

INITIAL STAGES AND QGP ANISOTROPY CONSTRAINED THROUGH HIGH- p_{\perp} DATA

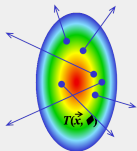
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QGP TOMOGRAPHY

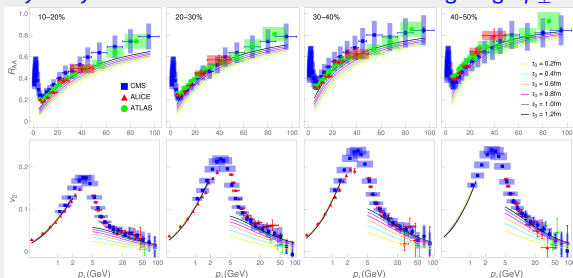
- **Main goal:** use high- p_{\perp} data to infer bulk properties of QGP.
- **Dynamical Radiative and Elastic ENergy Loss Approach:** our numerical framework



- High-energy particles lose energy when they traverse QGP.
- This energy loss is sensitive to QGP properties.
- **We can realistically predict this energy loss.**



- High- p_{\perp} probes are excellent tomography tools.
- We can use them to infer some of the bulk QGP properties.
- **Example: hydrodynamics initial time constrained through high- p_{\perp} data:**



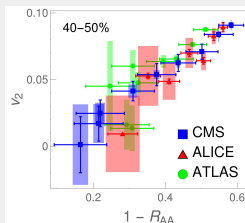
ANISOTROPY

- Initial spatial anisotropy: major limiting factor for QGP tomography.
- Still not possible to infer it from experimental data.
- We propose a novel approach: inference from high- p_{\perp} R_{AA} and v_2 measurements.
- Previous study: $v_2/(1 - R_{AA})$ saturates at high- p_{\perp} , and the saturation value reflects the geometry of the system

M. Djordjevic, S. Stojku, M. Djordjevic and P. Huovinen, Phys.Rev. C Rapid Commun. **100**, 031901 (2019).

- This study employed analytic considerations and a simple 1+1D medium expansion
- Here: we study the behavior of $v_2/(1 - R_{AA})$ in a system that expands in both longitudinal and transversal directions.

S. Stojku, J. Auvinen, L. Zivkovic, P. Huovinen, M. Djordjevic, Physics Letters B **835**, 137501 (2022)



- v_2 and $1 - R_{AA}$ are directly proportional at high p_{\perp} .
- Equivalent to a p_{\perp} -independent ratio of v_2 and $1 - R_{AA}$.
- Can fluid dynamical calculations reproduce such proportionality?
Can we relate this observation to the anisotropy of the system?

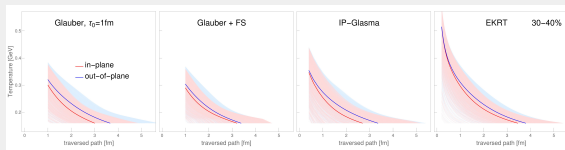
ANISOTROPY

- **DREENA-A: can accommodate any temperature profile and generate high- p_{\perp} R_{AA} and v_2 predictions.**

D. Zigic, I. Salom, J. Auvinen, P. Huovinen and M. Djordjevic, *Front. Phys.* **10**:957019 (2022)

- We visualize the temperatures partons experience in the **in-plane** and **out-of-plane** directions for different initializations and evolutions.

Stefan Stojku, Jussi Auvinen, Lidija Zivkovic, Pasi Huovinen, Magdalena Djordjevic, *Physics Letters B* **835**, 137501 (2022)



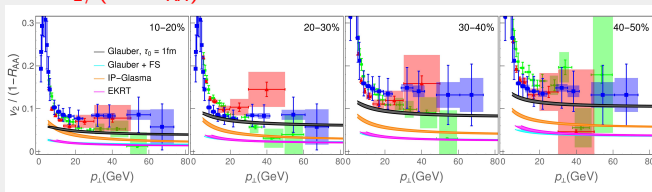
$$\langle T_x(t) \rangle = \frac{1}{N} \sum_{i=1}^N T(x_i + t, y_i, t)$$

$$\langle T_y(t) \rangle = \frac{1}{N} \sum_{i=1}^N T(x_i, y_i + t, t)$$

$v_2/(1 - R_{AA})$ RESULTS

- Does $v_2/(1 - R_{AA})$ saturate?
- Does this saturation carry information on the anisotropy of the system?
- What kind of anisotropy measure is revealed through high- p_{\perp} data?

We calculate $v_2/(1 - R_{AA})$ within DREENA-A framework:



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The phenomenon of $v_2/(1 - R_{AA})$ saturation is robust!

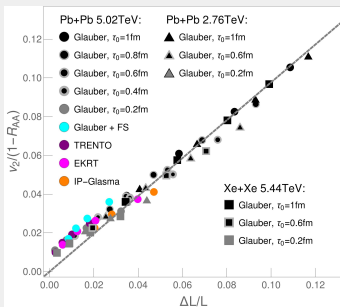
How to explore if it contains information on the system anisotropy?

CONNECTION TO ANISOTROPY

■ Next: Plot charged hadrons' $v_2/(1 - R_{AA})[100\text{GeV}]$ vs. $\Delta L/\langle L \rangle$

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- Centrality classes: 10-20%, 20-30%, 30-40%, 40-50%

- Surprisingly simple relation between $v_2/(1 - R_{AA})$ and $\Delta L/\langle L \rangle$.

- Slope ≈ 1 .

- $v_2/(1 - R_{AA})$ carries information on the system anisotropy, through $\Delta L/\langle L \rangle$.

JET-PERCEIVED ANISOTROPY

- Define a more direct measure of anisotropy? Explicit dependence on time evolution?
- We define jT :

$$jT(\tau, \phi) \equiv \frac{\int dx dy T^3(x + \tau \cos \phi, y + \tau \sin \phi, \tau) n_o(x, y)}{\int dx dy n_o(x, y)}$$

- jT is not azimuthally symmetric. We define its 2nd Fourier coefficient jT_2 :

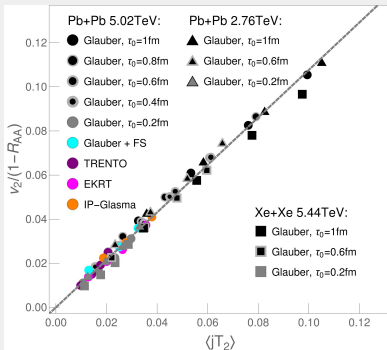
$$jT_2(\tau) = \frac{\int dx dy n_o(x, y) \int \phi \cos 2\phi T^3(x + \tau \cos \phi, y + \tau \sin \phi, \tau)}{\int dx dy n_o(x, y) \int \phi T^3(x + \tau \cos \phi, y + \tau \sin \phi, \tau)}$$

JET-PERCEIVED ANISOTROPY

■ A simple time-average of jT_2 : jet-perceived anisotropy:

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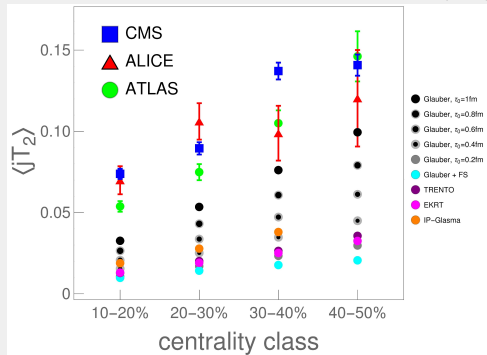
$$\langle jT_2 \rangle = \frac{\int_{\tau_0}^{\tau_{\text{cut}}} d\tau jT_2(\tau)}{\tau_{\text{cut}} - \tau_0}$$



- τ_{cut} : the time when the center of the fireball has cooled to critical temperature T_C .
- $v_2/(1 - R_{AA})$ shows a linear dependence on $\langle jT_2 \rangle$, with a slope close to 1.
- $v_2/(1 - R_{AA})$ carries information on this property of the medium.

JET-PERCEIVED ANISOTROPHY

- We evaluated $\langle jT_2 \rangle$ from experimentally measured $R_{AA}(p_\perp)$ and $v_2(p_\perp)$: the fitted ratio was converted to $\langle jT_2 \rangle$.



- All three experiments lead to similar values of $\langle jT_2 \rangle$.
- Jet-perceived anisotropy provides an important constraint on bulk-medium simulations - they should be tuned to reproduce it.

CONCLUSIONS AND ACKNOWLEDGEMENTS

- High- p_{\perp} theory and data - traditionally used to explore high- p_{\perp} parton interactions with QGP.
- High- p_{\perp} probes can become powerful tomography tools, as they are sensitive to global QGP properties (e.g. spatial anisotropy).
- A (modified) ratio of R_{AA} and v_2 - a reliable and robust observable for straightforward extraction of spatial anisotropy.
- The saturation is directly proportional to jet-perceived anisotropy.
- It will be possible to infer anisotropy directly from LHC Run 3 data: an important constraint to models describing the early stages of QGP formation.
- Synergy of more common approaches for inferring QGP properties with high- p_{\perp} theory and data.

ACKNOWLEDGEMENTS



European Research Council

Established by the European Commission



**МИНИСТАРСТВО ПРОСВЕТЕ,
НАУКЕ И ТЕХНОЛОШКОГ РАЗВОЈА**

BACKUP

2.1. Medium evolution

Our starting point and reference we used for all collision energies and systems is a simple optical Glauber model based initialisation. In Pb+Pb collisions at full LHC energy ($\sqrt{s_{NN}} = 5.02$ TeV) we used initial times $\tau_0 = 0.2, 0.4, 0.6, 0.8,$ and 1.0 fm, whereas the lower energy ($\sqrt{s_{NN}} = 2.76$ TeV) Pb+Pb and Xe+Xe ($\sqrt{s_{NN}} = 5.44$ TeV) calculations were carried out for $\tau_0 = 0.2, 0.6,$ and 1.0 fm. The initialisation and code used to solve viscous fluid-dynamical equations in 3+1 dimensions are described in detail in Ref. [22], and parameters to describe Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV in Ref. [23]. In particular, we use a constant shear viscosity to entropy density ratio $\eta/s = 0.12$ (Pb+Pb) or $\eta/s = 0.10$ (Xe+Xe), and the EoS parametrisation s95p-PCE-v1 [25].

Different initial state models lead to slightly different shapes of the initial state. To find if our findings are a feature of the Glauber model, or have broader significance, we did the Pb+Pb calculations at the full LHC energy using several different initial state models. The first option in this extended set, Glauber + Free streaming, is to use the Glauber model to provide the initial distribution of (marker) particles, allow the particles to stream freely from $\tau = 0.2$ to 1.0 fm, evaluate the energy-momentum tensor of these particles, and use it as the initial state of the fluid. We evolve the fluid using the same code as in the case of pure Glauber initialisation. The EoS is s95p-PCE175, i.e., a parametrisation with $T_{\text{chem}} = 175$ MeV [26], and temperature-independent $\eta/s = 0.16$. For further details, see Ref. [23].

As more sophisticated initialisations, we employ EKRT, IP-Glasma and T_RENTo. The EKRT model [27–29] is based on the NLO perturbative QCD computation of the transverse energy and a gluon saturation conjecture. We employ the same setup as used in Ref. [30] (see also [26]), compute an ensemble of event-by-event fluctuating initial density distributions, average them, and use this average as the initial state of the fluid dynamical evolution. We again use the code of Molnar et al., [22], but restricted to boost-invariant expansion. The shear viscosity over entropy density ratio is temperature dependent with favoured parameter values from the Bayesian analysis of Ref. [30]. Initial time is $\tau_0 = 0.2$ fm, and the EoS is the s83s₁₈ parametrisation from Ref. [30].

IP-Glasma model [31,32] is based on Color Glass Conden-

gluon fields by solving classical Yang-Mills equations. The calculated event-by-event fluctuating initial states [37] were further evolved [38] using the MUSIC code [39–41] constrained to boost-invariant expansion. We subsequently averaged the evaluated temperature profiles to obtain one average profile per centrality class. In these calculations, the switch from Yang-Mills to fluid-dynamical evolution took place at $\tau_{\text{switch}} = 0.4$ fm, shear viscosity over entropy density ratio was constant $\eta/s = 0.12$, and the temperature-dependent bulk viscosity coefficient over entropy density ratio had its maximum value $\zeta/s = 0.13$. The equation of state was based on the HotQCD lattice results [42] as presented in Ref. [43].

T_RENTO [44] is a phenomenological model capable of interpolating between wounded nucleon and binary collision scaling, and with a proper parameter value, of mimicking the EKRT and IP-Glasma initial states. As with the EKRT initialisation, we create an ensemble of event-by-event fluctuating initial states, sort them into centrality classes, average, and evolve these average initial states. Unlike in other cases, we employ the version of the VISH2+1 code [45] described in Refs. [46,47]. We run the code using the favoured values of the Bayesian analysis of Ref. [47]; in particular, allow free streaming until $\tau = 1.16$ fm, the minimum value of the temperature-dependent η/s is 0.081, and the maximum value of the bulk viscosity coefficient ζ/s is 0.052. The EoS is the same HotQCD lattice results [42] based parametrisation as used in Refs. [46,47].

It is worth noticing that the initial nuclear configuration in all these cases is similar Woods-Saxon parametrisation of nuclear matter density, which is either assumed to be continuous (optical Glauber), or Monte-Carlo sampled to create ensembles of nucleons (EKRT, IP-Glasma, T_RENTO). The differences in the fluid-dynamical initial state depend on the initial particle production, and subsequent evolution before fluid-dynamical stage (none, Yang-Mills, free streaming).

All these calculations were tuned to reproduce, in minimum, the centrality dependence of charged particle multiplicity, p_{\perp} distributions and $v_2(p_{\perp})$ in Pb+Pb collisions at both collision energies, and the centrality dependence of charged particle multiplicity