

Jets in hot nuclear matter (A unified picture of medium-induced radiation)

Adam Takacs (Heidelberg University, University of Bergen),

Johannes H. Isaksen (University of Bergen),

Konrad Tywoniuk (University of Bergen),

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UNIVERSITÄT
HEIDELBERG
ZUKUNFT
SEIT 1386



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Introduction to medium-induced emissions

[arXiv:2206.02811](https://arxiv.org/abs/2206.02811)

QCD in a background medium

[Zakharov, BDMPS]

[Blaizot, Dominiguez, Iancu, Mehtar-Tani]

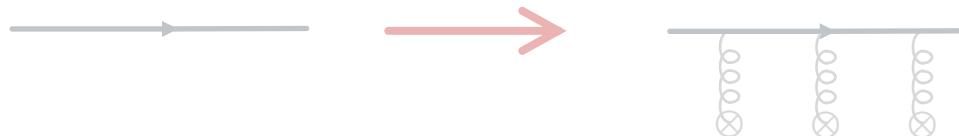
QCD with color bkg: $\mathcal{A}(t, \mathbf{x}) + \mathcal{A}_0(t, \mathbf{x})$

- Multiple scatterings

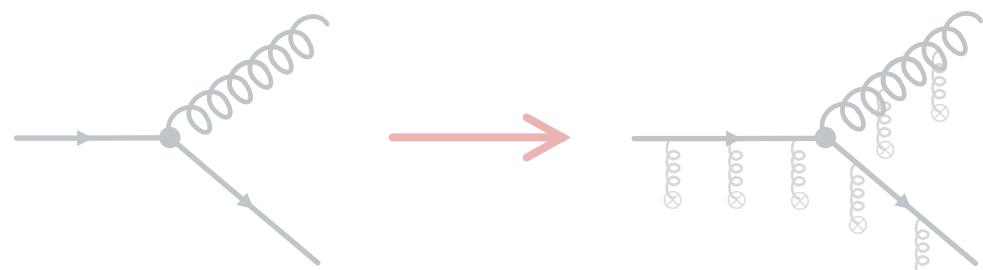


Medium Feynman rules:

- medium propagator:



- medium vertex:



- Medium average:

$$\langle \mathcal{A}_0^-(t, \mathbf{x}) \mathcal{A}_0^-(t', \mathbf{x}') \rangle_{med}$$

- weakly coupled, thermal plasma
- random fields
- “idk, evaluate later”

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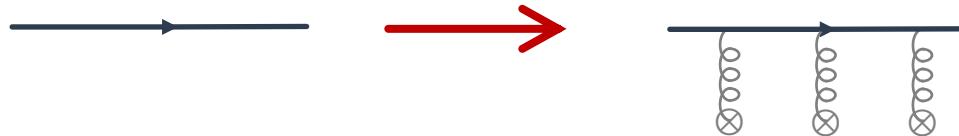
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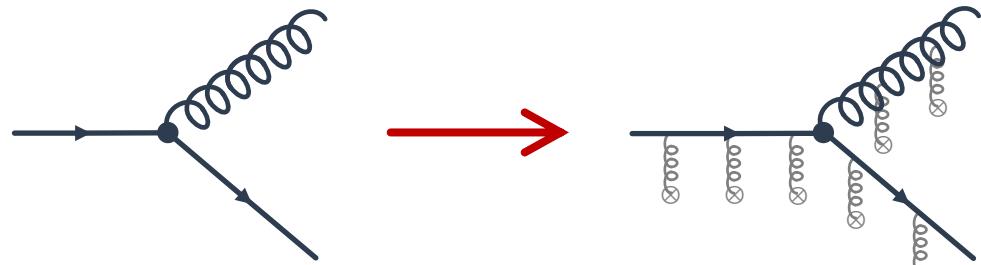


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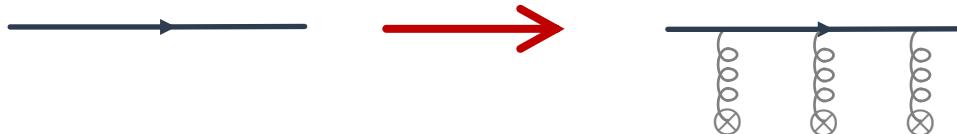
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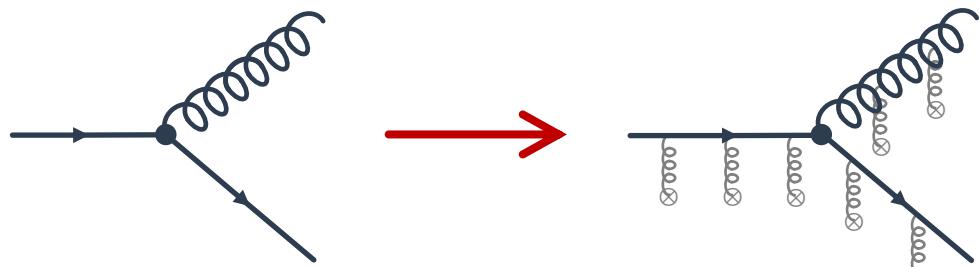


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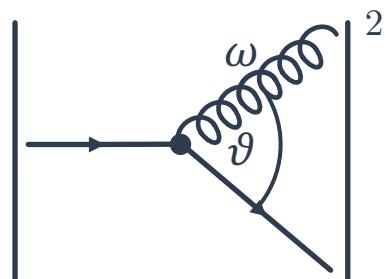
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Medium-induced emission

LO radiation in vacuum:



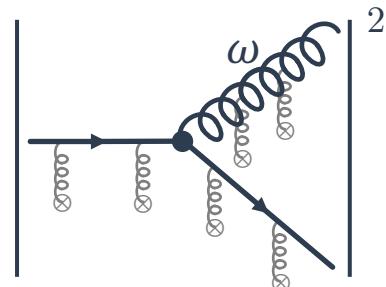
$$\frac{dI^{vac}}{d\omega d\vartheta} \sim \frac{2\alpha_s}{\pi} \frac{C_i}{\omega} \frac{1}{\vartheta}$$



soft and collinear singularity

Medium-induced emission

LO radiation:



$$\frac{dI}{d\omega d\vartheta} = \frac{dI^{vac}}{d\omega d\vartheta} + \frac{dI^{med}}{d\omega d\vartheta} = \text{complicated integral}$$

[Isaksen,Tywoniuk(2023)]

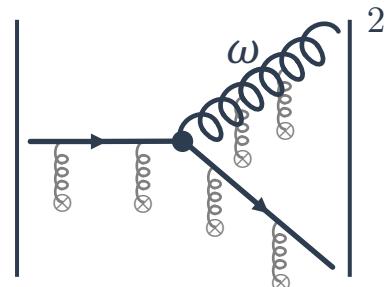
- $I^{med} > 0$: induced emissions
- I^{med} : no collinear divergence: ϑ is integrated

Approximate solutions:

1. analytic:
 - harmonic oscillator [BDMPSZ(1997)]
 - opacity expansion [GLV-Wiedemann(2000)]
 - improved opacity expansion [Mehtar-Tani(2020)]
2. numeric: [Feal,Vazquez(2018)]
[Andres,Aploinario,Dominiques(2020)]
[Schlichting,Soudi(2021)]

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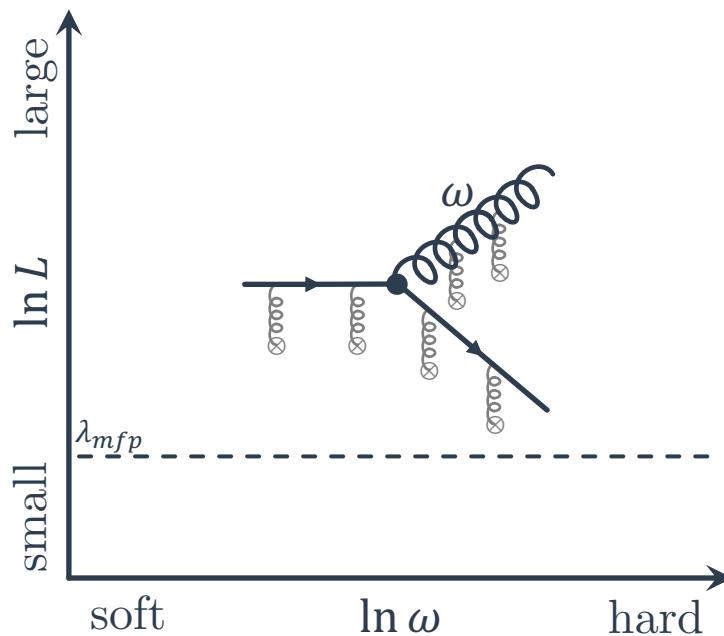
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Unified picture of MIE

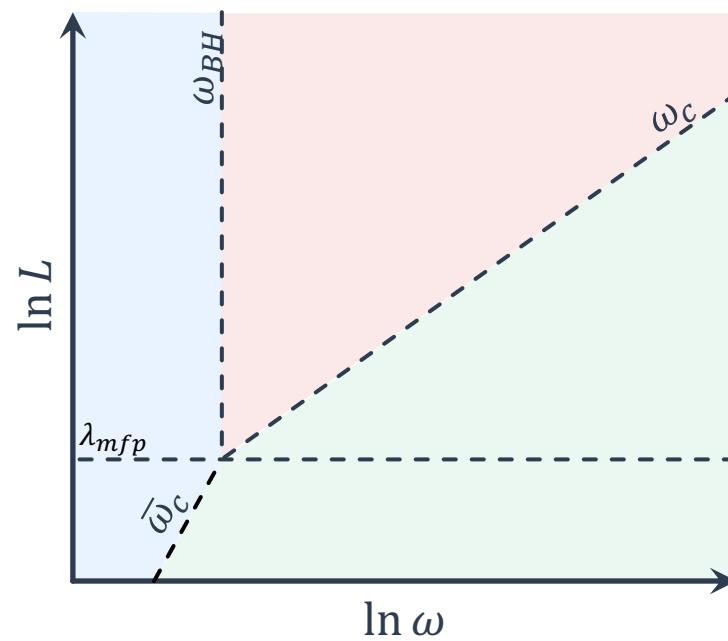
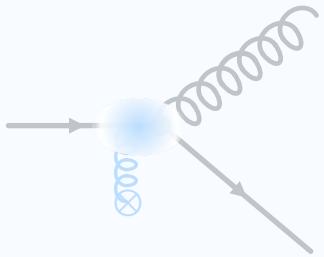
Leading physical picture of $\frac{dI^{med}}{d\omega} = \int d\vartheta \frac{dI^{med}}{d\omega d\vartheta}$:



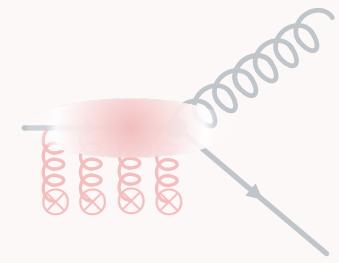
Unified picture of MIE

Physical picture of $\frac{dI^{med}}{d\omega} = \int d\vartheta \frac{dI^{med}}{d\omega d\vartheta}$:

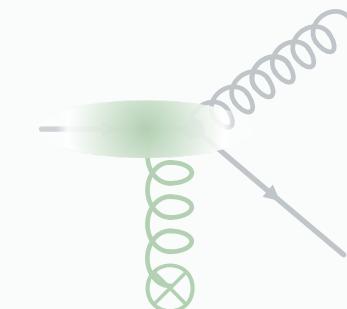
few soft scatterings



multiple soft scatterings



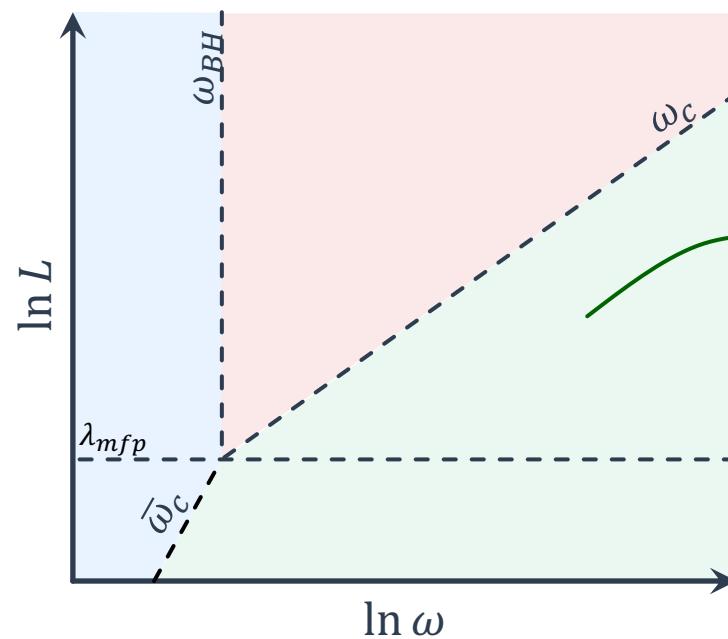
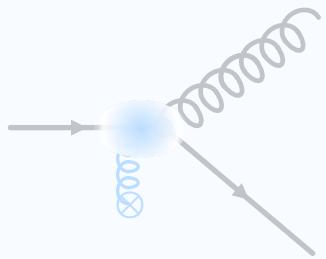
rare hard scatterings



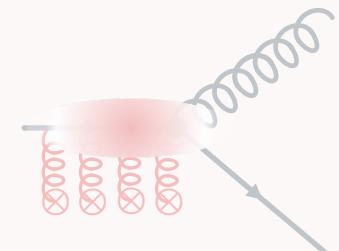
Unified picture of MIE

Physical picture of $\frac{dI^{med}}{d\omega} = \int d\vartheta \frac{dI^{med}}{d\omega d\vartheta}$:

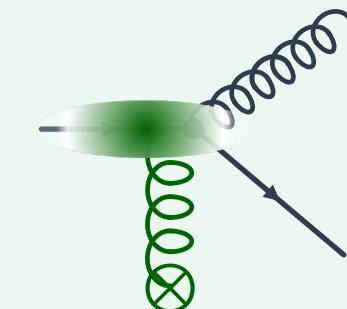
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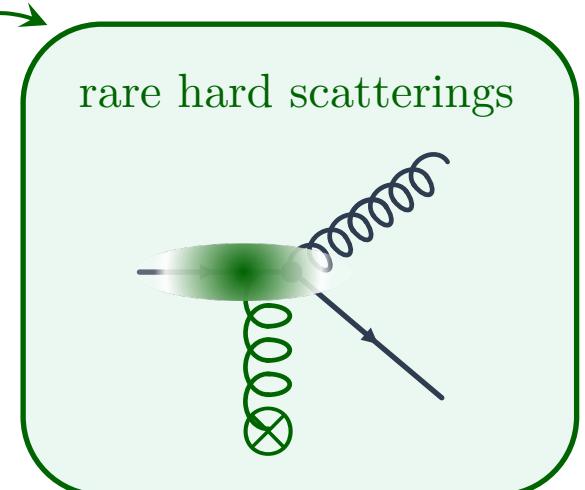
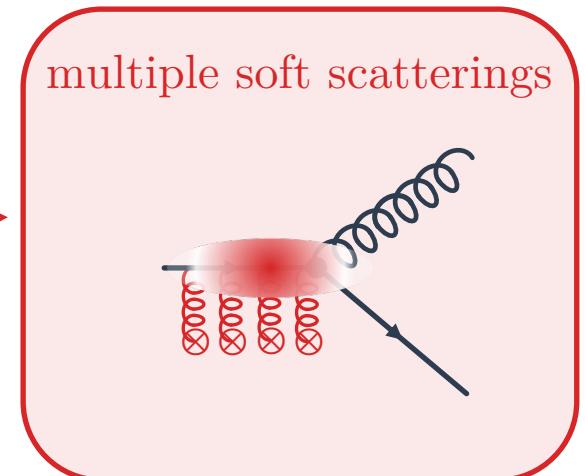
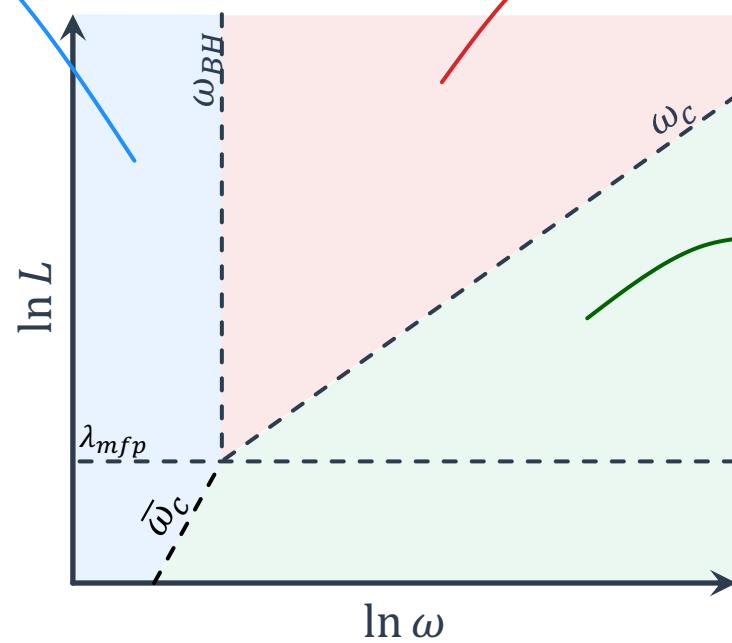
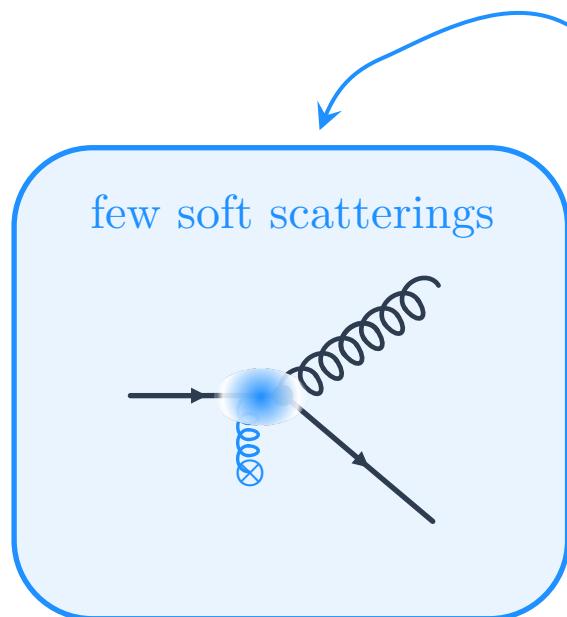


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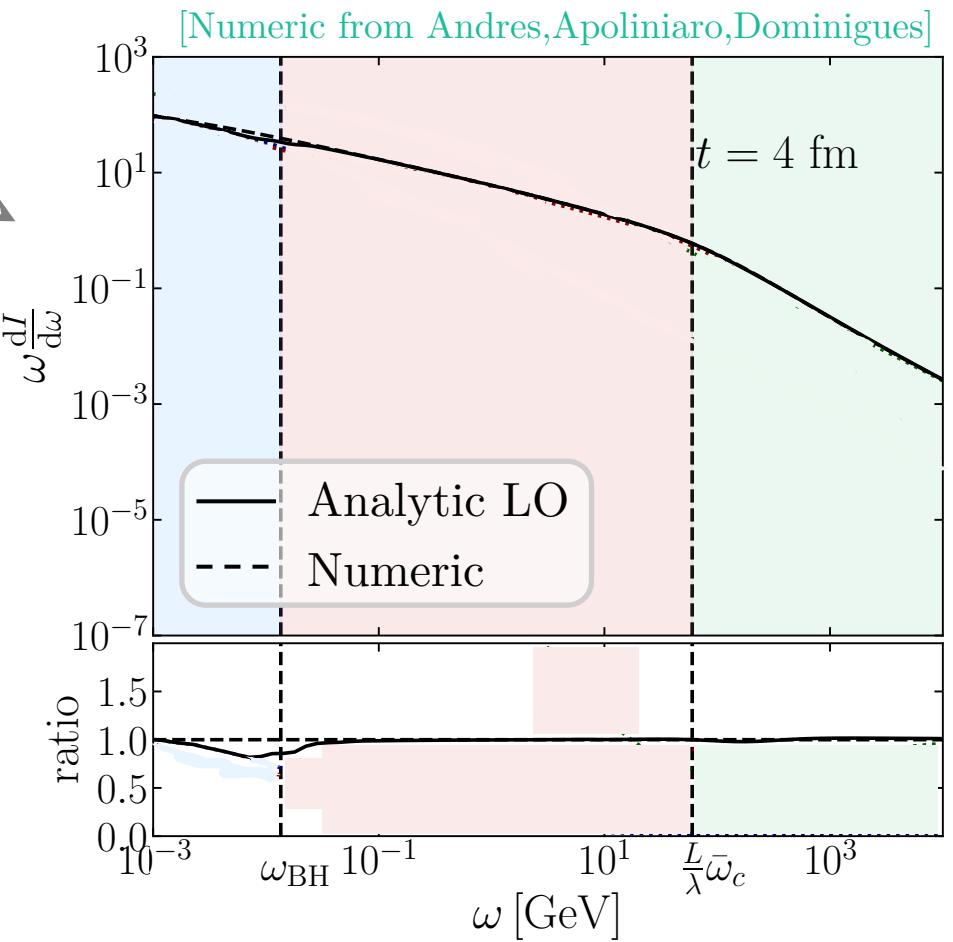
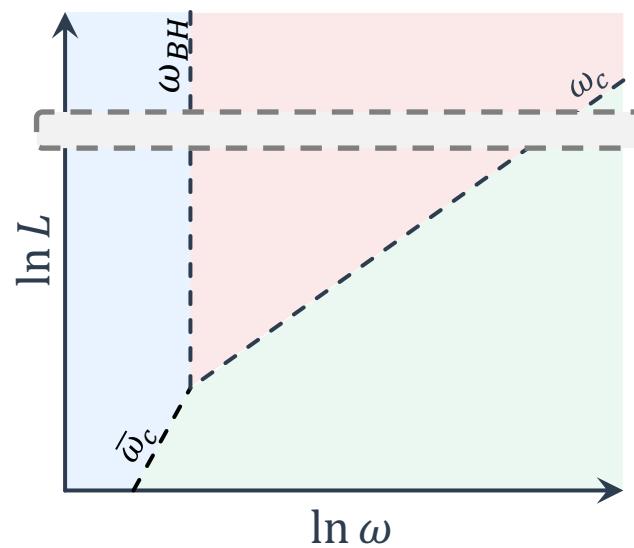
Unified picture of MIE

Physical picture of $\frac{dI^{med}}{d\omega} = \int d\vartheta \frac{dI^{med}}{d\omega d\vartheta}$:



Testing the unified picture

Comparing to numerical solution:



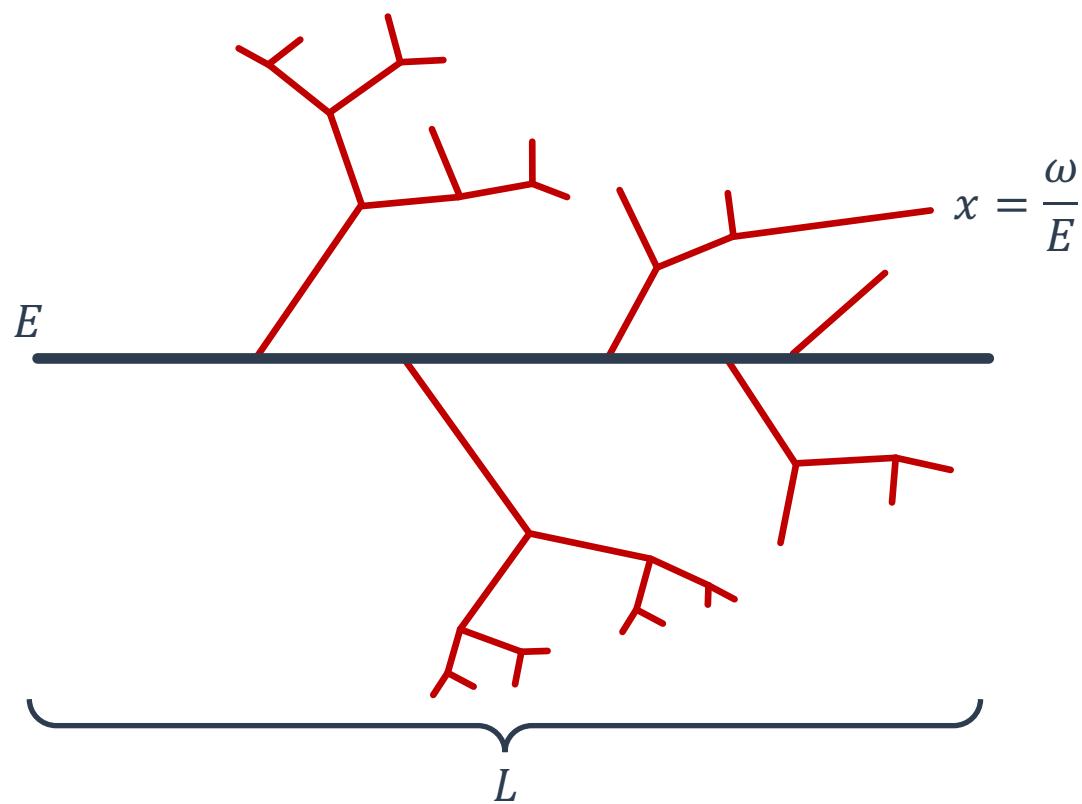
- Very good agreement.
- Computationally very effective.
- <https://github.com/adam-takacs/kernels>

Application of the unified picture

[arXiv:2206.02811](https://arxiv.org/abs/2206.02811)

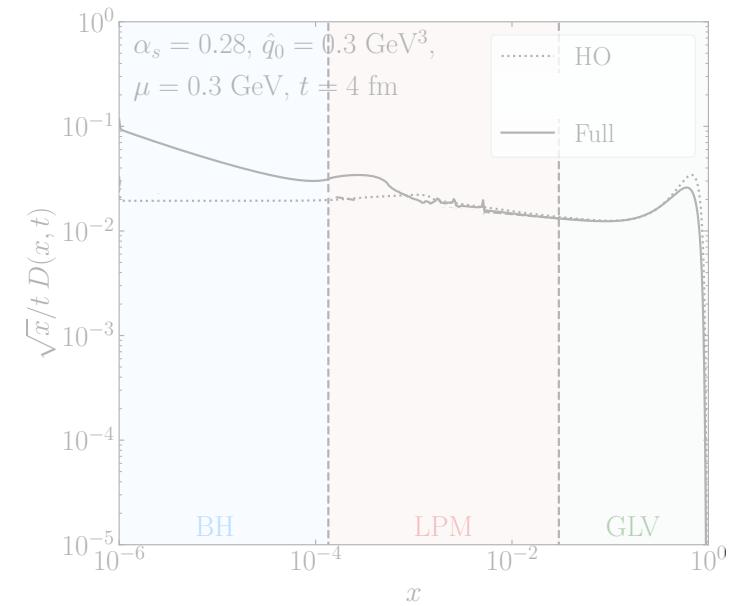
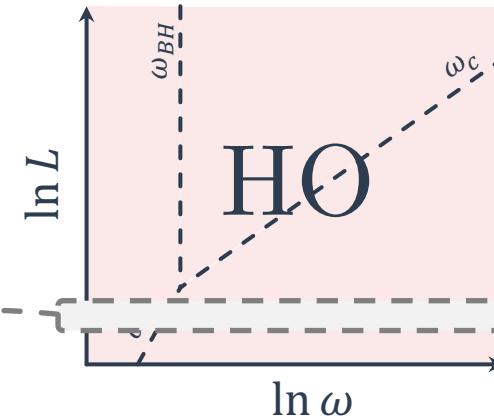
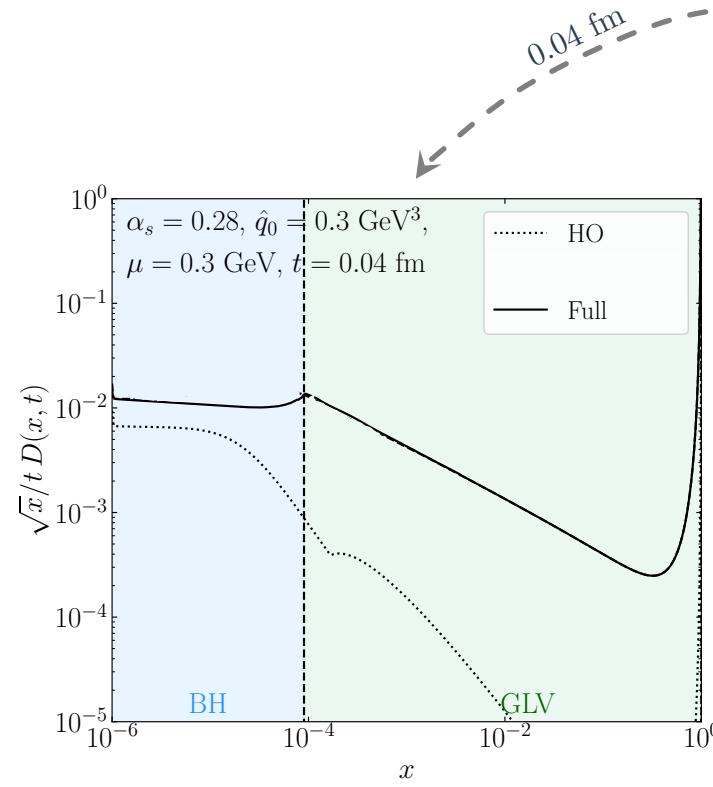
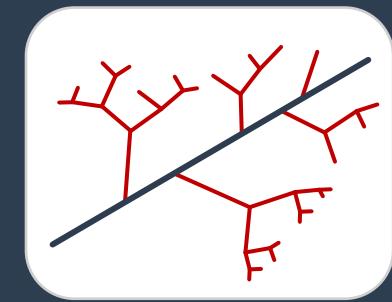
Application: MIE cascade

Medium-induced fragmentation function:

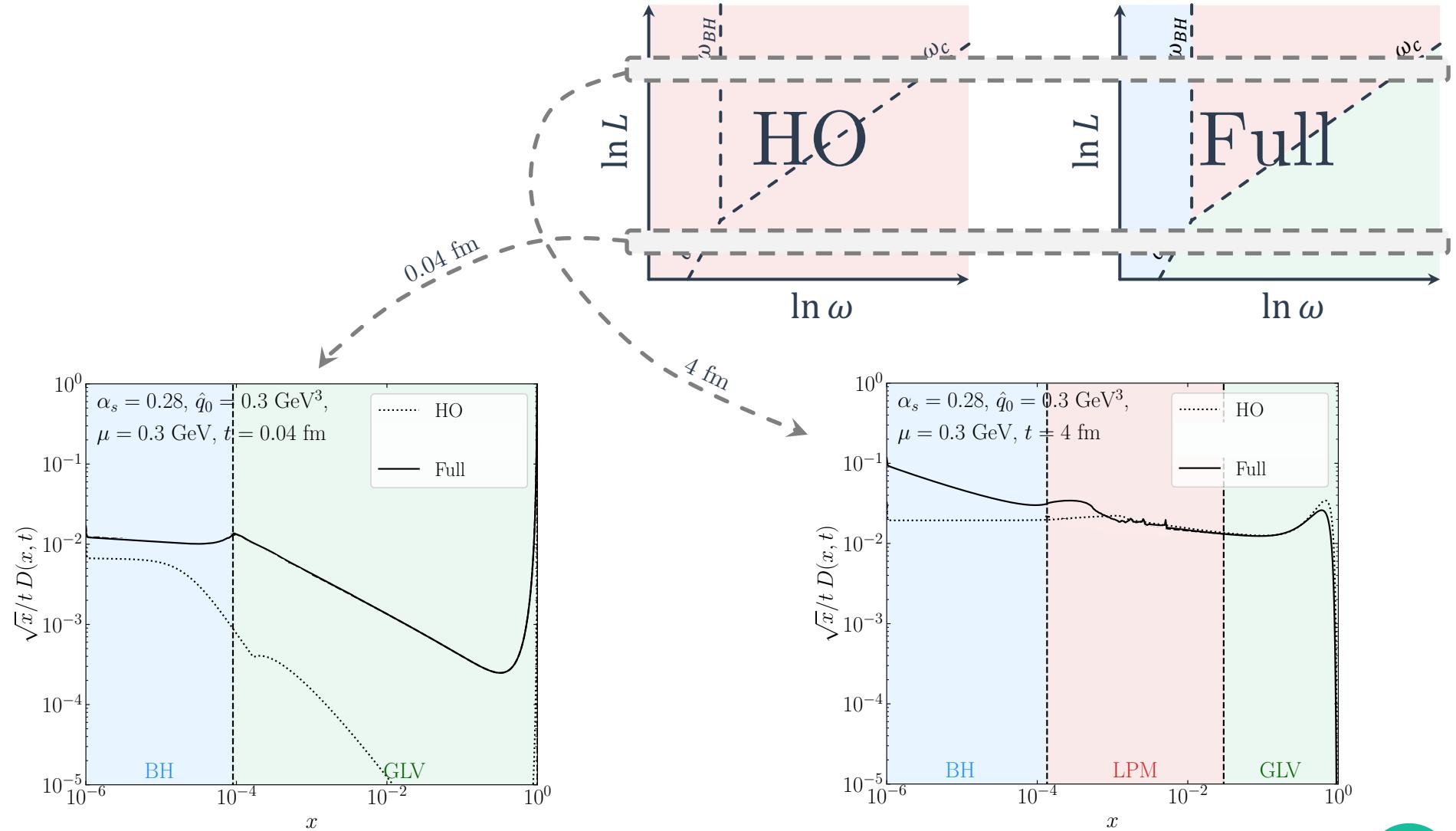
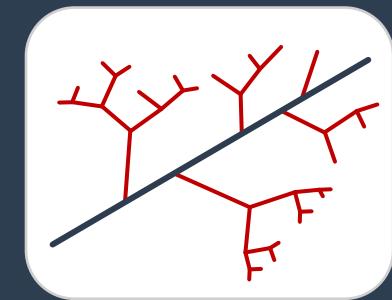


$$D(x, t) = x \frac{dN}{dx}$$

Application: MIE cascade



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Summary

- Understanding jet modification in medium
- In this work:
 - single induced emission $\frac{dI}{d\omega}$,
 - energy loss of a parton
- Outlook:
 - event generator implementation
 - energy loss of a whole jet
 - understanding pQCD in the plasma to all orders

Thank you for the attention!

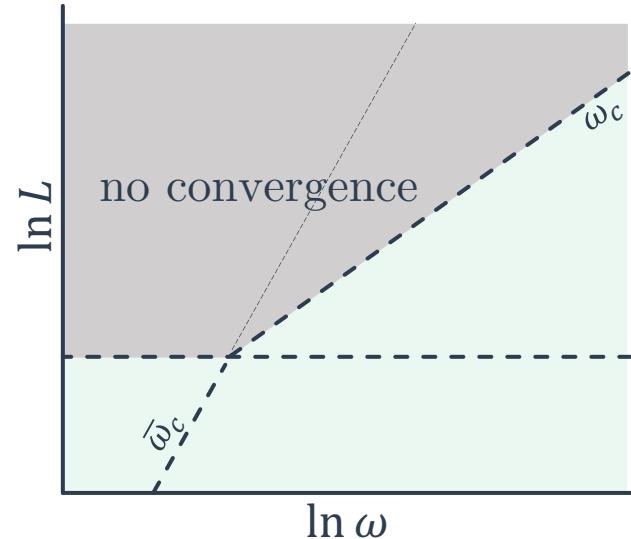


Unified picture of MIE

Expansion in scatterings:

$$\bar{\omega}_c = \frac{\mu^2 L}{2}$$

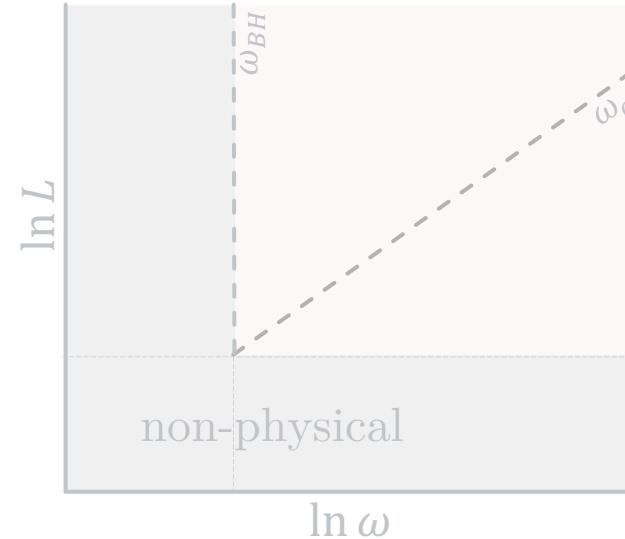
$$\omega \frac{dI}{d\omega} = \begin{cases} \bar{\alpha} \sum_{n=1}^{\infty} \left(\frac{L}{\lambda}\right)^n f_n\left(\frac{\omega}{\bar{\omega}_c}\right), & \omega \ll \bar{\omega}_c \\ \bar{\alpha} \sum_{n=1}^{\infty} \left(\frac{L \bar{\omega}_c}{\lambda \omega}\right)^n f'_n\left(\frac{\bar{\omega}_c}{\omega}\right), & \bar{\omega}_c \ll \omega \end{cases}$$



Multiple soft scatterings:

$$\omega_c = \frac{L}{\lambda} \bar{\omega}_c$$

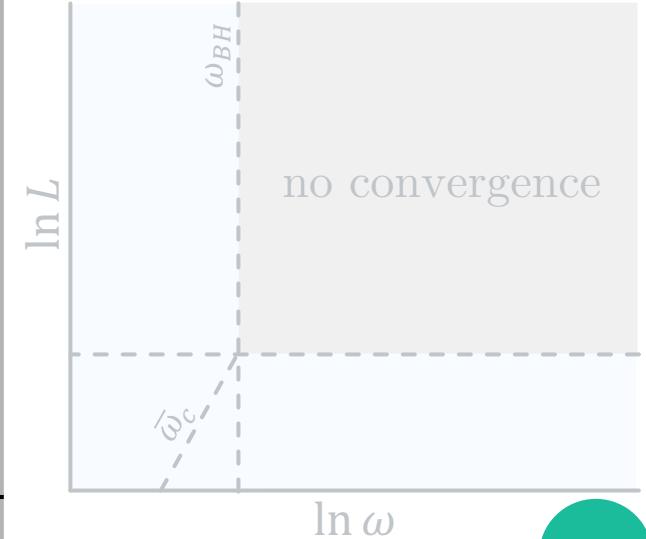
$$\omega \frac{dI^{HO}}{d\omega} = \text{Closed form}, \quad \omega_{BH} \ll \omega \ll \omega_c$$



Expansion in real scatterings:

$$\omega_{BH} = \frac{\bar{\omega}_c}{L/\lambda}$$

$$\omega \frac{dI}{d\omega} = \begin{cases} \text{Opacity expansion, } L \ll \lambda \\ \bar{\alpha} \frac{L}{\lambda} \sum_{n=1}^{\infty} g_n\left(\frac{\omega}{\omega_{BH}}\right), \omega \ll \omega_{BH} \end{cases}$$

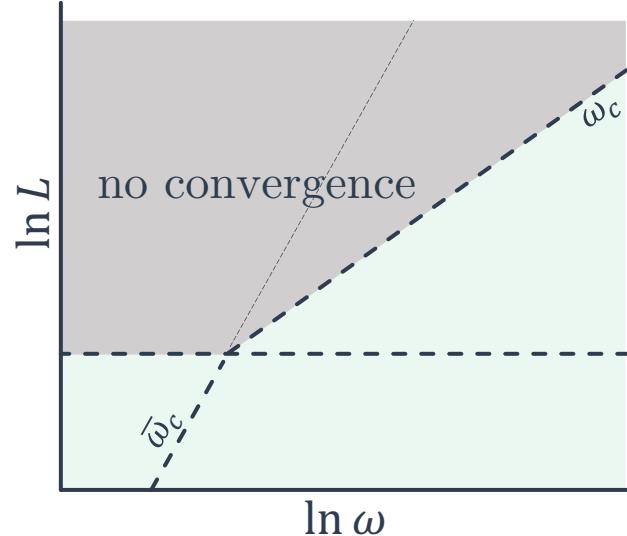


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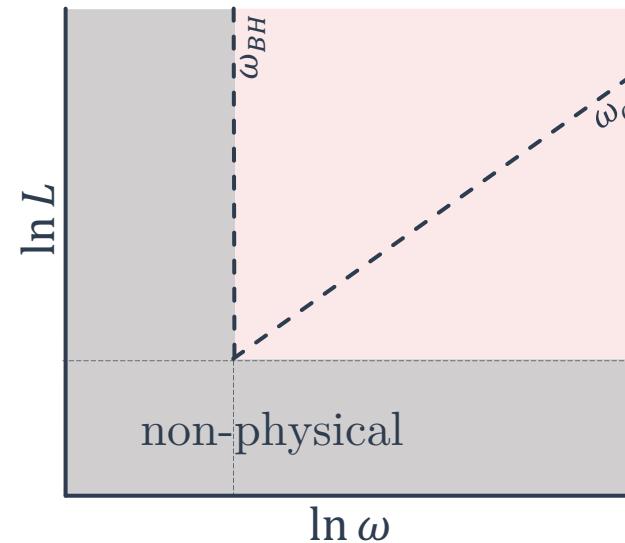
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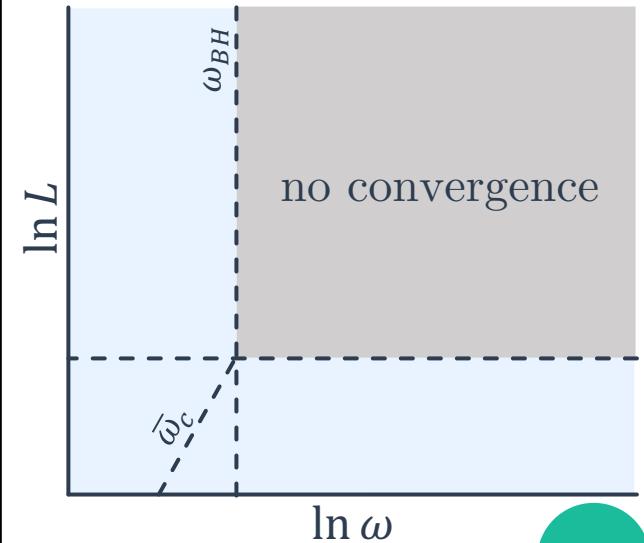
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Application: accuracy

Single emission in vacuum

$$\int dz \int d\vartheta P^{LO}(z, \vartheta) = \mathcal{O}(\alpha_s L^2) + \mathcal{O}(\alpha_s L) + \mathcal{O}(\alpha_s)$$

logarithmic enhancement: soft & collinear emission

Power counting in vacuum

$$\Sigma(v) = \int dv' \frac{1}{\sigma_0} \frac{d\sigma}{dv'}$$

$$= 1 + \alpha_s (\Sigma_{12} L^2 + \dots)$$

$$+ \alpha_s^2 (\Sigma_{24} L^4 + \Sigma_{23} L^3 + \Sigma_{22} L^2 + \dots)$$

$$+ \dots$$

$$+ \alpha_s^n (\Sigma_{n,2n} L^{2n} + \Sigma_{n,2n-1} L^{2n-1} + \dots)$$

DL NDL

Single emission in medium

$$\int dz \int d\vartheta I^{LO_m}(z, \vartheta) = \mathcal{O}(\alpha_s P_m) + \mathcal{O}(\alpha_s L_m) + \mathcal{O}(\alpha_s)$$

power & logarithmic enhancement: soft emission & large medium

Power counting in medium

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P_m NP_m