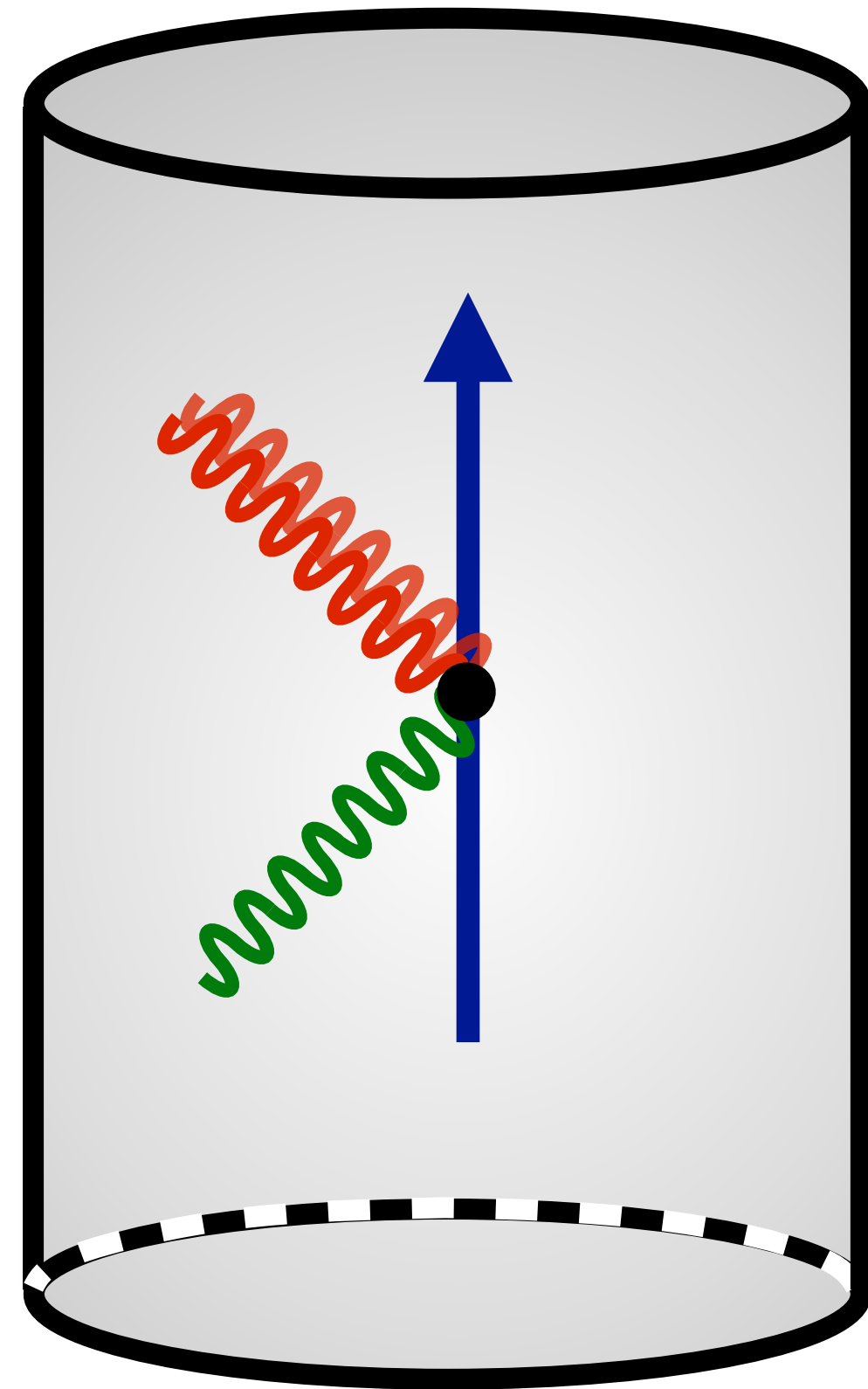

Resonant Cavities for Gravitational Waves

Sebastian A. R. Ellis

University of Geneva

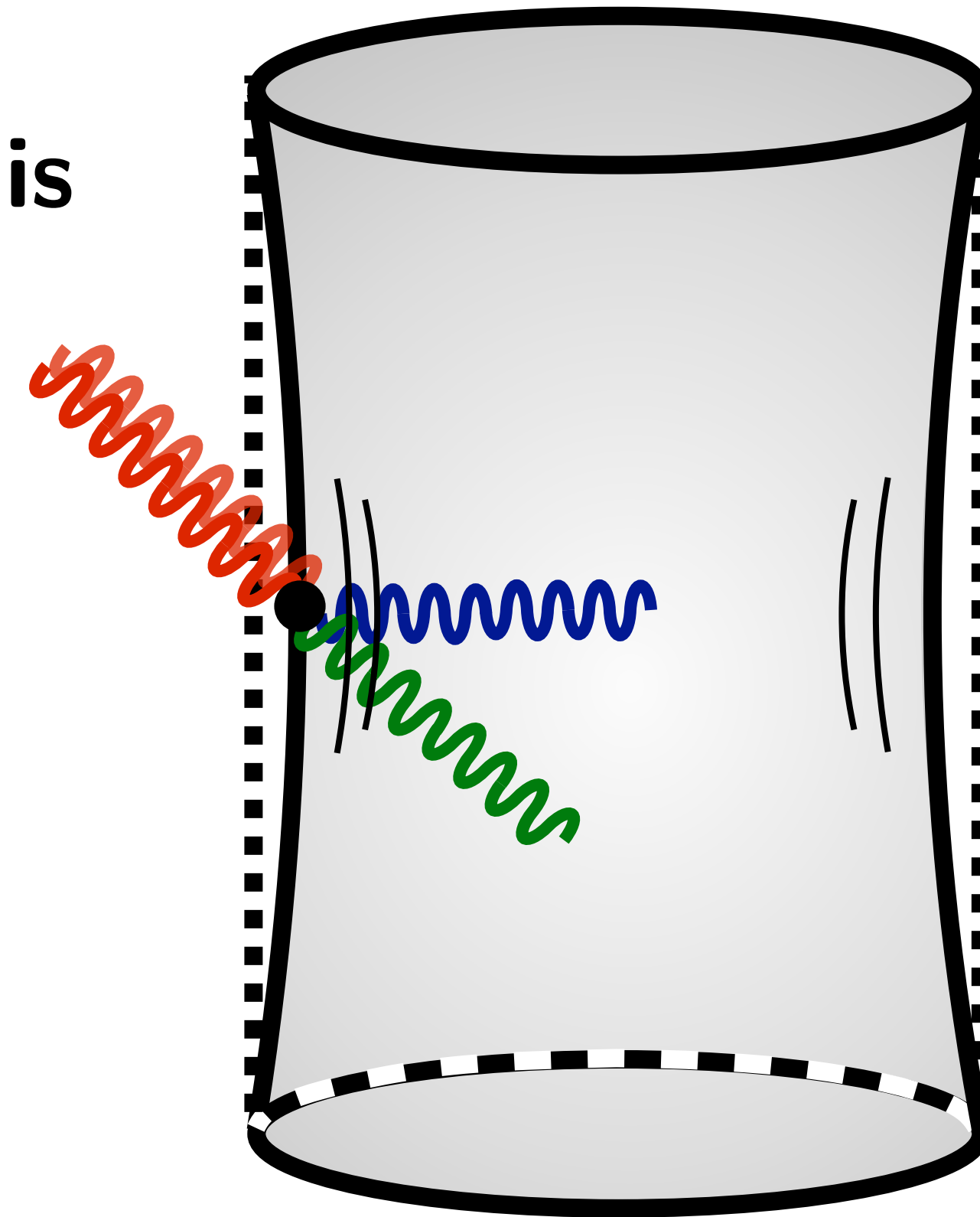
w/ A. Berlin, D. Blas, R. T. D'Agnolo, R. Harnik, Y. Kahn, J. Schütte-Engel & M. Wentzel
arXiv:2112.11465 arXiv:2303.01518

Resonant Cavities for Gravitational Waves



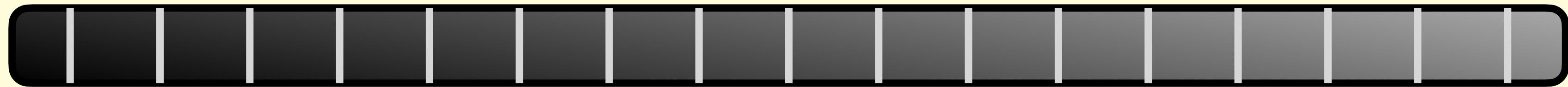
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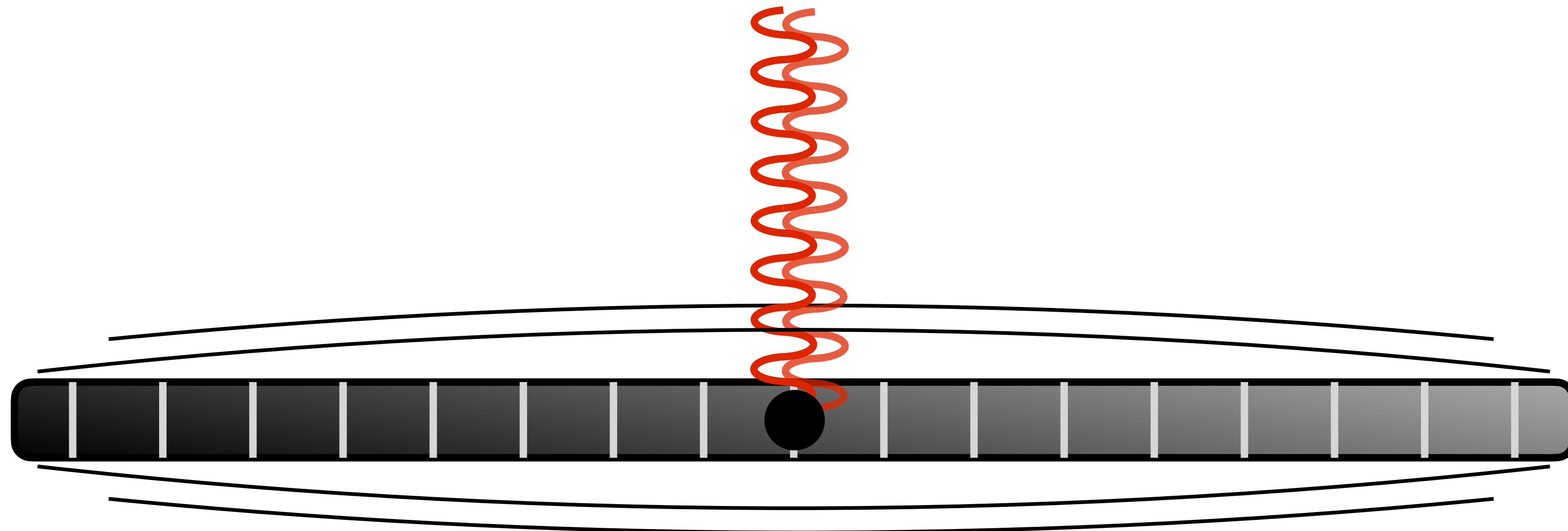
WARMUP



When is the ruler rigid?

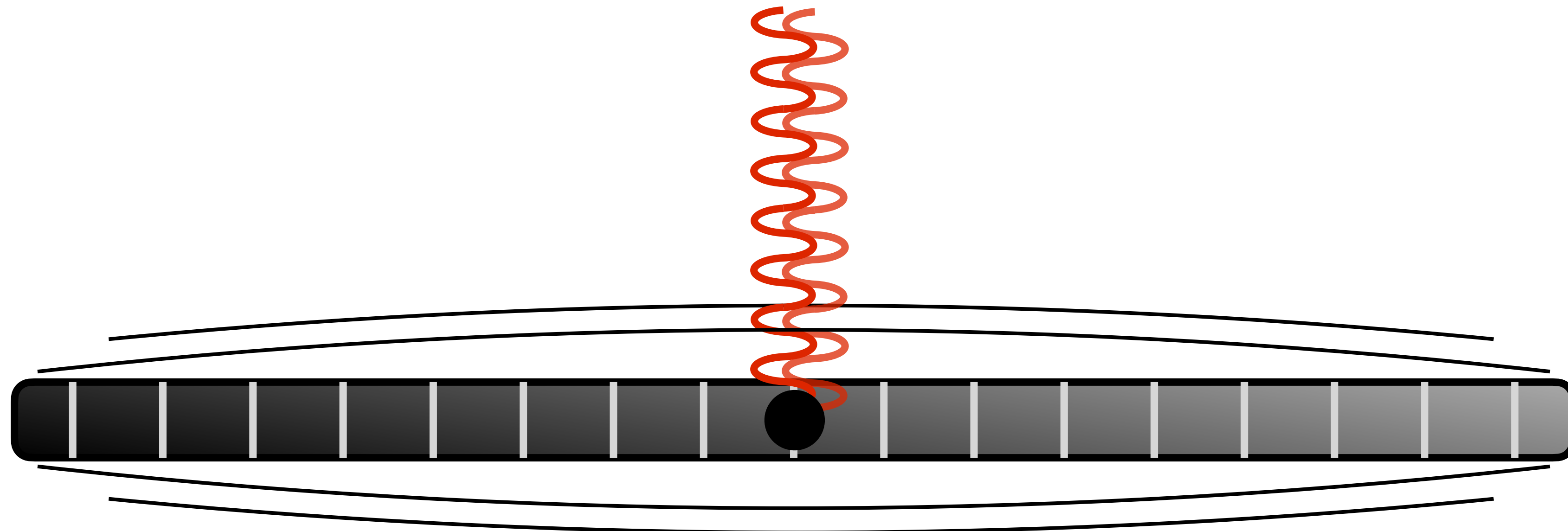
When is the ruler rigid?

Incoming GW: long wavelength



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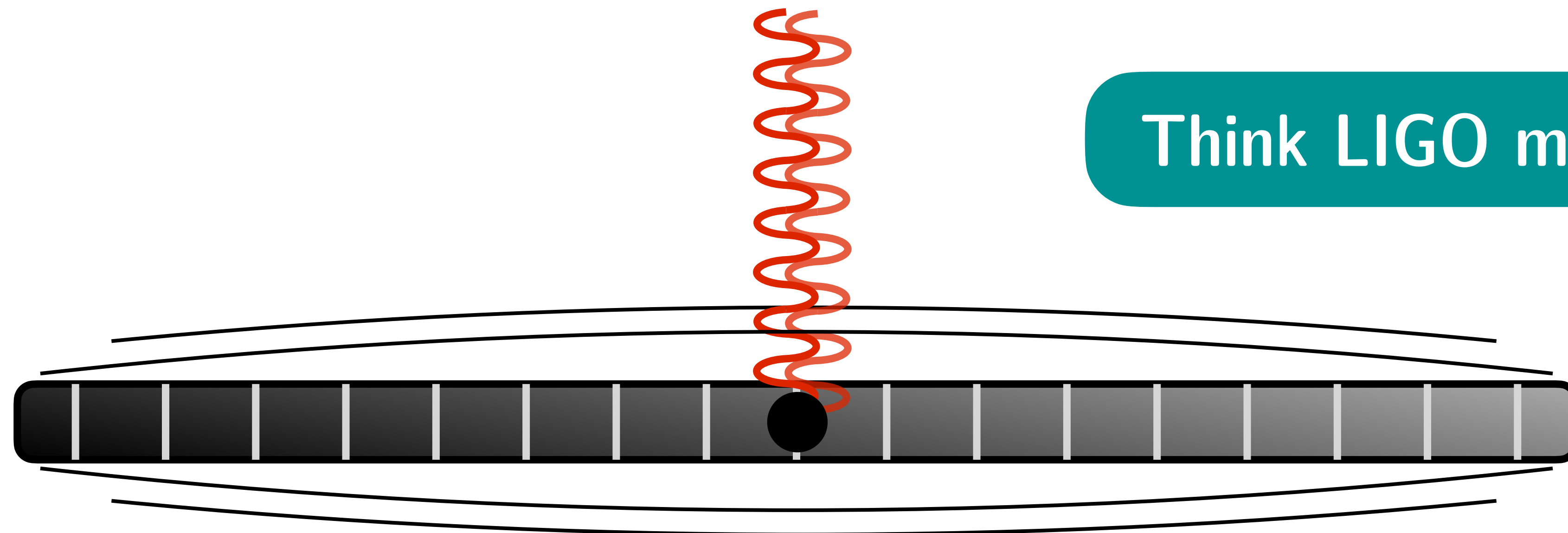
Incoming GW: long wavelength



The whole object moves back and forth

When is the ruler rigid?

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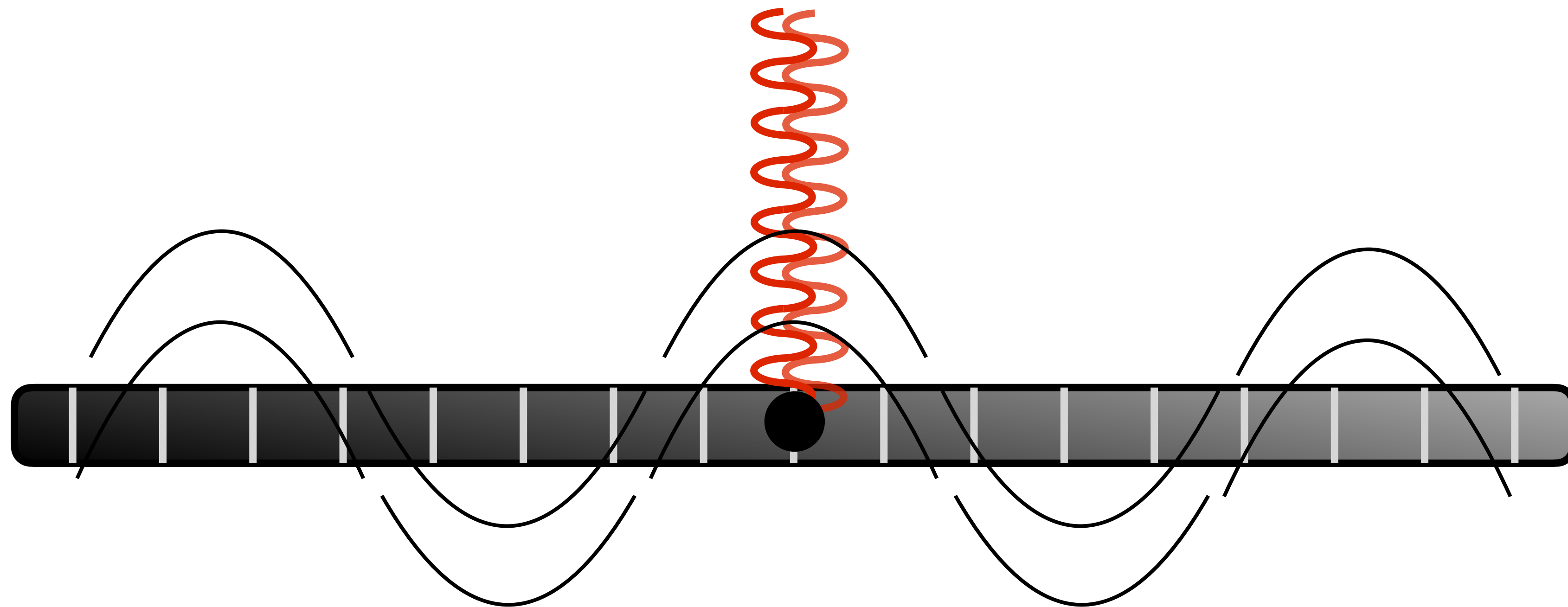


Think LIGO mirror

The whole object moves back and forth

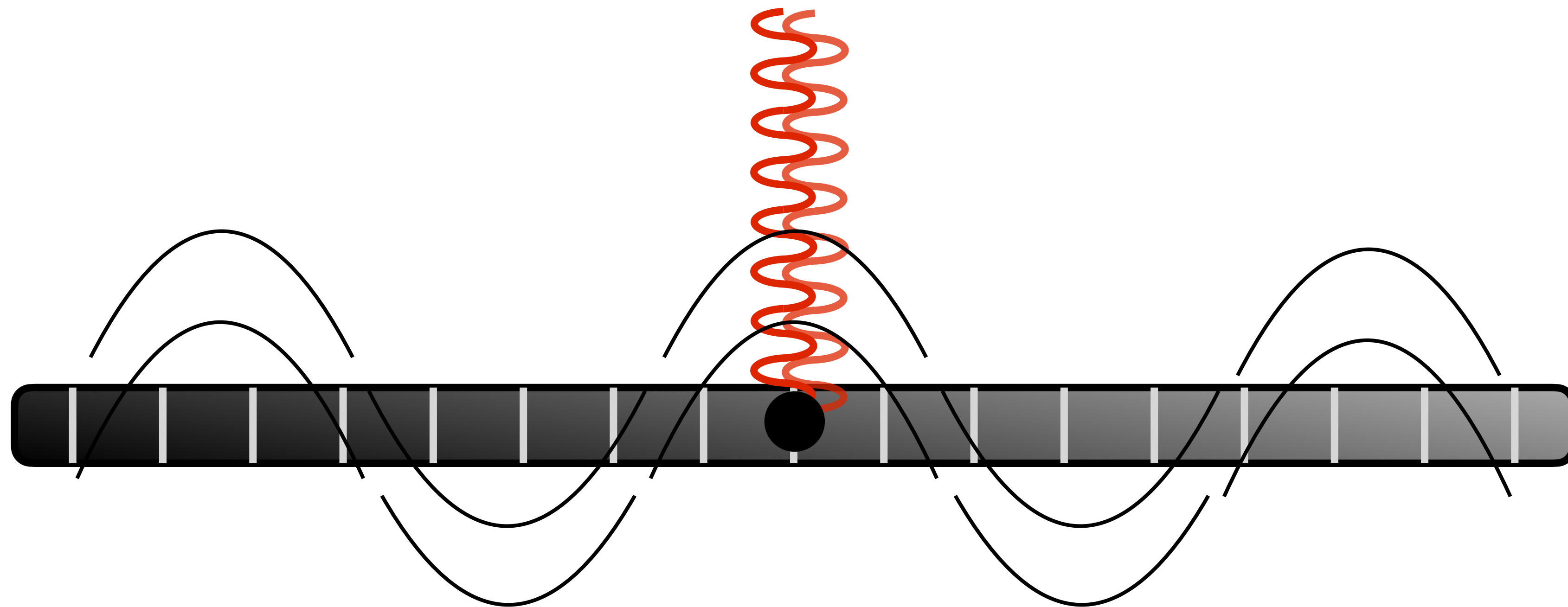
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Incoming GW: matched wavelength



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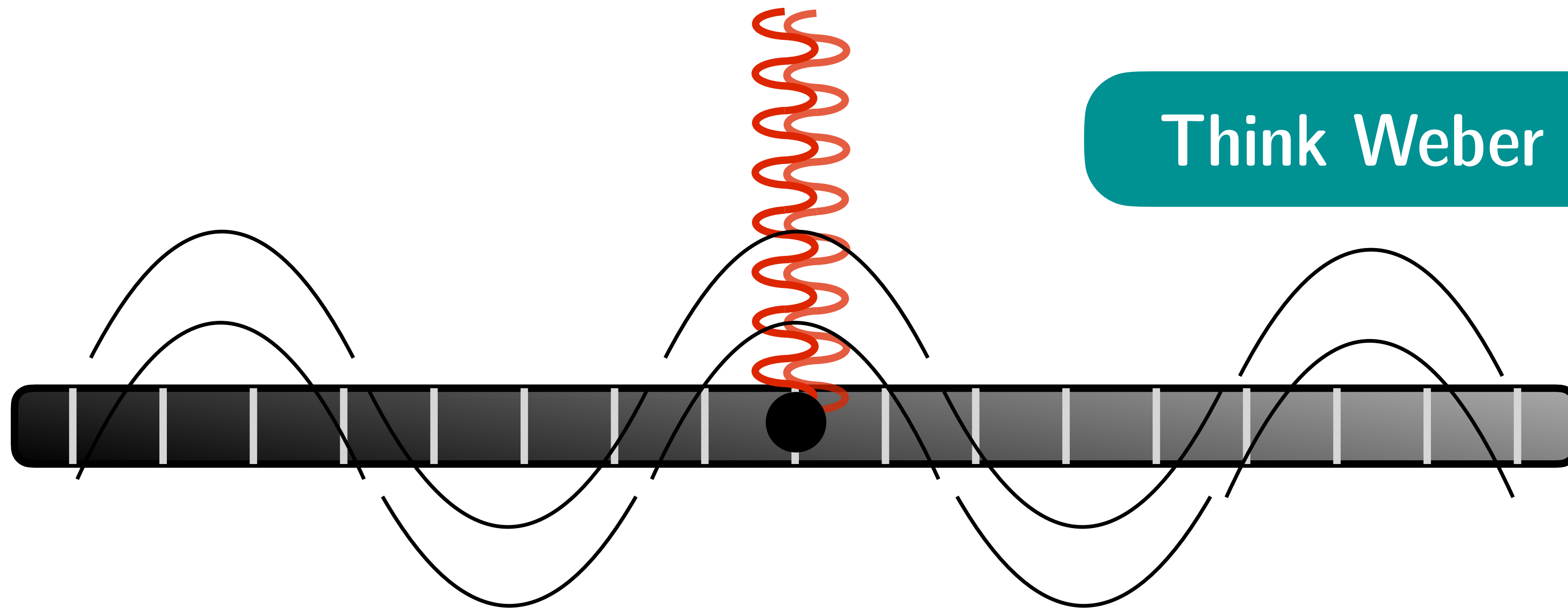
Incoming GW: matched wavelength



Resolve structure of the ruler, e.g. resonances

When is the ruler rigid?

Incoming GW: matched wavelength

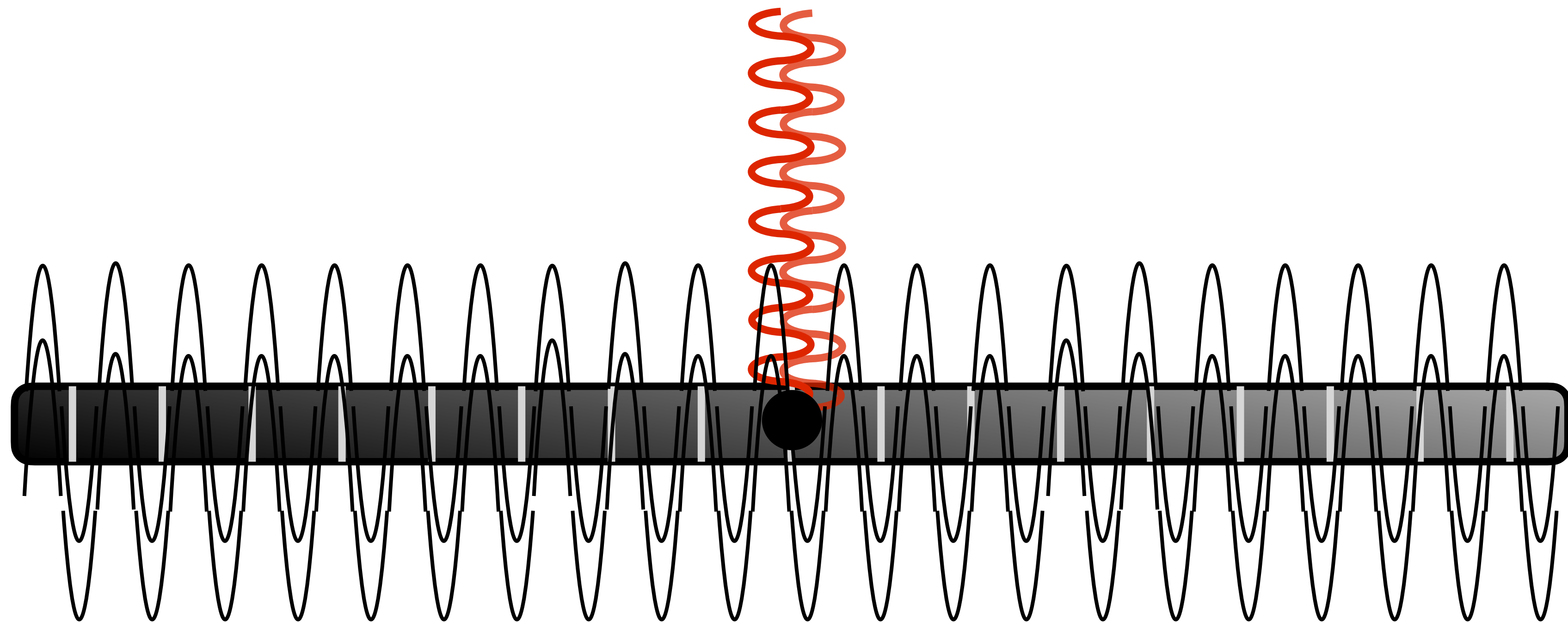


Think Weber Bar

Resolve structure of the ruler, e.g. resonances

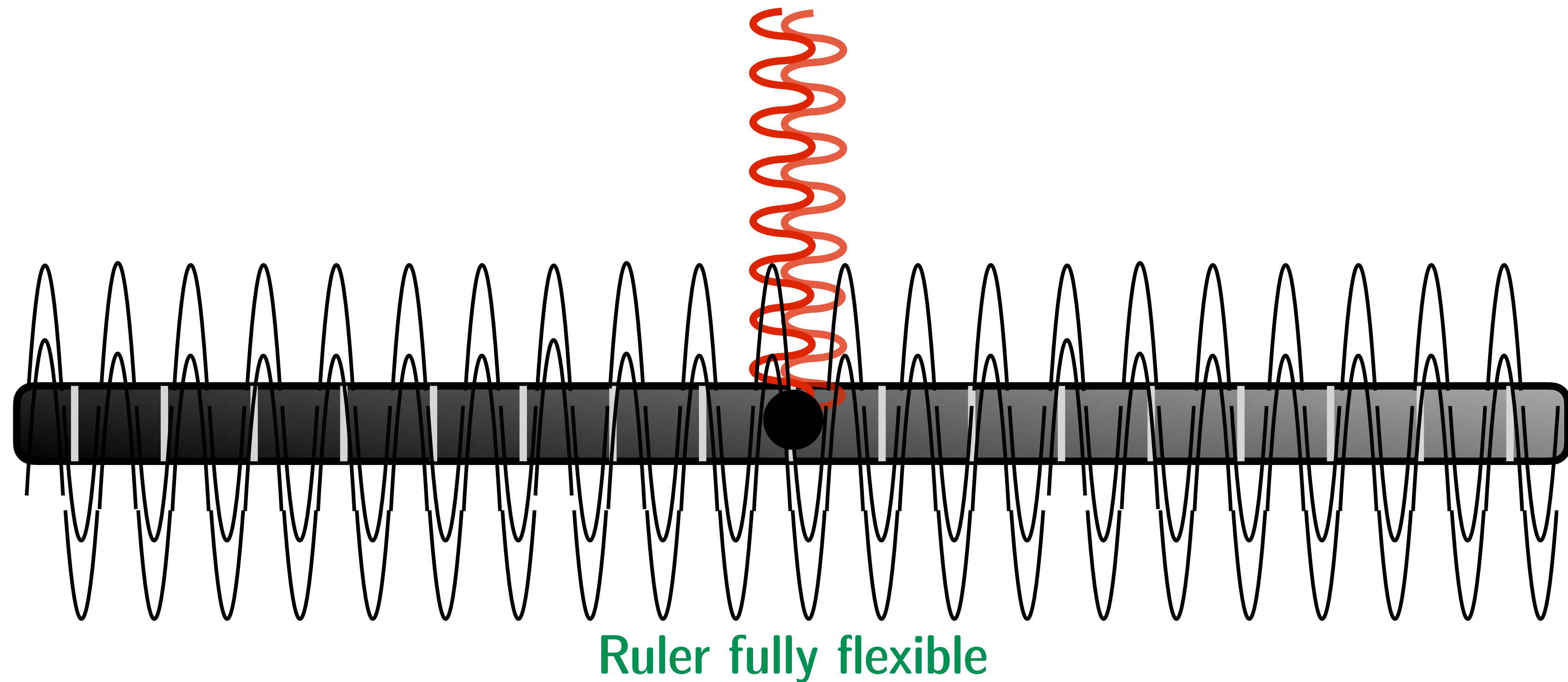
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Incoming GW: short wavelength



When is the ruler rigid?

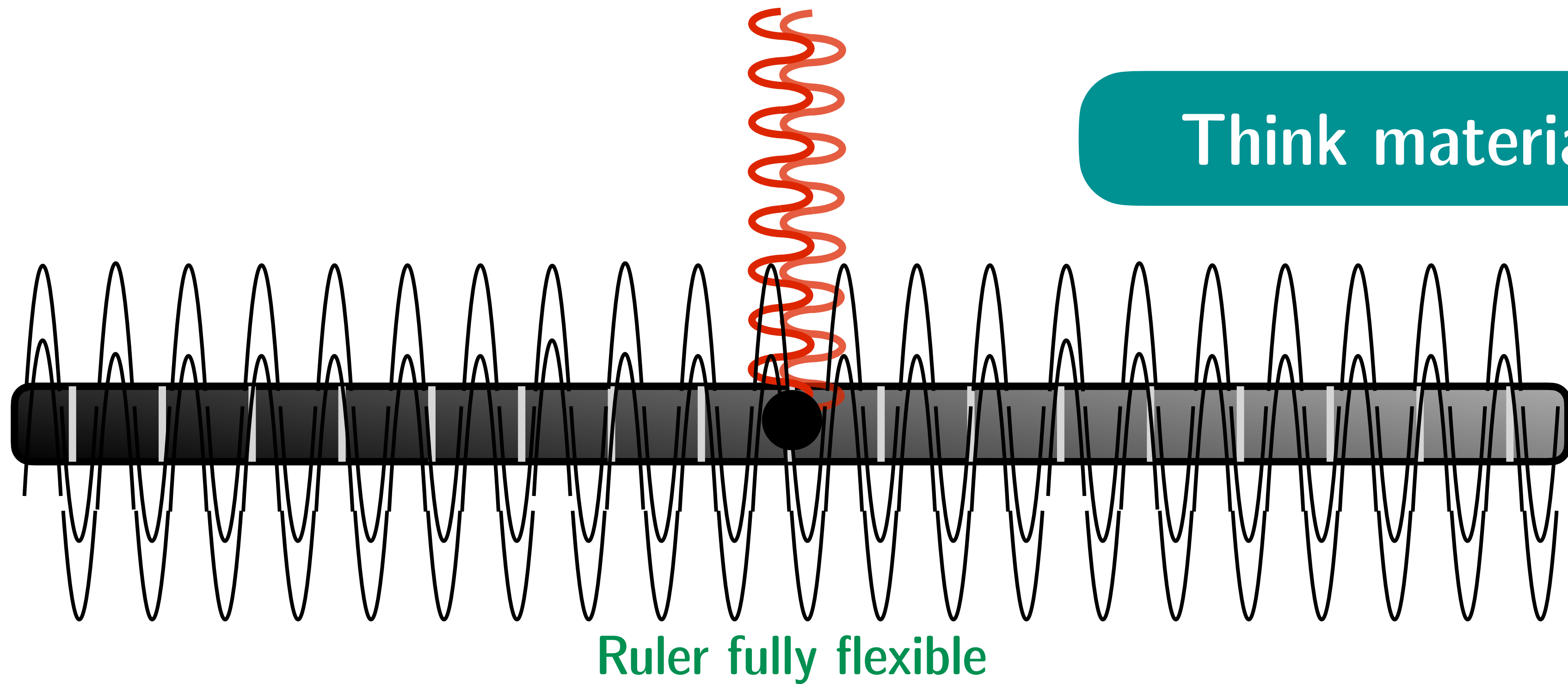
Incoming GW: short wavelength



When is the ruler rigid?

Incoming GW: short wavelength

Think materials



When is the ruler rigid?

At the level of equations:

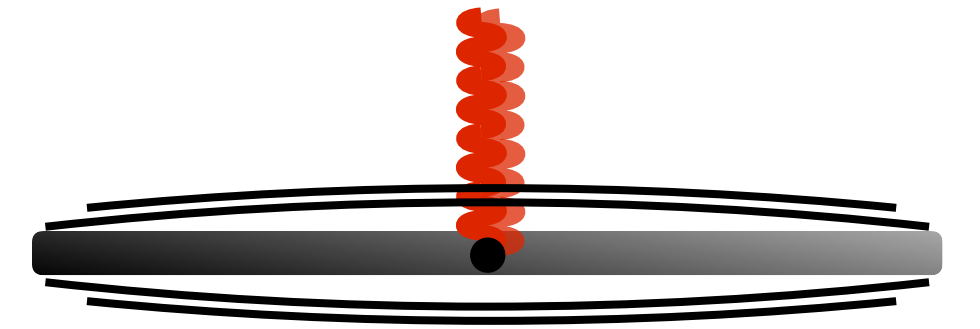
$$\partial_t^2 s_i + k^2 s_i \sim \partial_t^2 h_{ij} s_j$$

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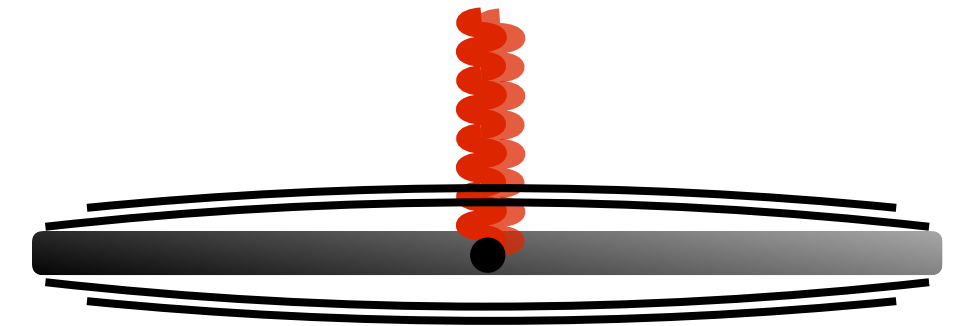
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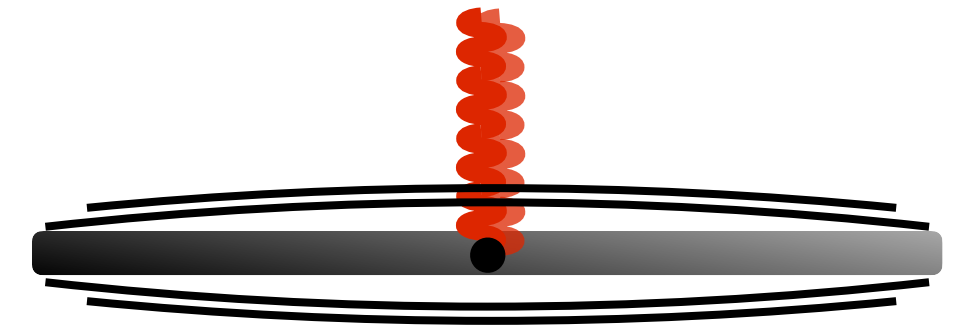
Looks like proper detector frame $k \sim 1/L$

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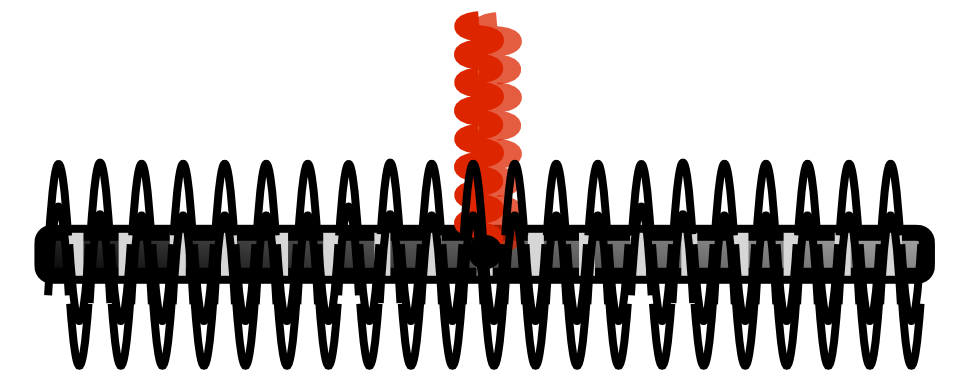
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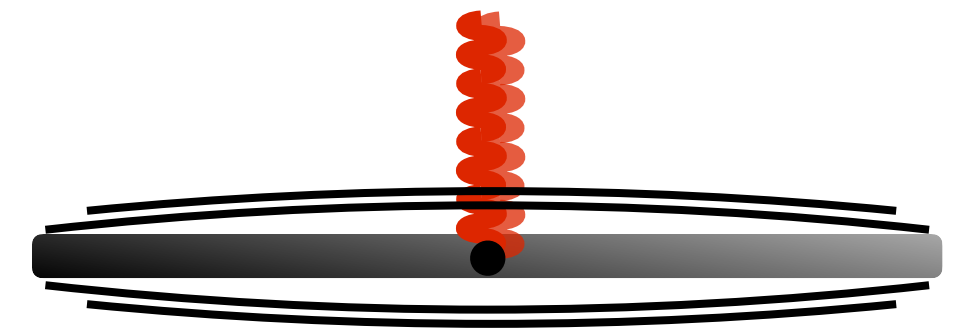
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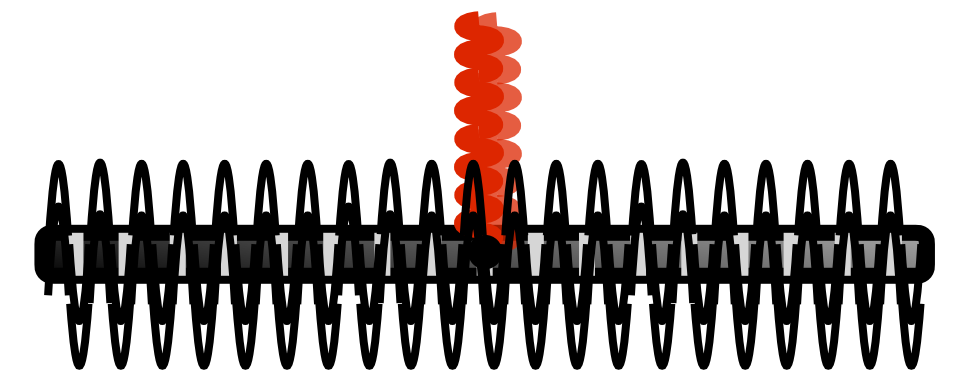
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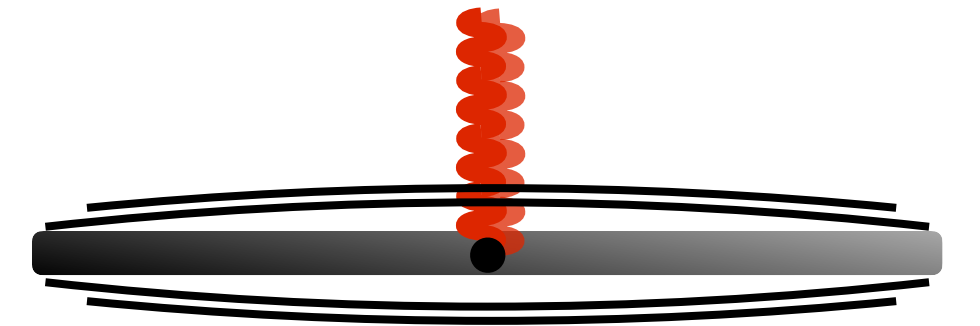
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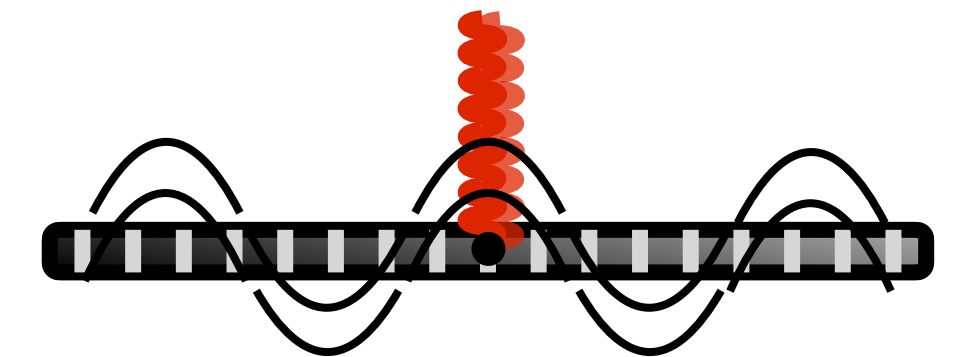
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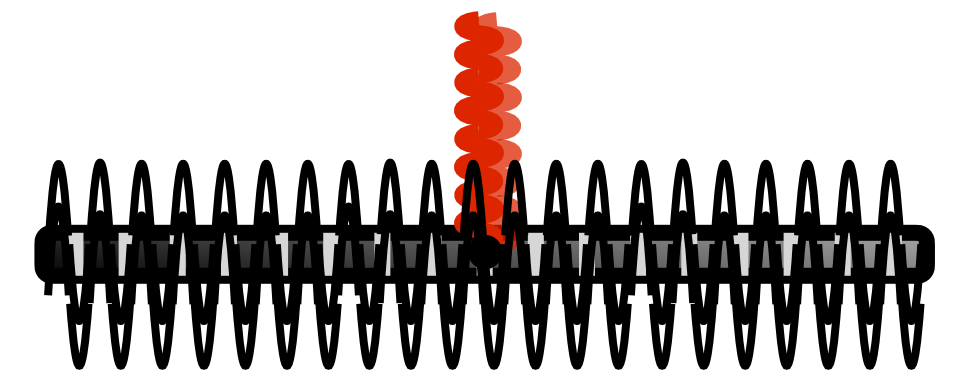


Case requires knowledge of detector



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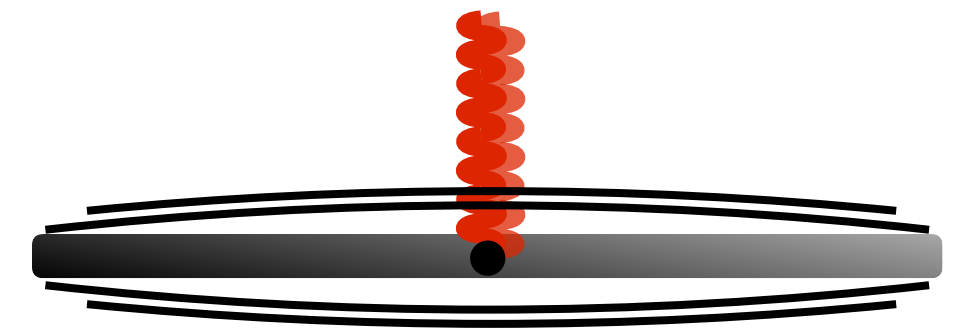
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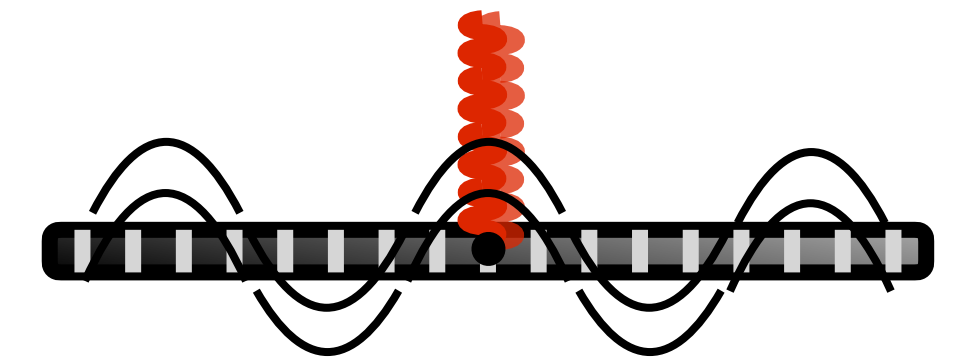
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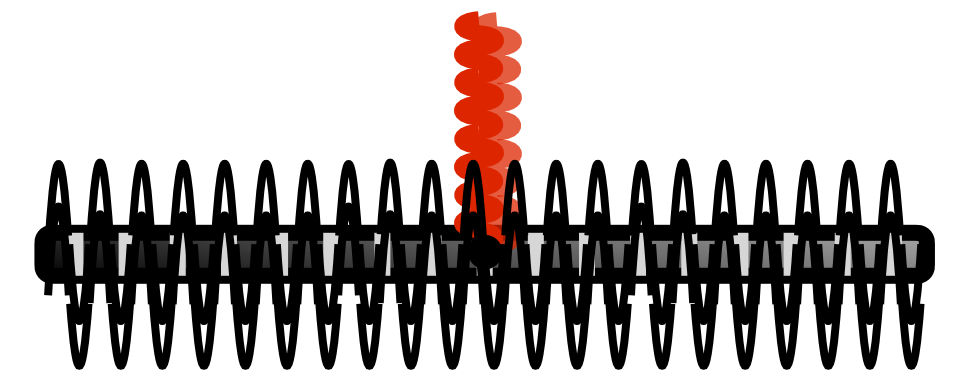
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Frame subtleties crucial!

Looks like TT frame



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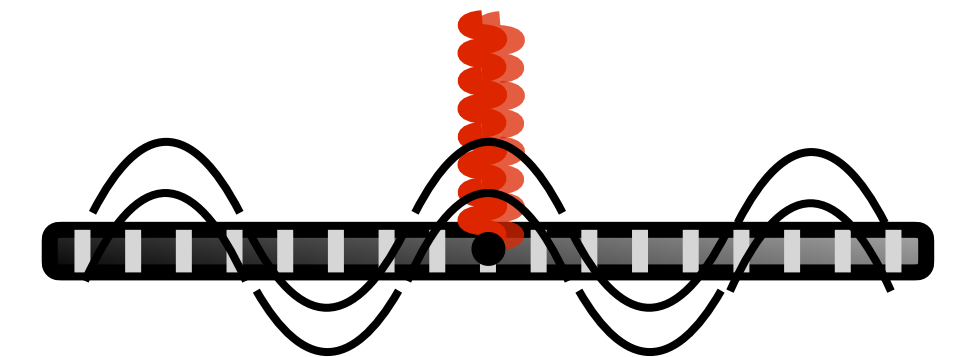
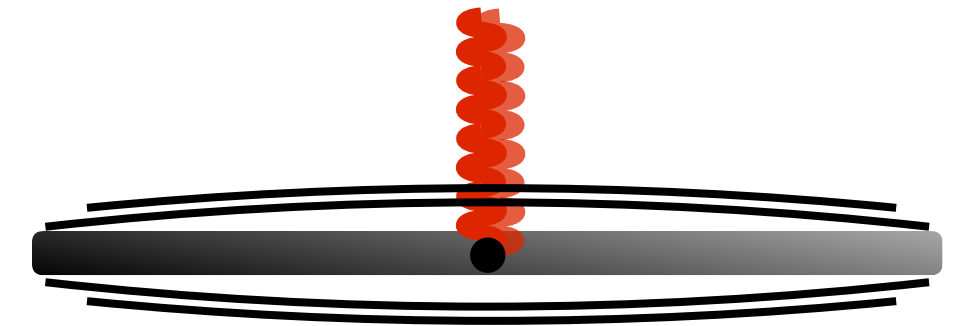
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NB: don't include for LIGO mirror

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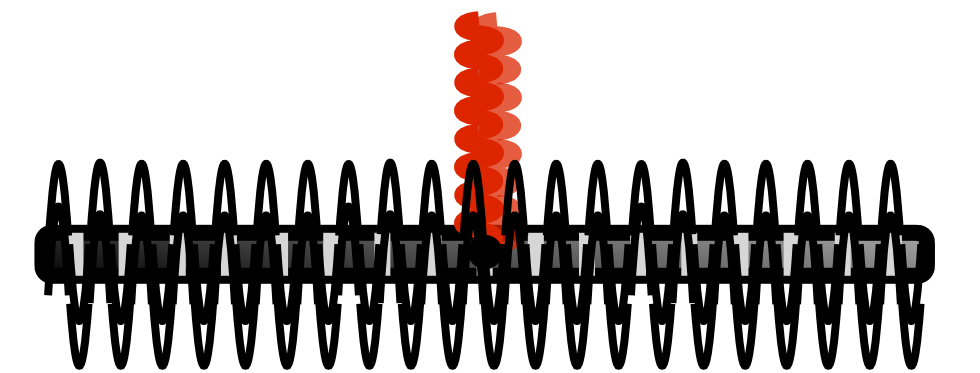
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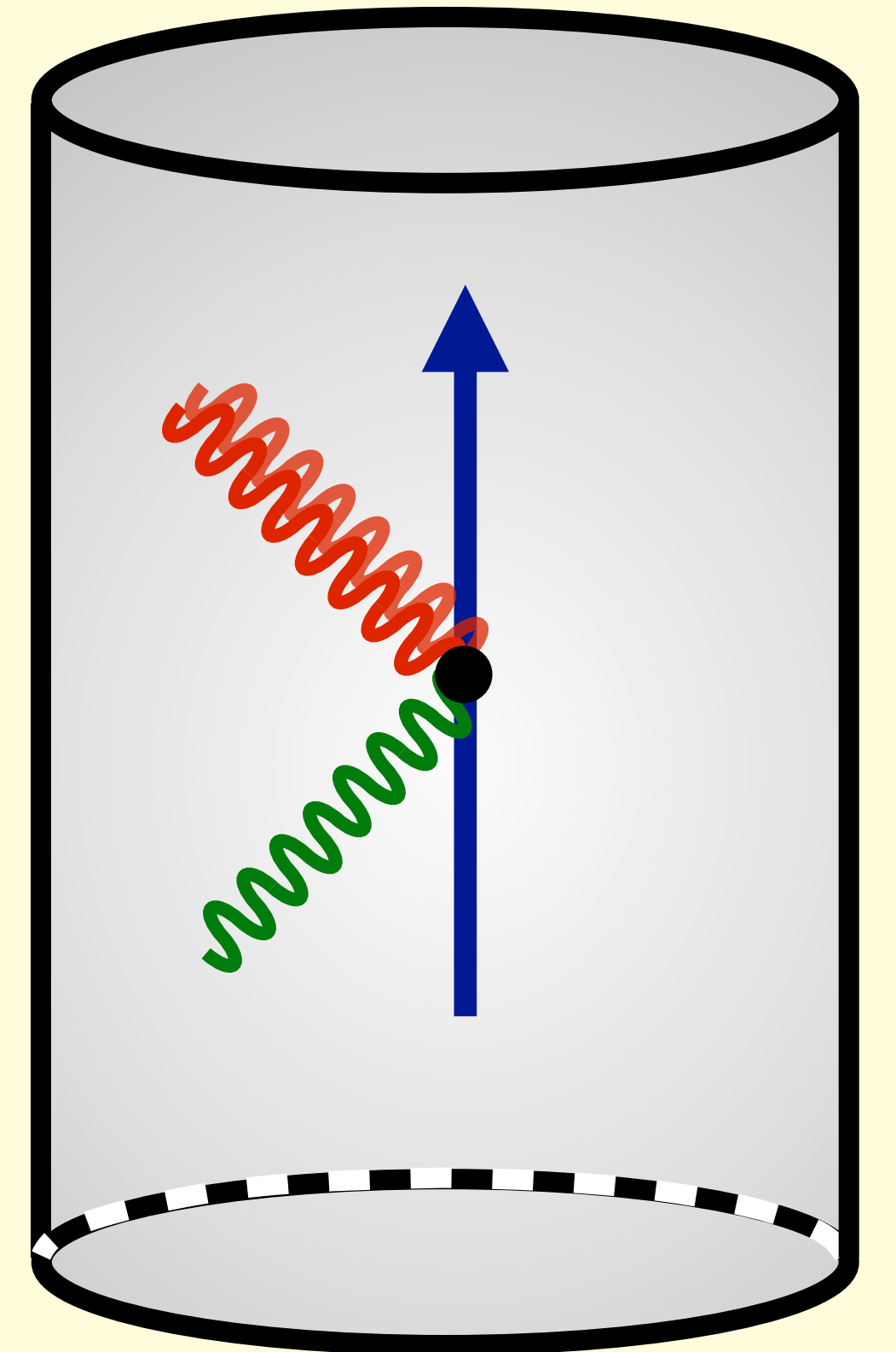
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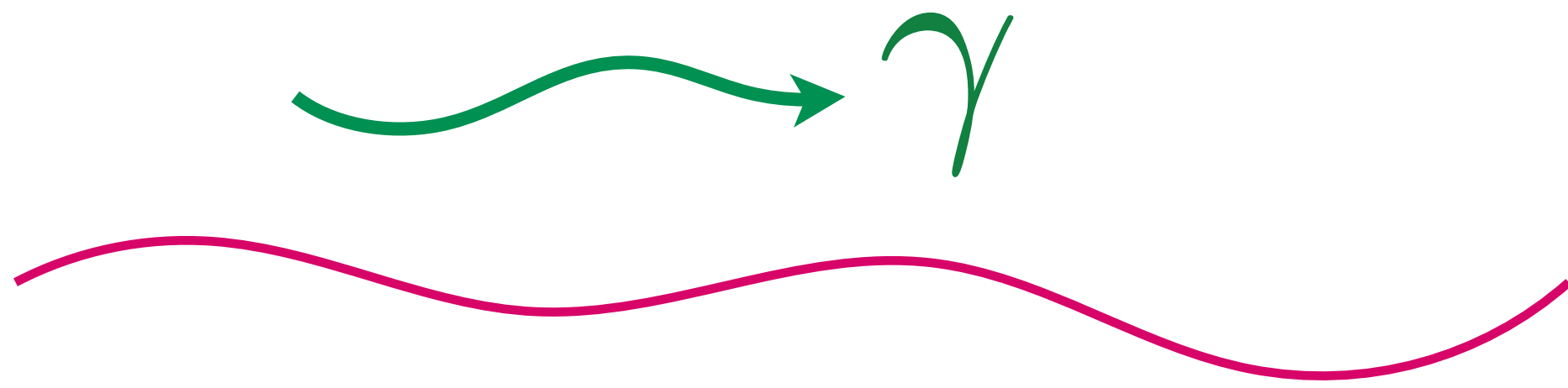
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ACTE I

Light is the Ruler:
Static B-field Cavities

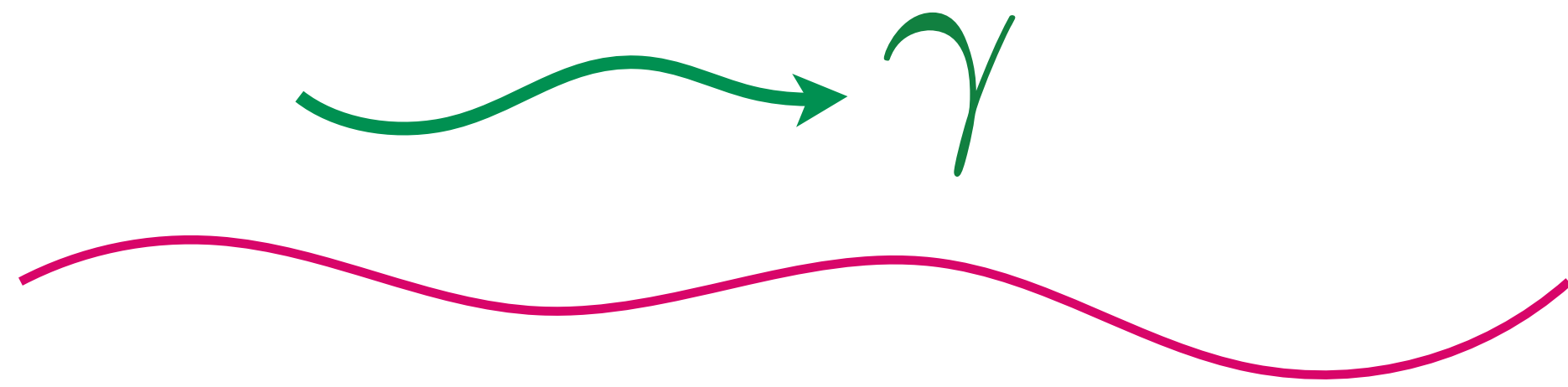


Interactions of Gravitational Waves *with light*



$$S_{\text{EM}} = \int d^4x \sqrt{-g} \left(-\frac{1}{4} g^{\mu\alpha} g^{\nu\beta} F_{\mu\nu} F_{\alpha\beta} + g^{\mu\nu} J_{\mu} A_{\nu} \right)$$

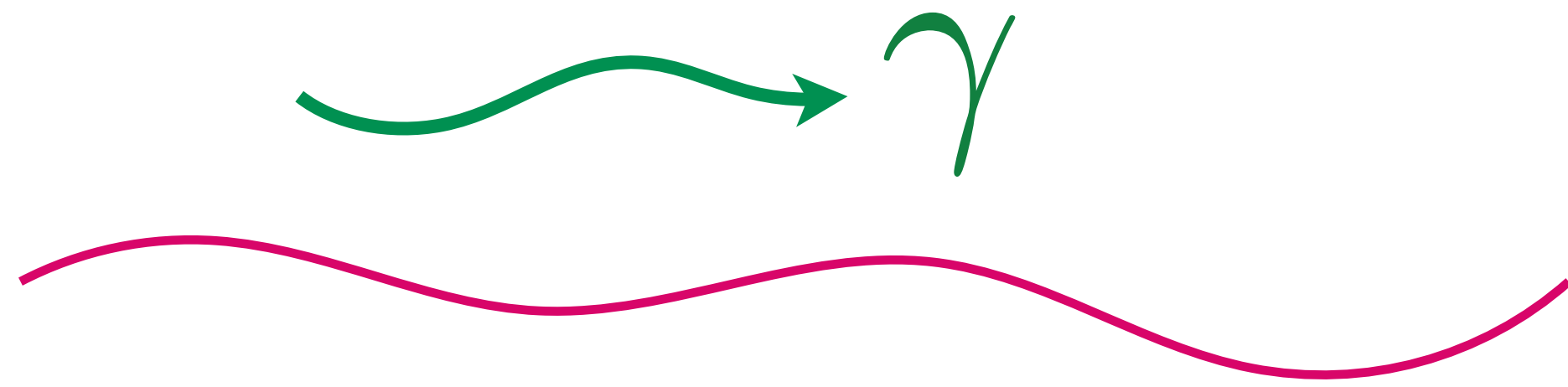
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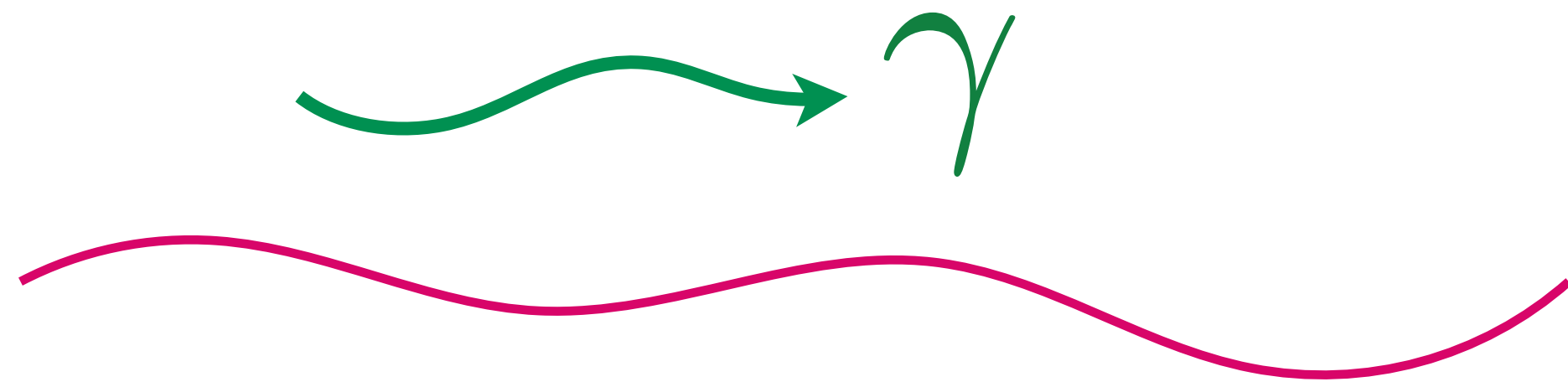


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Effective current from spatial or temporal variations of h or F

$$j_{\text{eff}}^{\mu} \equiv \partial_{\nu} \left(\frac{1}{2} h F^{\mu\nu} + h^{\nu}_{\alpha} F^{\alpha\mu} - h^{\mu}_{\alpha} F^{\alpha\nu} \right)$$

Cur Cavis?* Part I: Electromagnetic Ruler

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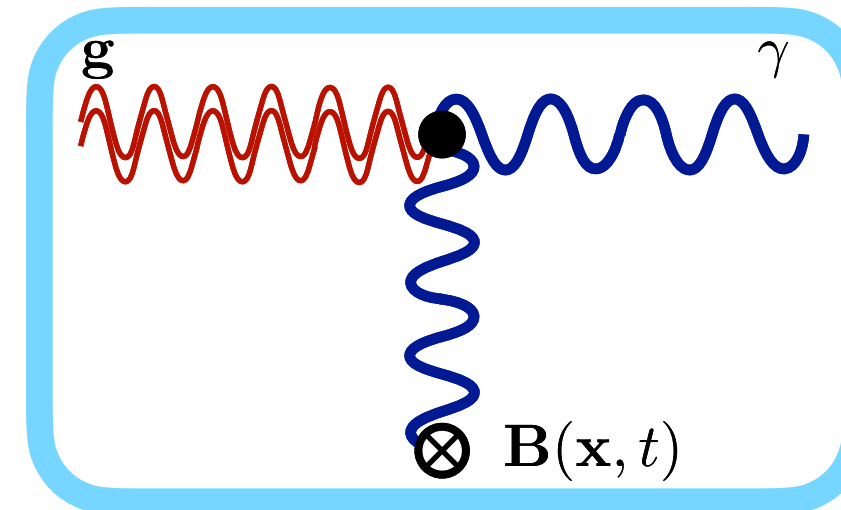
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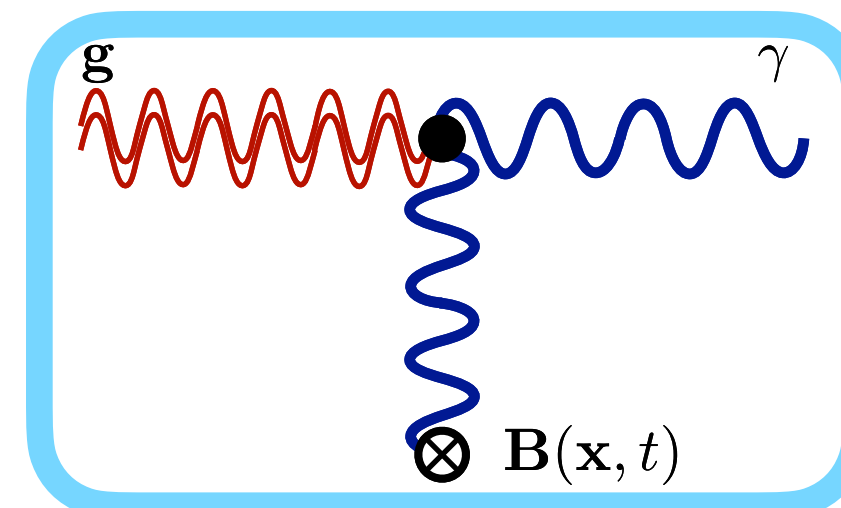
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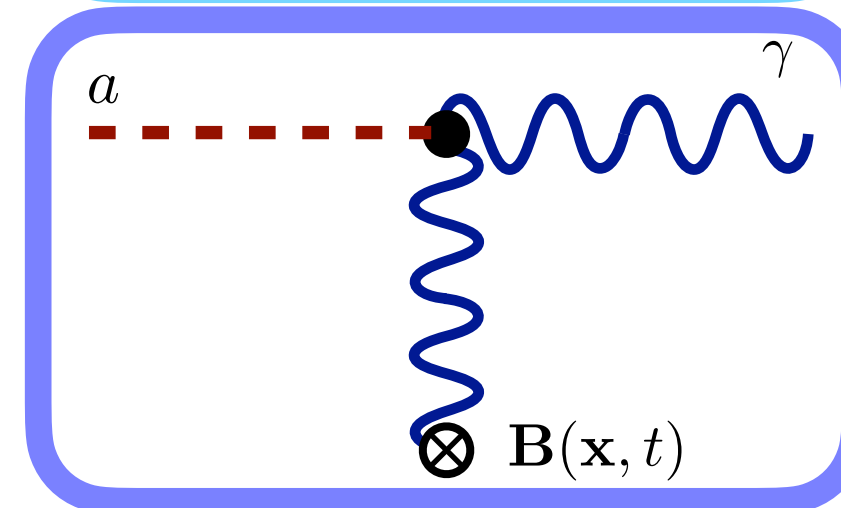
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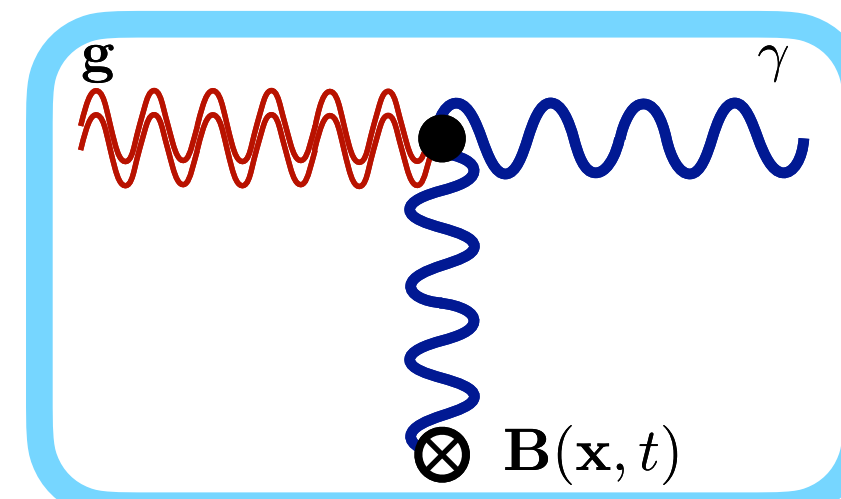
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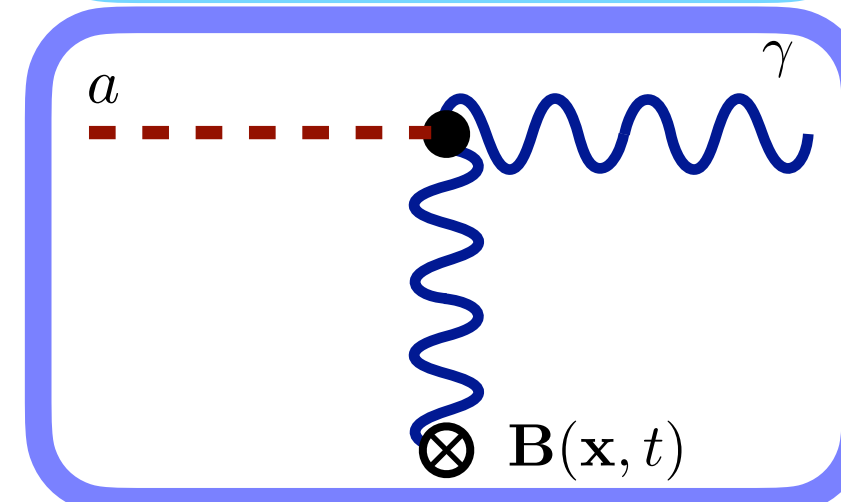
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Raffelt & Stodolsky (1988)

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Framing the Question

Crucial to work in appropriate reference frame!

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Detector in Local Inertial Frame (LIF)

$$\hat{n} \times \mathbf{E} = 0$$

$$\hat{n} \cdot \mathbf{B} = 0$$

Maxwell (19th century)

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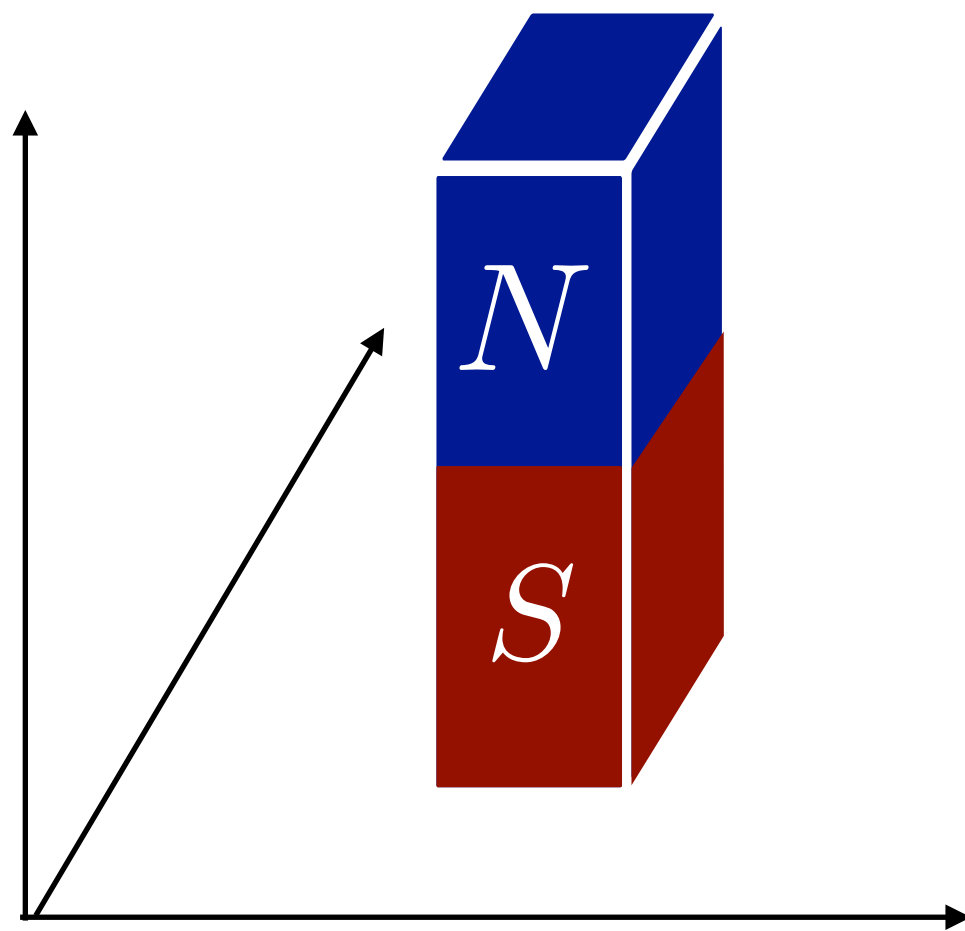
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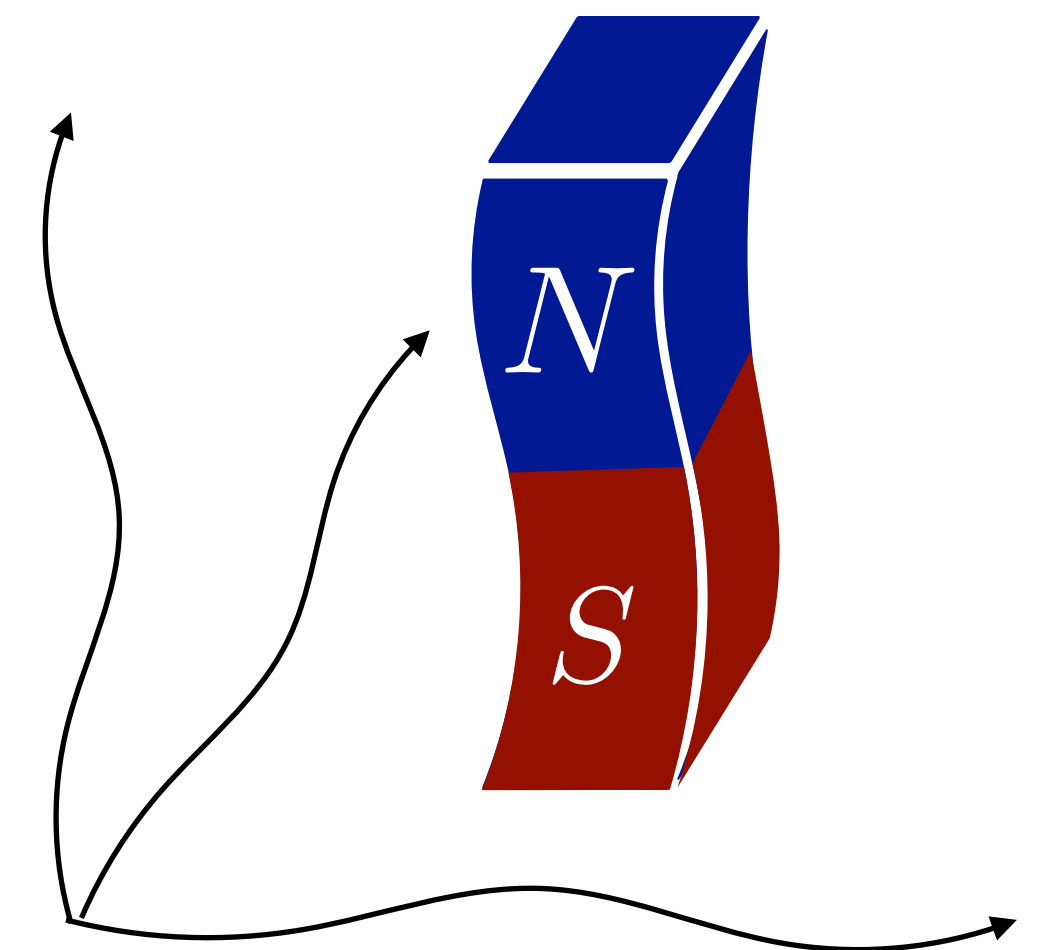
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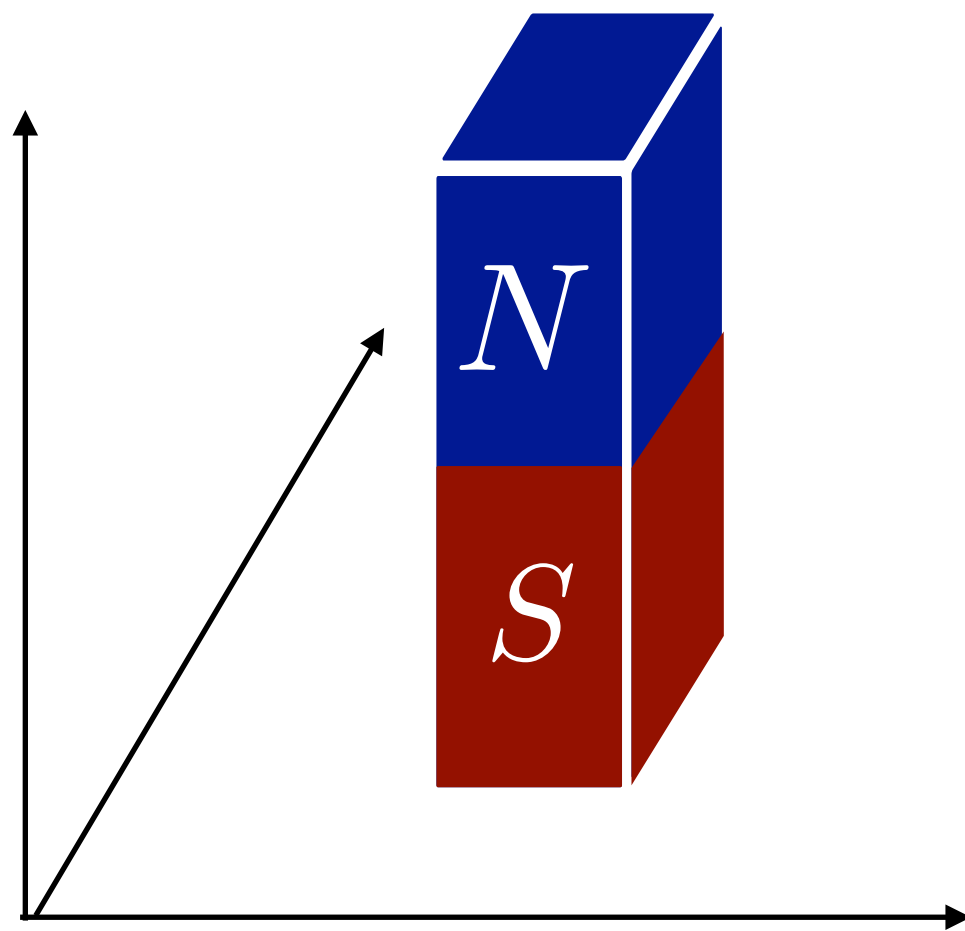
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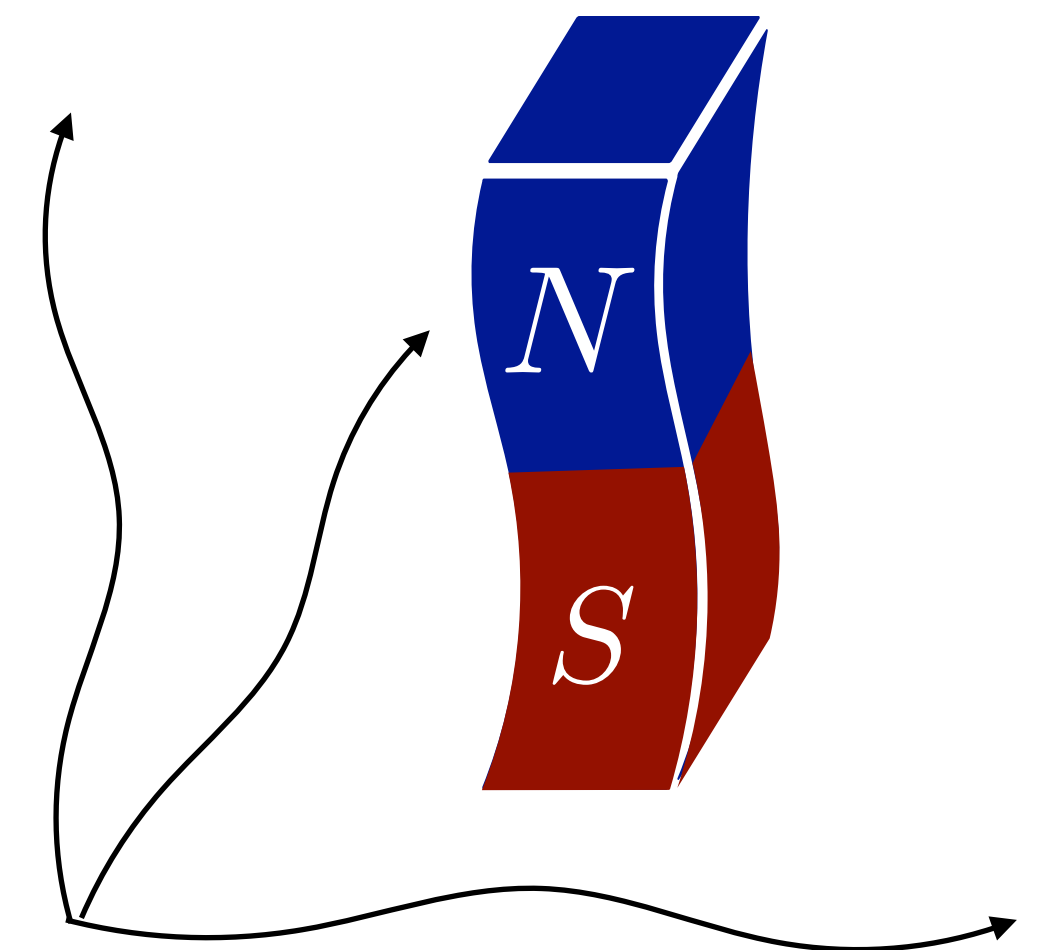
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Which frame is the right one to use?



Framing the Question

Proper Detector Frame — complication

Framing the Question

Proper Detector Frame — complication

Textbooks give long-wavelength approximation $\omega_g R_{\text{cav}} \ll 1$

$$ds^2 \simeq -dt^2(1 + R_{0i0j}x^i x^j) - \frac{4}{3} dt dx^i (R_{0ijk}x^j x^k) + dx^i dx^j \left(\delta_{ij} - \frac{1}{3} R_{ikjl}x^k x^l \right) \text{ e.g. Maggiore (2007)}$$

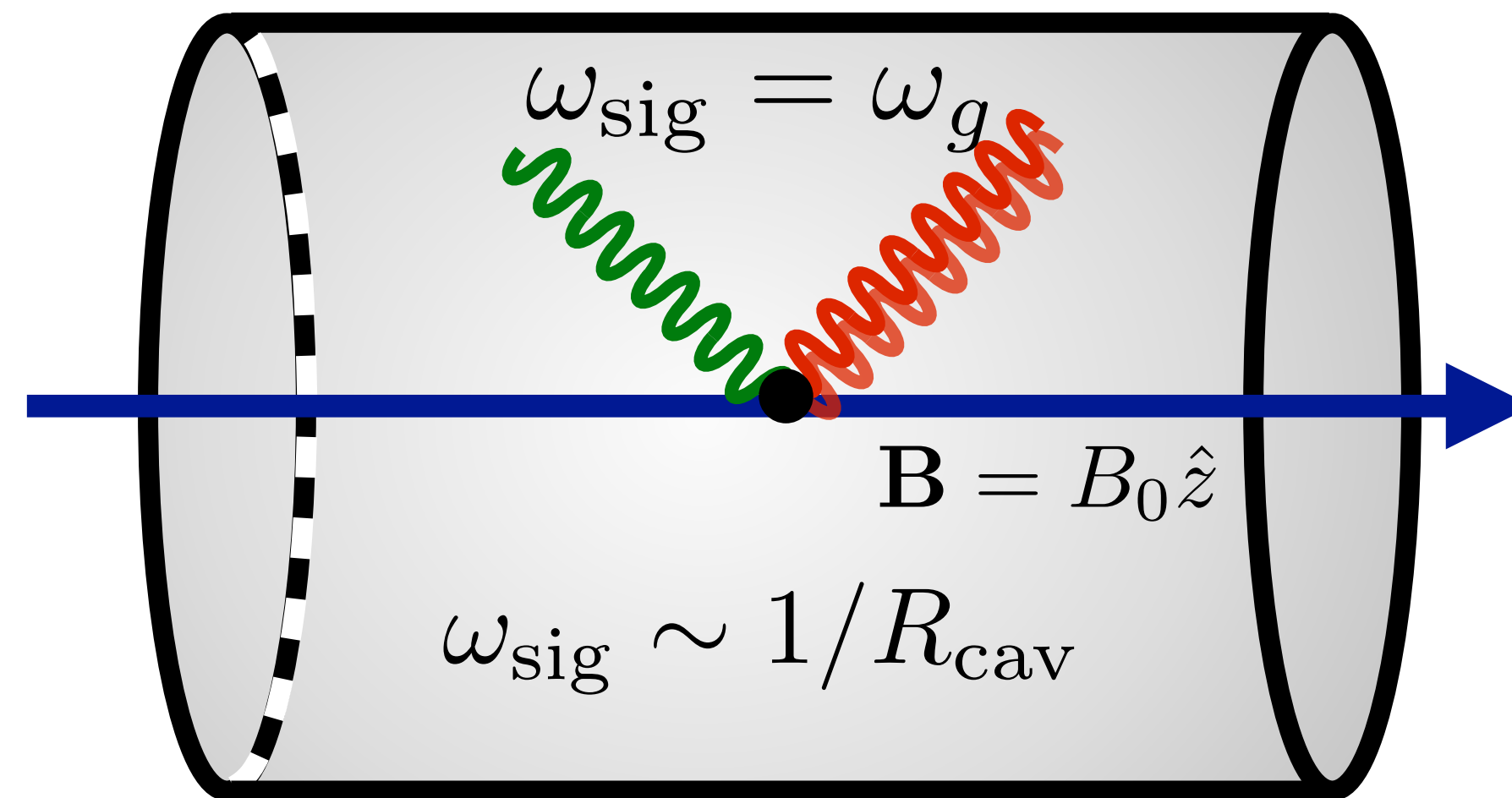
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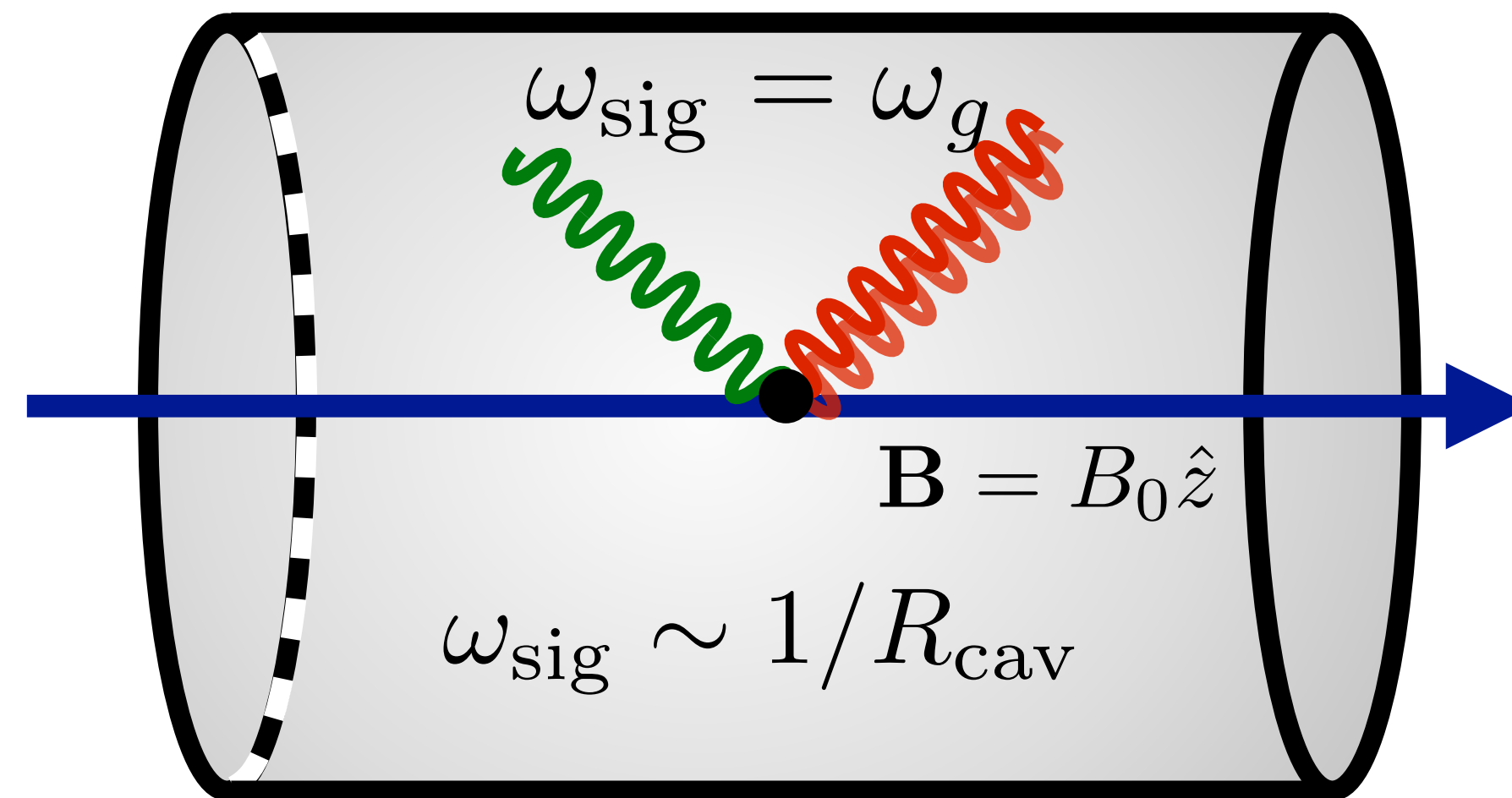
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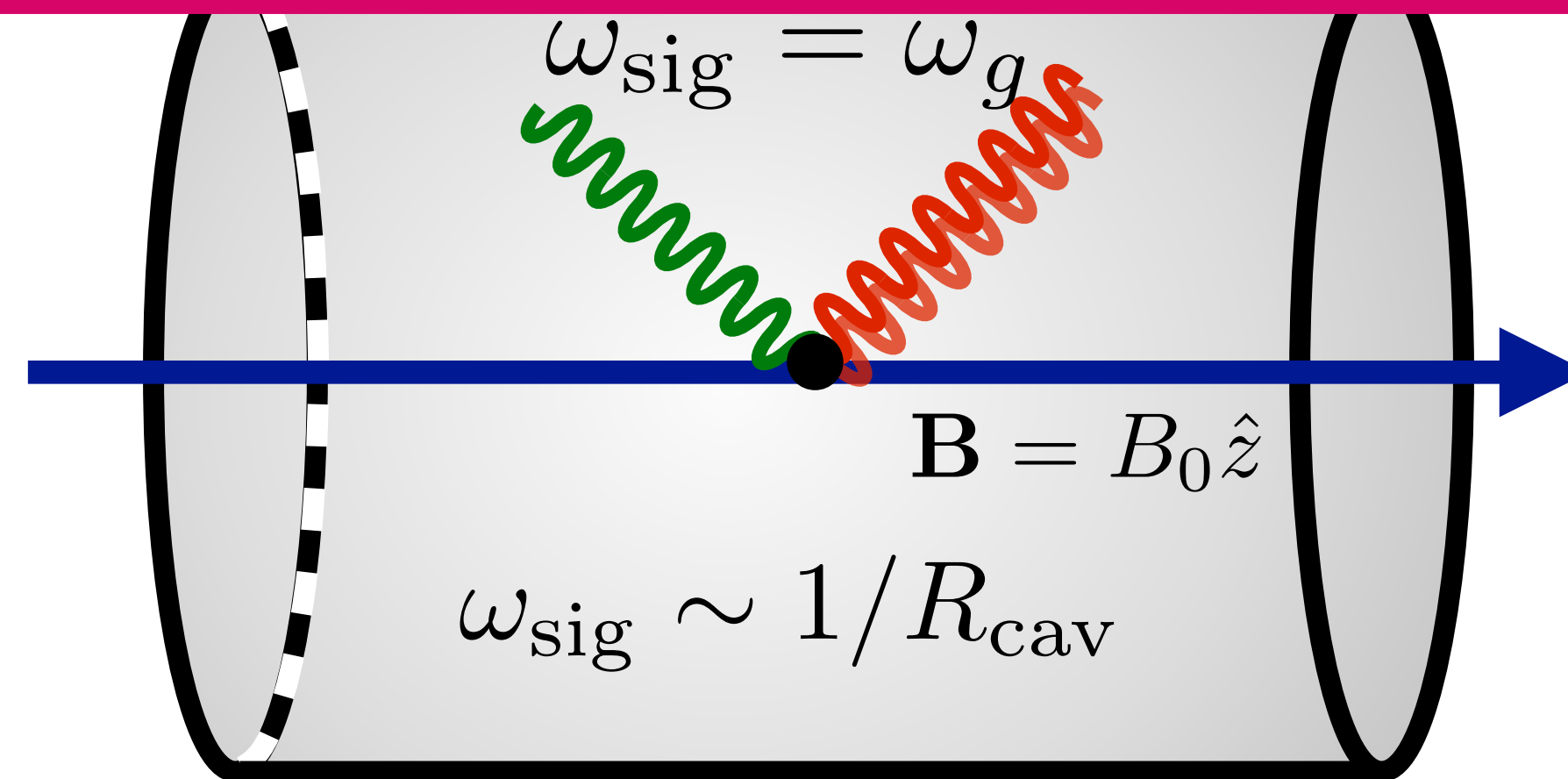
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Long-wavelength limit invalid!

Resonant Cavity:



Framing the Question

Solution — GW as sum of plane waves

$$h \propto e^{i\omega_g(t-z)} \longrightarrow \partial_i h_{jk}^{\text{TT}} \sim -\delta_{iz} \partial_t h_{jk}^{\text{TT}}$$

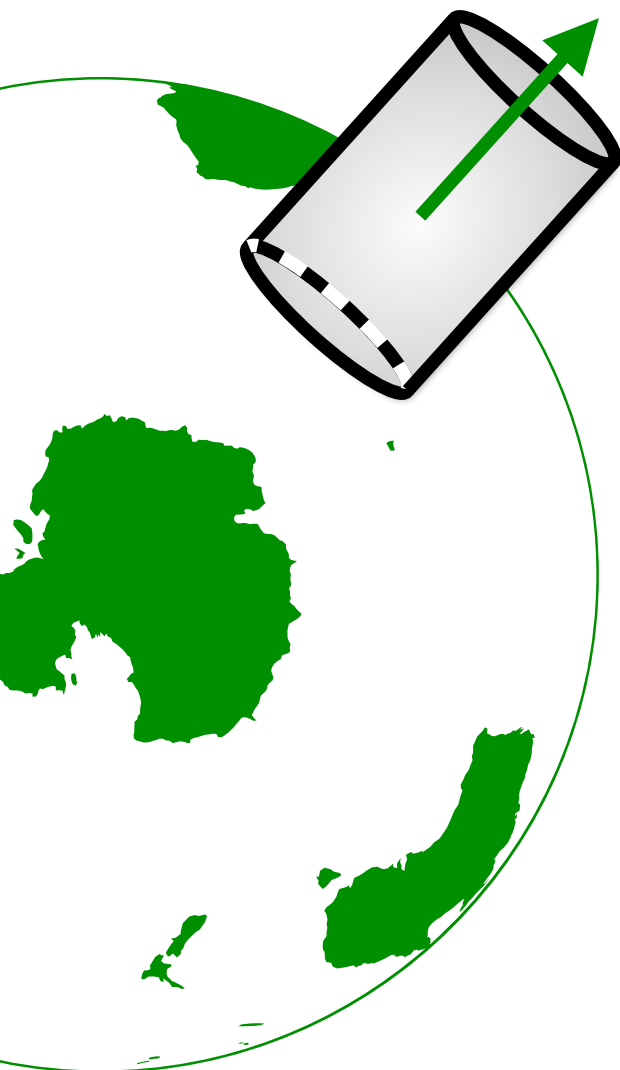
$$x^{k_1} \dots x^{k_r} R_{\mu\nu\rho\sigma, k_1 \dots k_r} = (-i\omega_g z)^r R_{\mu\nu\rho\sigma}$$

$$h_{00} = -2 \sum_{r=0}^{\infty} \frac{r+3}{(r+3)!} R_{0n0n, k_1, \dots, k_r} x^m x^n x^{k_1} \dots x^{k_r}$$

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Märzlin (1994)
Rakhmanov (2014)



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Solution — GW as sum of plane waves

$$h \propto e^{i\omega_g(t-z)} \longrightarrow \partial_i h_{jk}^{\text{TT}} \sim -\delta_{iz} \partial_t h_{jk}^{\text{TT}}$$

$$x^{k_1} \dots x^{k_r} R_{\mu\nu\rho\sigma, k_1 \dots k_r} = (-i\omega_g z)^r R_{\mu\nu\rho\sigma}$$

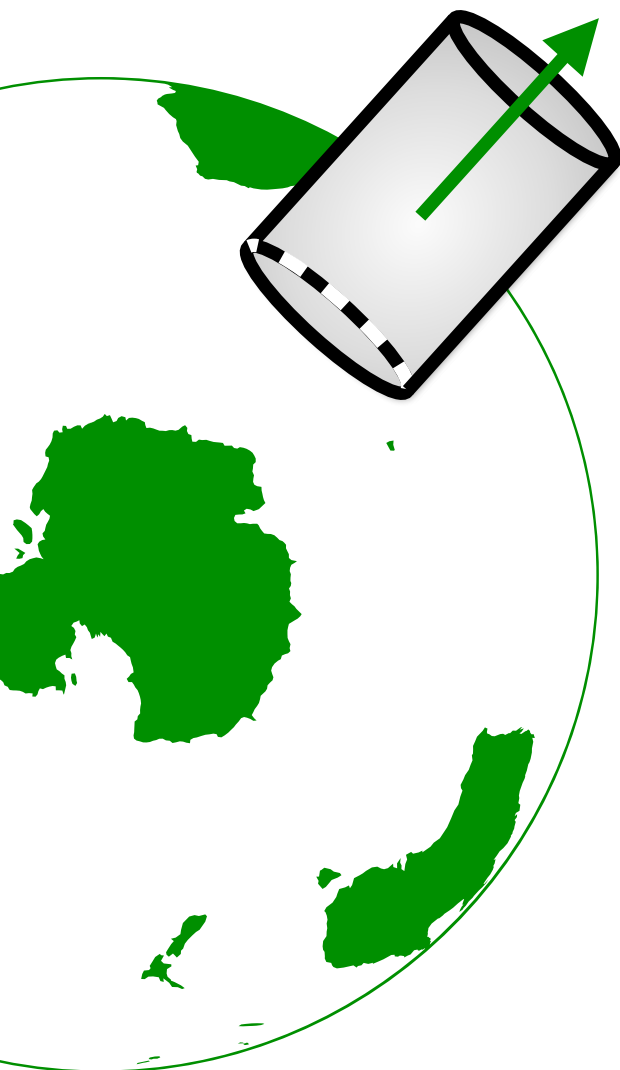
$$h_{00} = -2R_{0m0n} x^m x^n \left(-\frac{i}{\omega_g z} + \frac{1 + e^{-i\omega_g z}}{(\omega_g z)^2} \right)$$

$$h_{0i} = -2R_{0min} x^m x^n \left(-\frac{i}{2\omega_g z} - \frac{e^{-i\omega_g z}}{(\omega_g z)^2} - i \frac{1 - e^{-i\omega_g z}}{(\omega_g z)^3} \right)$$

$$h_{ij} = -2R_{imjn} x^m x^n \left(-\frac{1 + e^{-i\omega_g z}}{(\omega_g z)^2} - 2i \frac{1 - e^{-i\omega_g z}}{(\omega_g z)^3} \right)$$

Berlin, Blas, D'Agnolo, SARE, Harnik, Kahn, Schutte-Engel (PRD 2022)

Märzlin (1994)
Rakhmanov (2014)

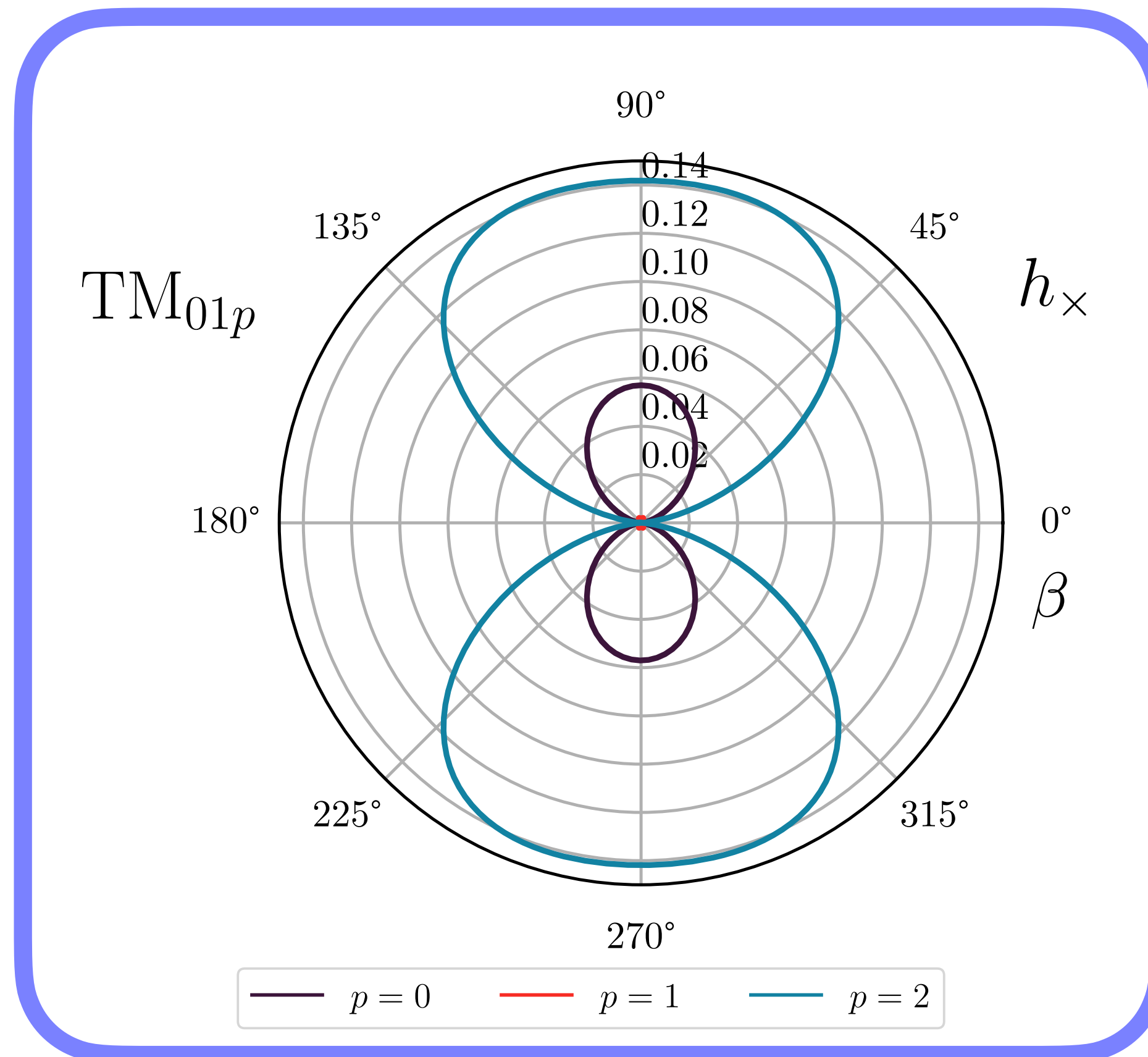


Axion Cavity Modes Couple to GWs

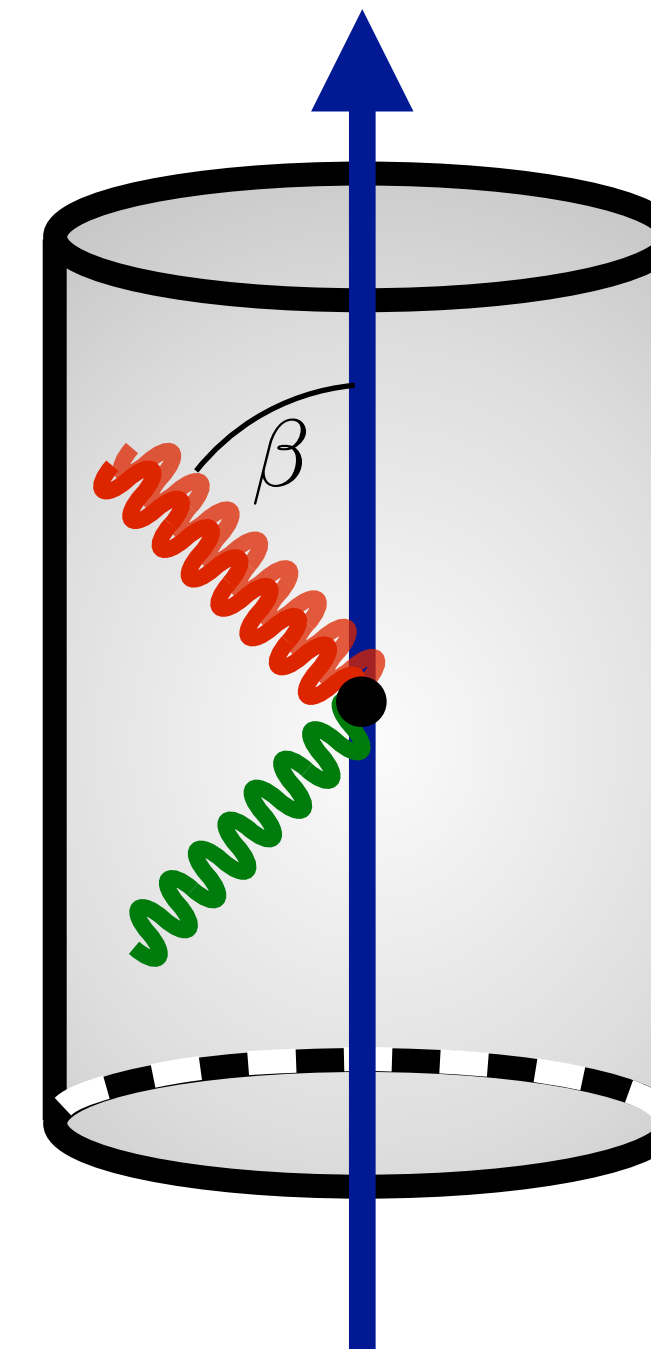
$$\eta \propto \int_V \mathbf{E}_{\text{cav}}^* \cdot \mathbf{J}_{\text{eff}}$$

Axion Cavity Modes Couple to GWs

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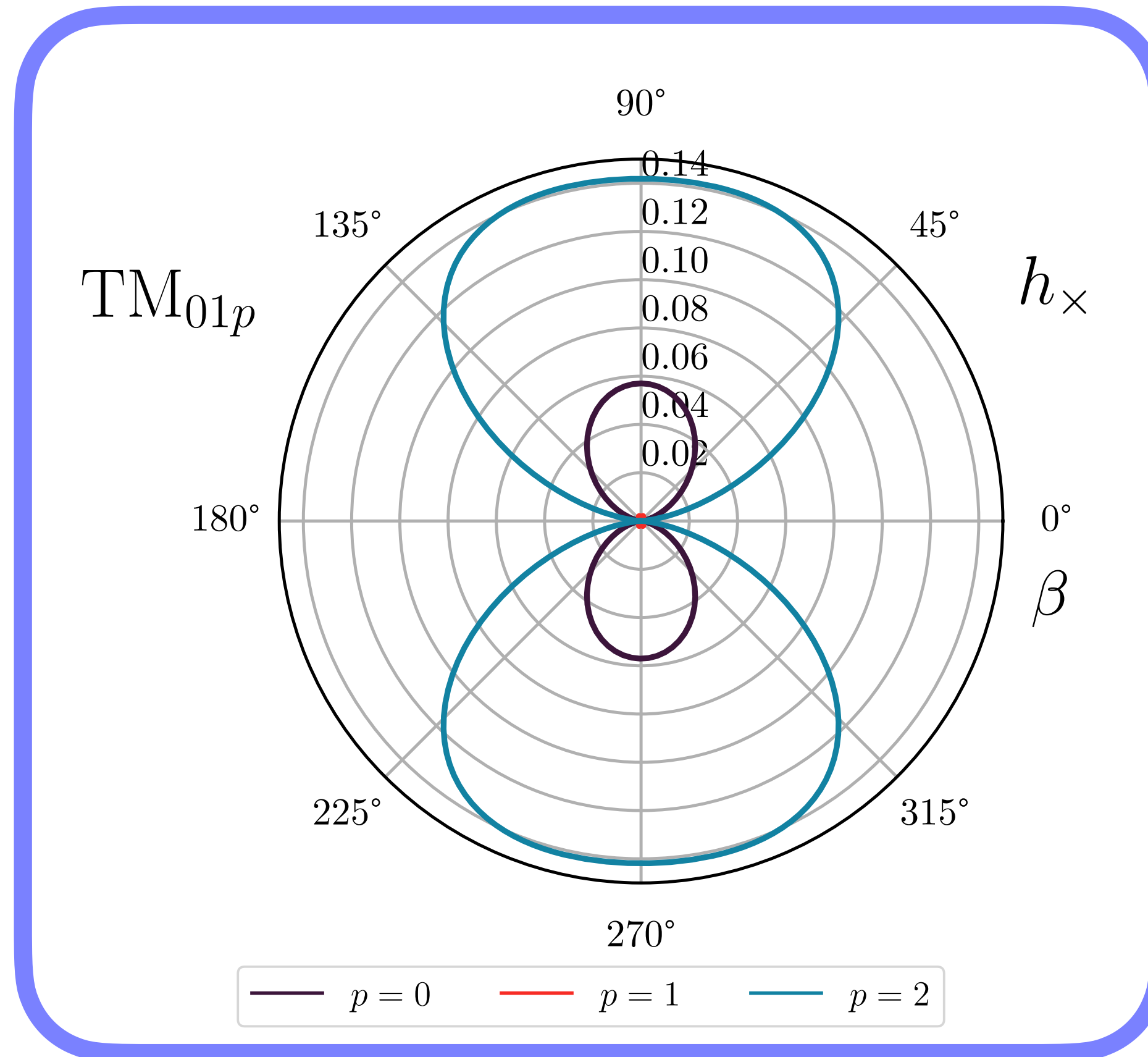


Berlin, Blas, D'Agnolo, SARE, Harnik, Kahn, Schutte-Engel (PRD 2022)



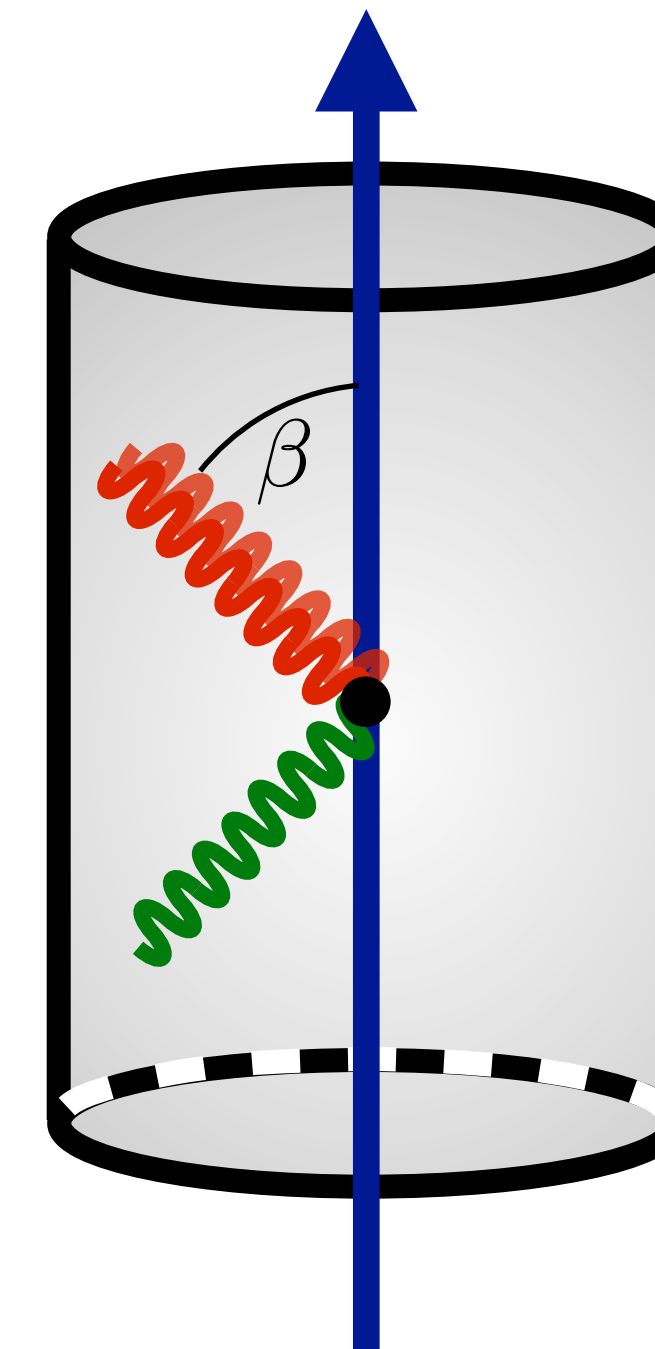
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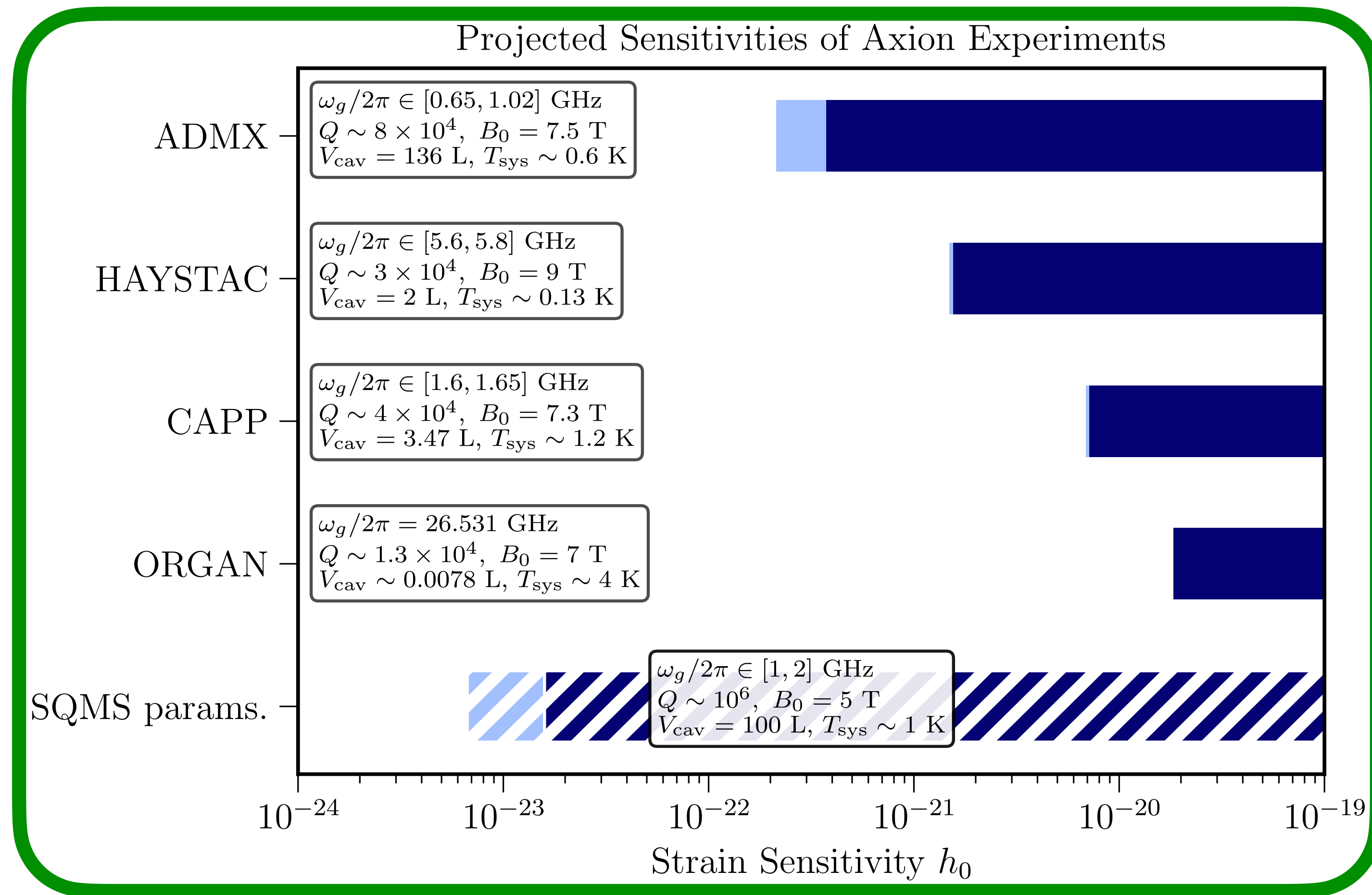


Berlin, Blas, D'Agnolo, SARE, Harnik, Kahn, Schutte-Engel (PRD 2022)

But TM modes not optimal...



Axion Cavity Sensitivity



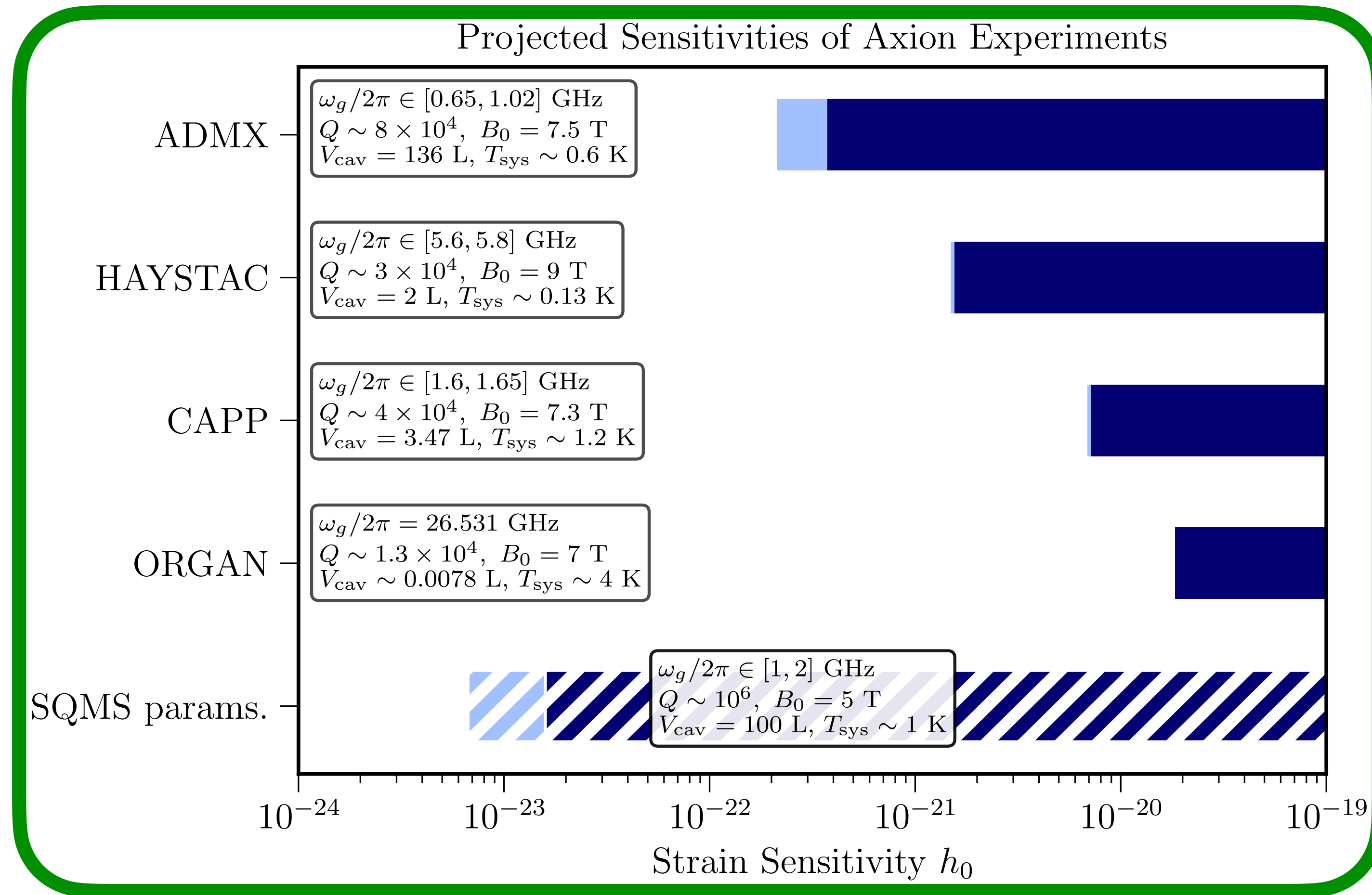
Berlin, Blas, D'Agnolo, SARE, Harnik, Kahn, Schutte-Engel (PRD 2022)

Coherent GW

$$P_{\text{sig}} = \frac{1}{2} Q \omega_g^3 V_{\text{cav}}^{5/3} (\eta_n h_0 B_0)^2$$

see Weds-Fri talks for more on sources

Axion Cavity Sensitivity



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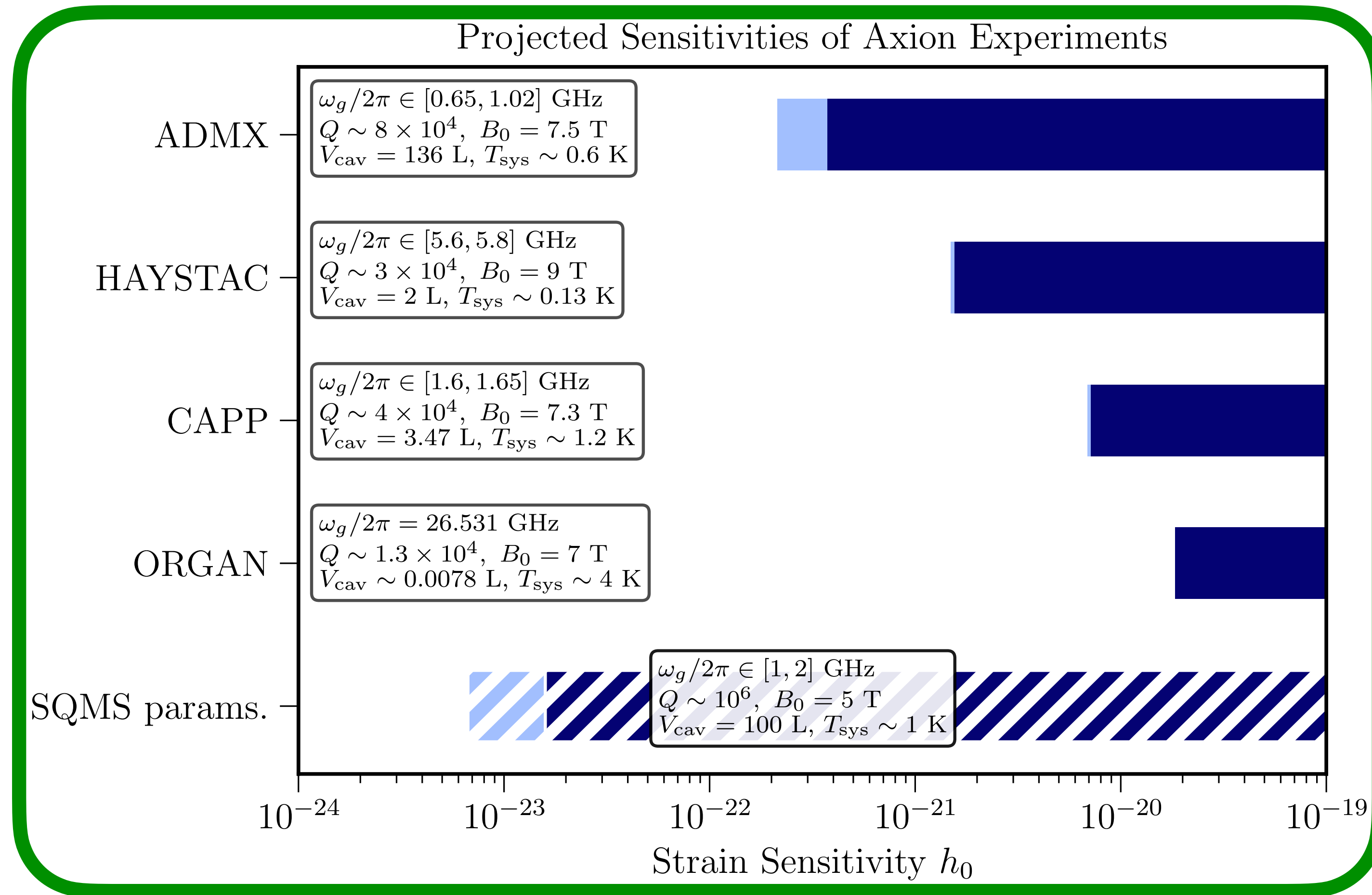
Stochastic GWs

$$\text{SNR} \sim Q \omega_g \eta_{\text{stoch}}^2 B_0^2 V_{\text{cav}} S_h(\omega_g) / T_{\text{sys}}$$

$$\Omega_g(\omega_g) \sim \omega_g^3 S_h(\omega_g) / H_0^2$$

see Weds-Fri talks for more on sources

Axion Cavity Sensitivity



Berlin, Blas, D'Agnolo, SARE, Harnik, Kahn, Schutte-Engel (PRD 2022)

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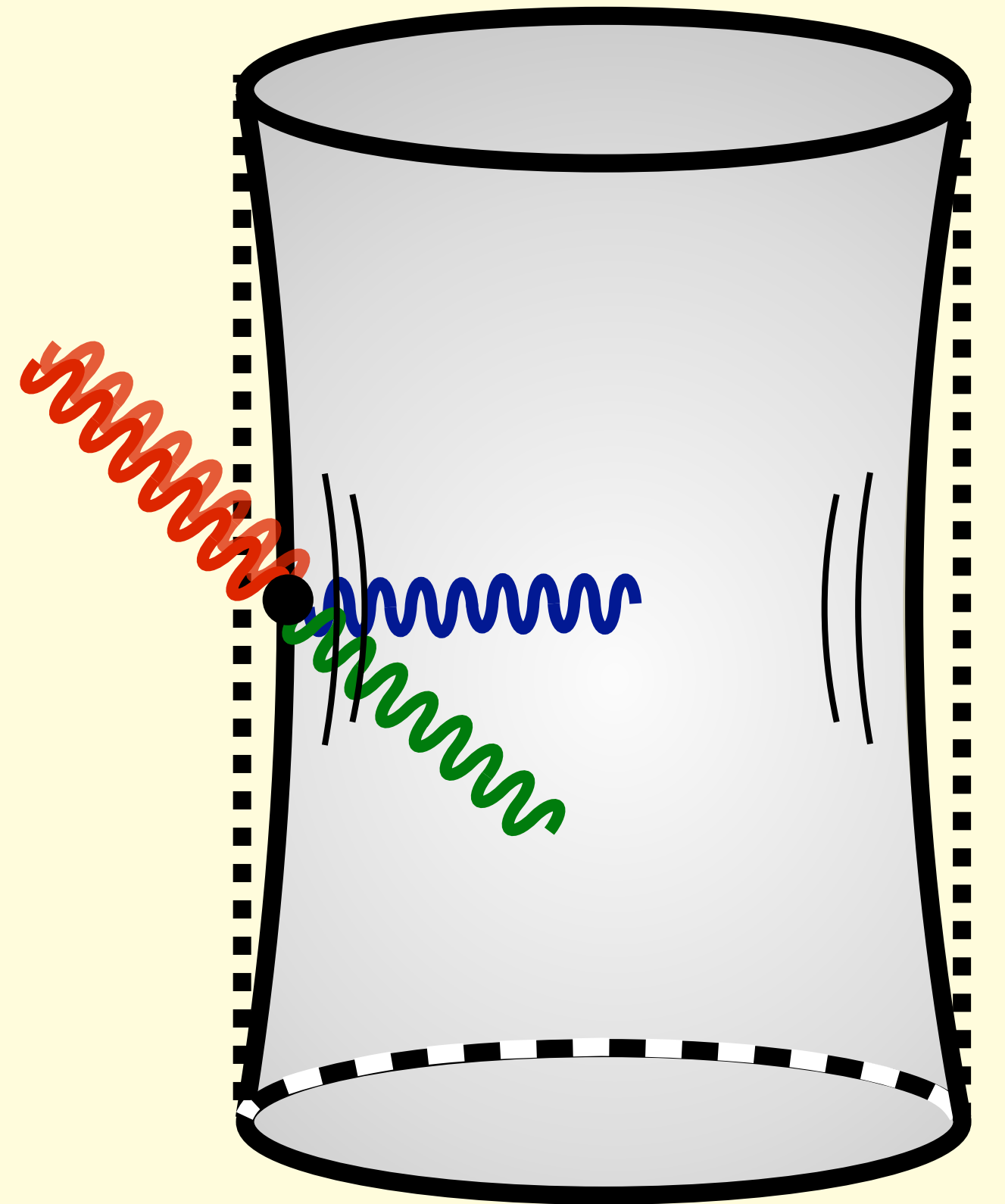
$$\Omega_g(\omega_g) \sim \omega_g^3 S_h(\omega_g) / H_0^2$$

Not beyond BBN bound...

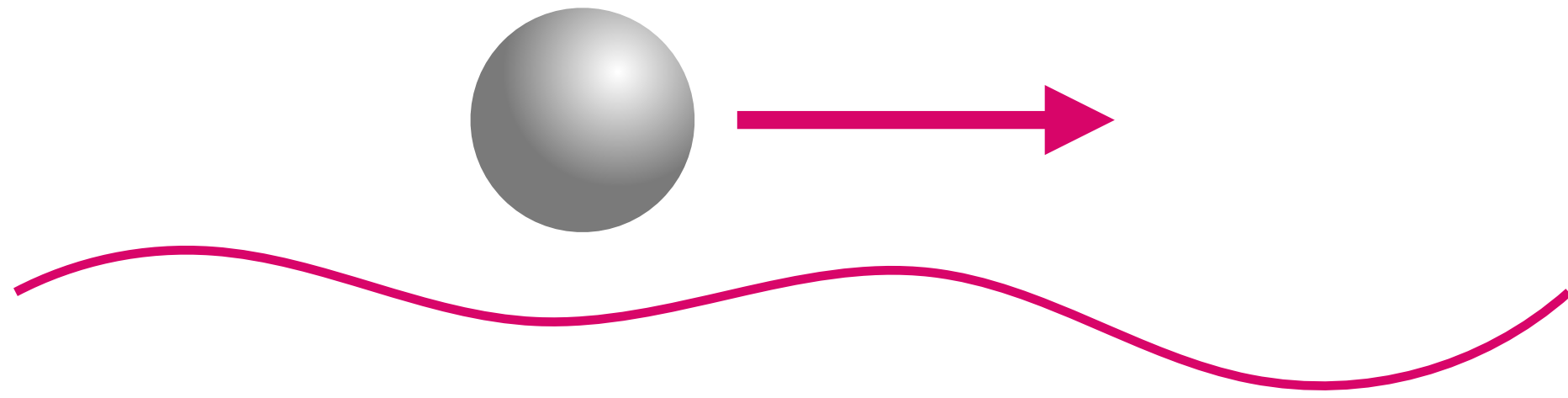
see *Weds-Fri talks for more on sources*

ACTE II

Matter is the Ruler:
Excited Cavities As
Weber Bars



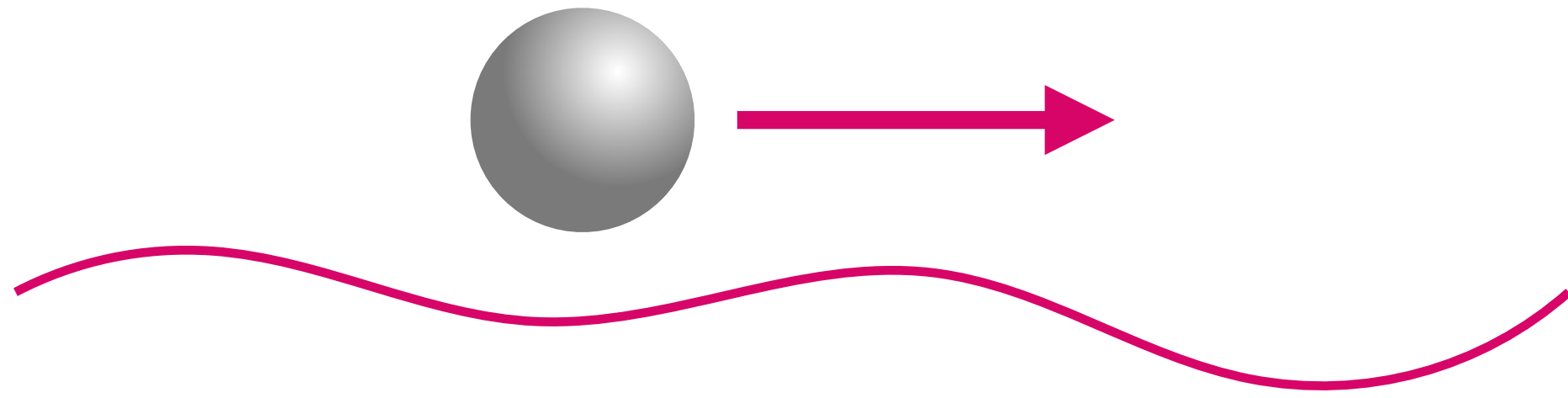
Interactions of Gravitational Waves *with masses*



$$S = - \int dt \, m \sqrt{-g_{\mu\nu} \frac{dx^\mu}{dt} \frac{dx^\nu}{dt}}$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad g^{\mu\nu} = \eta^{\mu\nu} - h^{\mu\nu}$$

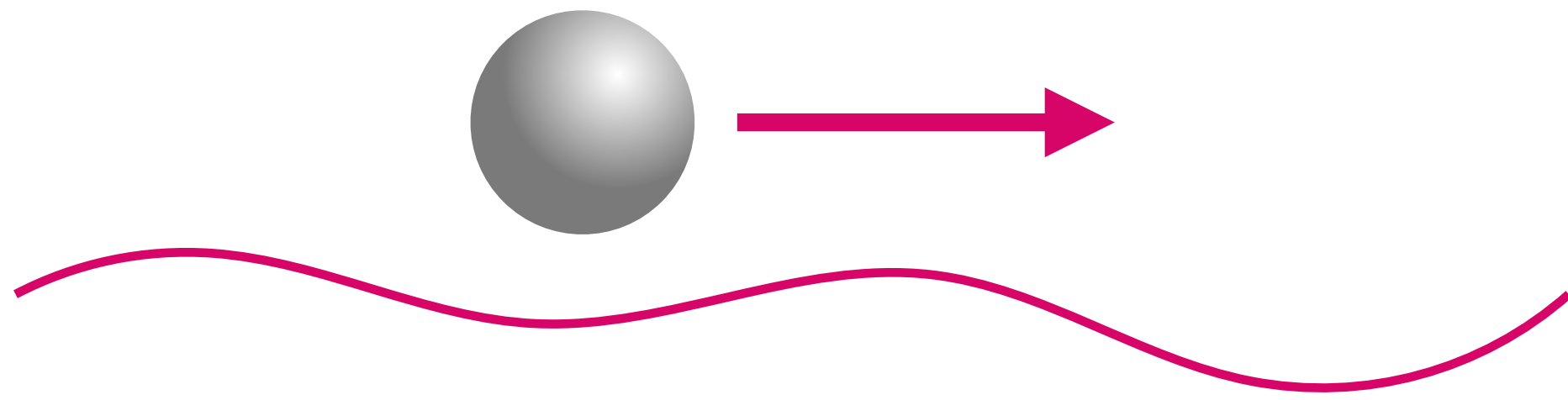
Interactions of Gravitational Waves *with masses*



$$S = - \int dt \, m \sqrt{-g_{\mu\nu} \frac{dx^\mu}{dt} \frac{dx^\nu}{dt}}$$

Equation of motion: $\frac{d^2 x^\mu}{d\tau^2} + \Gamma_{\nu\rho}^\mu(x) \frac{dx^\nu}{d\tau} \frac{dx^\rho}{d\tau} = 0$ $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad g^{\mu\nu} = \eta^{\mu\nu} - h^{\mu\nu}$

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Effect of GW encoded in Christoffel symbol $\Gamma \propto \partial h$

Encore: Framing the Question

Work in appropriate reference frame!

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Detector in Local Inertial Frame (LIF)

Encore: Framing the Question

Work in appropriate reference frame!

Detector in Local Inertial Frame (LIF)

Effect of GW in LIF is that of a Newtonian Force

Encore: Framing the Question

Work in appropriate reference frame!

Detector in Local Inertial Frame (LIF)

Effect of GW in LIF is that of a Newtonian Force

$$\frac{d^2 x_i}{d\tau^2} \simeq -\partial_i \Gamma_{00}^j x^i$$

Encore: Framing the Question

Work in appropriate reference frame!

Detector in Local Inertial Frame (LIF)

Effect of GW in LIF is that of a Newtonian Force

$$\frac{d^2 x_i}{d\tau^2} \simeq -\partial_i \Gamma_{00}^j x^i \quad \longrightarrow \quad \frac{d^2 x_i}{d\tau^2} \simeq -\frac{F_i}{m}$$

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$$\frac{d^2 x_i}{d\tau^2} \simeq -\partial_i \Gamma_{00}^j x^i \quad \longrightarrow \quad \frac{d^2 x_i}{d\tau^2} \simeq -\frac{F_i}{m} \quad \longrightarrow \quad F_i \simeq \frac{m}{2} \ddot{h}_{ij}^{\text{TT}} x^i$$

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Long-wavelength approximation valid because materials have $c_s \ll 1$

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Long-wavelength approximation valid because materials have $c_s \ll 1$

$$ds^2 \simeq -dt^2 (1 + R_{0i0j} x^i x^j) - \frac{4}{3} dt dx^i (R_{0ijk} x^j x^k) + dx^i dx^j \left(\delta_{ij} - \frac{1}{3} R_{ikjl} x^k x^l \right) \text{ e.g. Maggiore (2007)}$$

Cur Cavis* Pt 2: Mechanical and EM Signals

* “Why Cavities?” in Latin

Cur Cavis* Pt 2: Mechanical and EM Signals

On the operation of a tunable electromagnetic detector for gravitational waves

F Pegoraro[†], E Picasso[‡] and L A Radicati^{‡§}

[†]Scuola Normale Superiore, Pisa, Italy

[‡]CERN, Geneva, Switzerland

Received 6 December 1977, in final form 20 April 1978

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Microwave Apparatus for Gravitational Waves Observation

R. Ballantini, A. Chincarini, S. Cuneo, G. Gemme^{*}, R. Parodi, A. Podestà, and R. Vaccarone
INFN and Università degli Studi di Genova, Genova, Italy

Ph. Bernard, S. Calatroni, E. Chiaveri, and R. Losito
CERN, Geneva, Switzerland

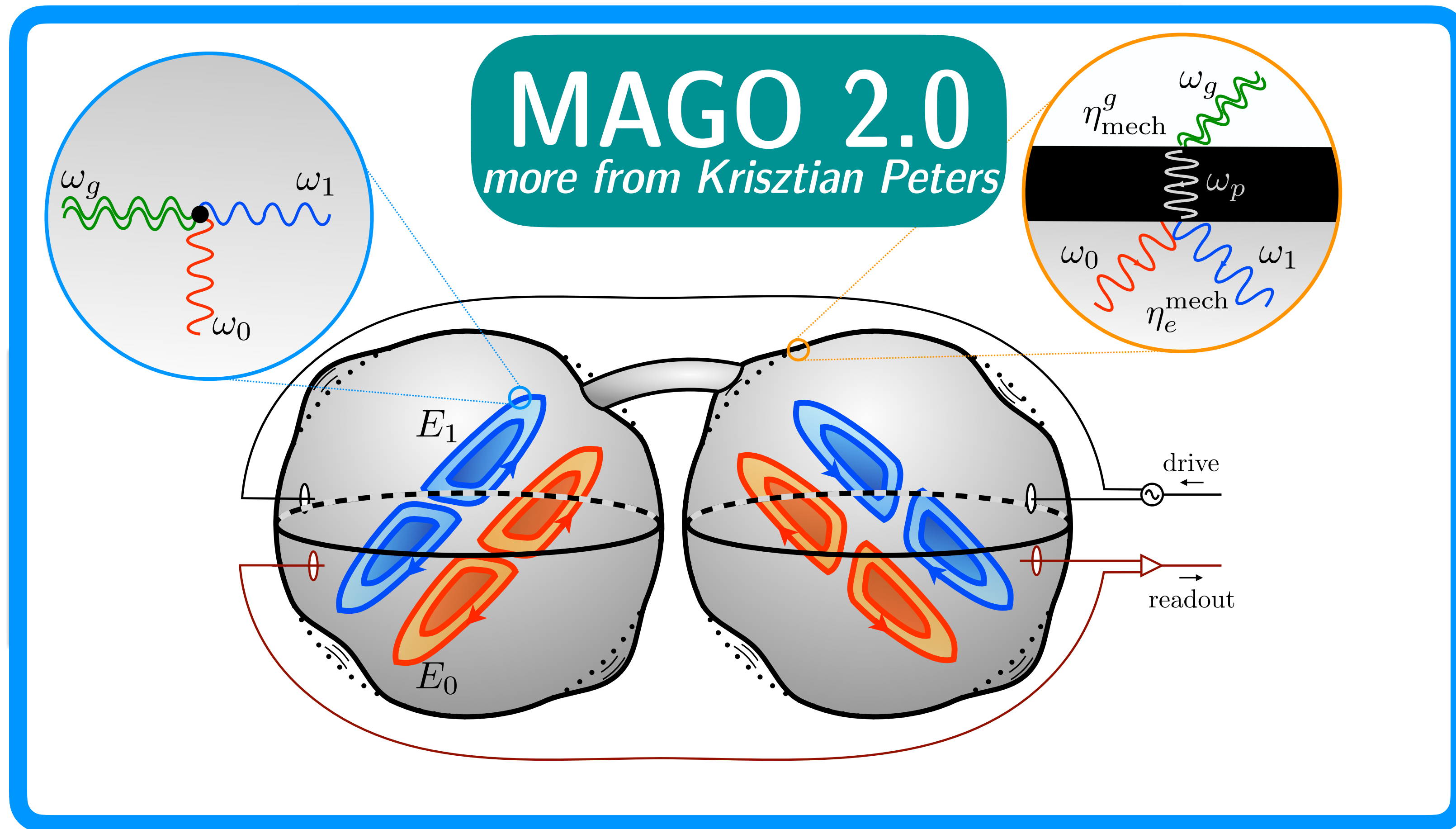
R.P. Croce, V. Galdi, V. Pierro, and I.M. Pinto
INFN, Napoli, and Università degli Studi del Sannio, Benevento, Italy

E. Picasso
*INFN and Scuola Normale Superiore, Pisa, Italy and
CERN, Geneva, Switzerland*



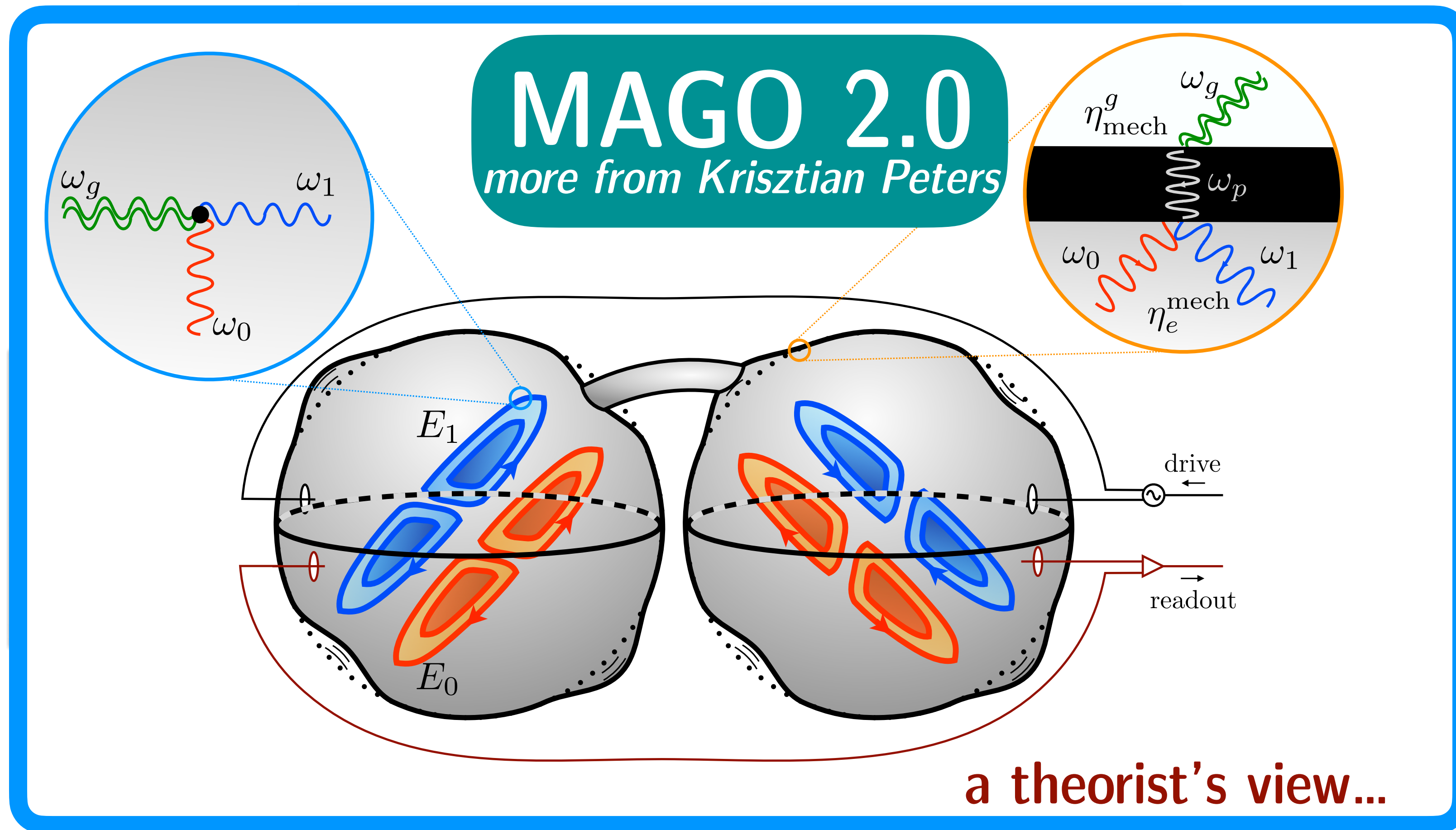
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Cur Cavis* Pt 2: Mechanical and EM Signals



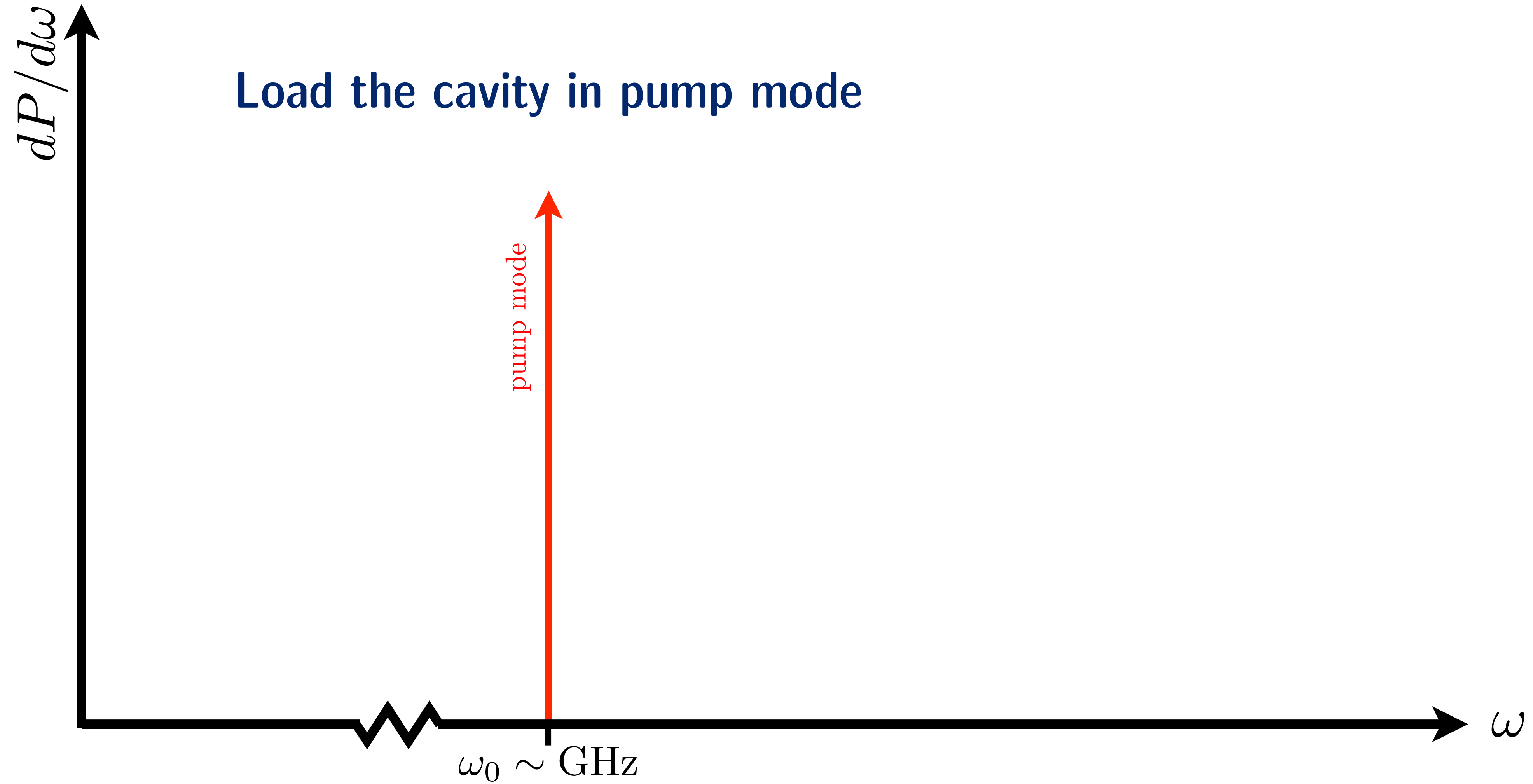
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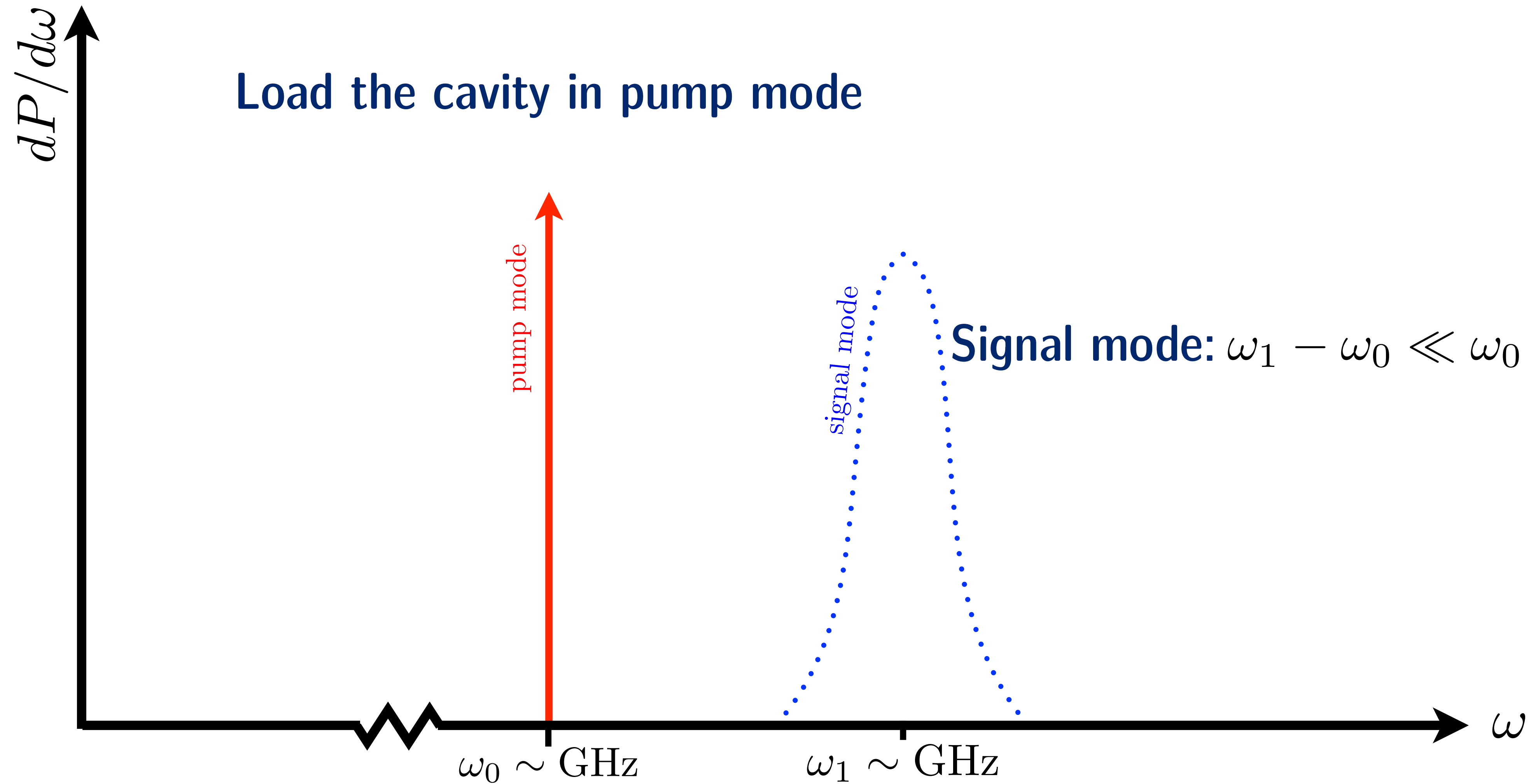


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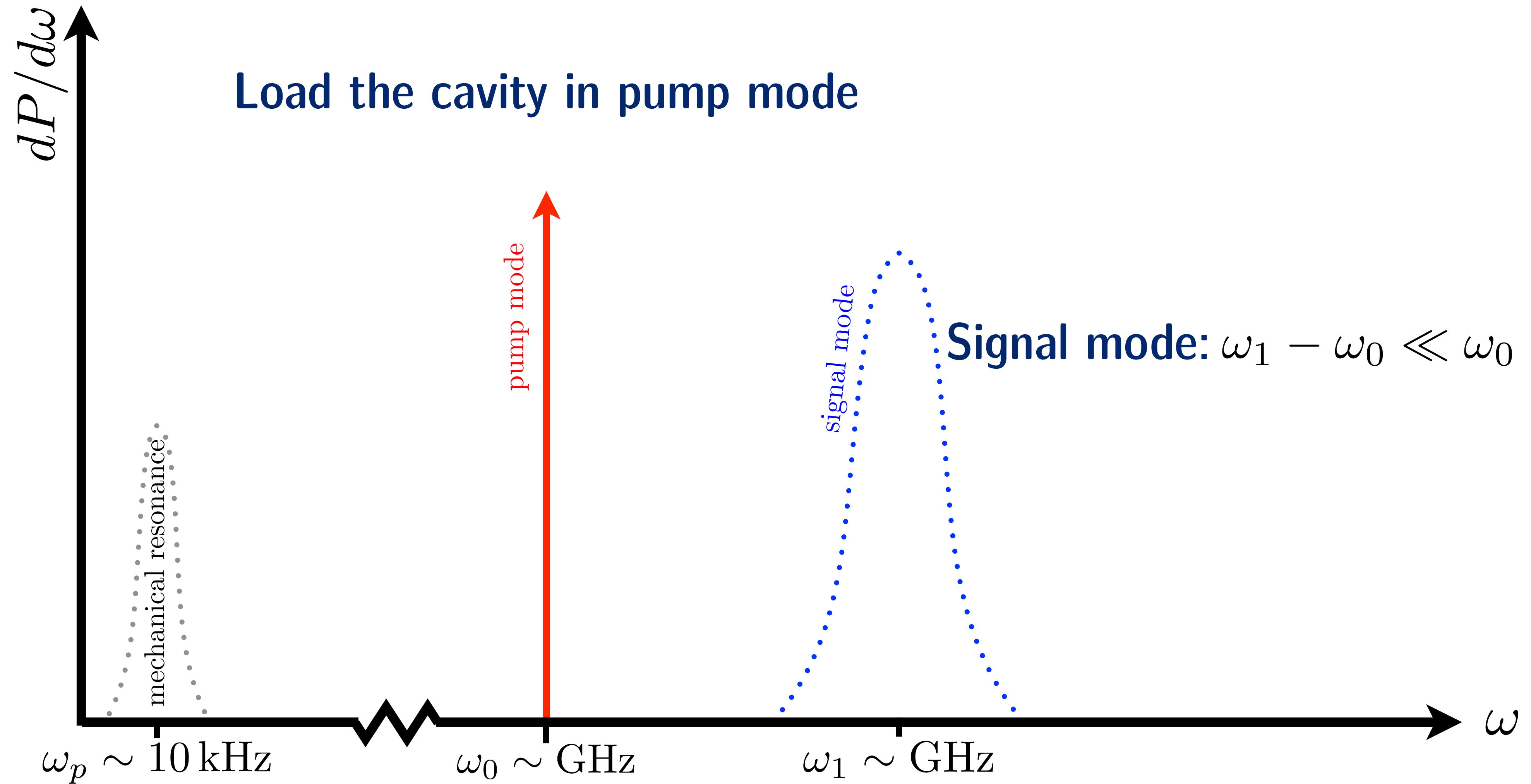
MAGO 2.0



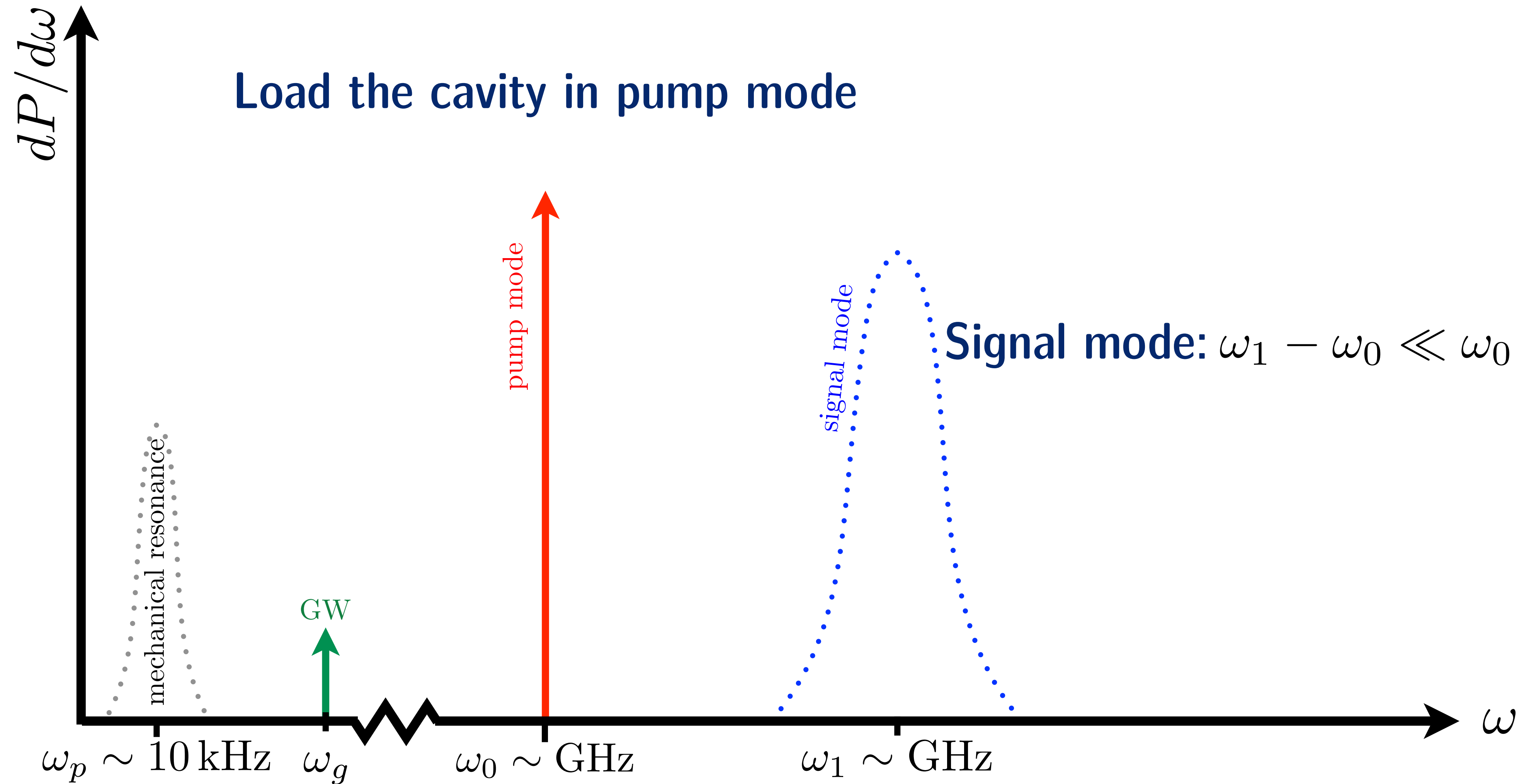
MAGO 2.0



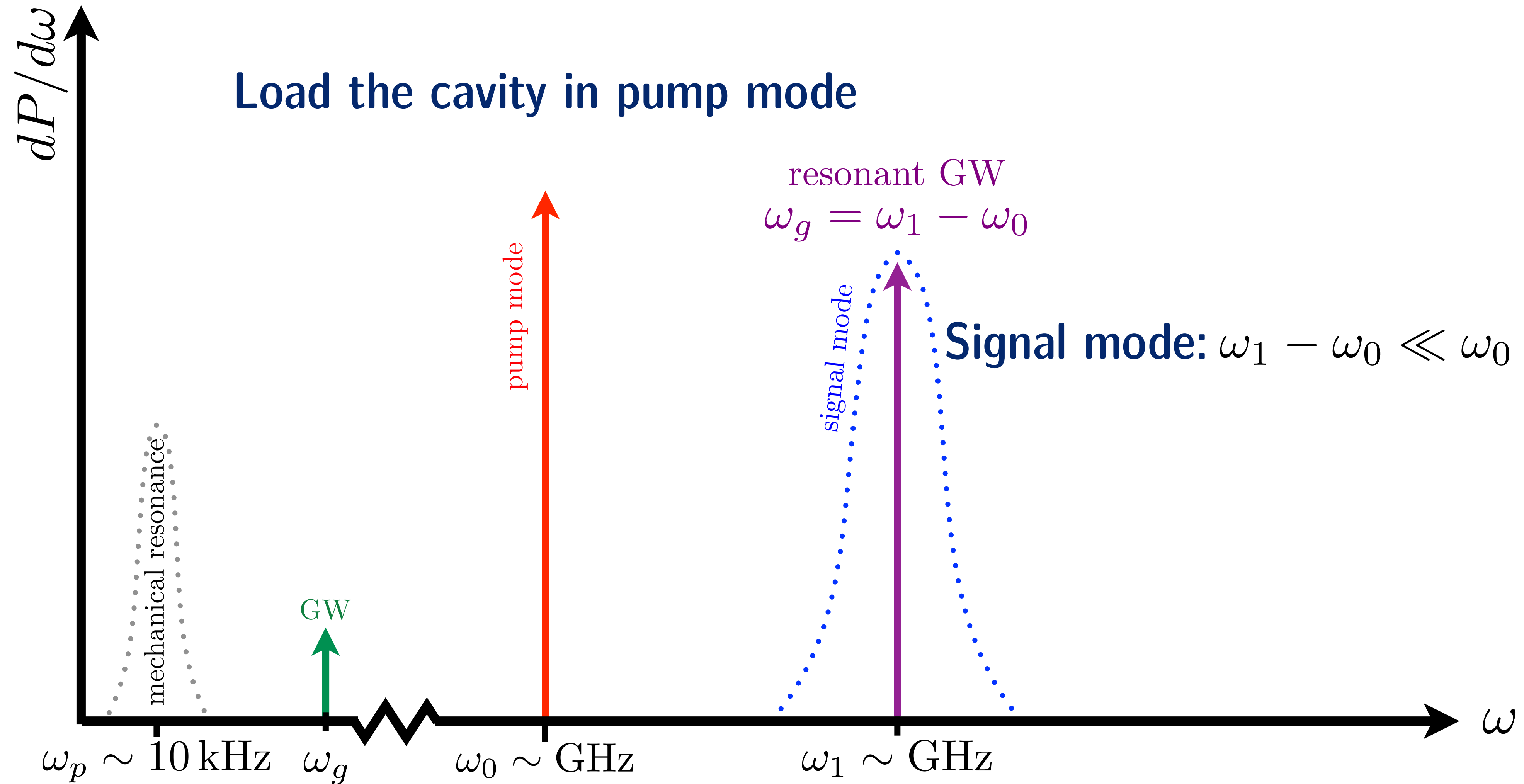
MAGO 2.0



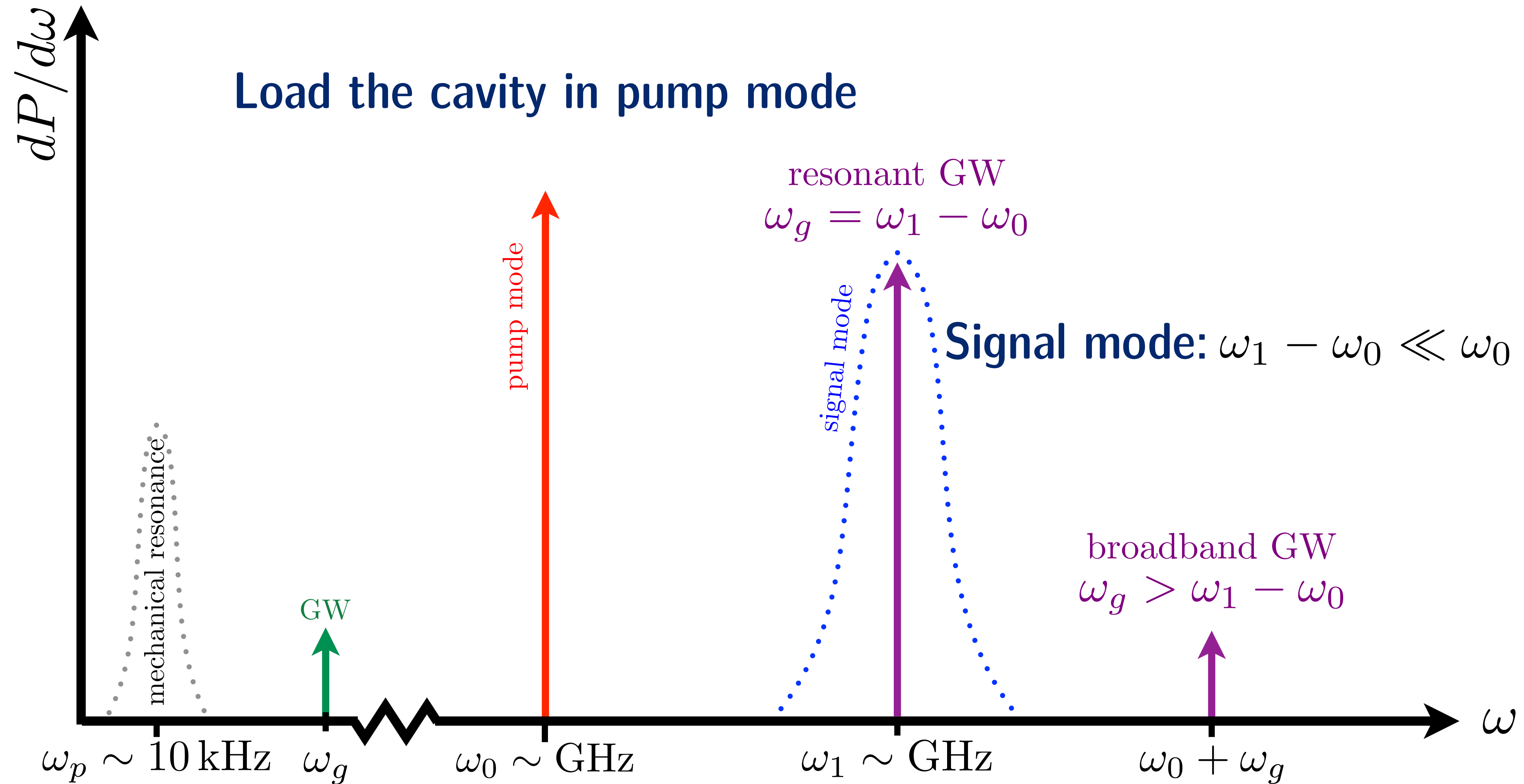
MAGO 2.0



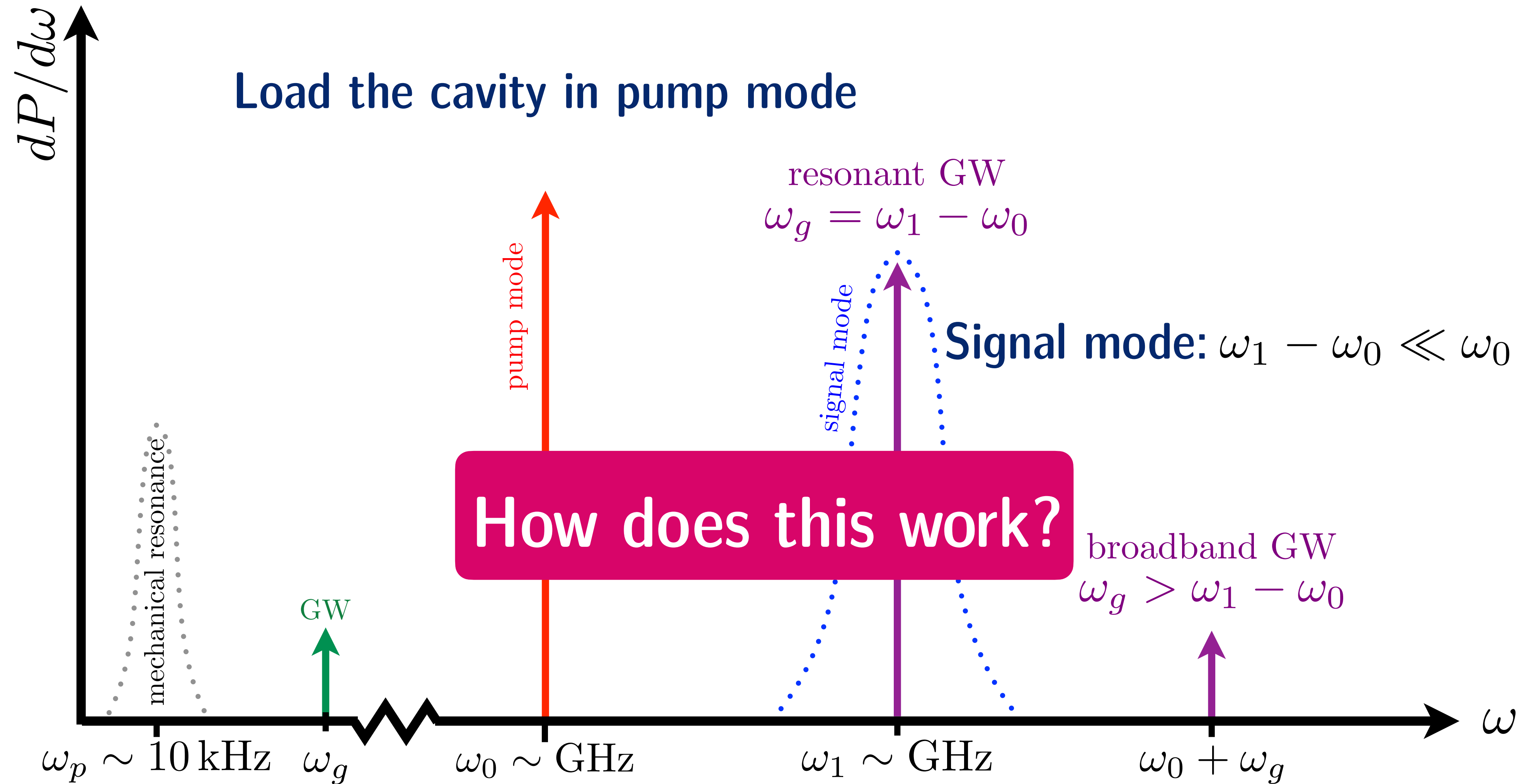
MAGO 2.0



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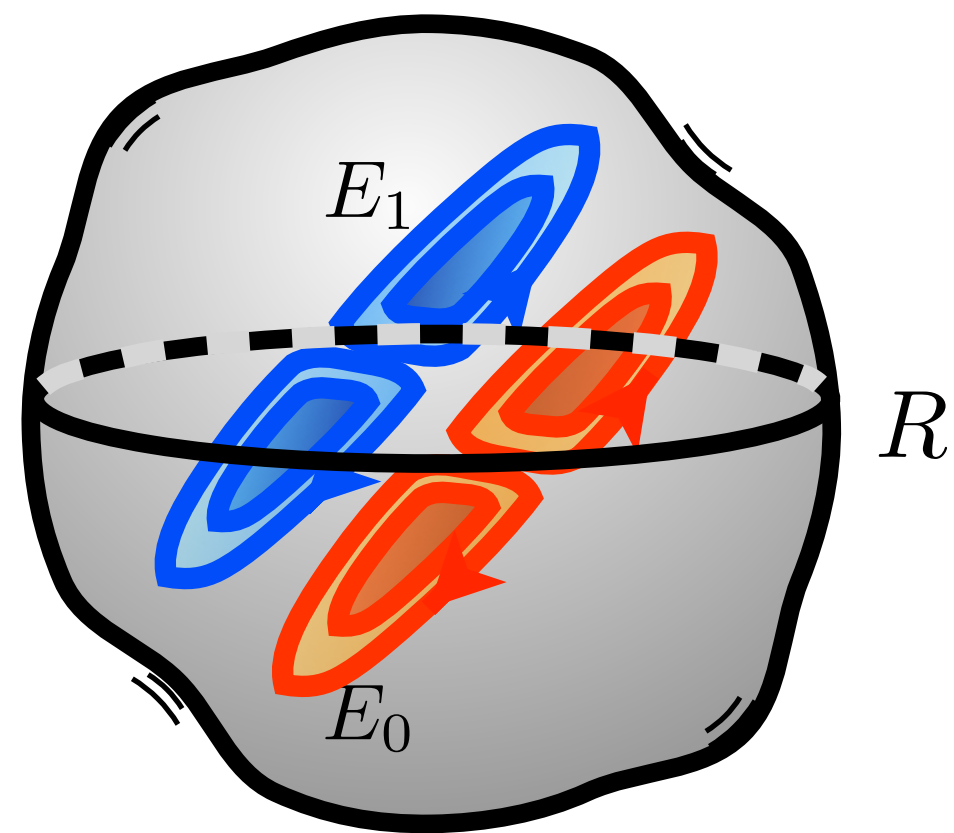


MAGO 2.0

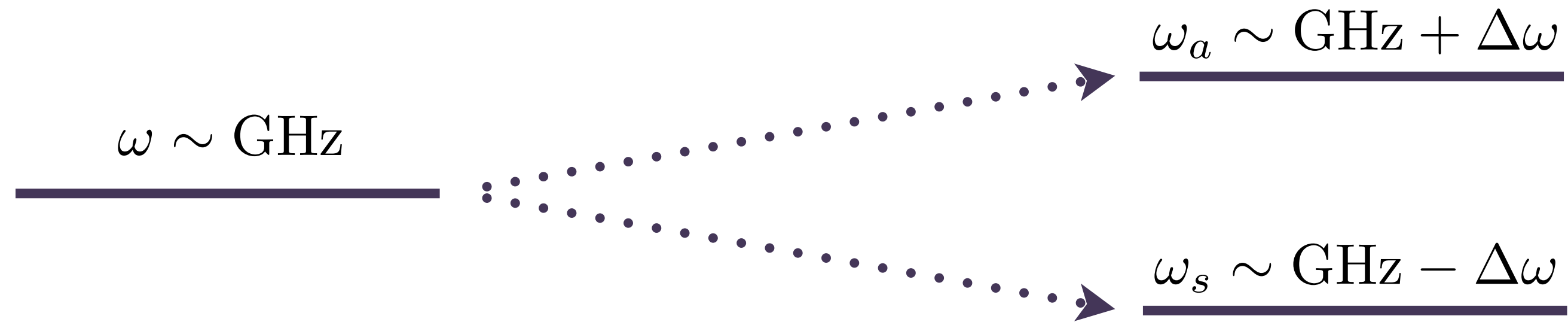


MAGO 2.0

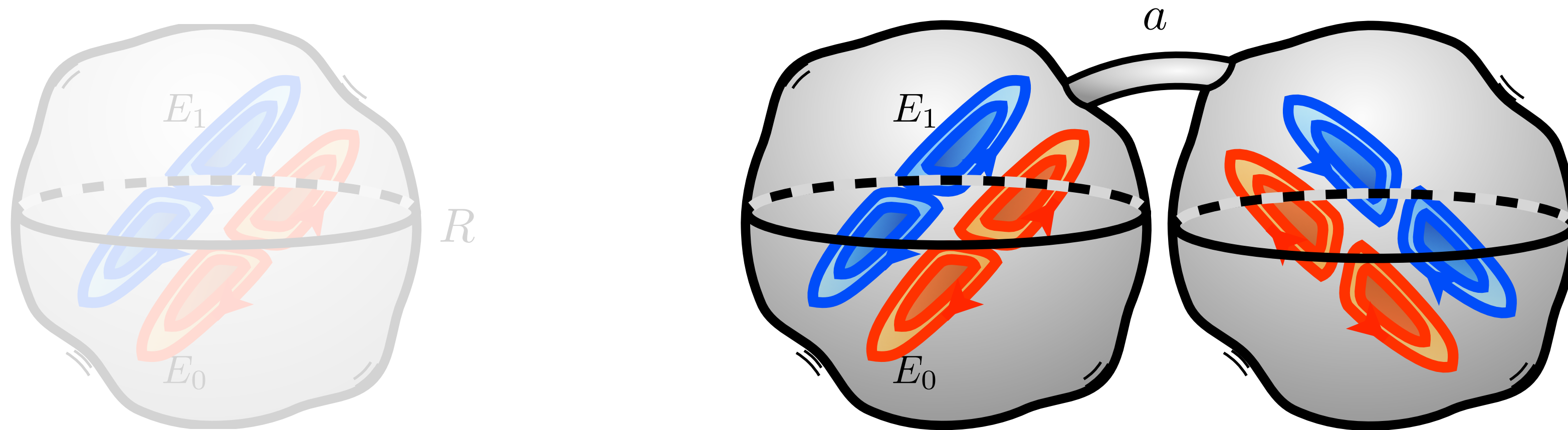
$\omega \sim \text{GHz}$



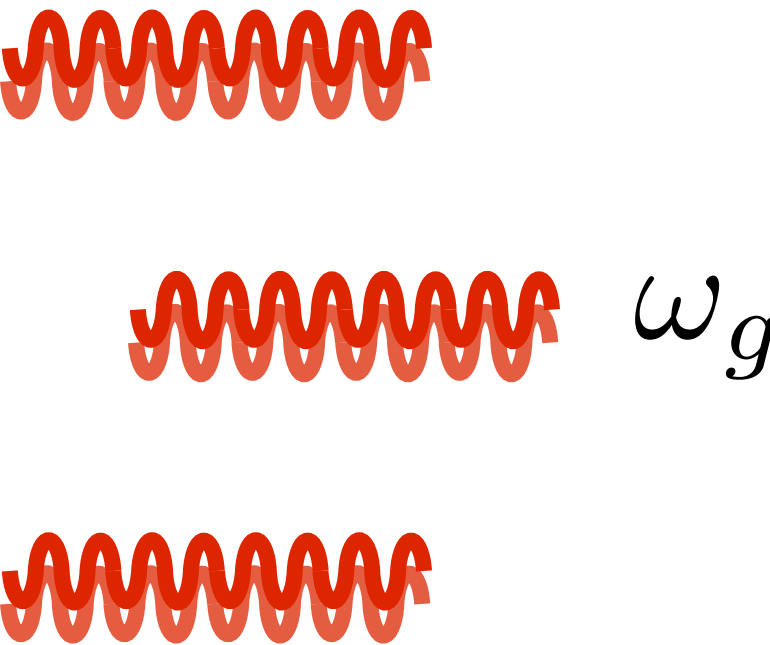
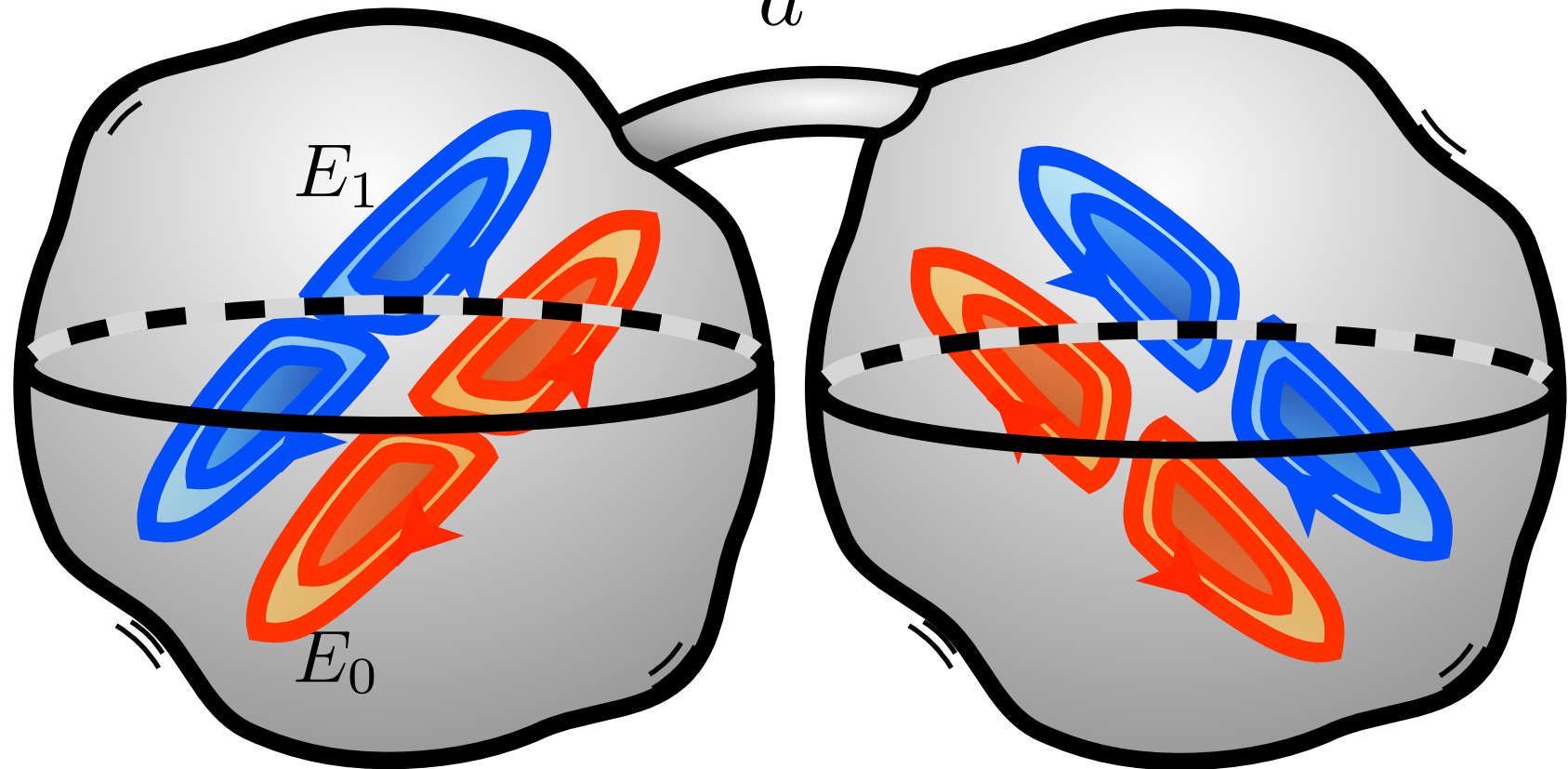
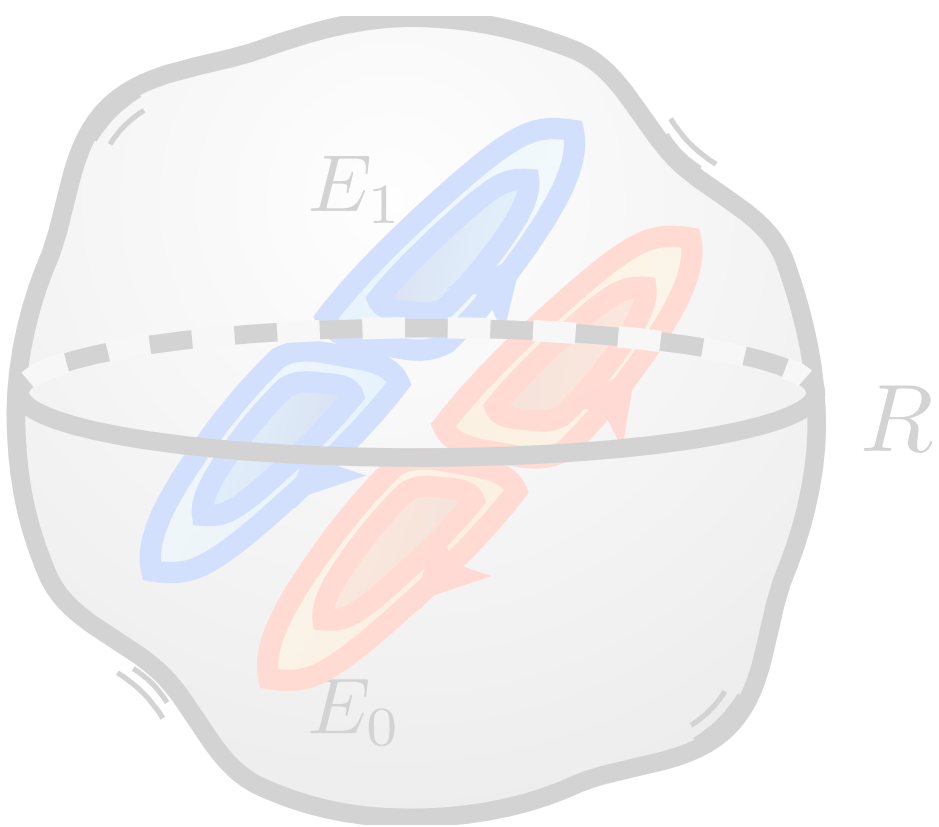
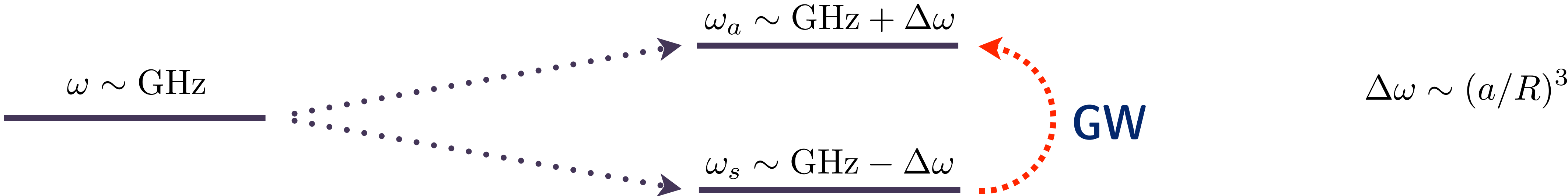
MAGO 2.0



$$\Delta\omega \sim (a/R)^3$$



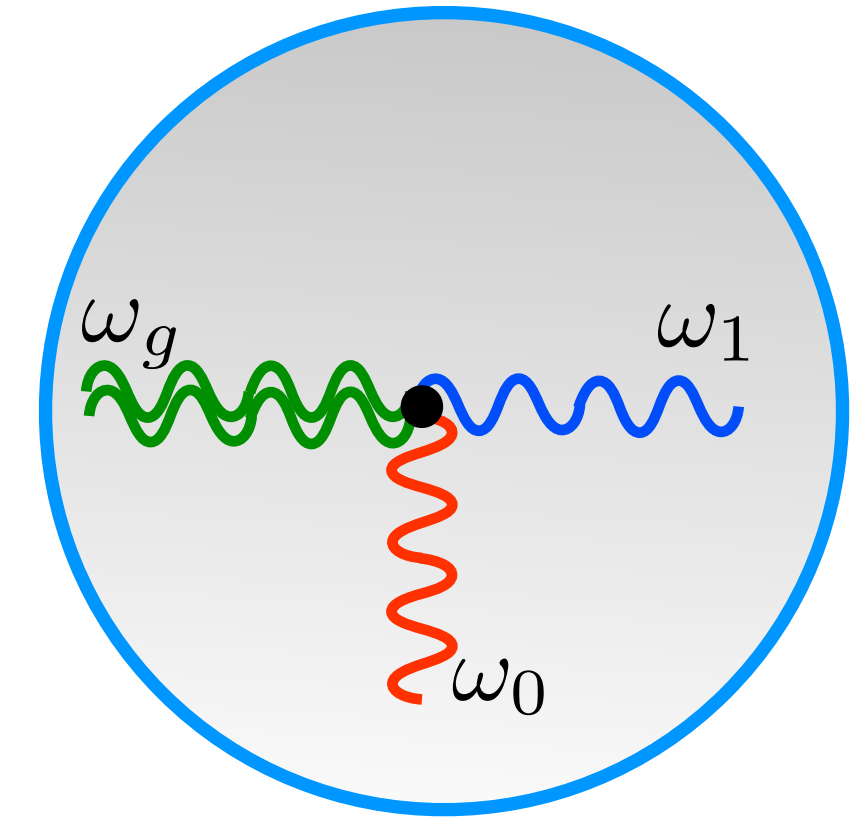
MAGO 2.0



EM and Mechanical signals

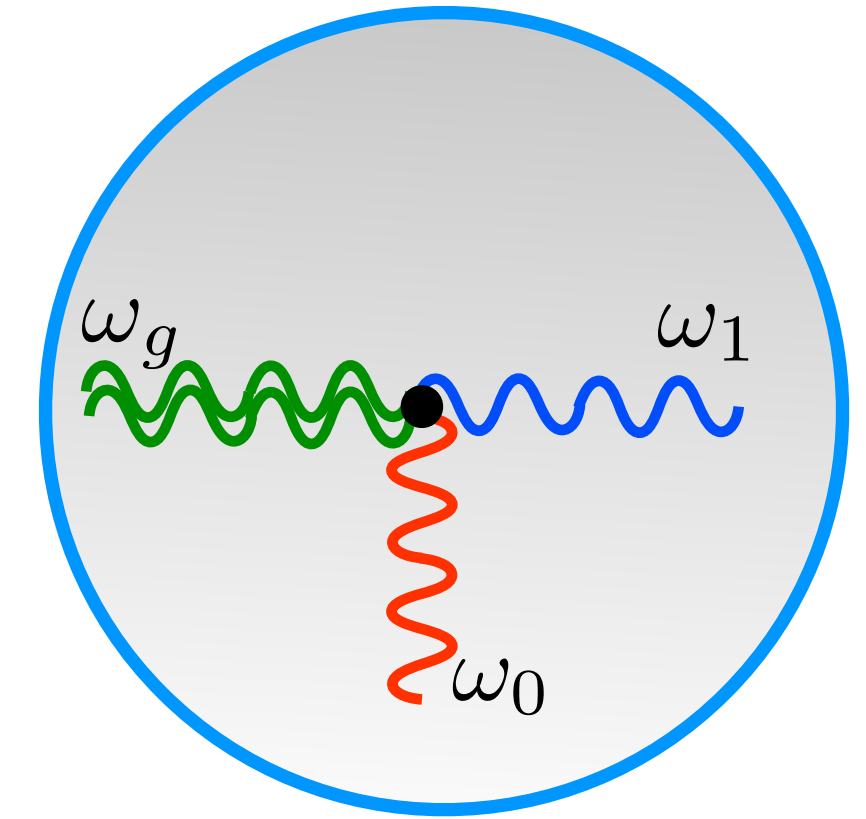
EM and Mechanical signals

Parametrics of the EM signal: $E_{\text{sig}}^{(\text{EM})} \sim Q_{\text{em}} (\omega_g L_{\text{cav}})^2 h^{\text{TT}} E_0$



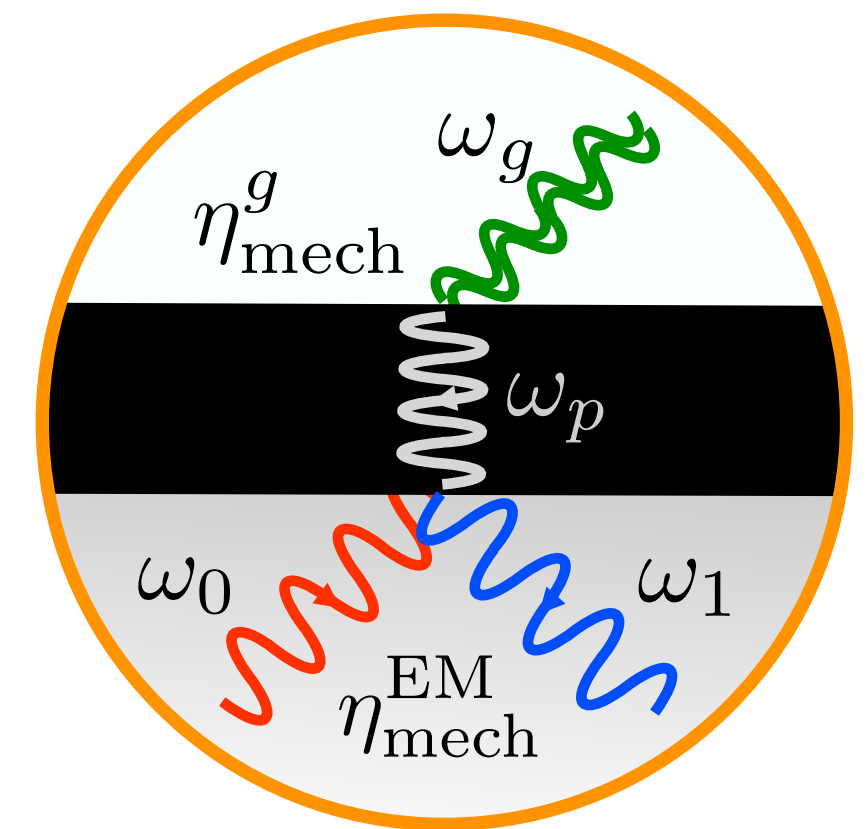
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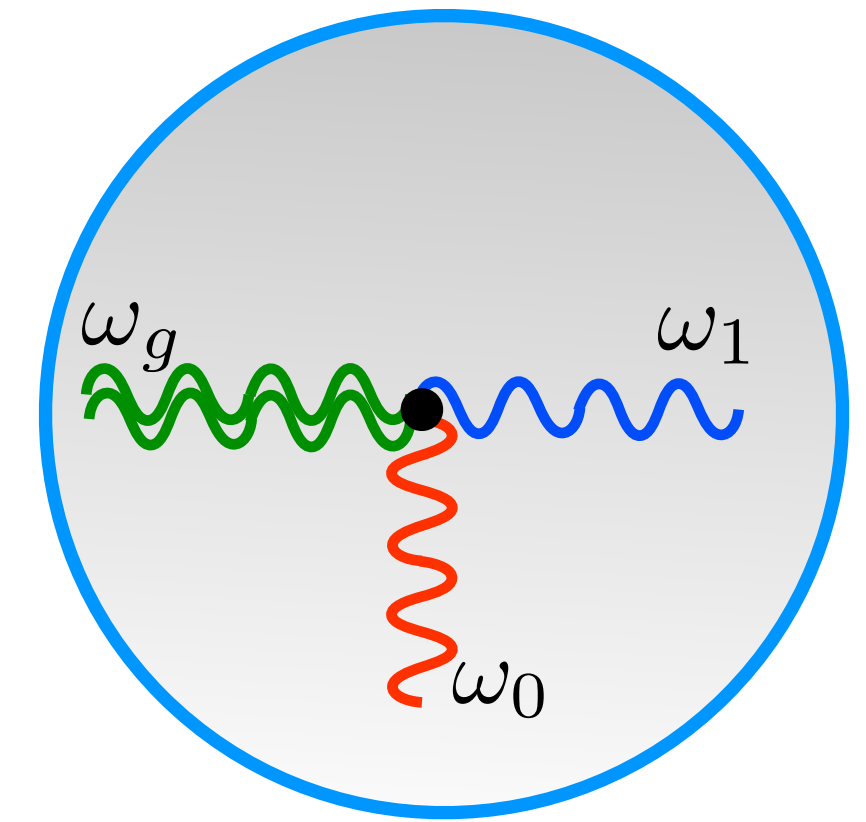
Mechanical signal:

$$E_{\text{sig}}^{(\text{mech})} \sim Q_{\text{em}} h^{\text{TT}} E_0 \min \left(1, \frac{\omega_g L_{\text{cav}}}{c_s} \right)^2$$



EM and Mechanical signals

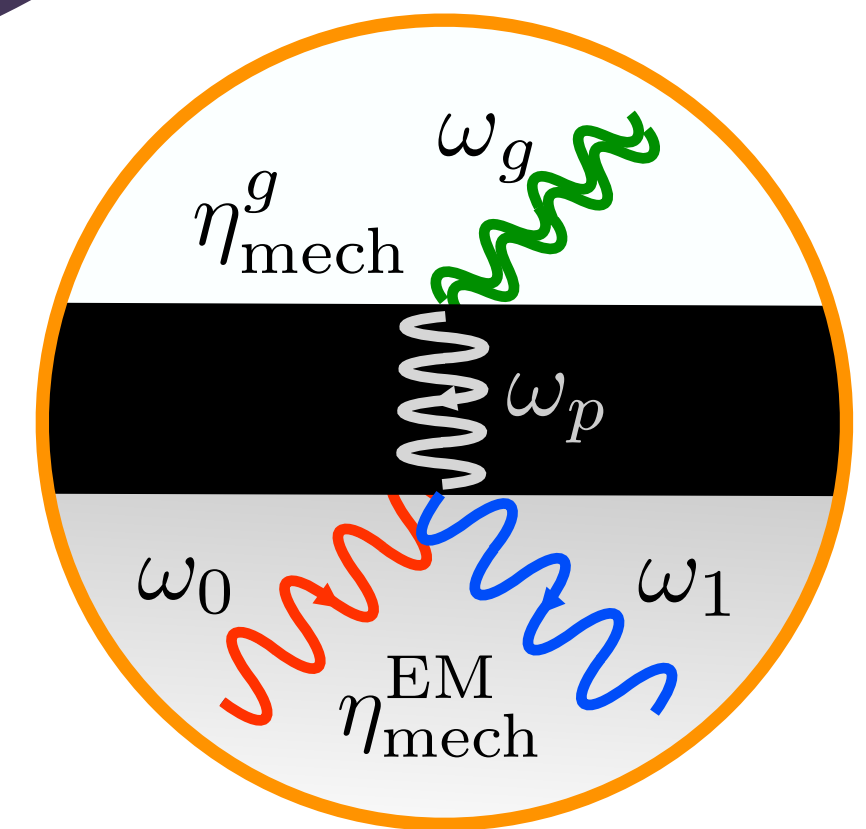
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Enhanced by $1/c_s^2 \gg 1$ (!)

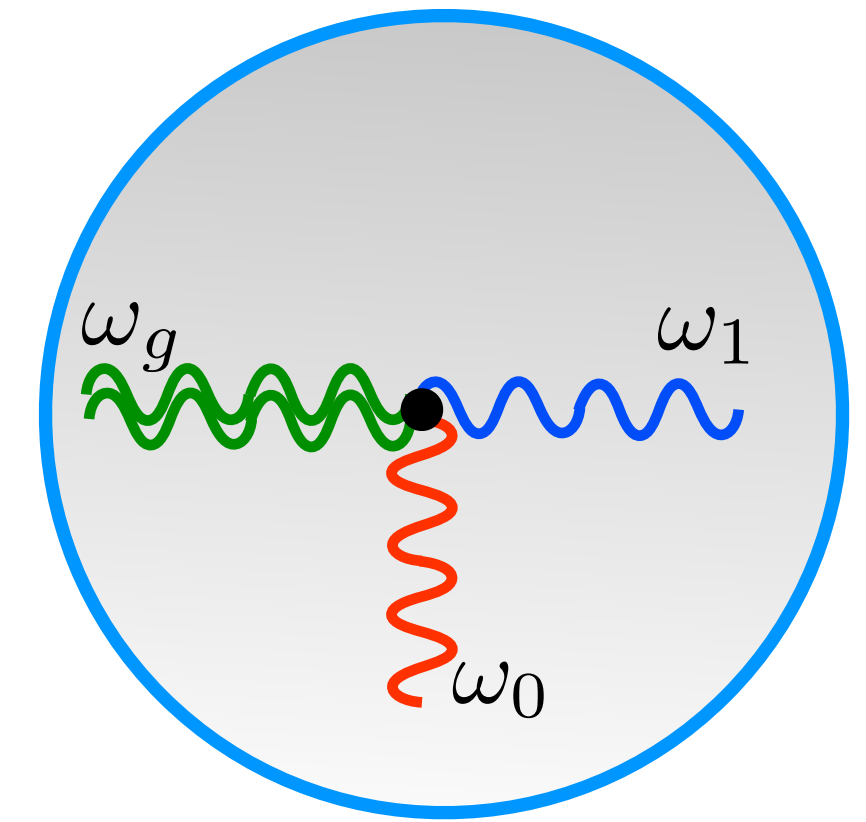
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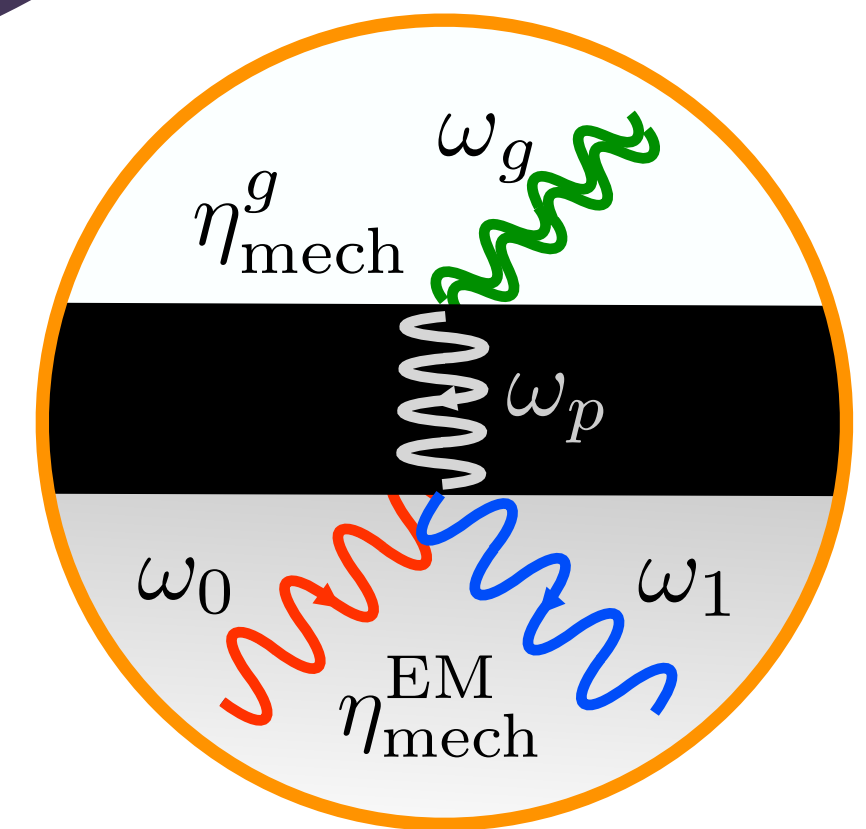


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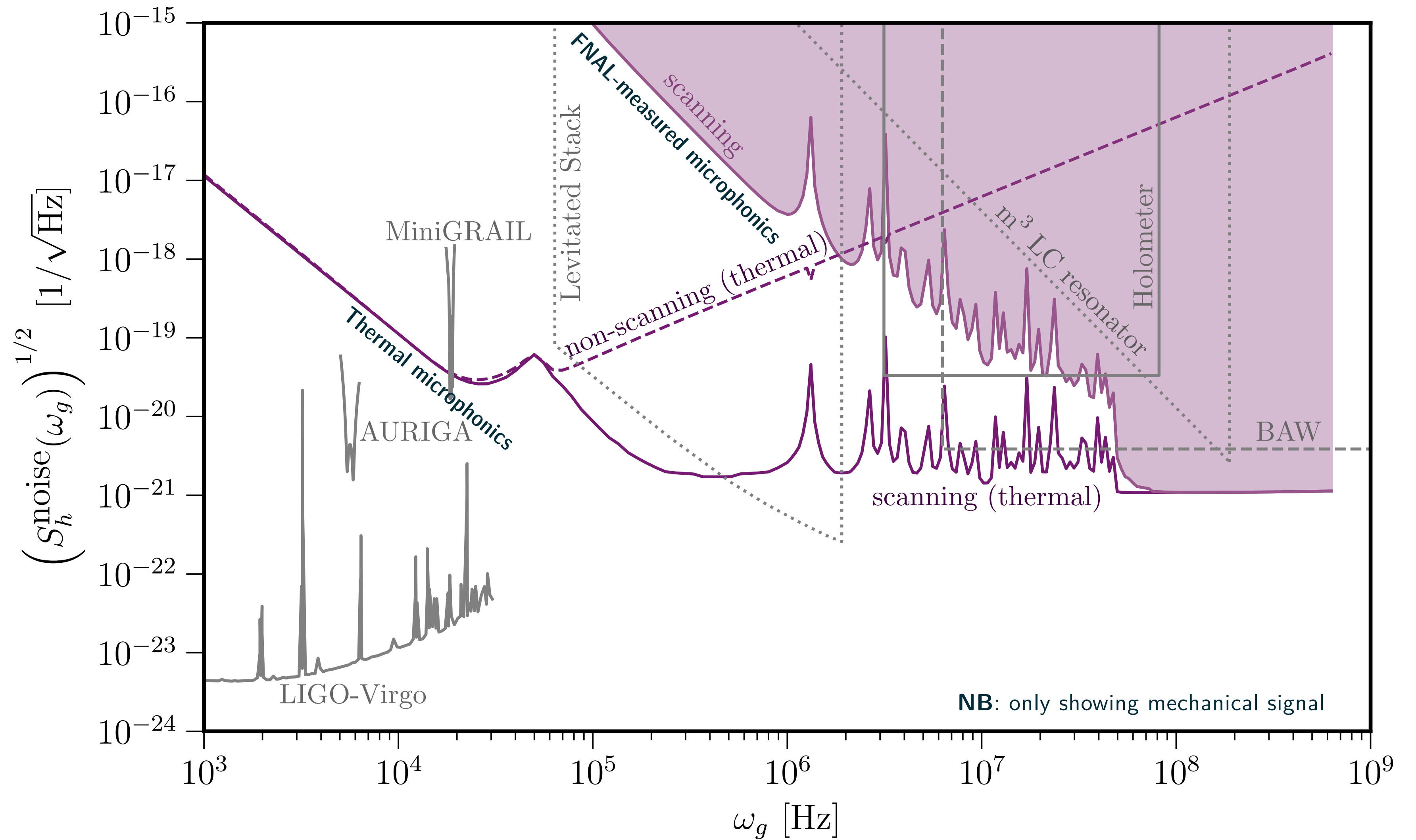
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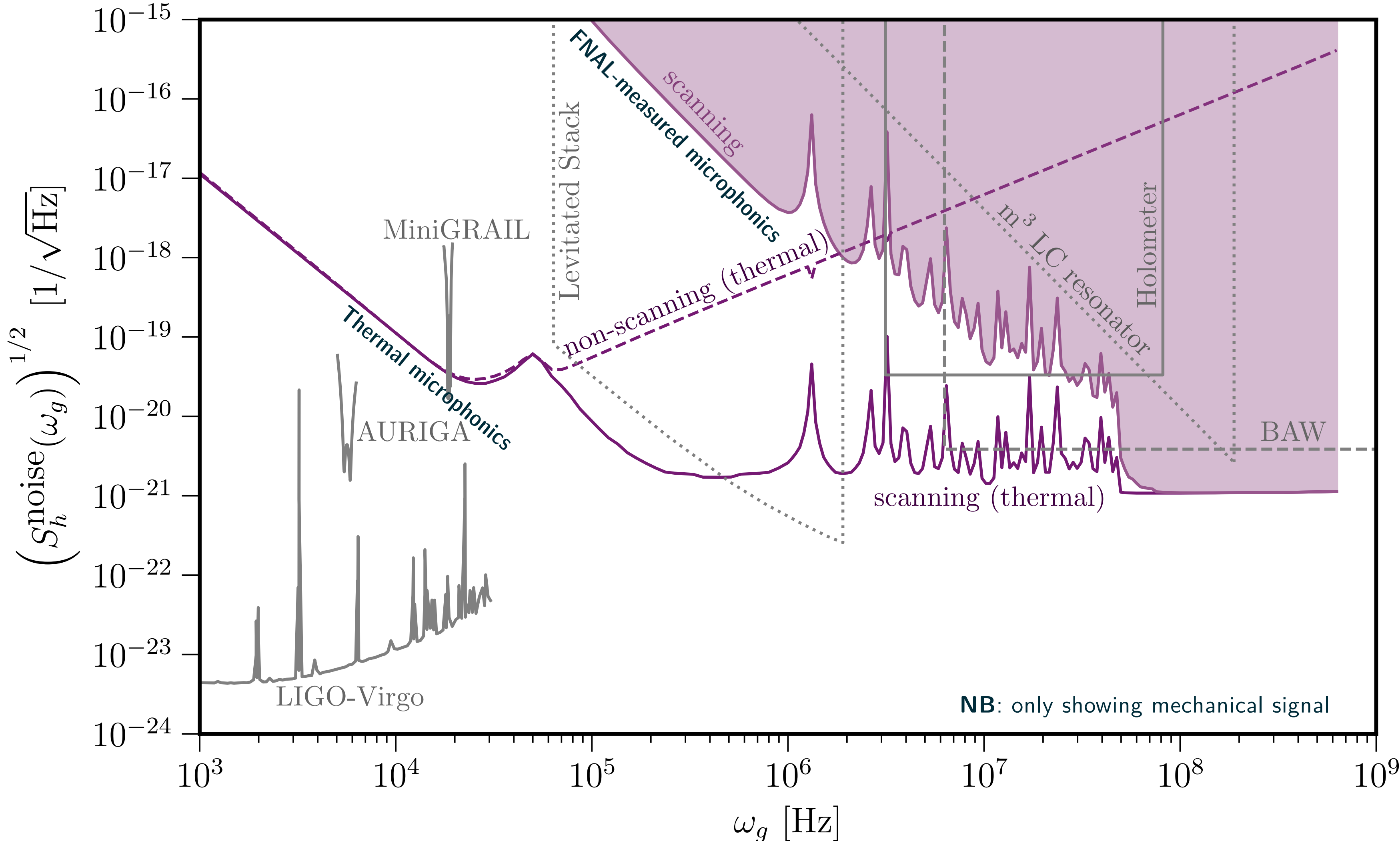
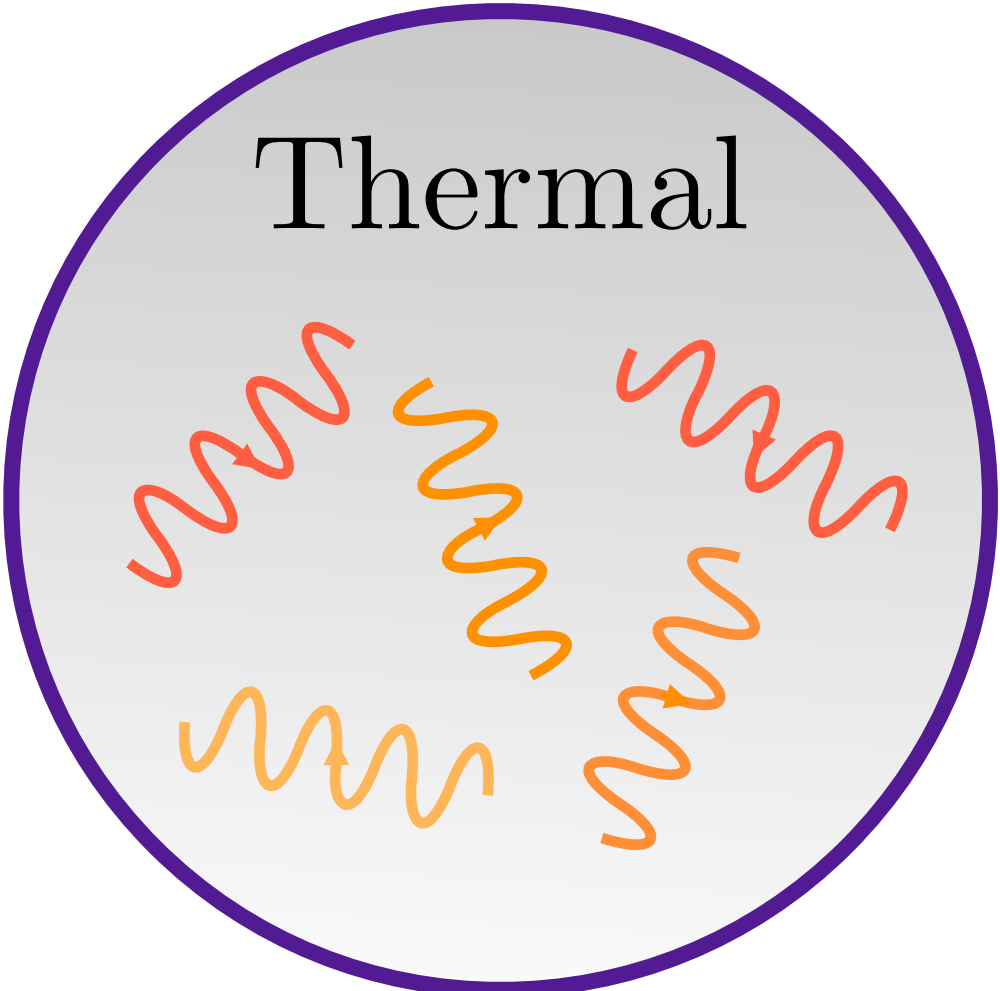
Mechanical modes less “rigid” than EM modes



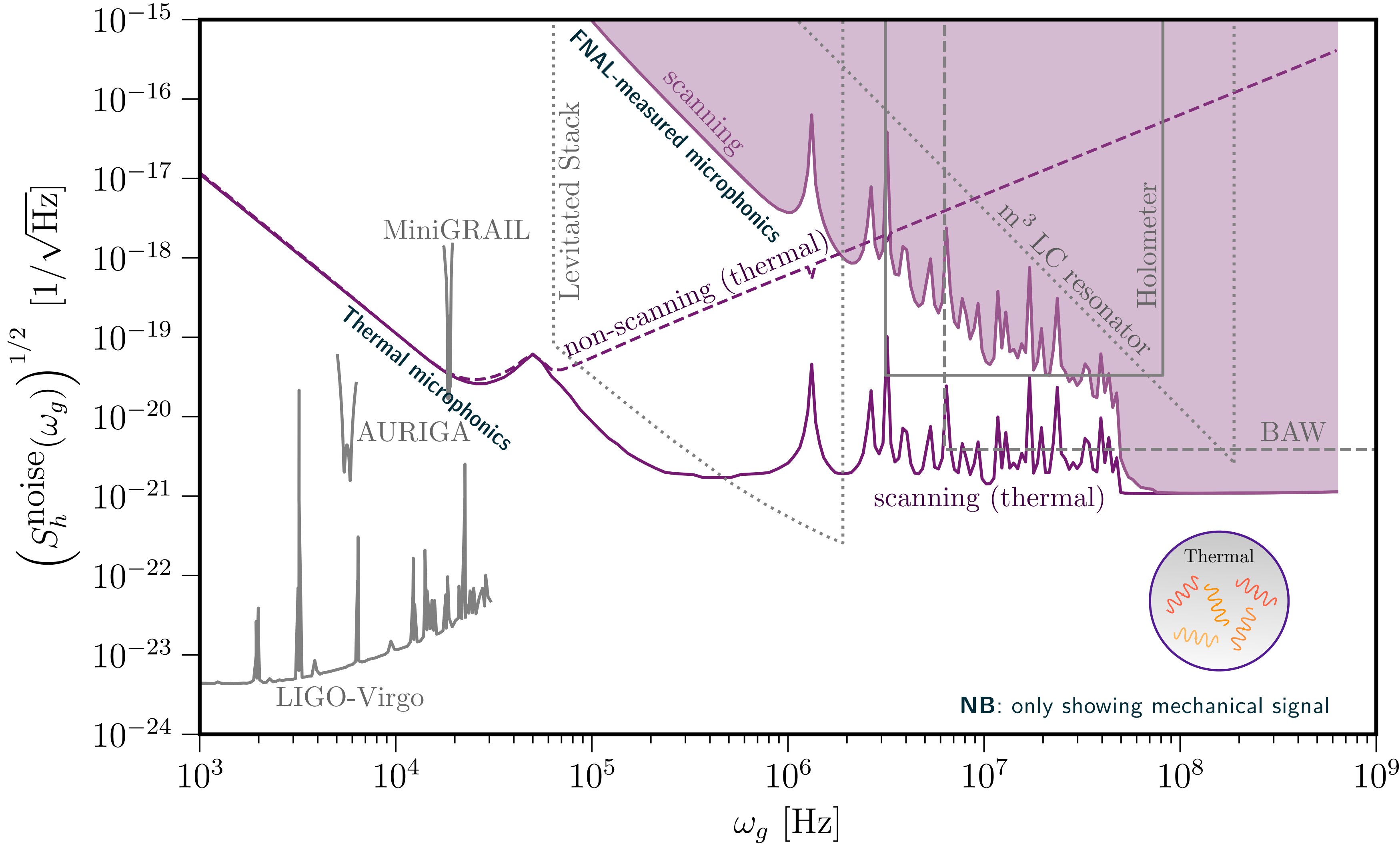
Noise in MAGO 2.0



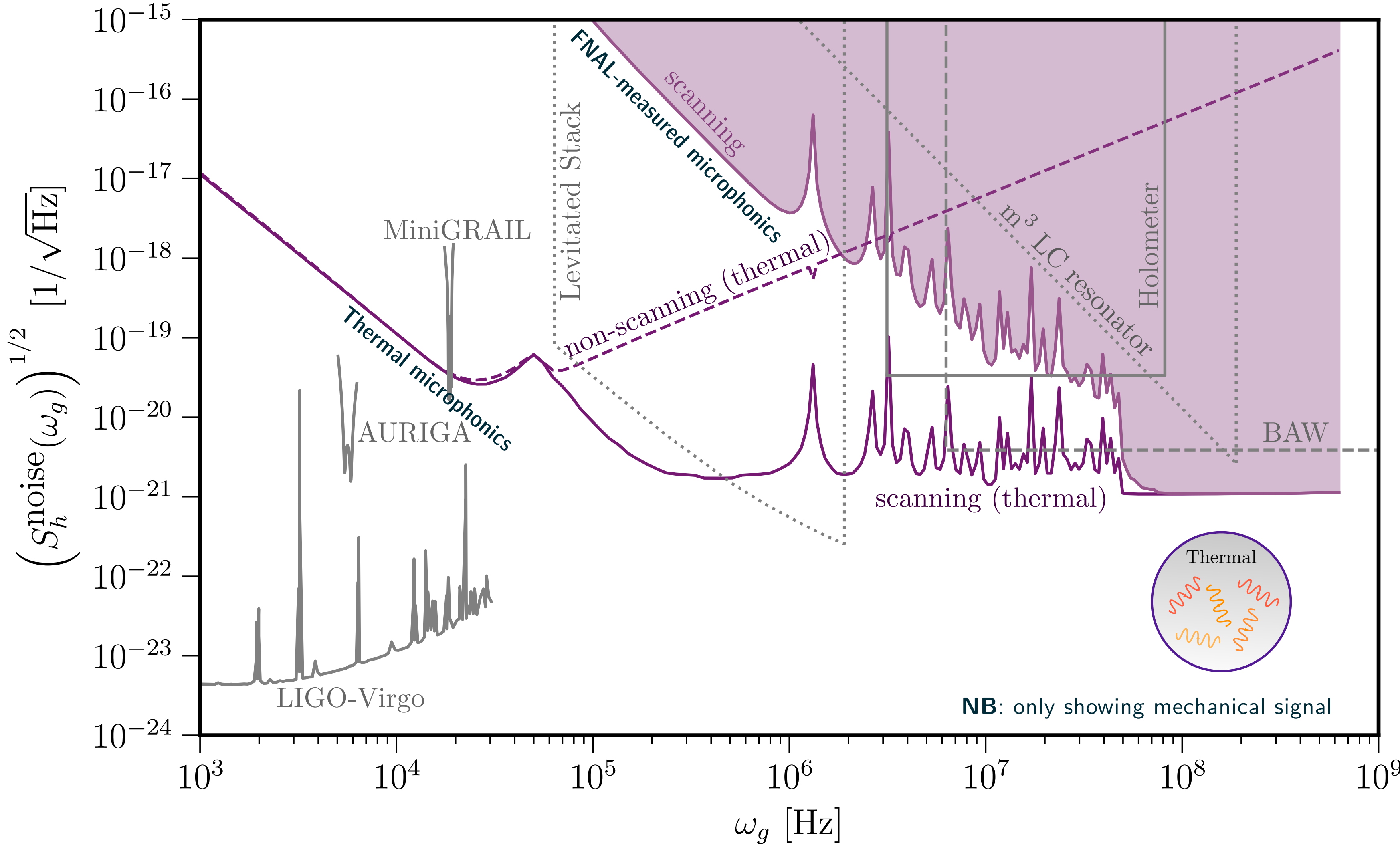
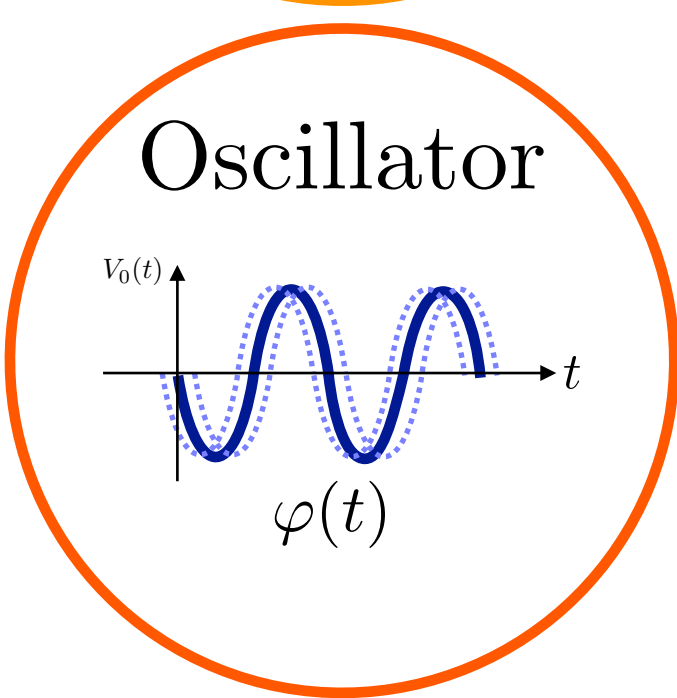
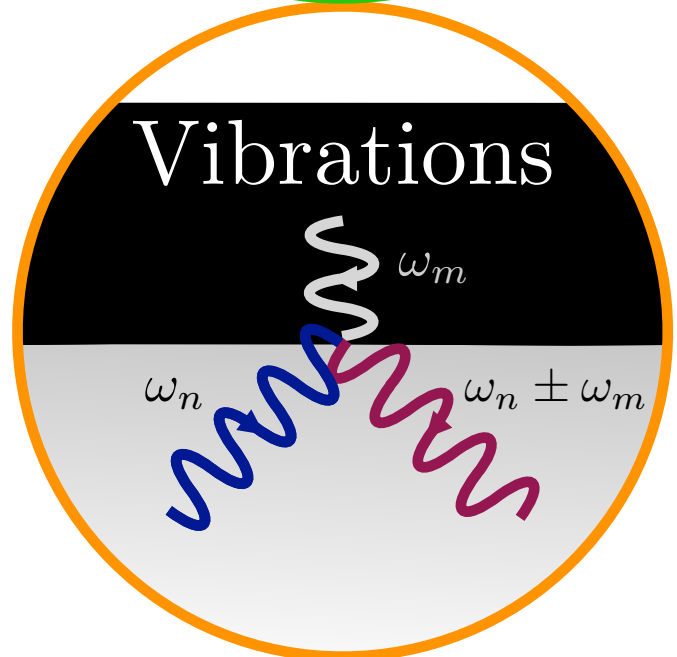
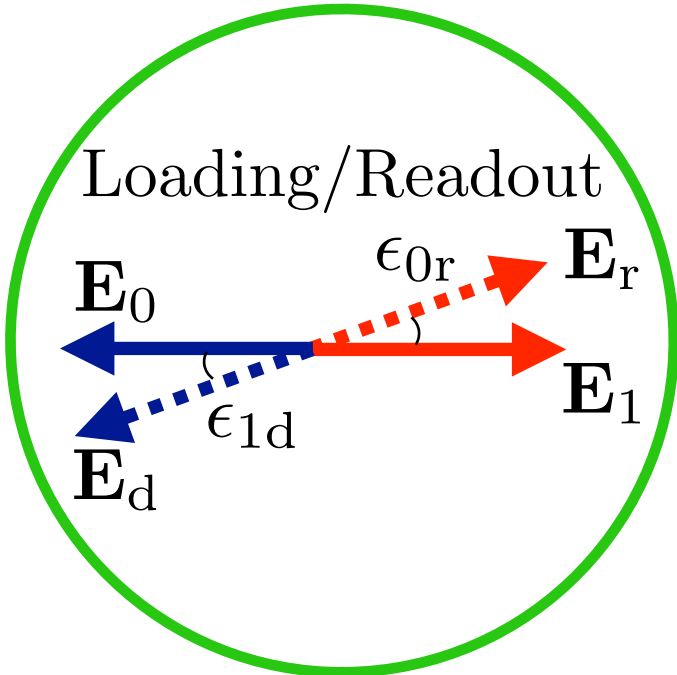
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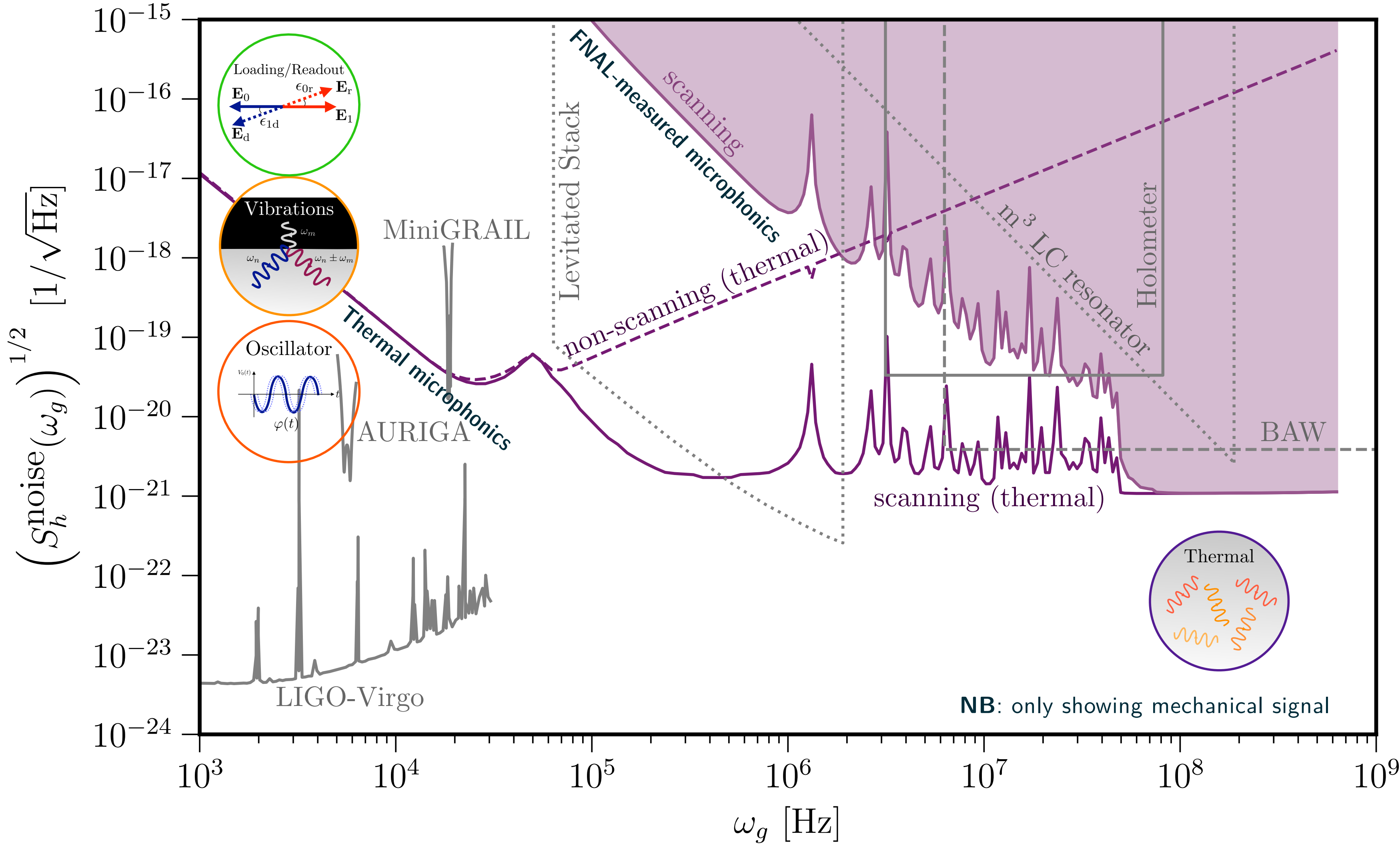
Noise in MAGO 2.0



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Why SRF Cavity as a Weber Bar?

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Thermal mechanical noise smaller for larger mass $S_{\text{noise}}^{(\text{mech.})} \propto \frac{4T\omega_p}{Q_p M}$

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Thermal EM noise-limited reach scales as $S_h^{(\text{EM})} \propto \frac{1}{Q} \frac{T}{\omega_0 U_0}$

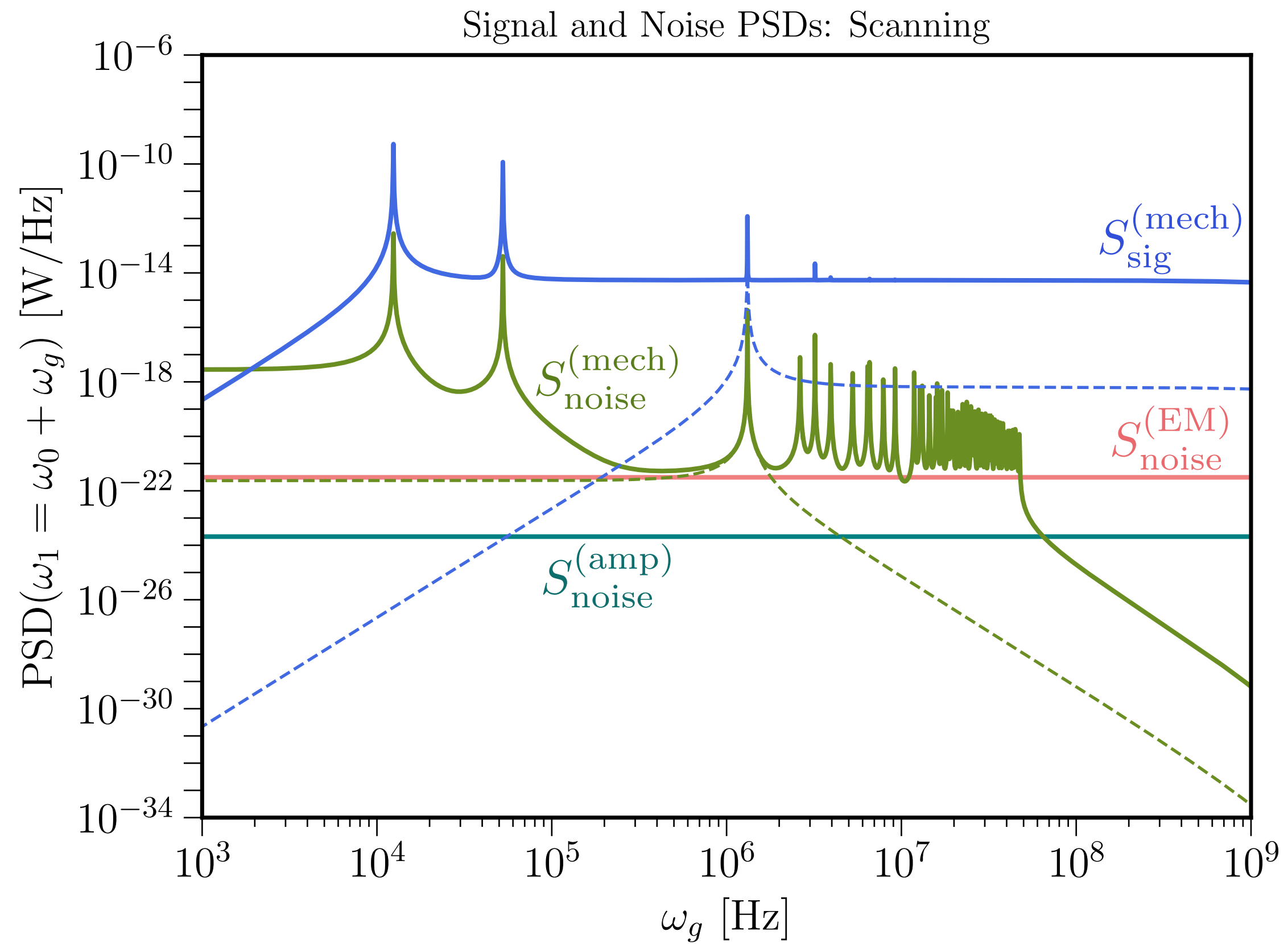
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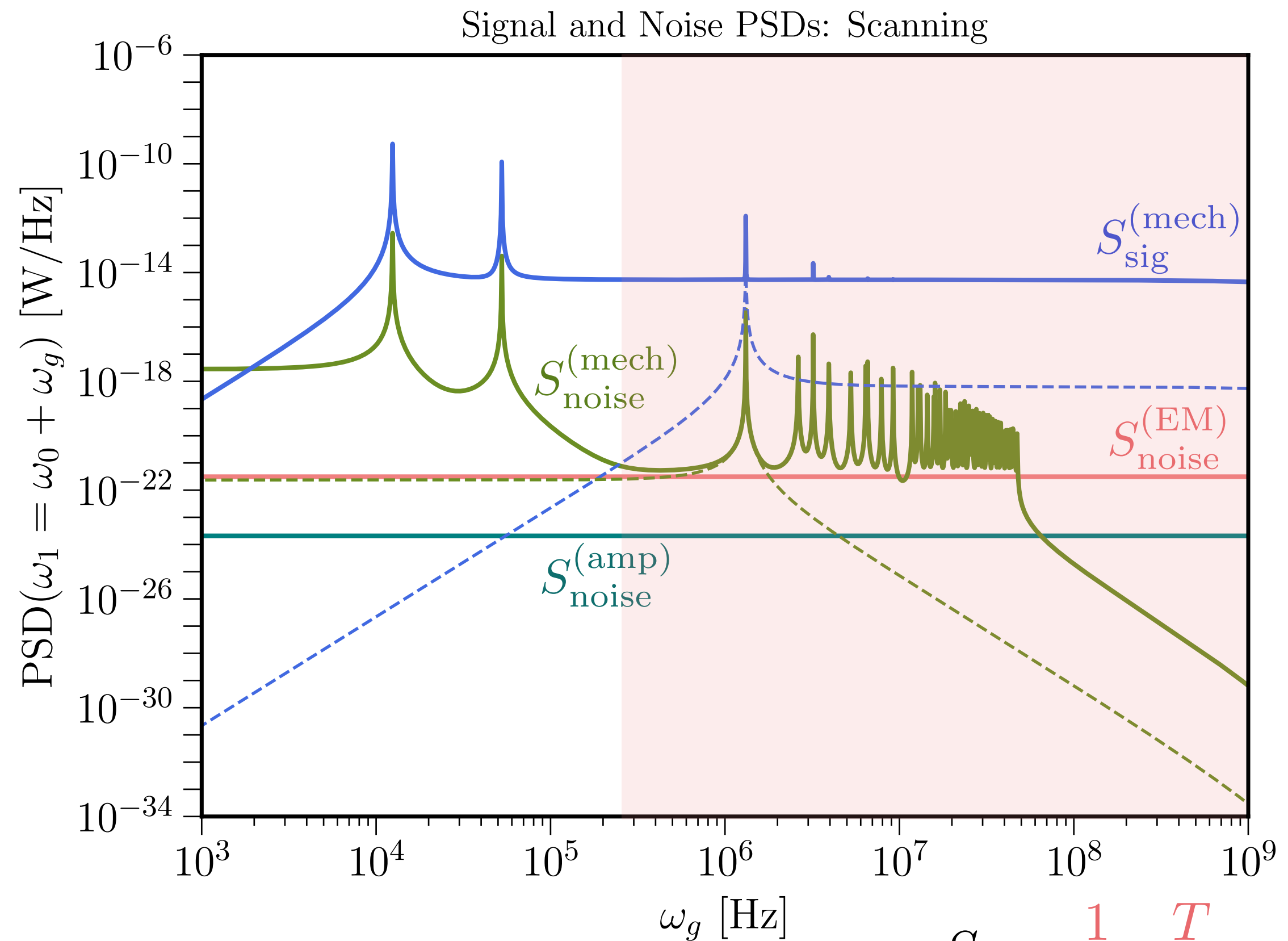
Thermal EM noise-limited reach scales as $S_h^{(\text{EM})} \propto \frac{1}{Q} \frac{T}{\omega_0 U_0}$

SRF Cavity w/ $Q \sim 10^{10}$ means better sensitivity

Noise in MAGO 2.0



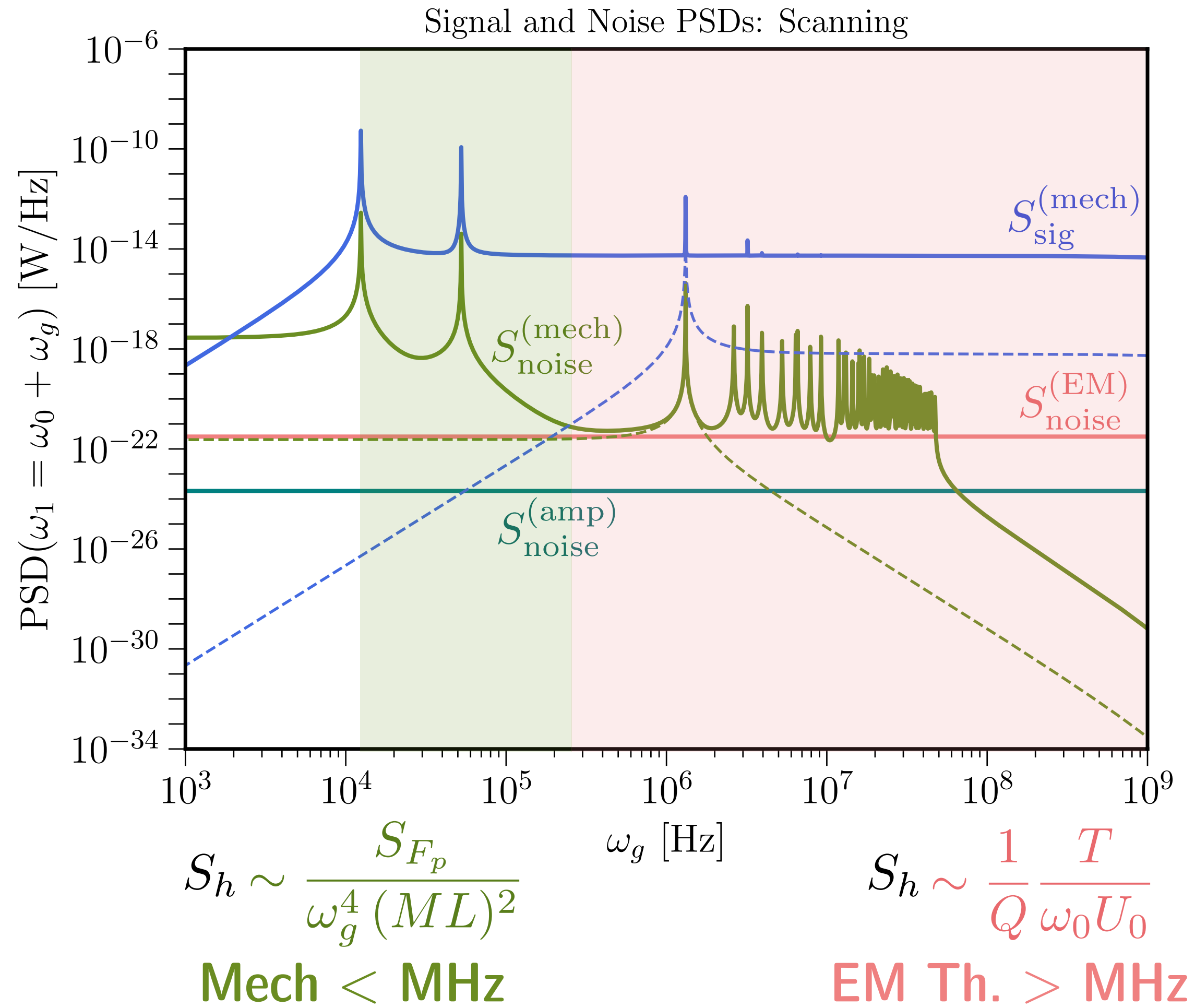
Noise in MAGO 2.0



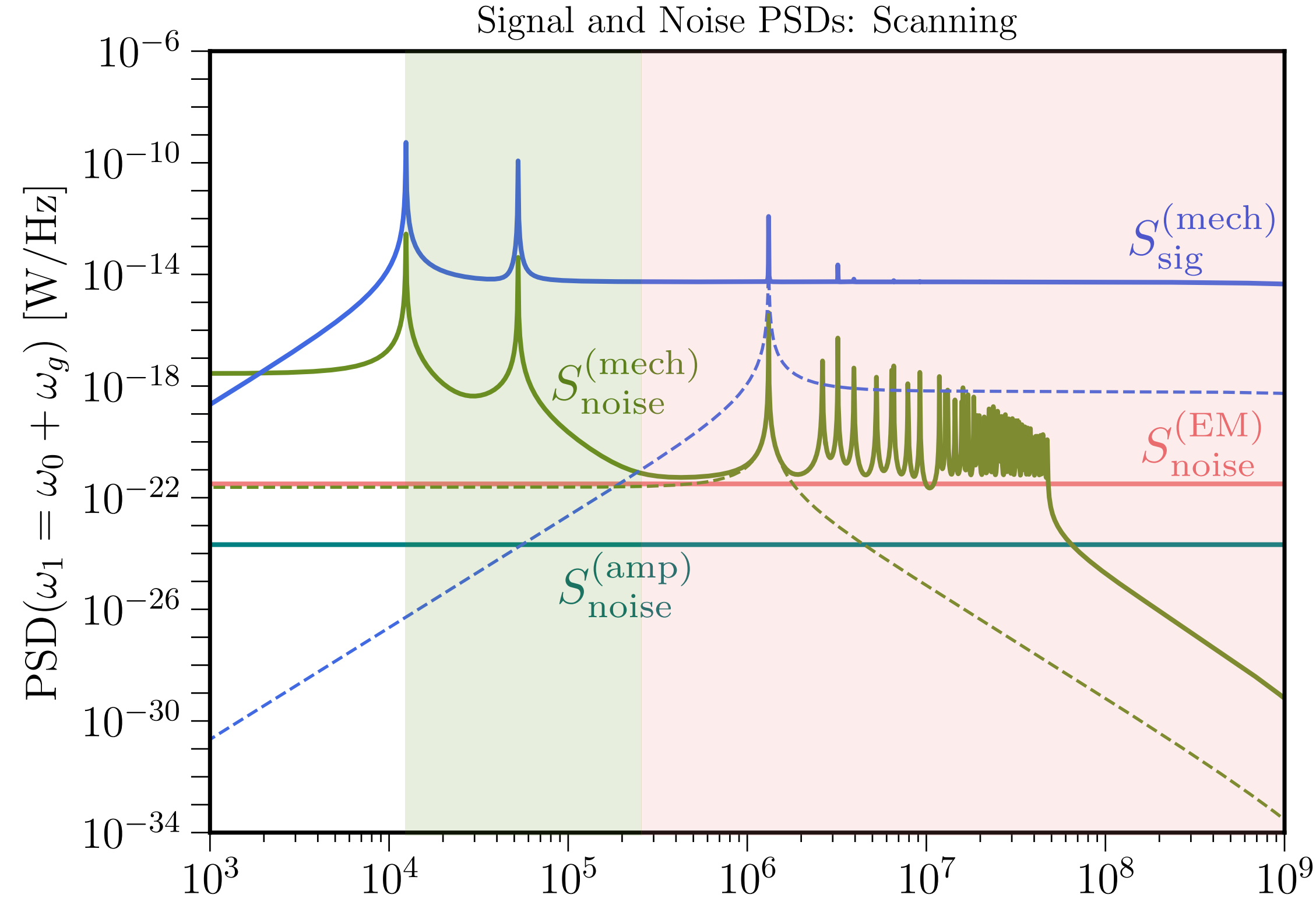
$$S_h \sim \frac{1}{Q} \frac{T}{\omega_0 U_0}$$

EM Th. > MHz

Noise in MAGO 2.0



Noise in MAGO 2.0

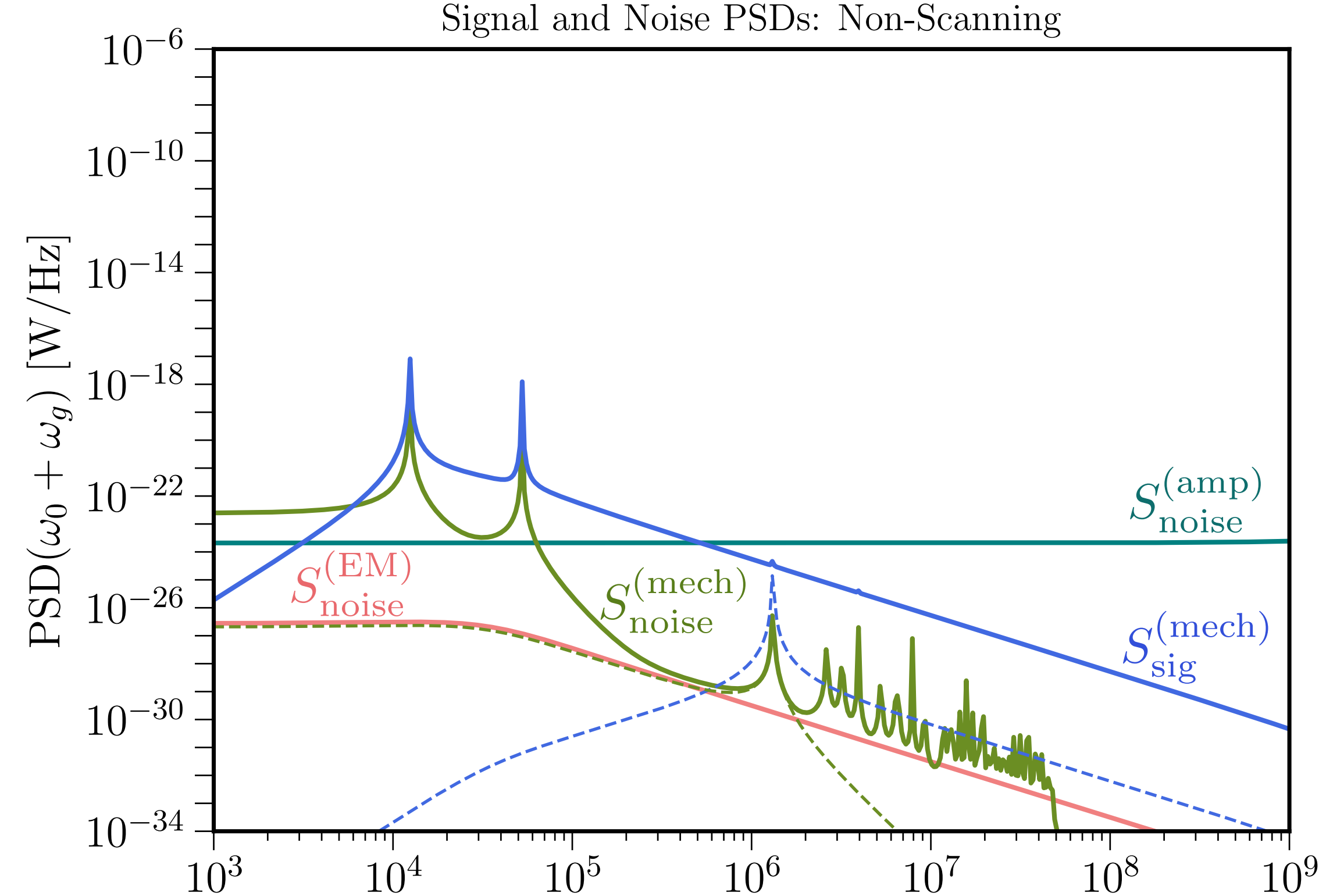


$$S_h \sim \frac{S_{F_p}}{\omega_g^4 (ML)^2}$$

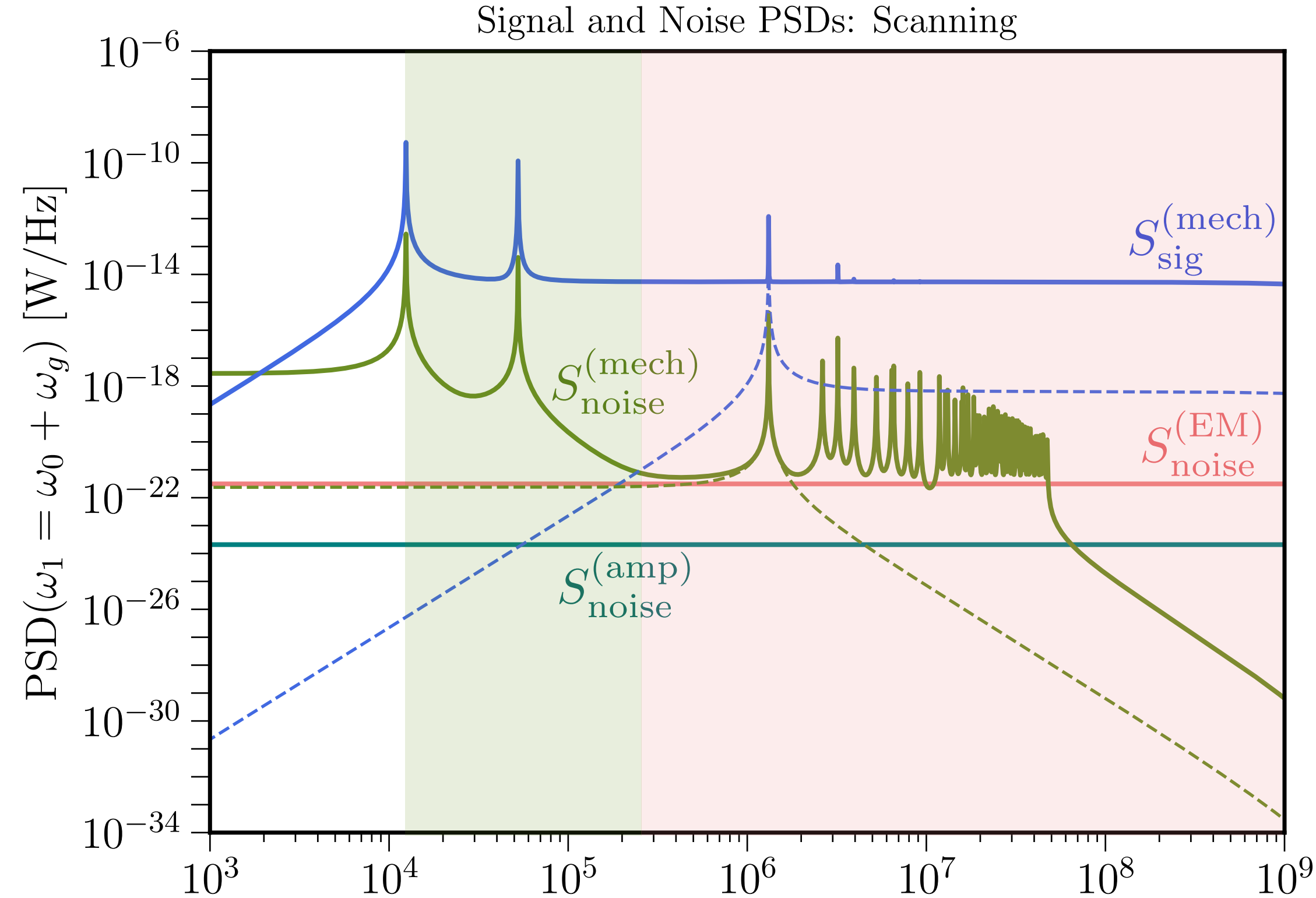
Mech < MHz

$$S_h \sim \frac{1}{Q} \frac{T}{\omega_0 U_0}$$

EM Th. > MHz



Noise in MAGO 2.0

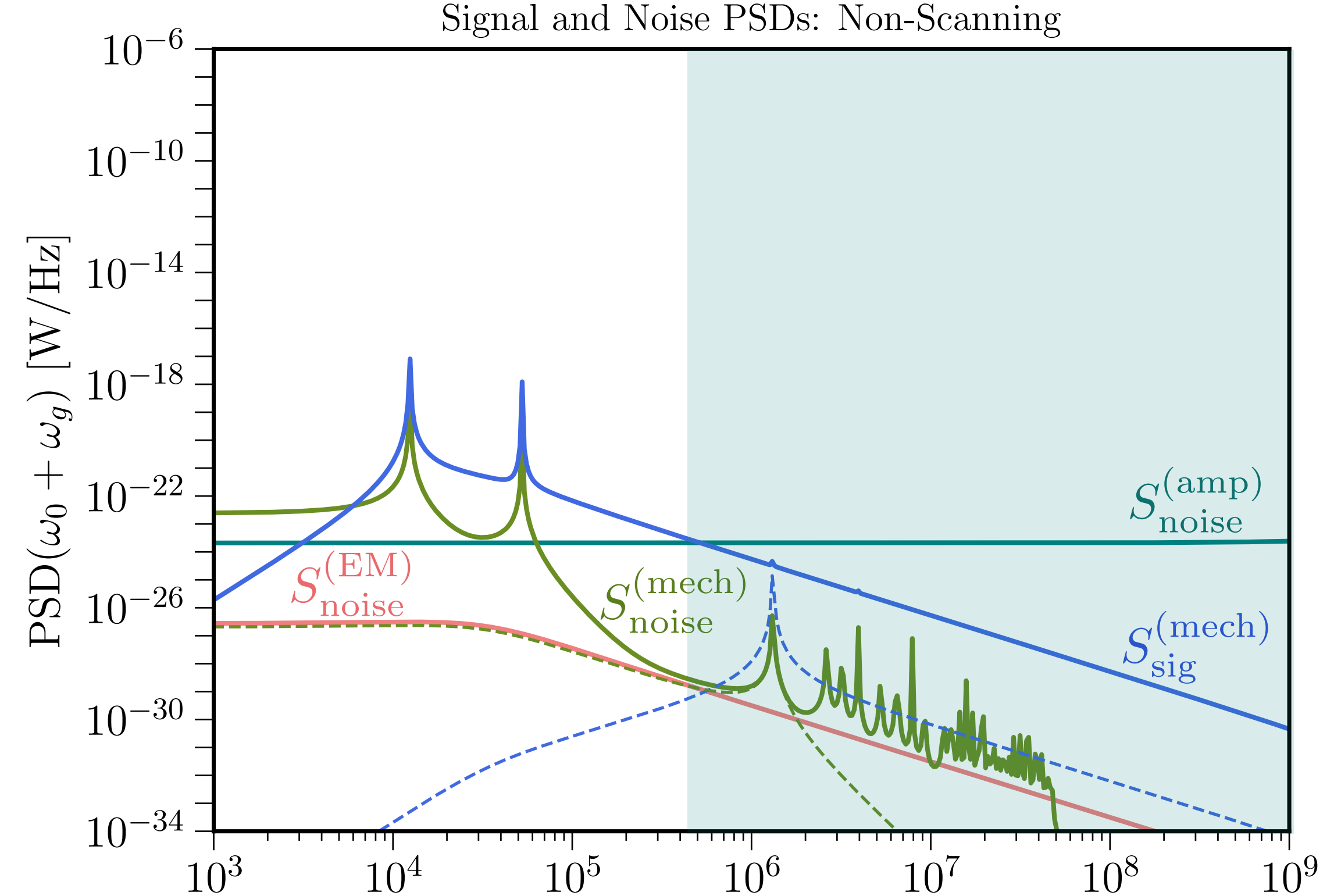


$$S_h \sim \frac{S_{F_p}}{\omega_g^4 (ML)^2}$$

Mech < MHz

$$S_h \sim \frac{1}{Q} \frac{T}{\omega_0 U_0}$$

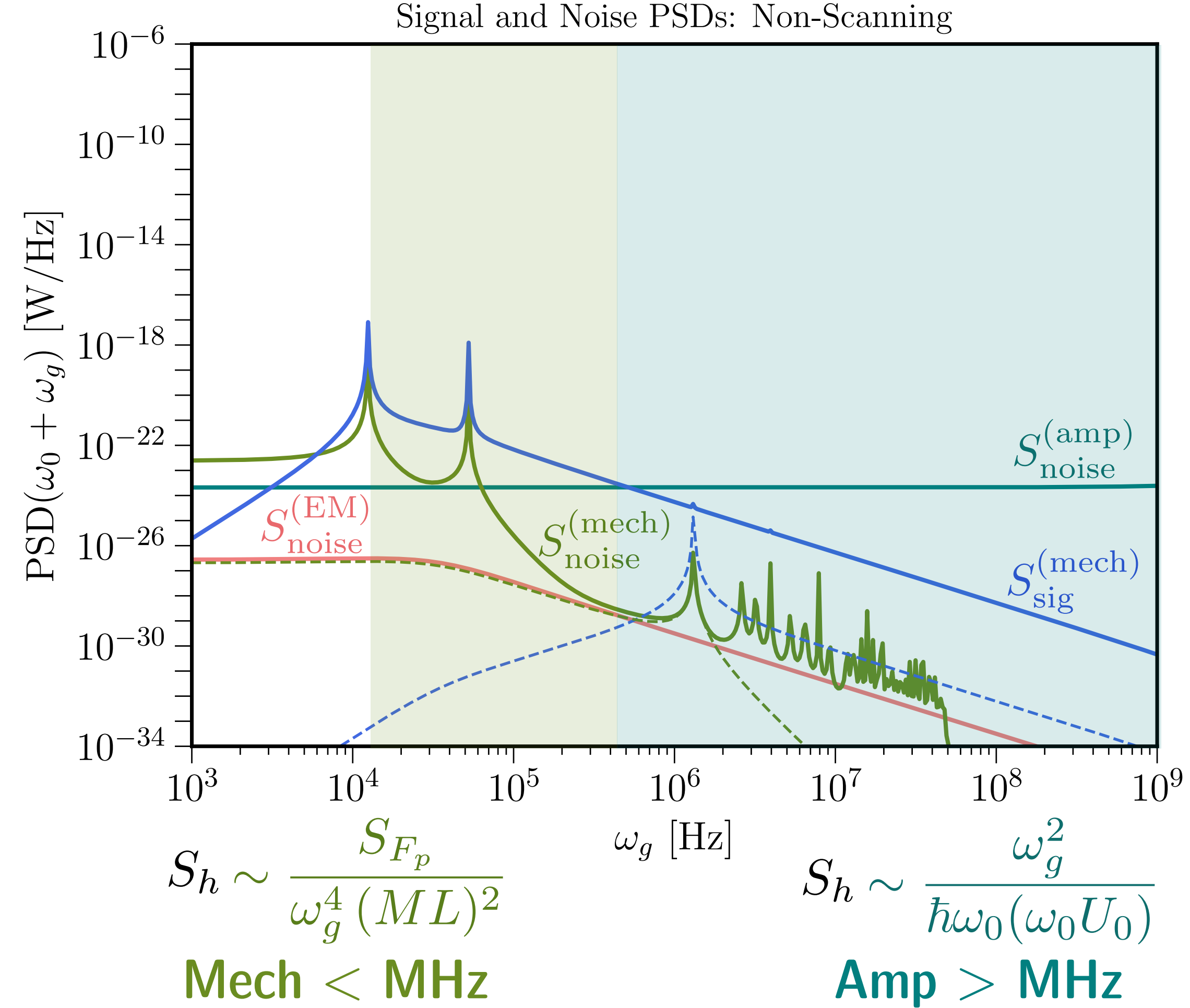
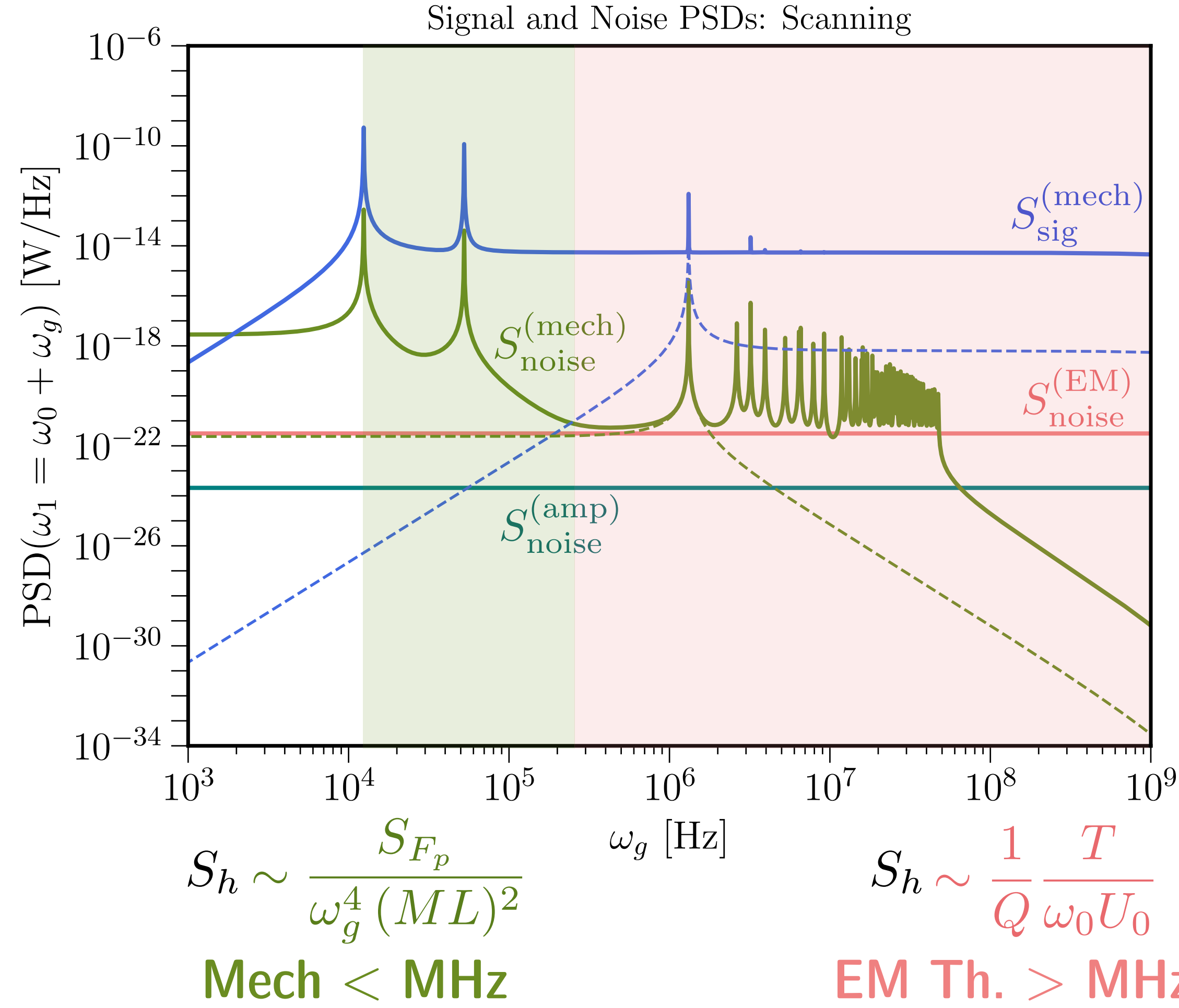
EM Th. > MHz



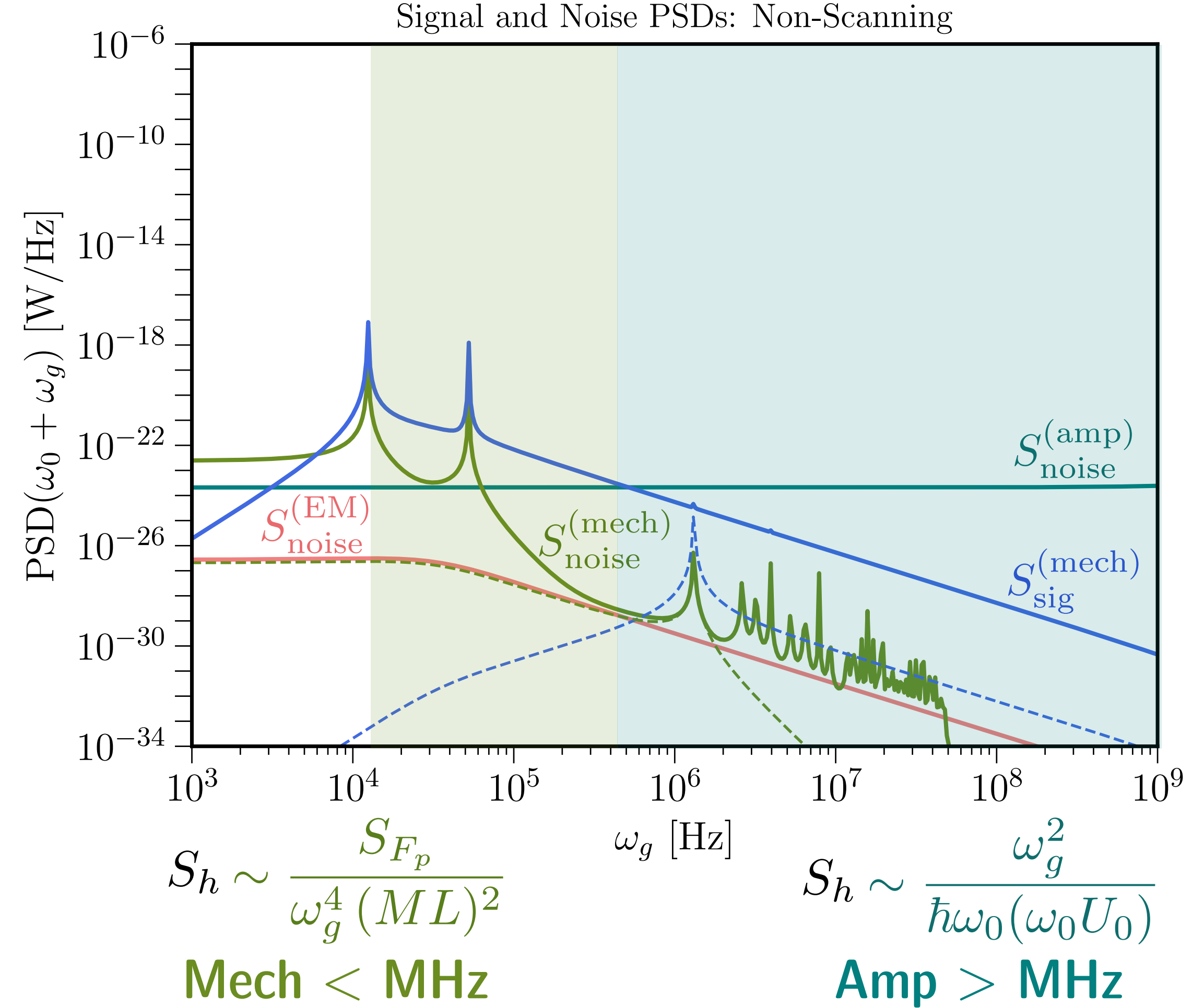
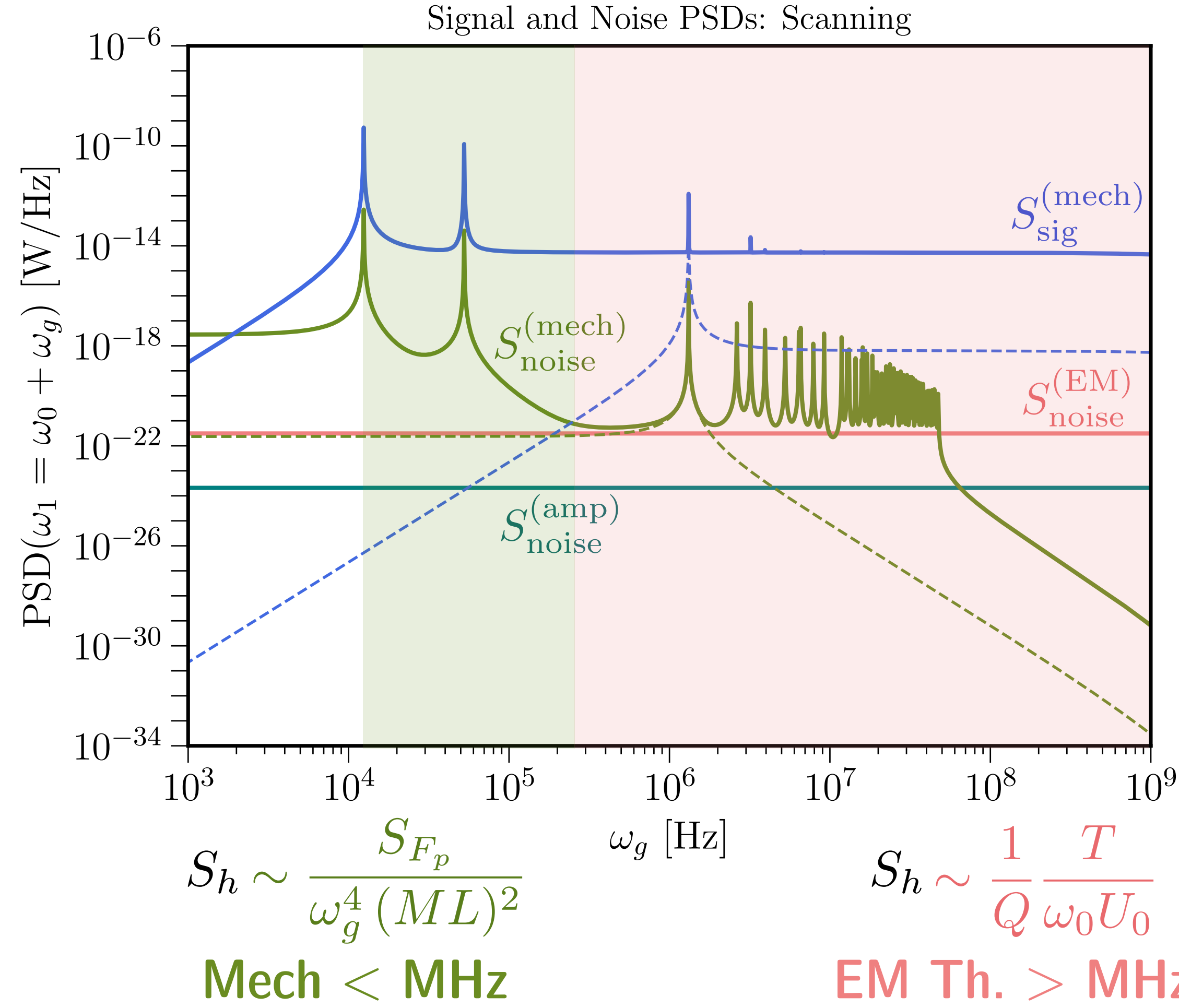
$$S_h \sim \frac{\omega_g^2}{\hbar \omega_0 (\omega_0 U_0)}$$

Amp > MHz

Noise in MAGO 2.0



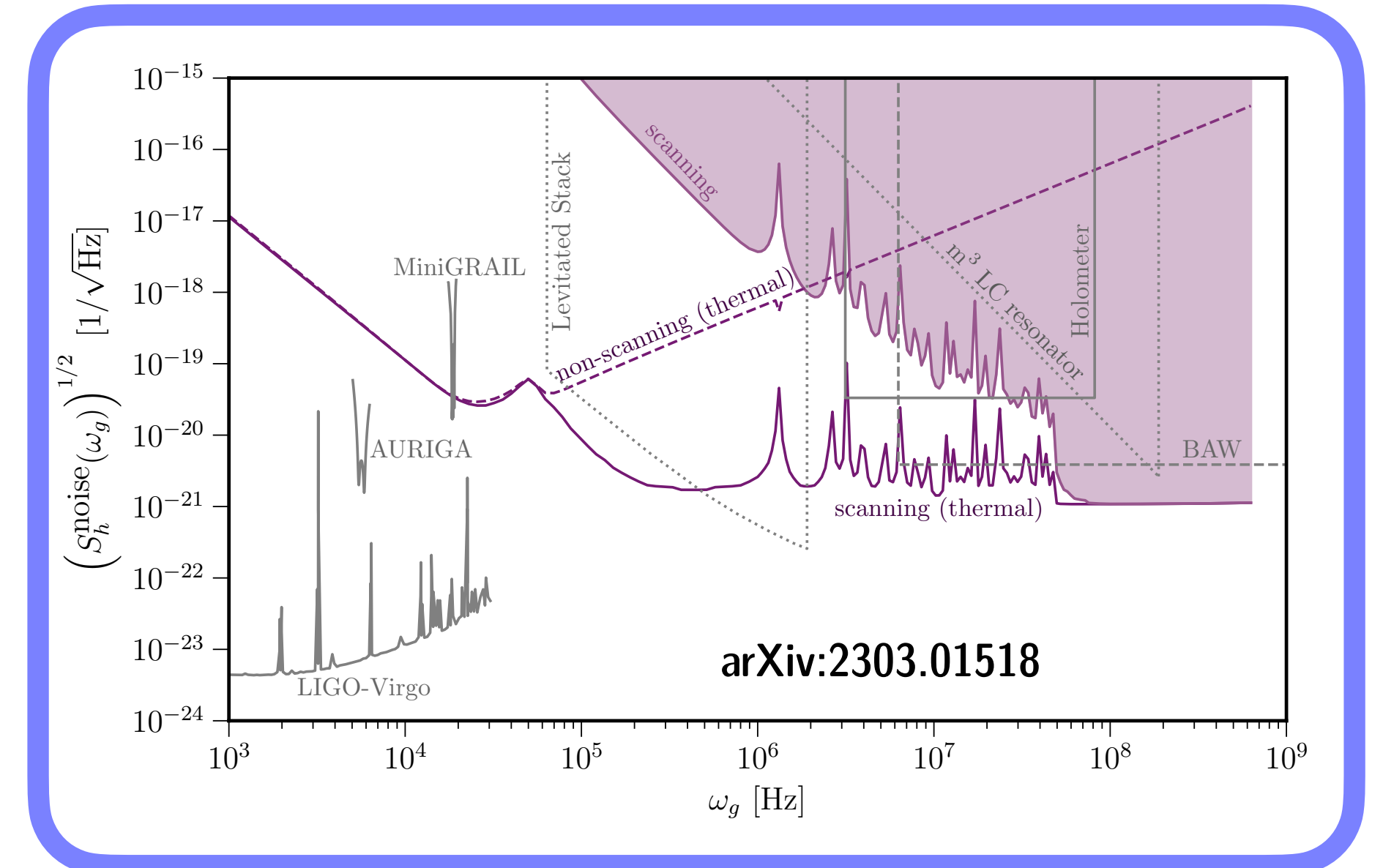
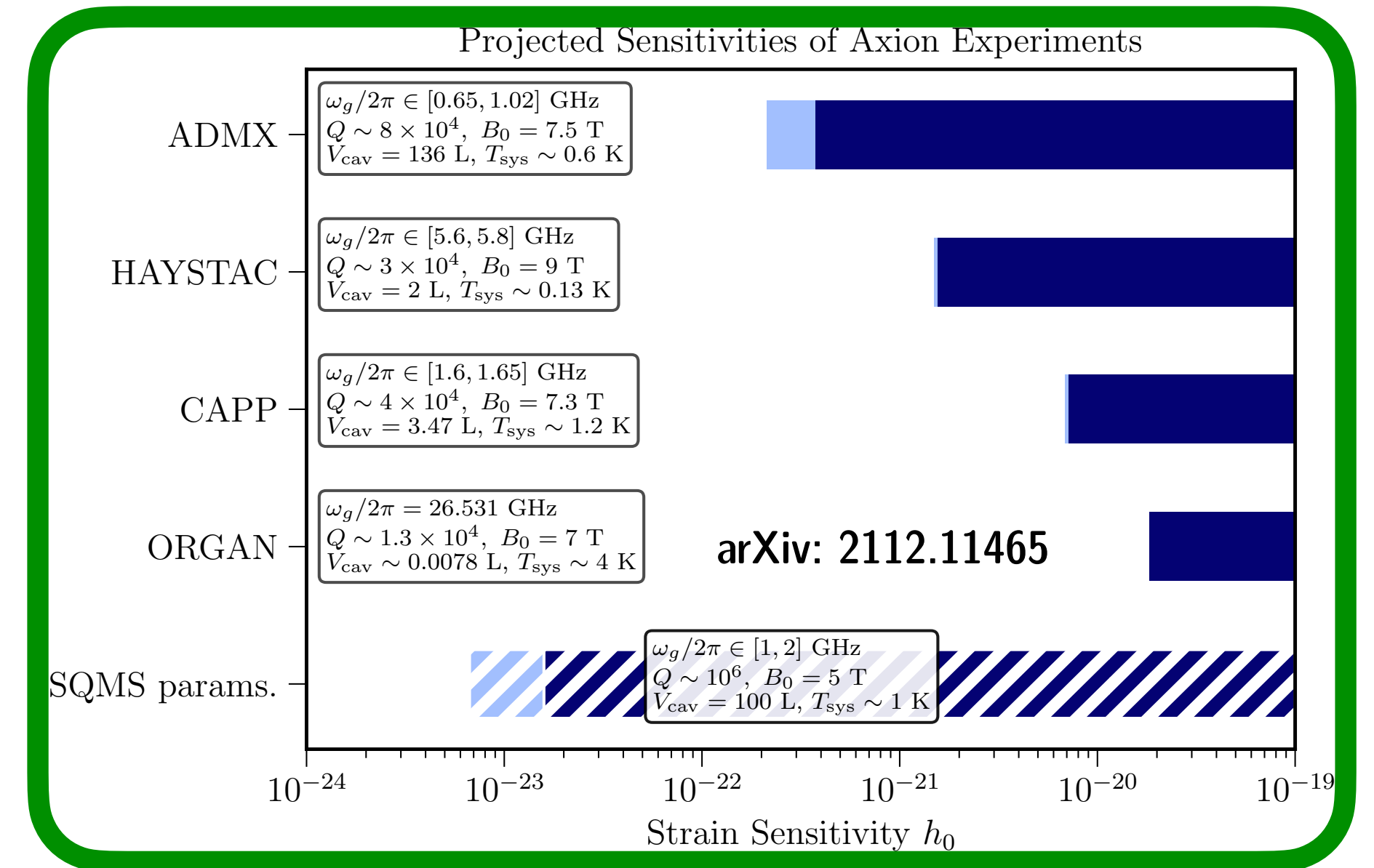
Noise in MAGO 2.0



NB: missing radiation damping effect studied in Löwenberg, Moortgat-Pick: 2307.14379

Open questions

Resonant cavities optimal in the radio band?

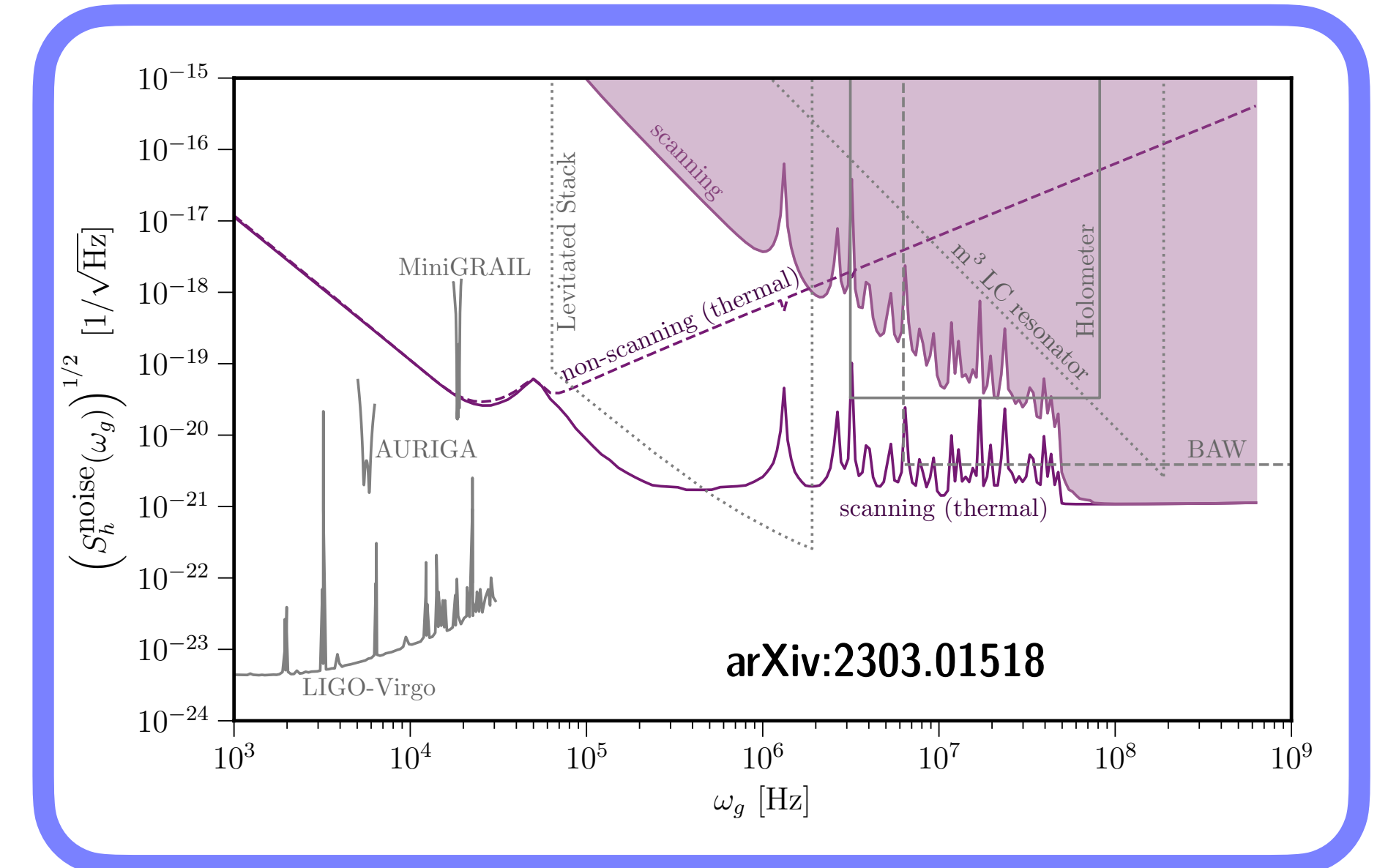
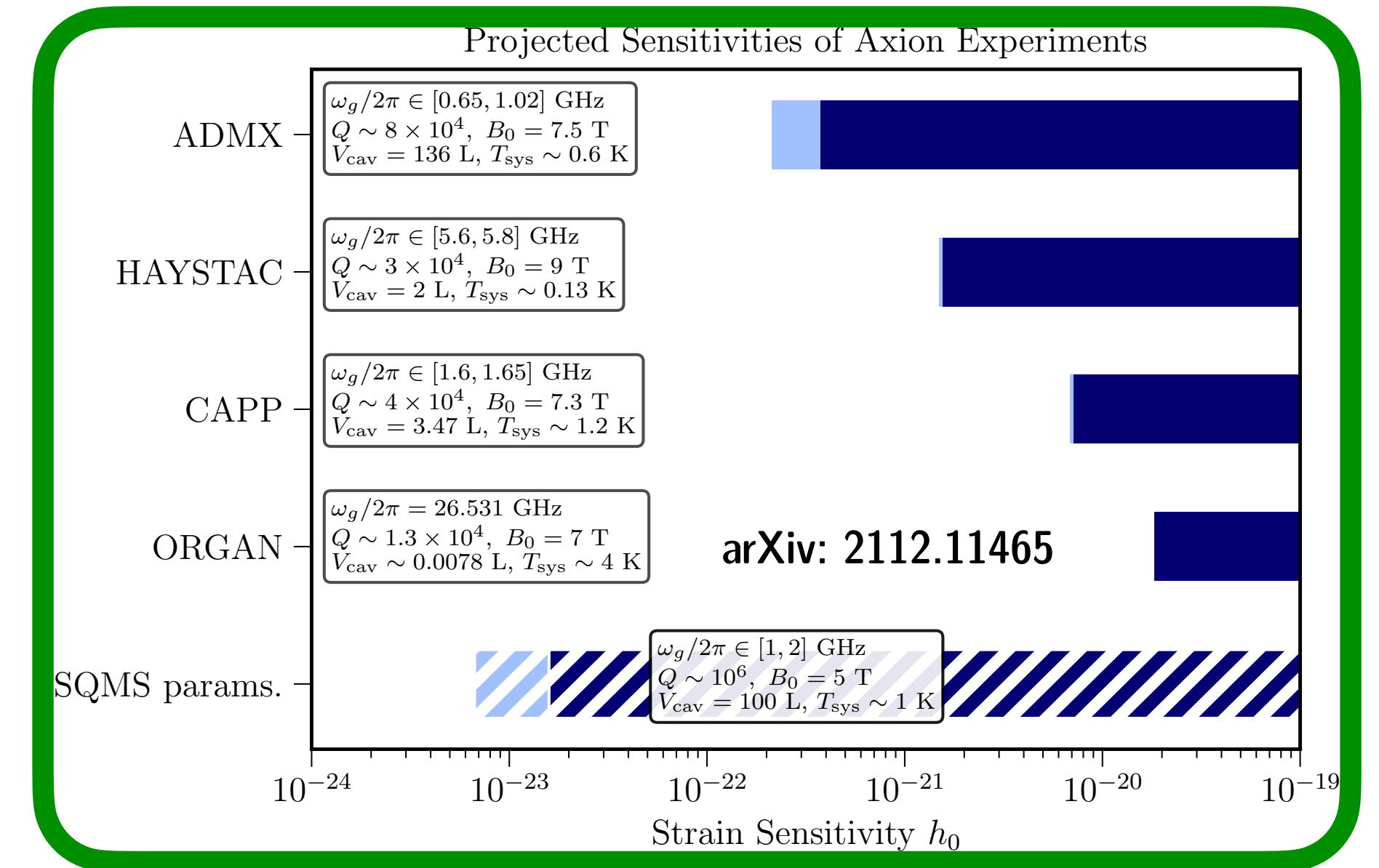


Open questions

Resonant cavities optimal in the radio band?

Advances in readout?

synergies w/ Axion searches, QC(?)



Open questions

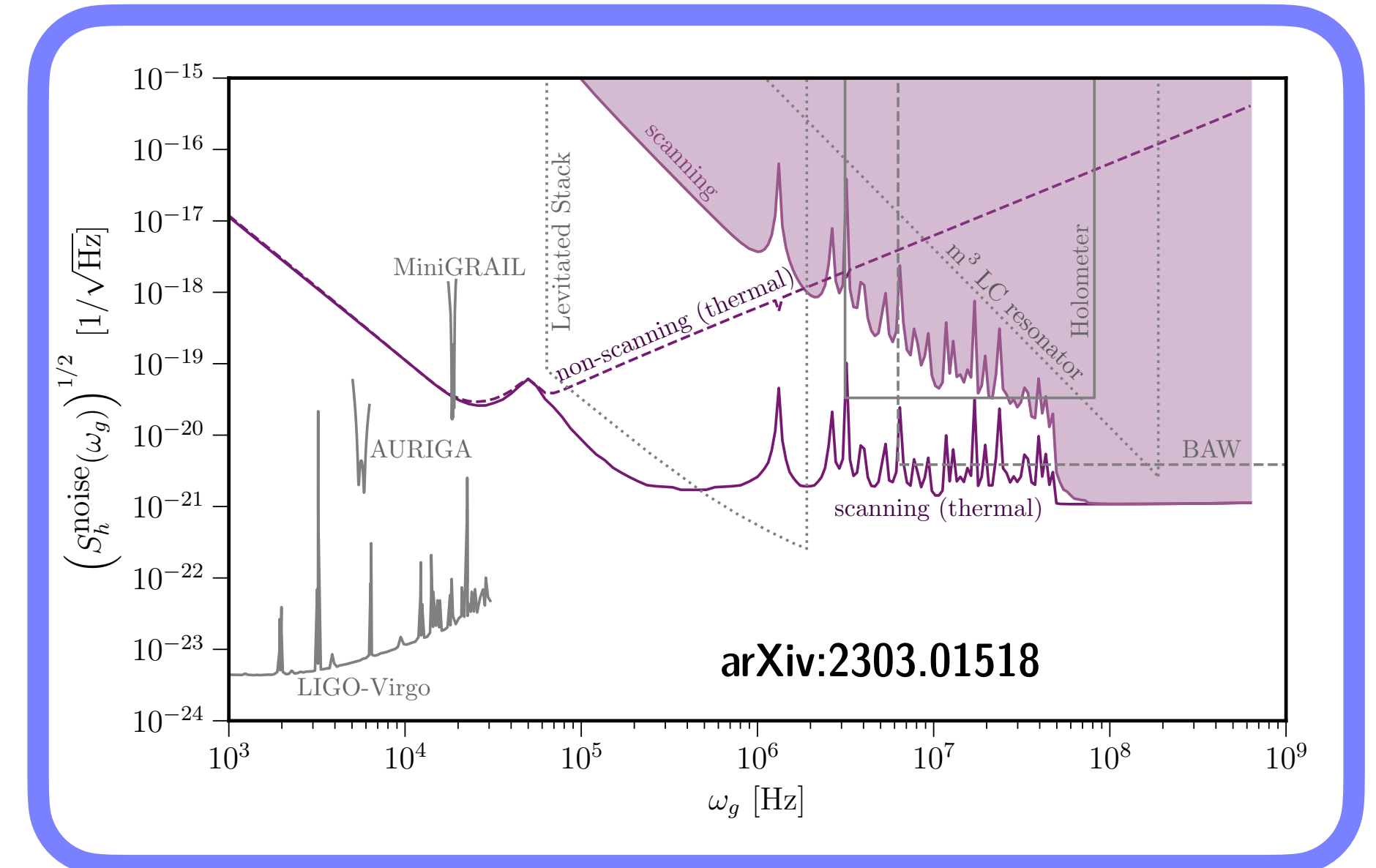
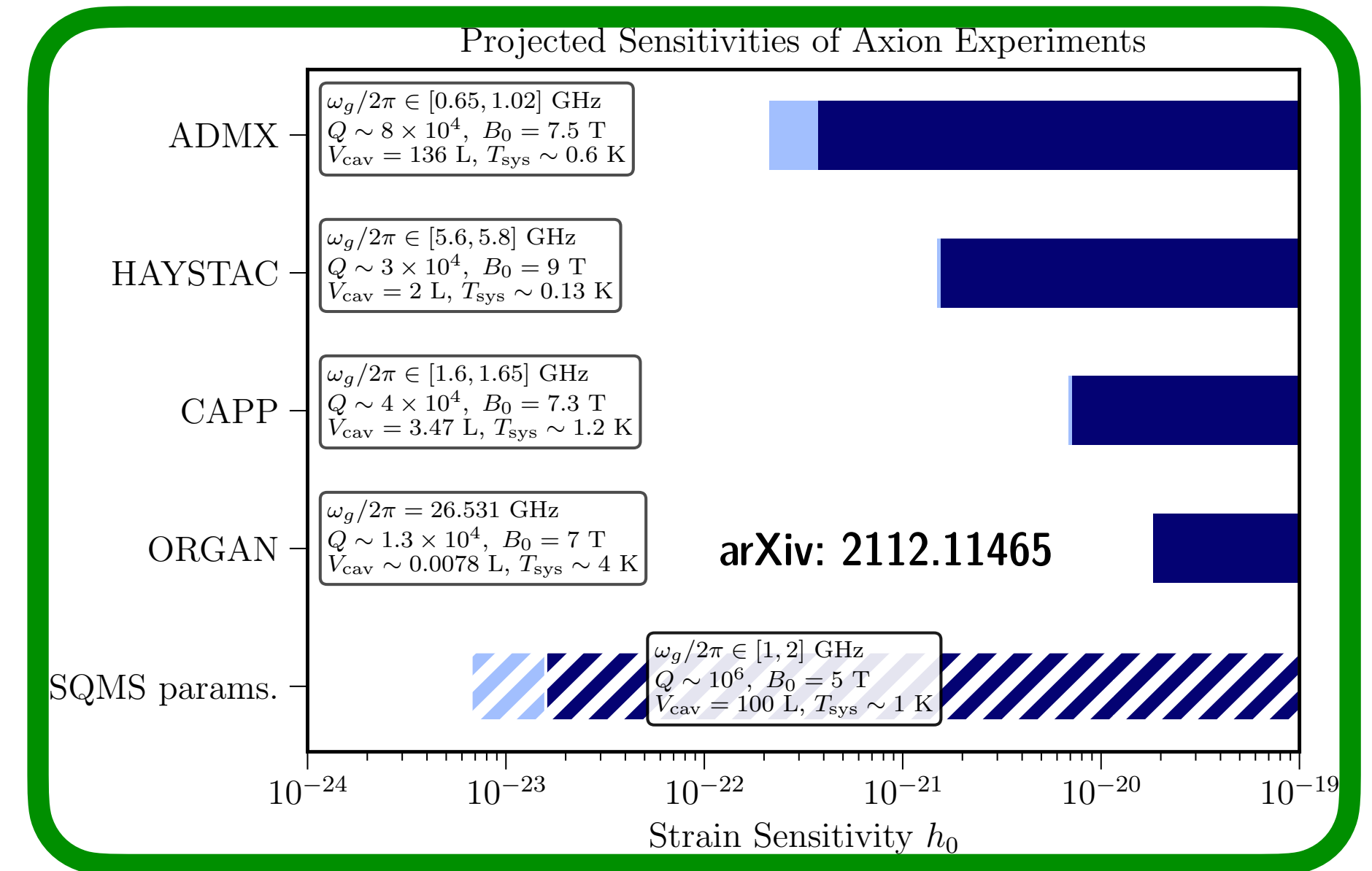
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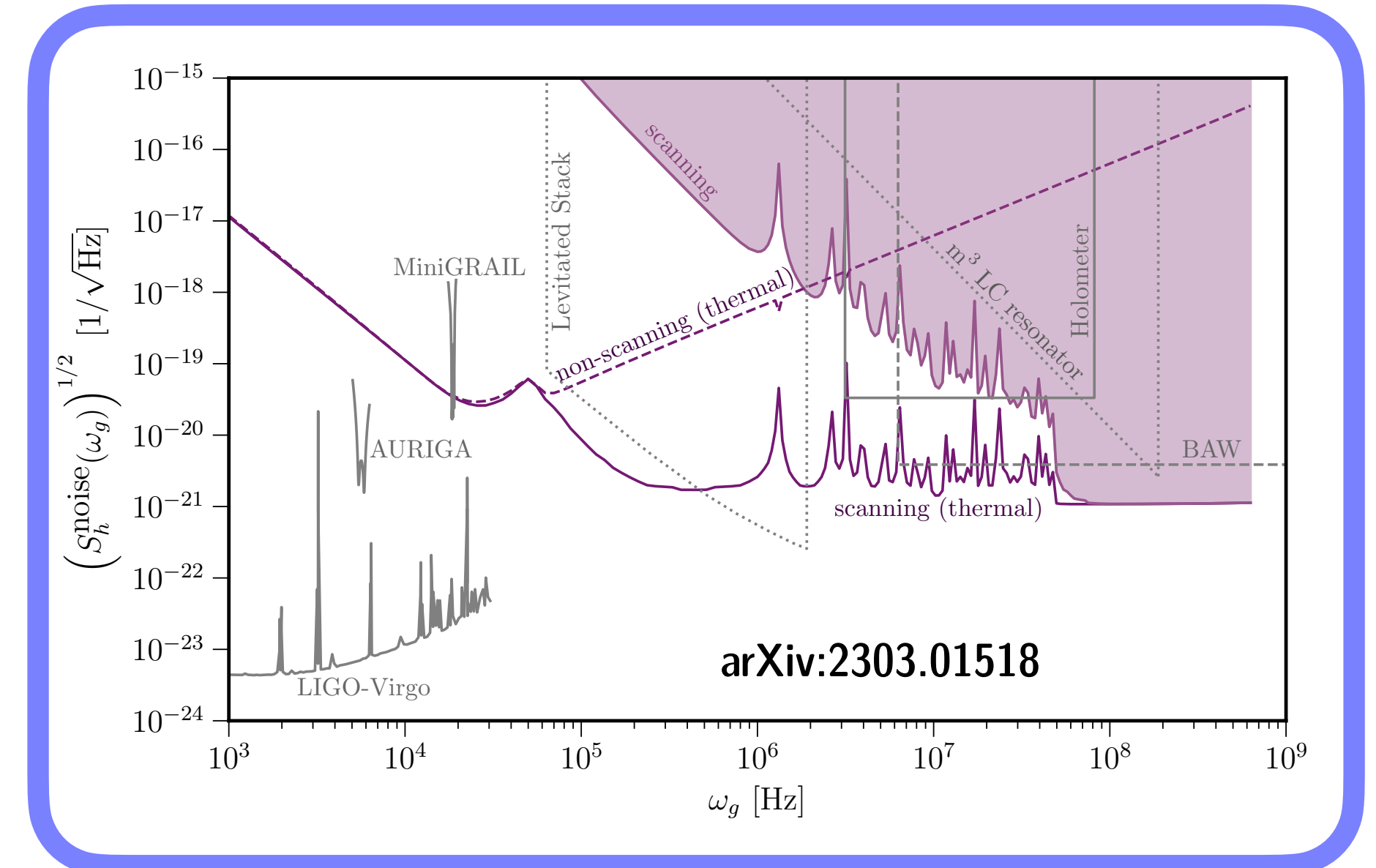
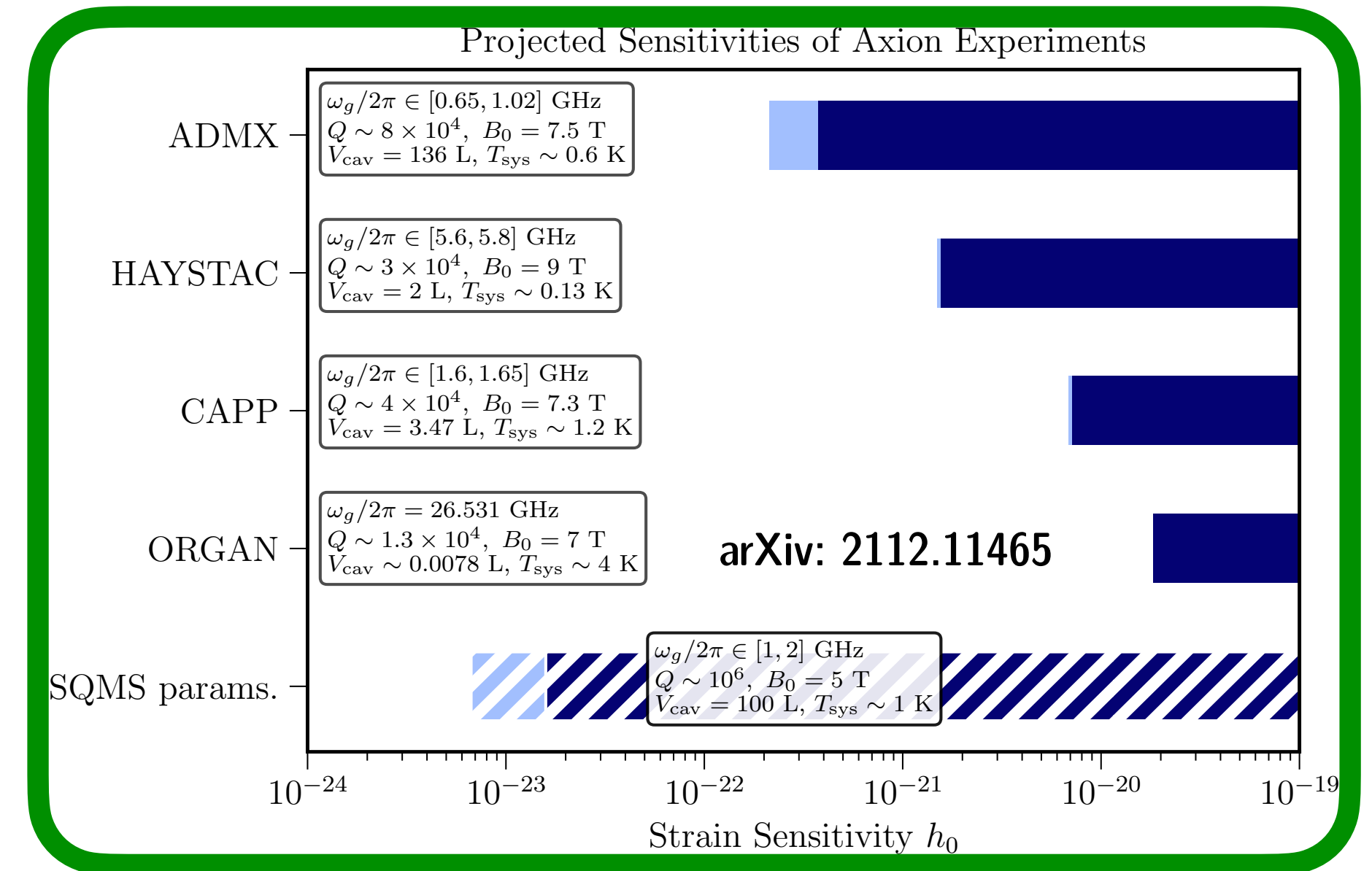
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Optimisation of MAGO-style setup?

FNAL + DESY



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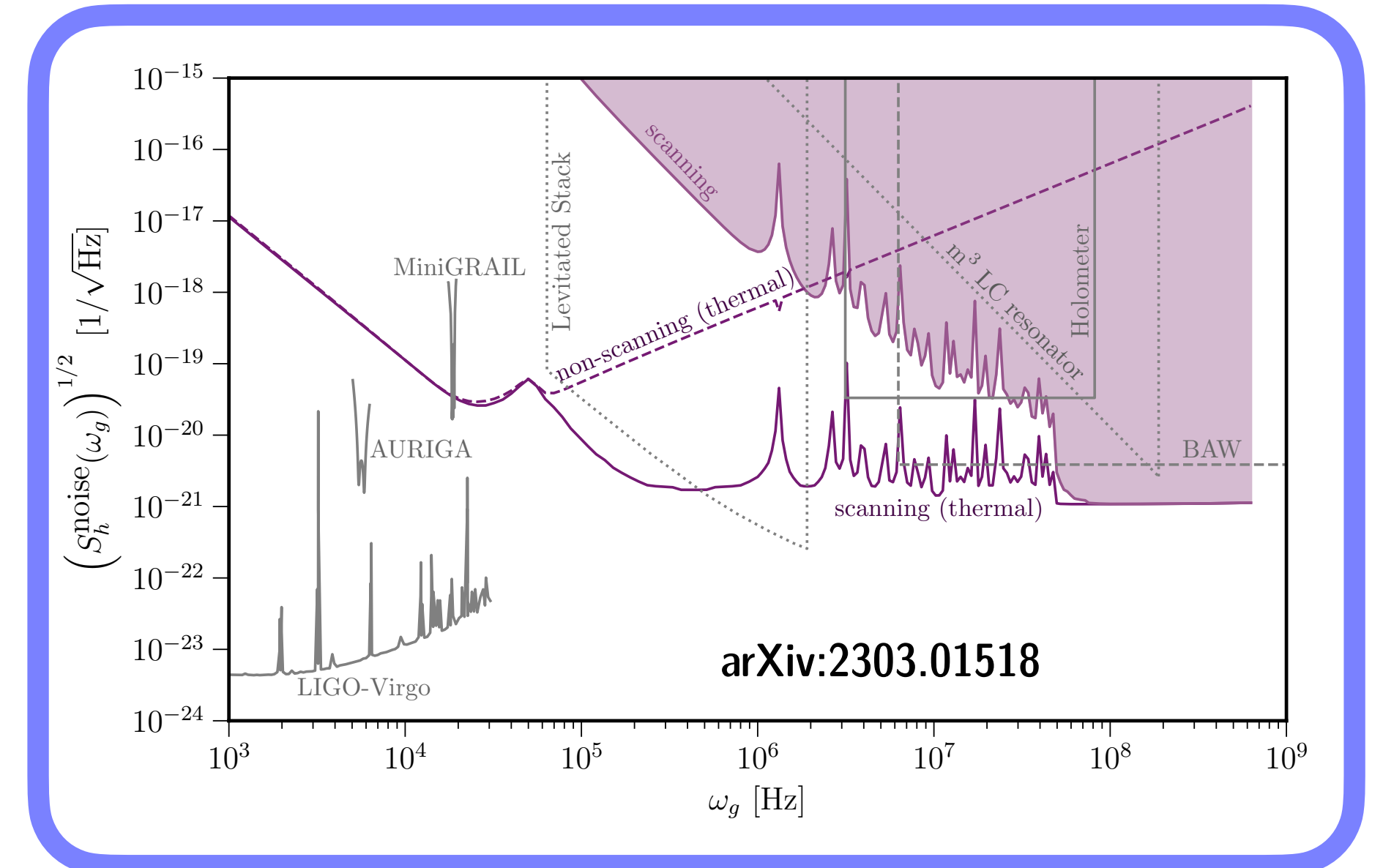
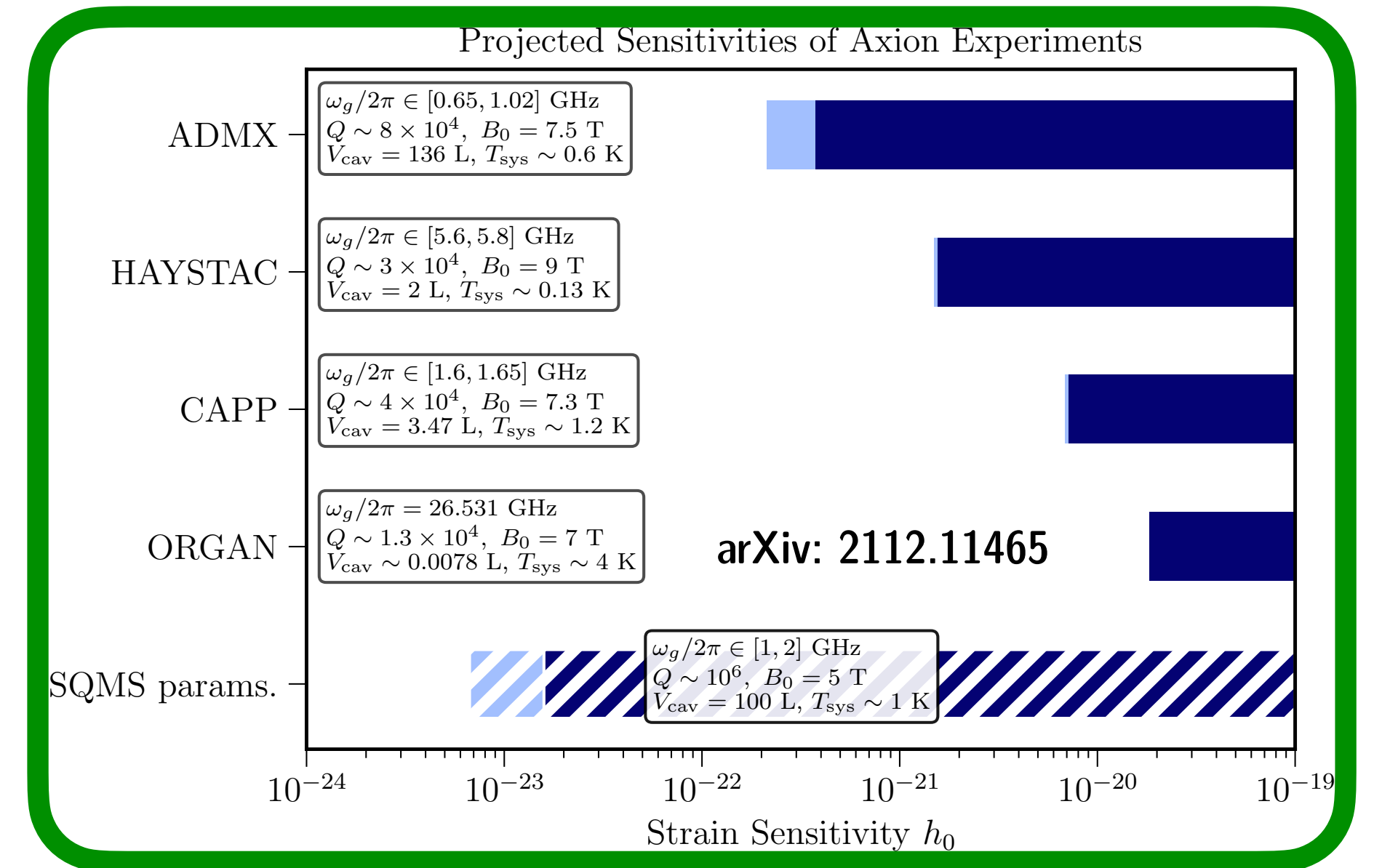
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Signals above kHz?



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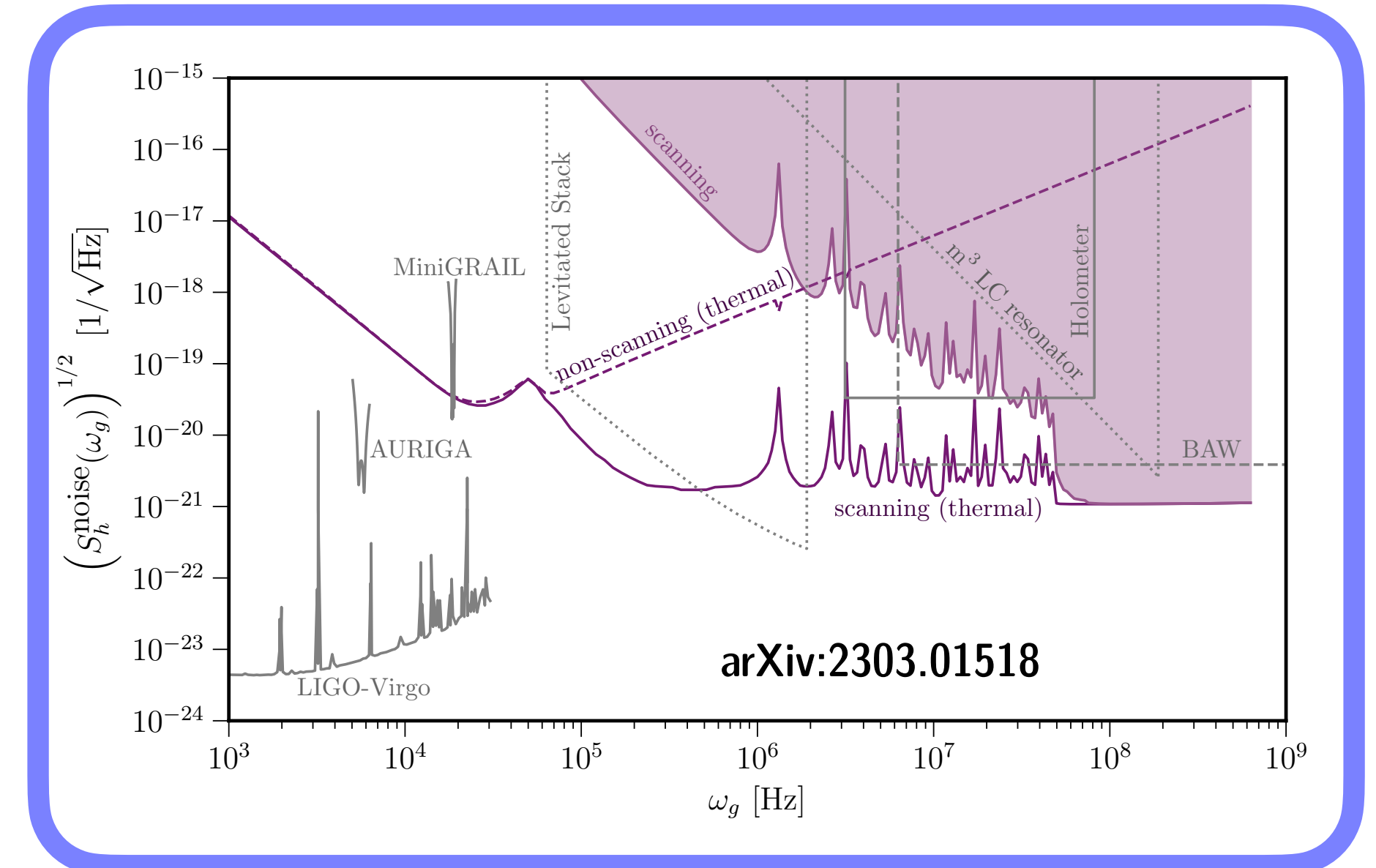
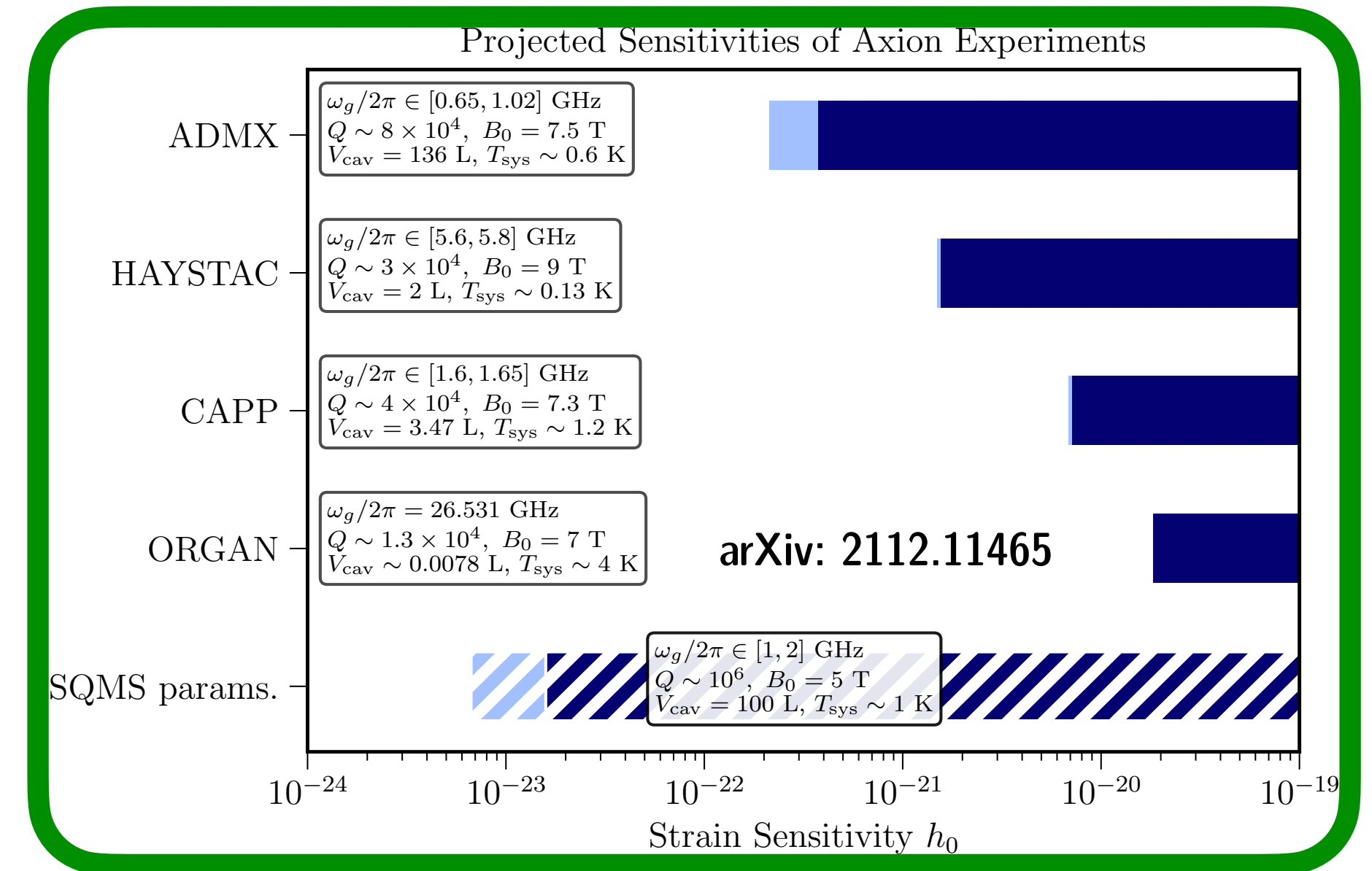
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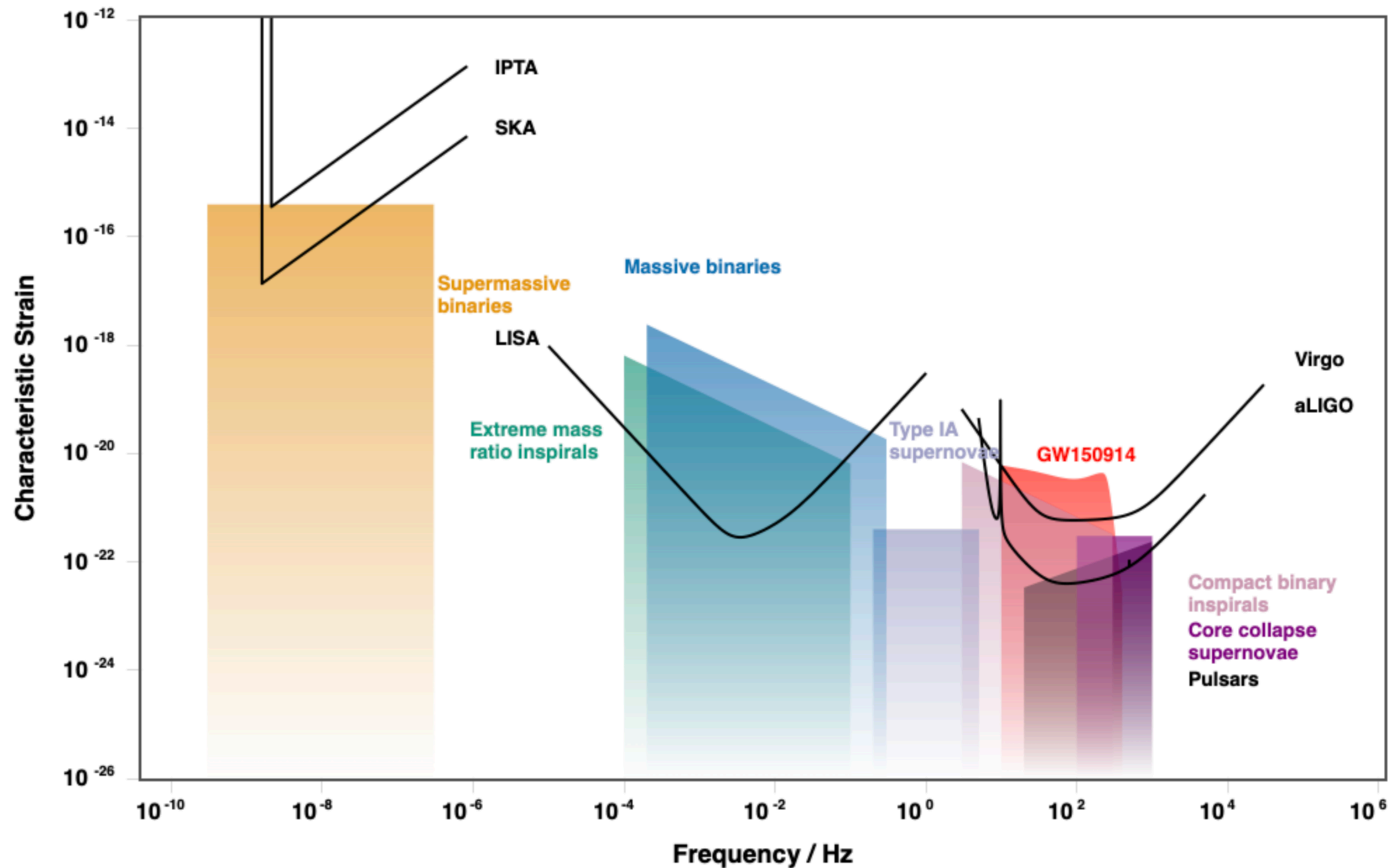
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BACKUP

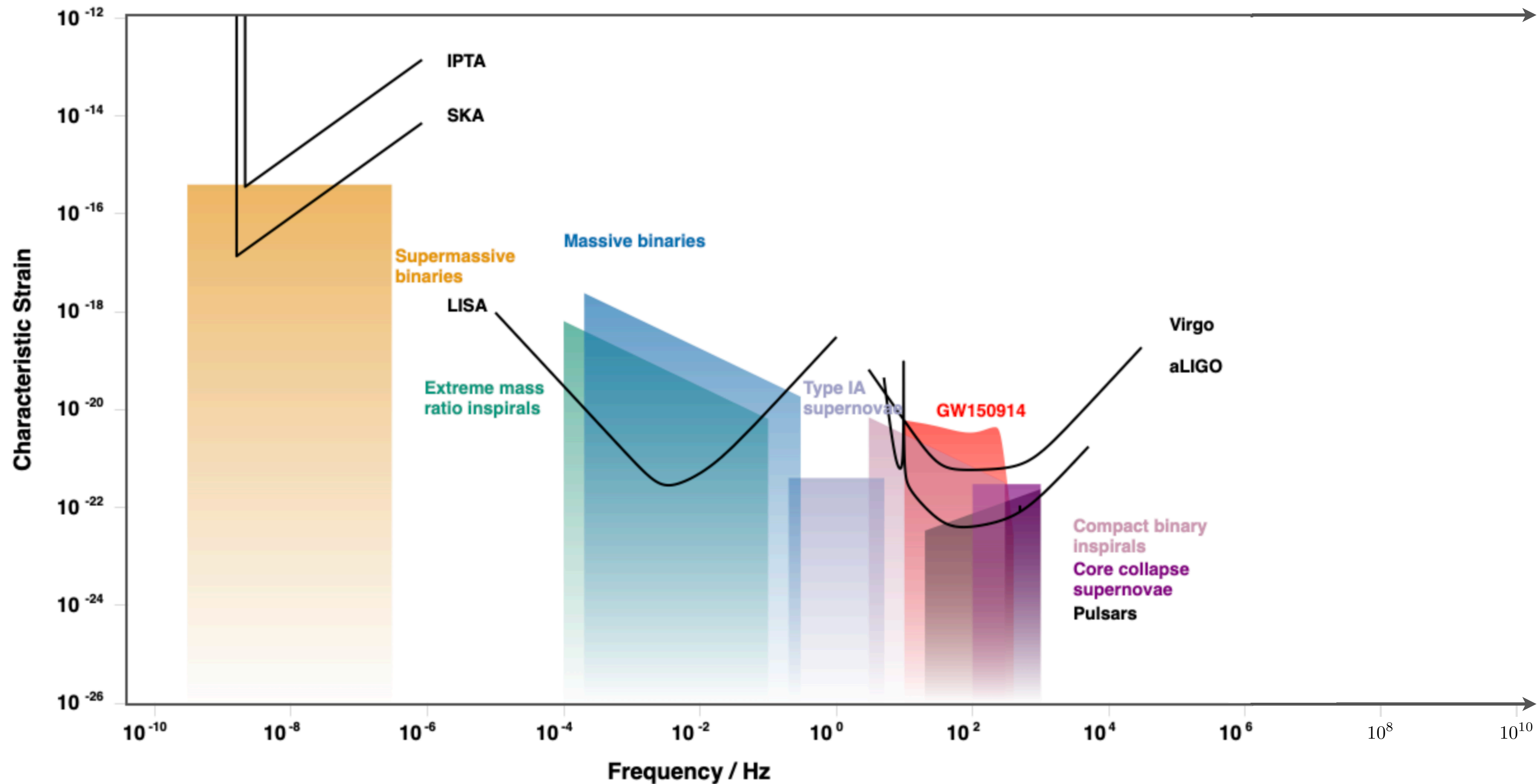
Gravitational Waves

By Christopher Moore, Robert Cole and Christopher Berry, formerly of the Gravitational Wave Group at the Institute of Astronomy, University of Cambridge



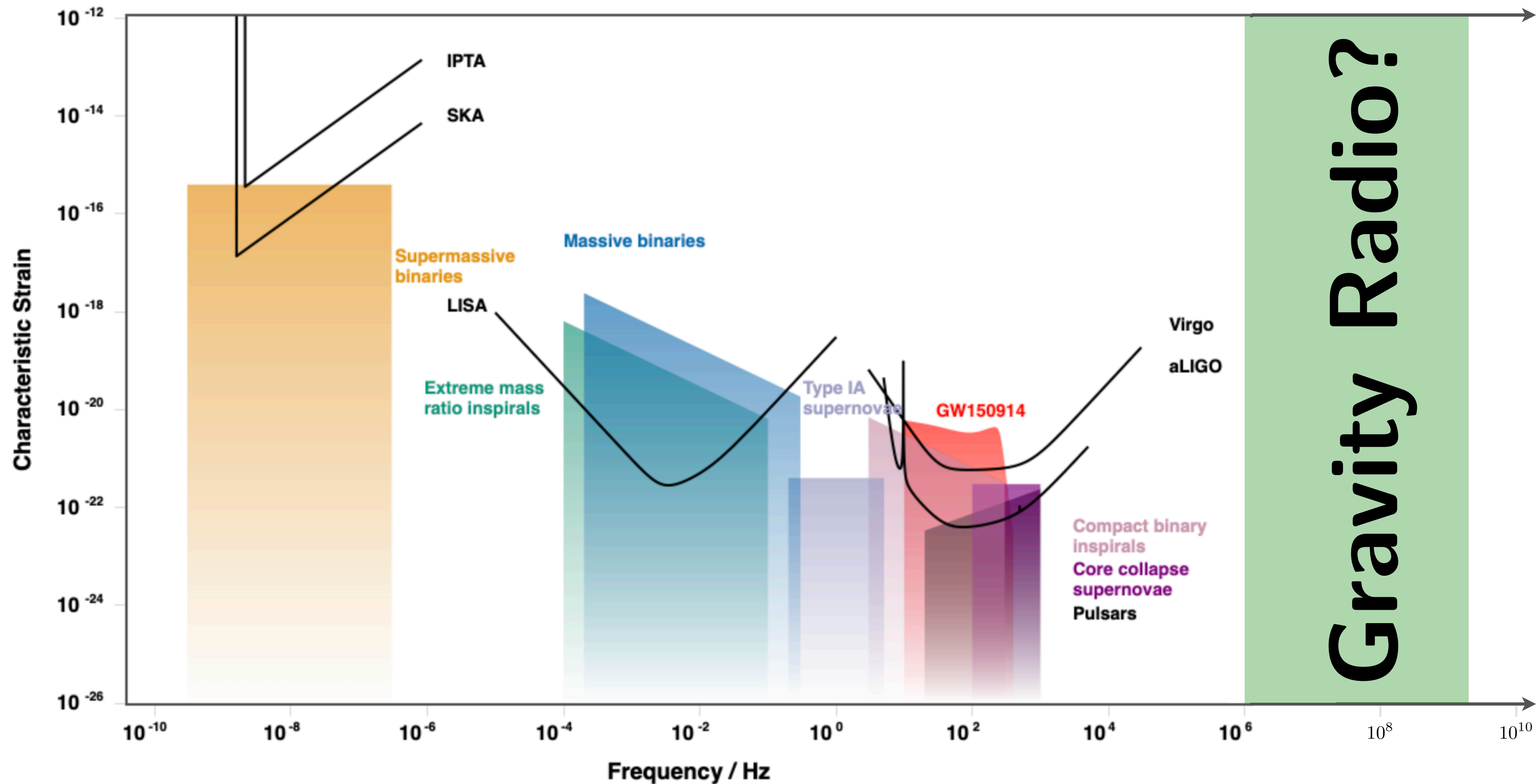
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Sources of HFGWs

Sources discussed in detail Weds-Fri

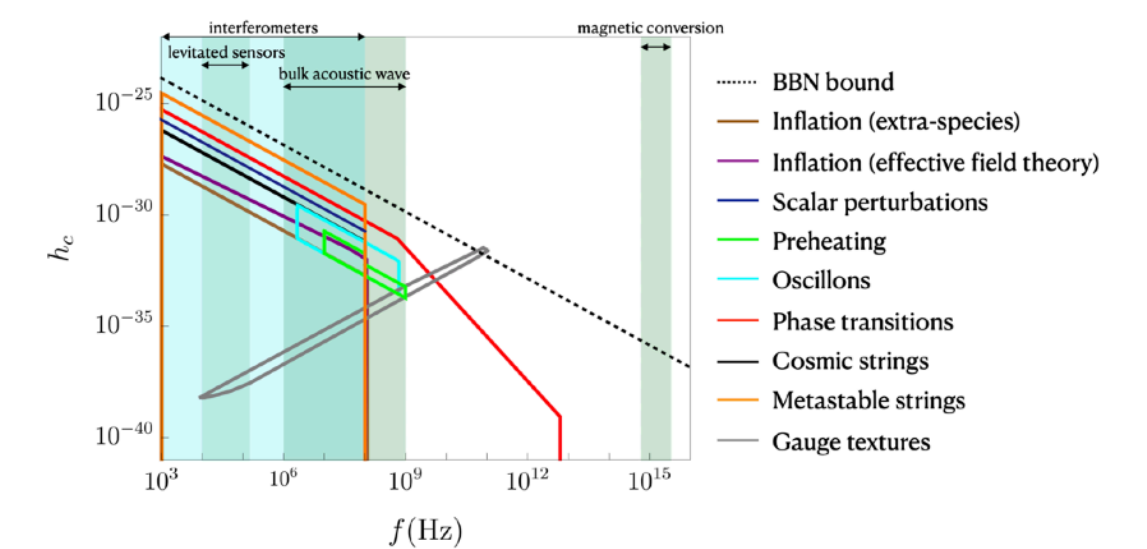
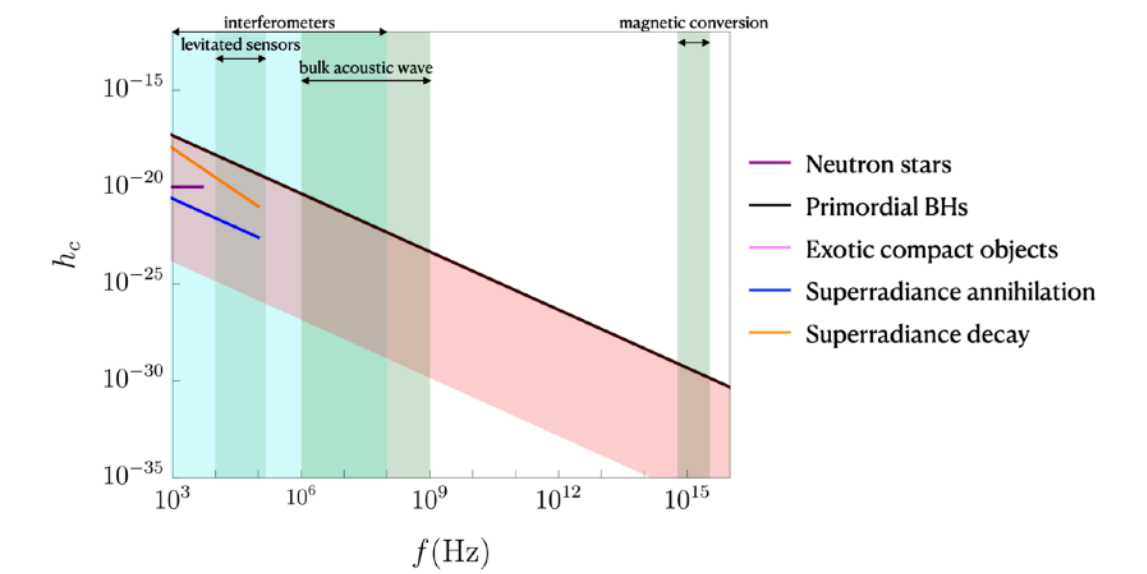
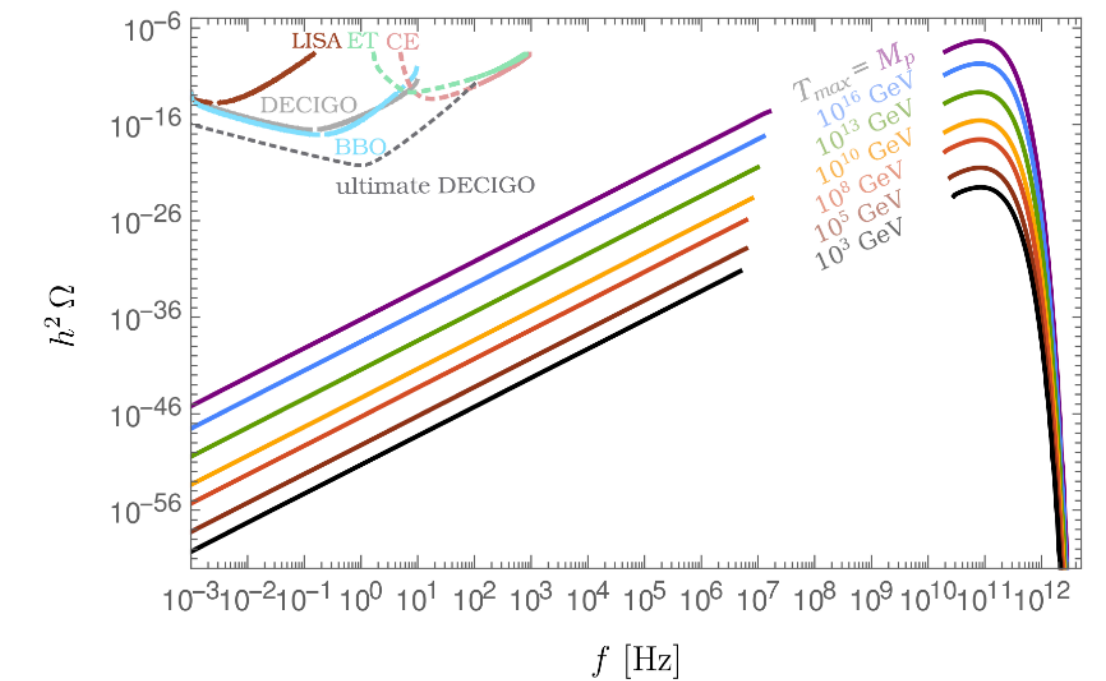
Stochastic

Coherent

Standard Model:

BSM:

Ringwald et al, 2011.04731



Aggarwal et al, 2011.12414

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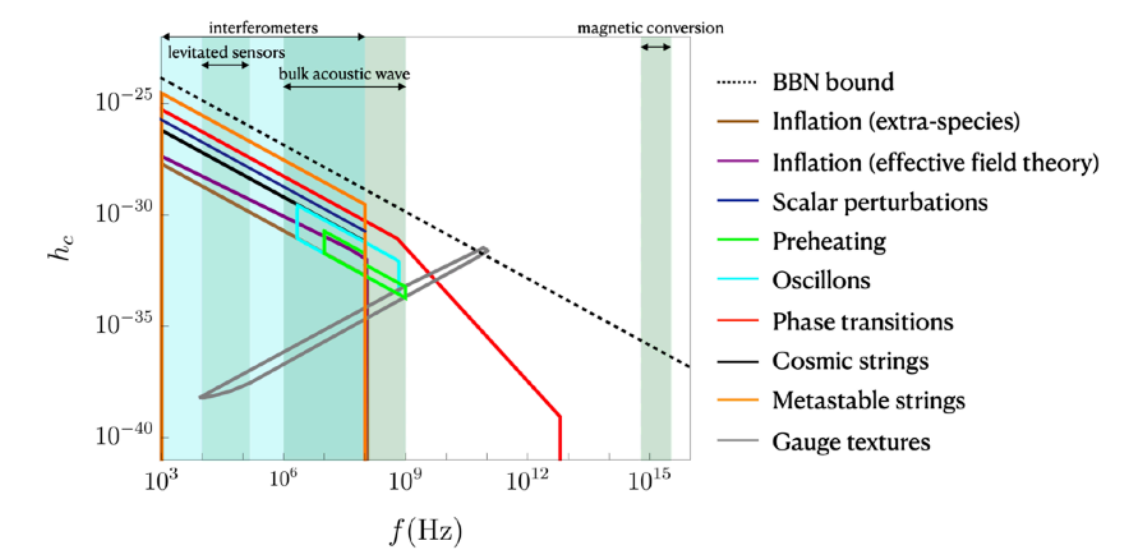
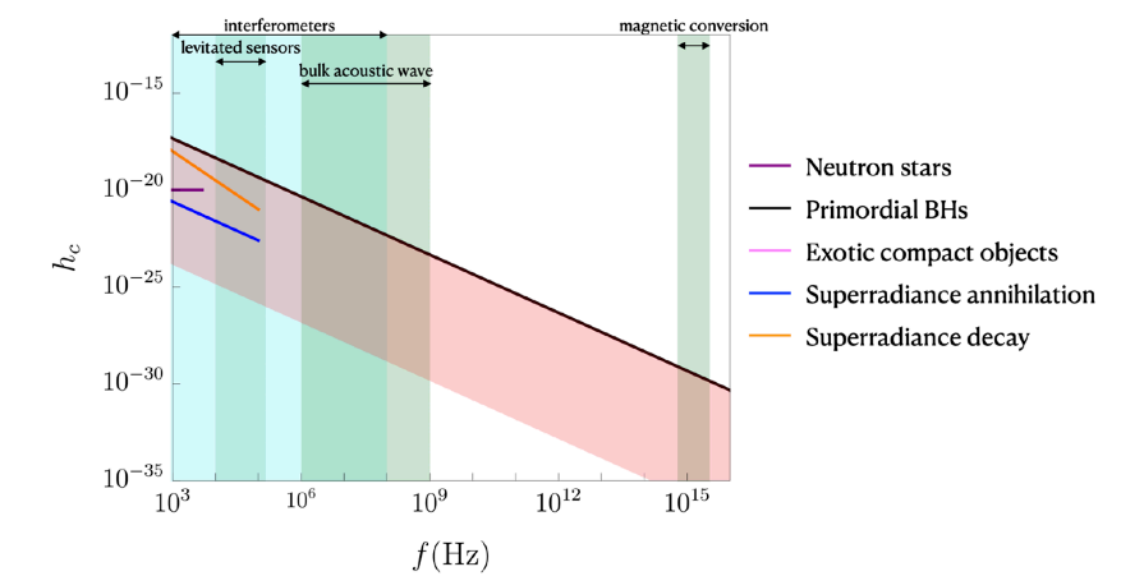
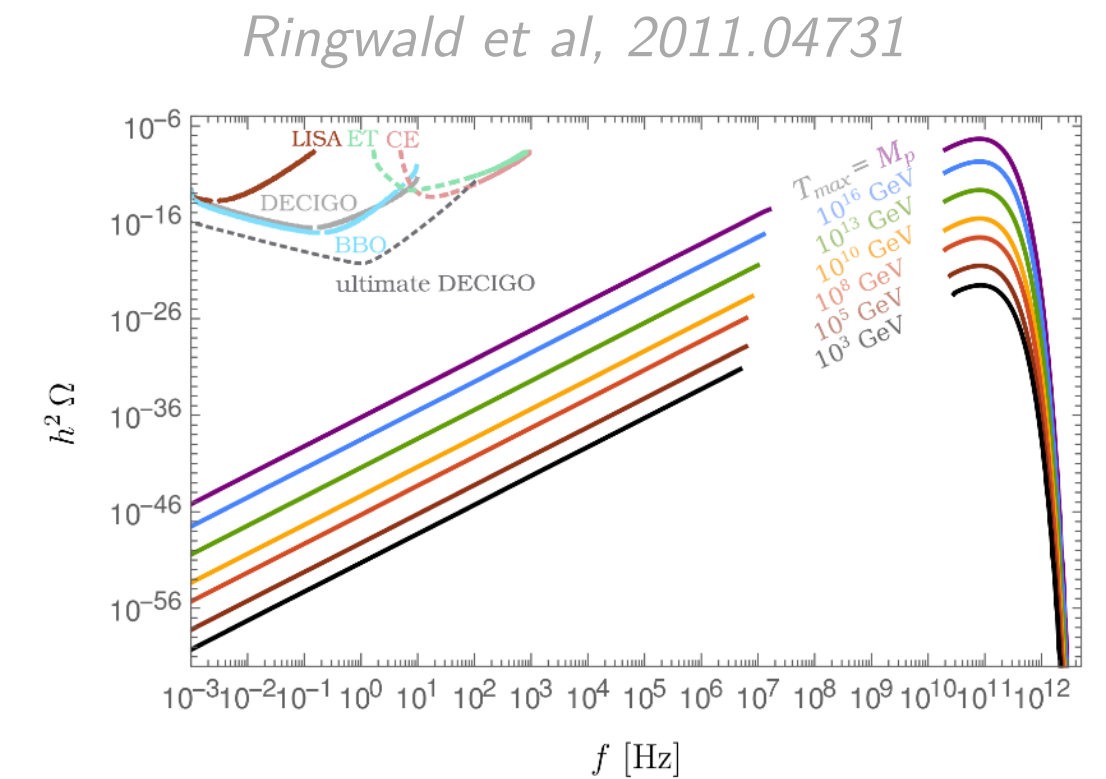
Thermal plasma fluctuations

Ghiglieri & Laine (2015)

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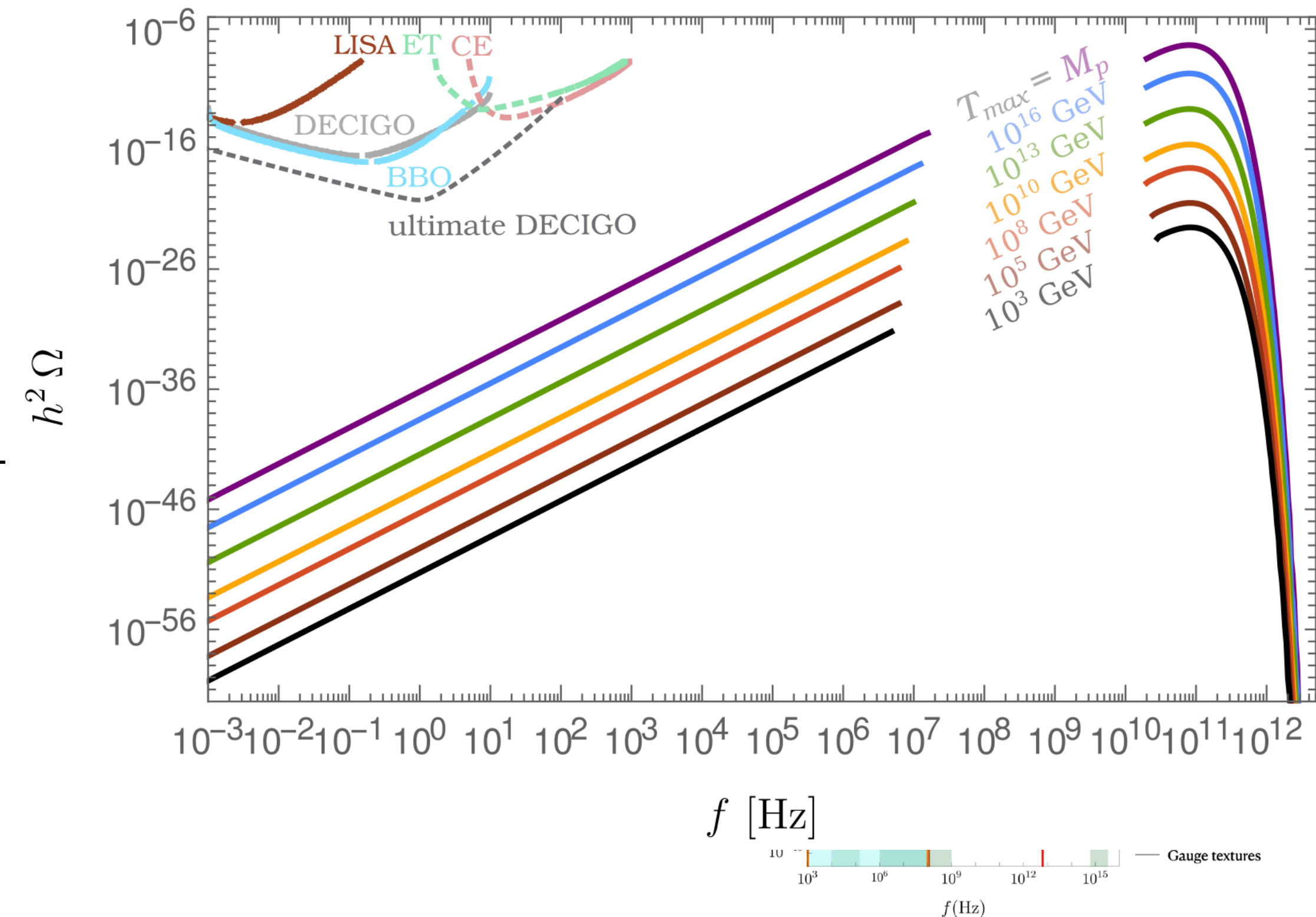
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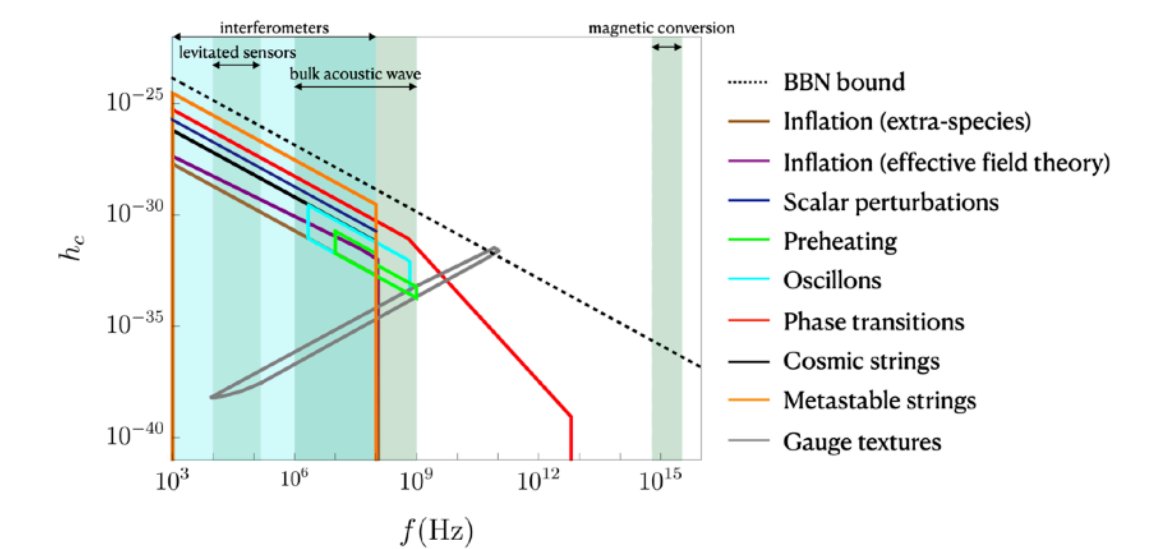
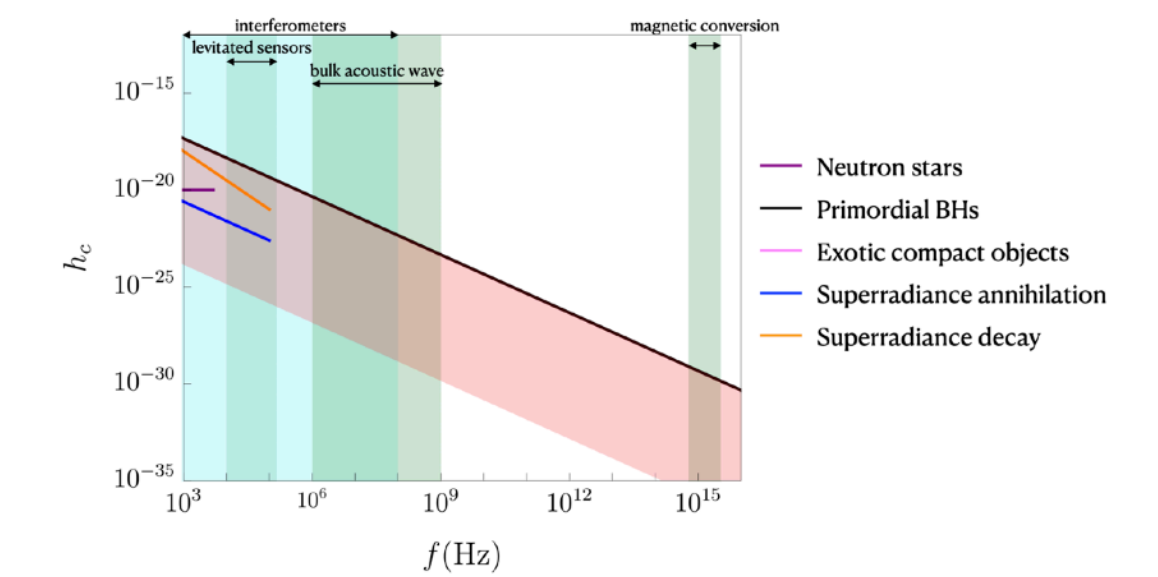
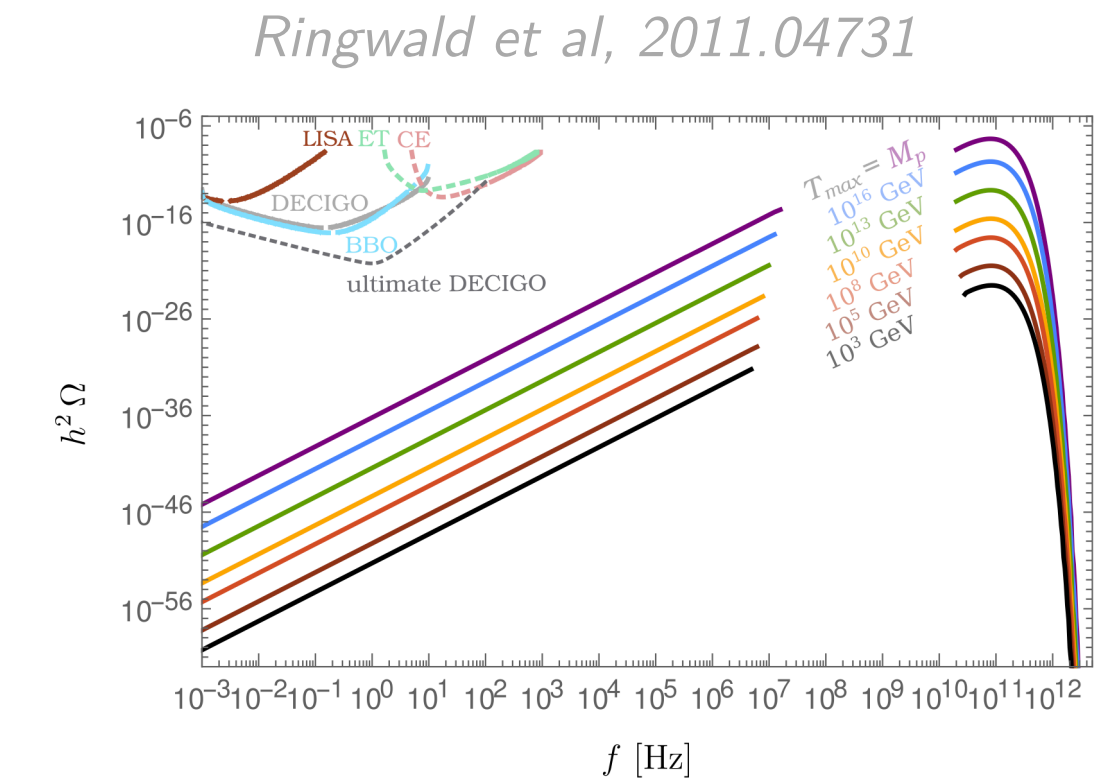
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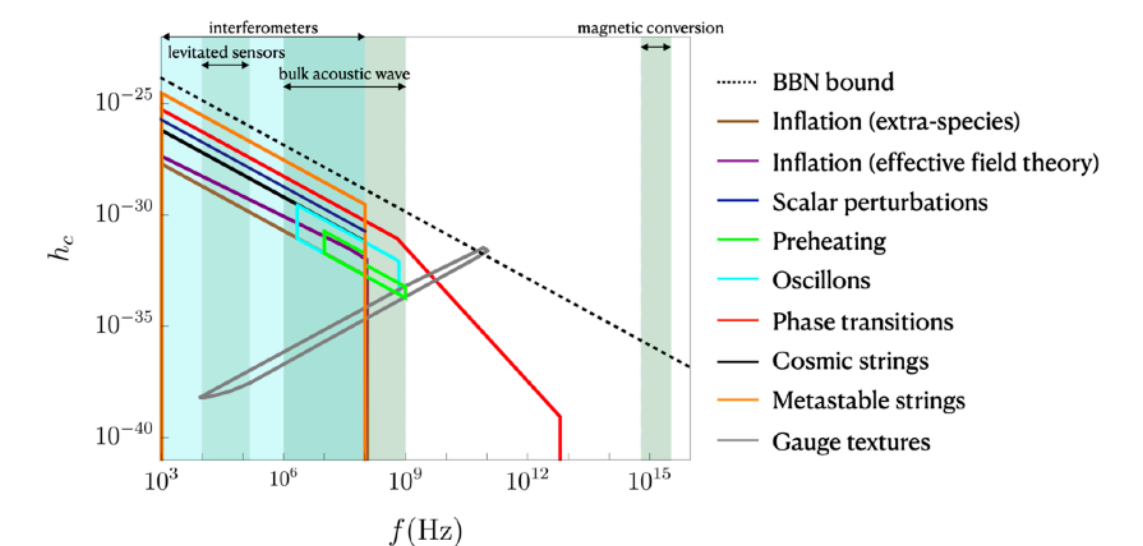
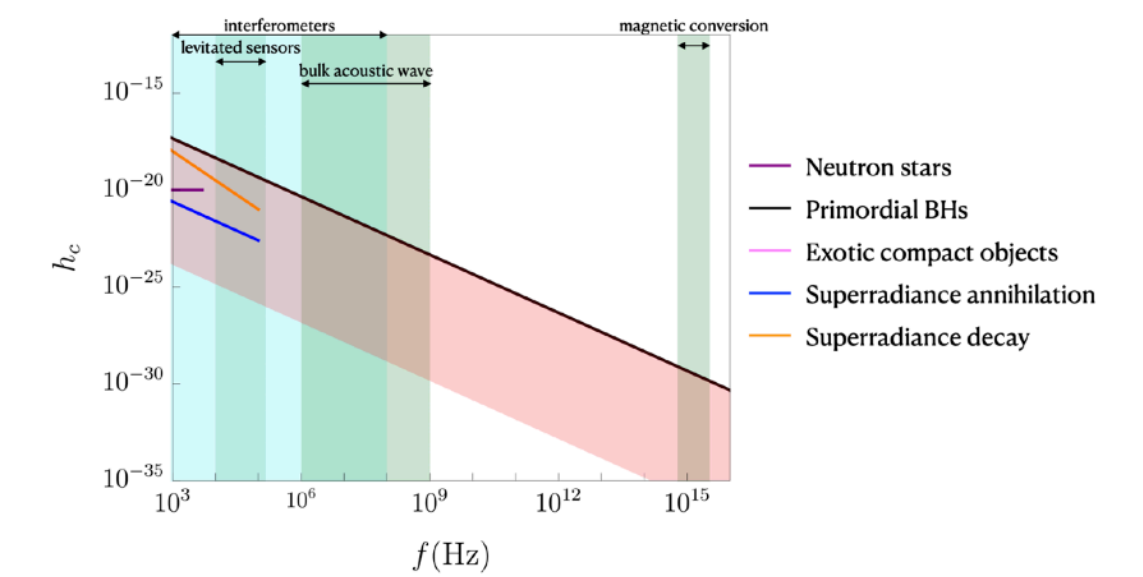
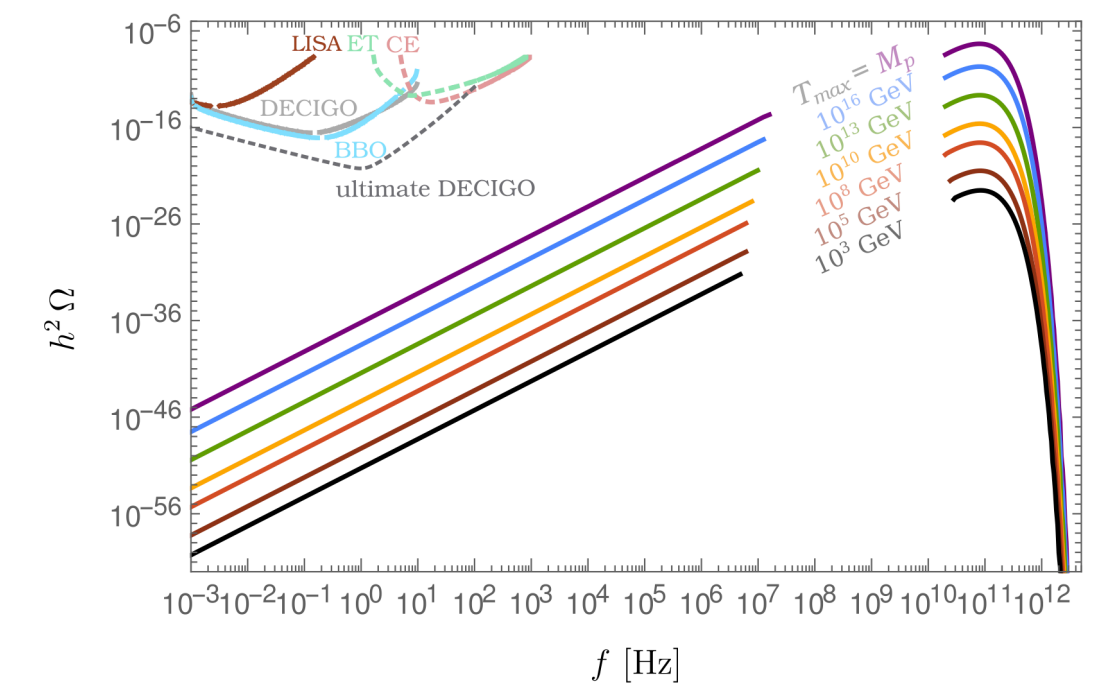
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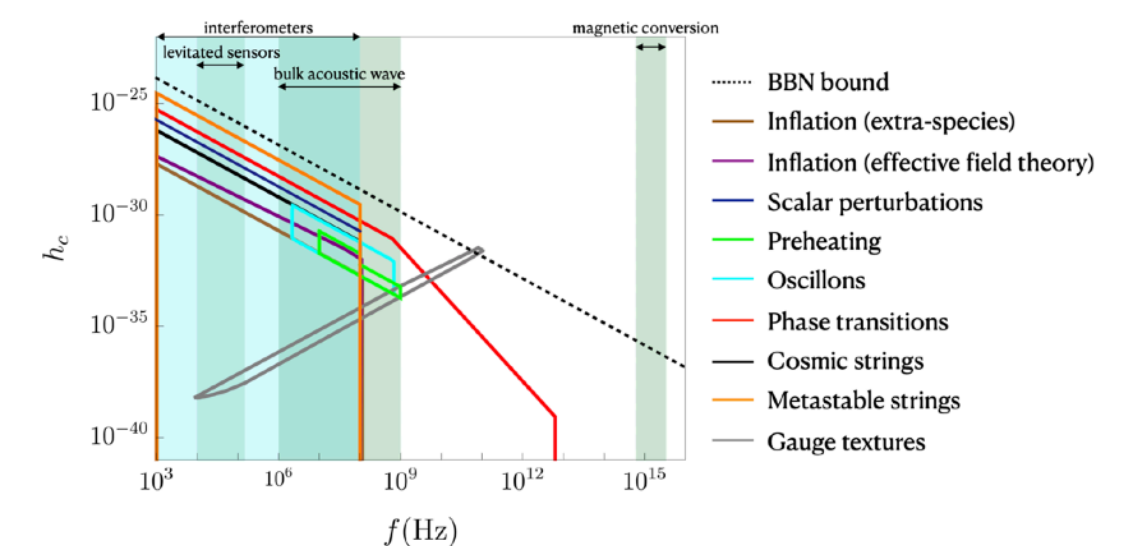
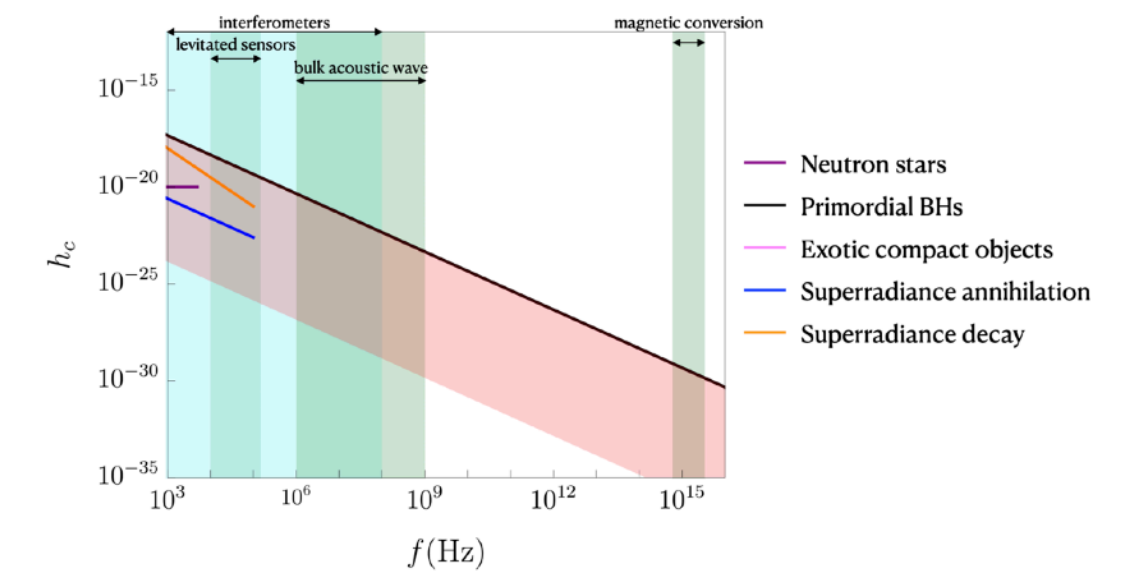
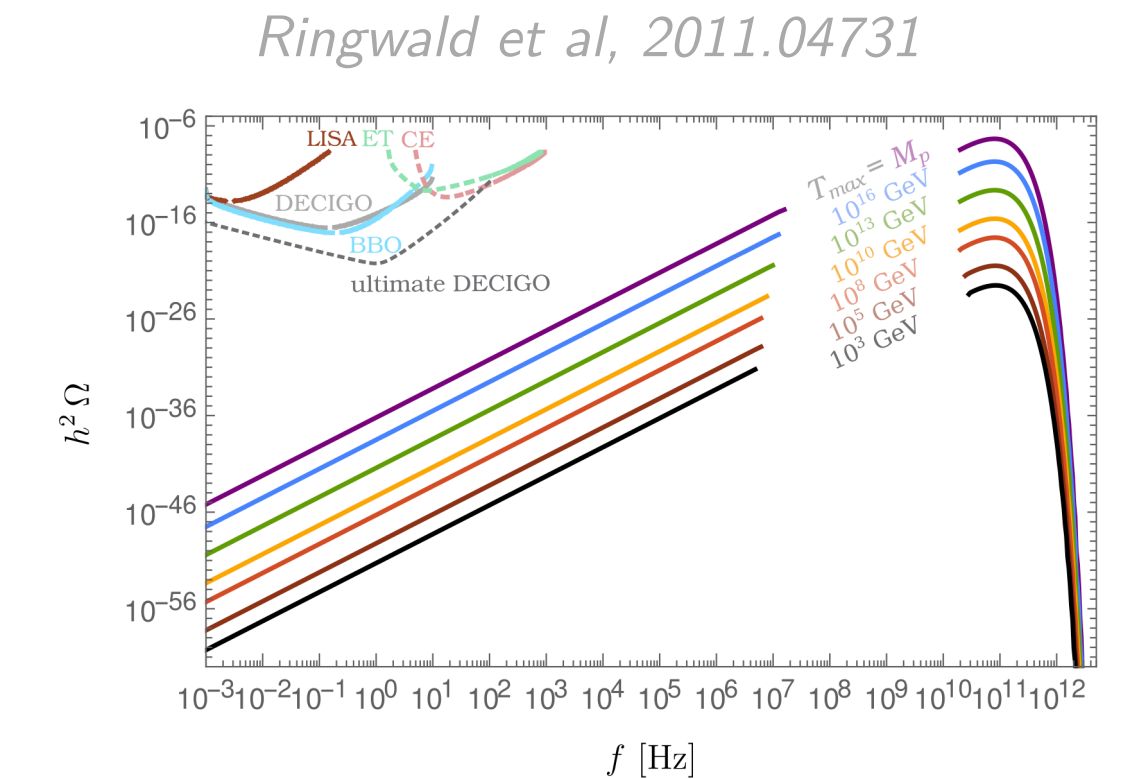


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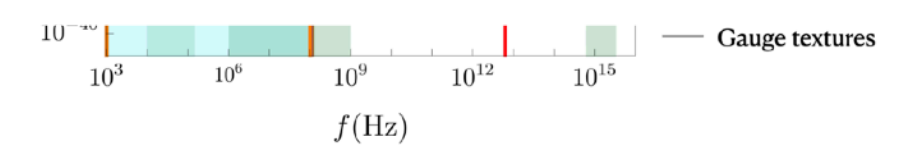
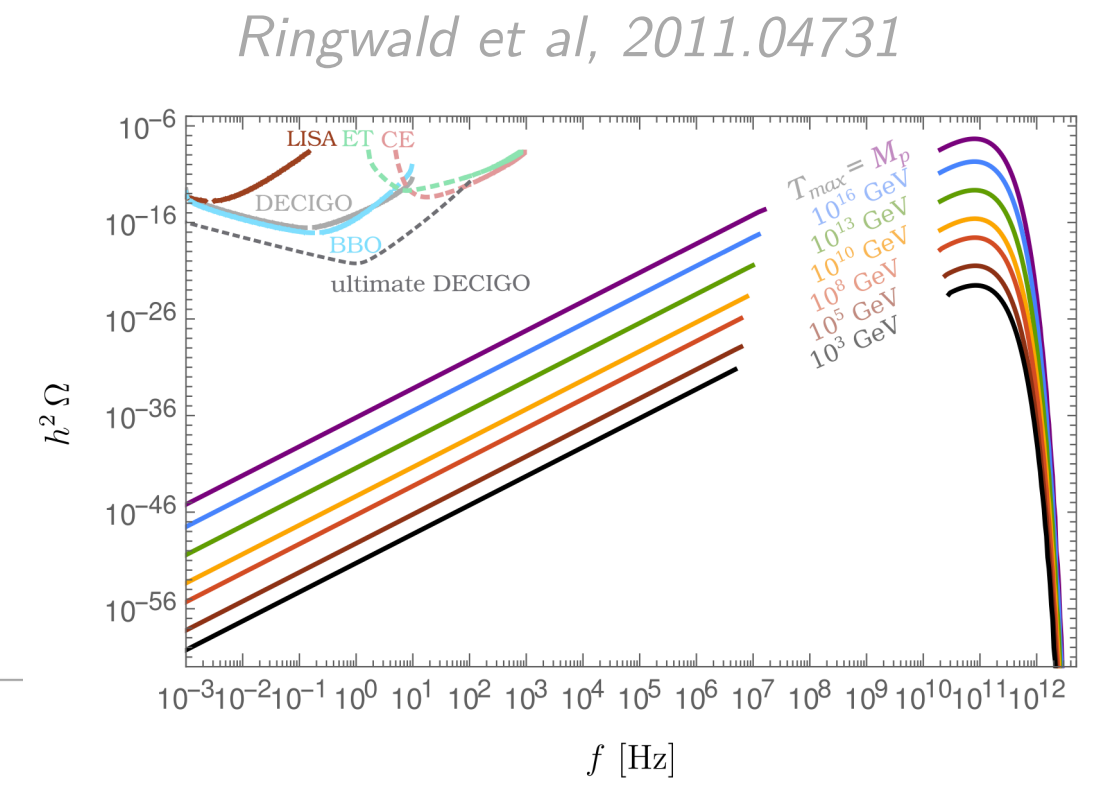
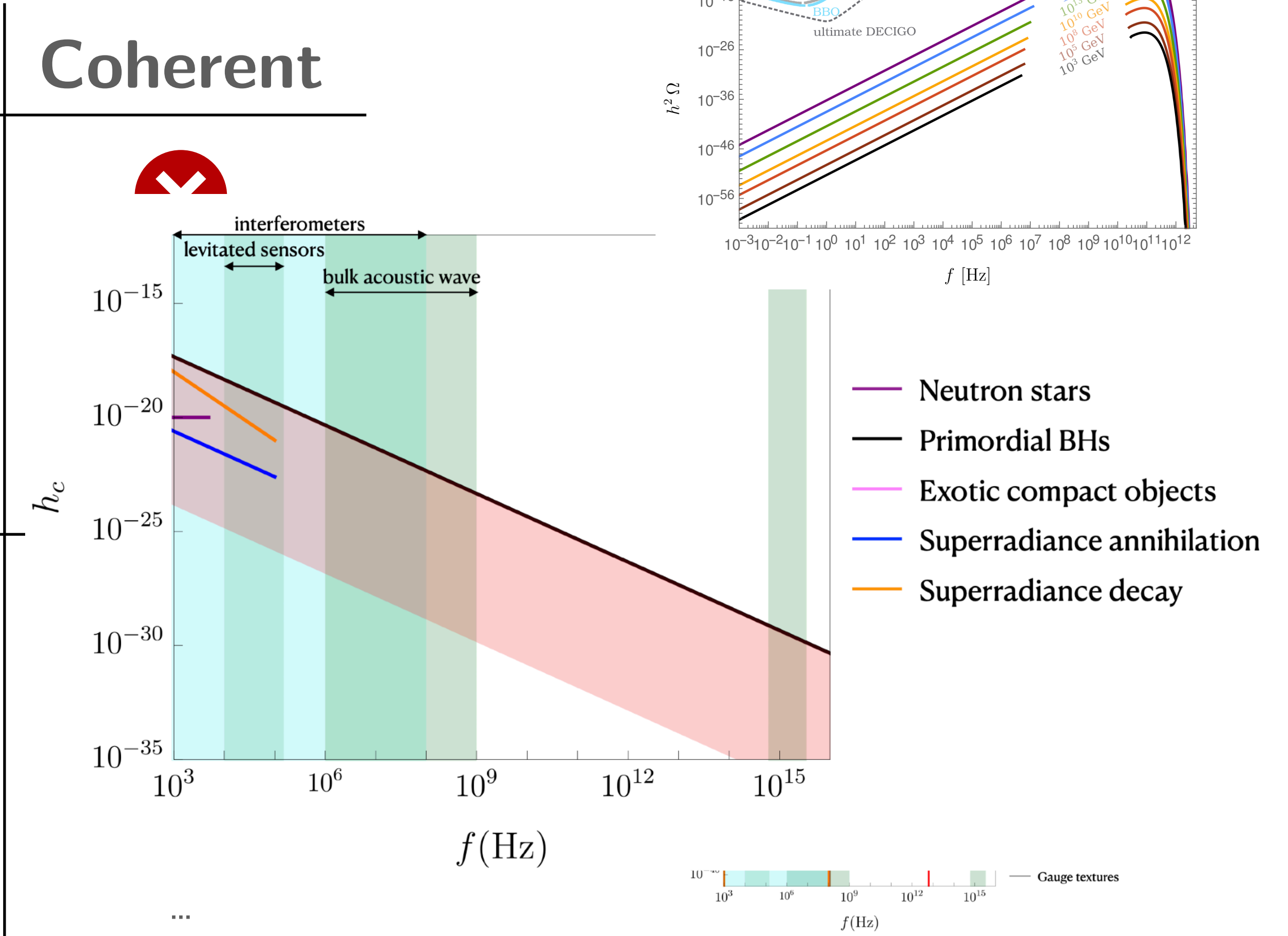
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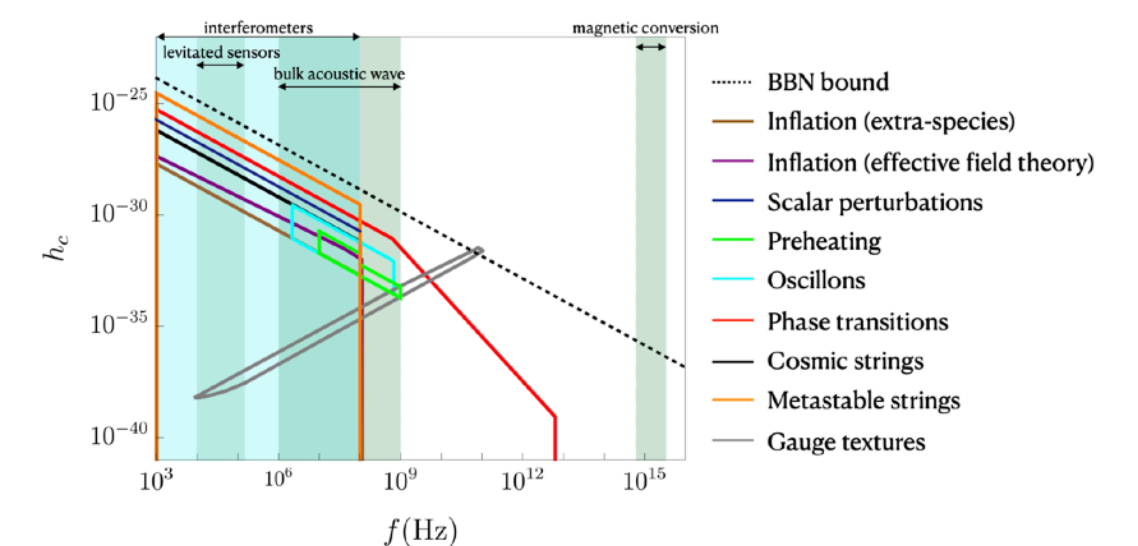
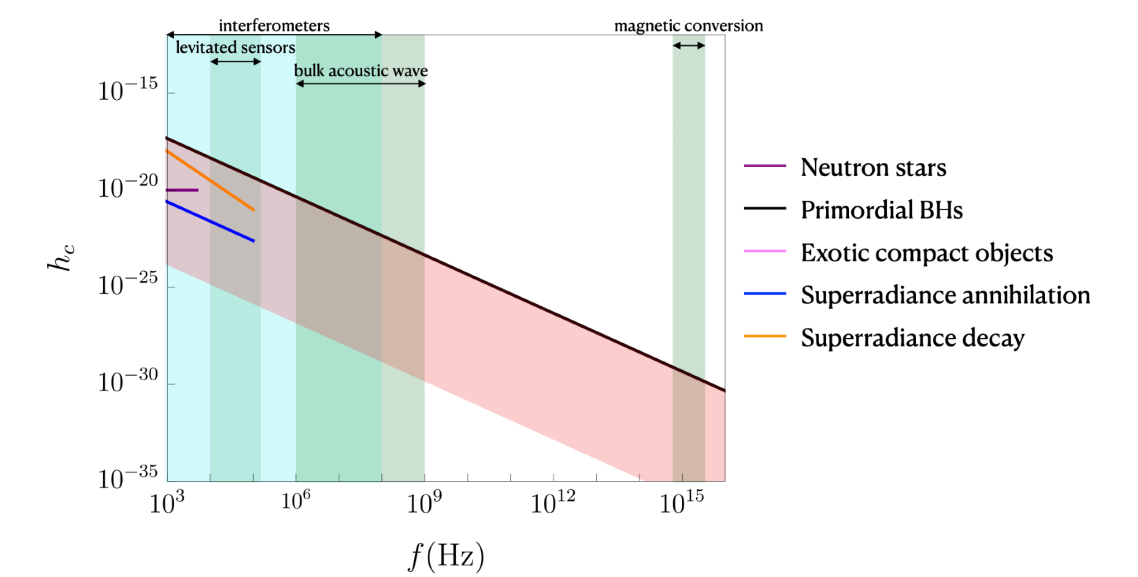
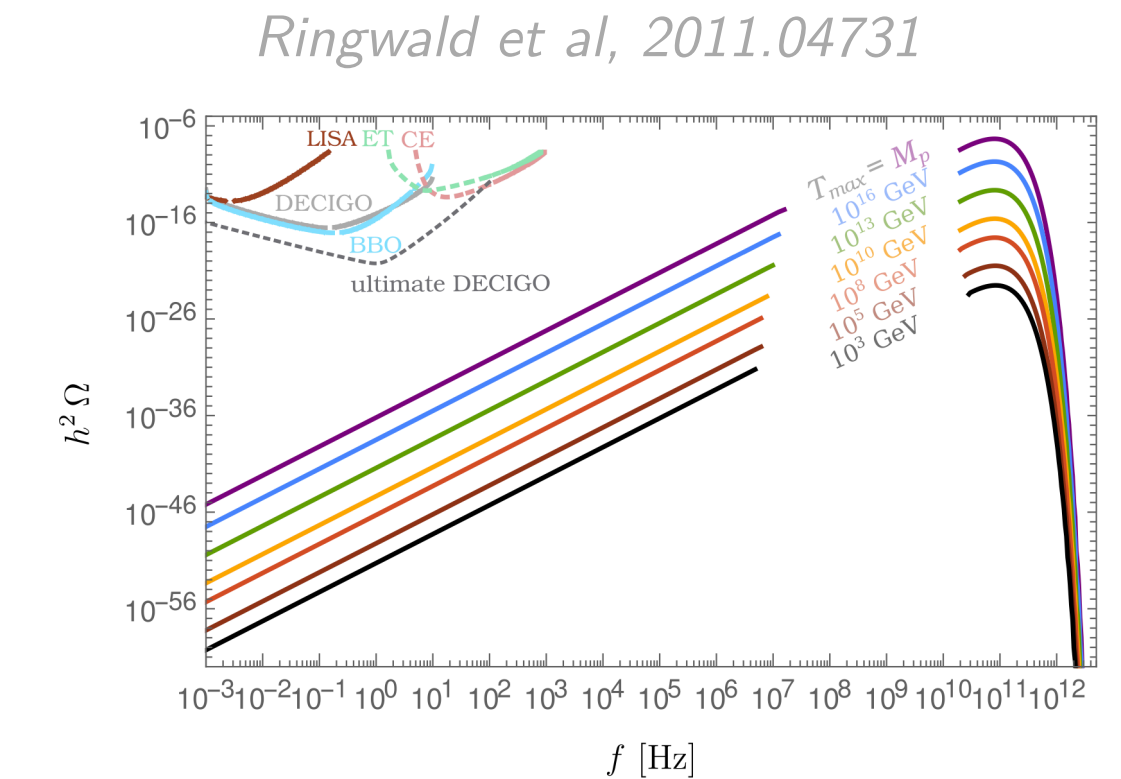


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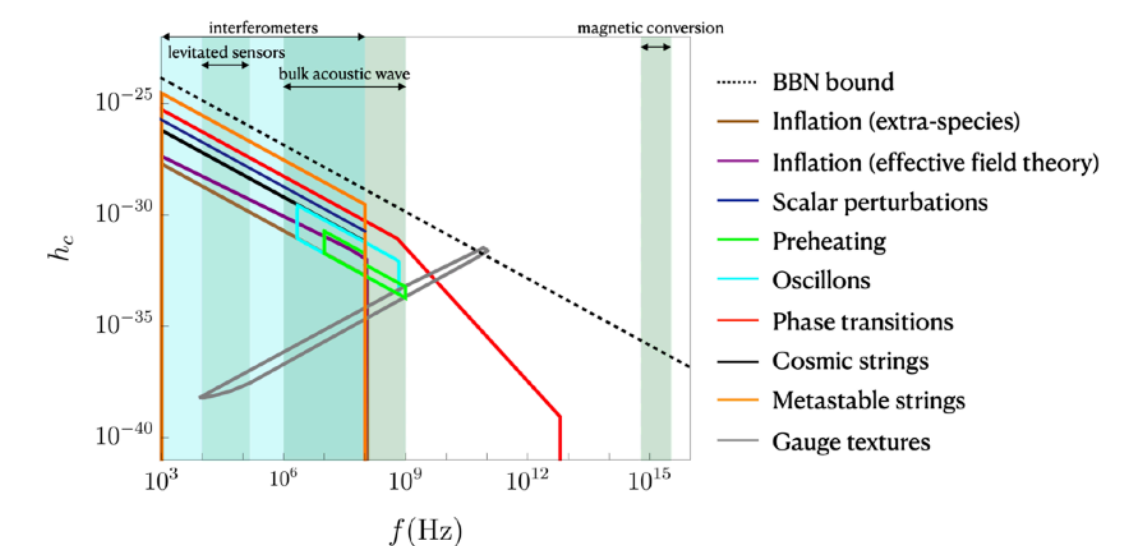
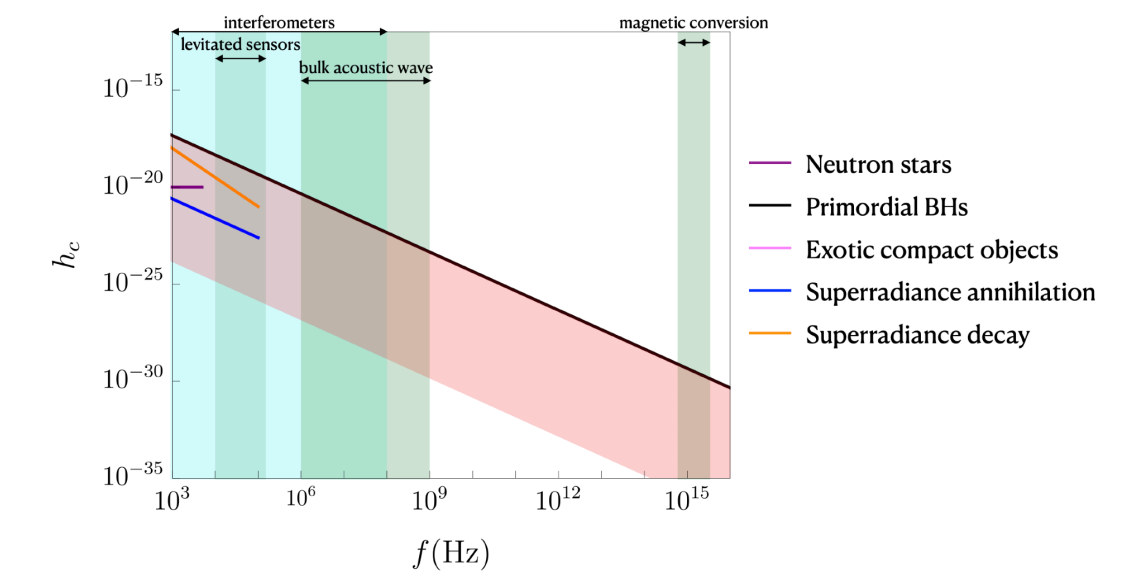
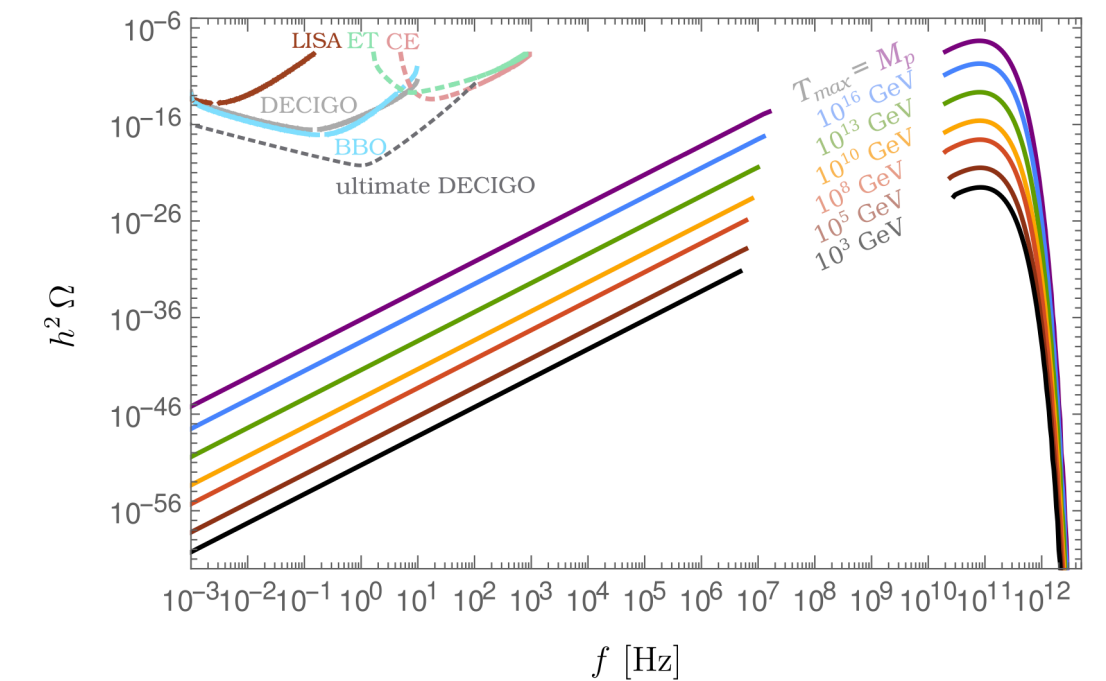
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Sources of HFGWs

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Standard Model: <div style="text-align: center;">  </div> <p style="text-align: center;">Thermal plasma fluctuations Ghiglieri & Laine (2015) Ghiglieri et al (2020) Ringwald et al (2020)</p>	<div style="text-align: center;">  </div> <div style="text-align: center;">  </div>	
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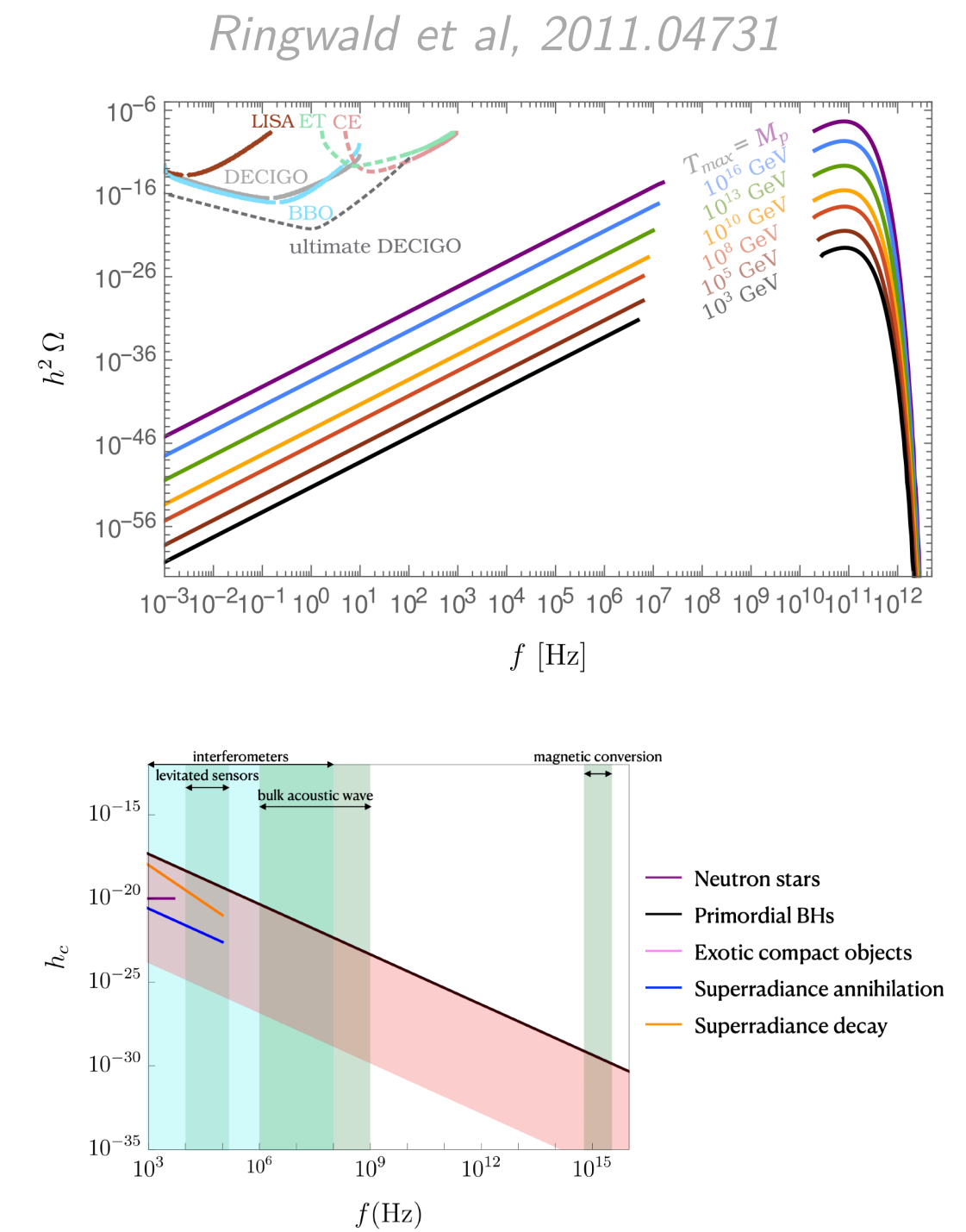
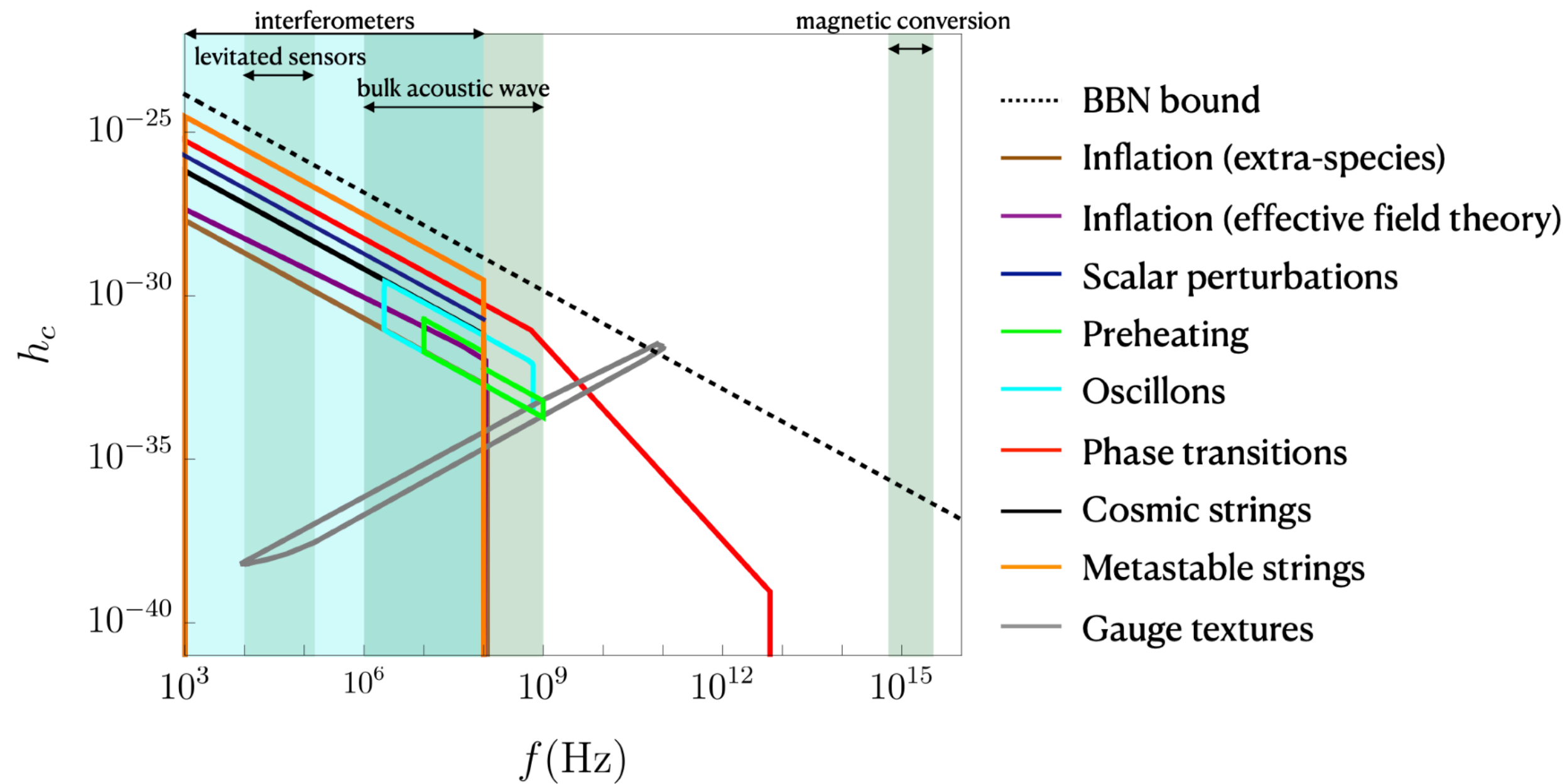


Aggarwal et al, 2011.12414

Sources of HFGWs

Sources discussed in detail Weds-Fri

Standard Model



BSM:



Inflation
Phase transitions
Cosmic Strings
...








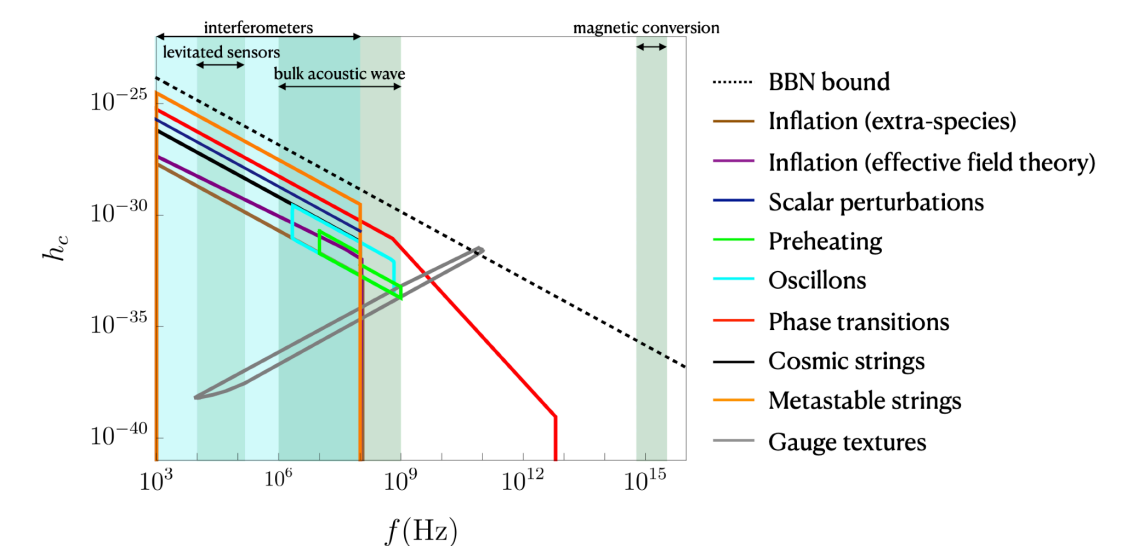
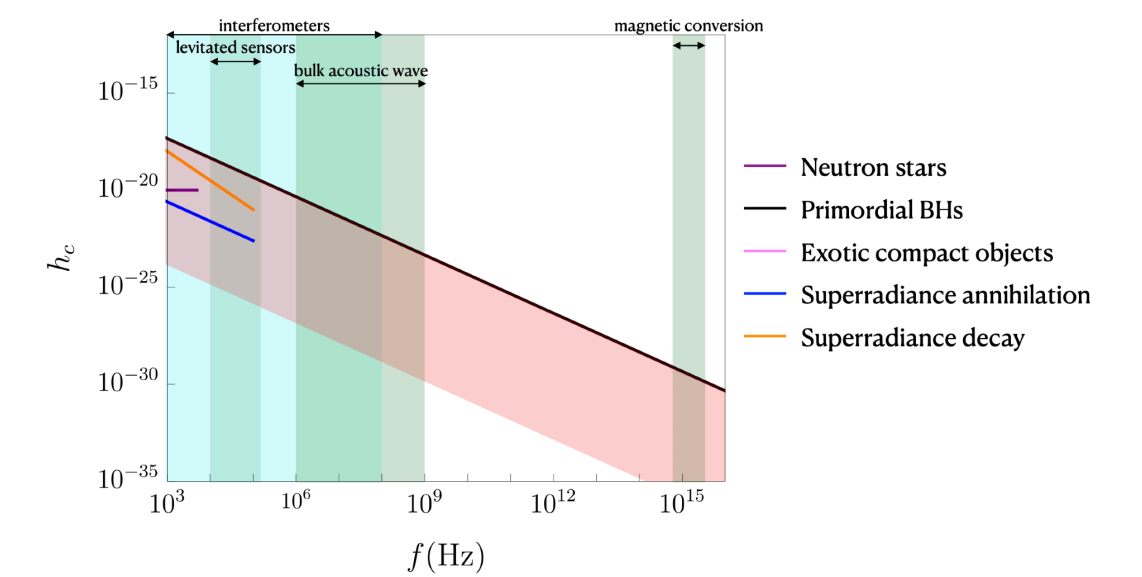
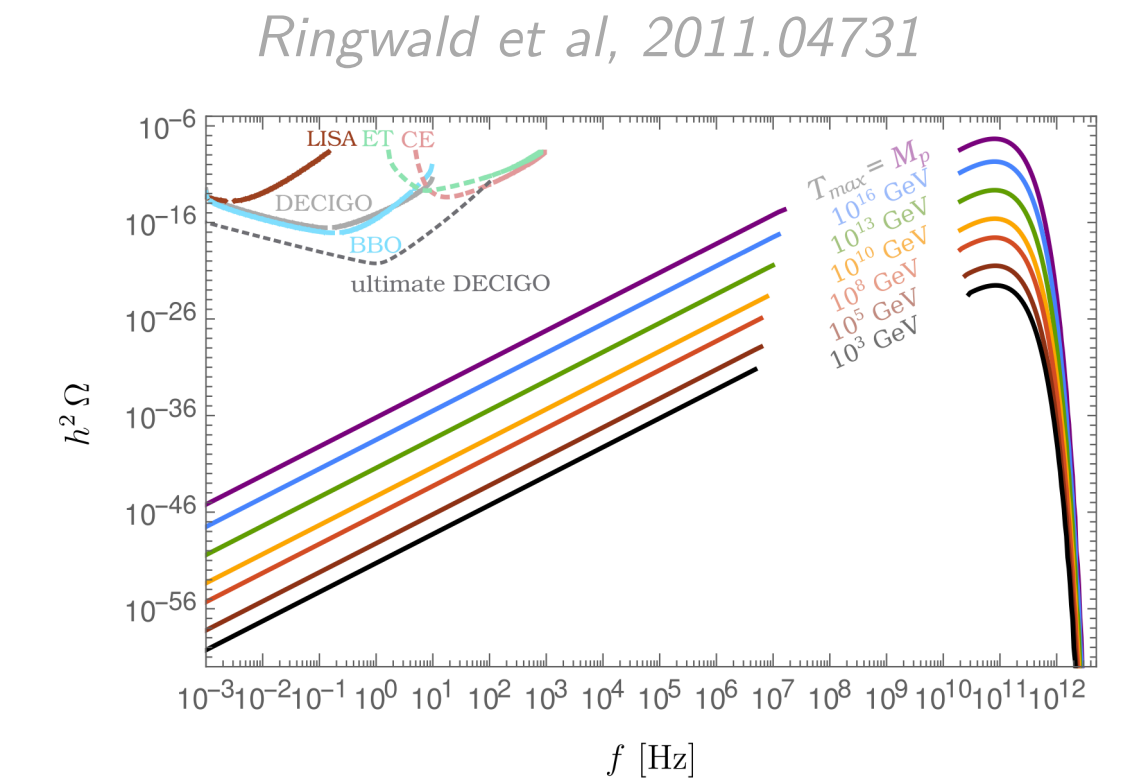
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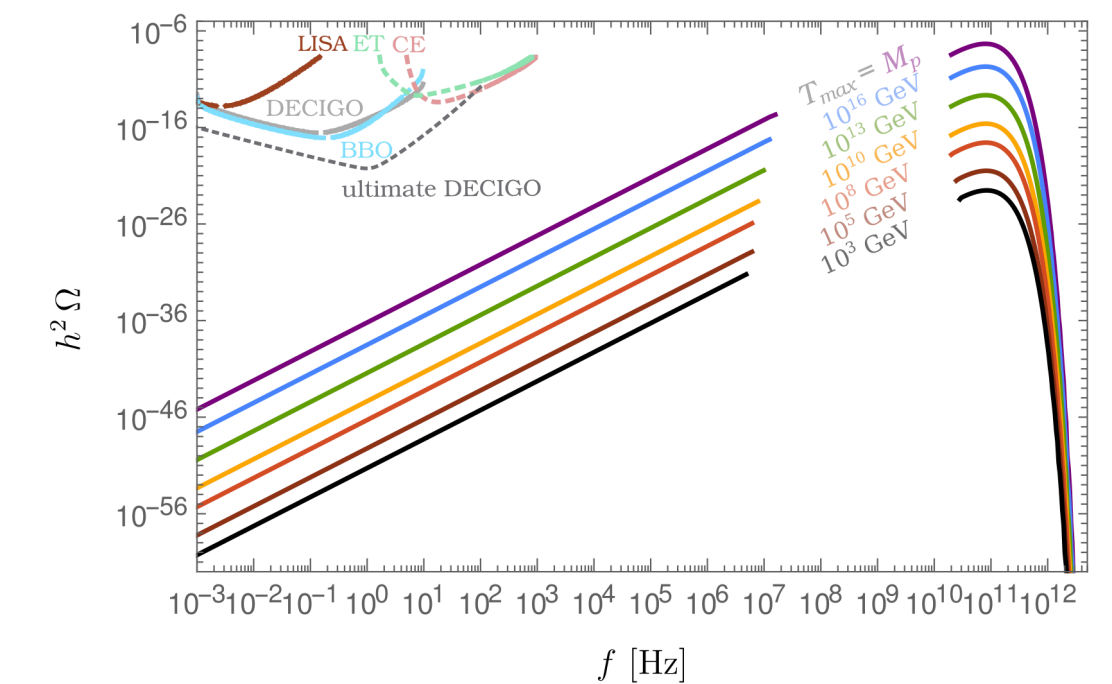
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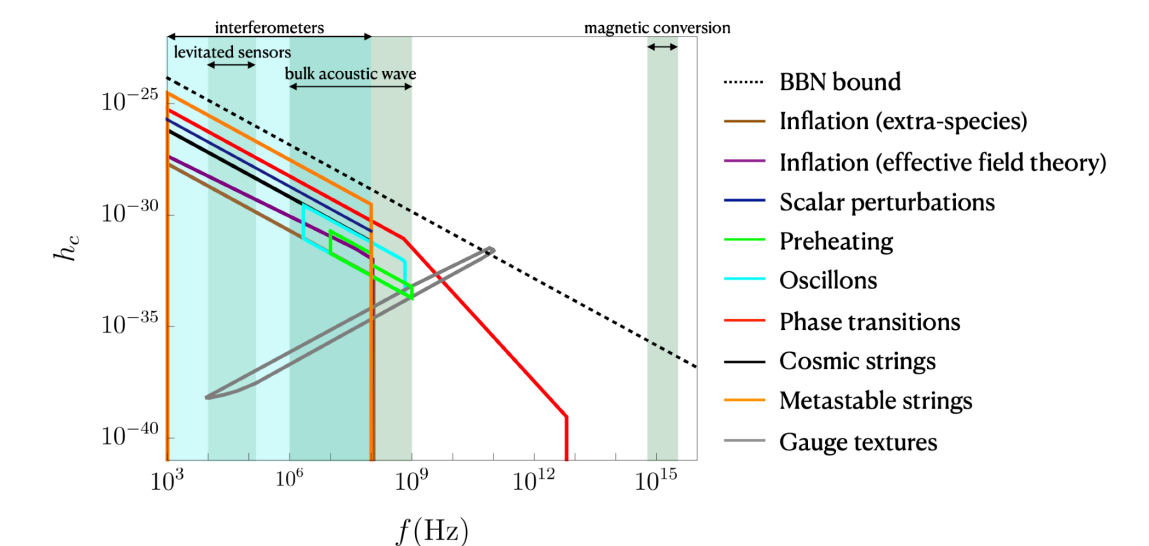
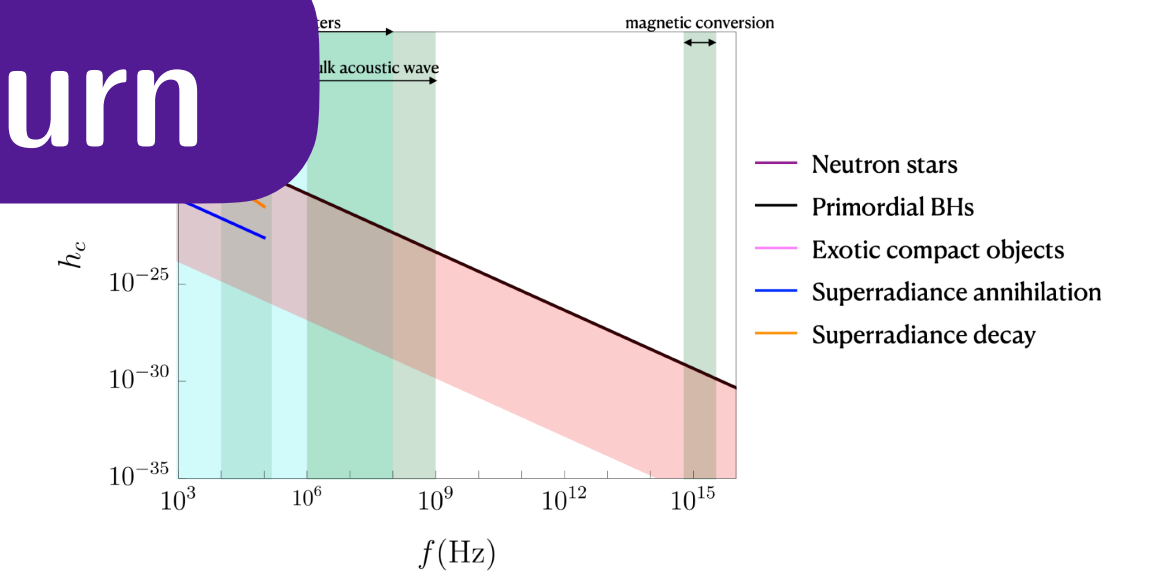
Sources discussed in detail Weds-Fri

	Stochastic	Coherent
Standard Model:	✔	✘
HFGW searches — high risk, high return		
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Ringwald et al, 2011.04731



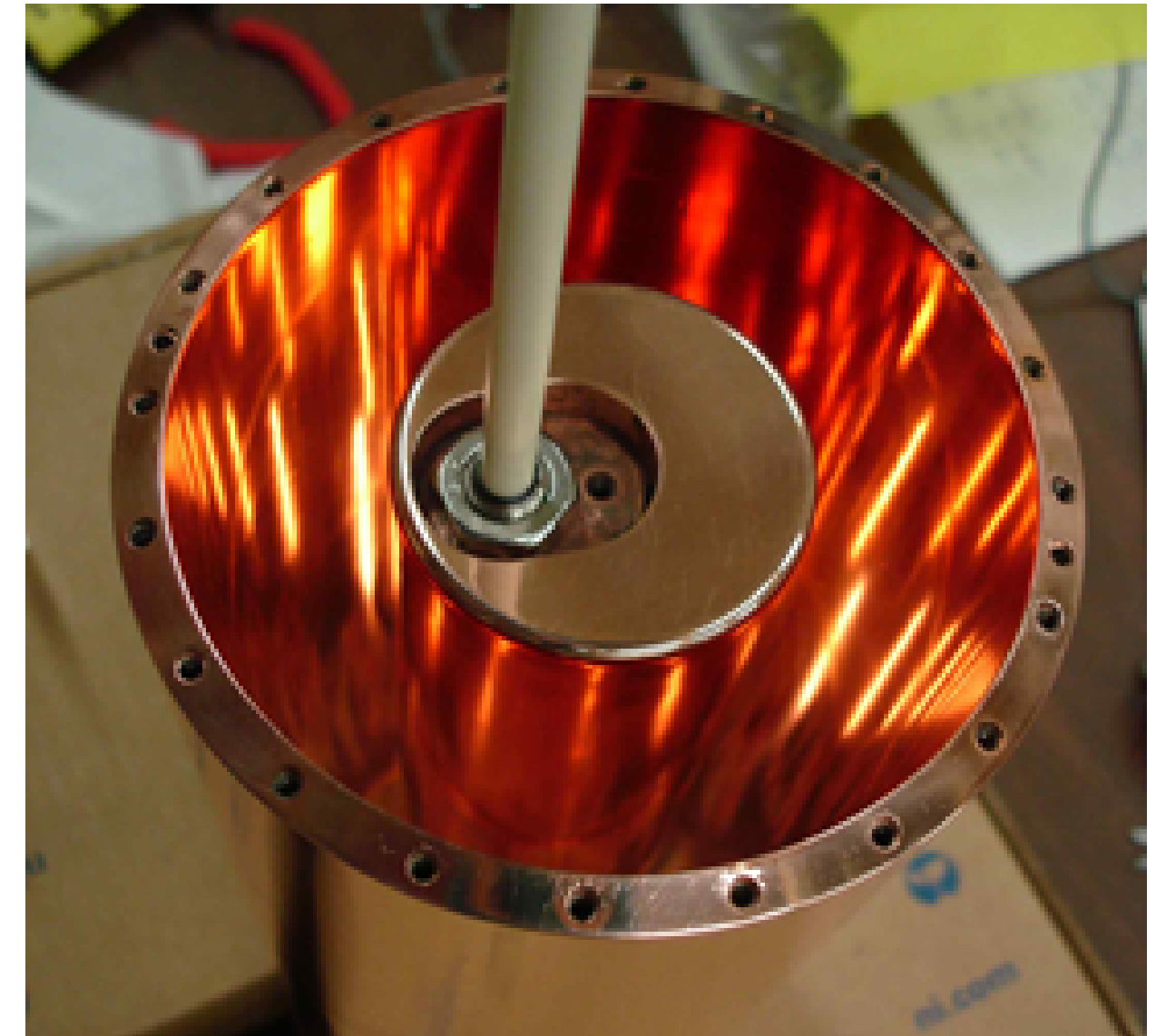
Ringwald et al (2020)



Aggarwal et al, 2011.12414

Resonant Cavities

Why?

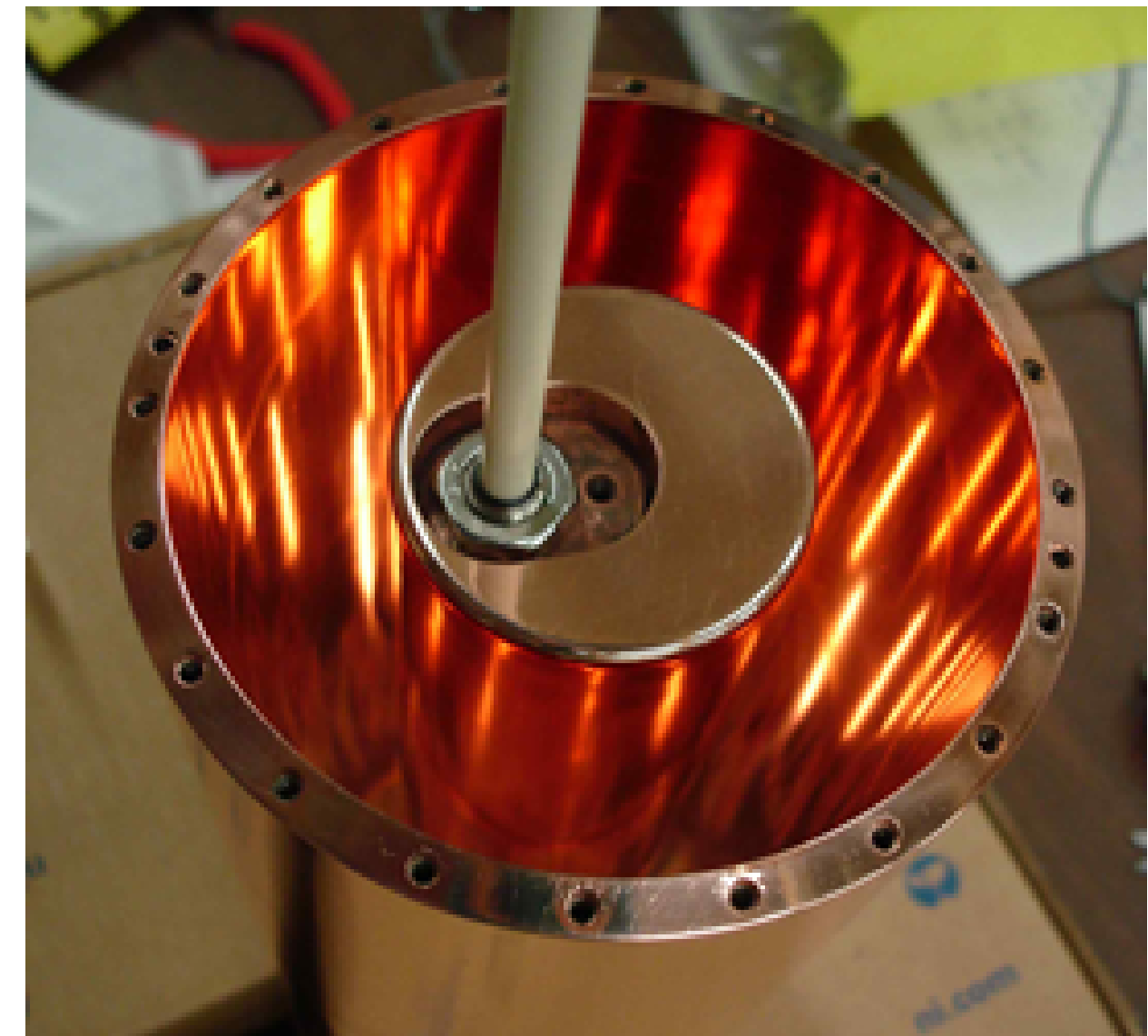


HAYSTAC

Resonant Cavities

Why?

Mature technology & constantly improving
Benefit from decades of development for accelerator use



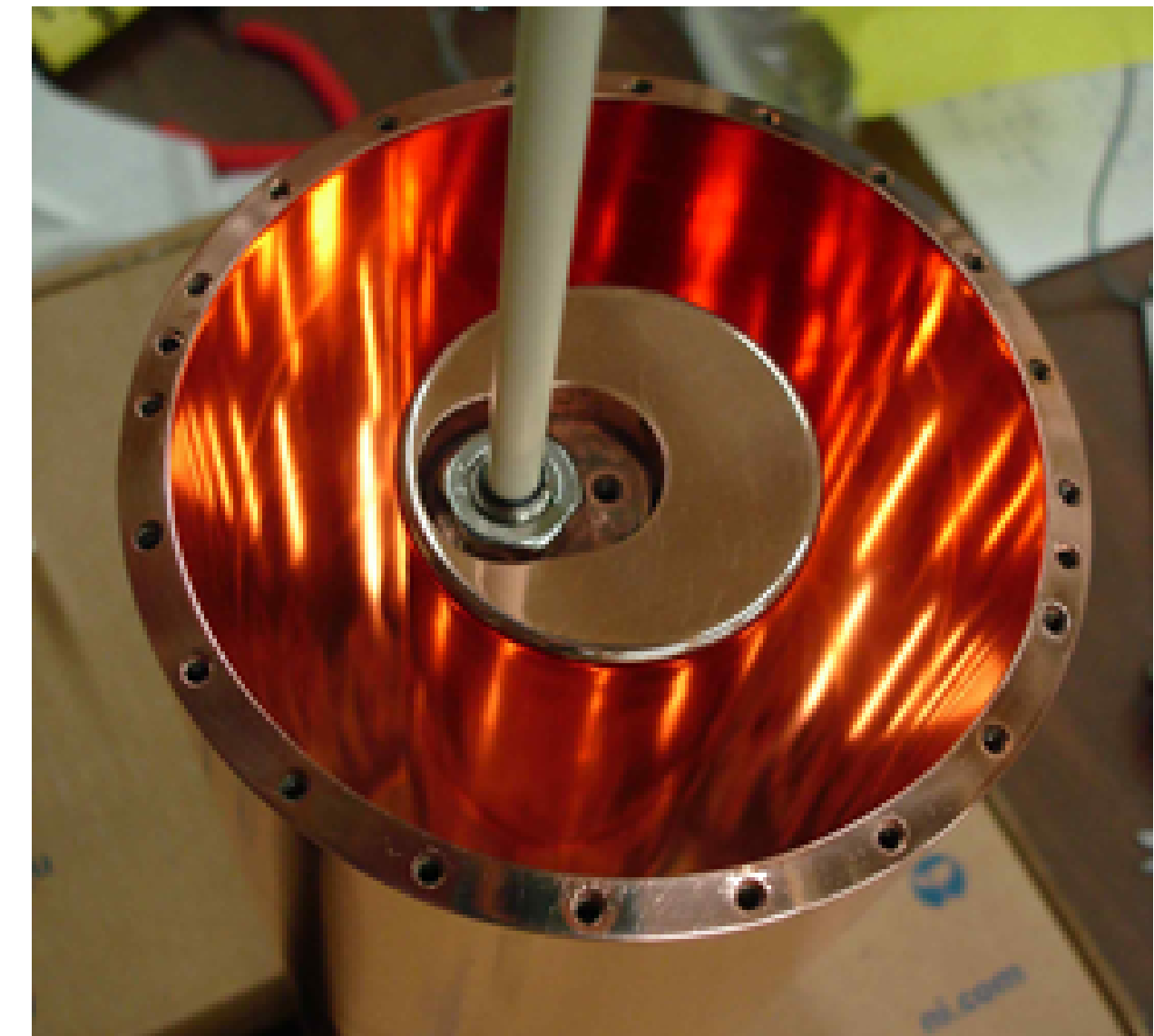
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Cavities for fundamental physics already in use:
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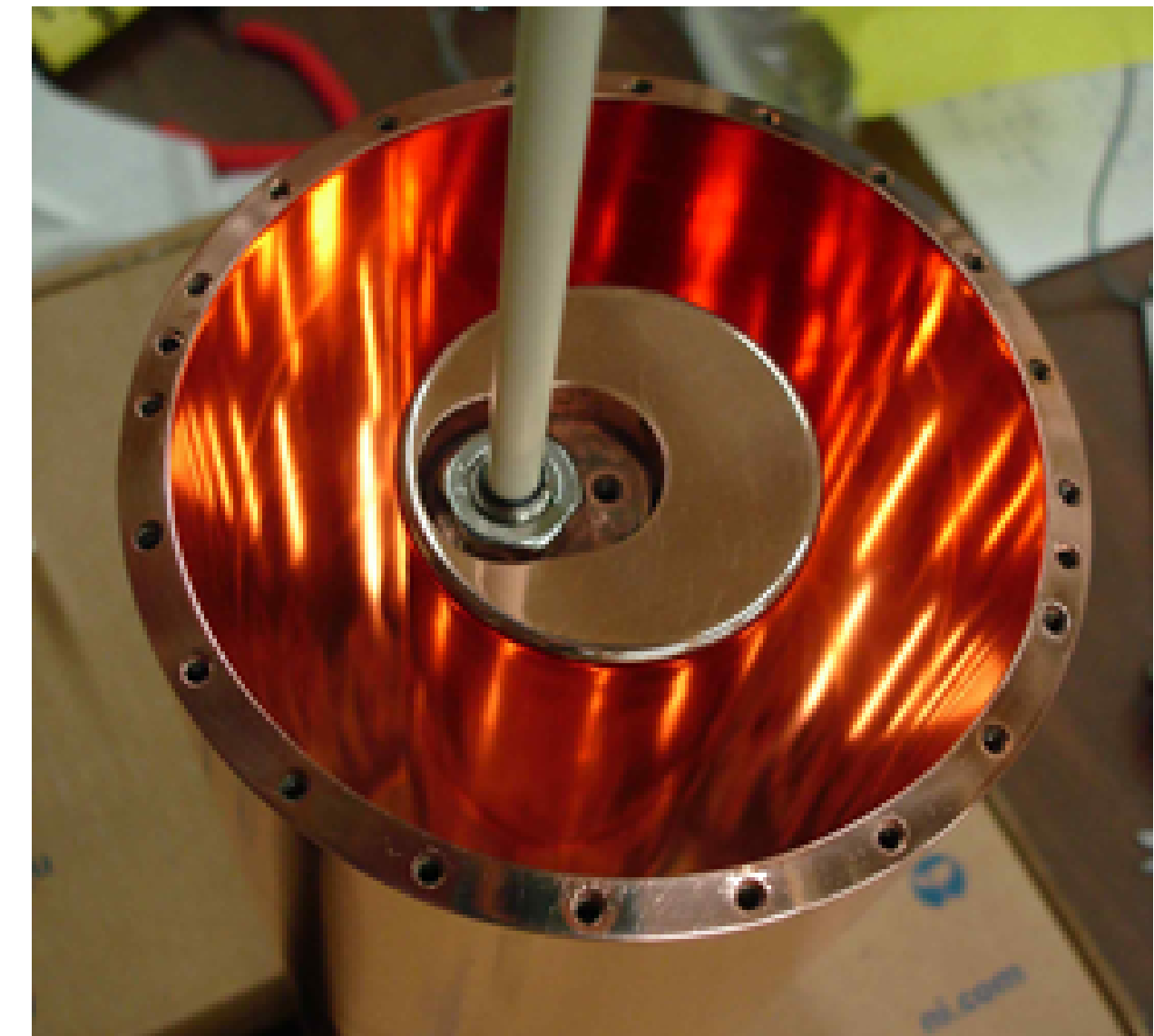
Resonant Cavities

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Domain where various subtleties arise...



HAYSTAC

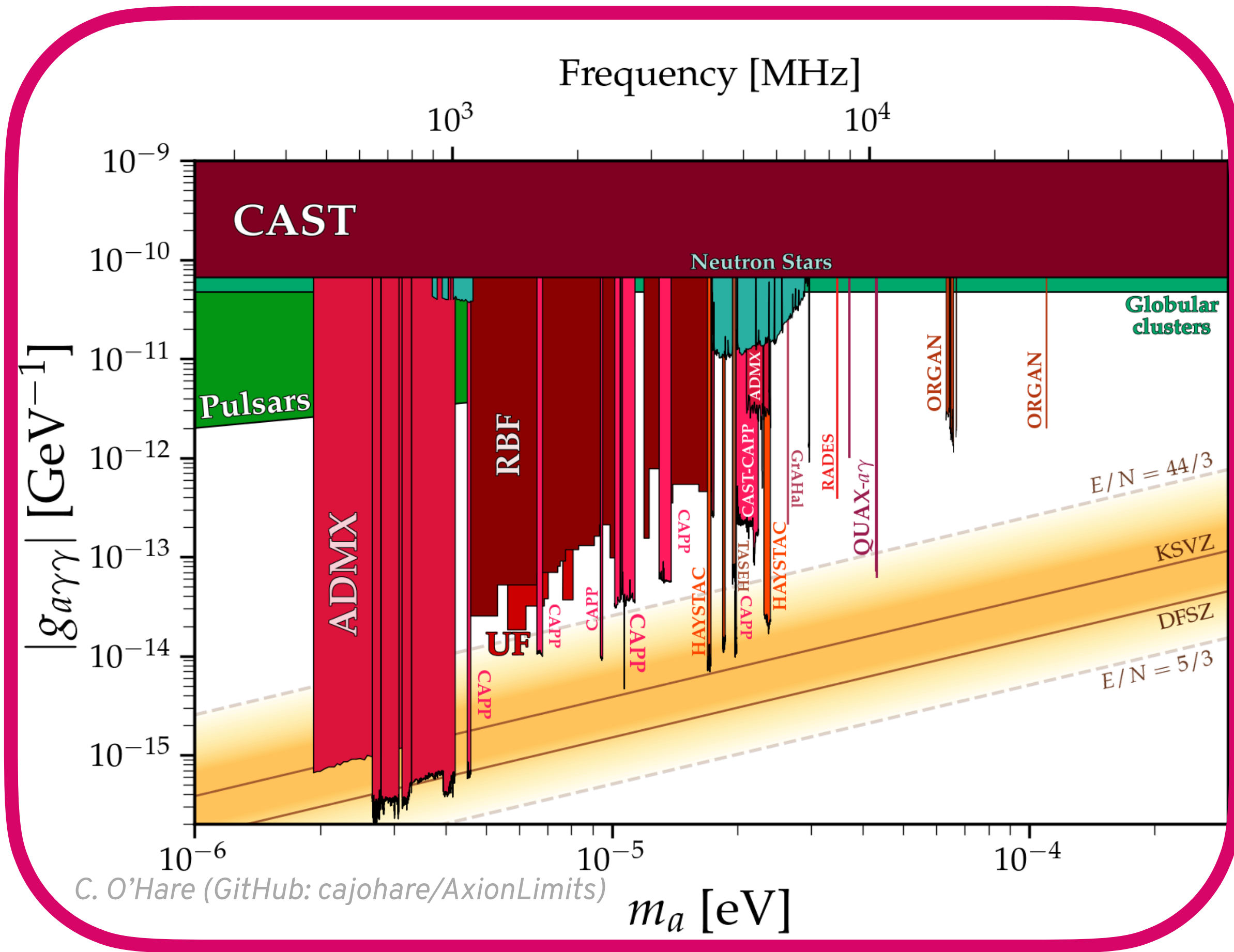
Intuition for EM signal

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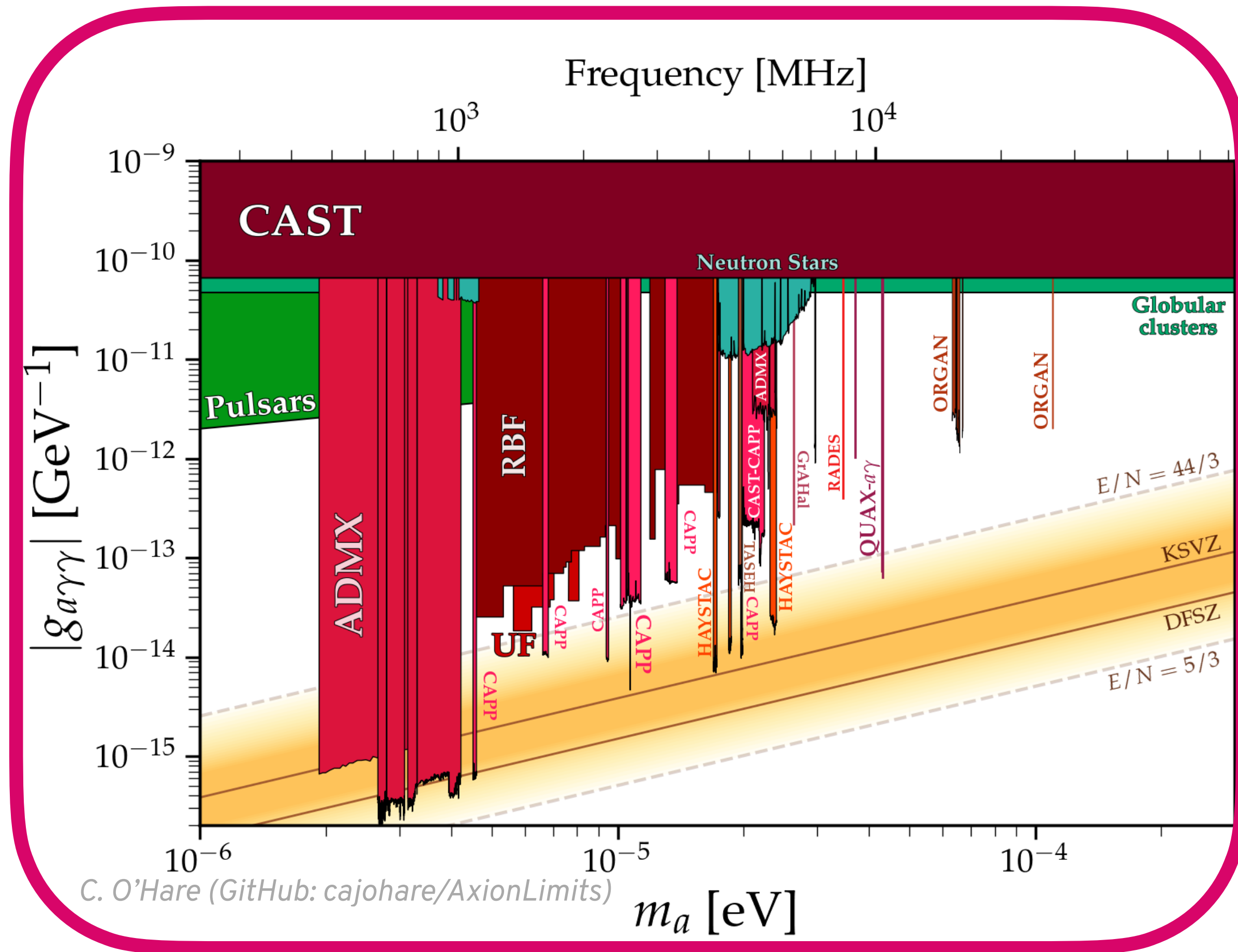
Estimate sensitivity to GWs by
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Intuition for EM signal

Estimate sensitivity to GWs by comparing sizes of currents



Intuition for EM signal

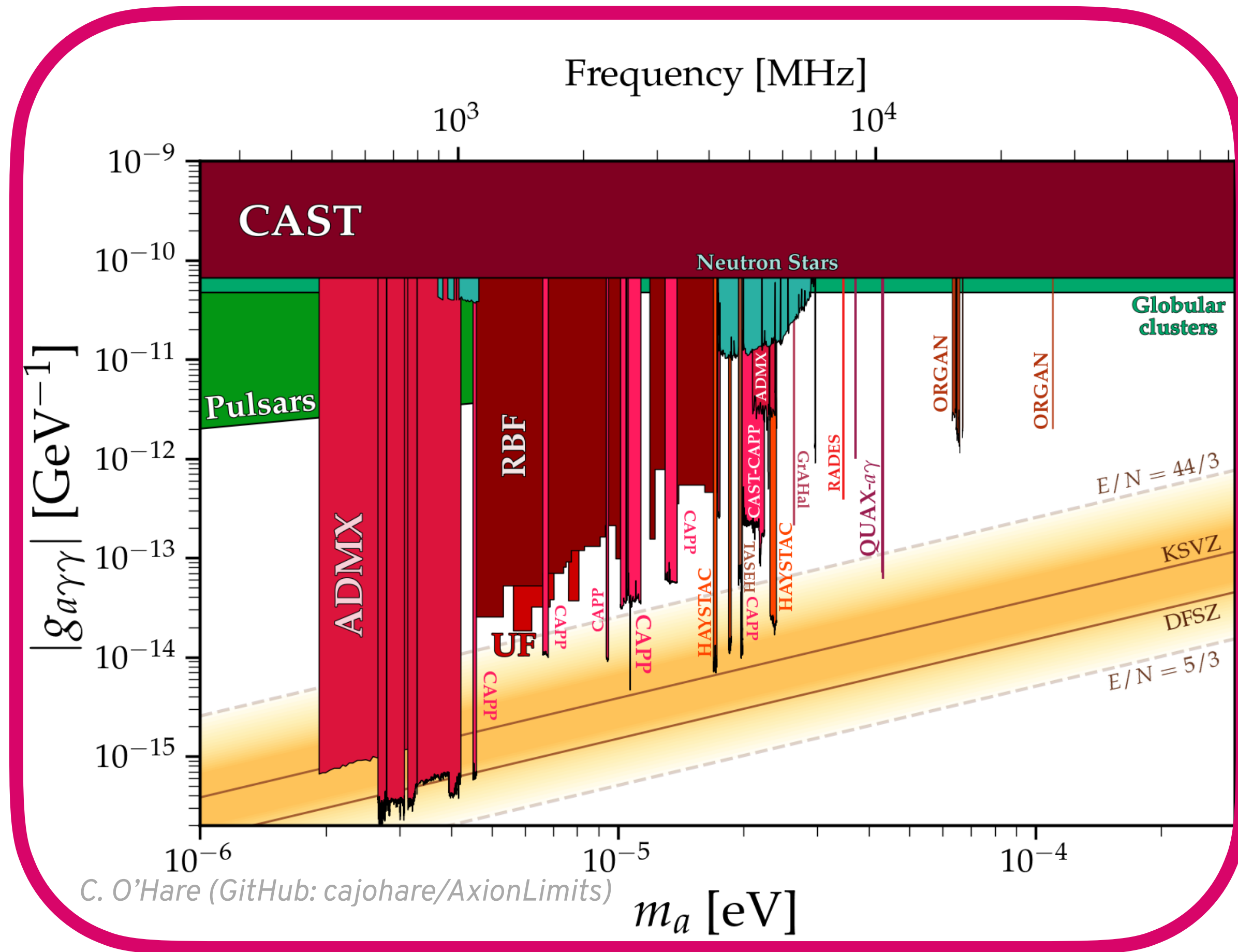


Estimate sensitivity to GWs by comparing sizes of currents

$$j_{\text{eff}}^{\text{axion}} \sim g_{a\gamma\gamma} \partial_t(a\mathbf{B}) + \mathcal{O}(v)$$

$$j_{\text{eff}}^{\text{axion}} \lesssim 10^{-19} \text{ T/m}$$

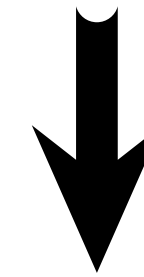
Intuition for EM signal



Estimate sensitivity to GWs by comparing sizes of currents

$$\dot{j}_{\text{eff}}^{\text{axion}} \sim g_{a\gamma\gamma} \partial_t(a\mathbf{B}) + \mathcal{O}(v)$$

$$\dot{j}_{\text{eff}}^{\text{axion}} \lesssim 10^{-19} \text{ T/m}$$



$$\dot{j}_{\text{eff}}^{\text{GW}} \sim \partial_t(h\mathbf{B}) + \dots$$

$$h \lesssim 10^{-21}$$

Framing the Question

A more detailed estimate requires some GR

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GW in TT gauge: $\partial_\mu h^{\mu\nu} = 0$, $h_\mu{}^\mu = 0$, $h_{00} = h_{0i} = 0$

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A more detailed estimate requires some GR

GW in TT gauge: $\partial_\mu h^{\mu\nu} = 0$, $h_\mu{}^\mu = 0$, $h_{00} = h_{0i} = 0$

Riemann tensor invariant at $O(h)$:

$$R_{0i0j} = -\frac{1}{2}\partial_t^2 h_{ij}^{\text{TT}},$$

$$R_{0ijk} = \frac{1}{2}\partial_t (\partial_k h_{ij}^{\text{TT}} - \partial_j h_{ik}^{\text{TT}}),$$

$$R_{ikjl} = \frac{1}{2} (\partial_k \partial_j h_{il}^{\text{TT}} + \partial_i \partial_l h_{jk}^{\text{TT}} - \partial_i \partial_j h_{kl}^{\text{TT}} - \partial_k \partial_l h_{ij}^{\text{TT}})$$

MAGO 2.0

Revive an old idea from the 1970s, and a prototype from the 2000s

On the operation of a tunable electromagnetic detector for gravitational waves

F Pegoraro[†], E Picasso[‡] and L A Radicati^{‡§}

[†]Scuola Normale Superiore, Pisa, Italy

[‡]CERN, Geneva, Switzerland

Received 6 December 1977, in final form 20 April 1978

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Microwave Apparatus for Gravitational Waves Observation

R. Ballantini, A. Chincarini, S. Cuneo, G. Gemme, R. Parodi, A. Podestà, and R. Vaccarone
INFN and Università degli Studi di Genova, Genova, Italy

Ph. Bernard, S. Calatroni, E. Chiaveri, and R. Losito
CERN, Geneva, Switzerland

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E. Picasso
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MAGO 2.0

Revive an old idea from the 1970s, and a prototype from the 2000s

On the operation of a tunable electromagnetic detector for gravitational waves

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Received 6 December 1977, in final form 20 April 1978

Microwave Apparatus for Gravitational Waves Observation

R. Ballantini, A. Chincarini, S. Cuneo, G. Gemme, R. Parodi, A. Podestà, and R. Vaccarone
INFN and Università degli Studi di Genova, Genova, Italy

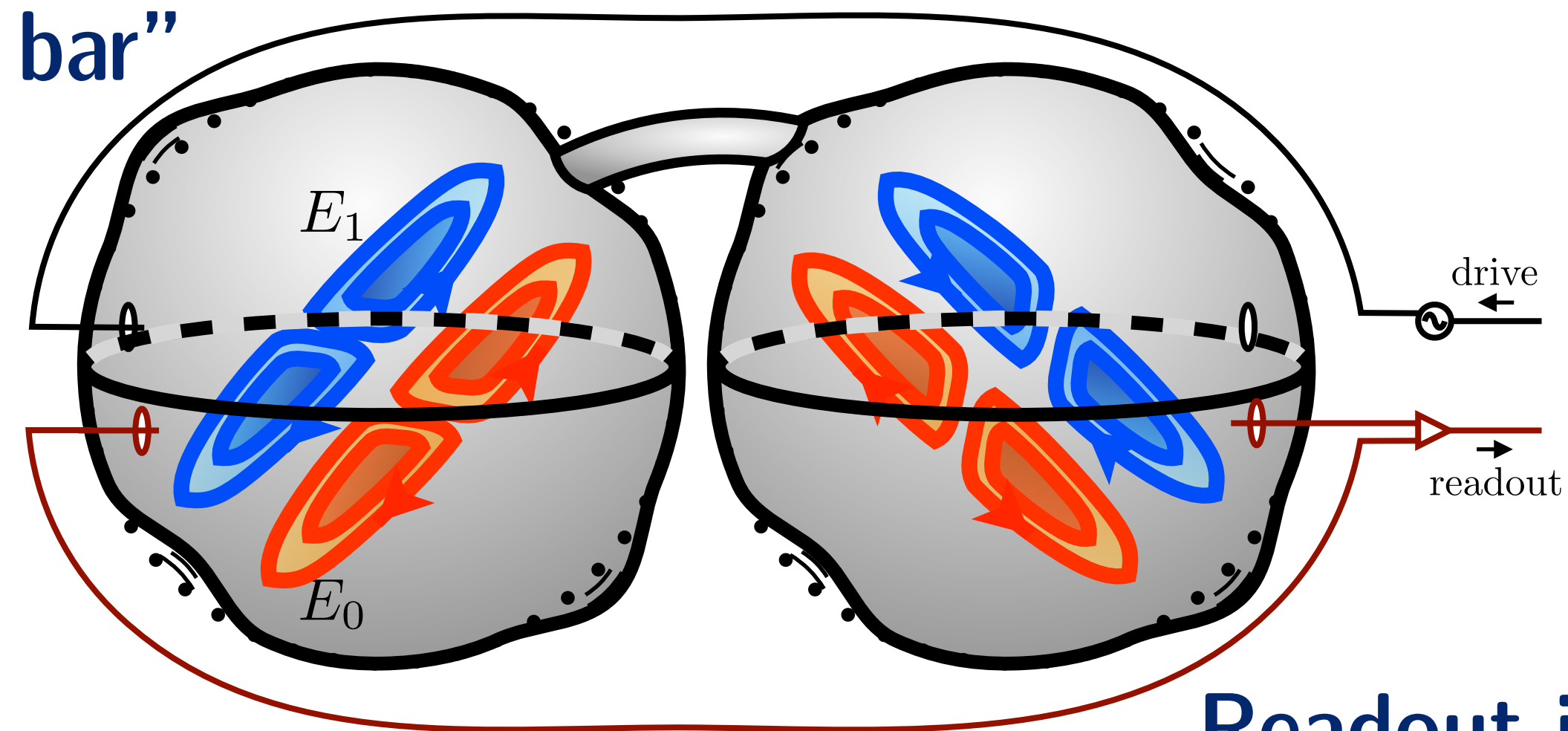
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Cavity walls are a “Weber bar”



Readout in a quiet mode of cavity

Gravitational Wave and a Hollow Sphere

Gravitational Wave and a Hollow Sphere

Mechanical modes of a sphere

$$\mathbf{U}_{lmn} = \nabla\phi_L + i\nabla \times \mathbf{L}\phi_{T_1} + i\mathbf{L}\phi_{T_2} .$$

$$\mathbf{U}(\mathbf{x}, t) = u_p(t) \mathbf{U}_p(\mathbf{x})$$

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Equation of motion

$$\ddot{u}_p + \frac{\omega_p}{Q_p} \dot{u}_p + \omega_p^2 u_p \simeq -\frac{1}{2} \omega_g^2 V_{\text{cav}}^{1/3} \eta_{\text{mech}}^g h_0 e^{i\omega_g t}$$

$$\eta_{\text{mech}}^g = \frac{\hat{h}_{ij}^{TT}}{V_{\text{cav}}^{1/3} V_{\text{shell}}} \int_{V_{\text{shell}}} d^3\mathbf{x} U_p^{*i} x^j$$

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Tiny displacement \ll nm

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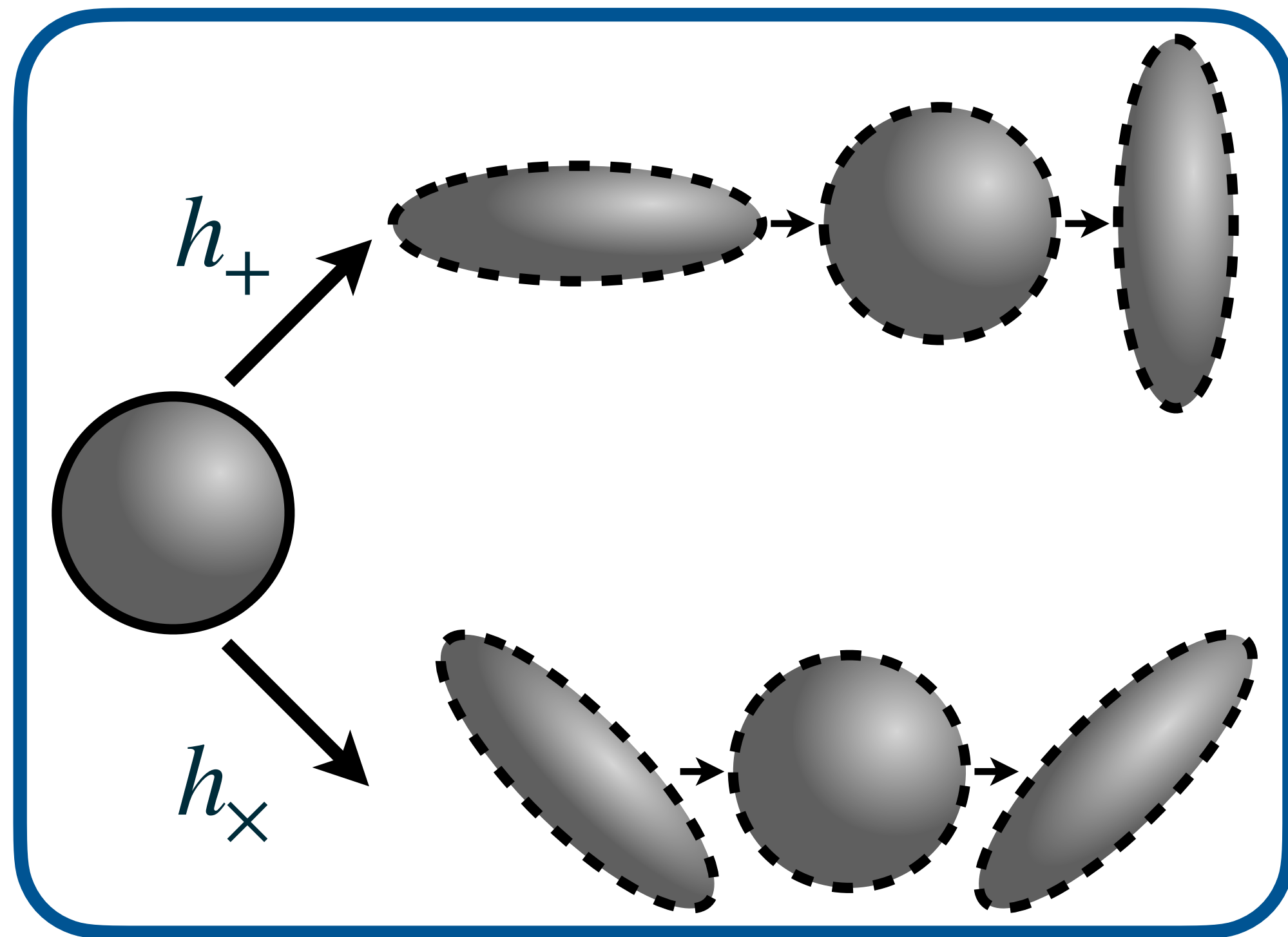
Tiny displacement \ll nm

Cur Cavis?* pt. 2

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* “Why Cavities?” in Latin

Gravitational Wave and a Hollow Sphere

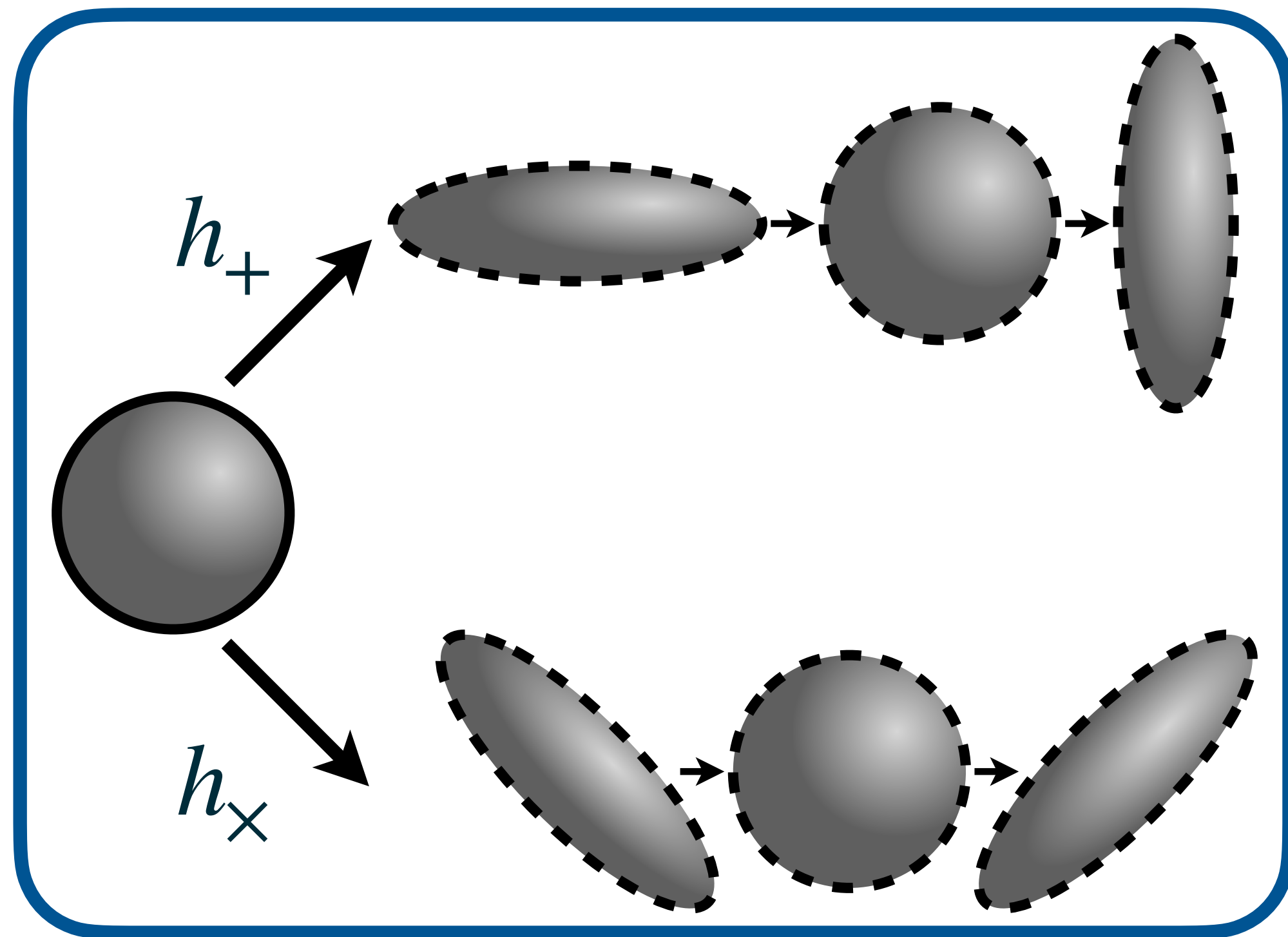


TT frame intuition

Gravitational Wave and a Hollow Sphere

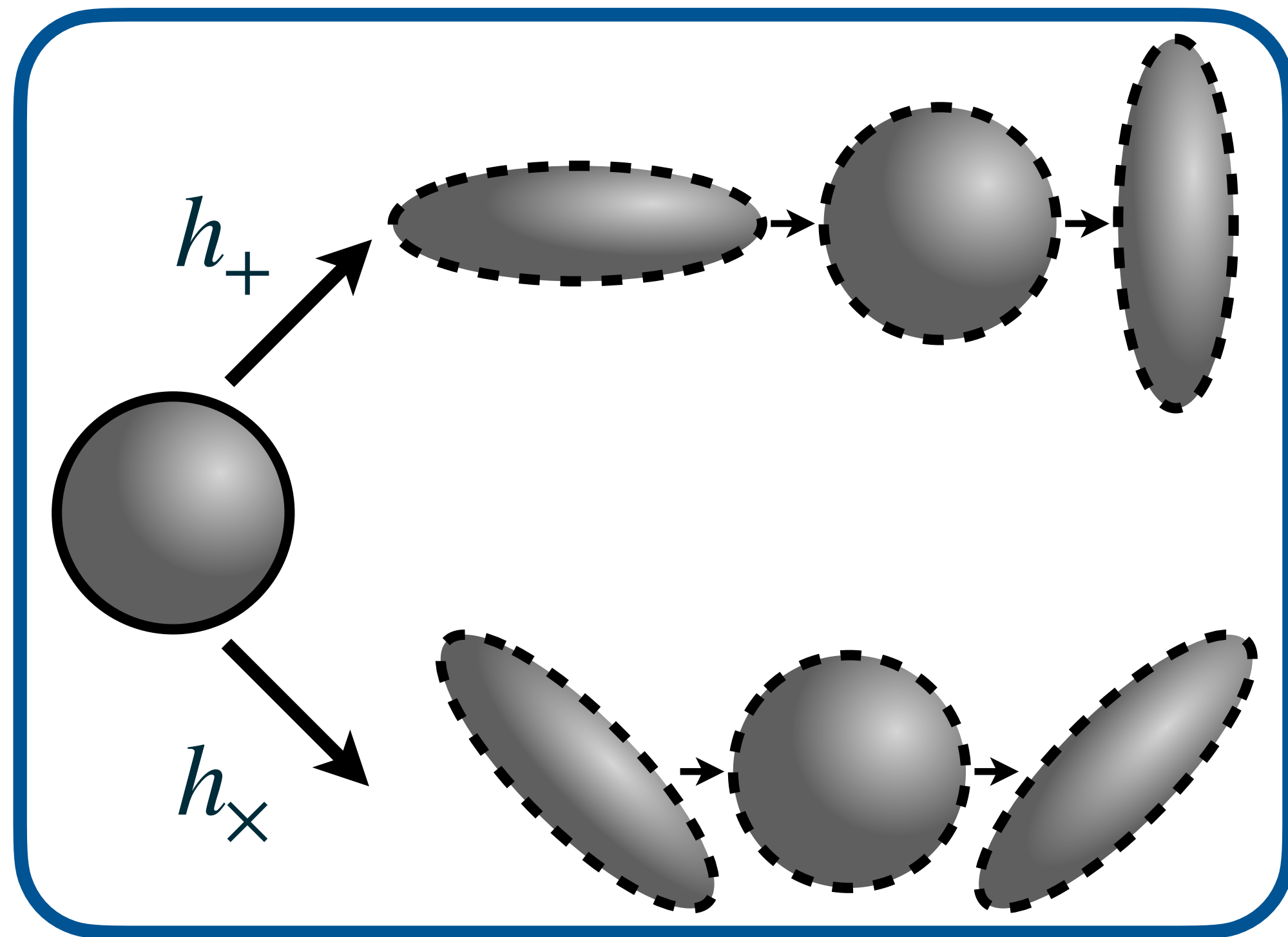
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TT frame intuition

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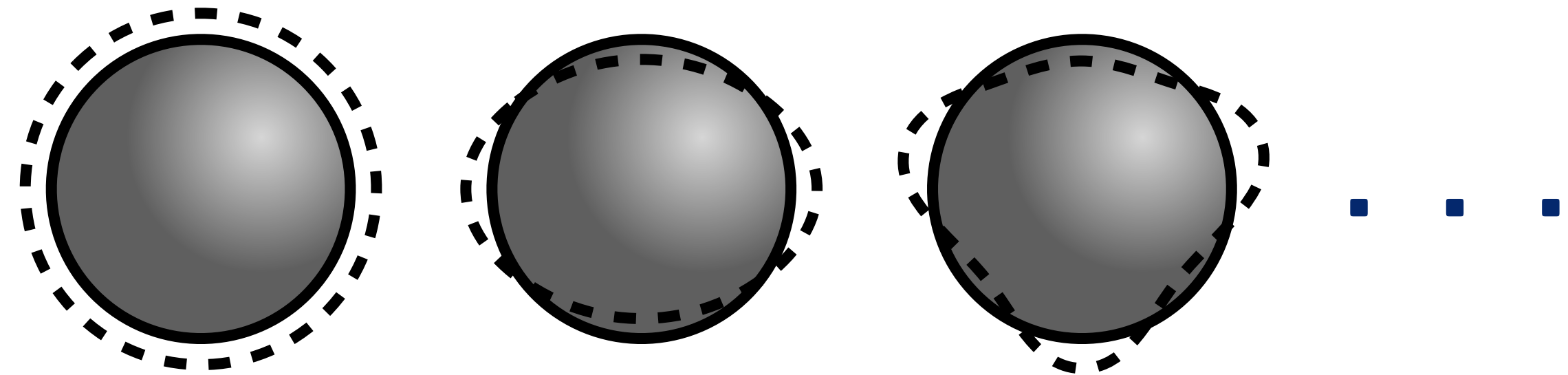


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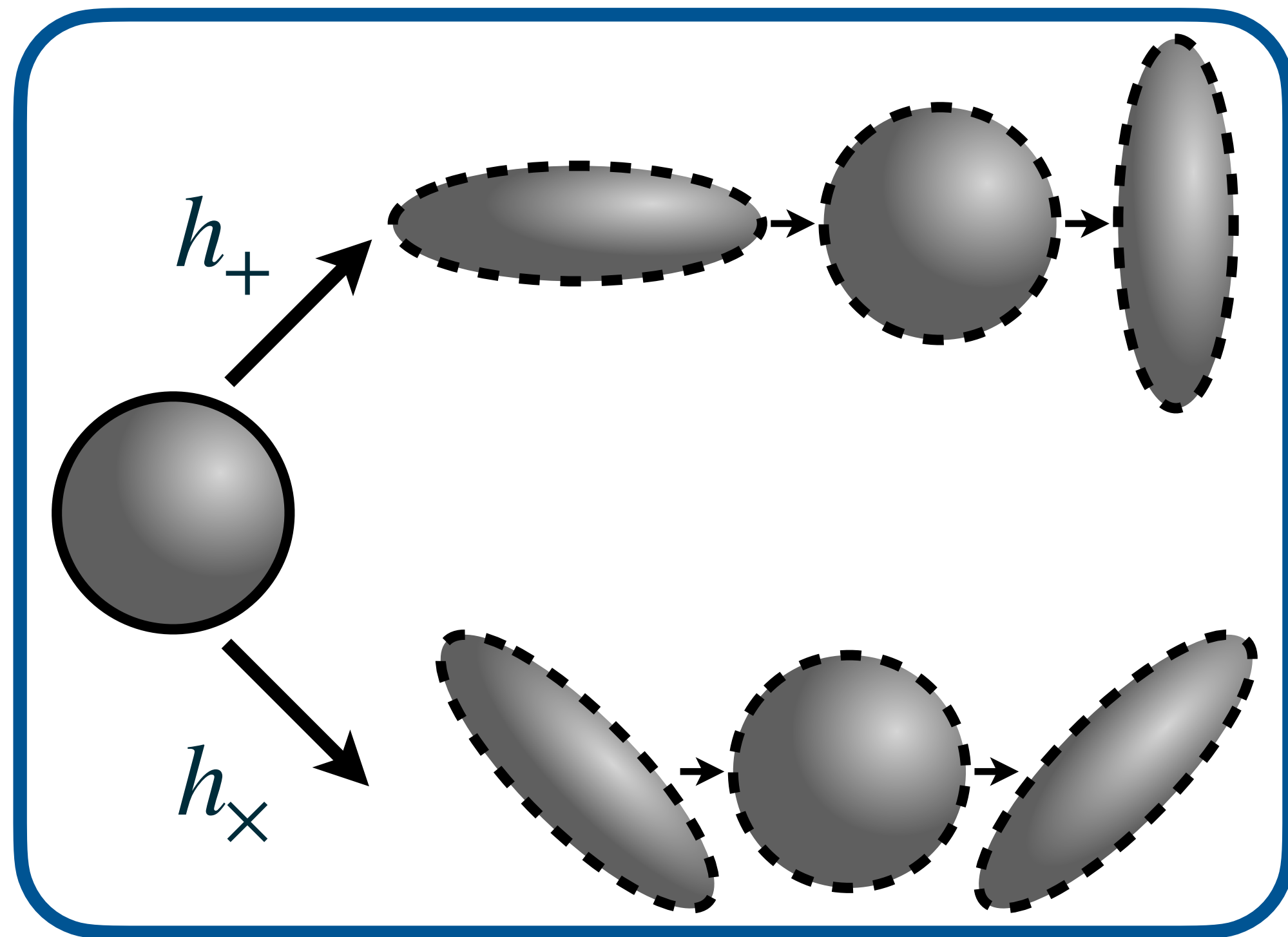
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Spheroidal



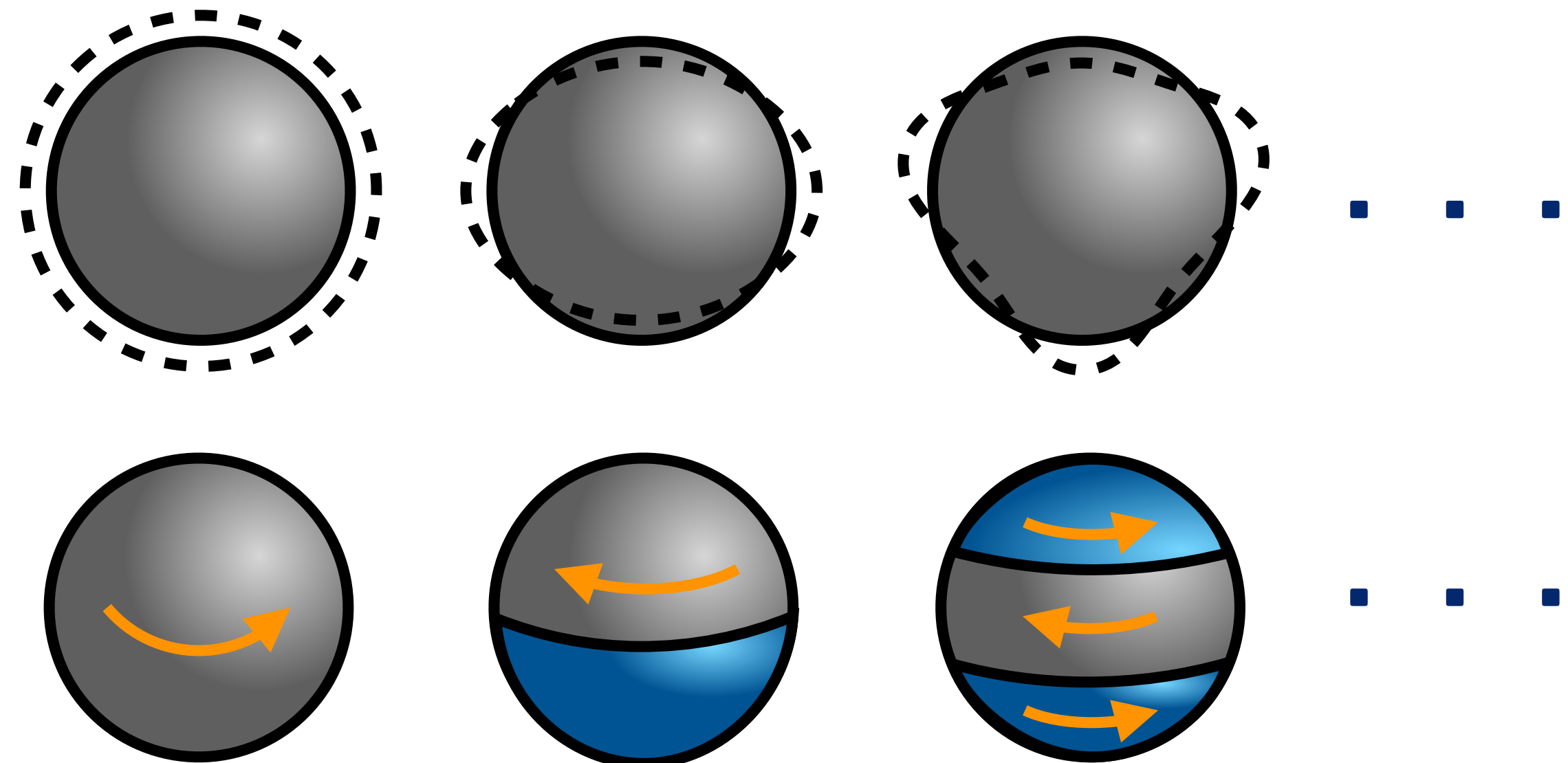
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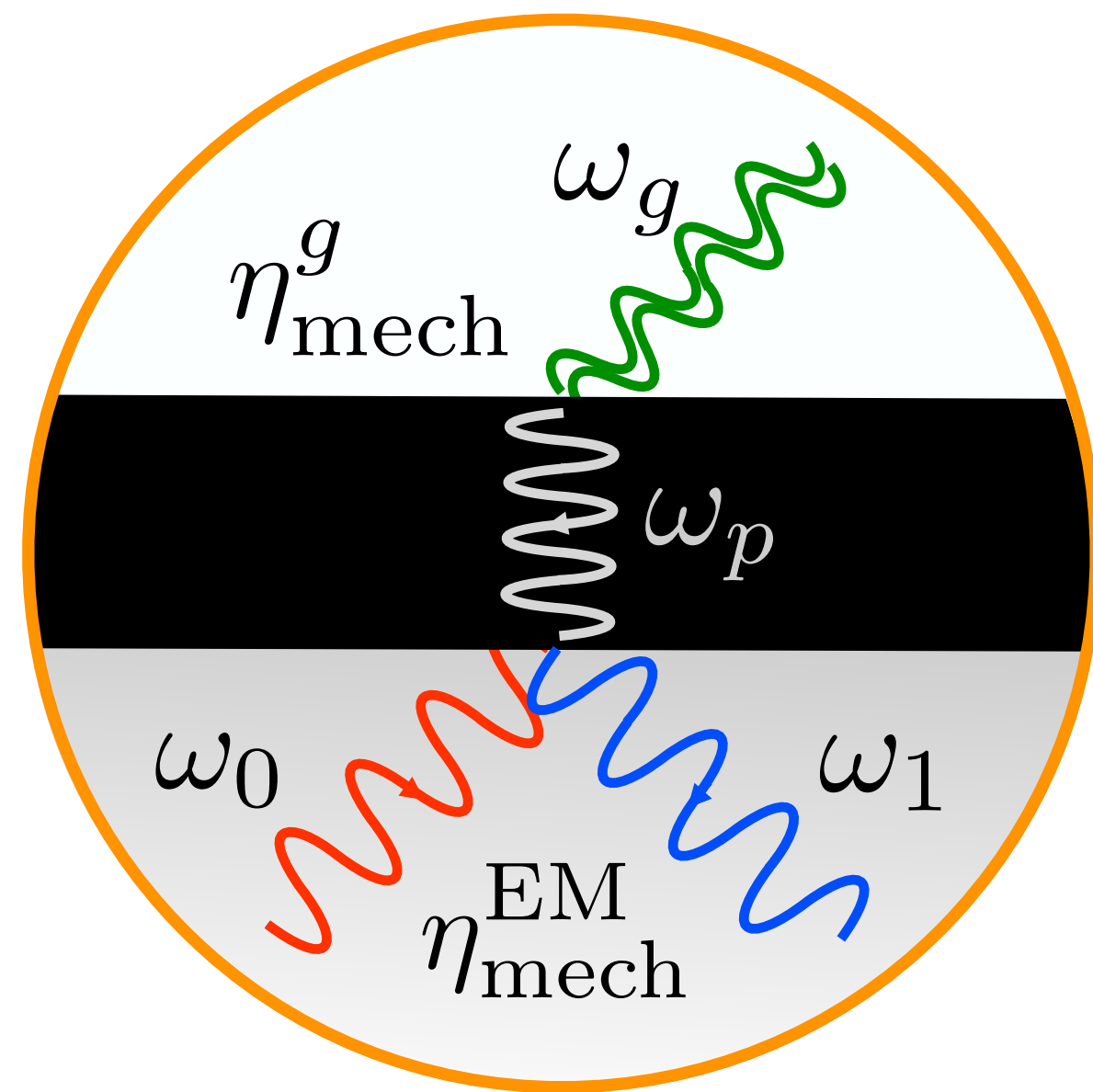
Mechanical modes of a sphere

$$U_{lmn} = \underbrace{\nabla\phi_L + i\nabla \times \mathbf{L}\phi_{T_1}}_{\text{Spheroidal}} + \underbrace{i\mathbf{L}\phi_{T_2}}_{\text{Toroidal}}.$$



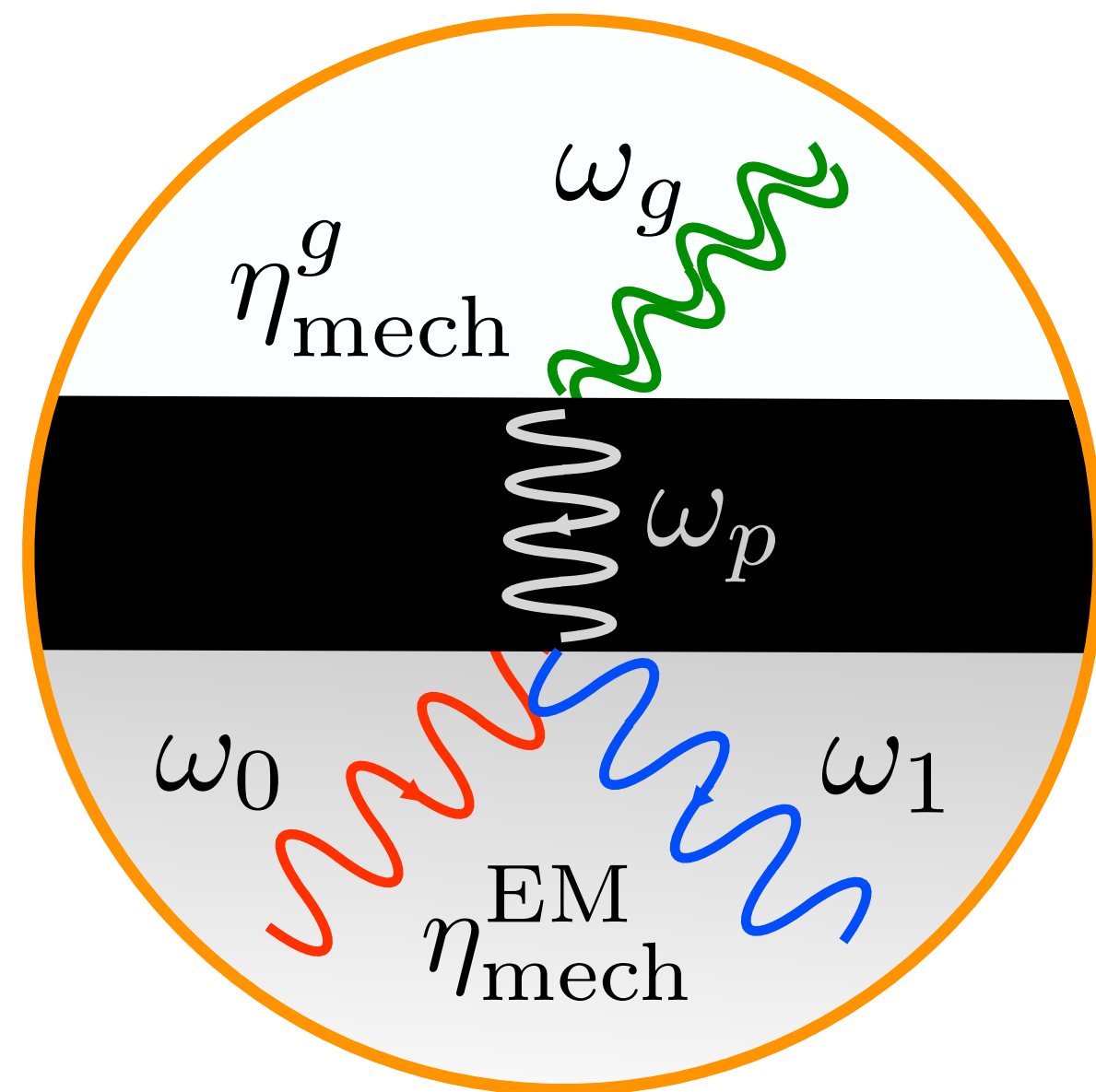
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Vibrations of cavity walls — modification of boundary conditions



MAGO 2.0

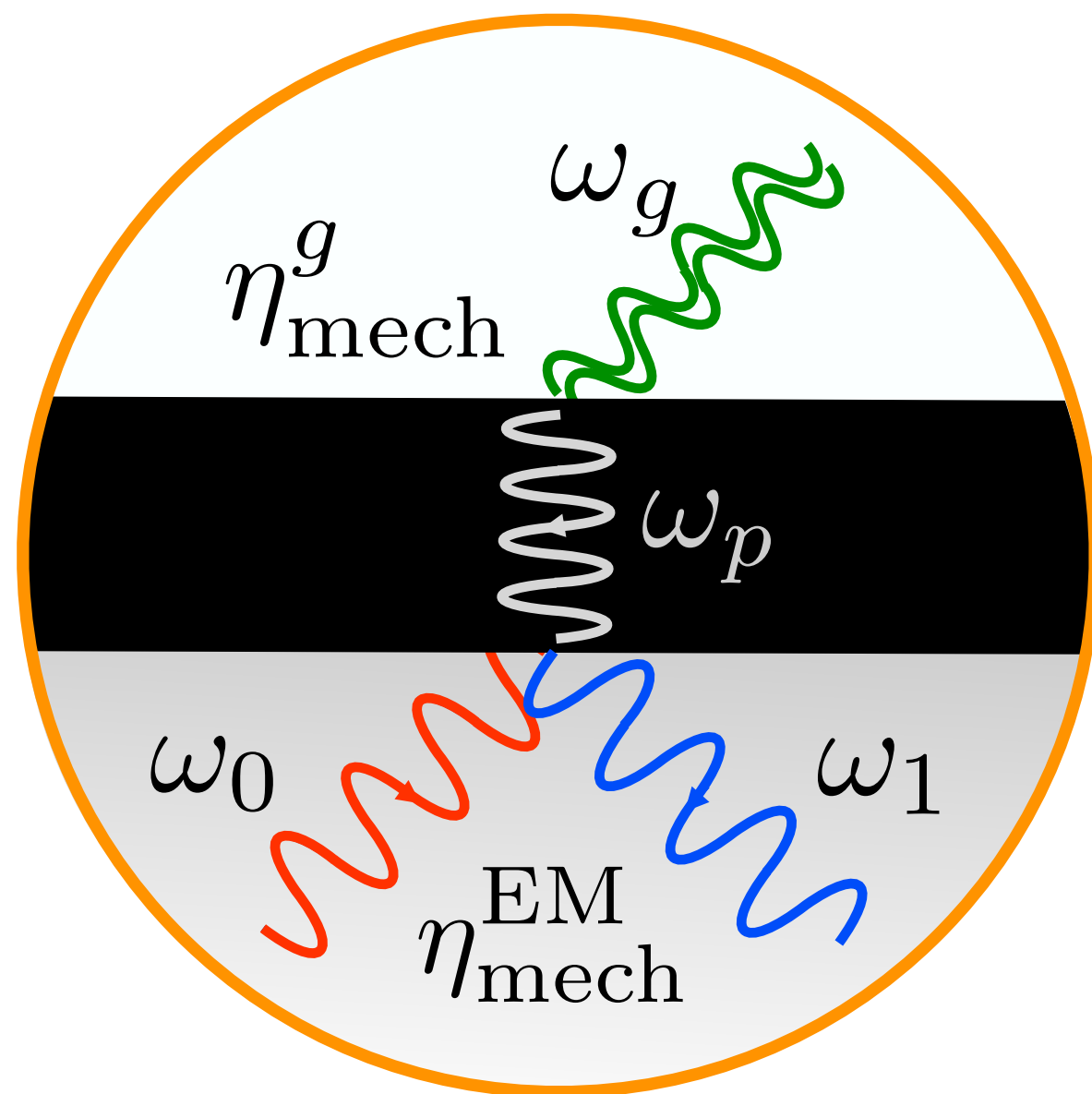
Vibrations of cavity walls — modification of boundary conditions



$$\eta_{\text{mech}}^{\text{EM}} = V_{\text{cav}}^{1/3} \frac{\int_{S_0} d\mathbf{A} \cdot \mathbf{U}_p (\mathbf{E}_0 \cdot \mathbf{E}_1^* - \mathbf{B}_0 \cdot \mathbf{B}_1^*)}{\int_{V_{\text{cav}}} d^3\mathbf{x} |\mathbf{E}_1|^2}$$

MAGO 2.0

Vibrations of cavity walls — modification of boundary conditions



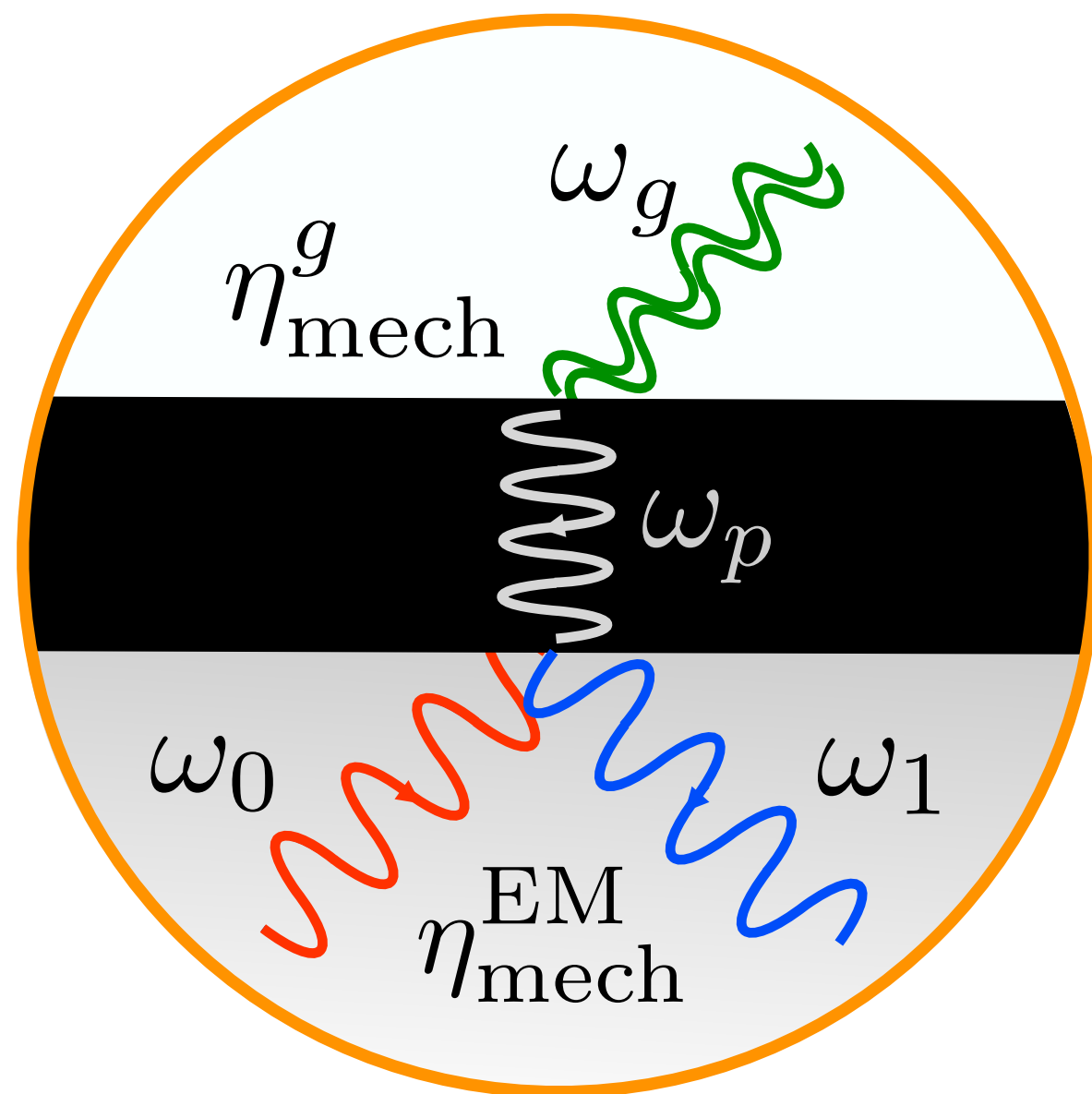
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$$\ddot{e}_1 + \frac{\omega_1}{Q_1} \dot{e}_1 + \omega_1^2 e_1 \simeq -2 \eta_{\text{mech}}^{\text{EM}} \omega_1^2 V_{\text{cav}}^{-1/3} u_p e_0$$

$$\mathbf{E}(\mathbf{x}, t) = e_i(t) \mathbf{E}_i(\mathbf{x})$$

MAGO 2.0

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Recall effect of GW force

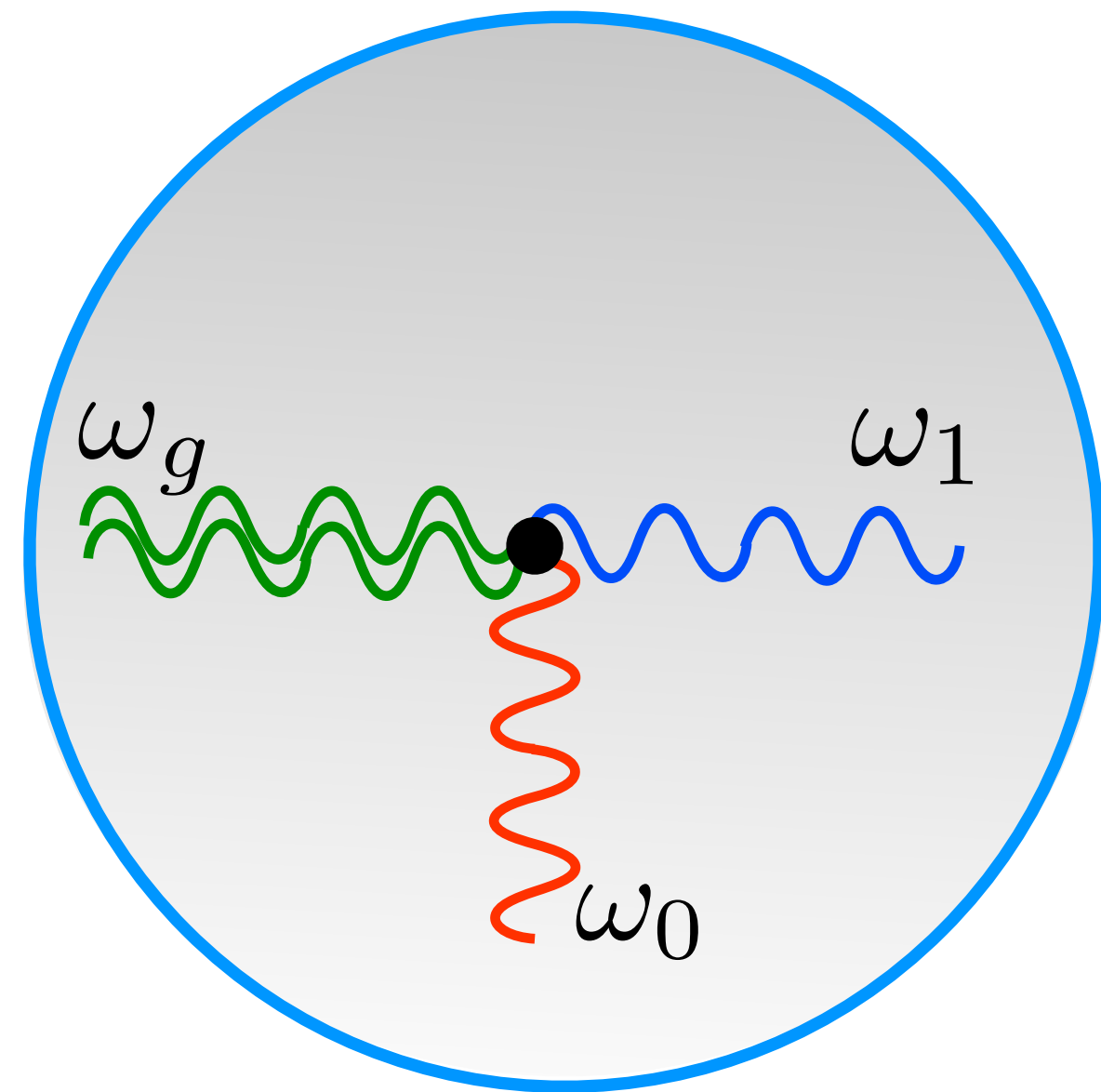
$$u_p \propto h_0$$

$$\ddot{e}_1 + \frac{\omega_1}{Q_1} \dot{e}_1 + \omega_1^2 e_1 \simeq -2 \eta_{\text{mech}}^{\text{EM}} \omega_1^2 V_{\text{cav}}^{-1/3} u_p e_0$$

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MAGO 2.0

Inverse Gertsenshtein effect — direct coupling to EM waves



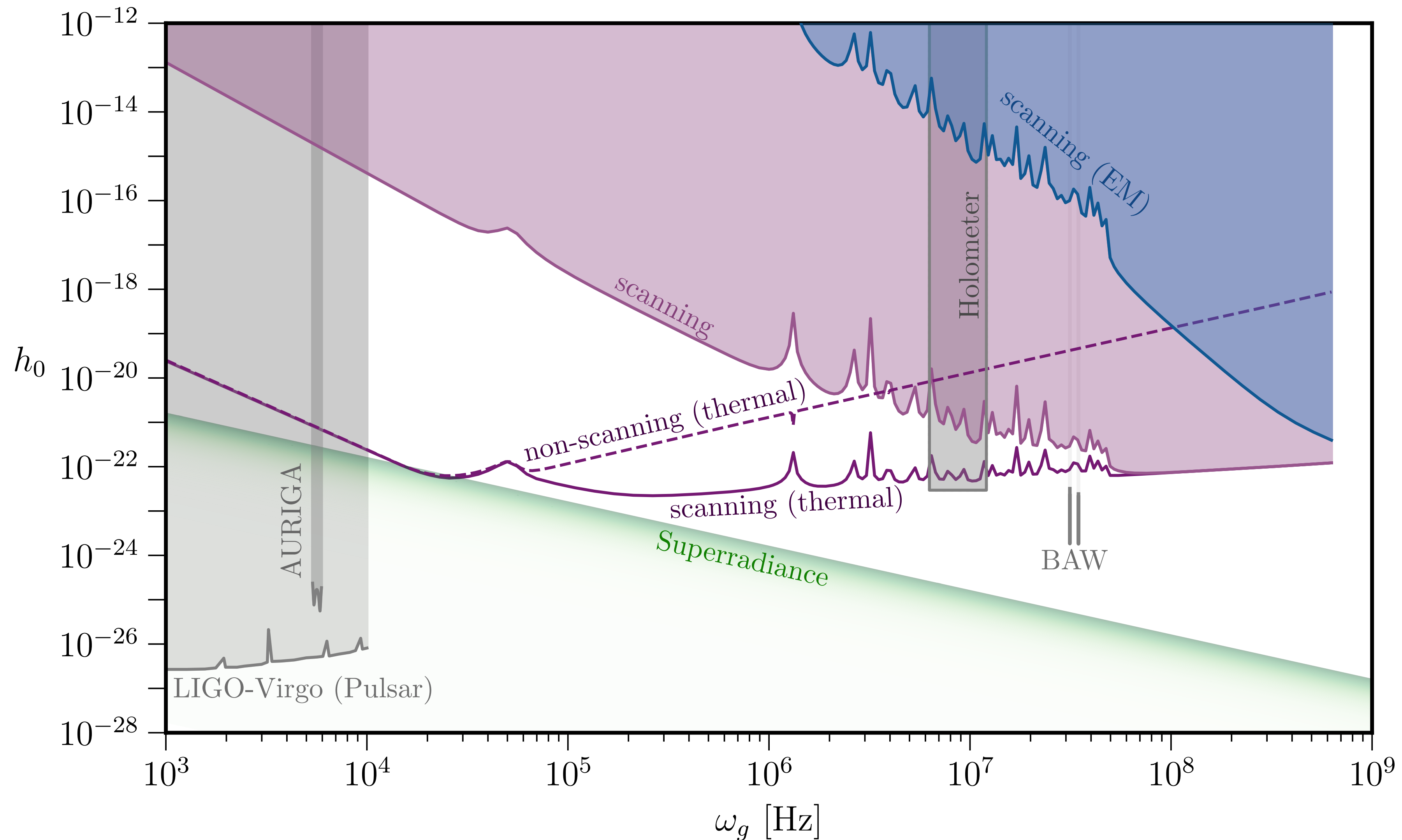
$$j_{\text{eff}}^{\mu} \equiv \partial_{\nu} \left(\frac{1}{2} h F^{\mu\nu} + h^{\nu}_{\alpha} F^{\alpha\mu} - h^{\mu}_{\alpha} F^{\alpha\nu} \right)$$

$$P_{\text{sig}} \sim \frac{Q_1^2 Q_{\text{int}}}{Q_{\text{cpl}}} |\eta_{\text{EM}}^g|^2 P_{\text{in}} h_0^2 (\omega_g V_{\text{cav}}^{1/3})^4$$

Recall discussion of rigid ruler

$$\eta_{\text{EM}}^g = \frac{\int d^3 \mathbf{x} \mathbf{E}_1^* \cdot \hat{\mathbf{j}}}{(V_{\text{cav}} \int d^3 \mathbf{x} |\mathbf{E}_1|^2)^{1/2}}$$

MAGO 2.0 sensitivity to coherent GWs



Optimal Scanning

