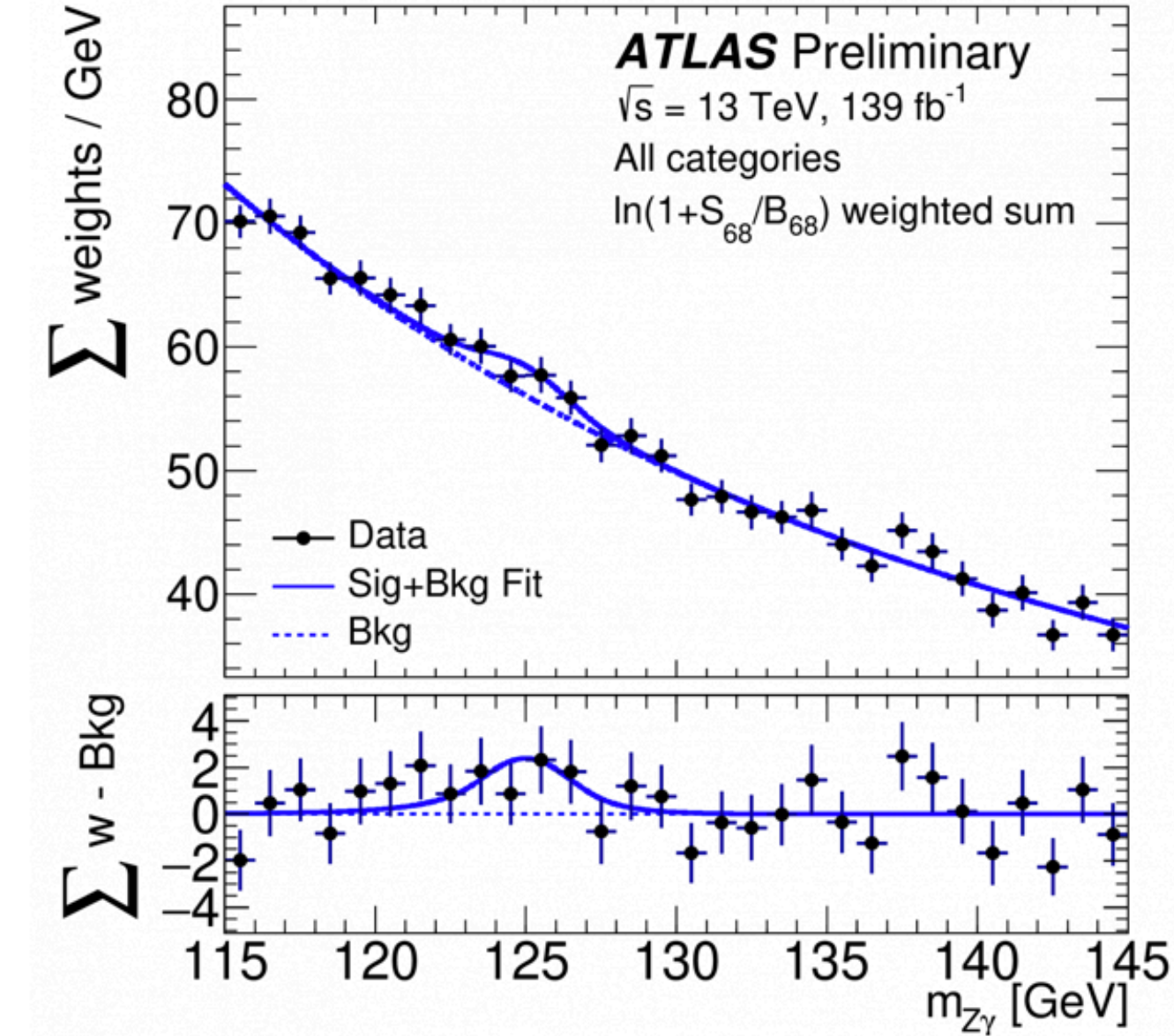
- **New Trigger:** Develop new trigger logic to enhance selection of your signature, sufficiently selective for minimal background (\downarrow Bandwidth = \downarrow Rate x Size).
Advantage: Full detector information, unique data.
- **Delayed Stream Strategy:** Store full event data on SFOs (storage at Point 1) to reduce TIER0 bottleneck.
Advantage: Full detector information
- **Trigger(-object) Level Analysis (TLA):** Reduce event size 100x by recording only HLT reconstructed objects - take advantage of offline-like reco algorithms at HLT (\downarrow Bandwidth = \uparrow Rate x $\downarrow\downarrow$ Size).
Advantage: Trigger thresholds no longer limited by HLT thresholds.
- **Partial Event Building (PEB):** Reduce event size by recording only raw detector data in Regions Of Interest (\downarrow Bandwidth = \uparrow Rate x $\downarrow\downarrow$ Size, $\downarrow\downarrow\downarrow$ CPU).
Advantage: CPU limitations lifted.

H → Zγ and γγ*

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The SM predicts 0.15% of Higgs to decay to Zγ comparable to the decay to two photons (Z BR bosons decay to leptons) makes this more challenging.

- significance of 2.2σ obs (1.2σ exp)
- 95%CL upper limit at 3.6xSM obs (2.6xSM exp)



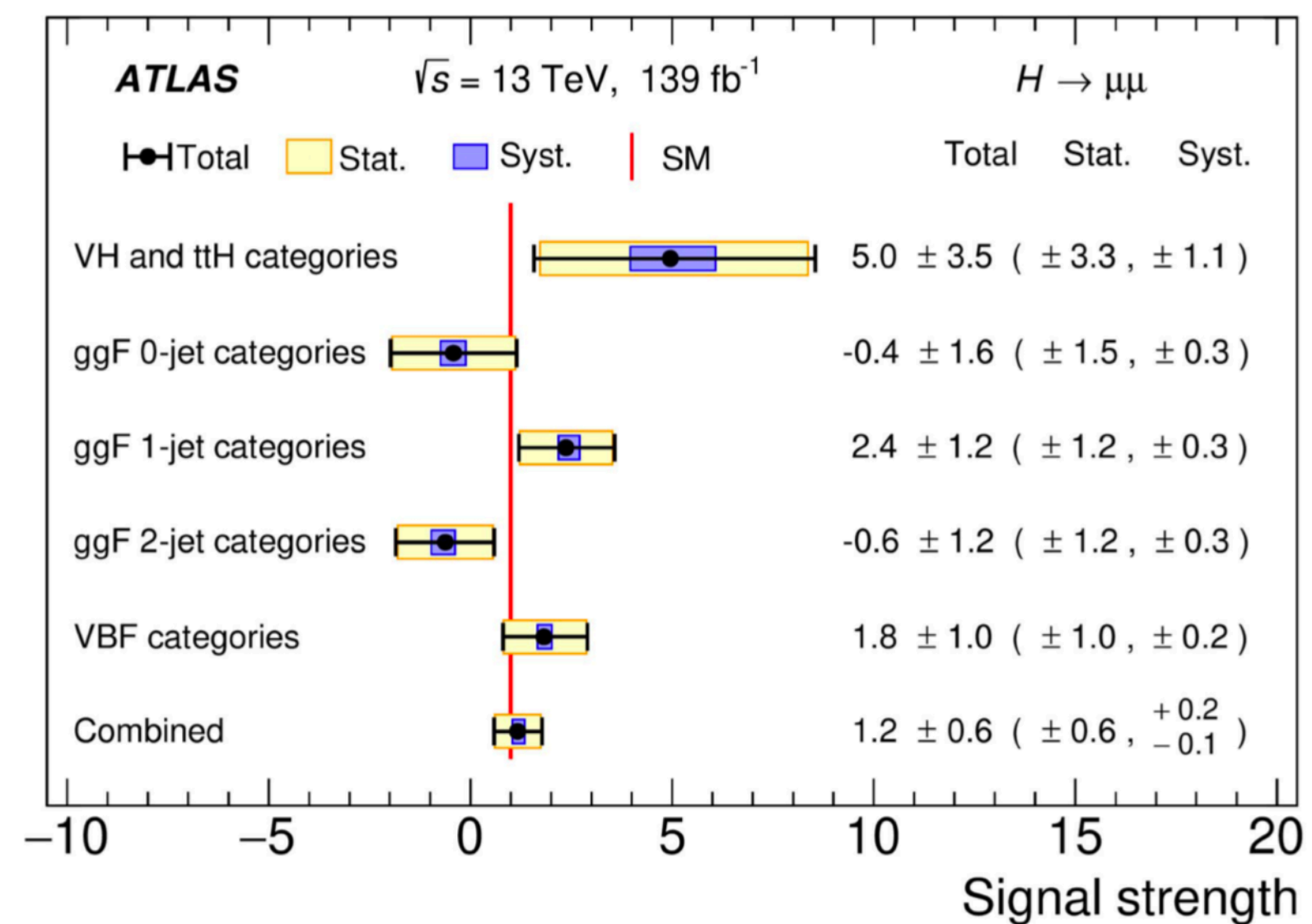
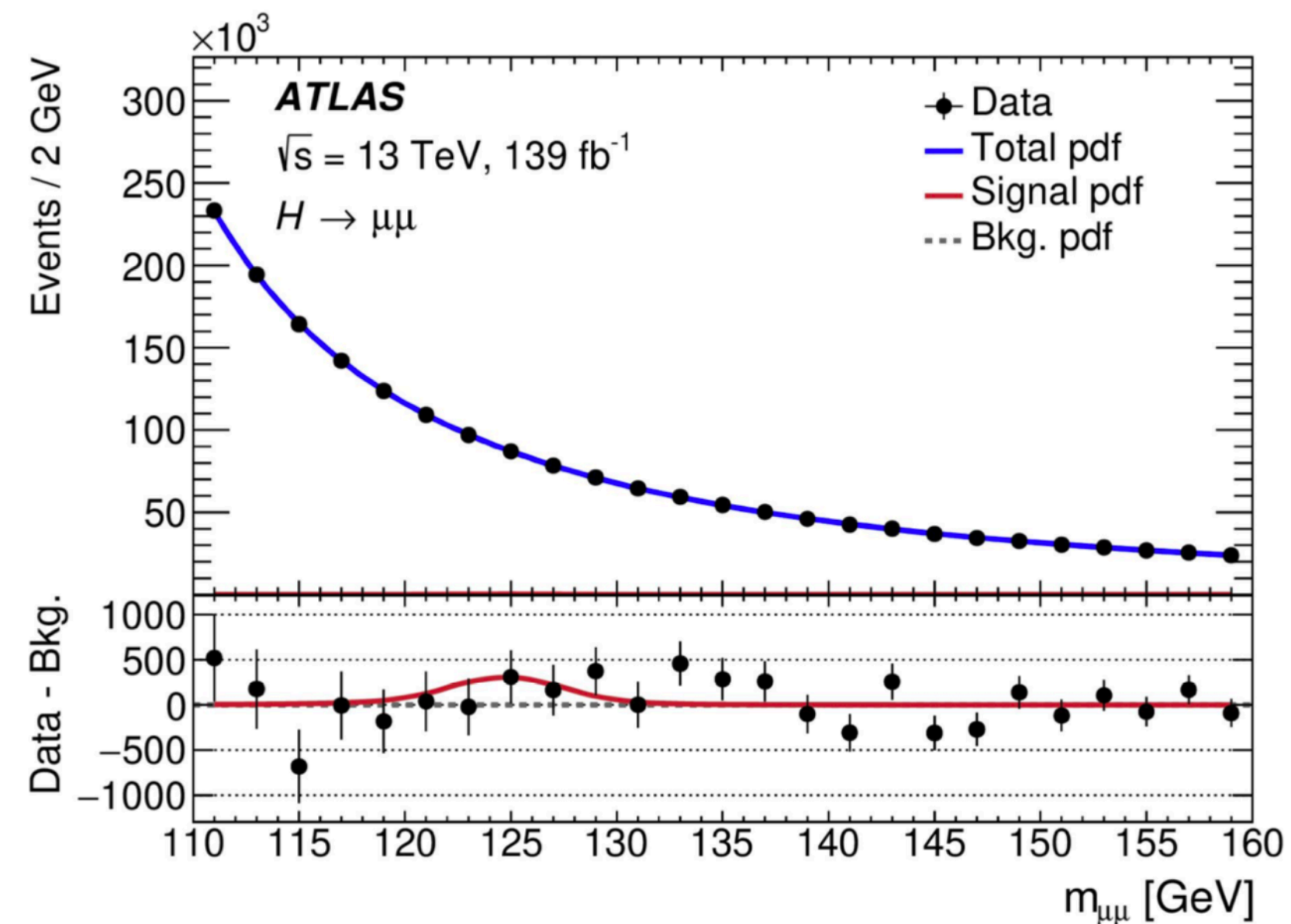
γ* is a virtual particle with (non zero) mass, decays instantly to two leptons
 $m_{\ell\ell} < 30$ GeV (typically < 1 GeV)
 high pT, small leptons separation ~cm
 (challenge for electrons)

significance of 3.2 σ obs (2.1 σ exp)

$H \rightarrow \mu\mu$

$Z \rightarrow \mu\mu$ main background
statistically limited

VBF category is the most powerful!



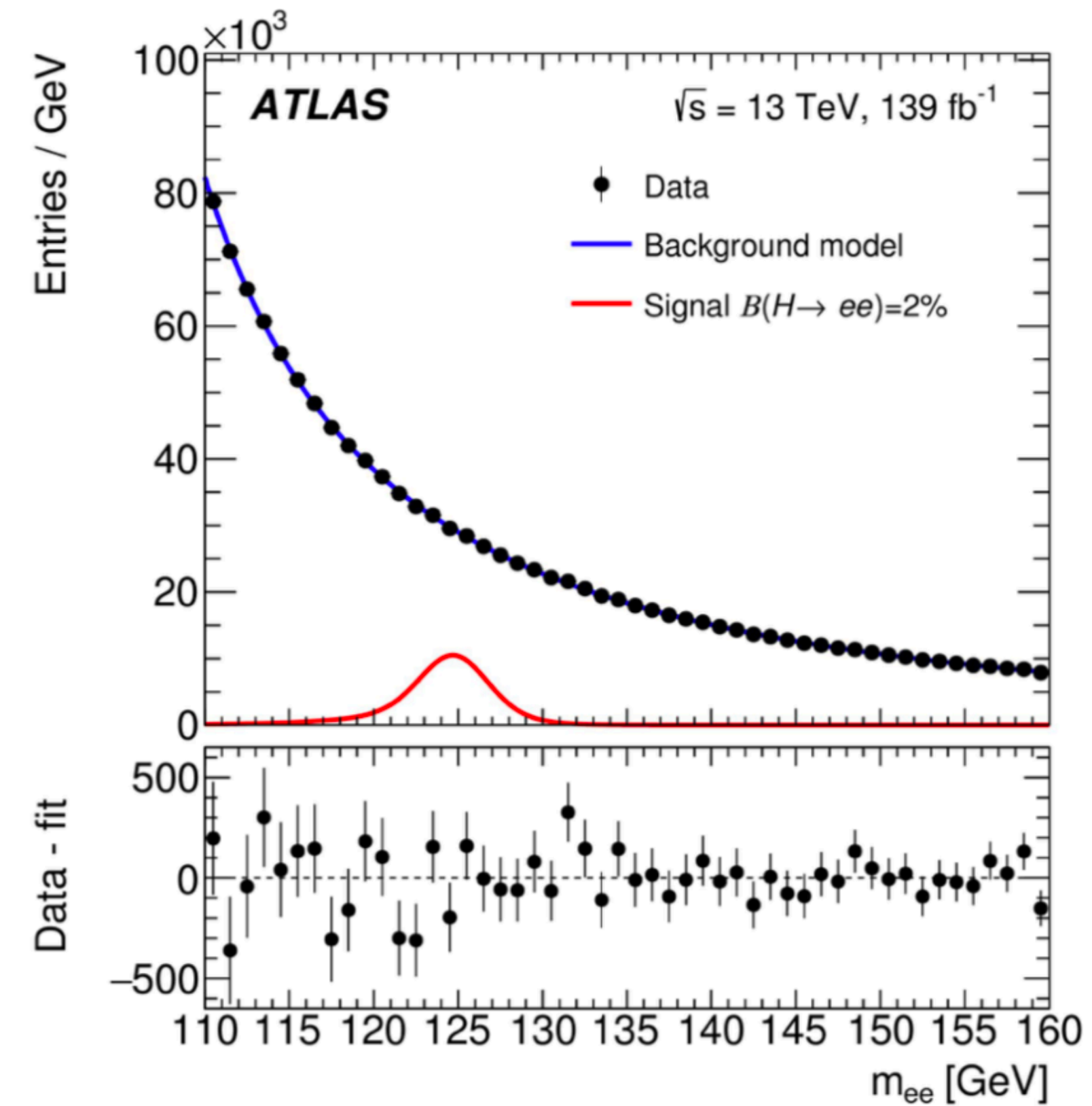
Observed (expected) significance of 2.0 (1.7) σ

Important in run 3: in reach 3 sigma
per experiment and 5 in combination.

H → ee

Z → ee main background
statistically limited
similar analysis strategy as H → μμ

Observed (expected) limit at 95% CL:
 $BR_{H \rightarrow ee} < 3.6 (3.5) \times 10^{-4}$

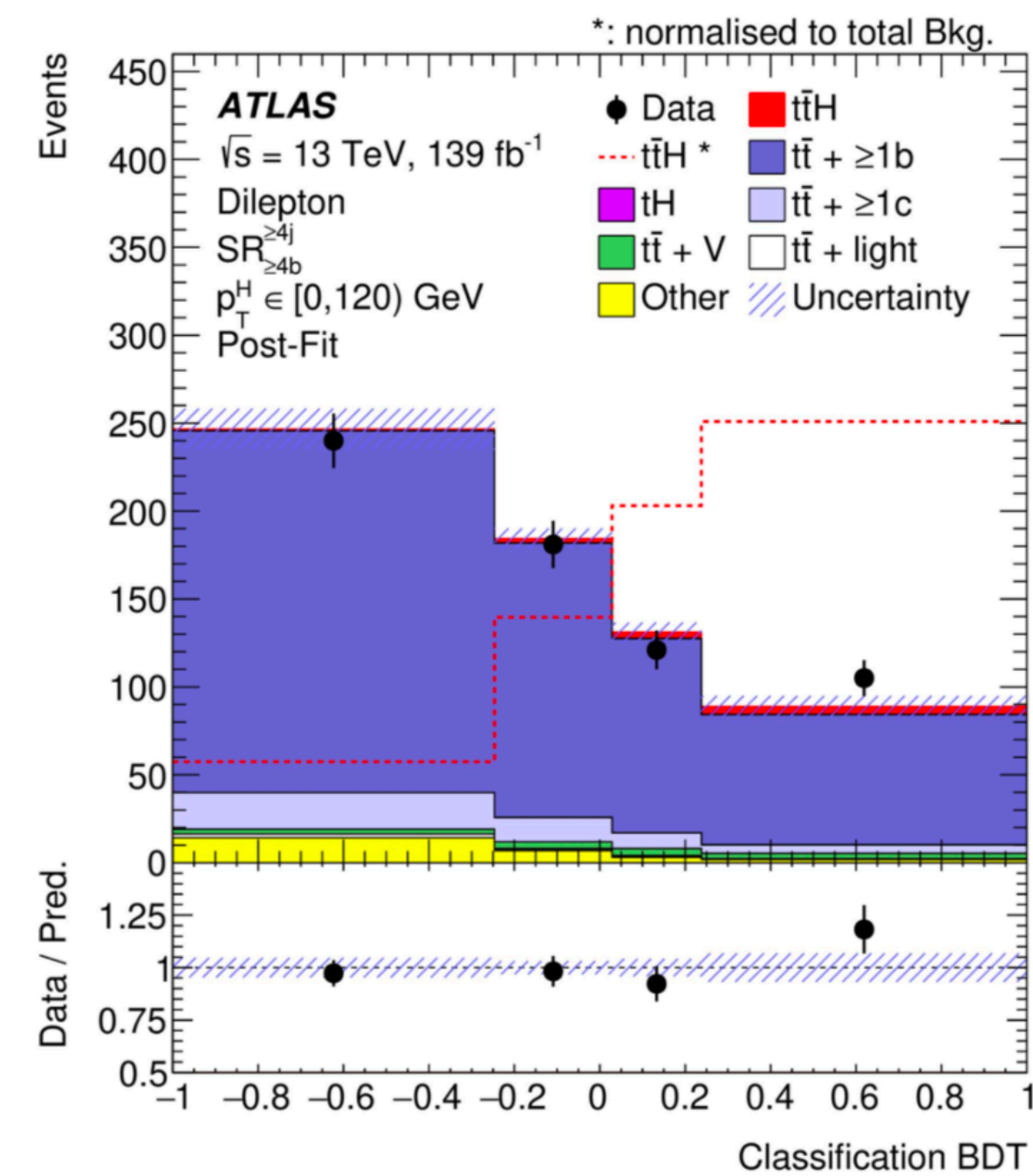
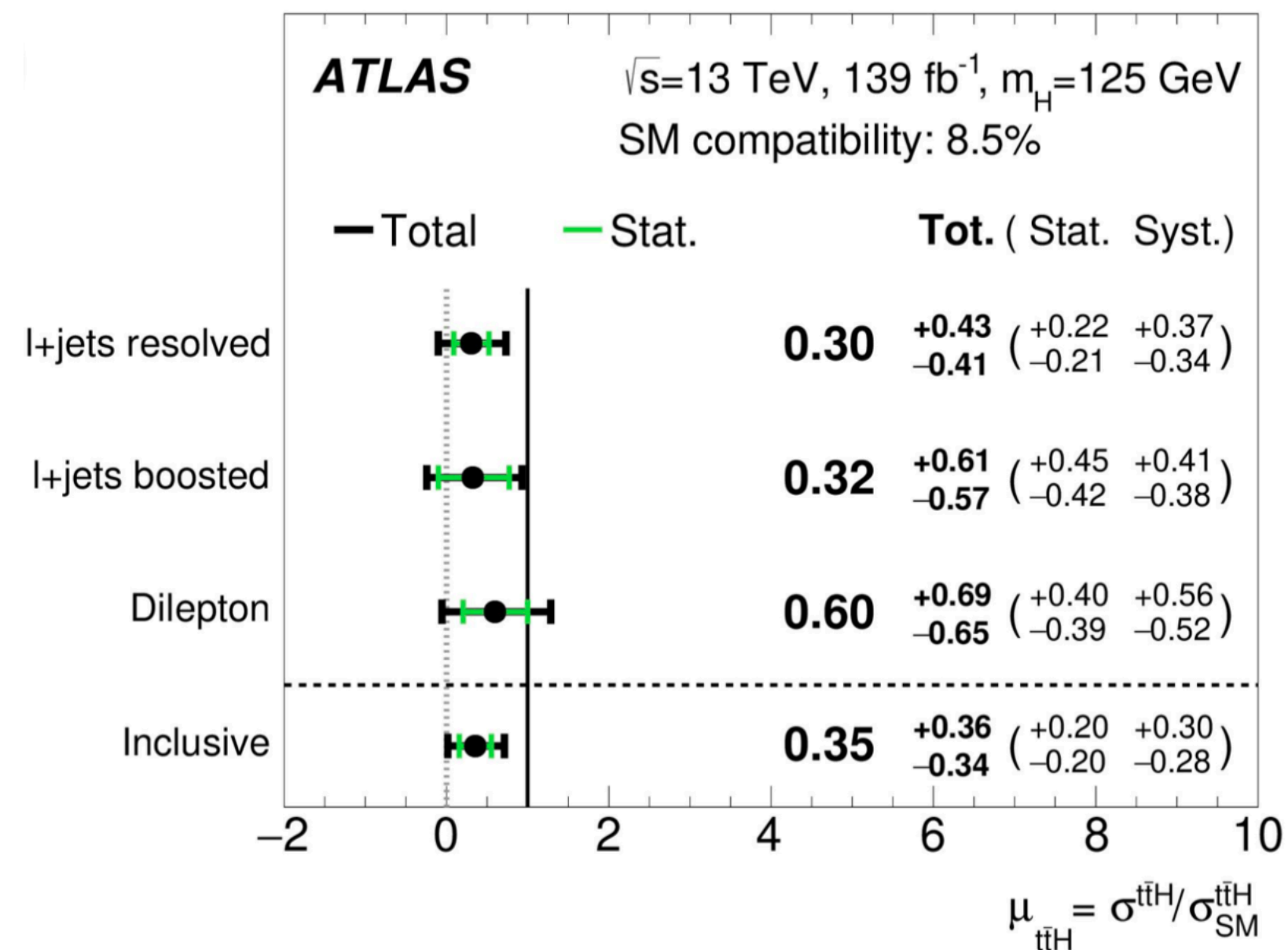


the Higgs boson is around 40,000 times less likely to decay into electrons as it is into muons

$t\bar{t}H(bb)$

- The top Yukawa coupling probed with this channel, 1 or both tops decaying leptonically.
- Events classified according to the number of leptons, jets and b-jets
- Machine learning techniques to aid the signal/background discrimination

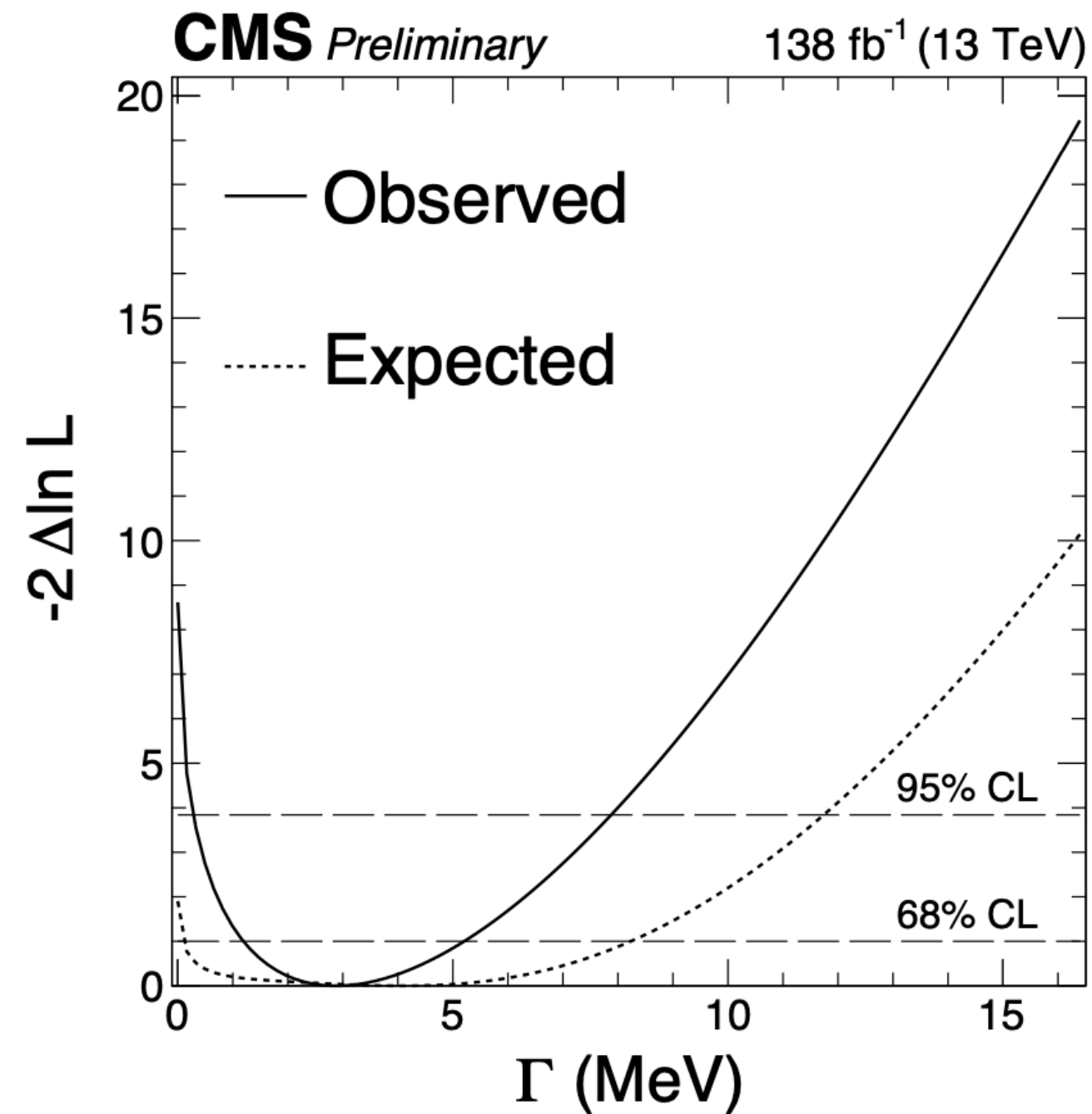
- significance of 1σ obs (2.1σ exp)



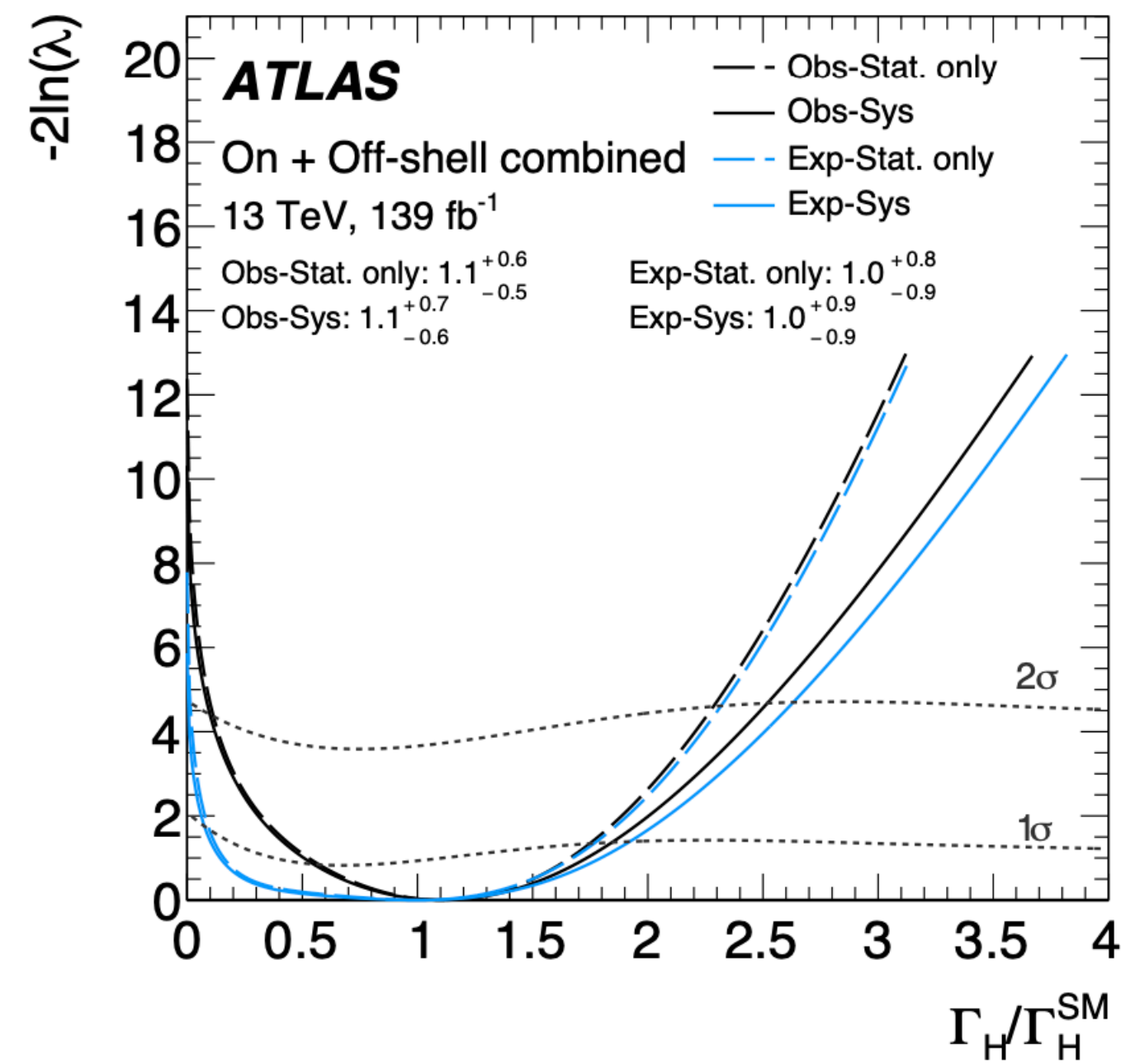


Higgs width

The best estimate of the Higgs boson total width is: $\Gamma_{HSM} = 4.07$ MeV
3 orders of magnitude smaller than our mass resolution



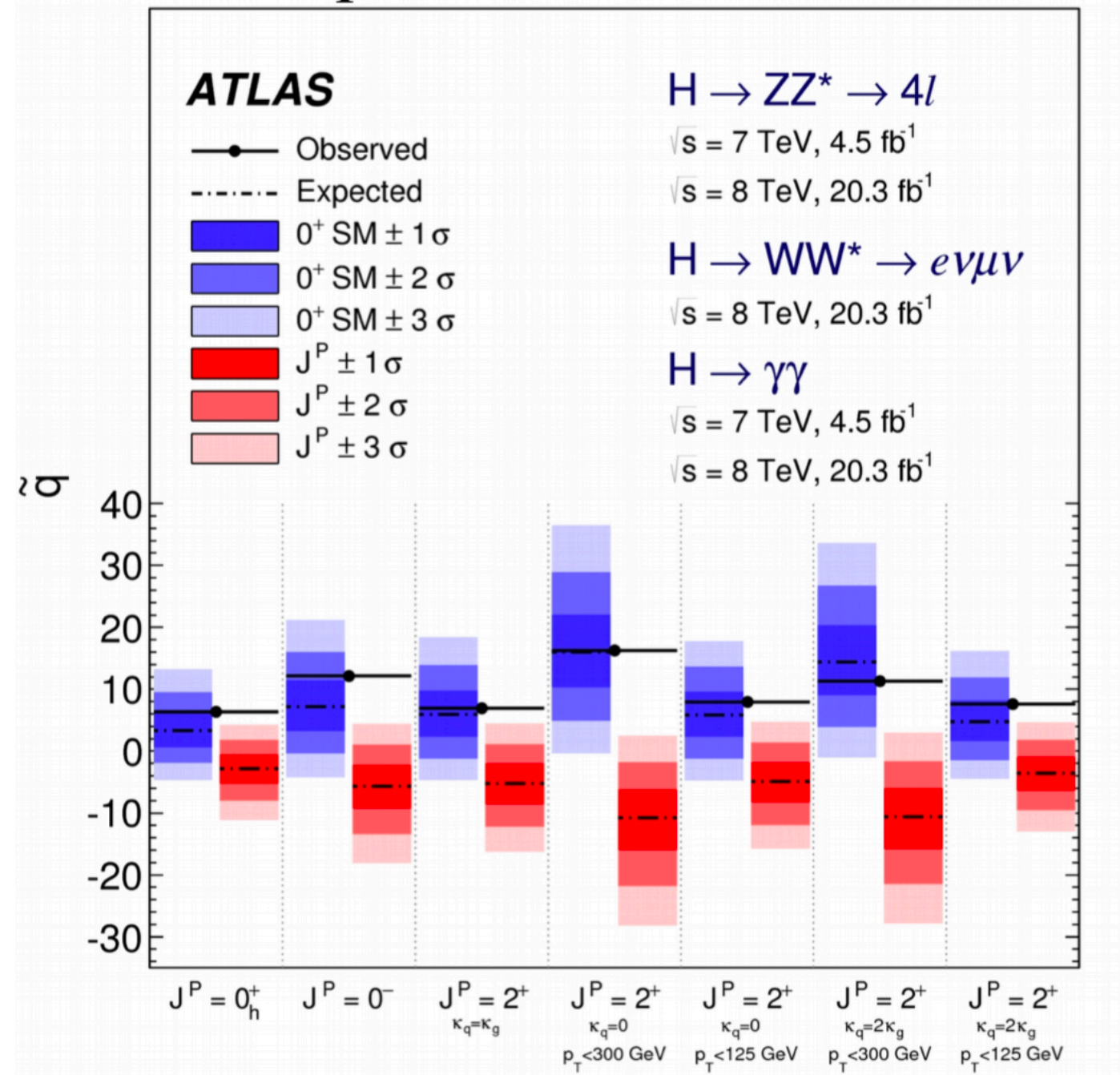
Extracted width: $\Gamma_H = 2.9^{+2.3}_{-1.7}$ MeV



$\Gamma_H = 4.5^{+3.3}_{-2.5}$ MeV

◆ Spin/parity: $J^{PC} = 0^{++}$

– spin 1 and 2 excluded at > 99% CL



In bosonic couplings parametrized with higher order terms suppressed by powers of Λ (scale of new physics)

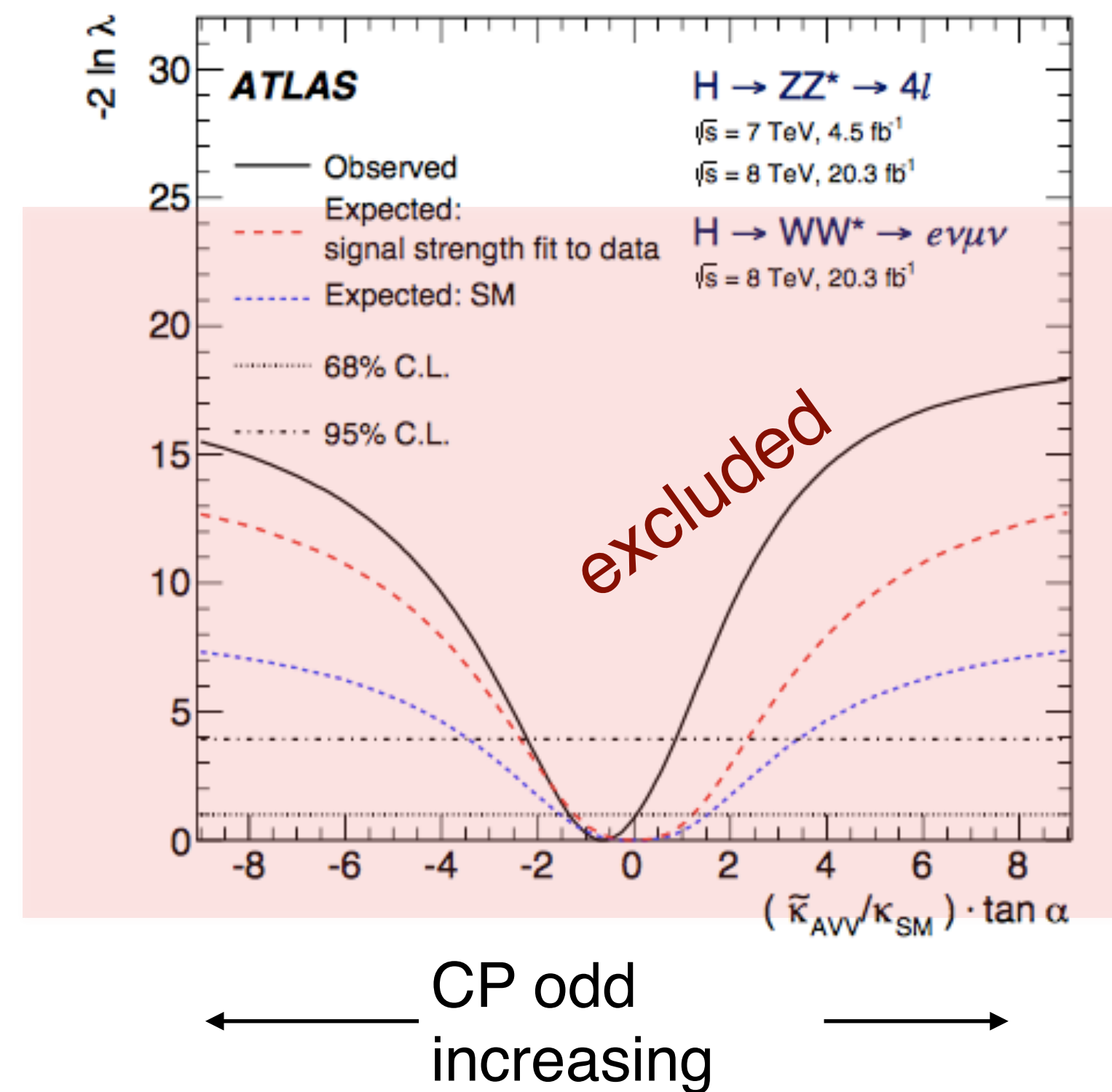
$$\mathcal{L}_{VVH} = \mathcal{L}_{VVH,SM} + \frac{1}{\Lambda^2} c \phi \tilde{V}_{\mu\nu} V^{\mu\nu} + \dots$$

CP violation

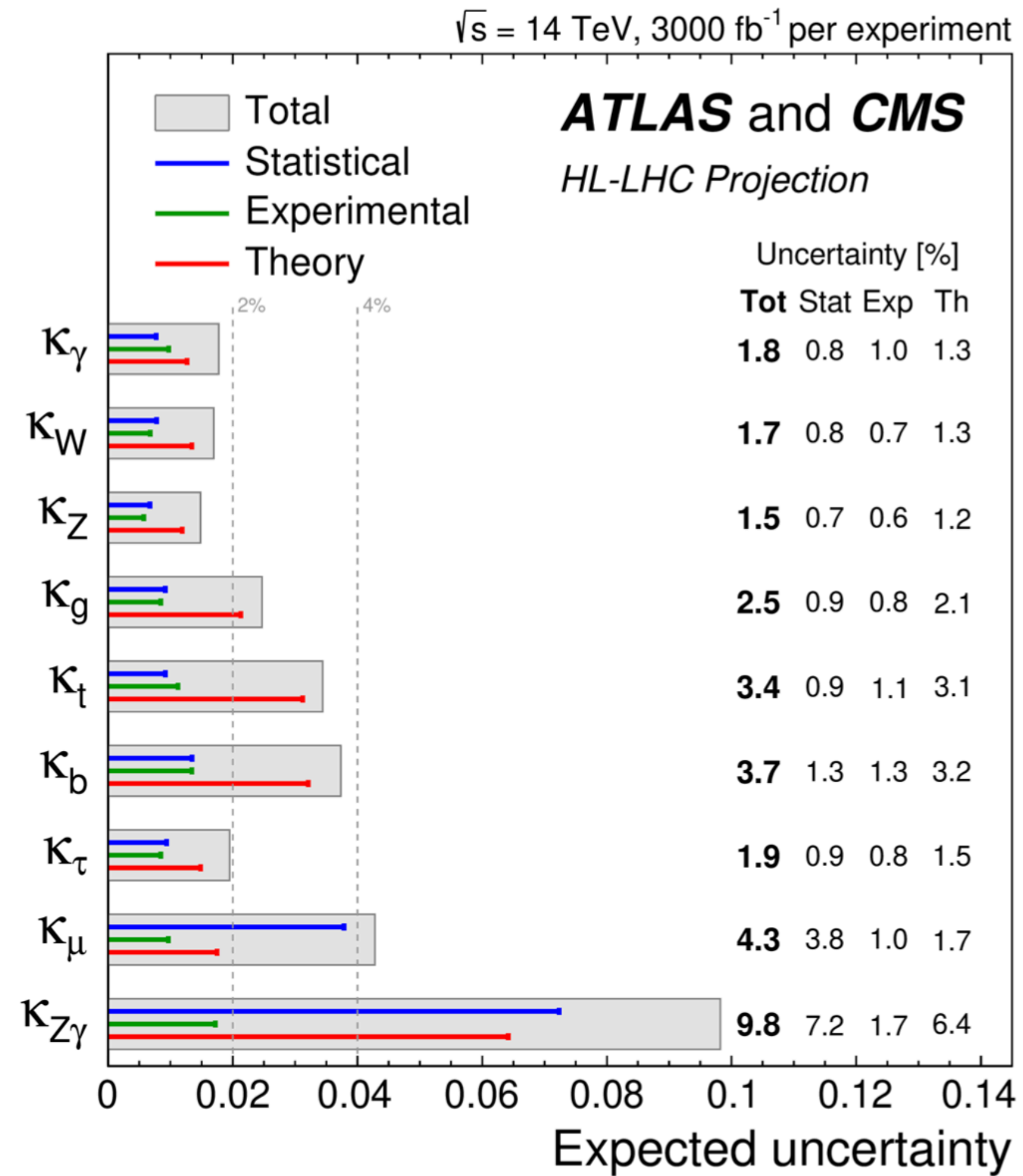
Fermionic couplings affected at tree level (more important for heavier fermions due to higher coupling)

α CP- even and CP-odd mixing angle

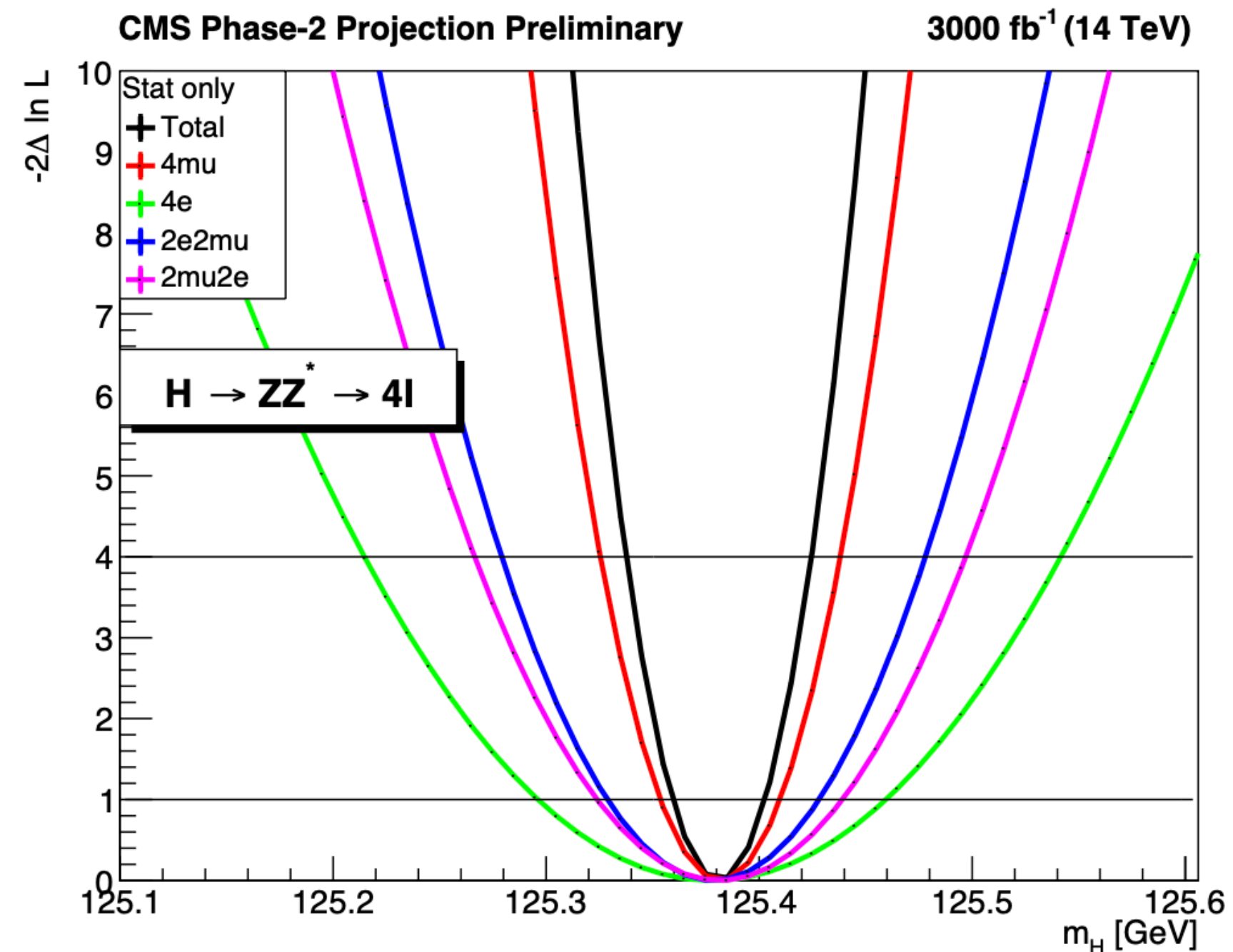
$$\mathcal{L}_{ffH} = \kappa'_f y_f \phi \bar{\psi}_f (\cos \alpha + i \gamma_5 \sin \alpha) \psi_f$$



HL-LHC projections



The projected expected result, for mass measurement, is $m_H = 125.38 \pm 0.03[0.022(\text{stat}) \pm 0.020(\text{syst})]$ GeV and for width is $\Gamma_H < 0.09(0.18)$ GeV at 68(95)% confidence level

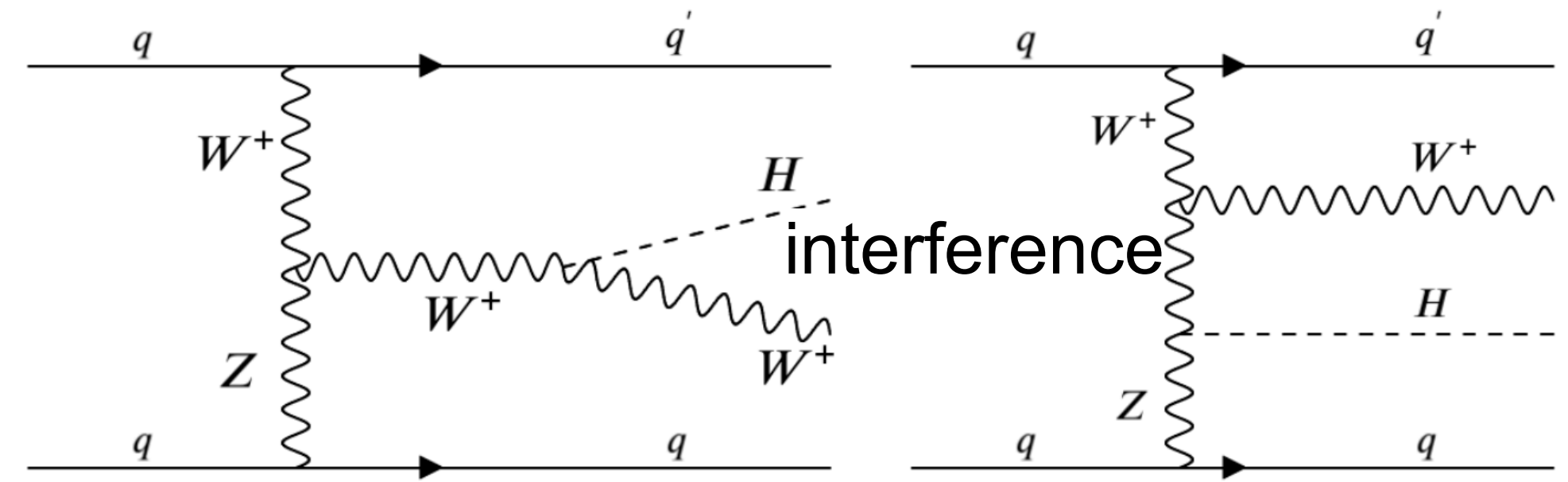


Higgs VBF, WH production sign of λ_{WZ} coupling

ATLAS-CONF-2023-057

New!

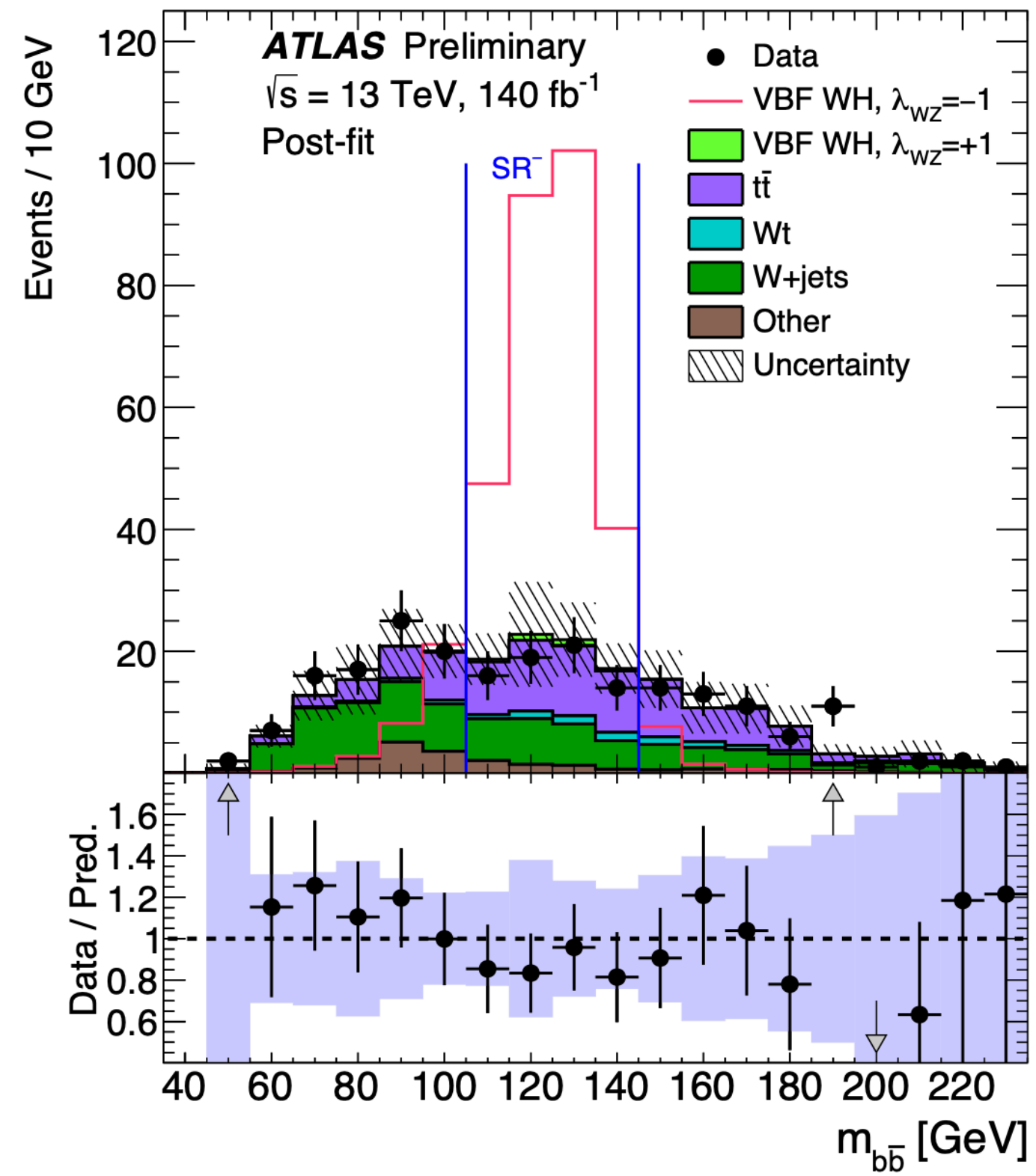
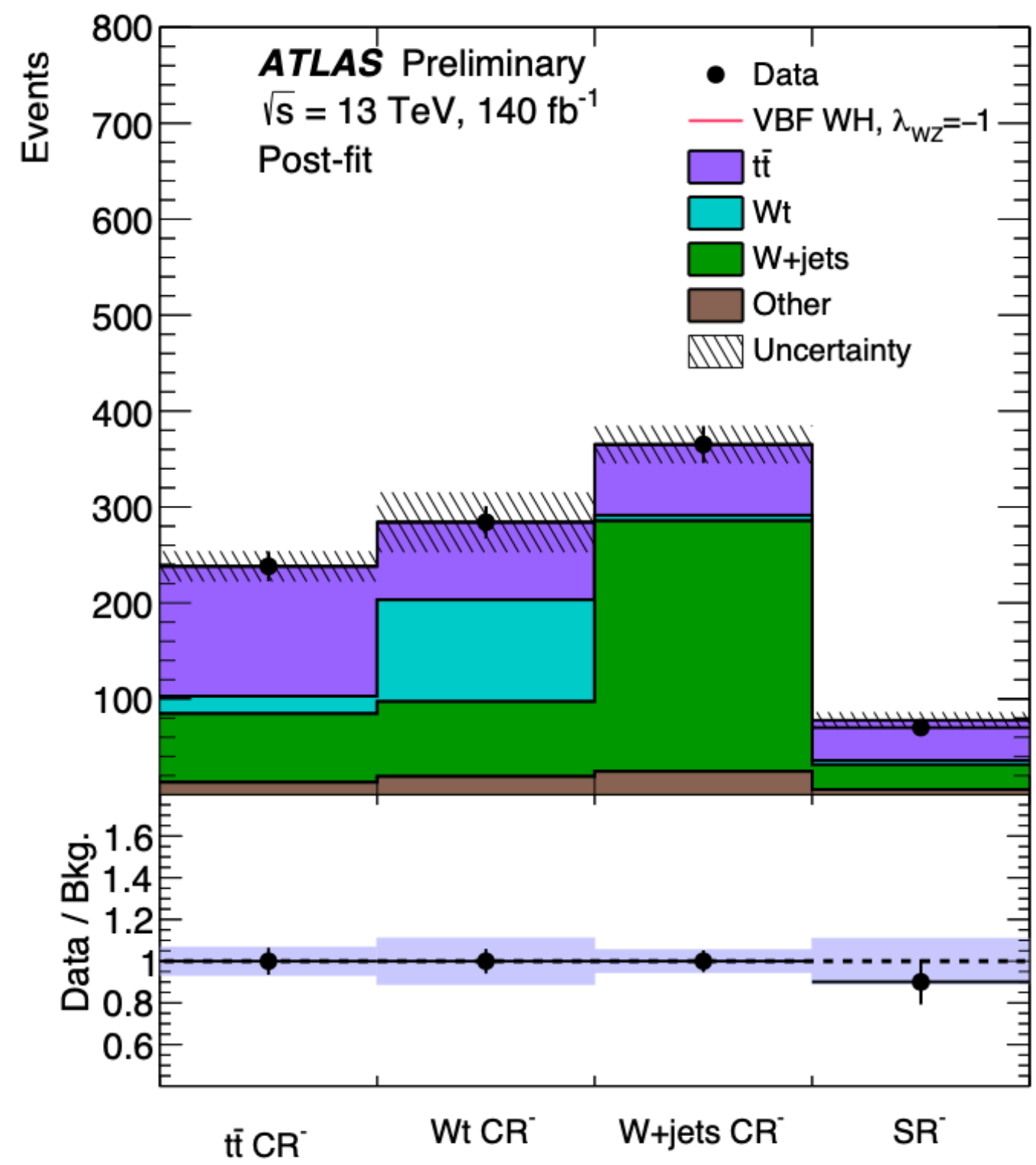
- New production mode: sensitive to sign of ratio of Higgs-W/Z couplings ($\lambda_{WZ} = \kappa_W / \kappa_Z$), ($H \rightarrow b\bar{b}$) for which we had no sensitivity before



- Excluded $\lambda_{WZ} = -1$ at $>8\sigma$

- Measure μ for $+\lambda_{WZ}$ signal

Fit: $\hat{\mu} = 2.6^{+4.6}_{-4.5}$



Events / 0.35