# Understanding the effect of strangeness and electric charge on the NCQ scaling of directed flow 

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$\square$ Introduction and Motivation
■A Multi-Phase Trasport Model (AMPT)
$\square$ Analysis Details
Bndependent set of equations

- Analytical solutions
$\square$ Testing the solution using Default-AMPT
『Summary


## Introduction

$$
\begin{gathered}
E \frac{d^{3} N}{d p^{3}}=\frac{1}{2 \pi} \frac{d^{2} N}{p_{T} d p_{T} d y}\left(1+2 \mathbf{v}_{\mathbf{1}} \cos \left(\phi-\psi_{R}\right)+2 \mathbf{v}_{\mathbf{2}} \cos 2\left(\phi-\psi_{R}\right)+\ldots\right) \\
\left.v_{1}=\left\langle\cos \left(\phi-\psi_{R}\right)\right\rangle \quad \phi=\tan ^{-1}\left(\frac{p_{y}}{p_{x}}\right) \right\rvert\,
\end{gathered}
$$

$\sqrt{ }$ Directed flow ( $\mathbf{v}_{\mathbf{1}}$ ): $1^{\text {st }}$ harmonic in the Fourier expansion of the
 azimuthal distribution of emitted particles relative to the reaction plane

- A collective sideward motion whose dominant component is an odd function of the particle rapidity (y)
$\sqrt{ }$ Probes early stage of collision dynamics
- Features a strong magnetic field $\left(\sim 10^{14}-10^{15} \mathrm{~T}\right)$ dominated by the passing spectator protons



## Motivation

$\checkmark$ Quarks have velocity perpendicular to B direction, because of strong longitudinal flow velocity
$\checkmark$ Lorentz force results in an electric current along x -axis $\rightarrow$ Hall effect
$\checkmark$ B decreases $\rightarrow$ Faraday current is induced in the direction opposite to Hall effect
$\checkmark$ Coulomb force exerted by the outgoing positively charged spectator on the produced plasma
$\checkmark$ Net electric current is sum of Faraday, Lorentz and Coulomb force resulting in a charge dependent $\mathbf{v}_{1}$

## Does $\mathrm{v}_{1}$-splitting $\left(\Delta \mathrm{v}_{1}\right)$ influence by the EM field (charge)?

## Does it depend on the strangeness?

$\checkmark$ A new method of testing of NCQ scaling using $\mathrm{v}_{1}$ of produced hadrons ( $K^{-}, \bar{p}, \bar{\Lambda}, \phi, \bar{\Xi}^{+}, \Omega^{-}, \bar{\Omega}^{+}$) in the same rapidity and $p_{\mathrm{T}} / \mathrm{N}_{\mathrm{cq}}$ phase space
$\checkmark$ Need to understand the dependence of change in electric charge, strangeness content, choice of equations and collision beam energies


## Default-AMPT Model

$\checkmark$ AMPT model provides a comprehensive kinetic description of essential stages of high energy heavy ion collisions


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## Analysis Details

$\sqrt{ } v_{1}$ of a hadron (H) and its constituent quarks satisfying NCQ scaling:

$$
\begin{aligned}
& v_{1, \mathrm{H}}\left(p_{\mathrm{T}, \mathrm{H}}\right)=\sum_{i} v_{1, i}\left(p_{\mathrm{T}}, i\right) \\
& v_{1, \mathrm{H}}\left(N_{c q} p_{\mathrm{T}, \mathrm{q}}\right)=N_{c q} v_{1, q}\left(p_{\mathrm{T}, \mathrm{q}}\right)
\end{aligned}
$$

Same quark masses
$\checkmark \Delta v_{1}$ values are obtained for two combinations:
Identical $(\Delta m=0, \Delta q=0$ and $\Delta S=0)$

- Non-identical $(\Delta m=0, \Delta q \neq 0$ and $\Delta S \neq 0)$
$\sqrt{ }$ Produced and transported quarks have different $\mathrm{v}_{1}$
$\boldsymbol{m}$ Interpretation is challenging when hadrons having transported quarks are included

For NCQ scaling testing, assumptions are:

- $v_{1}(\bar{u}) \approx v_{1}(\bar{d})$;
- $v_{1}(s) \approx v_{1}(\bar{s})$

$\sqrt{ }$ The same number of light, strange quarks and anti-quarks on Left \& Right hand side $(R)$ with proper weighting factors

$$
\frac{1}{2} v_{1}[\phi(s \bar{s})]=\frac{1}{3}\left[\Omega^{-}(s s s)\right]
$$

$$
\begin{aligned}
& \Delta S=S^{L}-S^{R}=1 \\
& \Delta q=q^{L}-q^{R}=1 / 3 \\
& \Delta q_{u d}=q^{L}-q^{R}=0
\end{aligned}
$$

Produced hadrons


|  | Set | Charge ( $\Delta q$ ) | $\Delta q_{u d}$ | Strangness ( $\Delta S$ ) | Expression |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1 | 0 | 0 | 0 | $v_{1}\left[K^{-}(\bar{u} s)\right]+v_{1}[\bar{\Lambda}(\bar{u} \bar{d} \bar{s})]=v_{1}[\bar{p}(\bar{u} \bar{u} \bar{d})]+v_{1}[\phi(s \bar{s})]$ |
| Identical | 2 | 0 | 0 | 0 | $v_{1}[\bar{\Lambda}(\bar{u} \bar{d} \bar{s})]=\frac{v_{1}}{2}\left[\bar{\Xi}^{+}(\bar{d} \bar{s} \bar{s})\right]+\frac{v_{1}}{2}[\bar{p}(\bar{u} \bar{u} \bar{d})]$ |
|  | 3 | 0 | 0 | 0 | $\frac{v_{1}}{3}\left[\Omega^{-}(s s s)\right]+\frac{v_{1}}{3}\left[\bar{\Omega}^{+}(\bar{s} \bar{s} \bar{s})\right]=v_{1}[\phi(s \bar{s})]$ |
|  | 4 | 1/3 | 0 | 1 | $\frac{1}{2} v_{1}[\phi(s \bar{s})]=\frac{1}{3}\left[\Omega^{-}(s s s)\right]$ |
| Non-identical | 5A | 2/3 | 1/3 | 1 | $\left.\frac{v_{1}}{2} \phi(s \bar{s})\right]+\frac{v_{1}}{3}[\bar{p}(\bar{u} \bar{u} \bar{d})]=v_{1}\left[K^{-}(\bar{u} s)\right]$ |
|  | 5B | 2/3 | 1/3 | 1 | $v_{1}[\bar{\Lambda}(\bar{u} \bar{d} \bar{s})]=\frac{v_{1}}{2}[\phi(s \bar{s})]+\frac{2}{3} v_{1}[\bar{p}(\bar{u} \bar{u} \bar{d})]$ |

## Solution for Independent Equations

$$
\begin{array}{cc}
\Delta N_{i}=N_{i}^{L}-N_{i}^{R} & \Delta S=S^{L}-S^{R} \\
\Delta N_{\bar{u}}+\Delta N_{d}=0 & \Delta q=q^{L}-q^{R} \\
\Delta N_{s}+\Delta N_{\bar{s}}=0 & \Delta q_{u d}=q^{L}-q^{R}
\end{array}
$$

- $\Delta q_{u d}=\Delta N_{\bar{d}} / 3-2 \Delta N_{\bar{u}} / 3=\Delta N_{\bar{d}}$
- $\Delta S=2 \Delta N_{\bar{s}}$
$\Delta v_{1}=\Delta q_{u d}\left(v_{1, \bar{d}}-v_{1, \bar{u}}\right)+\Delta S\left(\frac{v_{1, \bar{s}}-v_{1, s}}{2}\right) \Delta v_{1}=c_{q} \Delta q_{u d}+c_{s} \Delta S+c_{0}$
- $c_{s}$ reflects $v_{1}$ difference between $s$ and $\bar{s}$ quarks
- Similarly, $c_{\mathrm{q}}$ is the coefficient of $\Delta q \& c_{0}$ is the intercept

$$
\Delta v_{1}=\Delta q\left(v_{1, \bar{d}}-v_{1, \bar{u}}\right)+\Delta S\left(\frac{v_{1, \bar{s}}-v_{1, s}}{2}-\frac{v_{1, \bar{d}}-v_{1, \bar{u}}}{3}\right)
$$

$$
\Delta v_{1}=c_{q}^{*} \Delta q+c_{s}^{*} \Delta S+c_{0}^{*}
$$

## Analytical solution

$\sqrt{ }$ Let's us take the average of Set 1-3 and denote it as Set A, Similarly

- Average of Set 1-3: Set A
- Set 4: Set B
- Set 5B: Set C
$\sqrt{ }$ By considering these three sets (A, B, C) one can obtain three coefficients from Eq. 1 or Eq. 2 $\checkmark$ The $\Delta v_{1}$-slope also satisfy the same relation as $\left(\Delta v_{I}\right)_{A},\left(\Delta v_{1}\right)_{B}$ and $\left(\Delta v_{1}\right)_{C}$

$\checkmark$ Similarly using Eq. 2: $\cdot c_{0}{ }^{*}=(\Delta v)_{\mathrm{A}}=c_{0}$
- $c_{q}^{*}=3\left[\left(\Delta v_{1}\right)_{C}-\left(\Delta v_{1}\right)_{B}\right]=c_{q}$
- $c_{s}^{*}=2\left(\Delta v_{1}\right)_{B}-\left(\Delta v_{1}\right)_{A}-\left(\Delta v_{1}\right)_{C}=c_{s}-c_{q} / 3$


## Rapidity dependence of $\mathrm{v}_{1}$ and $\Delta \mathrm{v}_{1}$


$\checkmark \Delta q=0$ and $\Delta S=0($ Set 2$)$ in Au+Au collisions at $\sqrt{S N}=7.7,14.5,27,54.4$ and 200 GeV
$\checkmark$ Calculated in the same $p_{\text {T }}$ region: Baryon $0-2 \mathrm{GeV}$ and Meson $0-3 \mathrm{GeV}$

$\checkmark \Delta \mathrm{v}_{1}$-slope is sensitive to the $\Delta \mathrm{q}, \Delta \mathrm{S}$, choice of equations and collision beam energies
$\checkmark$ A non-zero $\Delta v_{1}$-slope is found in $\mathrm{Au}+\mathrm{Au}$ collisions at $\sqrt{S} \mathrm{SN}=7.7,14.5,27,54.4$ and $200 \mathrm{GeV}, 10-50 \%$ centrality in Default-AMPT

## 3D plane fitting of $\Delta v_{1}$-slope

Testing independent equations
Testing analytical solution

$\sqrt{ }$ Equation of the plane (surface):

$$
d\left(\Delta v_{1}\right) / \mathrm{dy}=c_{q} \Delta q_{u d}+c_{s} \Delta S+c_{0}
$$



$$
\begin{gathered}
\Delta v_{1}=\Delta q_{u d}\left(v_{1, \bar{d}}-v_{1, \bar{u}}\right)+\Delta S\left(\frac{v_{1, \bar{s}}-v_{1, s}}{2}\right)+c_{0} \\
c_{\mathrm{q}}
\end{gathered} c_{S}
$$

$$
\begin{gathered}
\Delta v_{1}=\Delta q\left(v_{1, \bar{d}}-v_{1, \bar{u}}\right)+\Delta S\left(\frac{v_{1, \bar{s}}-v_{1, s}}{2}-\frac{v_{1, \bar{d}}-v_{1, \bar{u}}}{3}\right)+c_{0}^{*} \\
c_{\mathrm{q}}^{*}
\end{gathered} c_{\mathrm{s}}^{*}
$$

$\checkmark$ The $c_{\mathrm{s}}$ and $c_{\mathrm{s}}{ }^{*}$ coefficients are different i.e more sensitive unlike $c_{0}\left(c_{0}{ }^{*}\right)$ and $c_{\mathrm{q}}\left(c_{\mathrm{q}}{ }^{*}\right)$


$$
\begin{gathered}
\Delta v_{1}=\Delta q_{u d}\left(v_{1, \bar{d}}-v_{1, \bar{u}}\right)+\Delta S\left(\frac{v_{1, \bar{s}}-v_{1, s}}{2}\right)+c_{0} \\
c_{\mathrm{q}}
\end{gathered}
$$

$\checkmark$ The analytical solution is in good agreement with the solution obtained using 5 independent equations

## Linear Coefficients



The coefficients (*) are also sensitive to the choice of independent set of equations

$$
\begin{gathered}
\Delta v_{1}=\Delta q\left(v_{1, \bar{d}}-v_{1, \bar{u}}\right)+\Delta S\left(\frac{v_{1, \bar{s}}-v_{1, s}}{2}-\frac{v_{1, \bar{d}}-v_{1, \bar{u}}}{3}\right)+c_{0}^{*} \\
c_{q}^{*} \\
c_{\mathrm{s}}^{*}
\end{gathered}
$$

## Summary

- A new method of testing NCQ scaling of $\mathrm{v}_{1}$ using produced hadrons is proposed
$\star$ Considering 7 produced hadrons only 5 independent set of equations are possible
The testing should at least be performed by two combinations of independent set of equation
$\Delta v_{1}$-slope is very sensitive to the change in electric charge, strangeness content, choice of equations and collision beam energies.
$\star$ The linear coefficient of $\Delta \mathrm{S}$ i.e. $c_{S}$ is found to be the most sensitive parameter for $\mathrm{v}_{1}$-splitting or testing the NCQ scaling as compared to the corresponding $\Delta q_{u d}(\Delta q)$ coefficient $c_{s}\left(c_{s}{ }^{*}\right)$

The non-zero $\Delta \mathrm{v}_{1}$ for non-identical set of equations suggests that $\mathrm{v}_{1}$-splitting not may only be driven by electromagnetic effect $\rightarrow$ Strangeness of hadron might play important role


## Backup


[^0]:    Zi-Wei Lin et al., Phys. Rev. C 72, 064901 (2005);

