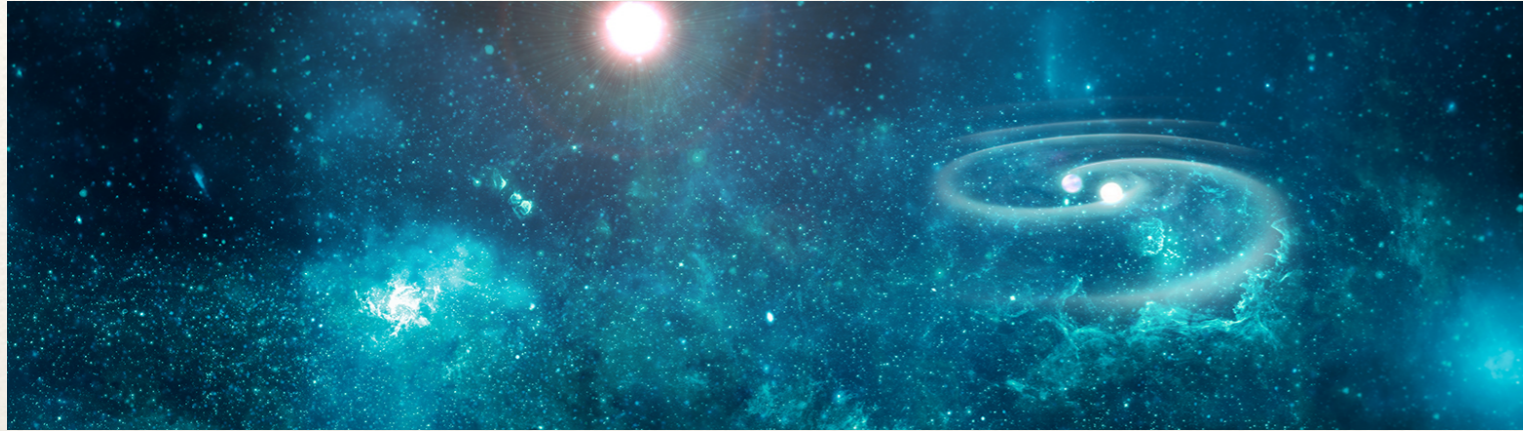


Stanislav (Stas) Babak.

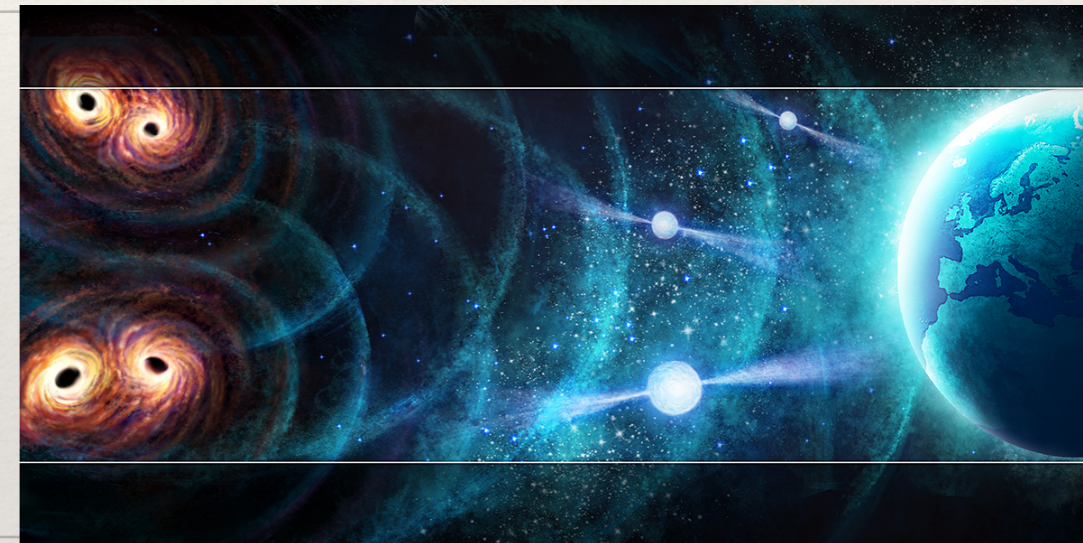
AstroParticule et Cosmologie, CNRS (Paris)



Université
Paris Cité



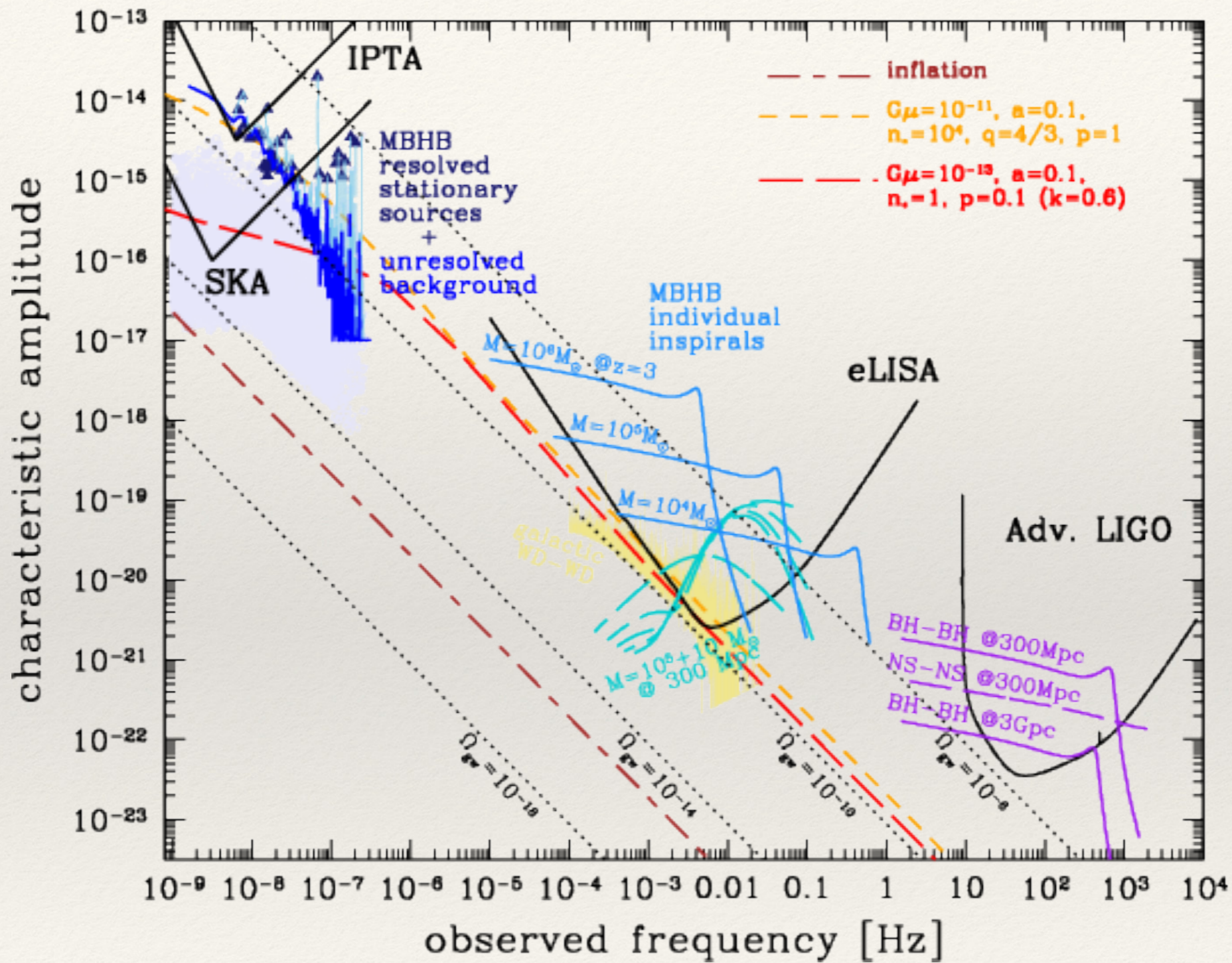
Pulsar Timing Array



CERN 19-21 July 2023



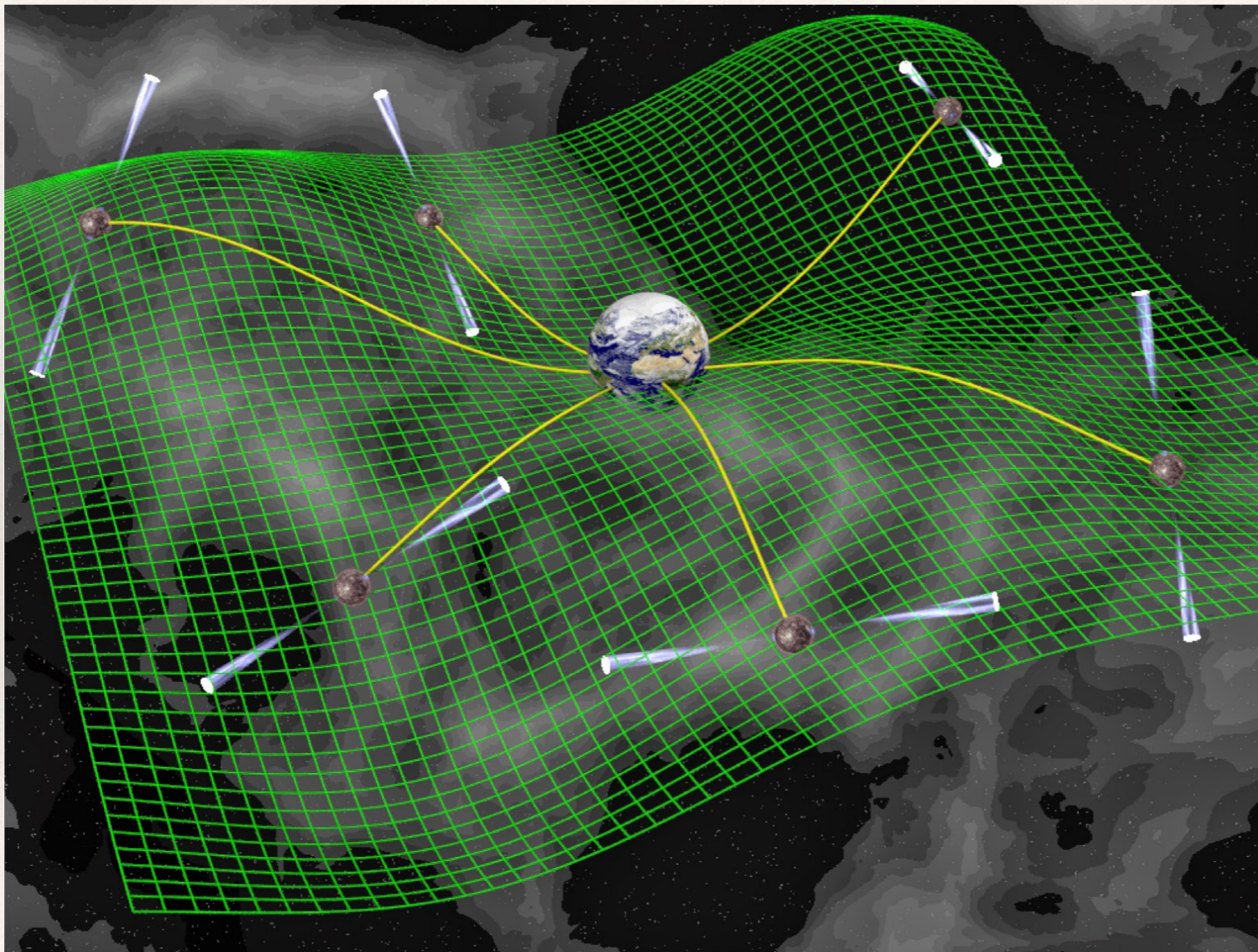
GW landscape



Pulsar Timing Array



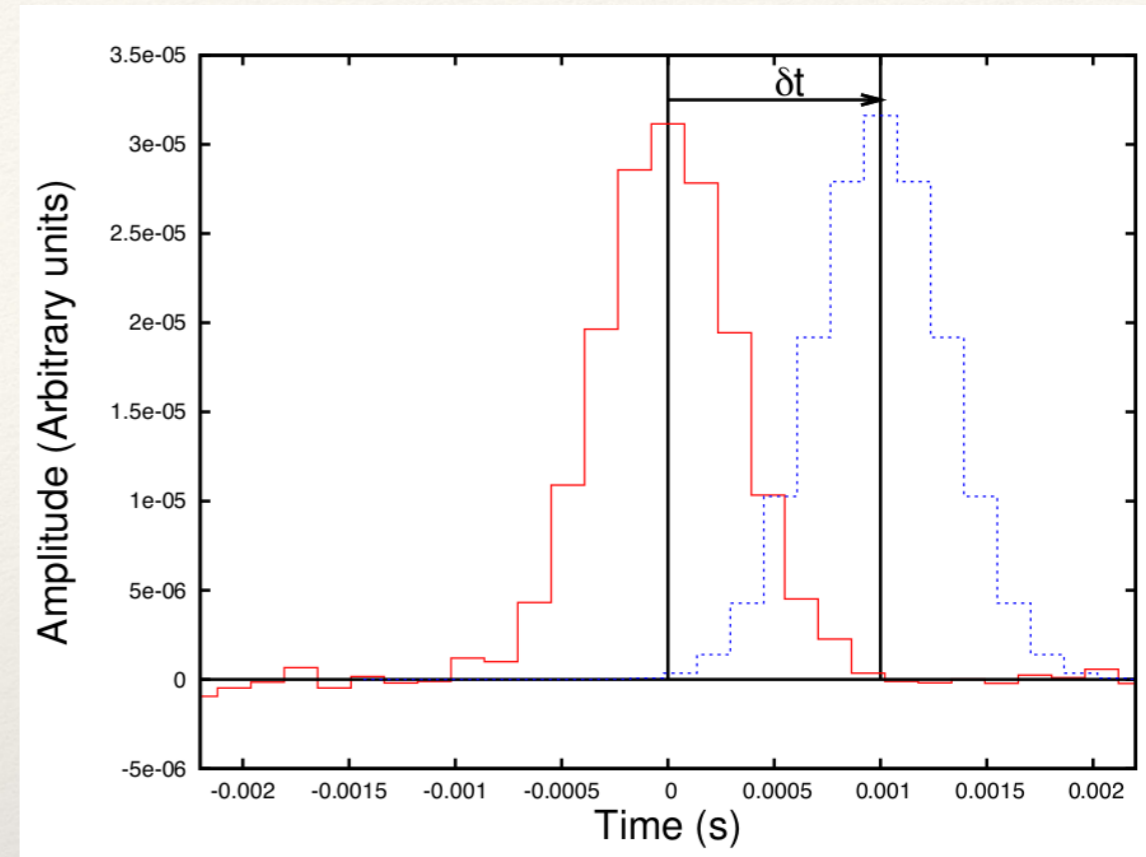
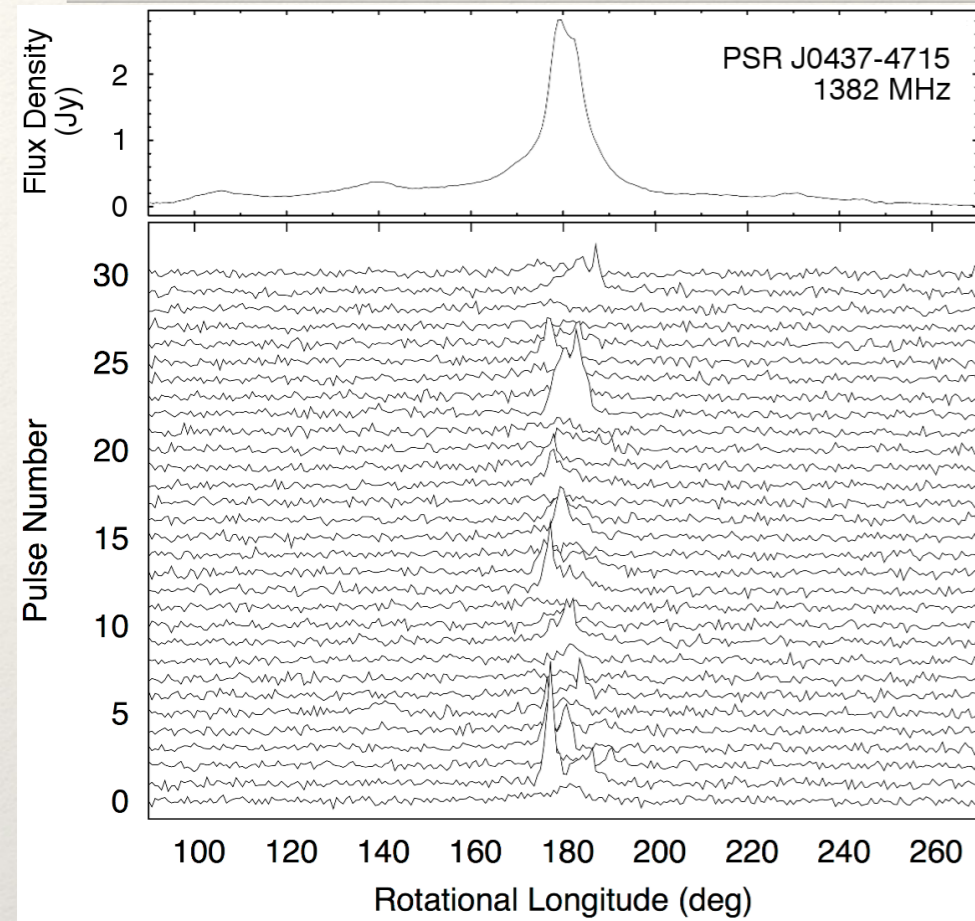
The main idea behind pulsar timing array (PTA) is to use ultra-stable millisecond pulsars as beacons (clocks sending signals) for detecting GW in the nano-Hz range (10^{-9} - 10^{-7} Hz).



[Credits: D. Champion]



Pulsar timing

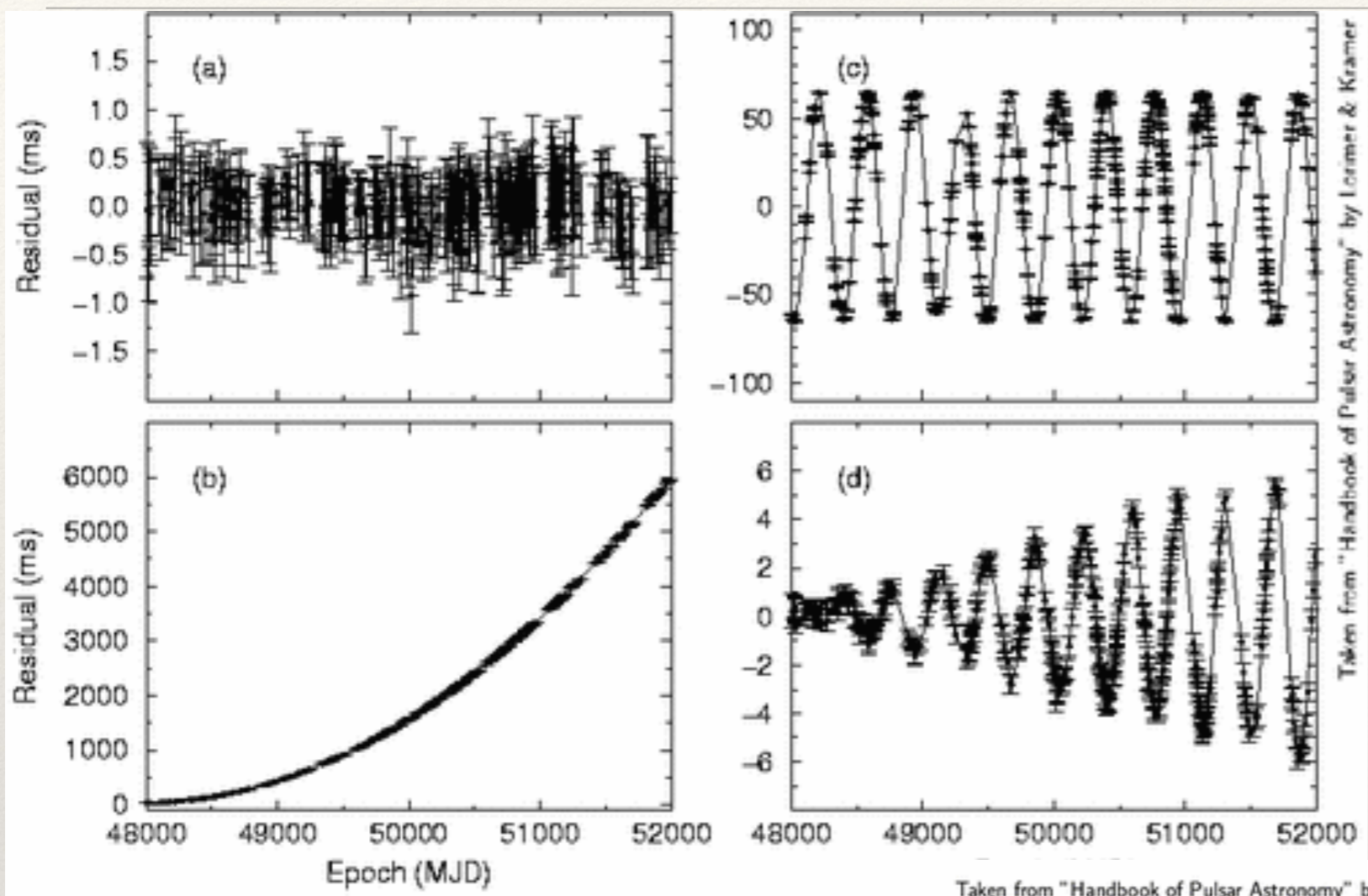


[Figs: credits
S. Burke-Spolar & L. Lentati]

- Each observed radio pulse profile has a lot of micro-structure. If we average over ~hour the (average) profile is very stable
- We can use the average pulse profile to estimate the time-of-arrival (TOA) of the pulses.
- The idea is to measure the TOA, and compare to the expected TOA. We know the spin of the pulsars, so we can predict the TOA. The difference between measure and expected TOA: *residuals*



Timing pulsars

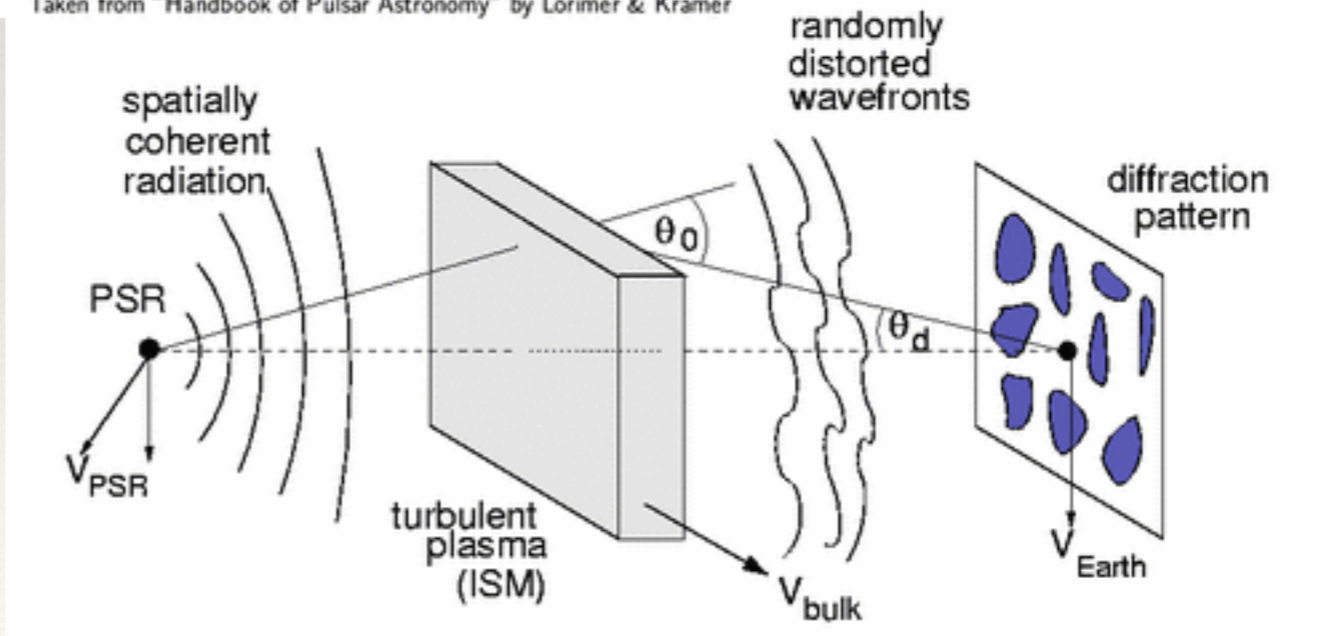


Taken from "Handbook of Pulsar Astronomy" by Lorimer & Kramer

- We need to build a timing model to make accurate prediction for TOAs - take into account various physical effects
- Dispersion of e/m wave and its time dependence
- Rate of change of rotation (b)
- Sky position of the pulsar (c)
- Proper motion of a pulsar (d)

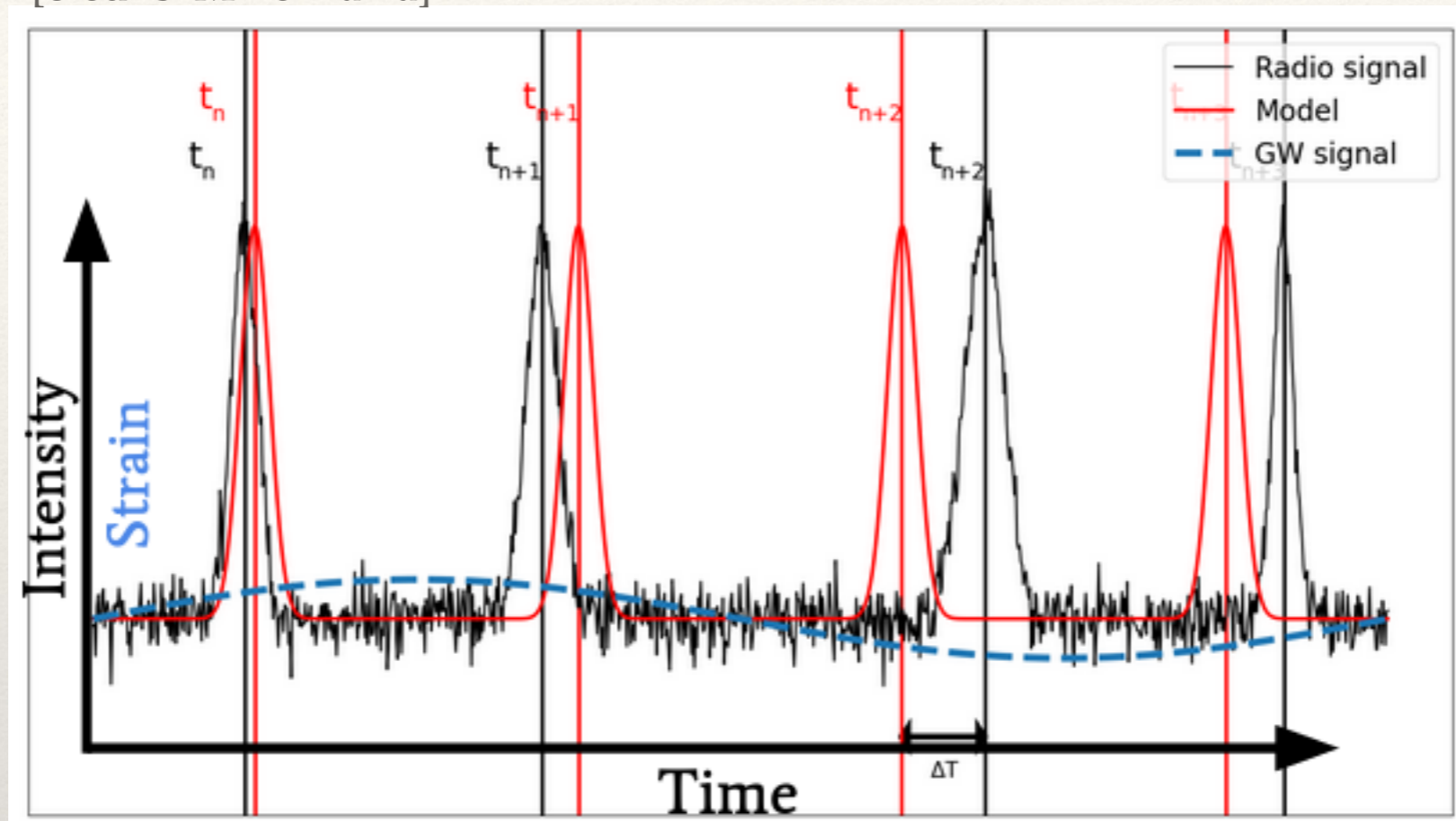
○ Timing model could be quite complex if pulsar is in the binary

Taken from "Handbook of Pulsar Astronomy" by Lorimer & Kramer



Timing Residuals

[credits: Mikel Falxa]



$$dt = t_{toa}^p - t_{toa}^o = dt_{errors} + \delta\tau_{GW} + noise$$

Errors in fitting the model \nearrow \nwarrow due to GWs

Detection statistic and search algorithm



- We assume that noise is Gaussian: the likelihood function (likelihood of the signal with given parameters) is

$$P(\vec{\delta t}, \vec{\theta}) = \frac{1}{\sqrt{(2\pi)^n \det(C)}} \exp\left(-\frac{1}{2}(\vec{\delta t} - \vec{s})^T C^{-1}(\vec{\delta t} - \vec{s})\right),$$

- $\vec{\delta t}$ - concatenated residuals from all pulsars in the array: total size n
- \vec{s} - is a model of deterministic signals (for example GW signals from individually resolvable SMBHBs)
- C is the noise variance-covariance matrix (size $n \times n$)

$$C_{\alpha i, \beta j} = C^{\text{wn}} \delta_{\alpha\beta} \delta_{ij} + C_{ij}^{\text{rn}} \delta_{\alpha\beta} + C_{ij}^{\text{dm}} \delta_{\alpha\beta} + C_{\alpha i, \beta j}^{\text{GW}} + \dots$$

white measurement noise	red noise spin noise	dispersion variation noise	stochastic GW signal
-------------------------------	----------------------------	----------------------------------	-------------------------



Noise modelling in PTA



- White noise — not very interesting. Two parameters per backend per pulsar: unaccounted noise.
- Red noise: very generic noise description in freq. domain

$$S(f) = A_{rn}^2 f^{-\gamma}$$

common, uncorrelated
red noise

$$S_{\alpha}(f) = A_{rn,\alpha}^2 f^{-\gamma_{\alpha}}$$

red noise in each pulsar

- DM (dispersion measurement variation) noise: depends on the radio-frequency of observation

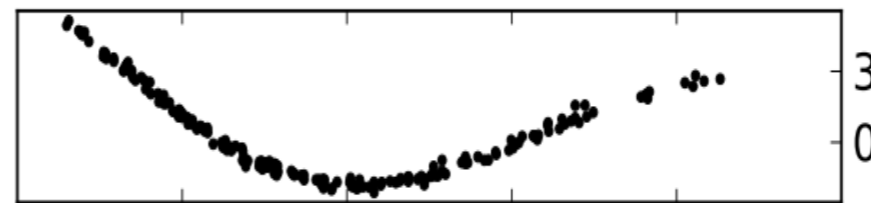
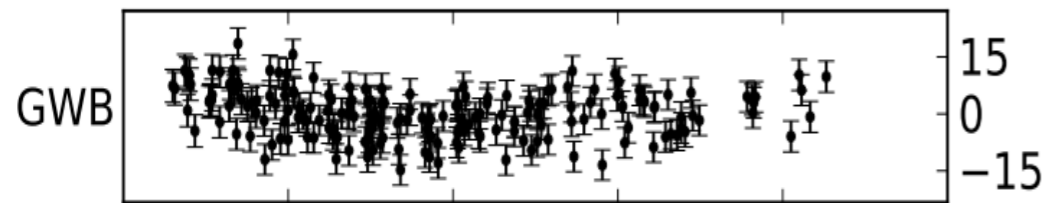
$$S_{DM}(f) \propto \frac{A_{dm}^2}{\nu^2} f^{-\gamma_{dm}}$$

- Correlated red noise processes

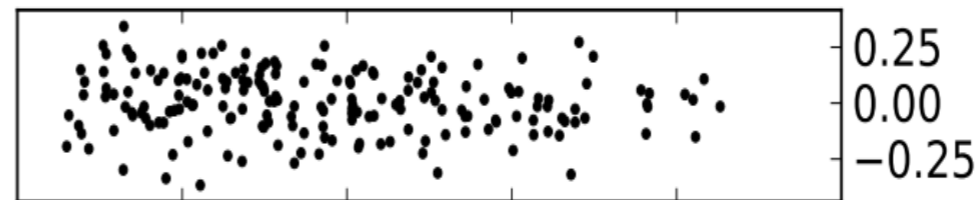
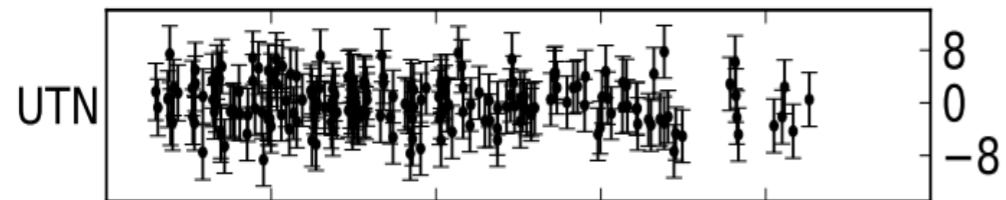
$$S_{\alpha\beta} = \Gamma_{\alpha\beta} A_{cor}^2 f^{-\gamma_{cor}} \quad \text{— includes also cross spectrum between each pair of pulsars: } \Gamma_{\alpha\beta} \text{ - spacial correlation coefficients}$$



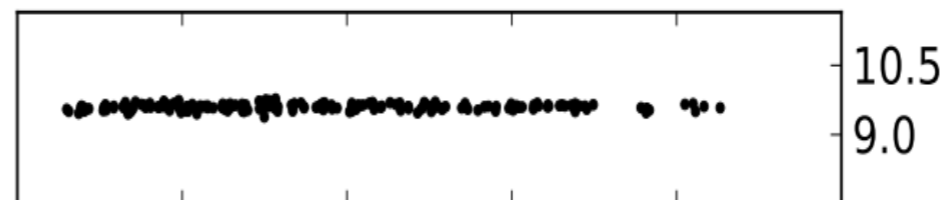
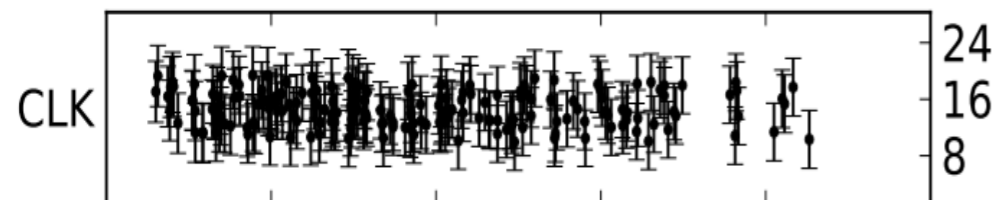
Correlated noise



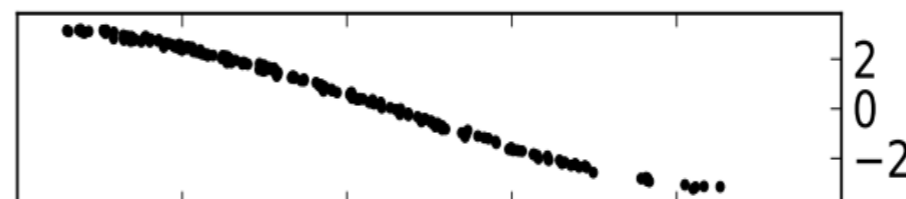
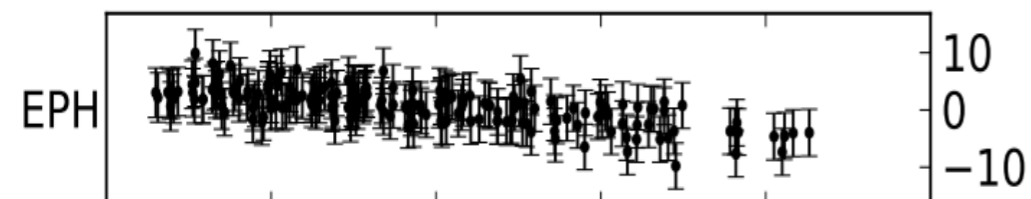
GWB (Quadrupole)
Hellings-Downs



Uncorrelated RN



Clock error (monopole)



Ephemeris (dipole)

TOAs, time

Correlation, angular separation

stochastic GW from population of SMBHBs:

$$S_{\alpha\beta}^{SMBHB} = \Gamma_{\alpha\beta}^{H-D} A_{GW}^2 f^{-13/3}$$



Gaussian-process approach to PTA: falling into a rabbit hole



- Applying GP to the PTA likelihood function:

$$p(\delta t|w_i, GP) = \frac{e^{-\frac{1}{2} \cdot \sum_{ij} (\delta t_i - \sum_a \phi_a(t_i) w_a) (C_{ij}^w)^{-1} (\delta t_j - \sum_a \phi_a(t_j) w_a)}}{\sqrt{(2\pi)^n \det(C^w)}} \times \frac{e^{-\frac{1}{2} \sum_{a,b} w_a (\Sigma_{ab})^{-1} w_b}}{\sqrt{(2\pi)^m \det(\Sigma)}}$$

GP
Gaussian prior on weights

white noise

weight-space approach

$$p(\delta t|w_i, GP) = \frac{e^{-\frac{1}{2} \cdot \sum_{ij} \delta t_i (C_{ij}^w + C_{ij}^{rn})^{-1} \delta t_j}}{\sqrt{(2\pi)^n \det(C^w + C^{rn})}}$$

with $C_{ij}^{rn} = k(t_i, t_j) = \sum_{a,b} \phi_a(t_i) \Sigma_{ab} \phi_b(t_j)$

red noise variance-covariance matrix

In time domain, uncorrelated red noise:

$$C_{ij}^{rn} = A^2 (f_L / \text{yr}^{-1}) \left\{ \Gamma(1 - \gamma) \sin\left(\frac{\pi\gamma}{2}\right) (f_L \tau_{ij})^{\gamma-1} - \sum_n \frac{(-1)^n (f_L \tau_{ij})^{2n}}{(2n)!(2n+1-\gamma)} \right\}$$

where $\tau_{ij} = |t_i - t_j|$ and f_L is low freq. cut-off



Gaussian-process approach to PTA: falling into a rabbit hole



- Alternatively we can use basis functions: based on the decomposition of residuals in the Fourier modes:

$$\delta t(t_i) \approx \sum_k a_k \sin 2\pi f t_i + b_k \cos 2\pi f t_i$$

weights
basis functions $\phi^F(f_a, t_i) = \phi_a^F(t_i)$

We use non-complete set of Fourier modes: covariance matrix can be approximated as

$$C_{ij}^{rn} \approx \sum_{a,b} \phi_a^F(t_i) \Sigma_{ab}^F \phi_b^F(t_j) \quad \text{where}$$

$$\Sigma_{ab}^F \propto (A_{rn}^2 f_a^{-\gamma}) \delta_{ab} / T \quad \text{— red noise PSD}$$

and for stochastic GW signal: $C_{i\alpha, j\beta}^{GW} = \sum_{i\alpha, j\beta} \Gamma_{\alpha\beta} \phi_a^F(t_{i\alpha}) \Sigma_{ab}^{F, GW} \phi_b^F(t_{j\beta})$, where

$$\Sigma_{ab}^{F, GW} = (A_{GW}^2 f_a^{-\gamma_{gw}}) \delta_{ab} / T$$



Gaussian-process approach to PTA: falling into a rabbit hole



Advantage of this description: again likelihood

$$p(\delta t | w_i, GP) = \frac{e^{-\frac{1}{2} \cdot \sum_{ij} \delta t_i (C_{ij}^w + C_{ij}^{rn})^{-1} \delta t_j}}{\sqrt{(2\pi)^n \det(C^w + C^{rn})}}$$

Data size: n - large, need to invert very large (covariance) matrices - $n \times n$

Can use Woodbury f-la

$$(C_w + C_{rn})^{-1} = (C_w + \Phi \Sigma \Phi^T)^{-1} = C_w^{-1} - C_w^{-1} \Phi \underbrace{(\Sigma^{-1} + \Phi^T C_w^{-1} \Phi)^{-1}}_{\text{inversion of } m \times m \text{ matrix}} \Phi^T C_w^{-1}$$

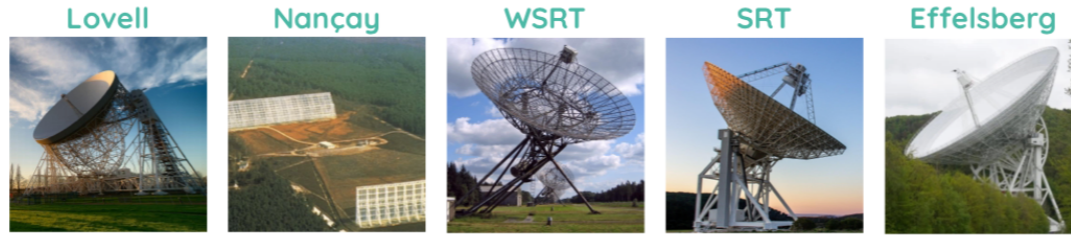
Number of modes: $m \ll n$ much faster and easier to invert, C_w is diagonal matrix

Bayesian analysis: model selection (hypothesis testing)

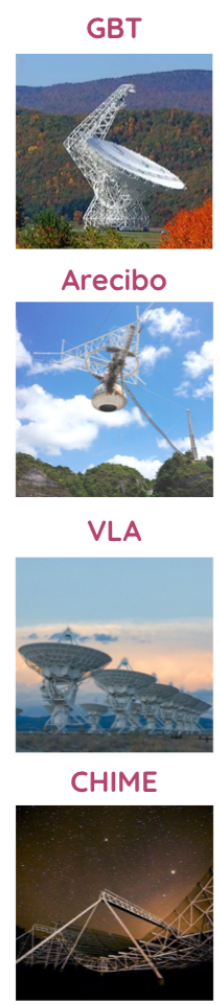
Odd ratio: $O(M_1, M_2) = \frac{p(M_1 | d)}{p(M_2 | d)} = \frac{p(d | M_1) \pi(M_1)}{p(d | M_2) \pi(M_2)}$



IPTA

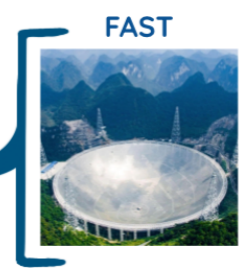


EPTA



NANOGrav

IPTA

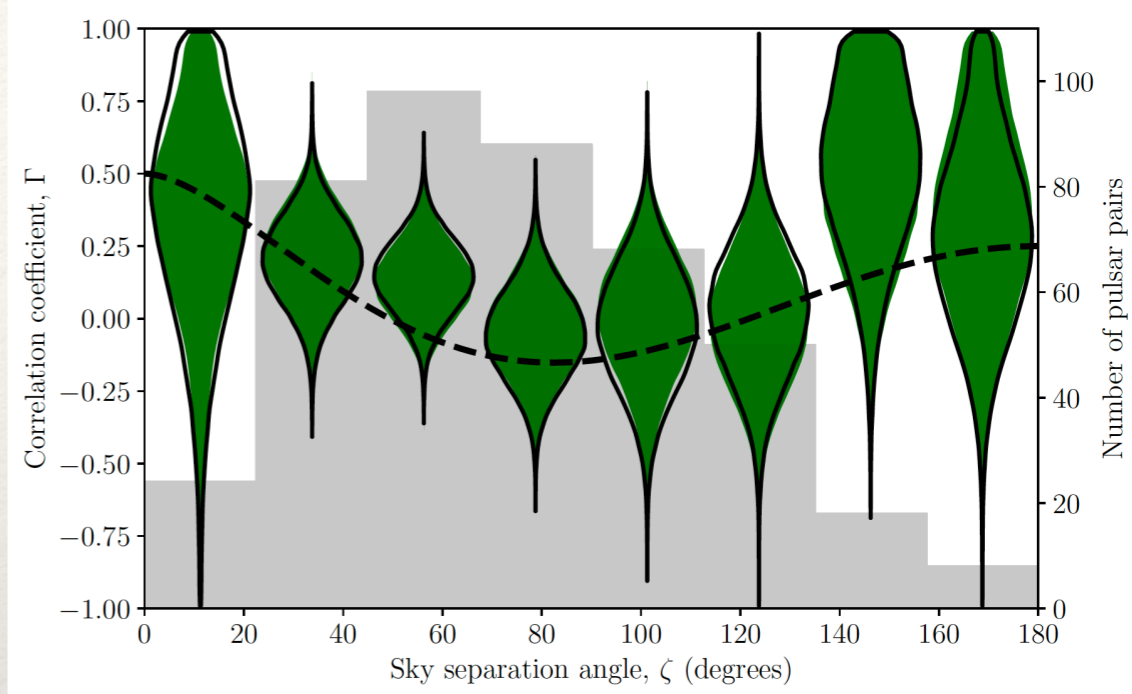
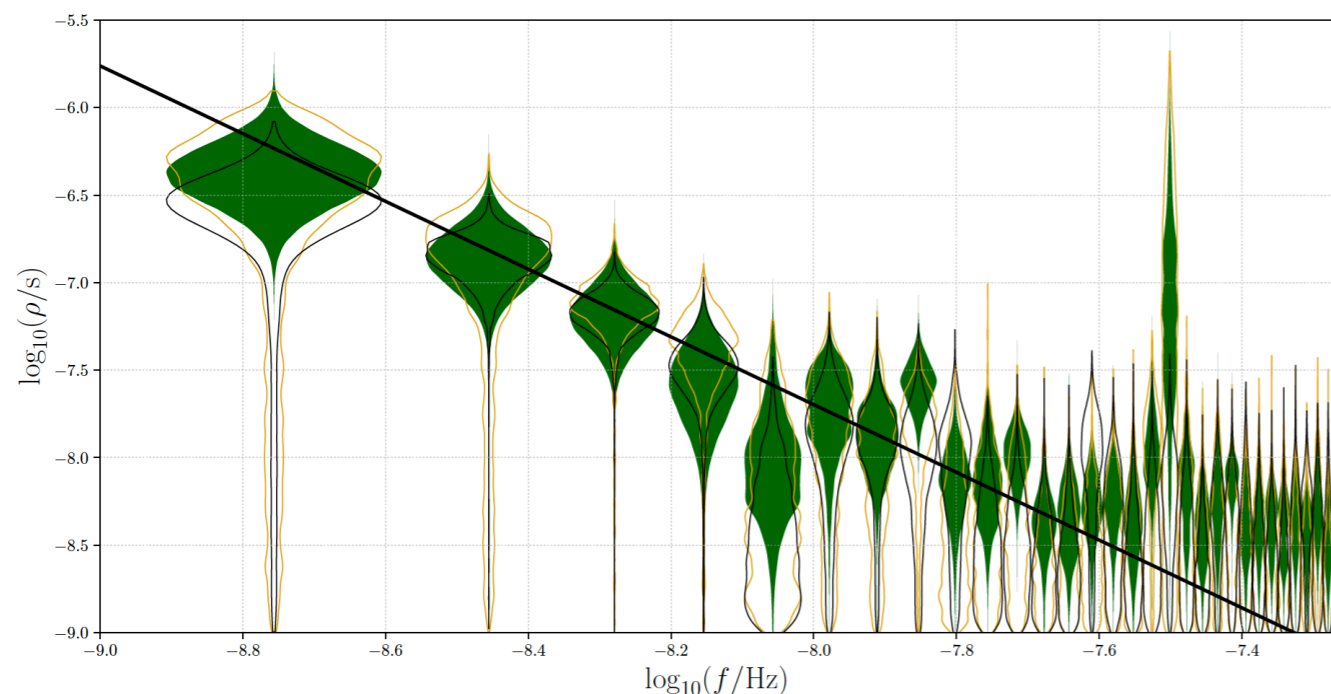


PPTA results

[PPTA 2306.16215]

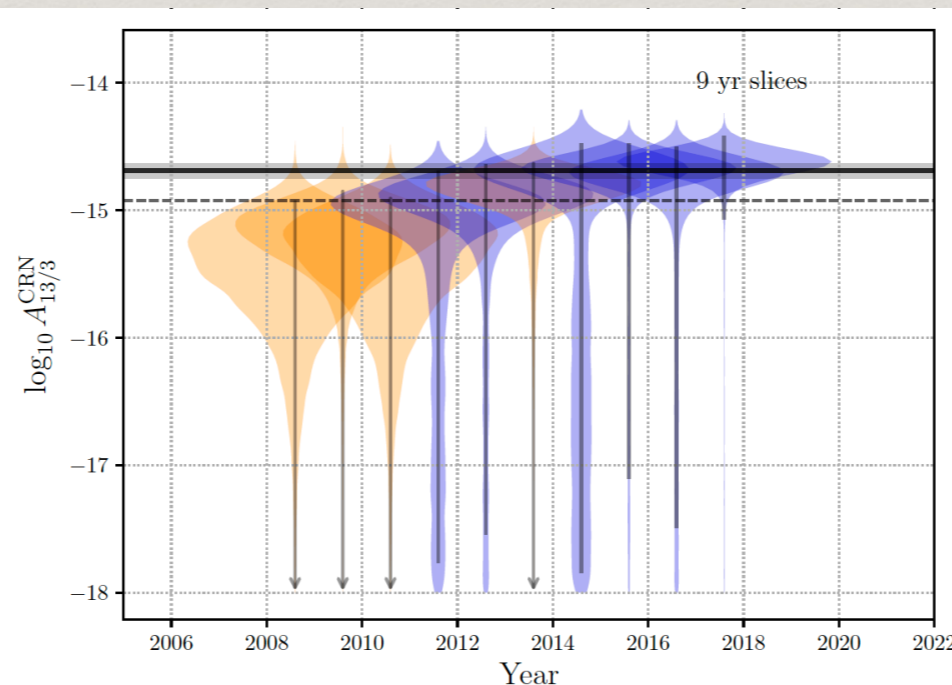
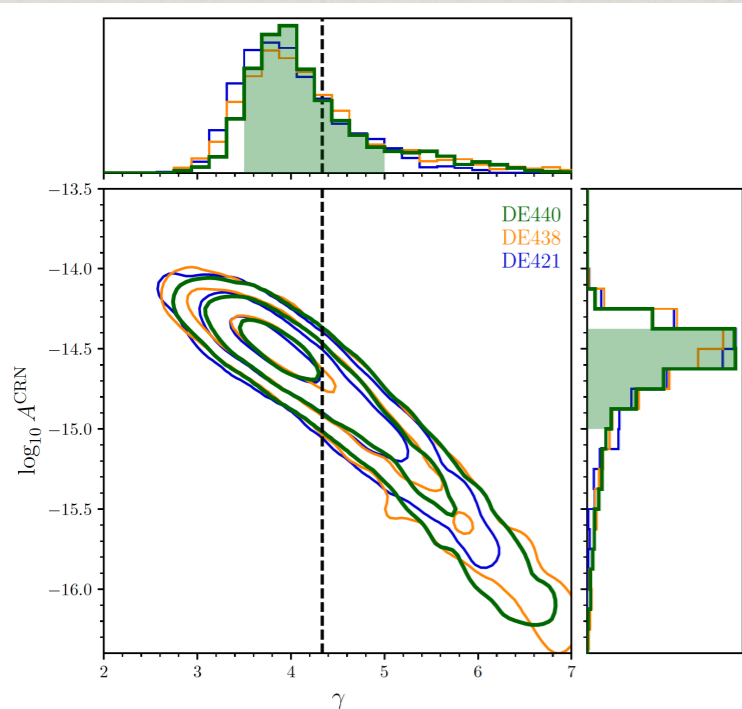


PPTA data: 18 years, 30 pulsars. 3 years of new ultra-widebandwidth radio observations



Estimating power at Fourier freq. (assuming independence).

Black: CURN, Gold: H-D



$$S_{\alpha\beta}^{SGWB} = \Gamma_{\alpha\beta}^{H-D} A_{GW}^2 f^{-\gamma}$$

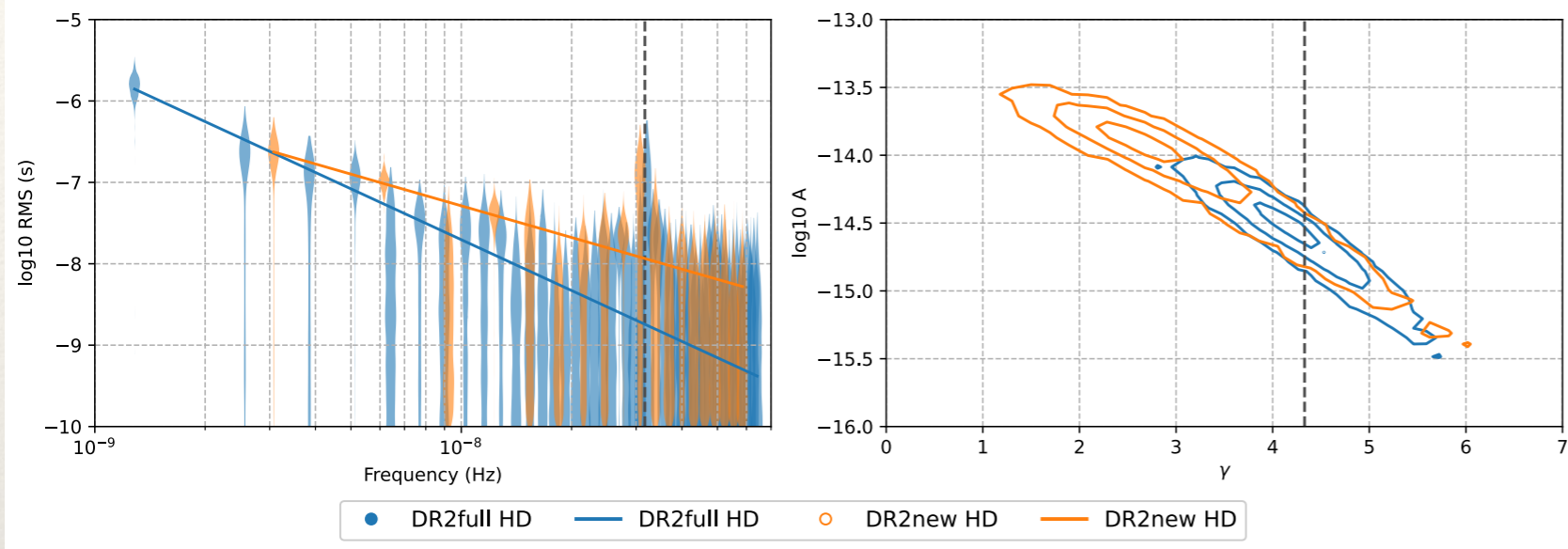


EPTA + InPTA



[EPTA+InPTA2306.16214]

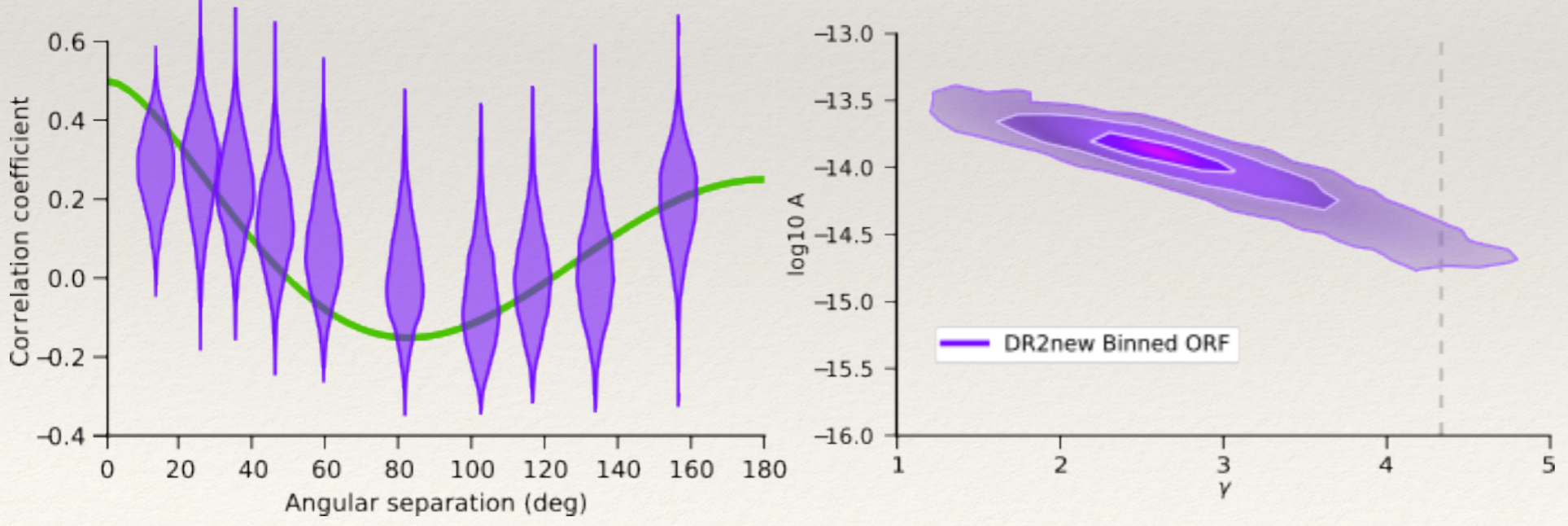
25 plsr, DR2full: up to 25yrs, DR2new: latest 14 yrs, DR2new+ Includes InPTA data (3.5 yrs)



$$S_{\alpha\beta}^{SGWB} = \Gamma_{\alpha\beta}^{H-D} A_{GW}^2 f^{-\gamma}$$

Statistic	DR2full	DR2new	DRnew+
\mathcal{B}_{CURN}^{HD}	4	60 33(+Ephem)	65 43(+Ephem)

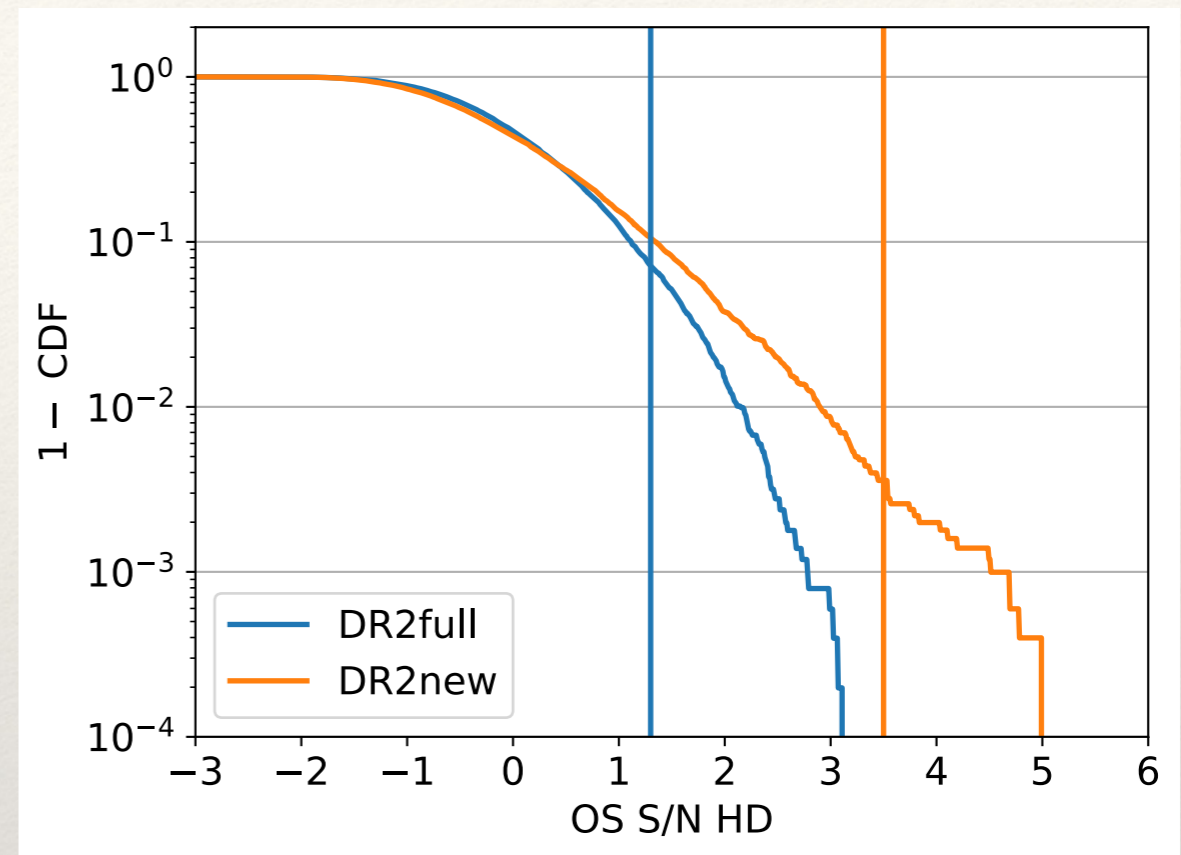
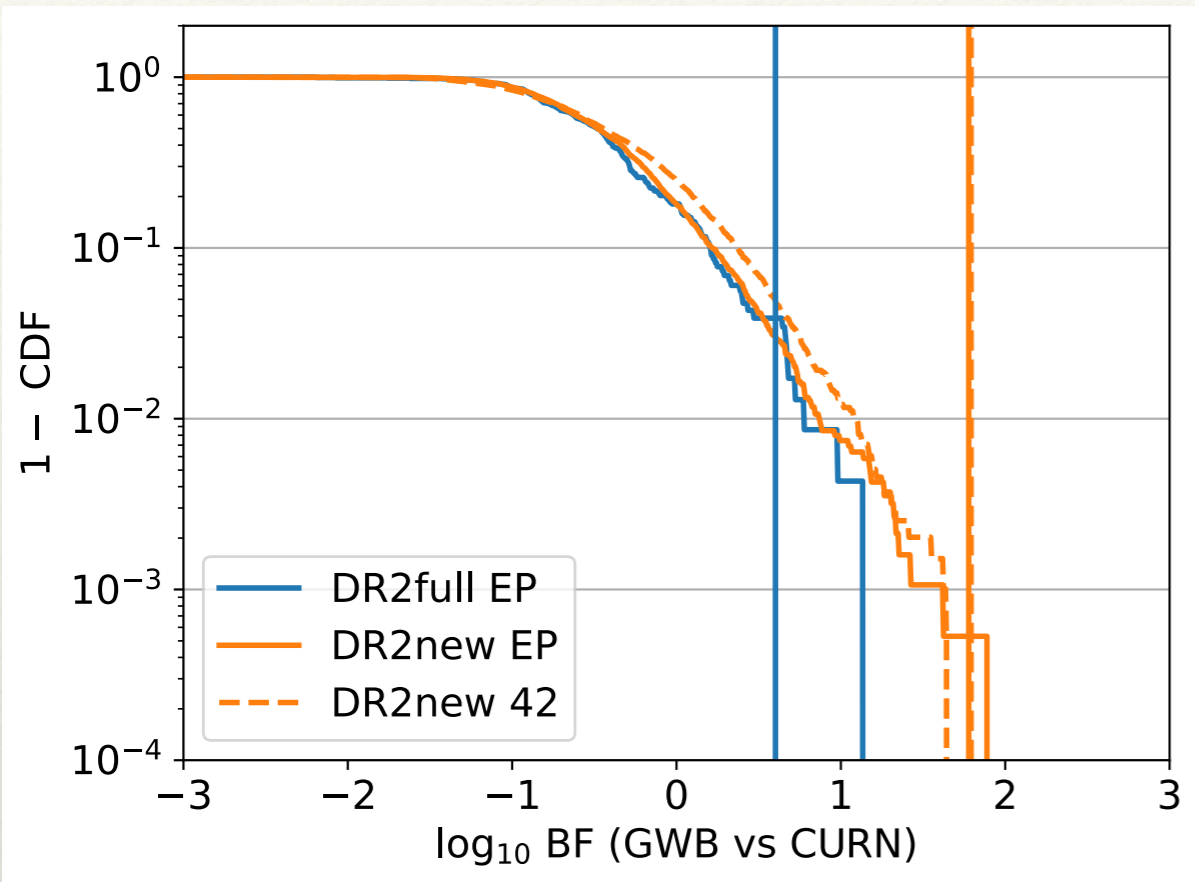
DR2new results: spatial correlations and amplitude-slope of power-law model



EPTA + InPTA



Significance: how likely to observe what we observe in absence (null hypothesis) of GW signal



We want [Cornish & Sampson 2015, Taylor+ 2016]

- Preserve properties of the noise (use observations)
- Data free of GW signal: not possible, instead we try to mimick measurements insensitive to GWs
 - Sky shuffling: change position of pulsars: observed correlation is not consistent with GW
 - Phase shift: introduce a random shift in phase at each frequency bin: destroy correlations

The question we are asking:

- how likely to get observed H-D pattern by randomly choosing pulsars on the sky
- How likely that the phases at low frequencies in all pulsars align to form observed H-D

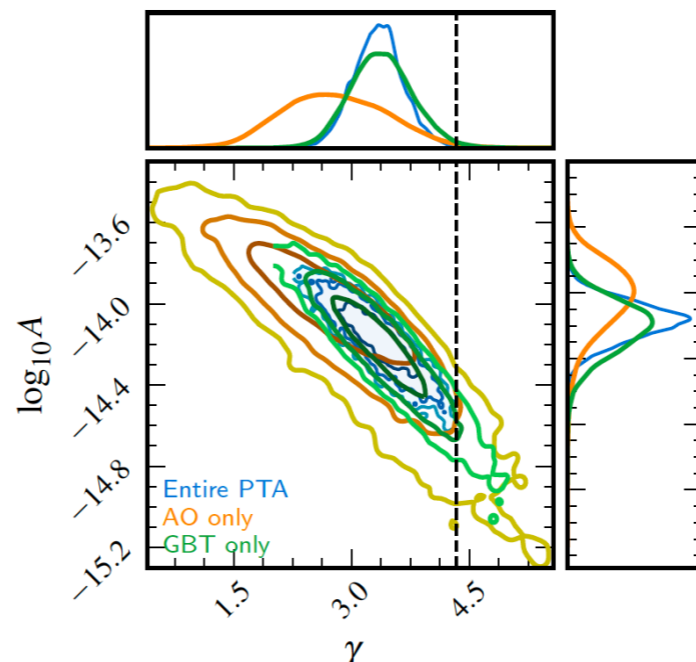
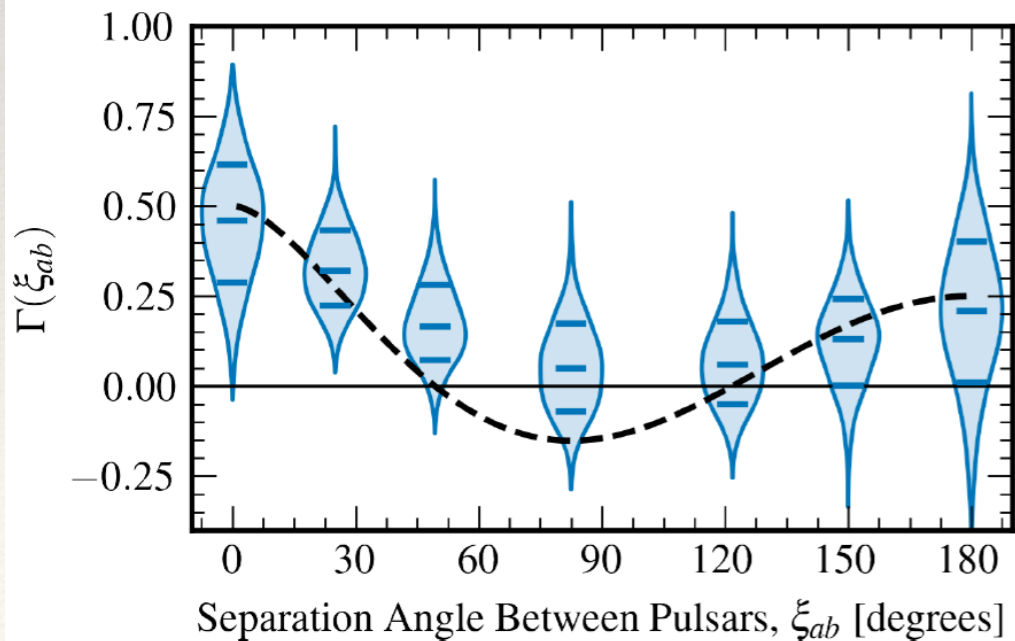
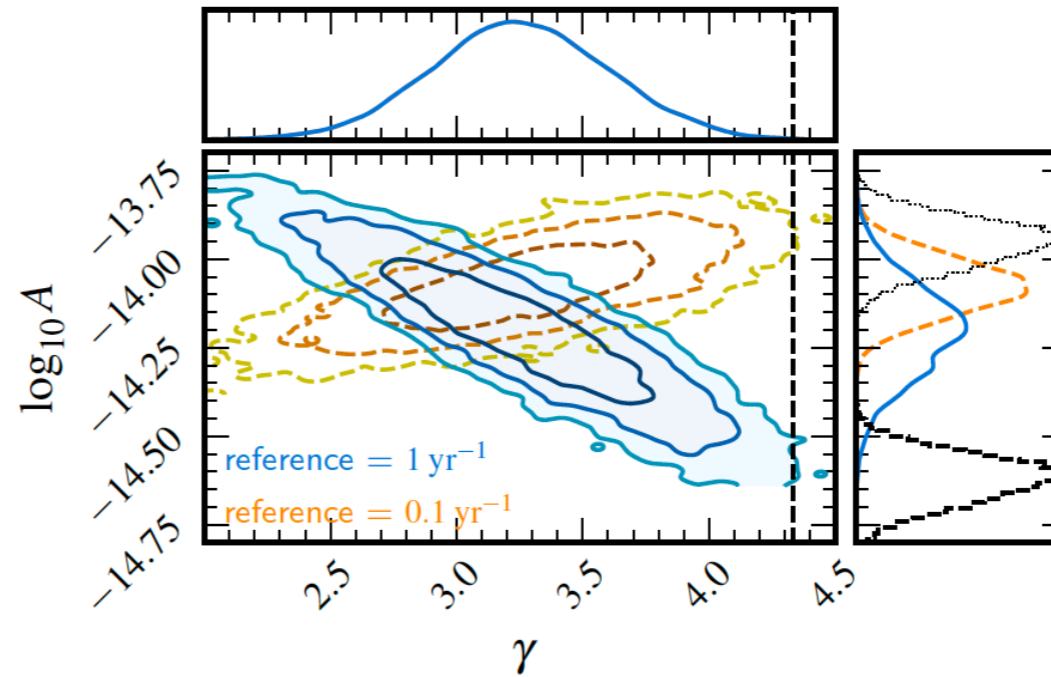
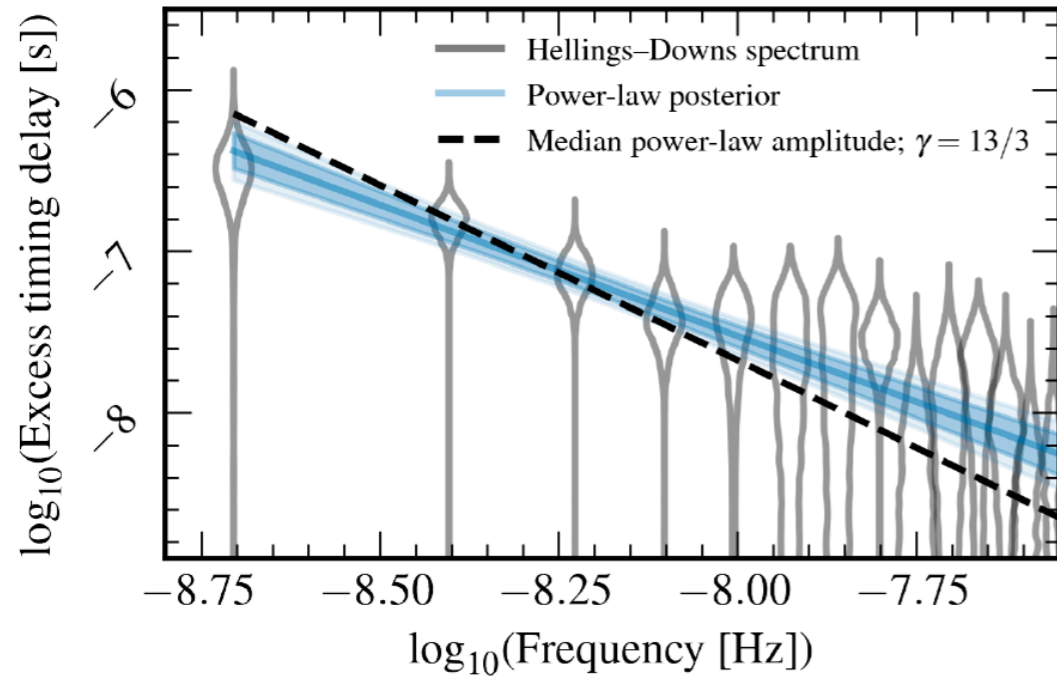


NanoGrav results

[NG 2306.16213]



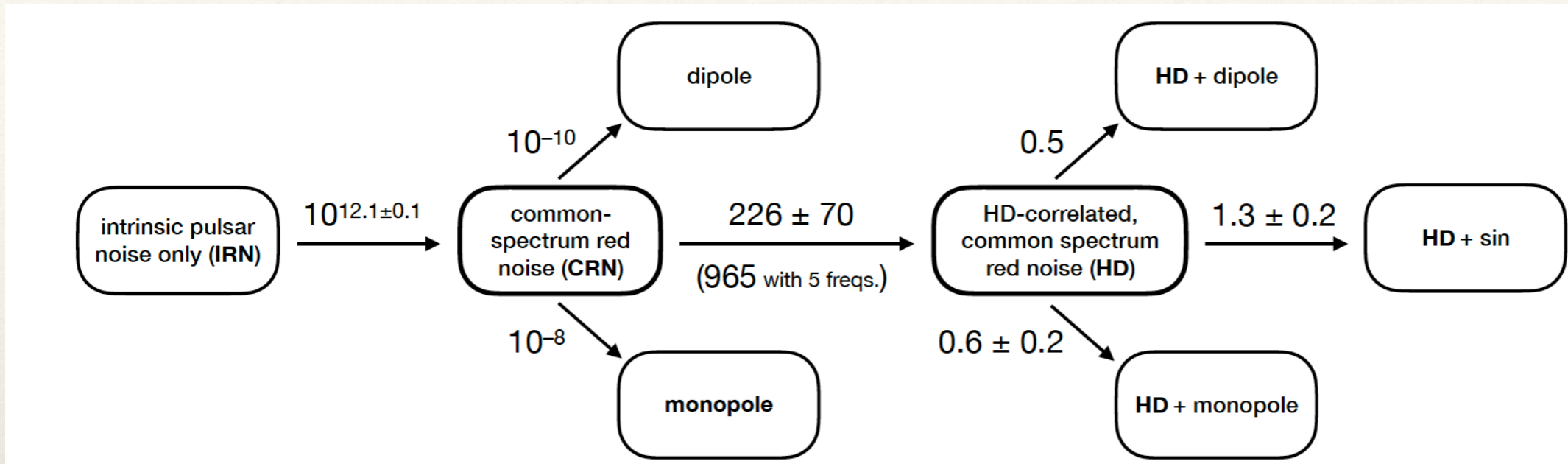
NG data: 15 years of data, 67 pulsars. Arecibo + Green Bank



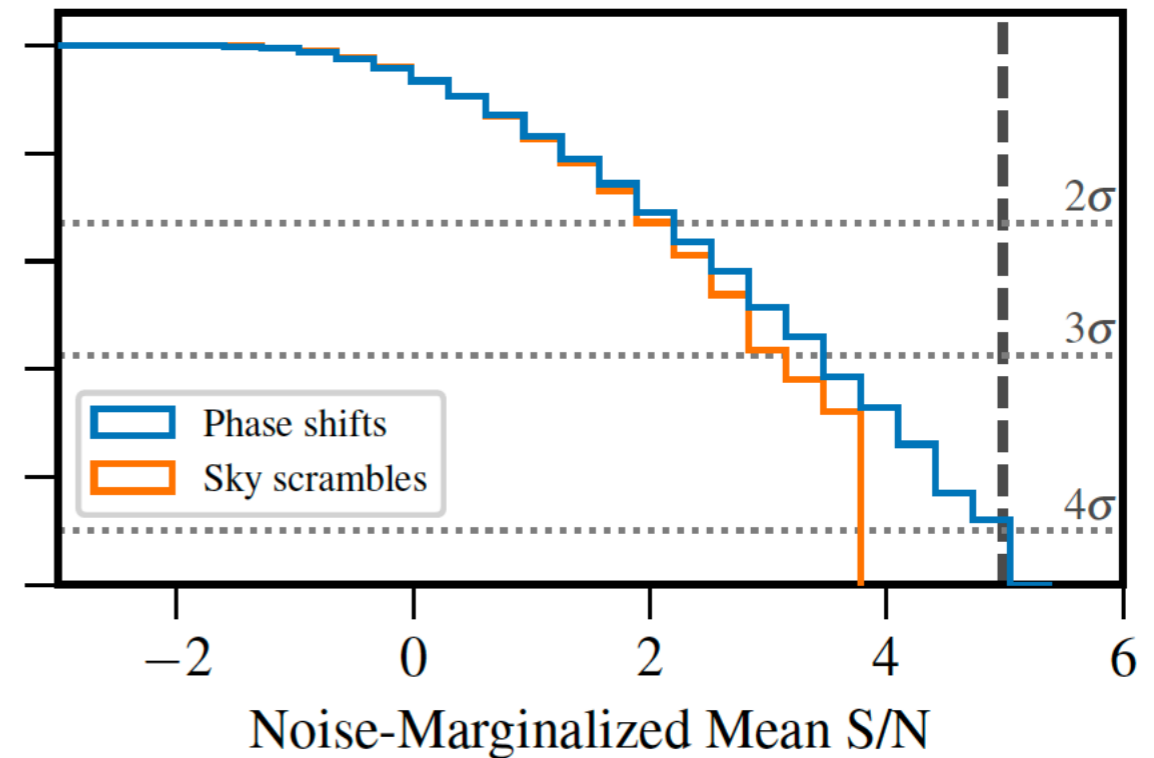
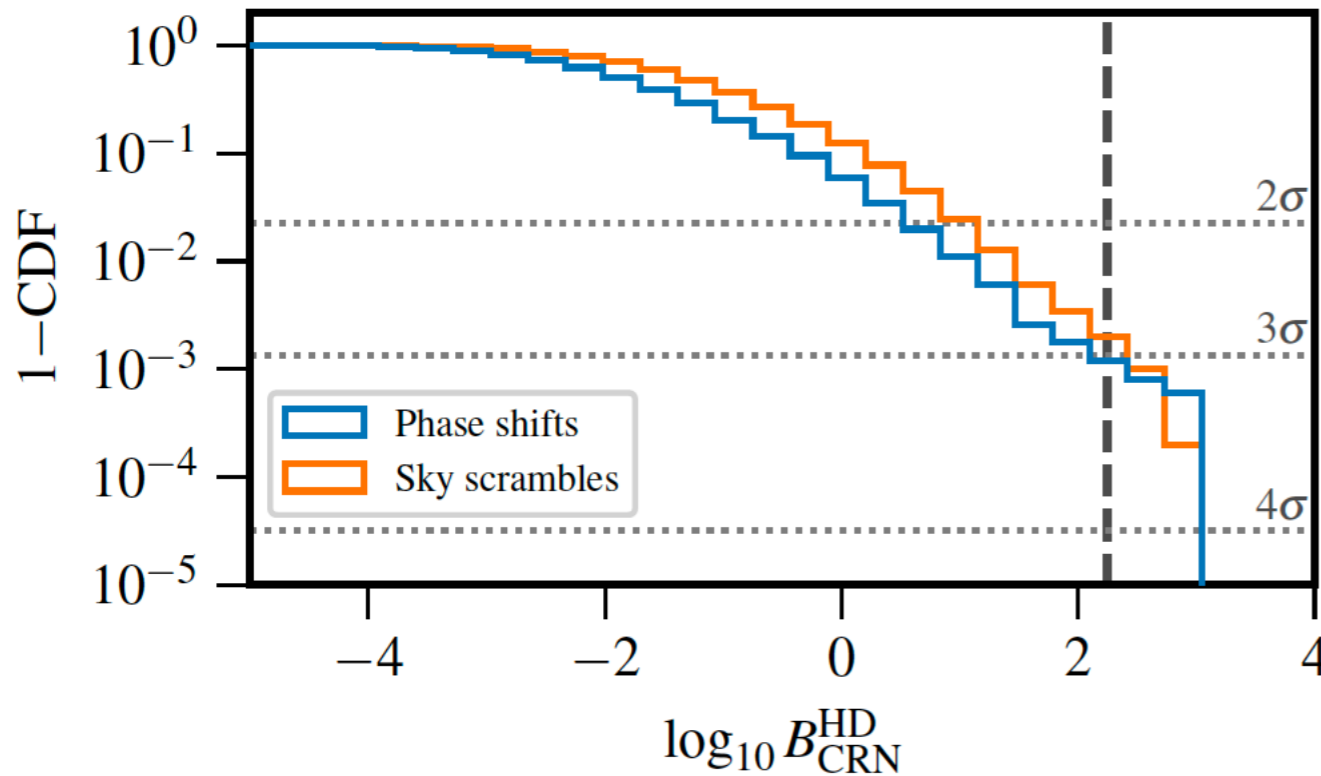
$$S_{\alpha\beta}^{SGWB} = \Gamma_{\alpha\beta}^{H-D} A_{GW}^2 f^{-\gamma}$$



NanoGrav results



Significance



Interpretation



LET US ASSUME THAT WHAT WE OBSERVE IS STOCHASTIC GW BACKGROUND
(SGWB)

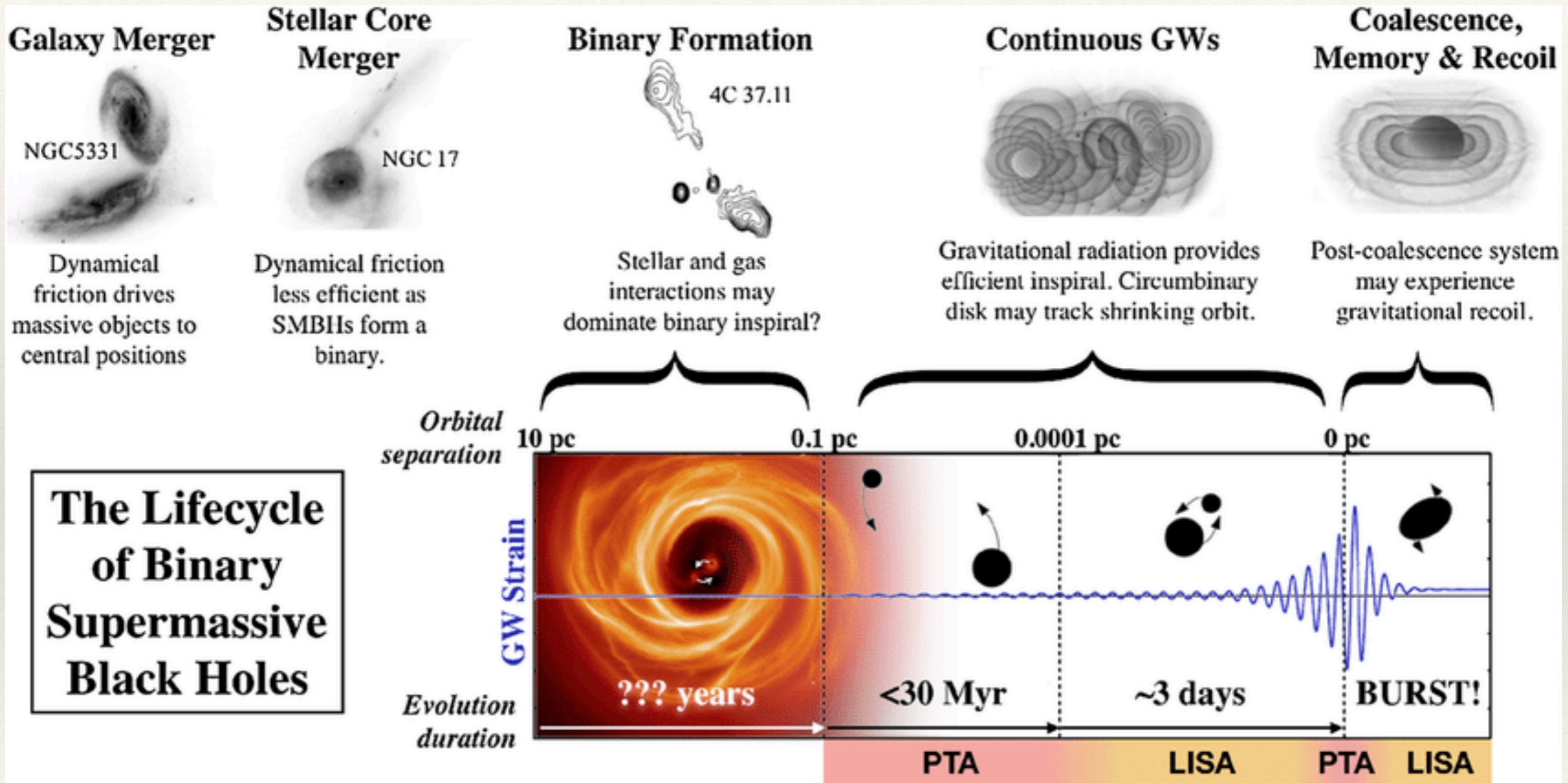
What could produce SGWB with the power-law-like spectrum?
Apparently almost anything that falls in nHz band... and even more
I'll give only few examples

DISCLAIMER

preference in interpretation of observed signal and its significance: my personal view



Massive black hole binaries



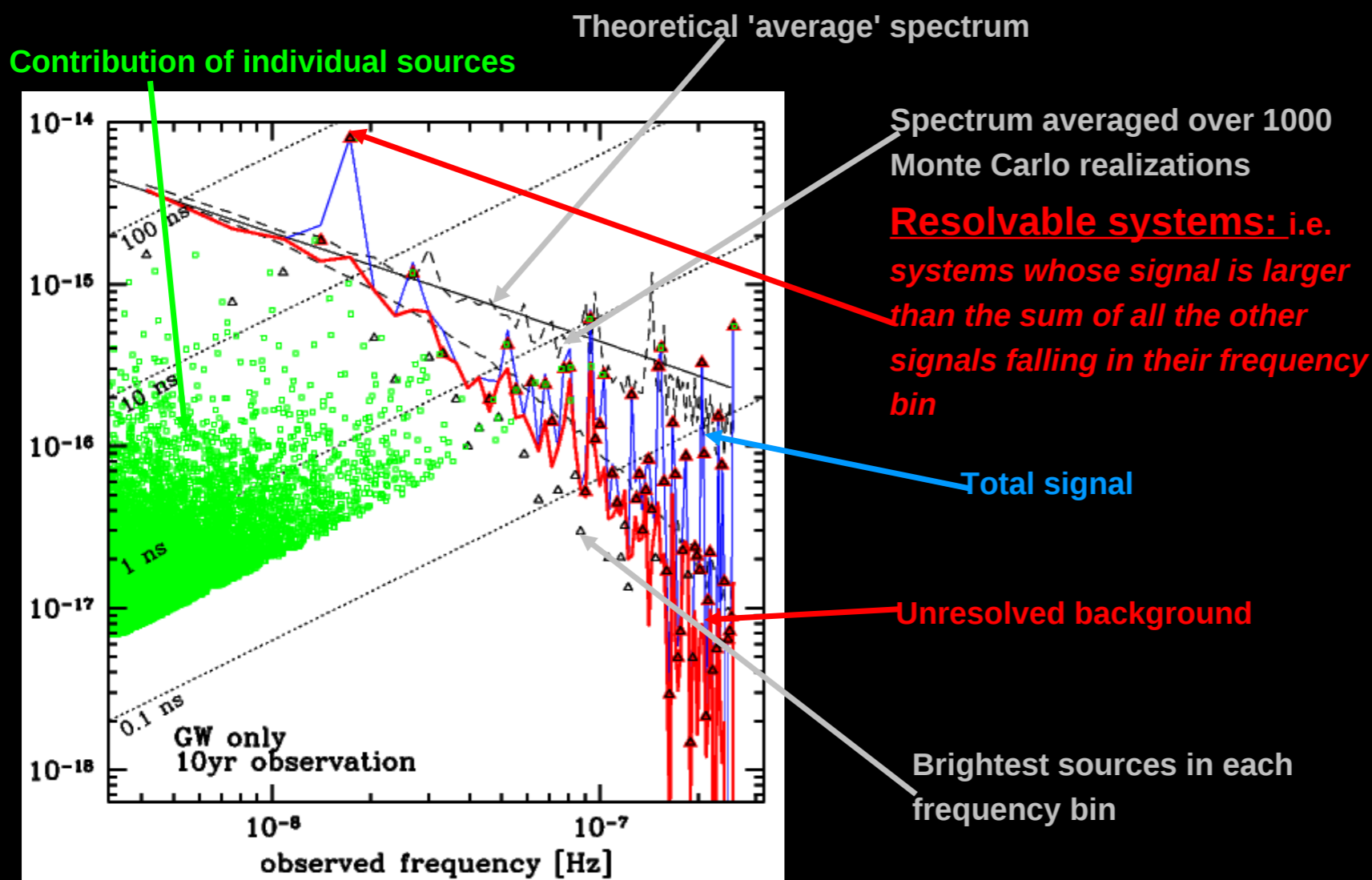
[S. Burke-Splolaor A&A review (2019)]

Supermassive black hole binaries



- Main sources are supermassive black hole binaries (mass 10^7 — 10^{10} solar) on very broad orbit (period \sim year(s))
- The orbital evolution due to GW emission is very slow: $\frac{dE}{dt} \propto \eta(M/r)^5$
signal is (almost) monochromatic over period of observations

Signal from a MBHB population



GW signal from the population of SMBH binaries: forms a stochastic signal at low freqs. (similar to Galactic binaries in LISA)

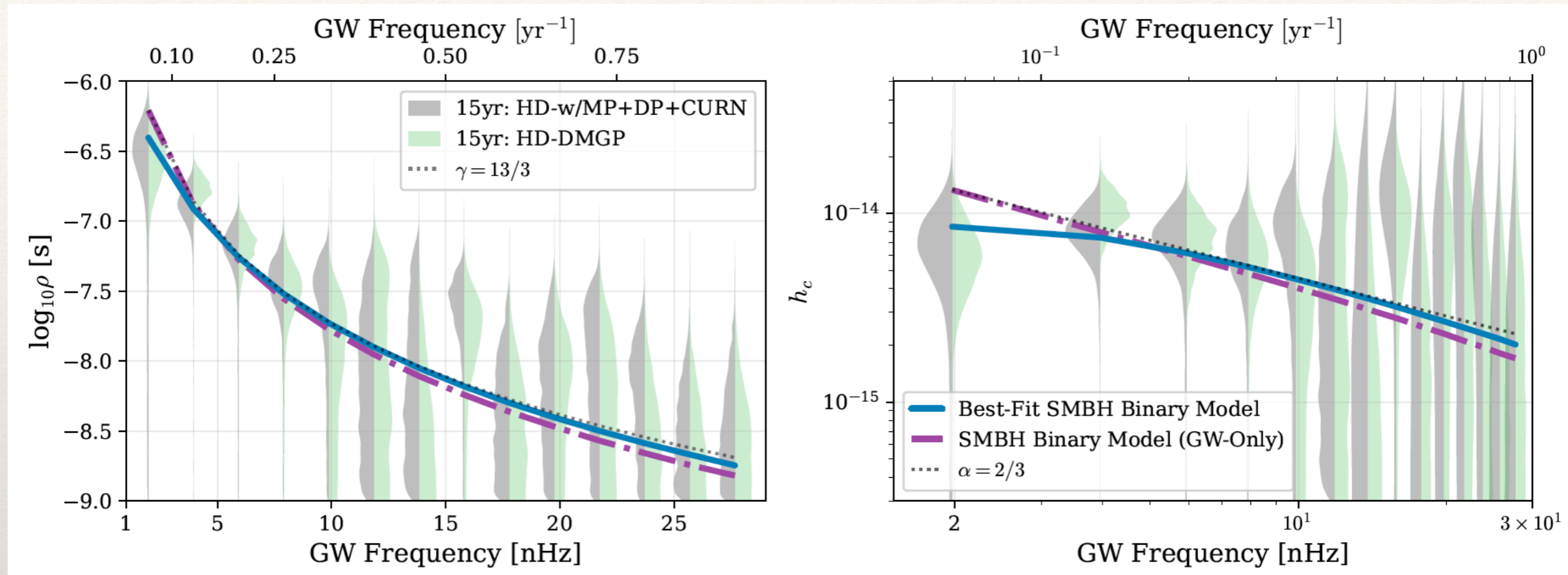
[Credits: A. Sesana]



SGWB from population of SMBHBs



[NG: 2306.16220]



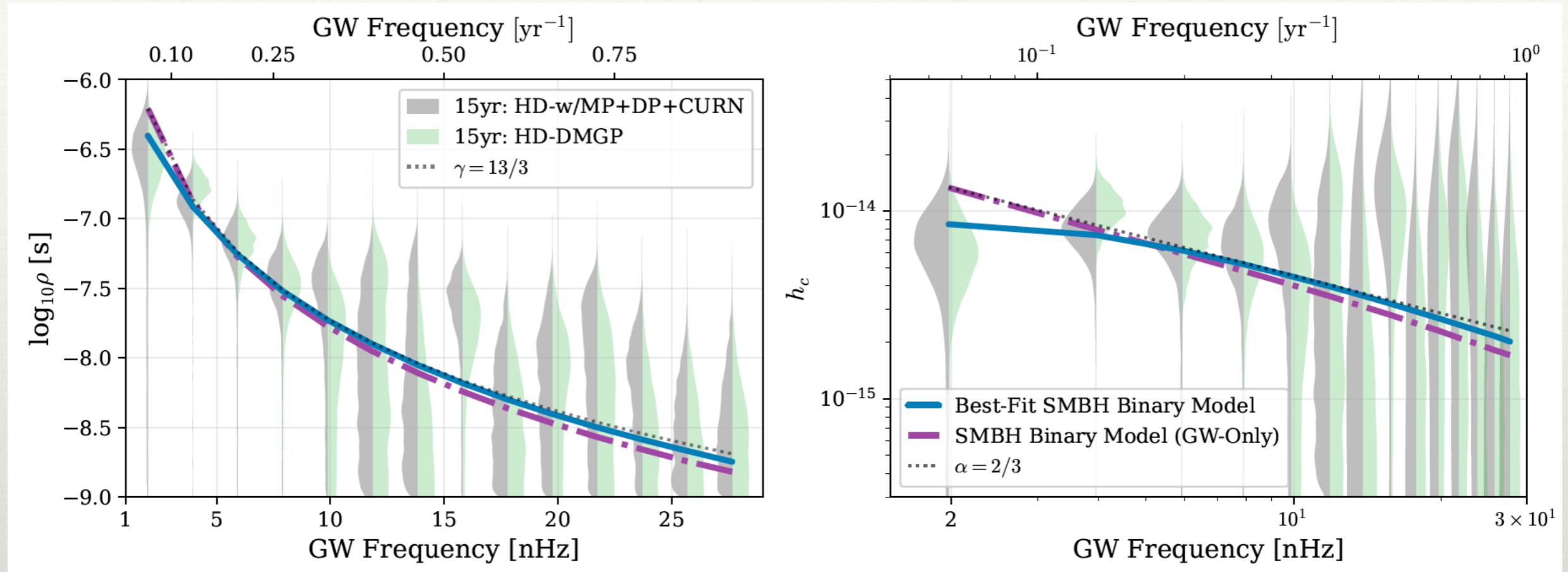
- Free spectrum (HD) : estimation of the ASD for the H-D correlated part of the noise
- SMBHB population: circular orbit + GW driven evolution: $\rho \propto f^{-13/3}$, $h_c \propto f^{-2/3}$ (black dots)
- Eccentric orbits: redistribution of GW energy towards higher modes (higher frequencies): lower amplitude, turn-over at low freq.
- Interaction with environment: dissipation of energy and angular momentum: turn-over at low freq.



SGWB from population of SMBHBs



[NG: 2306.16220]



Blue: uses self-consistent SMBH binary evolution model (slightly preferred)

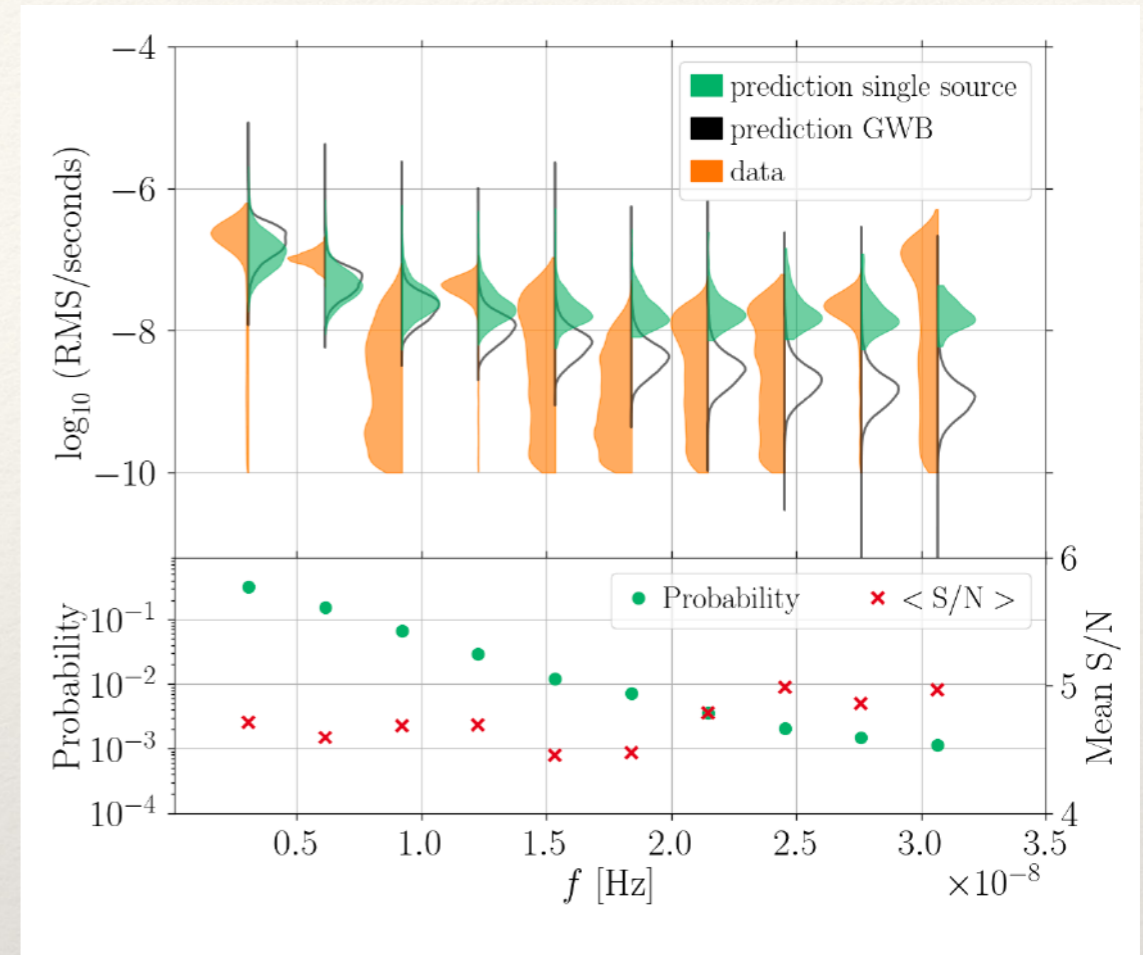
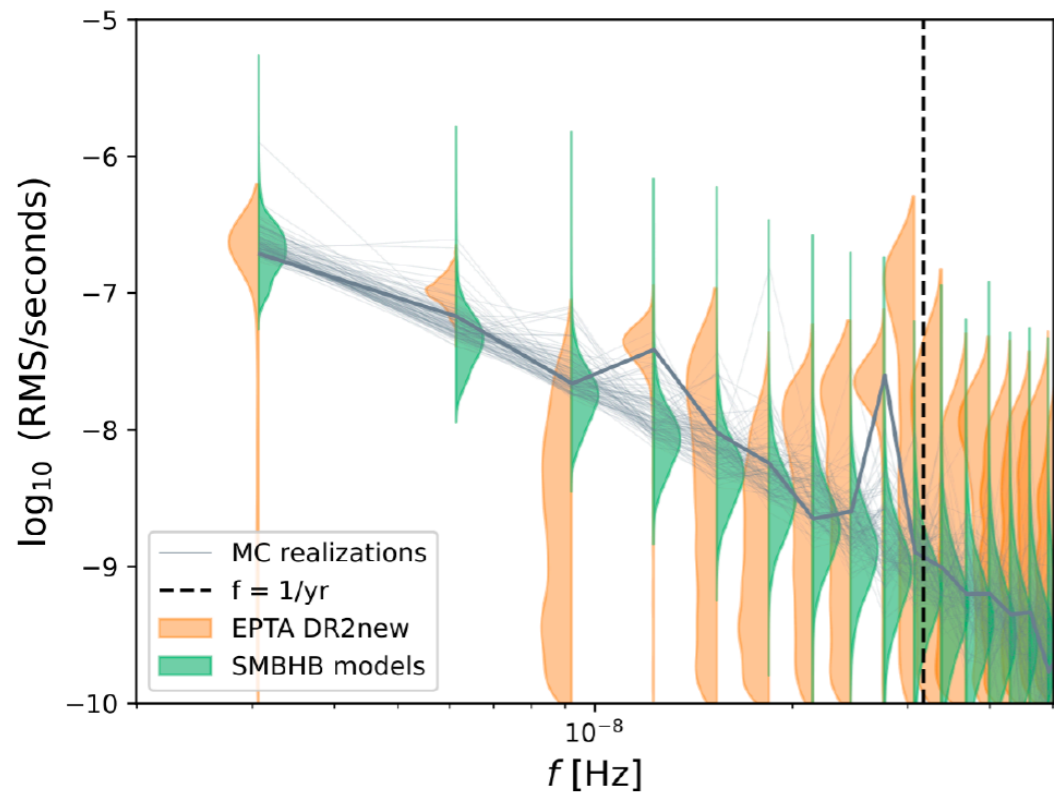
Purple: GW-only driven evolution model



SGWB from population of SMBHBs



[EPTA 2306.16227]



Astrophysically-informed SMBHB population model: interaction with environment, allows eccentric orbits (similar to “blue (phenom) model” of NG)

- Constrains SMBHB merger timescale
- Constrains SMBH-buldge mass relation

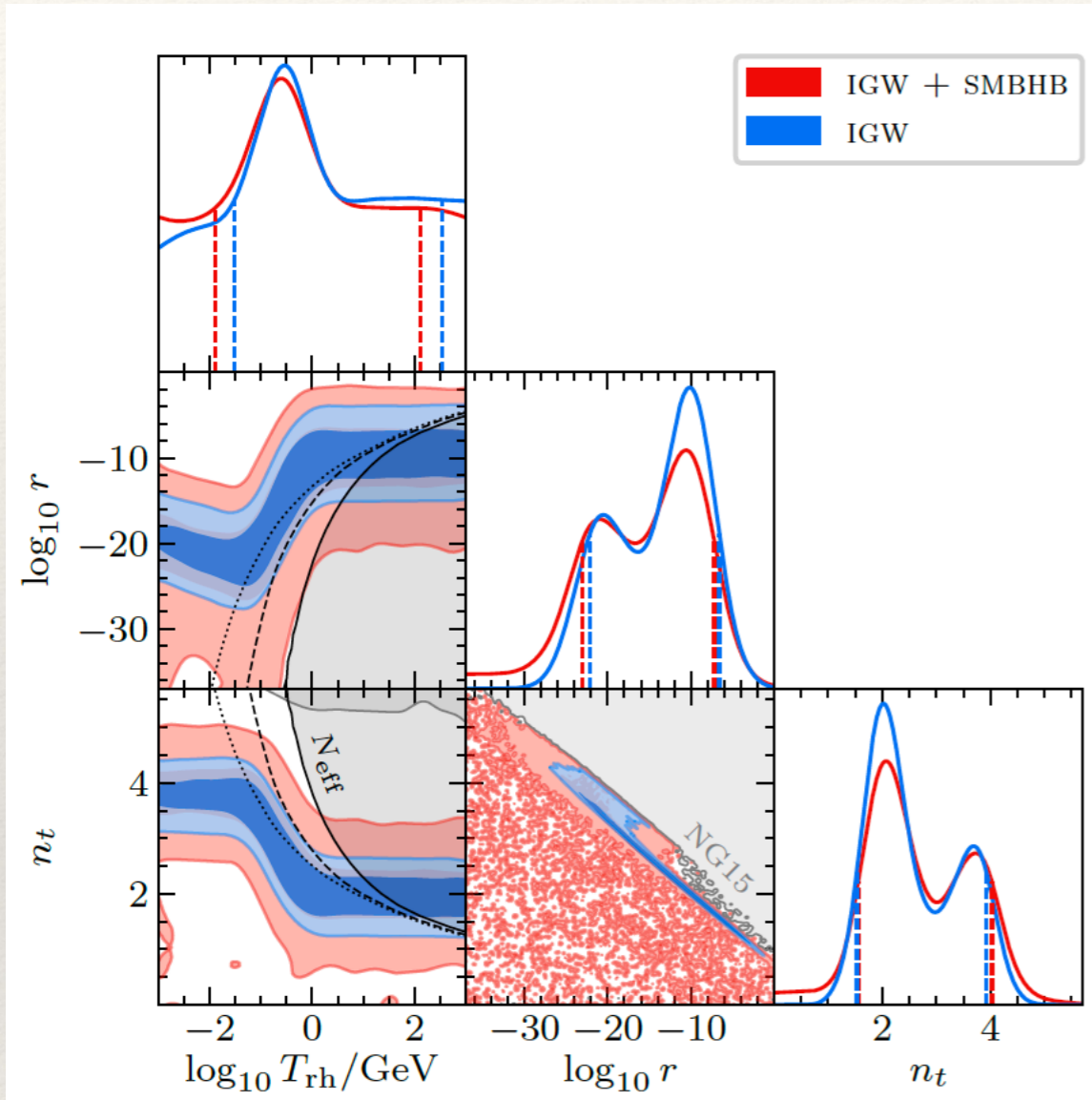
Both indicate efficient orbital decay



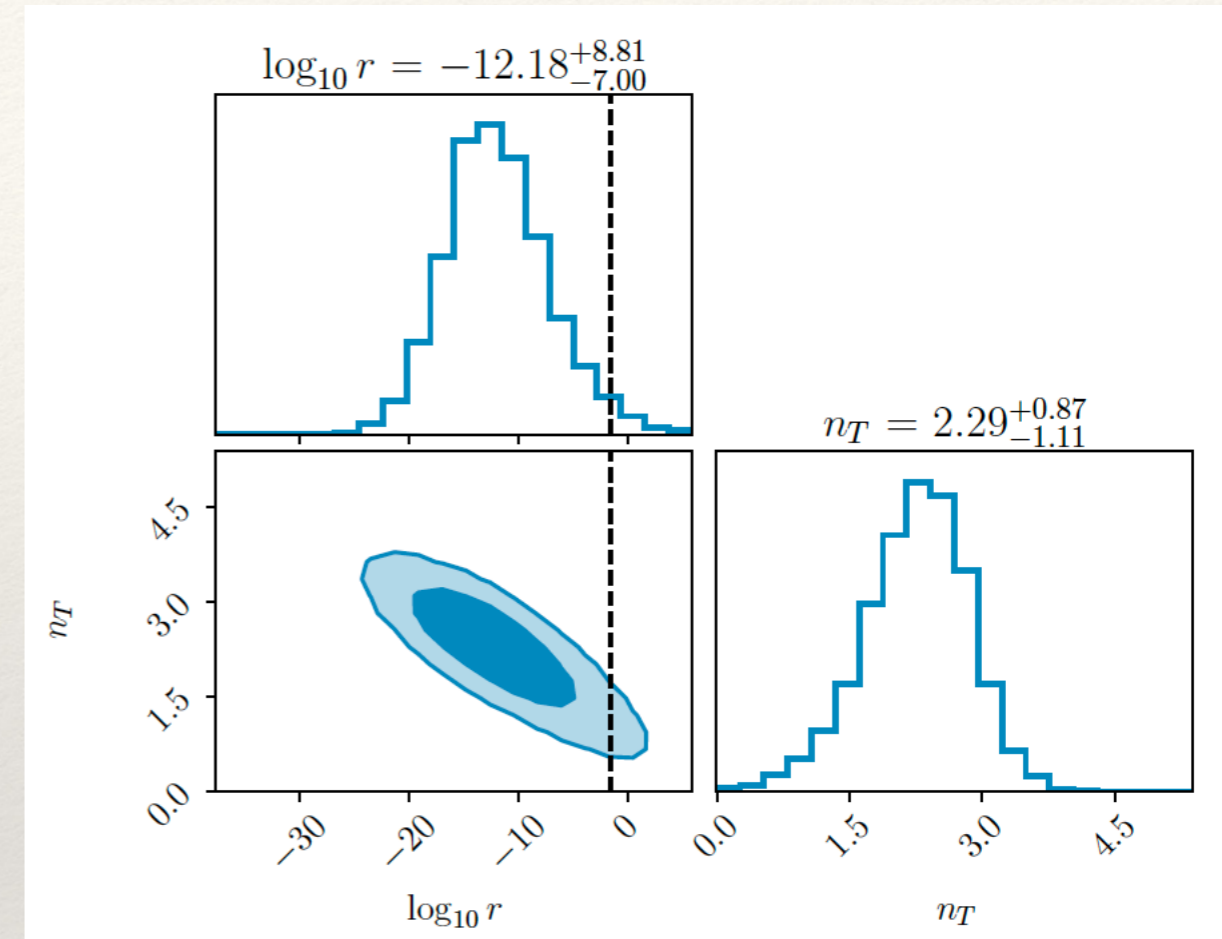
Relic SGWB



[NG, 2306.16219]



[EPTA, 2306.16227]



- agnostic about the microphysics of inflation and restrict ourselves to a model-independent analysis.
- the tensor-to-scalar ratio r and tensor spectral index n_t at the CMB pivot scale, reheating temperature T_{rh}

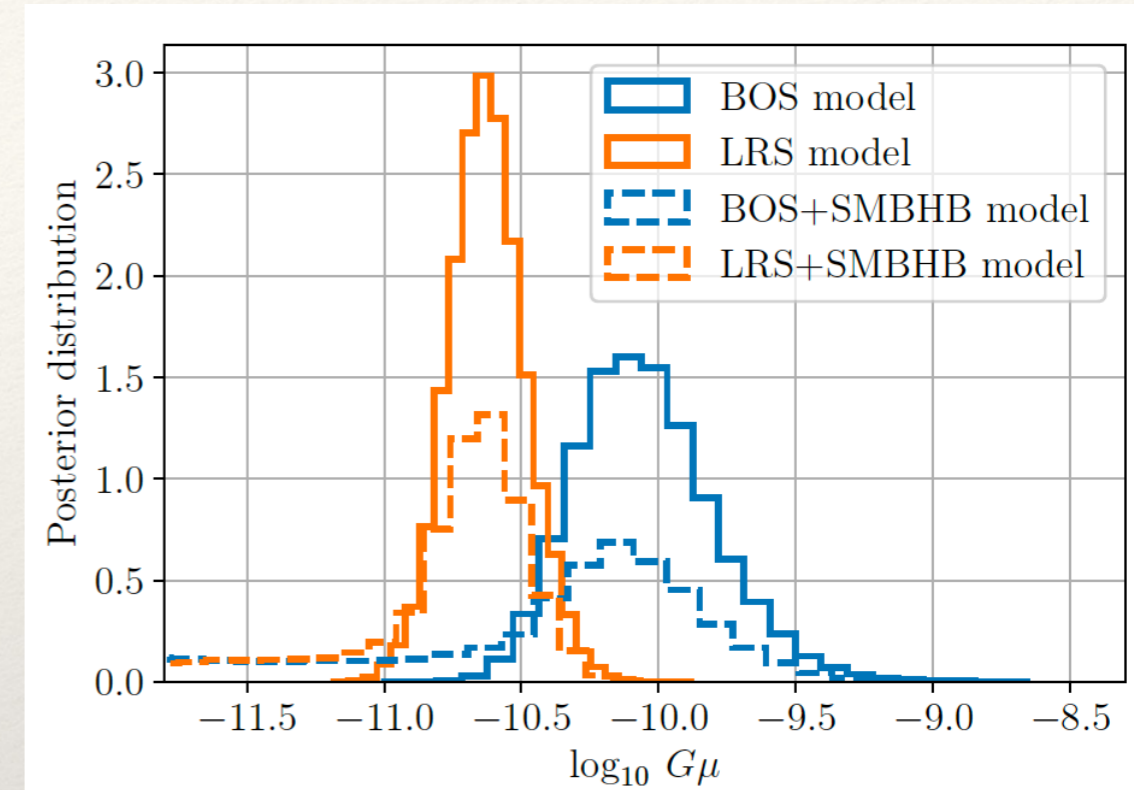
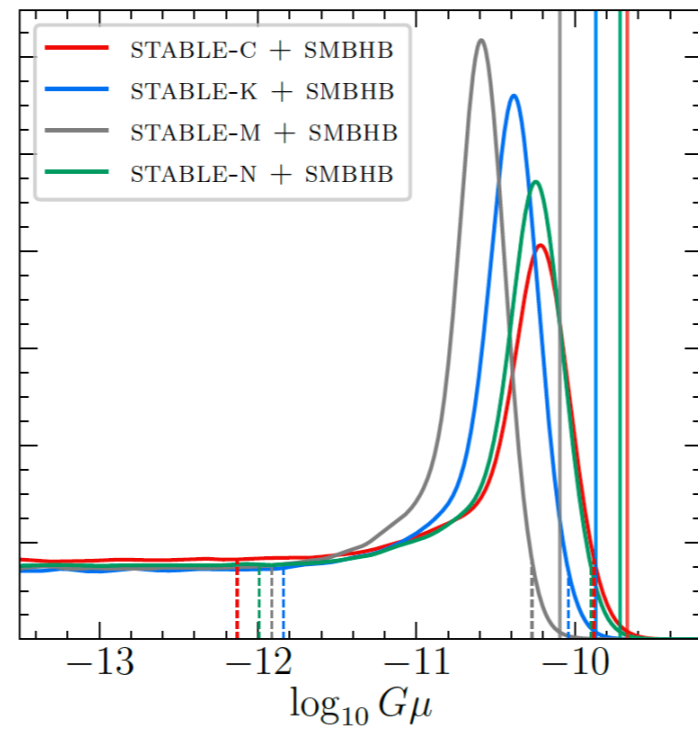
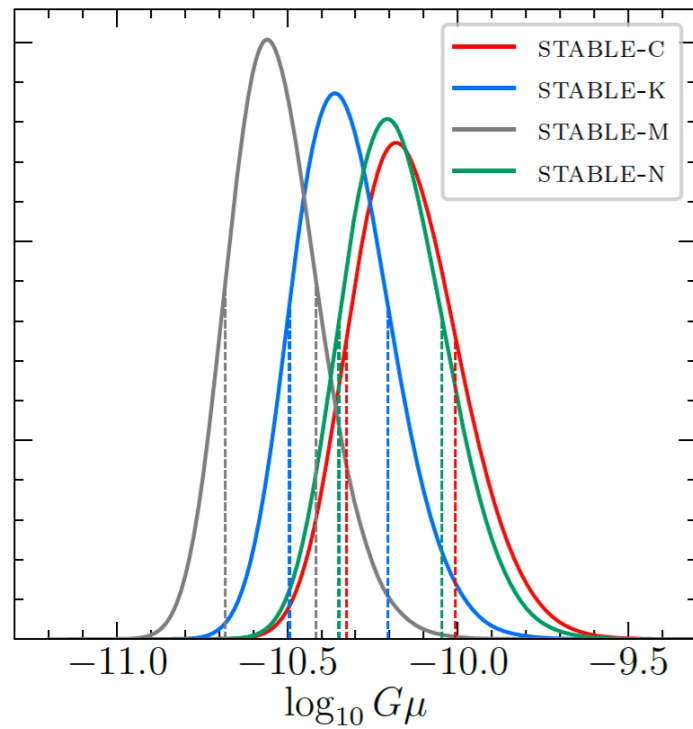


SGWB: Network of cosmic strings



[NG, 2306.16219]

[EPTA+InPTA, 2306.16227]



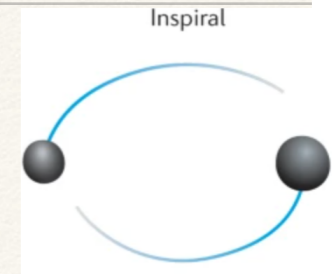
- We can constrain the tension of cosmic strings (model dependent) assuming the observed signal is entirely produced by the network of cosmic strings
- We can set an upper limit in two-component model of SGWB: CSs + SMBHBs



Search for individual MBHBs: continuous GW signal



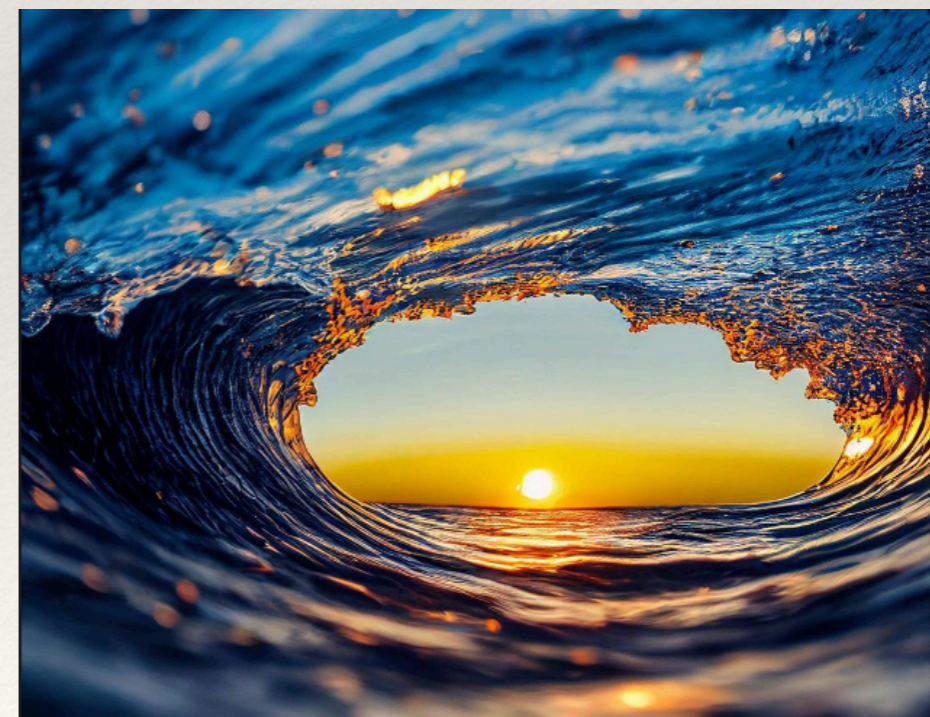
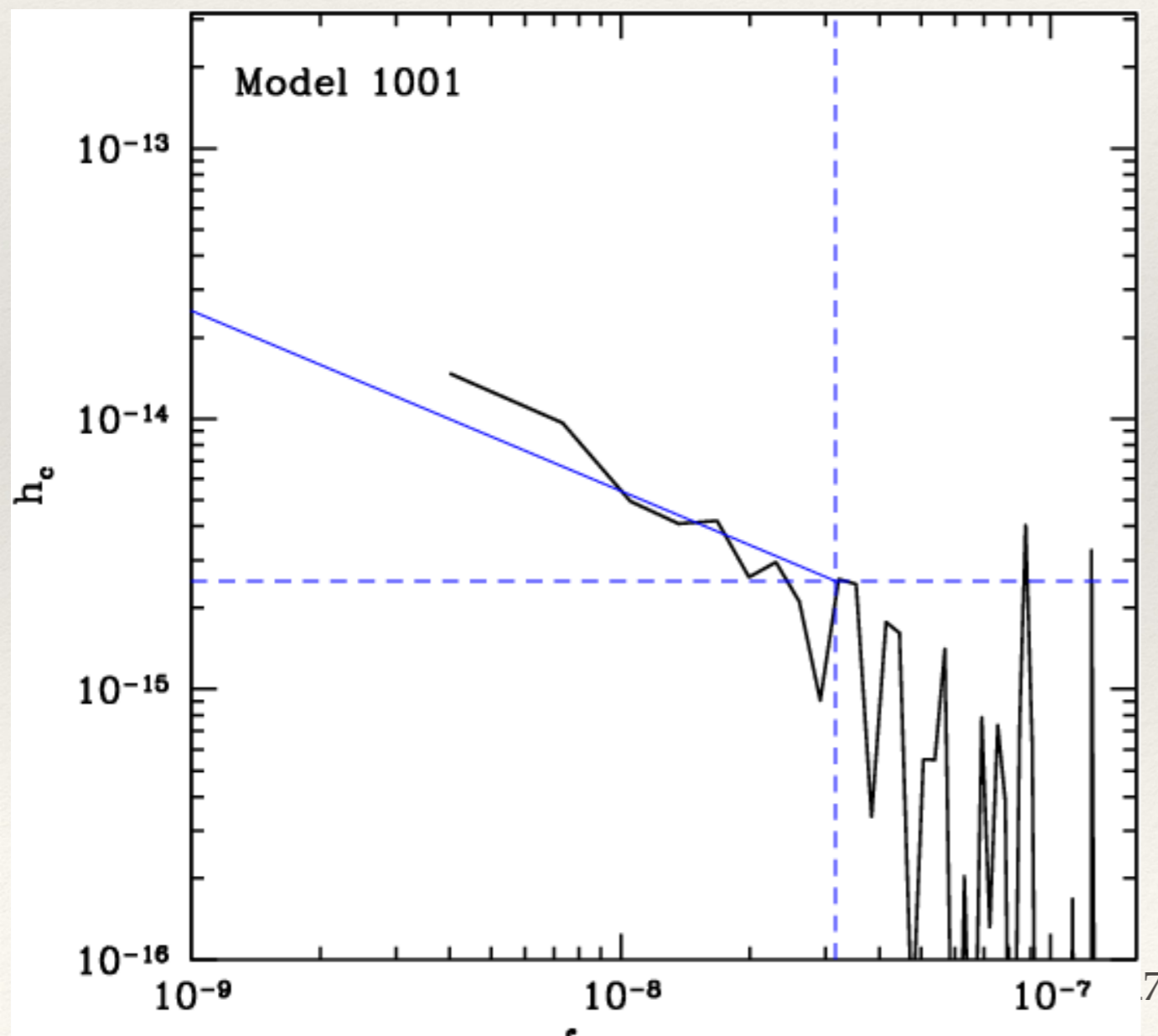
Searching for GW signal from individual SMBHB binary:



- Assume circular orbit
- Bayesian approach
- Strategy: all-sky search with simplistic model -> follow up candidates relaxing simplified assumptions on the reduced prior range

[NG: 2306.16222]

[EPTA+InPTA: 2306.16226]



CGW signal

Consider non-spinning SMBH binary in circular orbit

- pulsar and earth terms: each is monochromatic signal
- frequency. of pulsar term might or might not coincide with the earth term:
 $t_p = t - L(1 + \hat{n} \cdot \hat{k})$
- amplitude of the pulsar term is larger: $\sim \omega^{-1/3}$

$$s_\alpha = F_\alpha^+(\hat{k}, \hat{n}_\alpha) \left[\frac{h_+(t_p^\alpha, \omega_\alpha)}{2\pi f_\alpha} - \frac{h_+(t, \omega)}{2\pi f} \right] + \alpha - \text{pulsar index}$$

$$F_\alpha^\times(\hat{k}, \hat{n}_\alpha) \left[\frac{h_\times(t_p^\alpha, \omega_\alpha)}{2\pi f_\alpha} - \frac{h_\times(t, \omega)}{2\pi f} \right]$$

relative position pulsar and GW source

Pulsar term

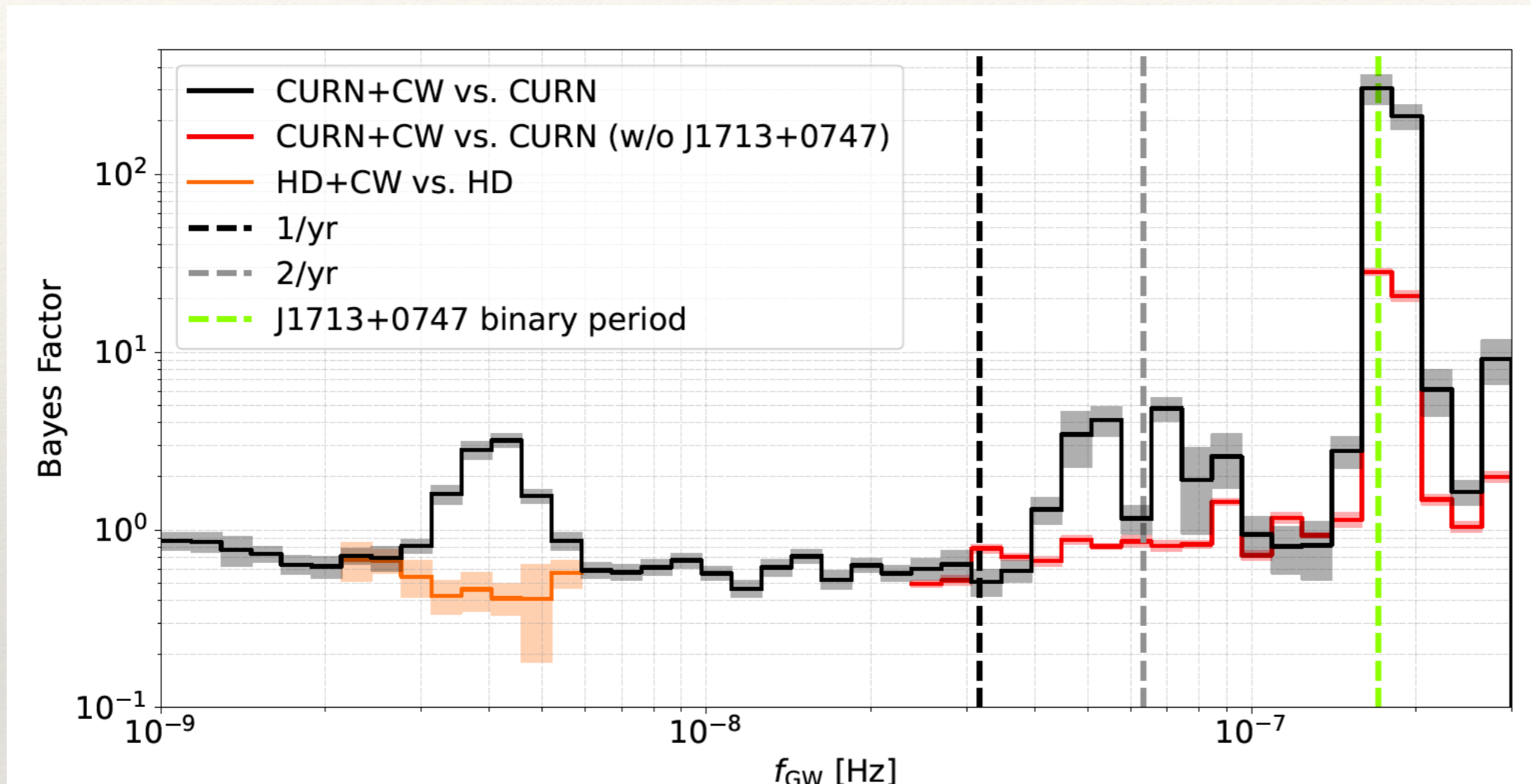
Earth term coherent across pulsars

$$\omega_\alpha = \omega(t - L_\alpha(1 + \hat{n}_\alpha \cdot \hat{k}))$$

CGW signal in NanoGrav



[NG: 2306.16222]



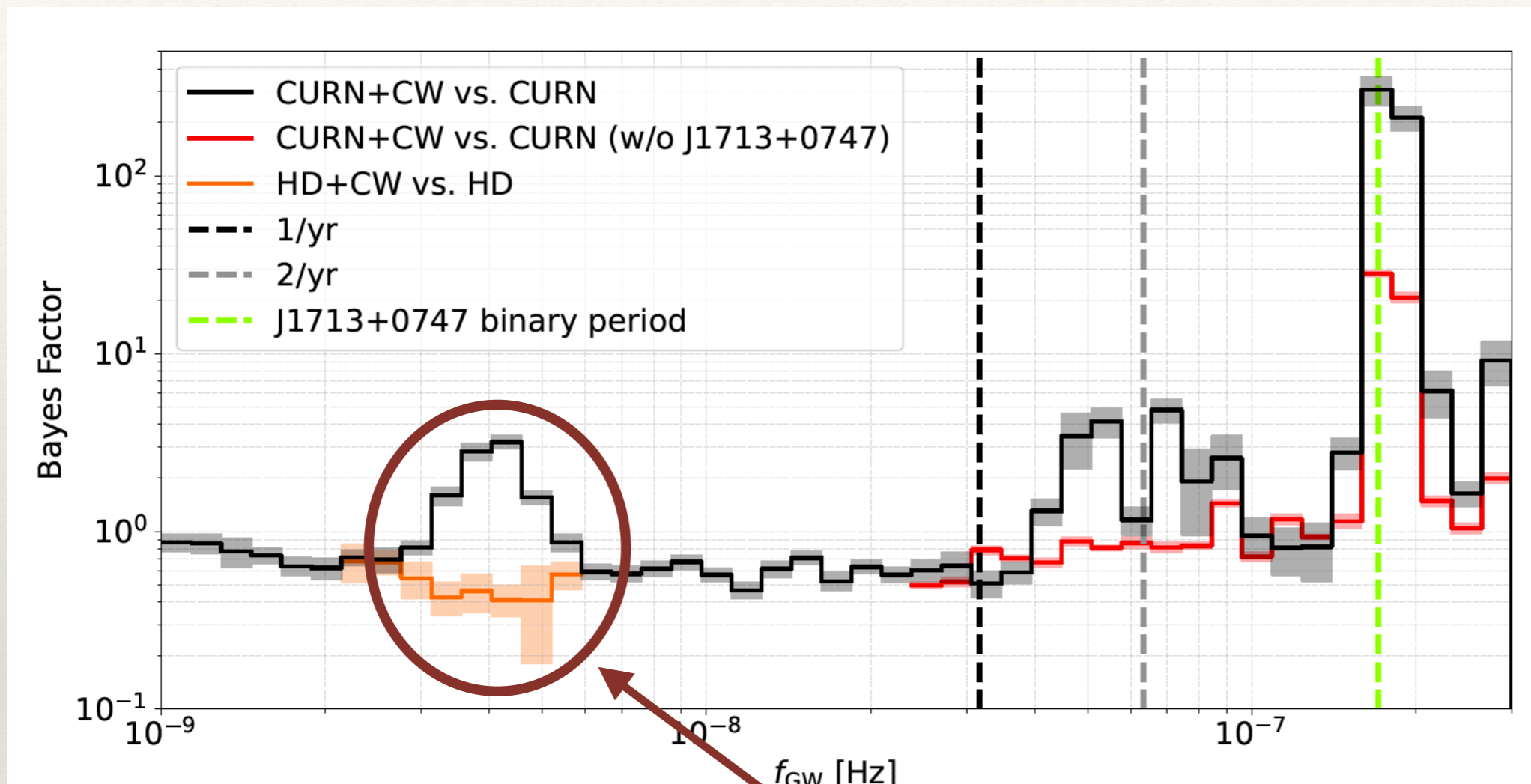
Bayesian all-sky search for a SMBHB in circular orbit: Bayes factor for presence of CGW



CGW signal



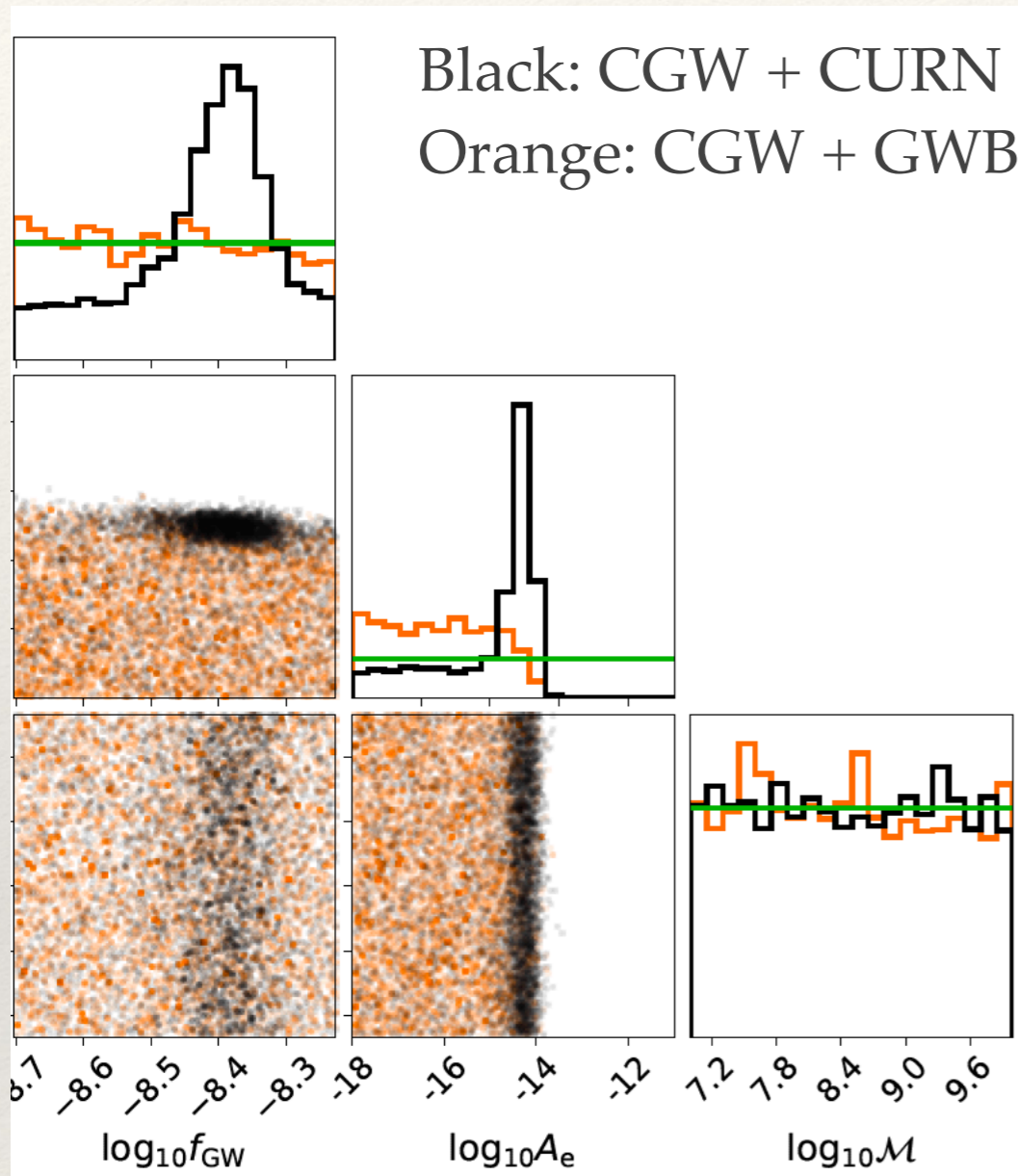
[NG: 2306.16222]



- Bayesian all-sky search for a SMBHB in circular orbit: Bayes factor for presence of CGW
- I concentrate on the low-frequency candidate

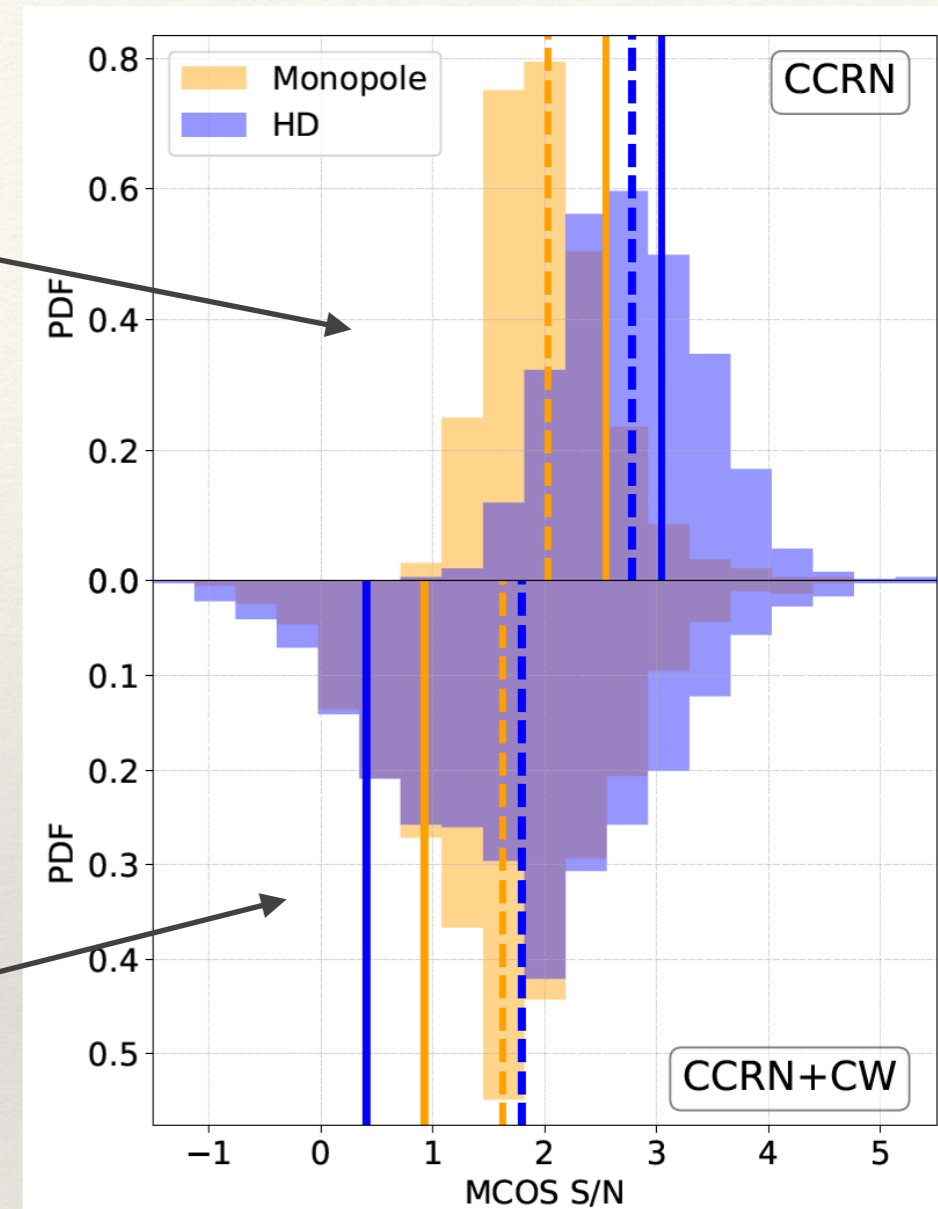


CGW in NG data

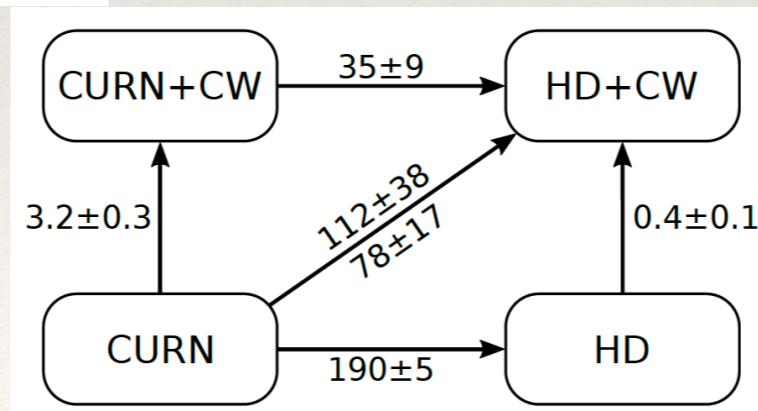


SNR of GWB and monopole in NG data

SNR of GWB and monopole in NG data after removing CGW



$f \in (3.2, 5.0)$ nHz

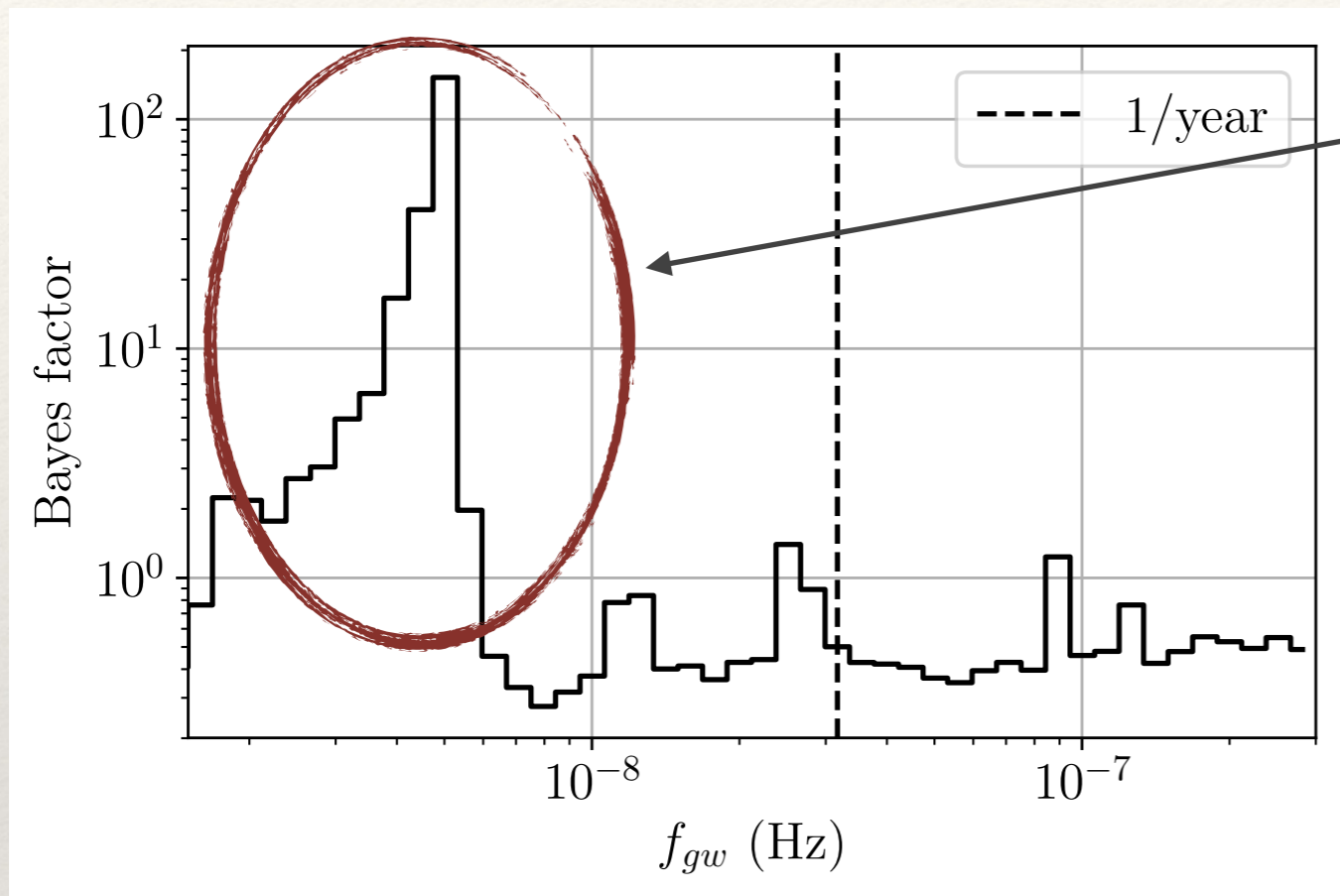


CGW signal in EPTA



BF: PSRN+CURN+CGW / PSRN+CURN

[EPTA+InPTA: 2306.16226]

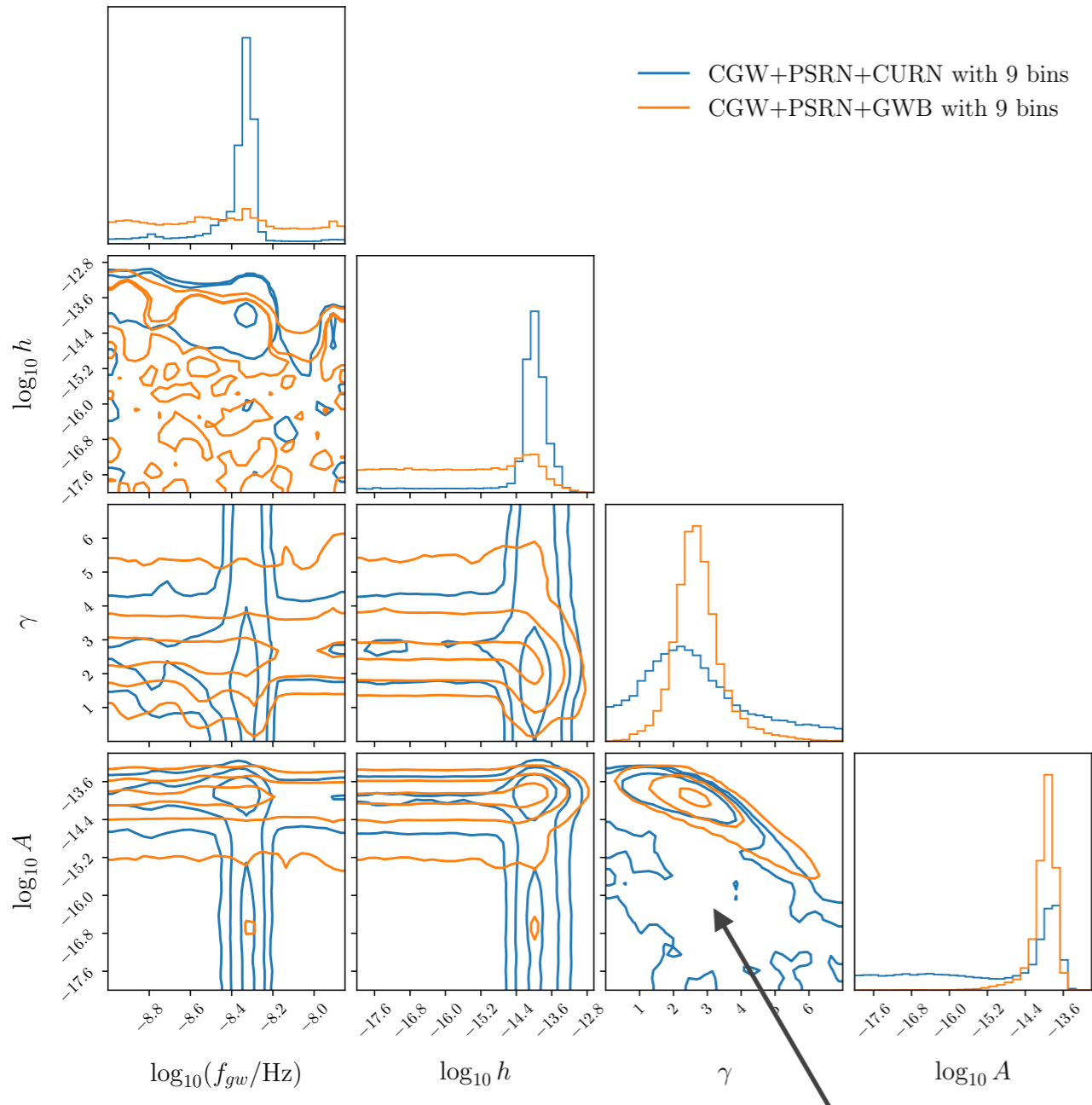


CGW candidate at low frequency

Bayesian search for CGW using Earth term only

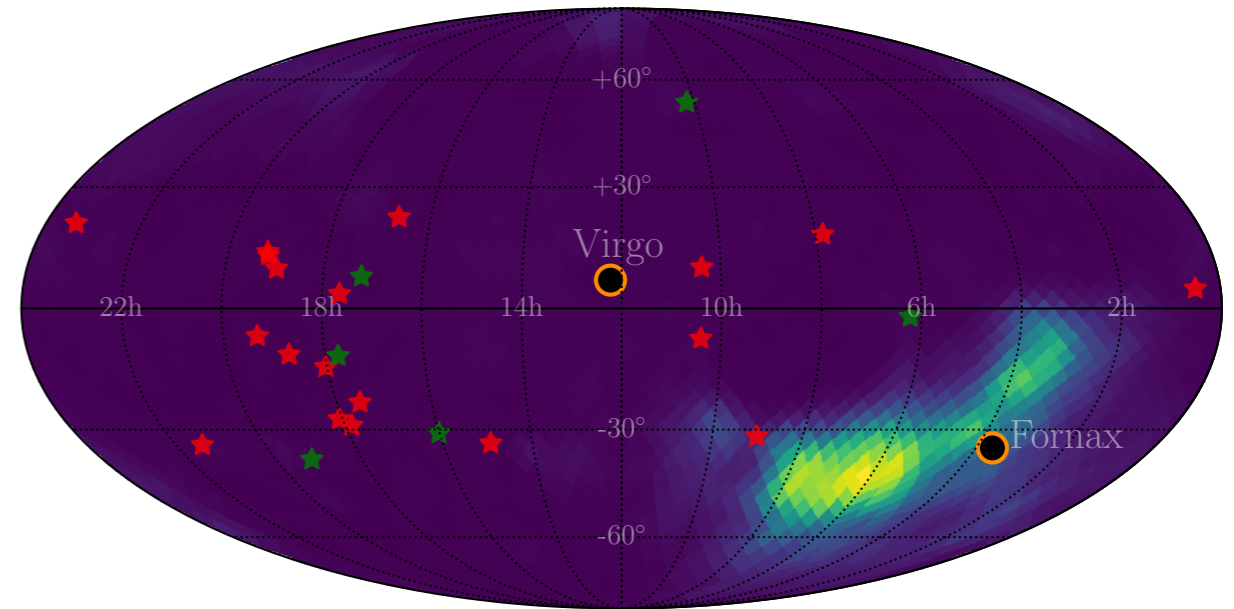


CGW signal in EPTA



$f \in (3.2, 6.0)$ nHz

CURN poorly constrained

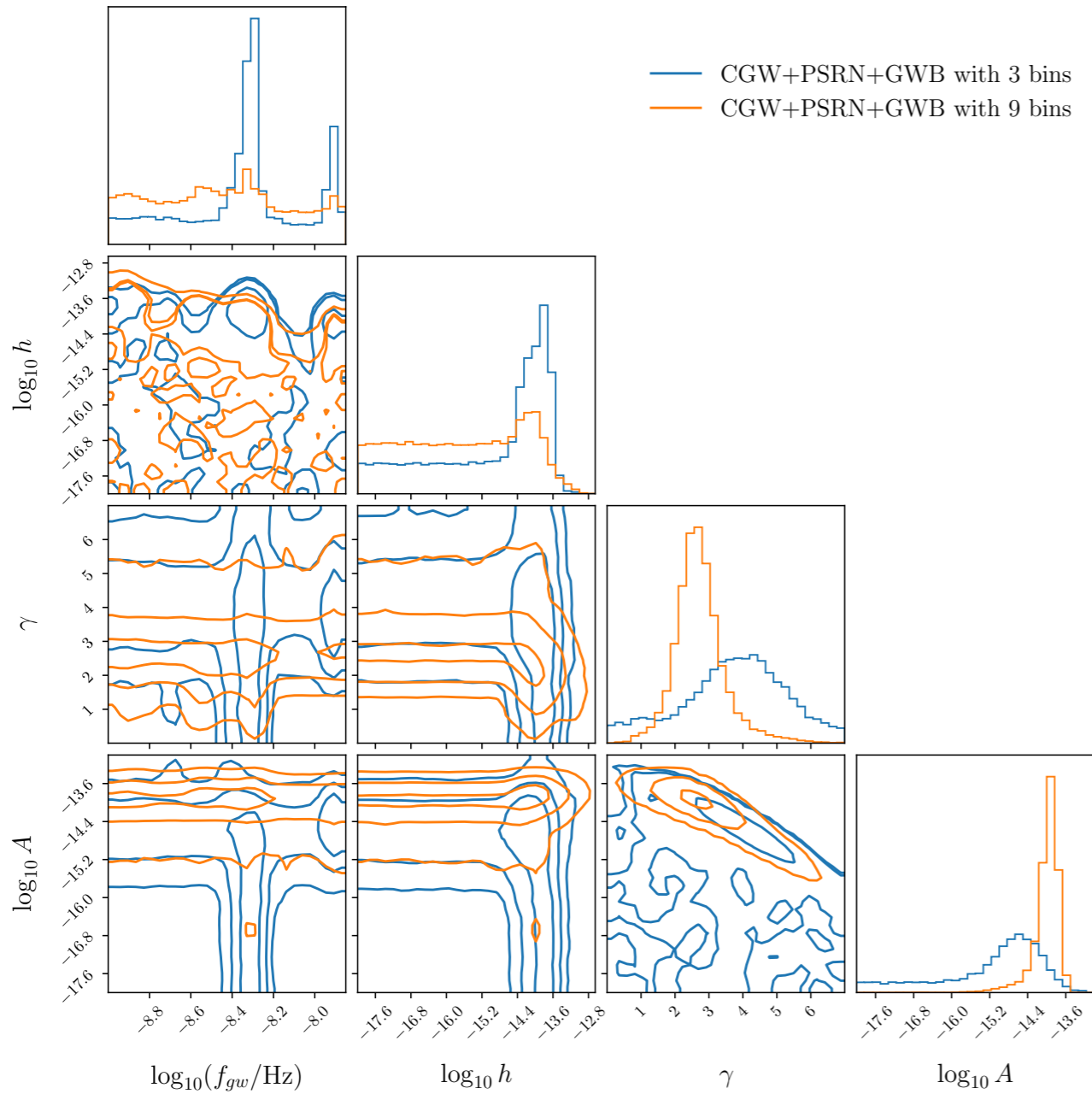


CGW: circular, Earth and Pulsar terms

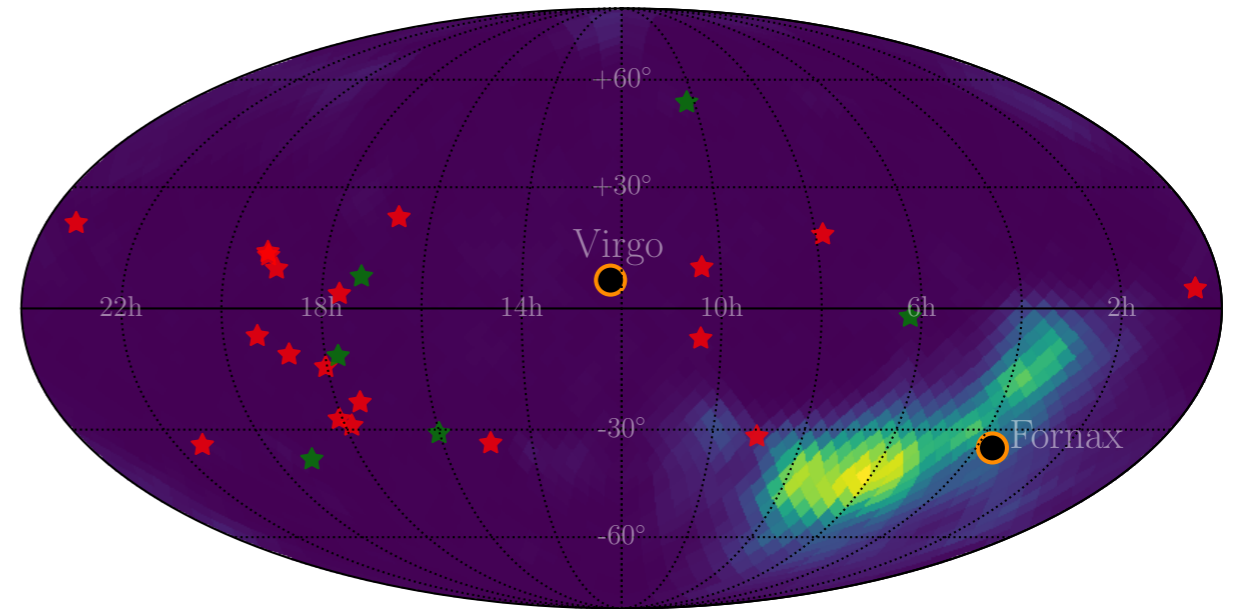
Model comparison	Bayes factor
CGW+PSRN vs PSRN	4000
CGW+PSRN+CURN vs PSRN+CURN, 3 bins	12
CGW+PSRN+CURN vs PSRN+CURN, 9 bins	4
CGW+PSRN+GWB vs PSRN+GWB, 3 bins	1
CGW+PSRN+GWB vs PSRN+GWB, 9 bins	0.7



CGW signal in EPTA



$f \in (3.2, 6.0)$ nHz

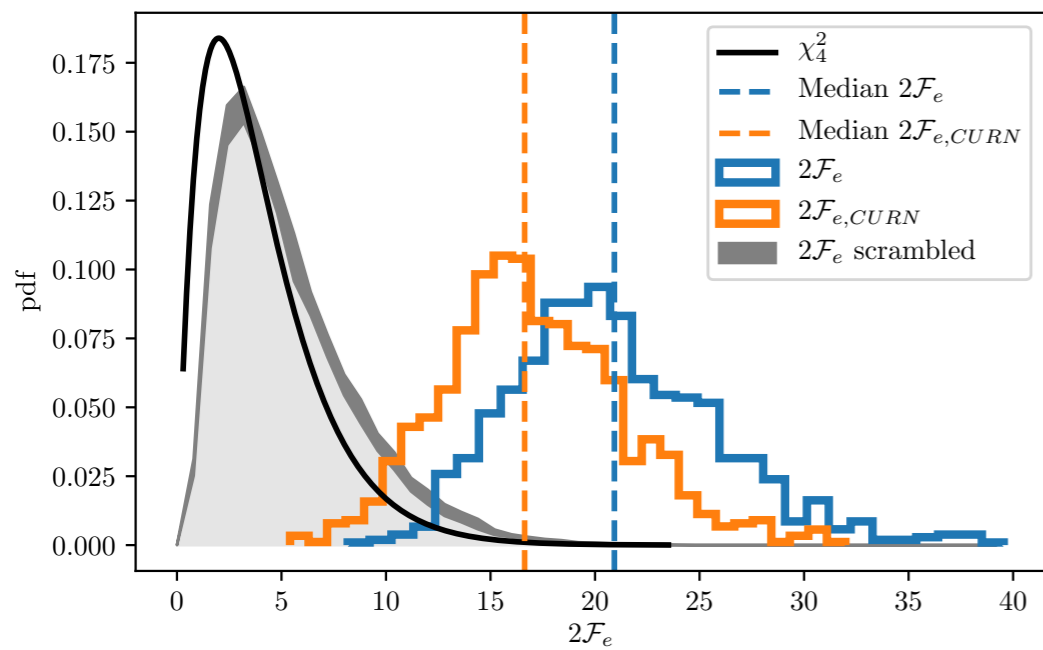


CGW: circular, Earth and Pulsar terms

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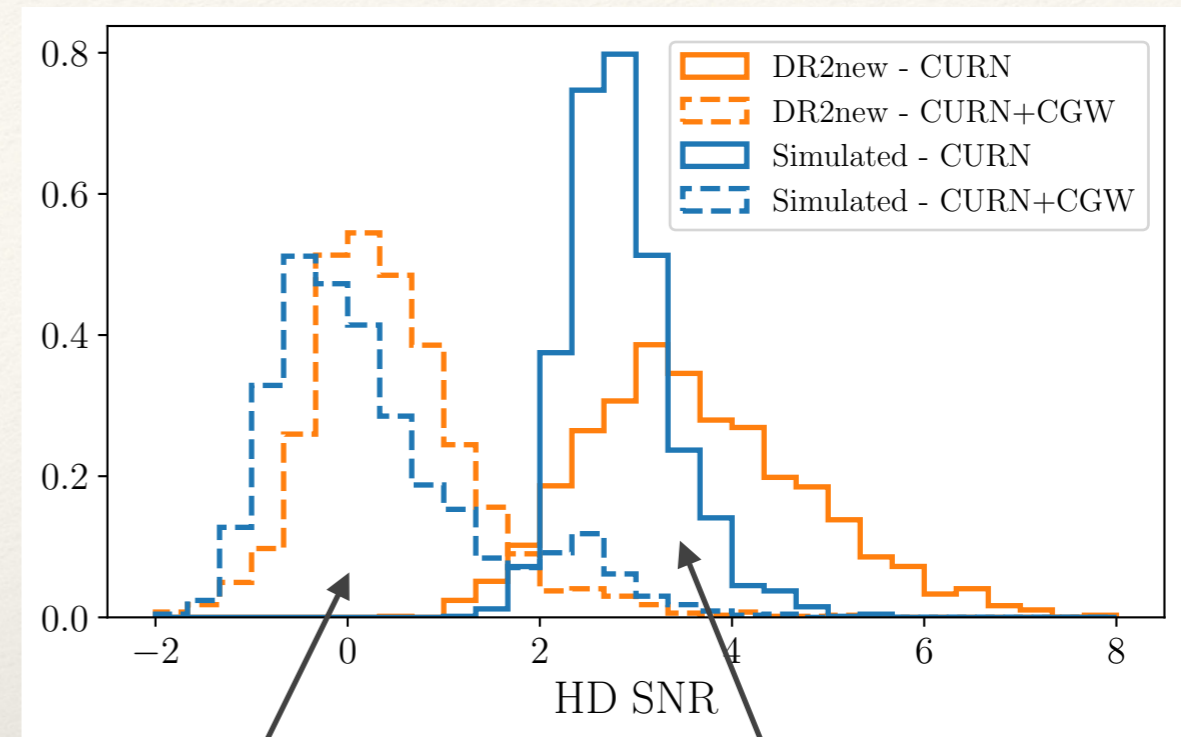
CGW signal in EPTA



Frequentist analysis (but taking into account large uncertainties in the noise)

Significance

	$p(\mathcal{F}_e)$	$p(\mathcal{F}_{e,CURN})$
χ^2_4	5×10^{-4}	1×10^{-3}
Sky scrambles	$(7 \pm 4) \times 10^{-4}$	$(6 \pm 1) \times 10^{-3}$



GWB SNR after removing CGW

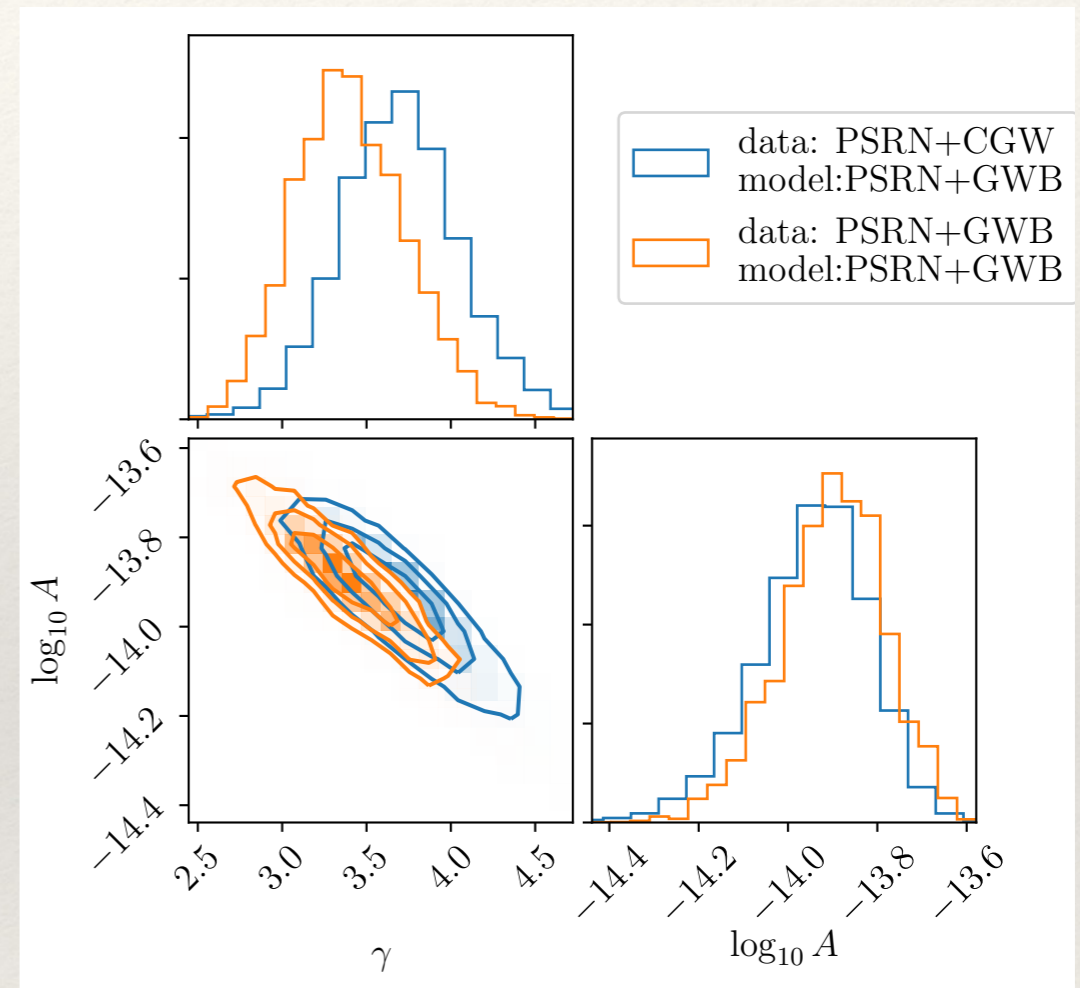
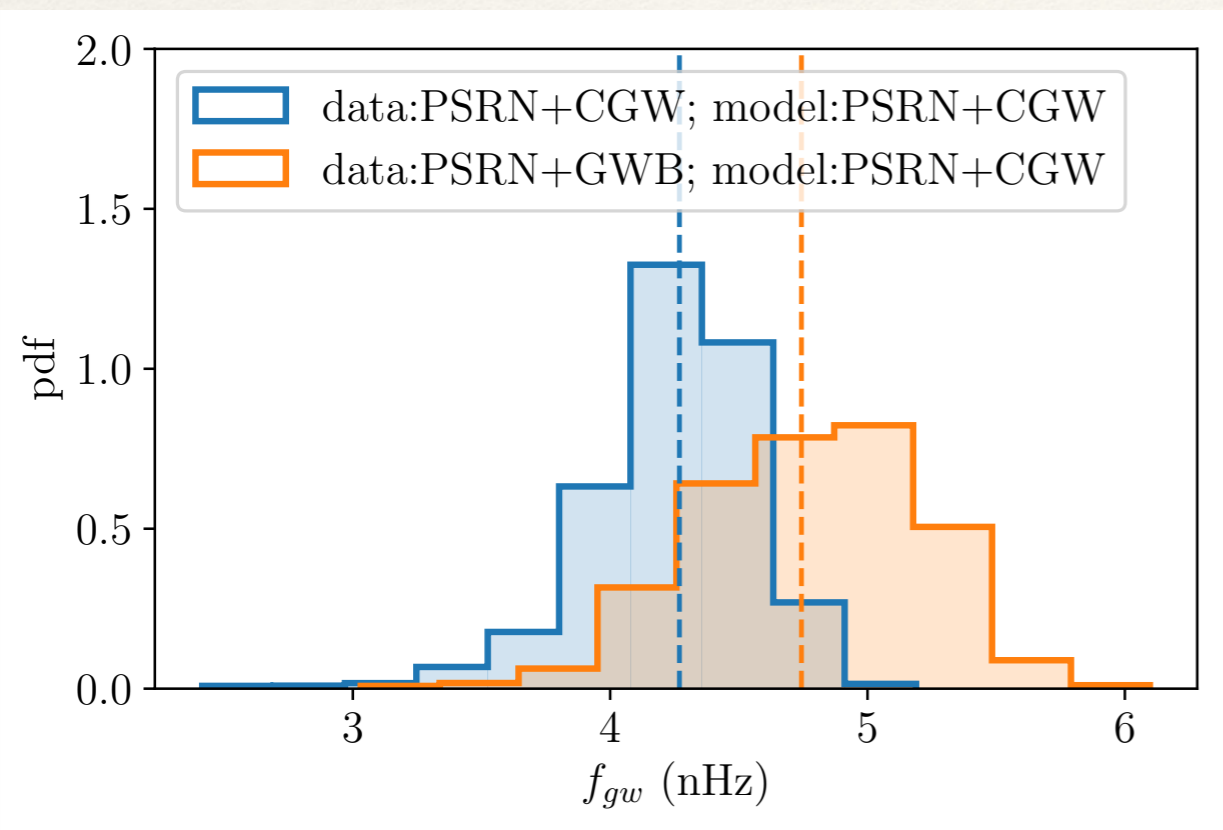
GWB SNR




CGW signal in PTA?



Simulated data: PSRN + GWB only, Model_1: GWB, Model_2 CGW
 Simulated data: PSRN + CGW only, Model_1: GWB, Model_2 CGW



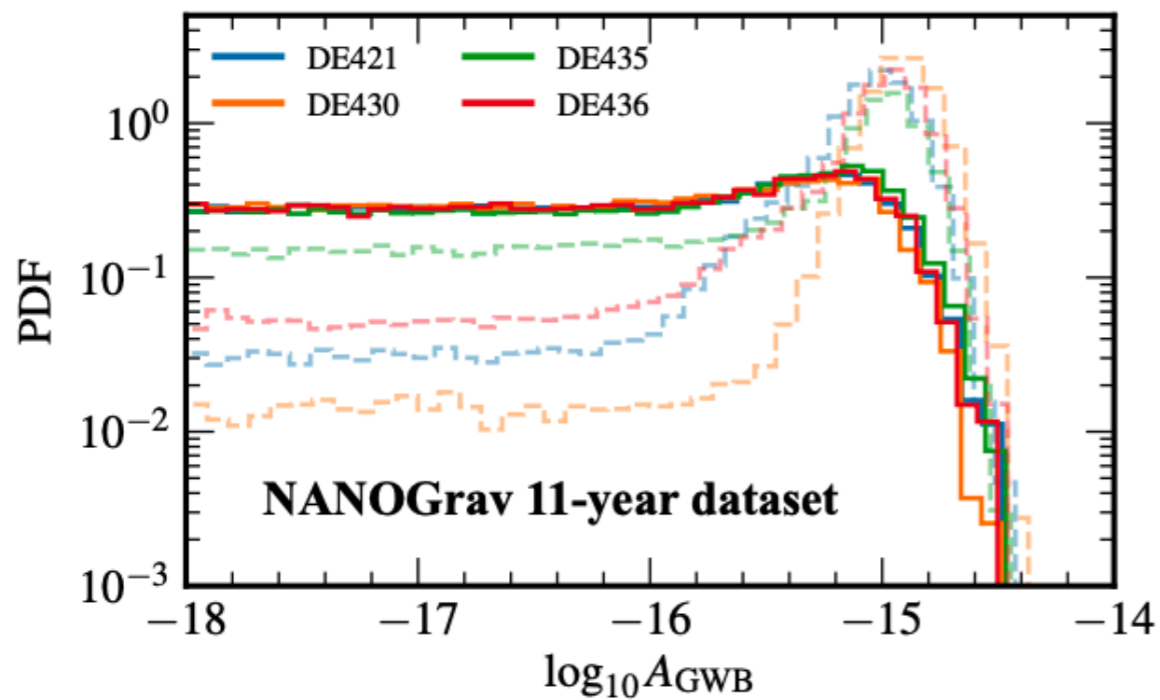
- It is easy to make mistake 
- However: GWB 2 parameters, CGW: N_p+8 pars



Is it really GW signal?



- Error in ephemerids: JPL ephemerids D440, good measurement of Jupiter



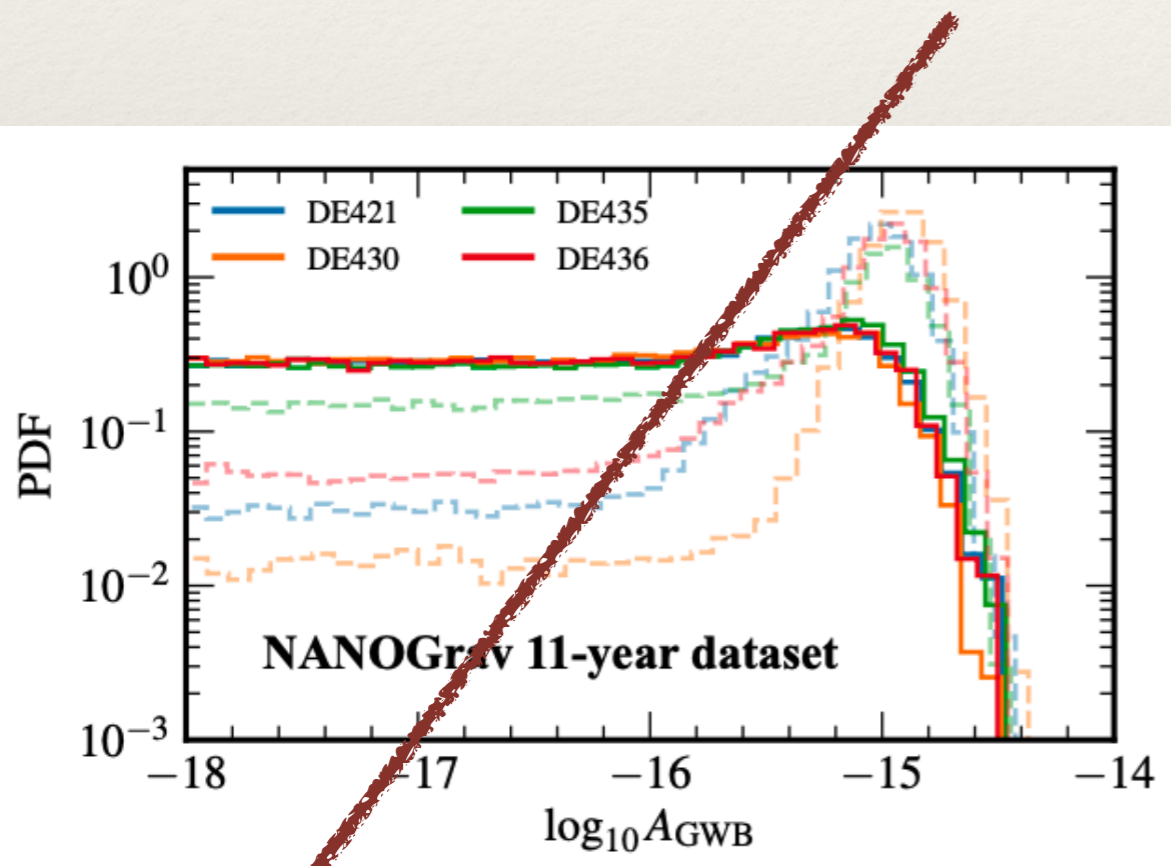
[Arzoumanian+ 2018]



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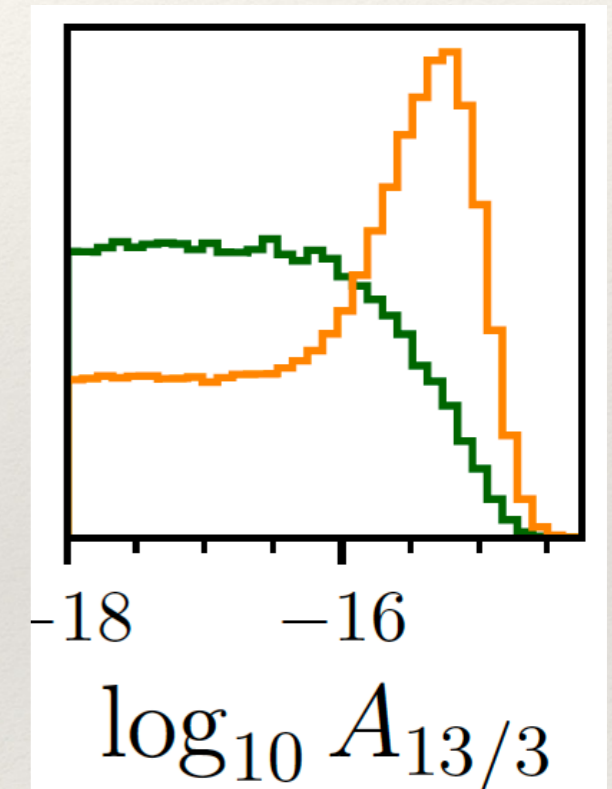
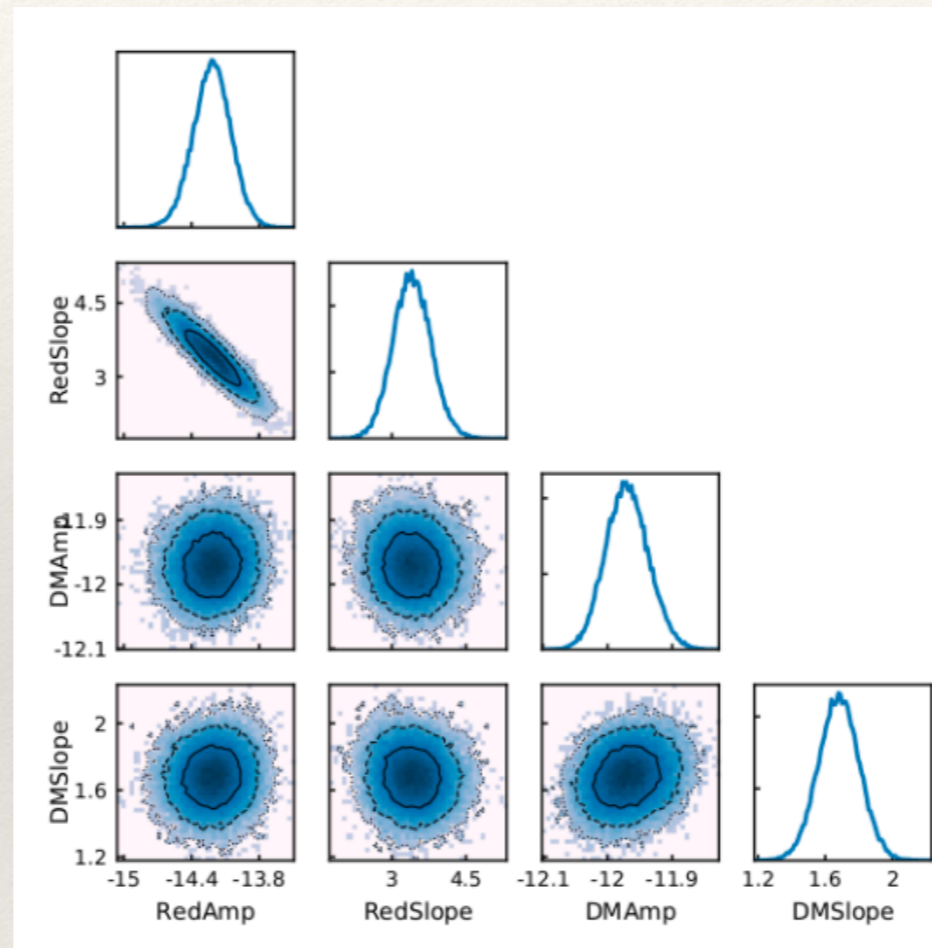
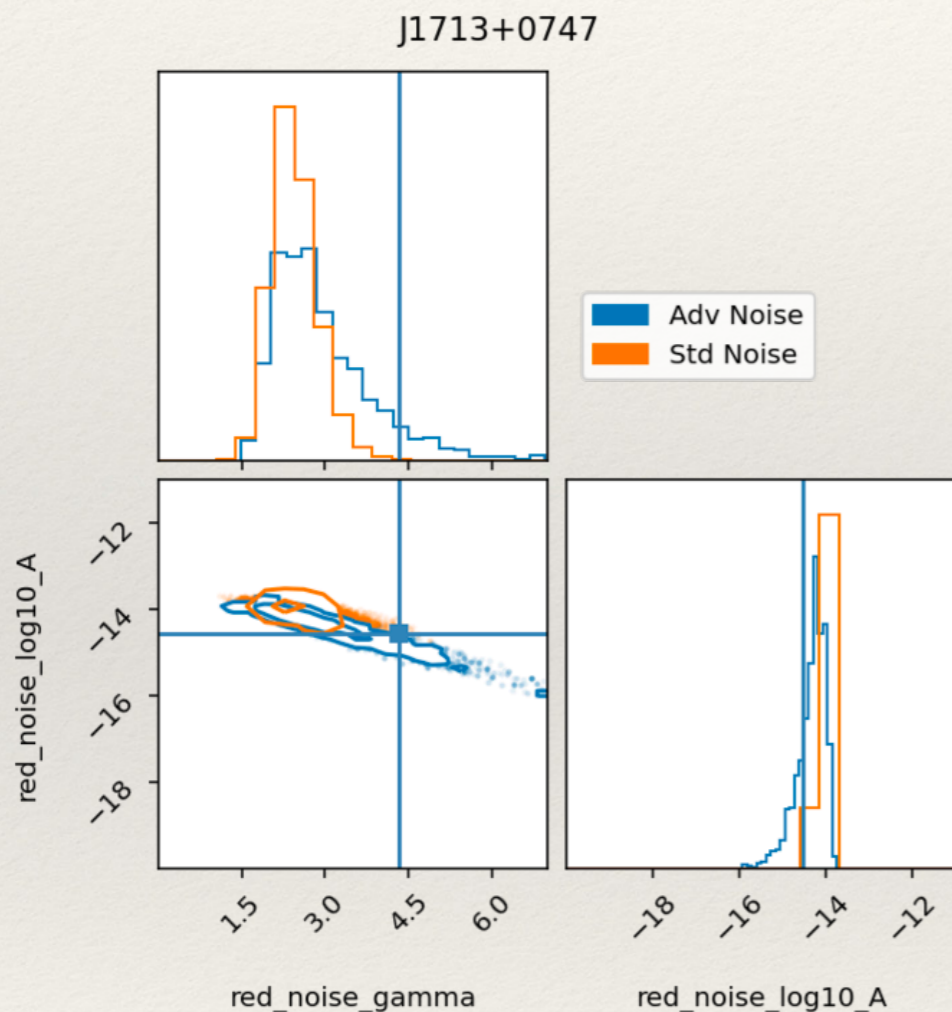


[Arzoumanian+ 2018]



Is it really GW signal?

- Error in ephemerids: JPL ephemerids D440, good measurement of Jupiter
- Modelling noise of each pulsar is very important: J1713+0747



Is it really GW signal?



- Error in ephemerids: JPL ephemerids D440, good measurement of Jupiter
- Modelling noise of each pulsar is very important: J1713+0747

Pulsar	Sel. model
J0613-0200	<i>RN10 DMv30 DMv-SN_NUP_1.4</i>
J1012+5307	<i>RN150 DMv30 DMv-SN_NUP_1.4 SN_NUP_2.5</i>
J1600-3053	<i>DMv30 Sv150 SN_LEAP_1.4</i>
J1713+0747	<i>RN15 DMv150 2 Exp. dips DMv-SN_NUP_1.4 SN_JBO_1.5 SN_LEAP_1.4 SN_BON_2.0 BN_Band.3</i>
J1744-1134	<i>RN10 DMv100 DMv-SN_NUP_1.4 BN_Band.2</i>
J1909-3744	<i>RN10 DMv100 Sv150</i>

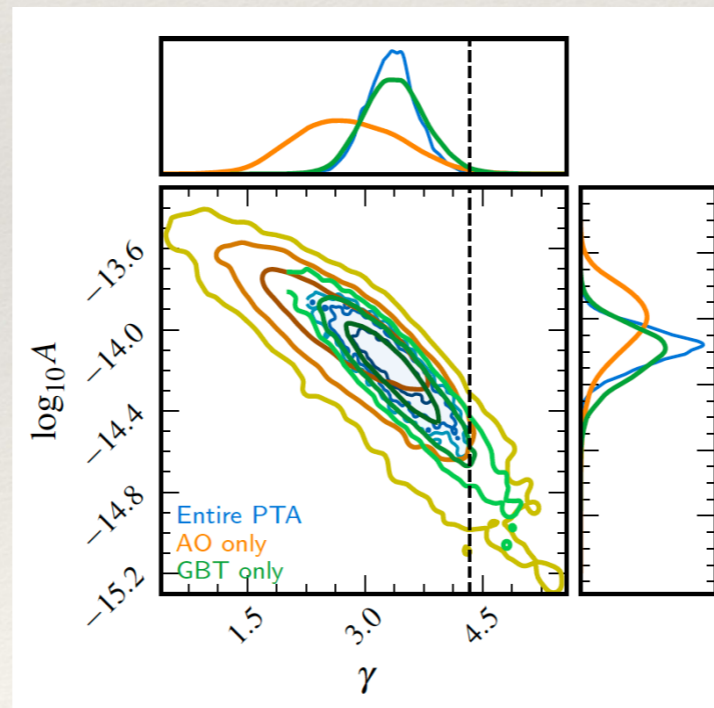
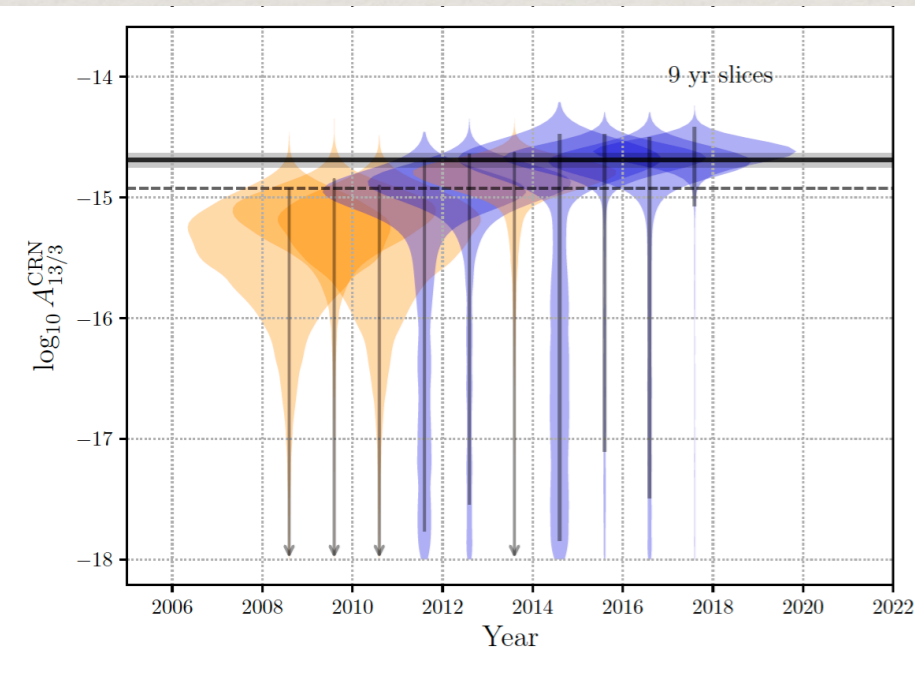
EPTA 6 best pulsars, custom noise models
[Chalumeau+ 2021]



Is it really GW signal?



- Error in ephemerids: JPL ephemerids D440, good measurement of Jupiter
- Modelling noise of each pulsar is very important: J1713+0747
- Quite different BF from each PTA: 1-2 (PPTA), 60-70 (EPTA), 230-950 (NG)
- EPTA “sees” the signal only in last 14 years, PPTA sees signs of non-stationarity
 - Is it non stationarity in the GWB?
 - or in the PSR noise model?
 - or evolution of how we deal with radioobservations?



Statistic	DR2full	DR2new	DRnew+
\mathcal{B}_{CURN}^{HD}	4	60 33(+Ephem)	65 43(+Ephem)



What's next?



IPTA data combination:

- We combine the data from IPTA: EPTA, NG, InPTA, PPTA
- We use additional data (MeerKAT, Chime)
- Better coverage (dense) in time (smaller cadence)
- Better coverage in radio freq: DM and scattering variations
- Not dominated by a single radiotelescope: should see/handle systematics

Kind of summary...

- We are pretty sure that the observed signal is GW
- We are not sure about its nature
- We got so excited that made a big press release
- In relativity we need to look at IPTA data, we need longer high quality data. It is "GW detection in slow motion"

