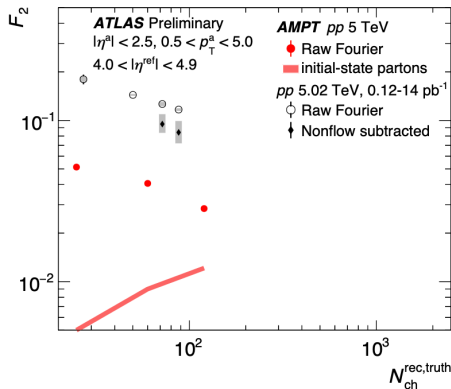




Multiparton picture of elliptic flow decorrelation in rapidity in pp collisions at the LHC

Piotr Bożek, Wojciech Broniowski

ISMD 2023, Gyöngyös, 21-26 August 2023



F_2 - (see following)

[ATLAS-CONF-2022-020]

flow decorrelation measure between forward and backward \vec{v}_2

$F_2 = 0$ - no decorrelation (perfect correlation)

F_2 large - large decorrelation (or small correlation)

Hitherto unexplained!

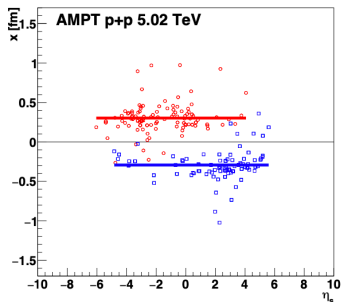
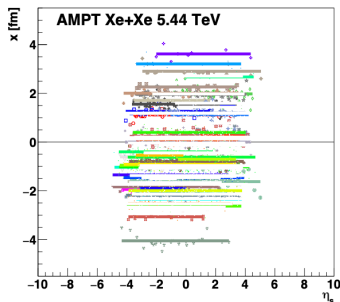
Outline

- Strings/flux tubes: FB picture
- Shape-flow transmutation
- FB flow correlation
- Model: multiple strings + fluctuations in hadronization
- F_2 in pp

Strings

Dual parton, Lund

AMPT [Wu et al. 2018]

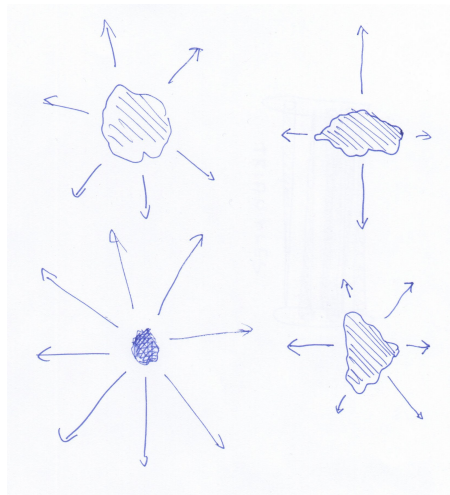
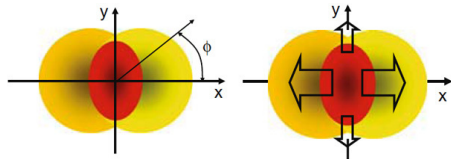


String end-points fluctuate in η_S , shorter strings \rightarrow larger FB decorrelation production of particles from the string, rescattering (!)

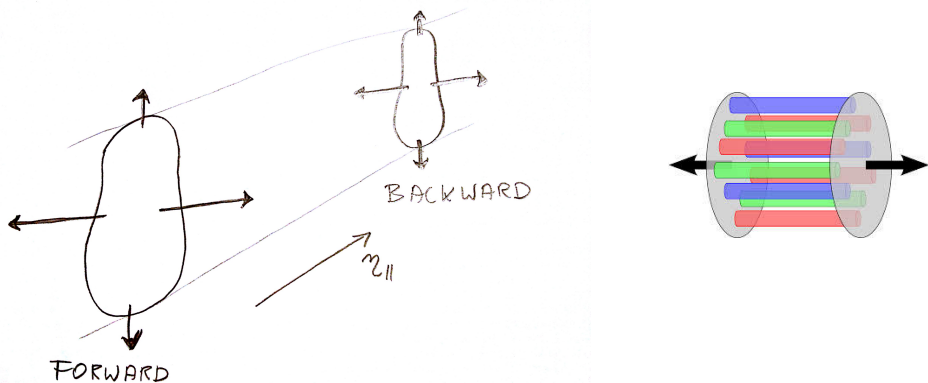
Shape-flow transmutation

Many particles, **rescattering**, generation of flow (not necessarily hydro!),

$$v_2 \sim \epsilon_2$$



FB shape similarity \rightarrow flow similarity \rightarrow ridges



FB similarity is a **result of early dynamics** and is not obtained from hydro or transport

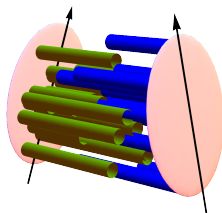
FB correlations

- One expects substantial FB correlations, e.g., for FB event plane angles (direction of flow), harmonic flow magnitude, or $\langle p_T \rangle$
- Focus on departures from perfect correlations, the **torque effect**:
twisted event-plane angles [Bożek, WB, Moreira 2010]

Fluctuations in energy deposition from each string

[Brodsky, Gunion, Kuhn, PRL 39 (1977) 1120]

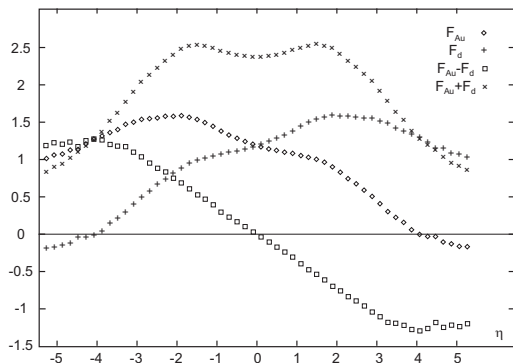
[PB, WB 2015, 2017]



- Position (in spatial rapidity) of one of the string end-points is random
- Results in FB decorrelation

Phenomenological average emission profiles from d-Au

[Białaś, Czyż, Acta Phys. Polon. B36, 905 (2005)]



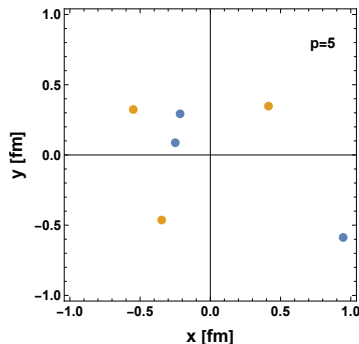
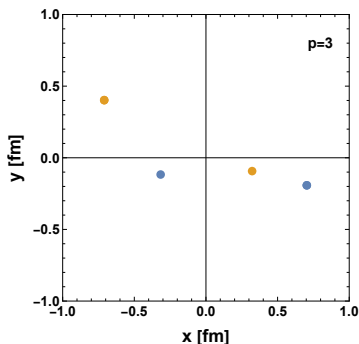
We use **uniformly** distributed end-points, which yields on the average the **triangular distribution**, supported phenomenologically

Wounded parton model

A model of the proton is needed. Recall the [near-side ridge](#) in high-multiplicity pp, interpreted as a sign of shape-flow transmutation

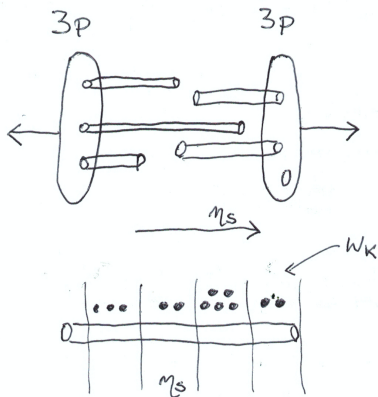
GLISSANDO 3 with [wounded partons](#) for the transverse shape pp cross section reproduced

We take exploratory cases of 3 and 5 partons per nucleon



Full picture

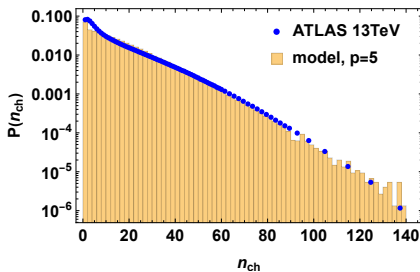
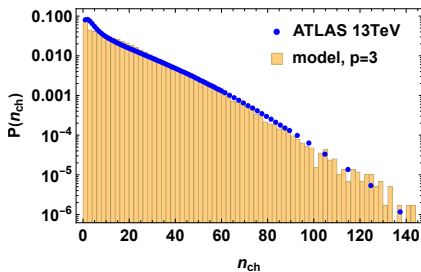
GLISSANDO 3 + strings with fluctuating end points



+ overlaid multiplicity fluctuation in each η bin ($\eta \simeq \eta_s$)

Multiplicity distribution in p-p

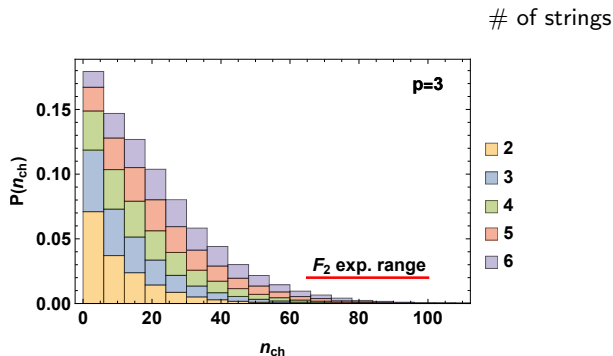
We overlay the distribution $\Gamma(x; \alpha, \beta) \sim x^{\alpha-1} e^{-\beta x}$, with parameters α, β (depending on the collision energy) adjusted to reproduce the height and slope of the [tail](#) of the data:



[data: $n_{\text{ch}} \geq 1$, $p_T > 500$ MeV, $|\eta| < 2.5$]

Anatomy of the model multiplicity

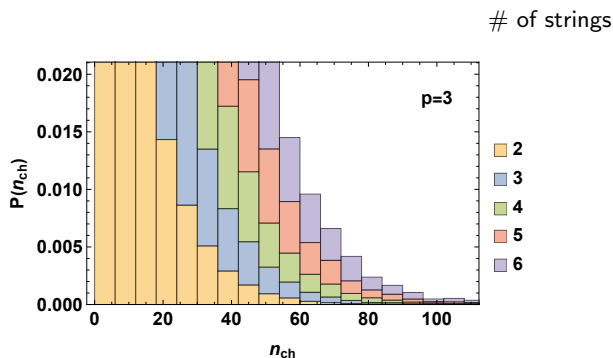
Tail dominated by numerous strings



higher multiplicity – more strings

Anatomy of the model multiplicity

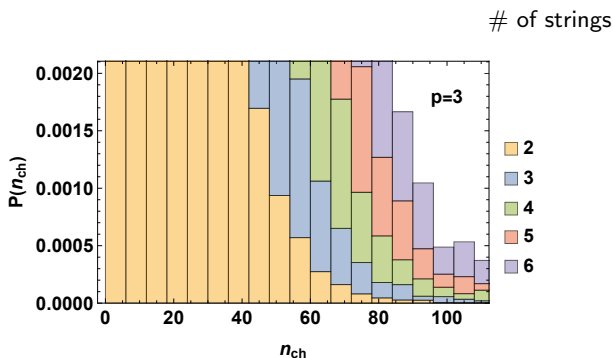
Tail dominated by **numerous strings**



higher multiplicity – more strings

Anatomy of the model multiplicity

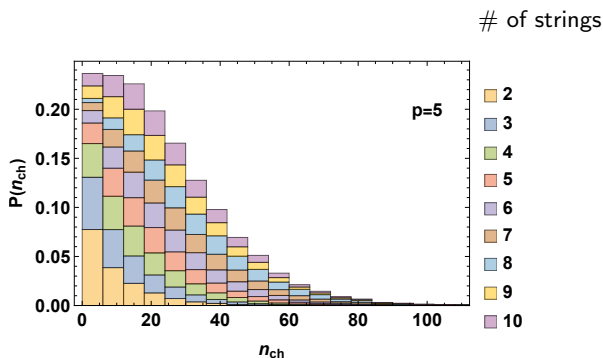
Tail dominated by numerous strings



higher multiplicity – more strings

Anatomy of the model multiplicity

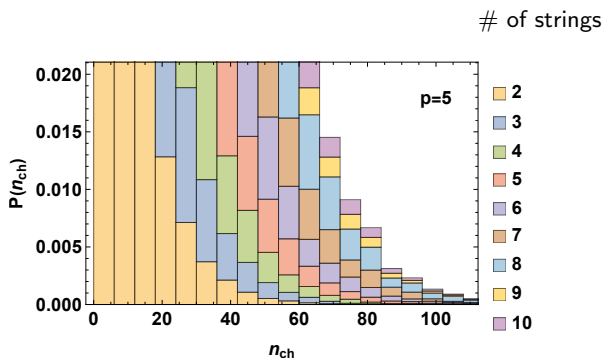
Tail dominated by numerous strings



higher multiplicity – more strings

Anatomy of the model multiplicity

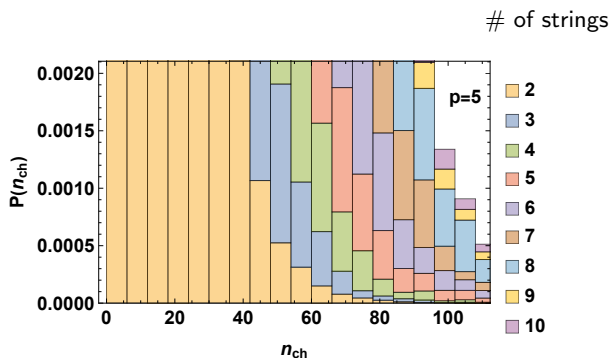
Tail dominated by **numerous strings**



higher multiplicity – more strings

Anatomy of the model multiplicity

Tail dominated by numerous strings



higher multiplicity – more strings

3-bin measure (CMS) of elliptic FB correlations

Need to use measures that cancel trivial decorrelations

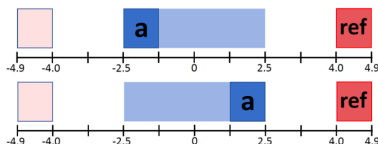
x_k, y_k - transverse coord. (relative to CM) of a piece of string k in an η bin

w_k - weight from overlaid fluctuations

$$\epsilon_2(\eta) = \frac{\sum_k w_k (x_k + iy_k)^2}{\sum_k w_k (x_k^2 + y_k^2)} = \frac{\sum_k w_k r_k^2 e^{2i\phi_k}}{\sum_k w_k r_k^2}$$

$$c_2(\eta, \eta_{\text{ref}}) \equiv \text{Re cov}_{ev} [\epsilon_2(\eta) \epsilon_2^*(\eta_{\text{ref}})], \quad r_2(|\eta_a|) = \frac{c_2(-|\eta_a|, \eta_{\text{ref}})}{c_2(|\eta_a|, \eta_{\text{ref}})}$$

see [CMS 2015]

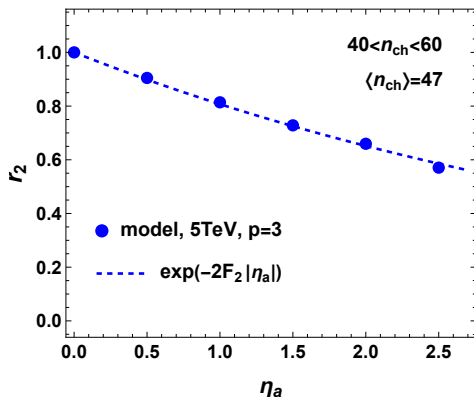


ATLAS measurement

r_2 from v_2 expected to be similar to r_2 from ϵ_2

The slope F_2

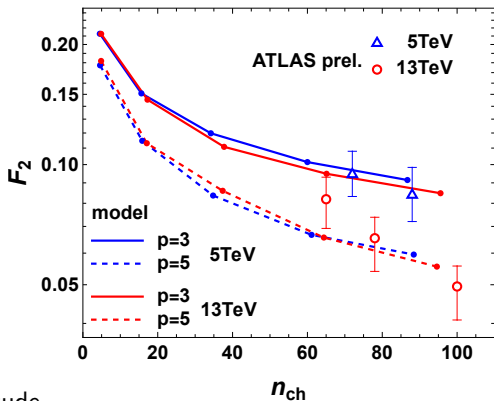
$$r_2(|\eta_a|) \simeq e^{-2F_2|\eta_a|}$$



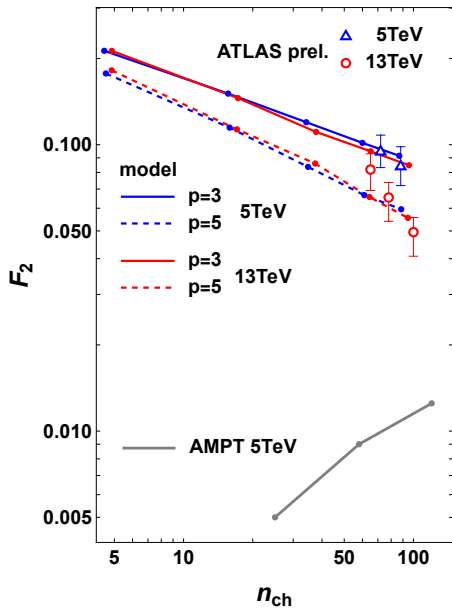
F_2 depends on multiplicity

Model vs ATLAS (preliminary)

(sophisticated corrections in the data analysis to get rid of non-flow)



- Right magnitude
- Proper monotonicity in n_{ch}
- Essentially no dependence on the collision energy between 5 and 13 TeV
- More partons \rightarrow reduced F_2 , or weaker decorrelation, or more correlation



Conclusions

Basic understanding:

- 1 Shorter strings \rightarrow more decorrelation (F_2 up)
- 2 More strings \rightarrow less decorrelation (F_2 down)
- 3 Fluctuations in hadronization must reproduce multiplicity

Message to take home:

F_2 in high-multiplicity pp collisions hints to multiparton dynamics

Possible improvements of our model:

- 1 Collisions of n on m partons (Fock components of the proton wave function)
- 2 More accurate profile of string fluctuations, reference to a_2 coefficient of multiplicity correlations [Rohrmoser, WB, 2019]
- 3 Robust conclusions should remain

Thanks!

BACKUP

Multiplicity in pp vs models

[ATLAS 2016]

