Simulating heavy-ion collisions at BES energies

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52nd International Symposium on Multiparticle Dynamics, Gyöngyös, Hungary Based on Phys. Rev. C 106 (2022) 5, 054906

Beam Energy Scan



Standard model of HIC needs to be extended to finite baryon densities. In this work we

- Present a new LEXUS based initial state model
- Use a new crossover equation of state
- Calculate departure functions using a quasiparticle model at finite baryon densities
- Use these new developments to do a baseline calculation for observables at RHIC BES energies

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LEXUS based initial conditions

- Original LEXUS (Linear EXtrapolation of Ultrarelativistic Scattering) had no free parameters S. Jeon and J. Kapusta, Phys. Rev. C. 56 (1997)
- LEXUS linearly extrapolated nucleon-nucleon scattering data to nucleus nucleus collisions. Goal was to provide a baseline which would work if there was no new physics beyond nucleon scattering
- Colliding nucleons are assumed to propagate in straight lines colliding with nuclei that cross its path
- Inelastic collisions produced mesons which were assumed to not further interact

• Probability of momentum loss in a collision is parametrized as

$$P(y_{loss}) = \lambda \frac{\cosh(2y_0 - y_{loss})}{\sinh(2y_0) - \sinh(y_0)}$$

• Rapidity width of produced mesons is given by the Landau model

$$\sigma_L = \frac{8}{3} \frac{c_s^2}{1 - c_s^4} \log\left(\frac{\sqrt{s}}{2m_N}\right) \approx \log\left(\frac{\sqrt{s}}{2m_N}\right)$$

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LEXUS inspired 3D initial state

- Nucleon sampling \rightarrow Lorentz contraction \rightarrow Binary collisions
- Nucleons are sampled from a Woods-Saxon distribution and Lorentz contracted in the beam direction
- Nucleons propagate in straight lines in beam direction
- Two nucleons from opposing nuclei collide with probability λ when their cylindrical trajectories overlap. Area of the cylinder is given by the total collision cross-section
- Hydrodynamic solver

LEXUS inspired 3D initial state



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• Perturbative QCD EOS in QGP phase is matched to a hadron resonance gas EOS at low temperatures M. Albright, J. Kapusta and C. Young, Phys. Rev. C 90 (2014)

$$\textit{P}(\textit{t},\mu) = \textit{S}(\textit{T},\mu)\textit{P}_{\text{QGP}}(\textit{T},\mu) + [1-\textit{S}(\textit{T},\mu)]\textit{P}_{\text{HRG}}(\textit{T},\mu)$$



It shows good match with Lattice EOS at $\mu = 0 \dots$



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... and with Lattice EOS calculated at non-zero μ using Taylor expansion...



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... and could be extended beyond.



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• Departure functions account for non-equilibrium corrections at freezeout

$$f_{a} = f_{a}^{eq} (1 + \phi_{a})$$

$$\phi_{a} = -A_{a} \partial_{\rho} u^{\rho} - B_{a} \rho_{a}^{\nu} D_{\nu} \left(\frac{\mu_{B}}{T}\right) + C_{a} \rho_{a}^{\mu} \rho_{a}^{\nu} \left(D_{\mu} u_{\nu} + D_{\nu} u_{\mu} + \frac{2}{3} \Delta_{\mu\nu} \partial_{\rho} u^{\rho}\right)$$

- Numerical studies using linear sigma model suggests that the relaxation time $\tau(E_a) \propto E_a \equiv \tau' E_a$ P. Chakraborty and J. Kapusta, Phys. Rev. C 95 (2017)
- For shear correction, we get $C_a = rac{ au'}{2 au}$ with $au' = rac{\eta}{2 au w}$

Thermal conductivity

• Unlike shear viscosity, the expression for thermal conductivity associated with baryon current is not so simple

$$B_a = \tau' \left(b_a - \frac{n_B}{w} E_a \right) - b(T, \mu_B)$$

where

$$b = \frac{\tau' T}{w^2} \Big[Ts(T\chi_{T\mu} + \mu_B \chi_{\mu\mu}) - Tn_B(T\chi_{TT} + \mu_B \chi_{T\mu}) - n_B w \Big]$$

and we have susceptibilities

$$\chi_{xy} = \frac{\partial^2 P(T,\mu)}{\partial x \partial y}$$

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- No overall normalization used
- No initial flow. Initial flow will necessitate use of higher η/s which will lead to more entropy production
- No bulk used here. That will also contribute to multiplicity
- No hadronic cascade. Pion winds shift proton spectra to higher p_T

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Rapidity spectra







- We present our results for simulations of HICs at non-zero baryon densities. A new baseline calculation for RHIC-BES energies using a LEXUS based initial state, a crossover EOS without critical point and departure functions at finite baryon densities
- A LEXUS inspired 3D initial state model was used
- A crossover EOS from matching HRG and pQCD results was used at finite baryon densities
- Departure functions are calculated at finite baryon densities
- Get reasonable matching to data without an overall normalization

This work was supported by the U.S. DOE Grant No. DE-FG02-87ER40328.