BSQ Hydrodynamics and Challenges with a 4D equation of state

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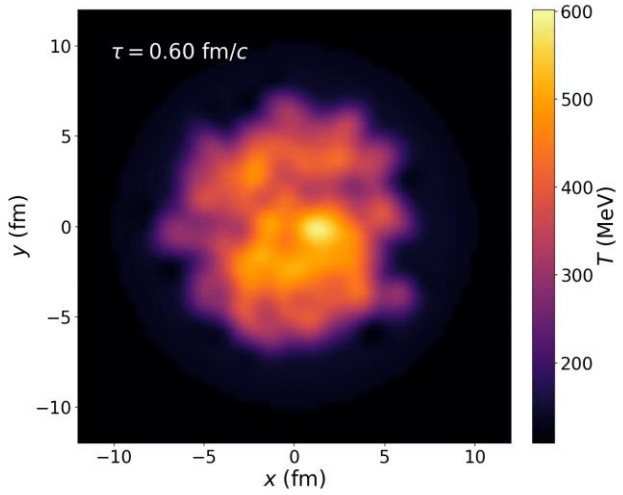
BSQ Simulations

– **Equations of motion**:

$\nabla_\mu T^{\mu\nu} = 0$

$(energy\text{-}momentum\ conservation)$



– **Propagate densities**:

e $(energy\ density)$ or

s $(entropy\ density)$

– Equations of motion closed by **equation of state**: $P = P(T)$

– **Need to know** T coordinate for a **given** e/s point

BSQ Simulations

– **Equations of motion**:

$\nabla_\mu T^{\mu\nu} = 0$

$(energy\text{-}momentum\ conservation)$

$\nabla_\mu J_i^\mu = 0 \quad (i = B, S, Q)$

$(charge\ conservation)$

– **Propagate densities**:

e $(energy\ density)$ or

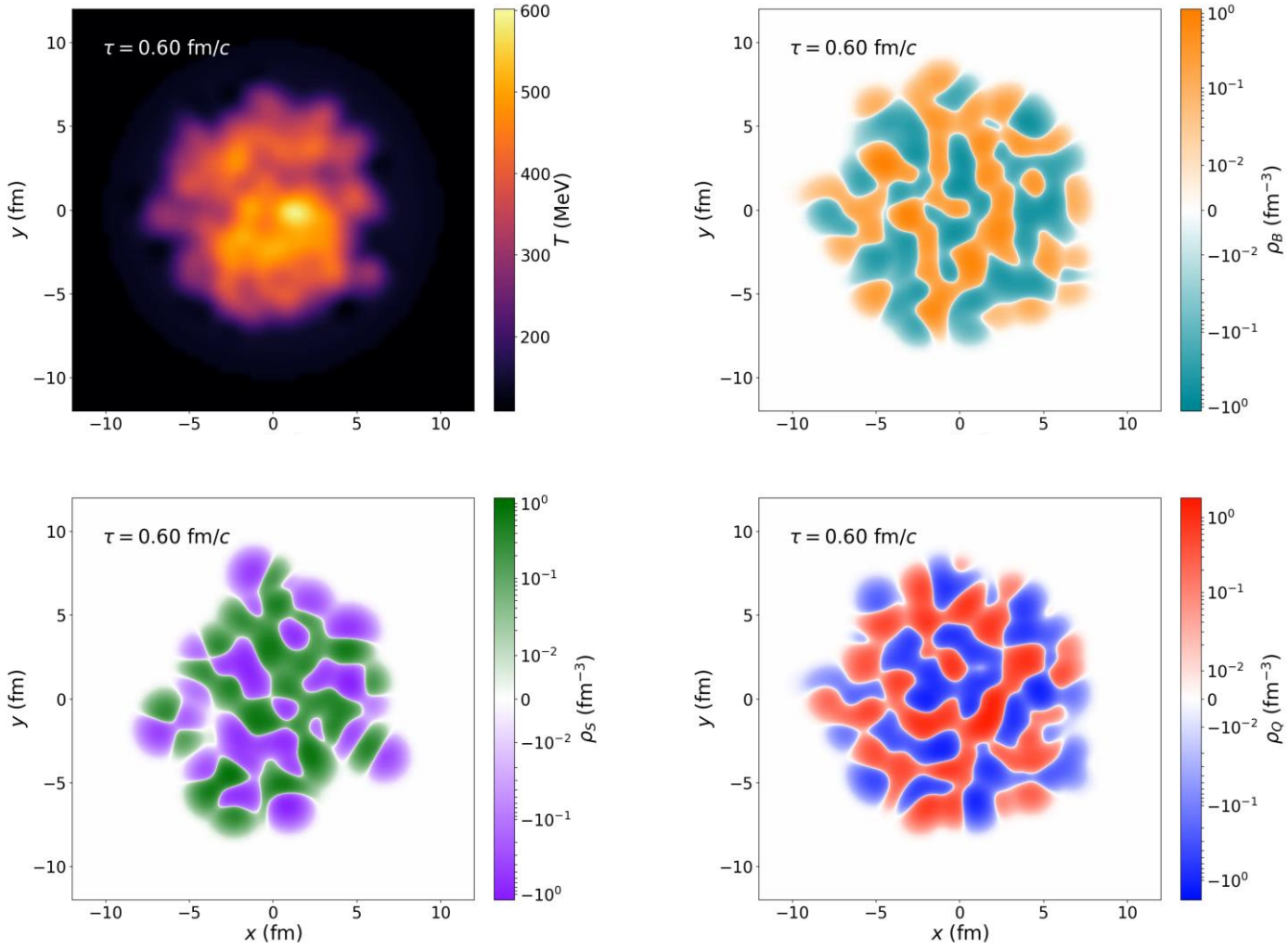
s $(entropy\ density)$

ρ_B $(baryon\ density)$

ρ_S $(net\ strangeness\ density)$

ρ_Q $(electric\ charge\ density)$

– Equations of motion closed by **equation of state**: $P = P(T, \{\mu_i\})$



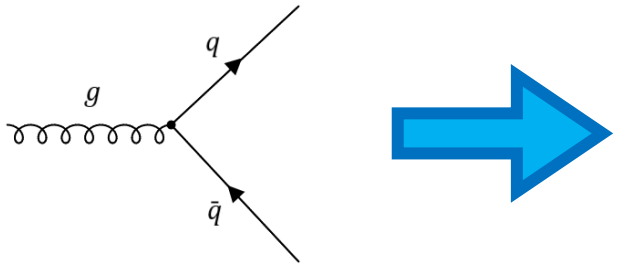
– **Need to know** (T, μ_B, μ_S, μ_Q) coordinates for a **given** $(e/s, \rho_B, \rho_S, \rho_Q)$ point

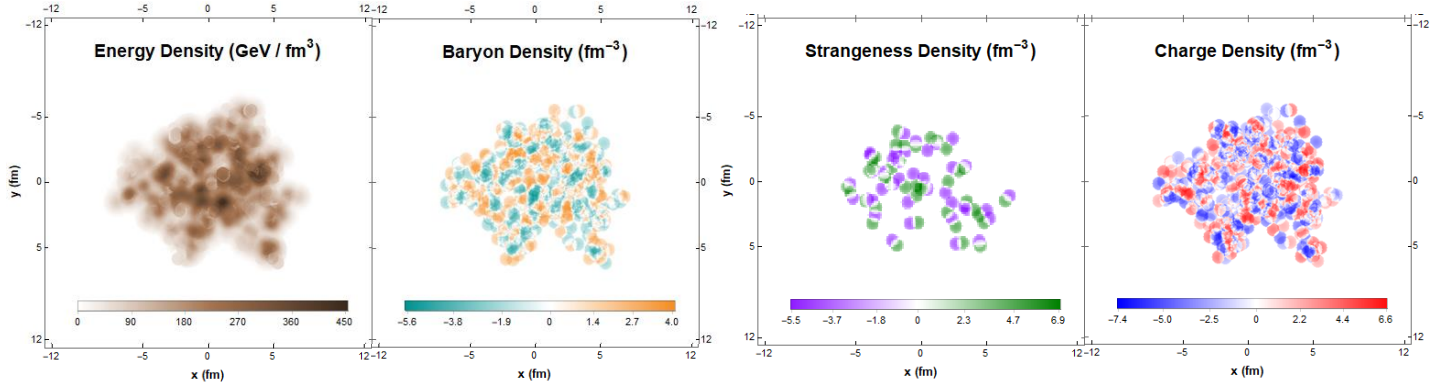
BSQ Initial Conditions and evolution

- ICCING initial conditions (Initial Conserved Charges In Nuclear Geometry)

- ICCING uses gluon-splitting probabilities color-glass condensate (CGC) framework to generate *local charge fluctuations*

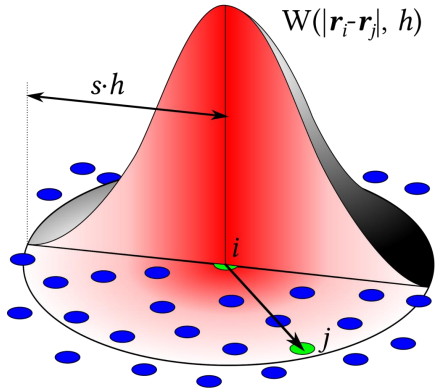
[arXiv:1911.12454 [nucl-th]]





- CCAKE (Conserved ChArges with hydrodynamiK Evolution)

- CCAKE employs *smoothed particle hydrodynamics* (SPH) to model Israel-Stewart theory with charge diffusion
  - Many new transport coefficients
  - Additional relaxation equations



BSQ Thermodynamics

This work: lattice QCD EoS given by Taylor expansion of pressure in powers of chemical potentials

$$\frac{P(T,\mu_B,\mu_Q,\mu_S)}{T^4} = \sum_{i,j,k} \frac{1}{i!j!k!} \chi^{BQS}_{i,j,k} \left(\frac{\mu_B}{T}\right)^i \left(\frac{\mu_Q}{T}\right)^j \left(\frac{\mu_S}{T}\right)^k$$

Claudia Ratti 2018 *Rep. Prog. Phys.* **81** 084301

J. Noronha-Hostler, P. Parotto, C. Ratti, J. Stafford PRC 100 (2019),
A. Monnai et al., PRC 100 (2019)

"Susceptibilities" $\chi^{BQS}_{i,j,k}$ functions of temperature; matched to lattice QCD at high T and hadron resonance gas at low T

$\rho_{B/S/Q}$ and s densities obtained by taking derivatives w.r.t. P

Energy density $e = sT + \sum_{i=B,S,Q} \mu_i \rho_i - p$

How do we invert given set of densities for corresponding phase diagram coordinates?

Root-finding in the equation of state

Goal: obtain $(T_0, \mu_{B,0}, \mu_{S,0}, \mu_{Q,0})$ from $(e_0, \rho_{B,0}, \rho_{S,0}, \rho_{Q,0})$

Given:

$$e = e(T, \vec{\mu})$$
$$\rho_B = \rho_B(T, \vec{\mu})$$
$$\rho_S = \rho_S(T, \vec{\mu})$$
$$\rho_Q = \rho_Q(T, \vec{\mu})$$

Solve:

$$e_0 = e(T_0, \vec{\mu}_0)$$
$$\rho_{B,0} = \rho_B(T_0, \vec{\mu}_0)$$
$$\rho_{S,0} = \rho_S(T_0, \vec{\mu}_0)$$
$$\rho_{Q,0} = \rho_Q(T_0, \vec{\mu}_0)$$

Construct interpolants from table of equation-of-state (e.g., LQCD) data

Couple to multi-dimensional rootfinder (e.g., via GSL library)

Current default functionality of `CCAKE`

Alternative: Delaunay interpolation + k-d trees (see back-up slides)

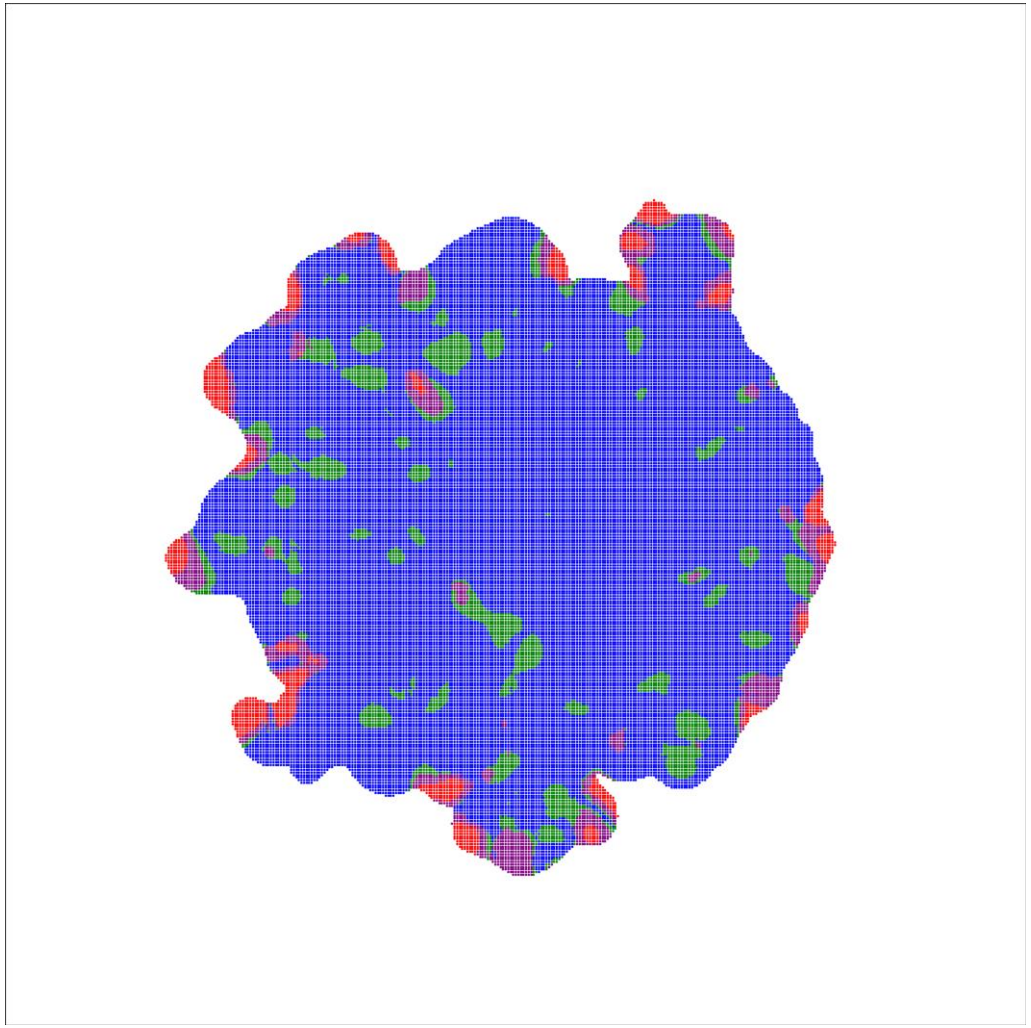
What if the numerical inversion fails?

- Obviously the ideal is that this never happens
- This can happen for any of a number of reasons:
  - The true solution **may exist outside** the current grid
  - There may not be ***any* solution** for the chosen equation of state
  - There may be **multiple "correct" solutions**
- Q: *When this happens in hydrodynamics, how should we close the equations of motion?*
- A: Supplement with an alternative "back-up" equation of state
  - "Tanh-conformal" EoS provides better approximation to lattice at $\mu = 0$
  - Conformal
  - Conformal-diagonal
  - Explicit parametrizations in backup slides

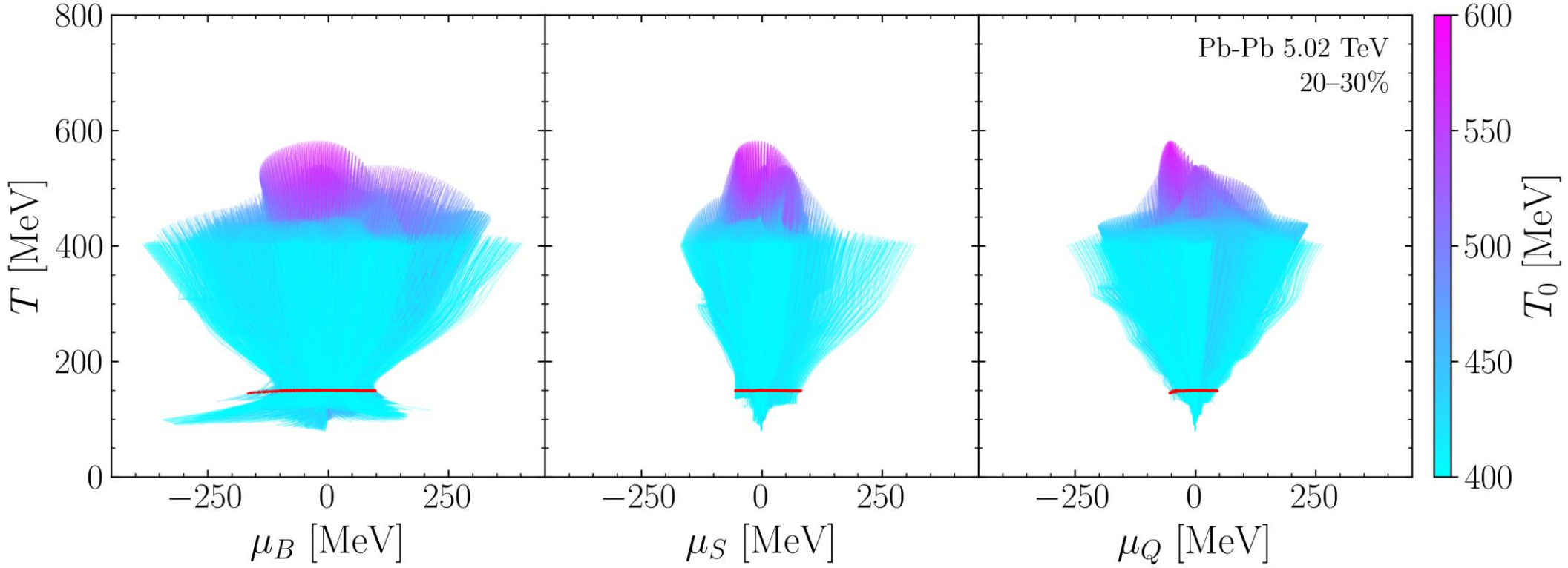
Animation of different types in hydro

- Typical (central) Pb+Pb event showing EoS used in each fluid cell

- **Blue: Lattice QCD**
- **Green: Tanh-conformal**
- **Purple: Conformal**
- **Red: Conformal-diagonal**

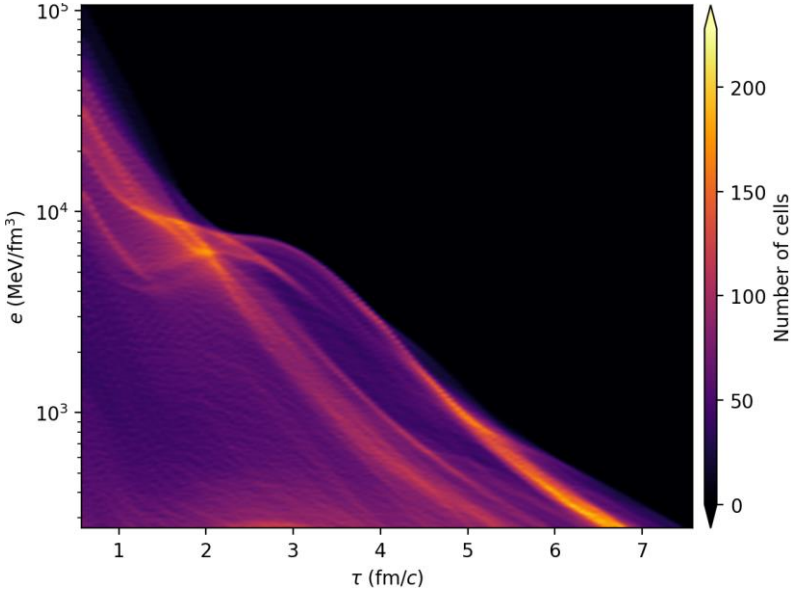


Phase diagram trajectories (20-30%)

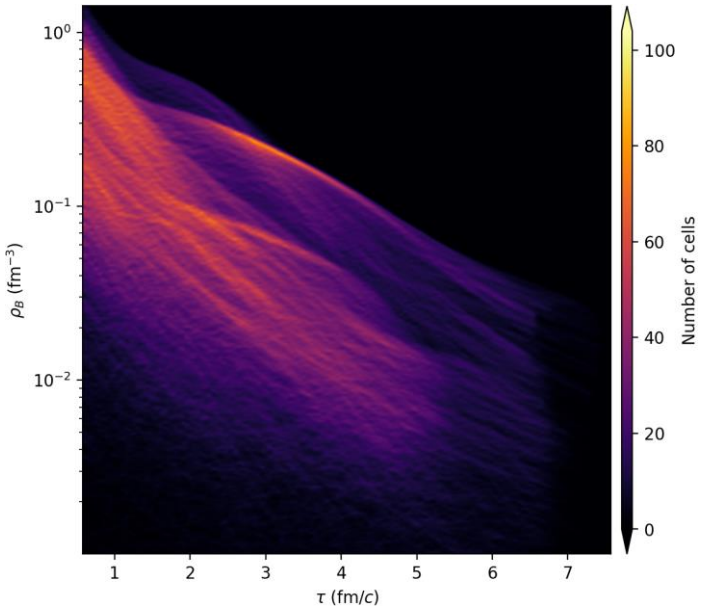


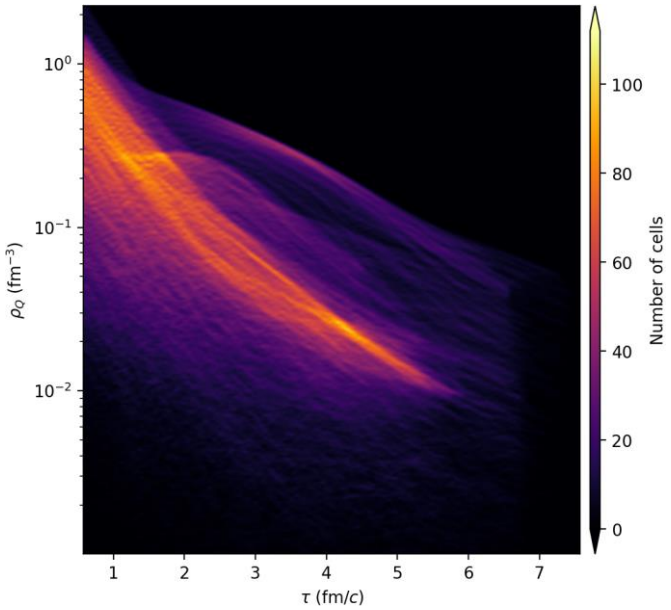
- Fluid cell trajectories in the QCD phase diagram, colored by initial temperature T_0
- Solid red lines represent a constant freeze out energy density $\approx 266\ \mathrm{MeV/fm}^3$

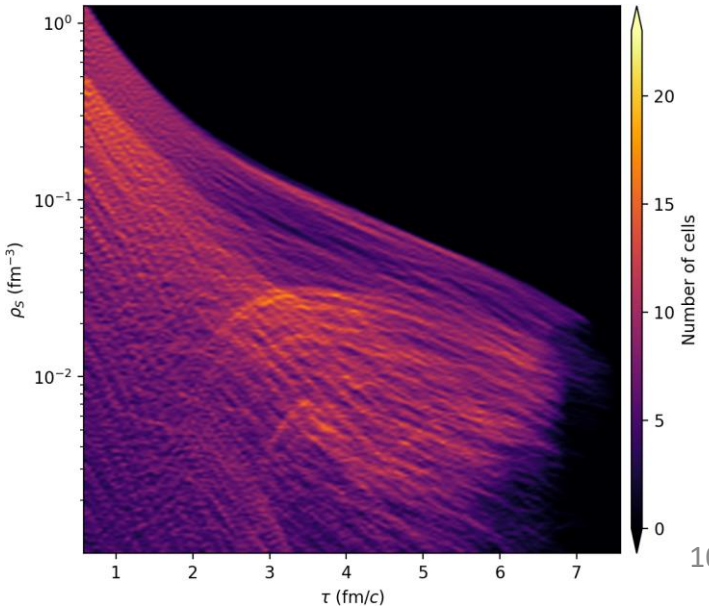
Density distributions (20-30)%



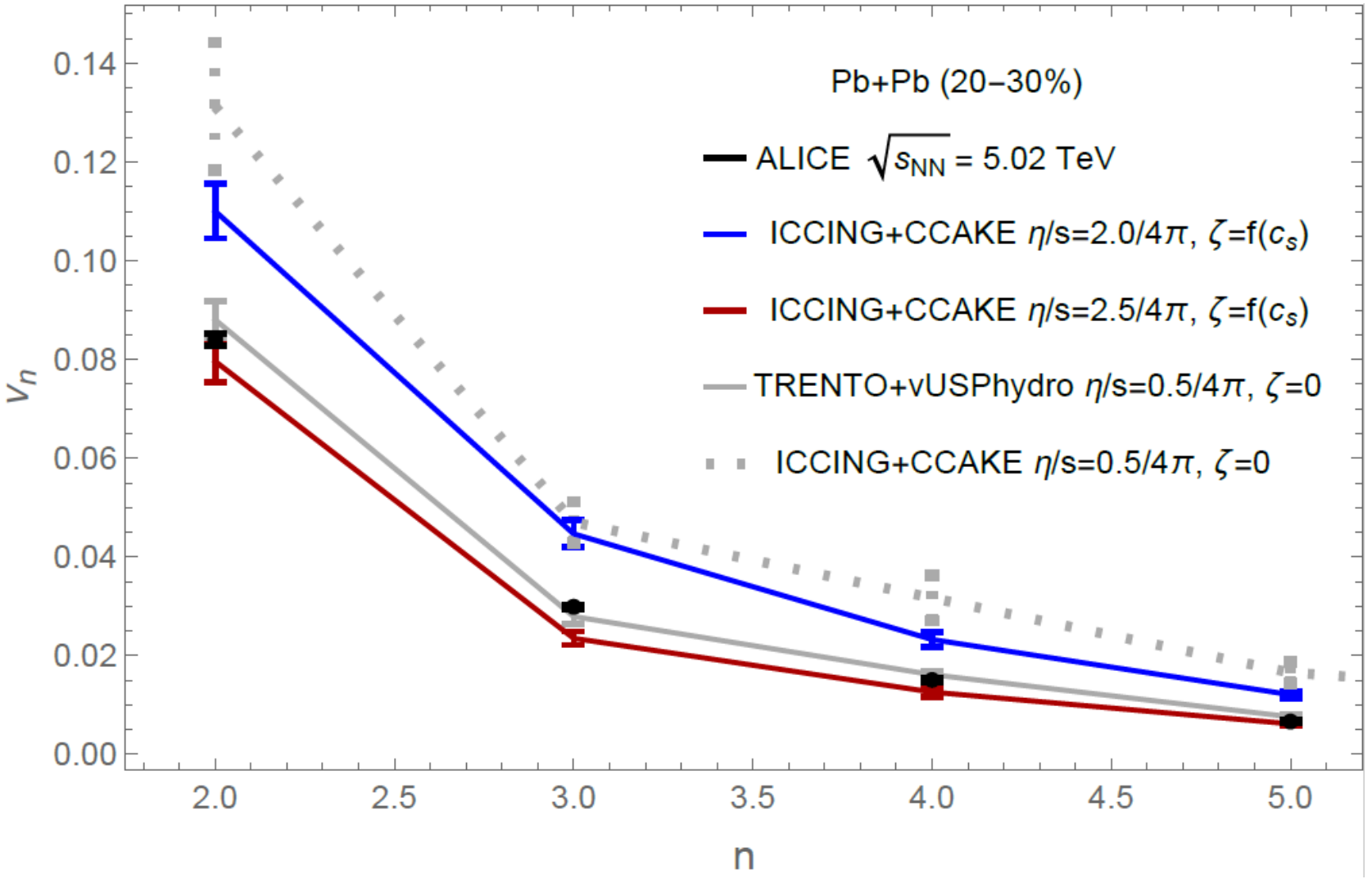
- Particle trajectories sample wide range of chemical potentials and densities, even at LHC energies
- Holds for central and mid-central collisions
- Prospect of constraining wide swath of QCD phase diagram using current (and future) HI experiments







Integrated flow



****Preliminary results****

- Quantitative consistency with ALICE flow measurements
- Calculation of additional observables underway
  - Differential flow, multiplicities, HBT, …
  - Novel charge-dependent observables

Summary

- ***New*** hydrodynamics and initial conditions codes with conserved charges: ICCING + CCAKE
  - Multiple charges require knowledge of multi-dimensional (>=4D) EoS
  - Charge fluctuations can reach large values even at LHC energies, requiring more complete coverage of multi-dimensional space
- Open challenges
  - Finding fast and stable ways to implement multi-dimensional EoS
    - Back-up equations of state
    - Delaunay interpolation + *k*-d trees
  - Improved treatment of BSQ initial conditions, transport coefficients, etc.

Thank you!

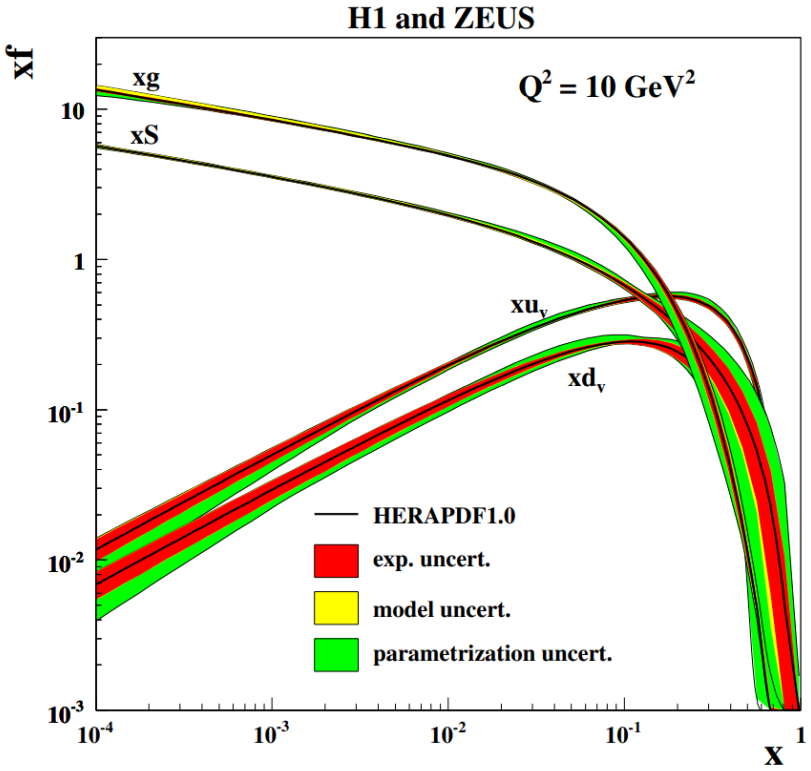
Backup slides

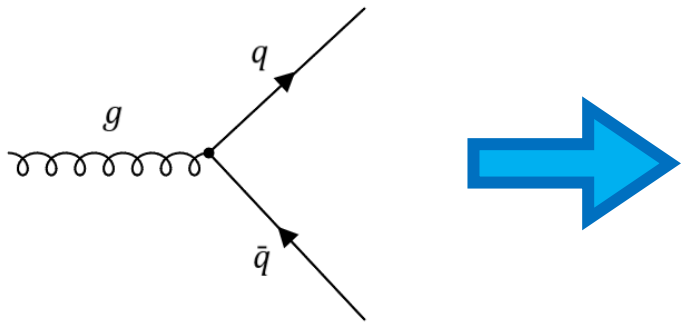
BSQ Initial Conditions - *ICCING*

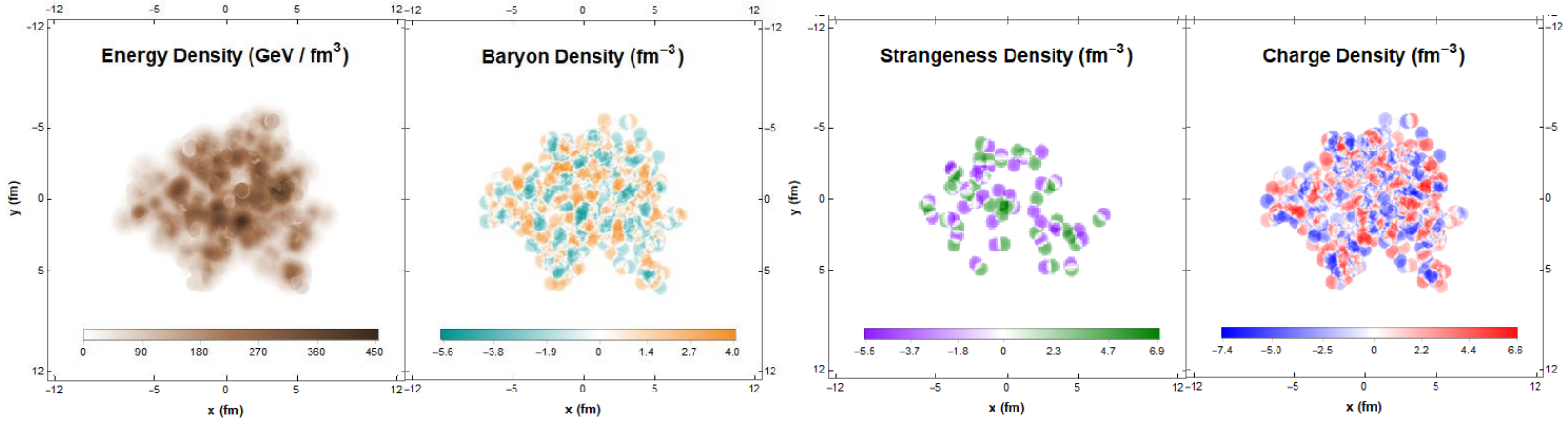
- ICCING initial conditions
  (Initial Conserved Charges In Nuclear Geometry)

- ICCING relies on the gluon-saturated initial state at mid-rapidity to determine probabilities for gluon splitting to quark pairs

- Use color-glass condensate (CGC) framework to generate *local charge fluctuations*







[arXiv:1911.12454 [nucl-th]]

BSQ Evolution - *CCAKE*

Israel-Stewart fluid dynamics

Dekrayat Almaalol, Travis Dore, Jacquelyn Noronha-Hostler [arXiv:2209.11210 [hep-th]]

$$S^\mu = su^\mu - \sum_q^{B,S,Q} \alpha_q n_q^\mu - \frac{1}{2}u^\mu \left(\beta_\Pi \Pi^2 + \beta_\pi \pi^{\mu\nu}\pi_{\mu\nu} + \sum_q^{B,S,Q} \beta_n^{qq'} n_q^\mu n_\mu^{q'} \right) - \sum_q^{B,S,Q} \left(\gamma_{n\Pi}^q n_q^\mu \Pi + \gamma_{n\pi}^q n_q^\nu \pi_\nu^\mu \right) - \frac{1}{2}(u^\nu \beta_{\Pi\pi} \Pi \pi_{\mu\nu})$$

equilibrium | 1^{st}-order term | 2^{nd}-order terms | 2^{nd}-order coupling terms

Second law of thermodynamics

$$\partial_\mu S^\mu = \frac{\beta_0}{2\eta}\pi_{\mu\nu}\pi^{\mu\nu} + \frac{\beta_0}{\zeta}\Pi^2 + \frac{1}{\kappa_{qq'}} n_\mu^q n_{q'}^\mu \geq 0$$

NS Transport coefficients

$$\eta, \zeta, \kappa_{qq'}$$

2^{nd} order Transport coefficients

$$\beta_\Pi, \beta_\pi, \gamma_{n\Pi}, \gamma_{n\pi}, \beta_{qq'}$$

Fotakis et al, 2203.11549 [nucl-th]

CCAKE (Conserved ChArges with hydrodynamiK Evolution)

Slide credit: Dekrayat Almaalol

“Back-up” EoS #1: “Tanh-conformal”

- Definition:

$$p_{\rm tc}(T,\mu_B,\mu_S,\mu_Q) = c\tanh\left(\frac{T-T_c}{T_s}\right)\left(\left(\frac{T}{T_0}\right)^2+\left(\frac{\mu_B}{\mu_{B,0}}\right)^2+\left(\frac{\mu_S}{\mu_{S,0}}\right)^2+\left(\frac{\mu_Q}{\mu_{Q,0}}\right)^2\right)^2$$

- Scale parameters determined to mimic tabulated EoS at high T as closely as possible:

$$\begin{aligned}
c &\equiv p_{\rm table}(T_{\rm max},0,0,0)/T_{\rm max}^4 \\
T_0 &\equiv 1 \\
p_{\rm c}(T_{\rm max},\mu_{B,\rm max},0,0;\mu_{B,0}) &\equiv p_{\rm table}(T_{\rm max},\mu_{B,\rm max},0,0) \\
p_{\rm c}(T_{\rm max},0,\mu_{S,\rm max},0;\mu_{S,0}) &\equiv p_{\rm table}(T_{\rm max},0,\mu_{S,\rm max},0) \\
p_{\rm c}(T_{\rm max},0,0,\mu_{Q,\rm max};\mu_{Q,0}) &\equiv p_{\rm table}(T_{\rm max},0,0,\mu_{Q,\rm max})
\end{aligned}$$

- Two additional parameters in tanh() chosen to mimic transition to HRG:

$$T_c = 220\ \text{MeV},\ T_s = 120\ \text{MeV}$$

“Back-up” EoS #2: “Conformal”

- Definition:

$$p_{\rm c}(T,\mu_B,\mu_S,\mu_Q) = c\left(\left(\frac{T}{T_0}\right)^2 + \left(\frac{\mu_B}{\mu_{B,0}}\right)^2 + \left(\frac{\mu_S}{\mu_{S,0}}\right)^2 + \left(\frac{\mu_Q}{\mu_{Q,0}}\right)^2\right)^2,$$

- Not the most general (any quartic combinations are acceptable)
- Scale parameters determined as in “Tanh-conformal”
- Overall factor c determined by

$$c \equiv \frac{\pi^2}{90}\left(2\left(N_c^2-1\right) + \frac{7}{2}N_cN_f\right)$$

where

$$N_c = 3 \text{ and } N_f = 2.5$$

“Back-up” EoS #3: “Conformal-diagonal”

- Definition:

$$p_{\rm cd}(T,\mu_B,\mu_S,\mu_Q) = c\left(\left(\frac{T}{T_0}\right)^4 + \left(\frac{\mu_B}{\mu_{B,0}}\right)^4 + \left(\frac{\mu_S}{\mu_{S,0}}\right)^4 + \left(\frac{\mu_Q}{\mu_{Q,0}}\right)^4\right),$$

- Scale parameters determined as in “Tanh-conformal” and “Conformal”
- Overall factor c same as “Conformal”
- One can prove

$$e \geq e_{\rm min}\left(\vec{\rho}\right) = \frac{3}{4\cdot 2^{2/3} c^{1/3}}\left(\left(\mu_{B,0}\left|\rho_B\right|\right)^{4/3} + \left(\mu_{S,0}\left|\rho_S\right|\right)^{4/3} + \left(\mu_{Q,0}\left|\rho_Q\right|\right)^{4/3}\right)$$

  is a necessary and sufficient condition for given set of (e,ρ_B,ρ_S,ρ_Q) to have a real solution

- If one propagates (s,ρ_B,ρ_S,ρ_Q), then a real solution is always guaranteed

Code demo: Delaunay $(e, \rho_B, \rho_S, \rho_Q)$ interpolator

```
//=============================================
// read path to input file from command line
string path_to_file = string(argv[1]);

//=============================================
// set up EoS object
cout << "Initializing equation of state "
                "interpolator:" << endl;
cout << "  --> reading in equation of state "
                "table from: " << path_to_file << endl;
eos_delaunay EoS( path_to_file );

//=============================================
// vectors to store input densities and
// interpolated result for (T,muB,muQ,muS)
vector<double> result(4, 0.0);
vector<double> point({ 5754.35,  0.00231029,
                       0.351709, 0.378919   }); // just an example
```

Code demo: Delaunay $(e, \rho_B, \rho_S, \rho_Q)$ interpolator

```
//=============================================
// call the interpolator
const size_t n_repeat = 1000;
cout << "Calling the interpolator "
     << n_repeat << " times for test point: \n"
                 "  --> {e,B,S,Q} = {"
     << point[0] << " MeV/fm^3," << point[1] << " 1/fm^3, "
     << point[2] << " 1/fm^3, "  << point[3] << " 1/fm^3}" << endl;

// multiple calls to improve timing estimate
for (size_t i = 0; i < n_repeat; i++)
  EoS.interpolate(point, result);
```

Invocation: `$ ./interpolate_ebsq eos.dat`

```
================================================================================
= Code:    Equation of state interpolator
= Purpose: Performance and closure tests of Delaunay interpolator
= Author:  Christopher Plumberg
= Contact: plumberg@illinois.edu
= Date:    April 28, 2022
================================================================================

Initializing equation of state interpolator:
  --> reading in equation of state table from: eos.dat
        - read in 1000000 lines.
        - read in 2000000 lines.
        - read in 3000000 lines.
        - read in 4000000 lines.
        - read in 5000000 lines.
        - read in 6000000 lines.
        - read in 7000000 lines.
  --> check minima and maxima:
        - e:    0.0151087    729992      [MeV/fm^3]
        - B:    -25.56       25.56       [1/fm^3]
        - S:    -65.1234     65.1234     [1/fm^3]
        - Q:    -37.5927     37.5927     [1/fm^3]
  --> setting up kd-trees: finished in 7.63744 seconds!

Calling the interpolator 1000 times for test point:
  --> {e,B,S,Q} = {5754.35 MeV/fm^3,0.00231029 1/fm^3, 0.351709 1/fm^3, 0.378919 1/fm^3}
  --> exact result (units MeV):          252.5    52.5    52.5    52.5
  --> interpolated result (units MeV):   252.448   52.5715   52.571   52.5597

Average time to interpolate: 0.00512986 s.
Total runtime: 56.6377 s.
```

Moderate grid size

Reasonable accuracy

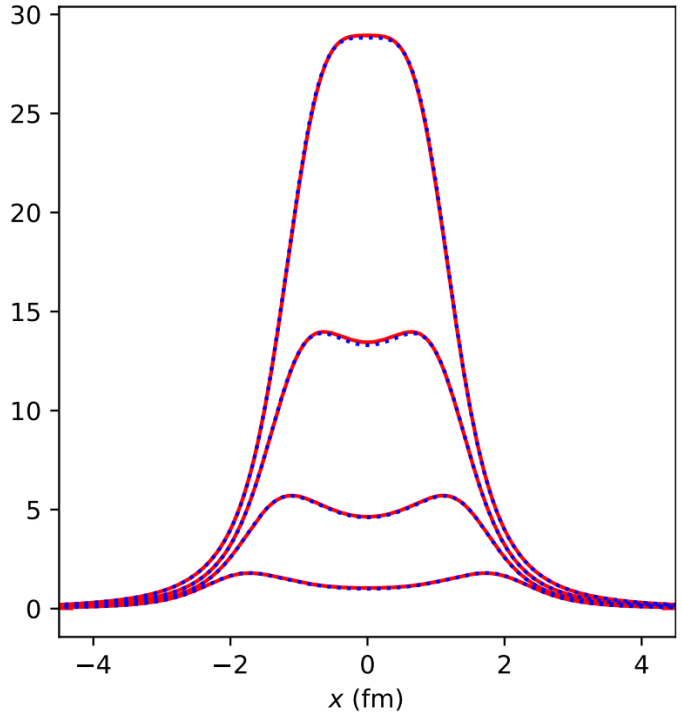
$\sim$ 200 solutions/s

Gubser checks

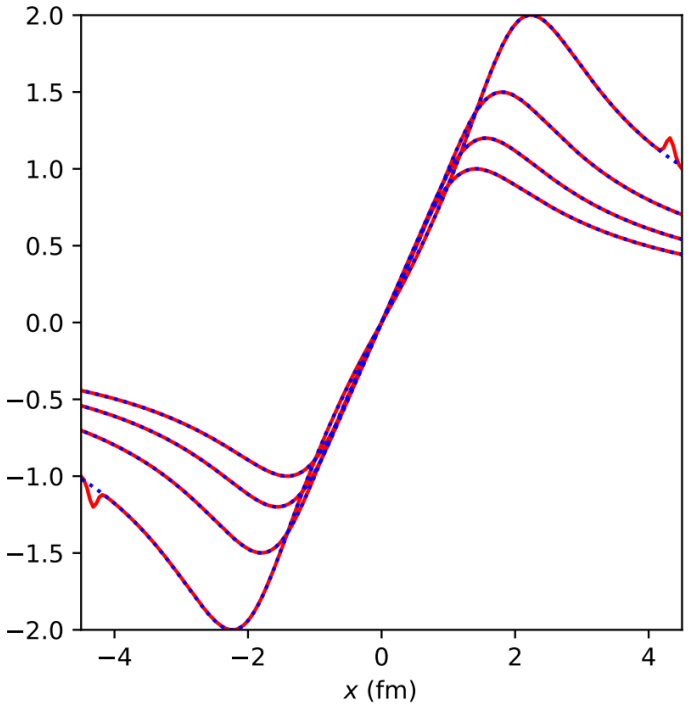
$\tau = 1.0,\ 1.2,\ 1.5,\ 2.0\ \mathrm{fm}/c$

Blue (dotted): exact

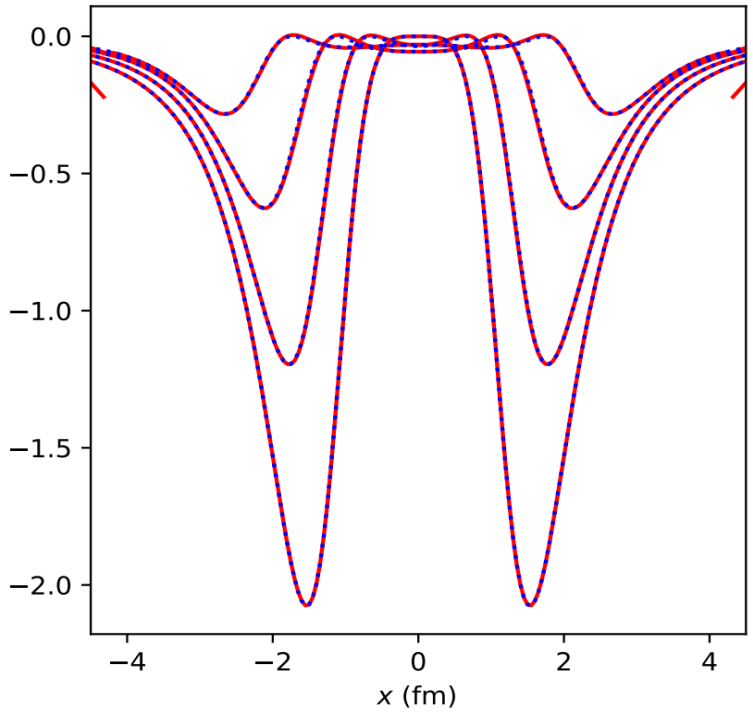
Red (solid): CCAKE



Energy density: e (fm^{-4})



Flow velocity: u^r



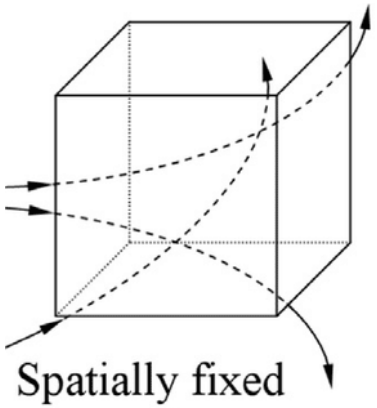
Shear stress: π^{xx} (fm^{-4})

Smoothed Particle Hydrodynamics (SPH)

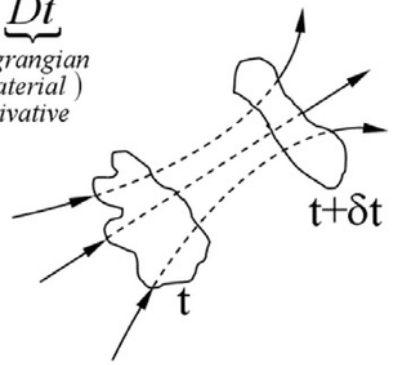
Eulerian

$$\underbrace{\frac{\partial}{\partial t}}_{\text{Eulerian derivative}} + \mathbf{v}\cdot\nabla = \underbrace{\frac{D}{Dt}}_{\text{Lagrangian (Material) derivative}}$$

Lagrangian



Spatially fixed volume element

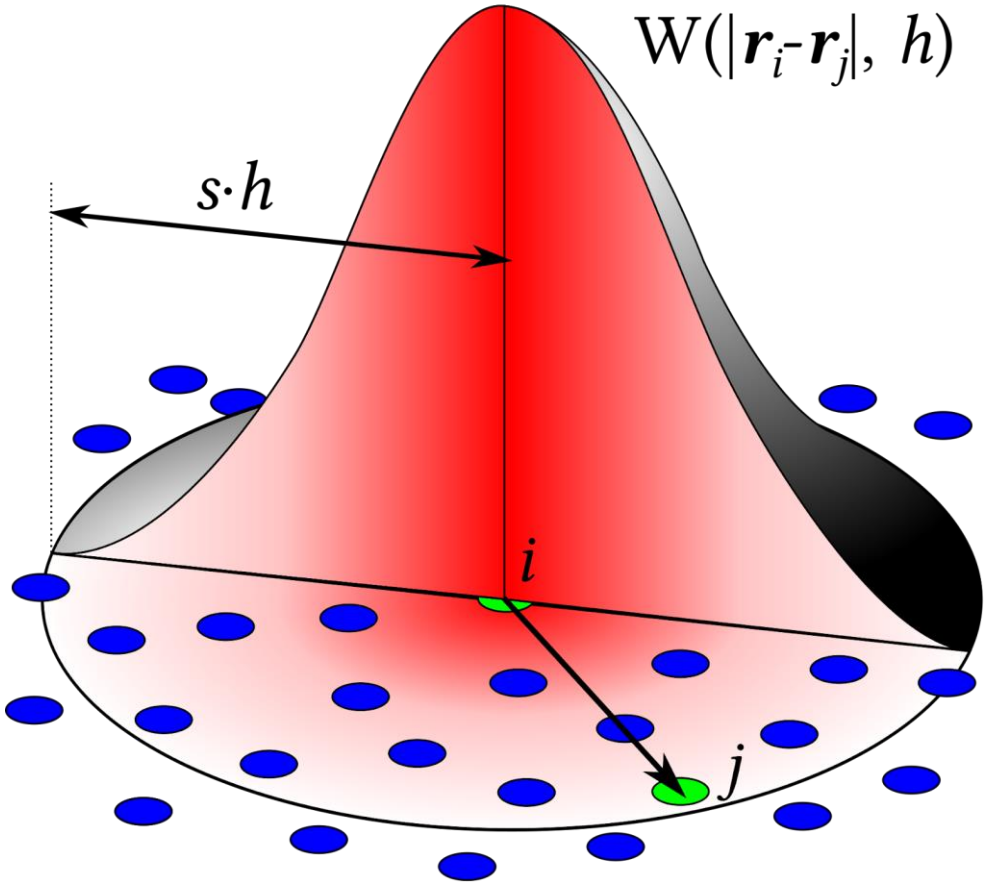


Following the motion of the fluid element

Grid-based hydrodynamics

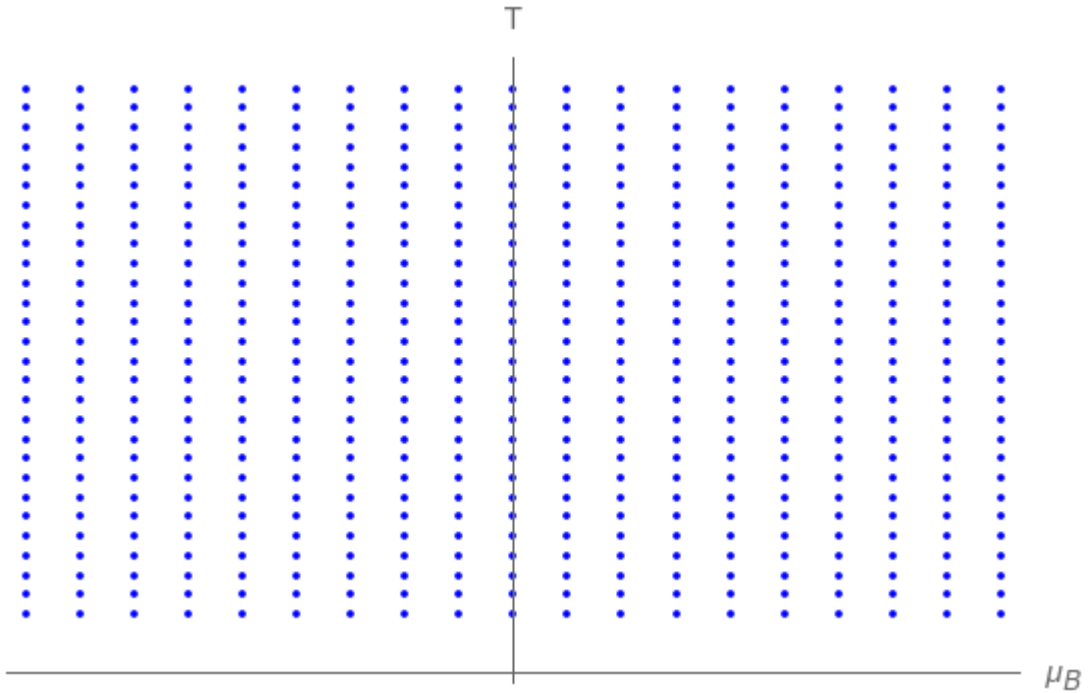
Smoothed particle hydrodynamics

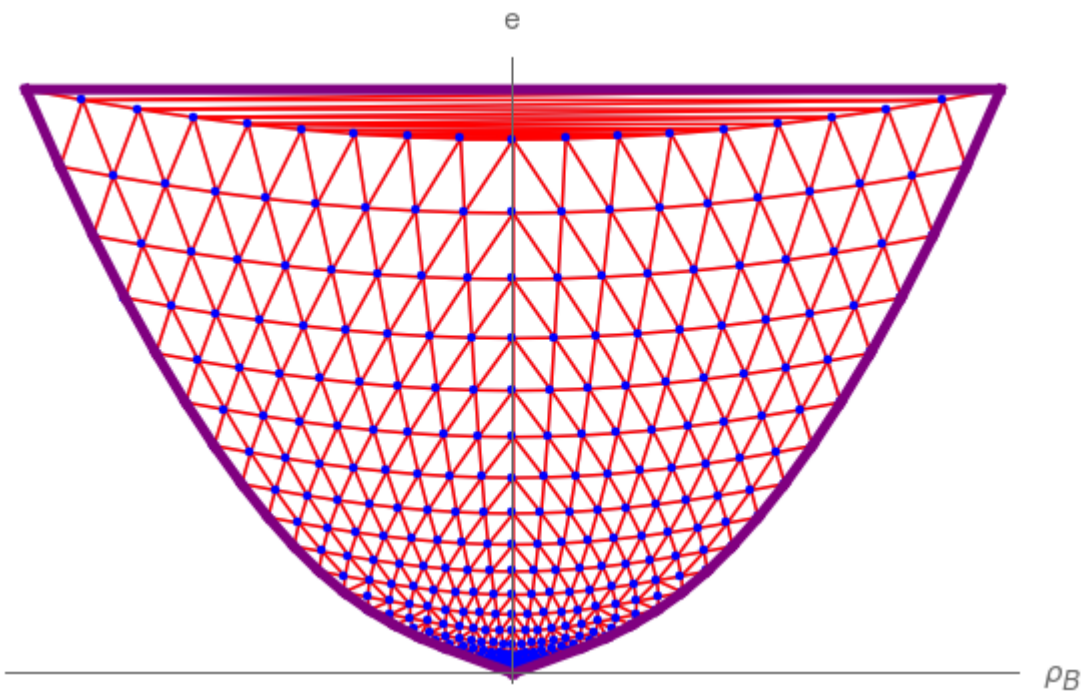
Kernel function W imposes coarse-graining onto set of fictitious 'SPH particles'



Conservation laws built-in by construction

Strategy #2: Delaunay interpolation



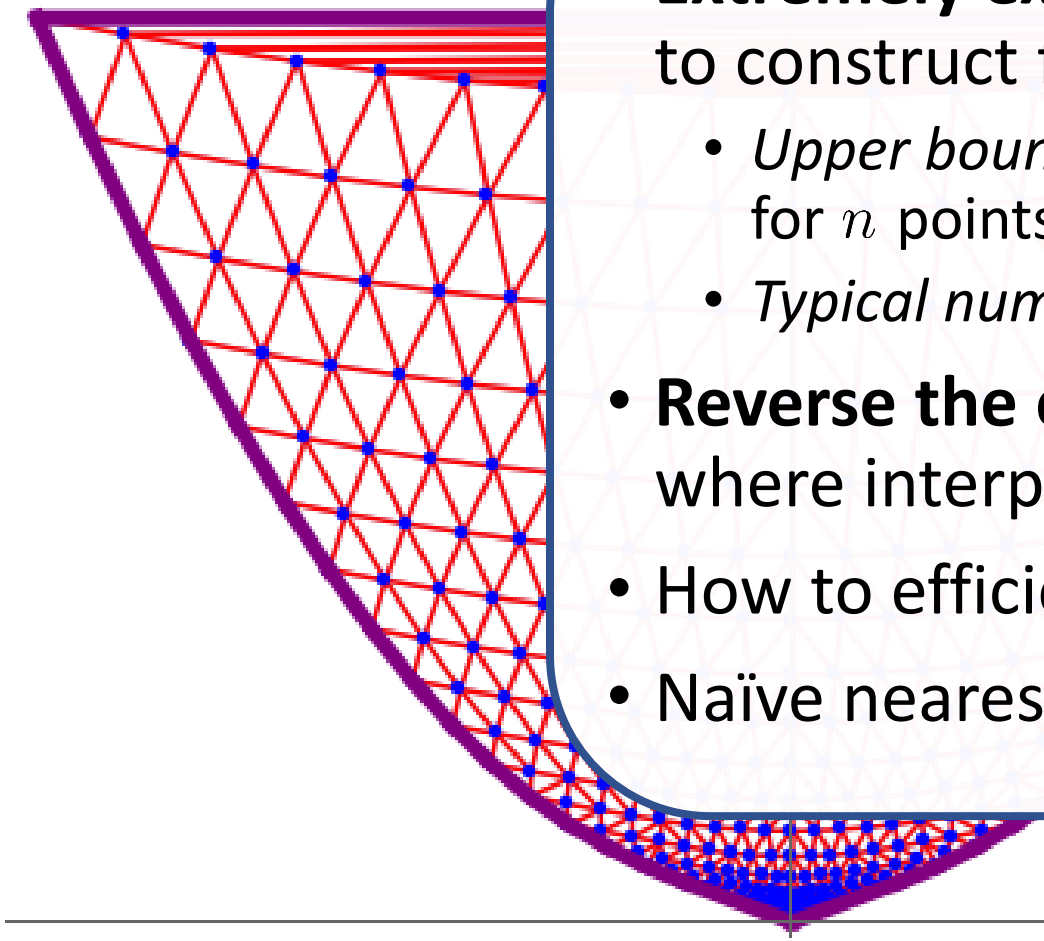


$$\text{Uniform } T\text{–}\mu_B \text{ grid} \quad \not\Rightarrow \quad \text{uniform } e\text{–}\rho_B \text{ grid}$$

– Perform linear interpolation on Delaunay triangulation of scattered density points

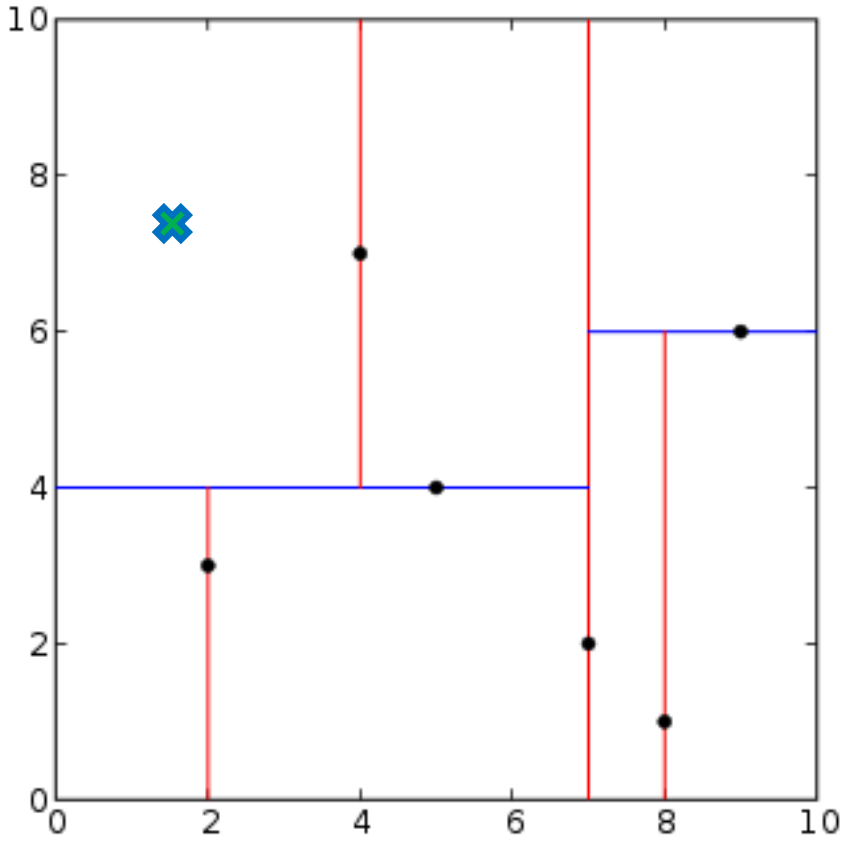
– Only defined inside **convex hull** (bold line)

Constructing the Delaunay mesh

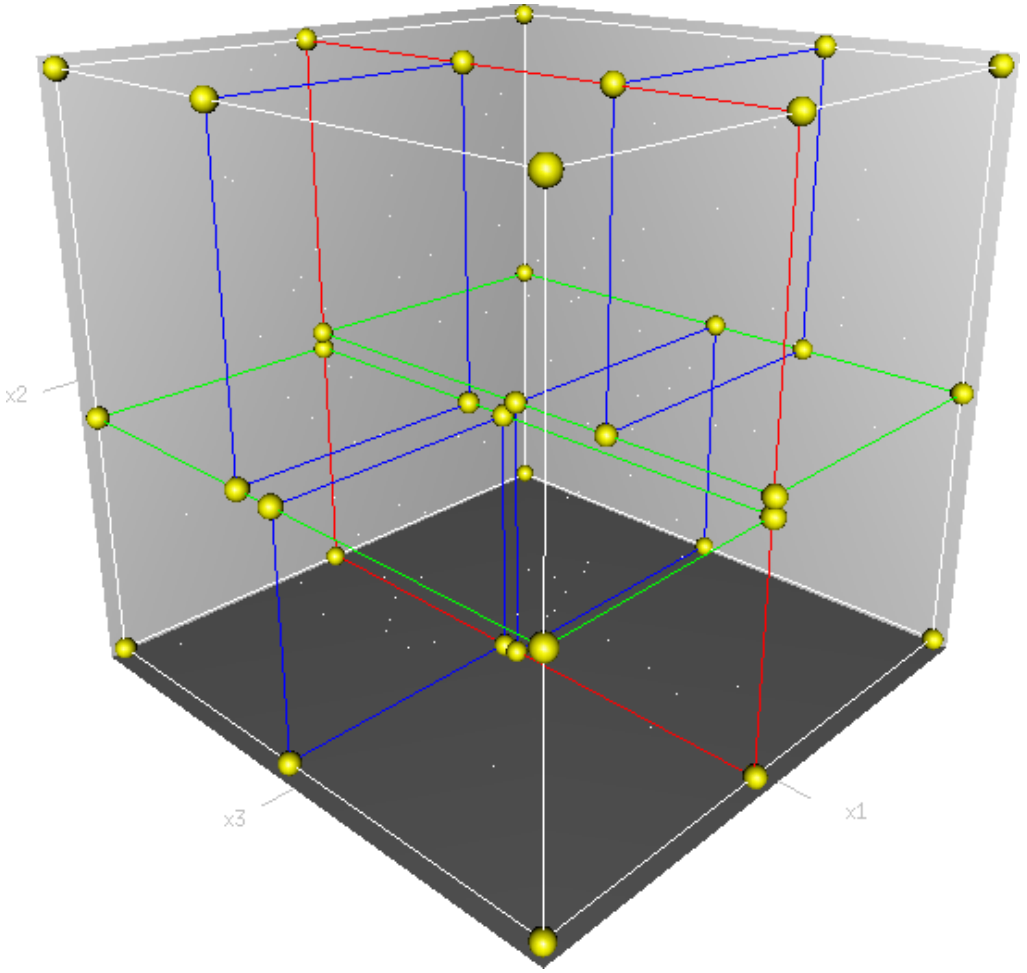


- **Extremely expensive** (memory/CPU time) to construct full mesh of EoS in advance
  - *Upper bound* on number of simplices grows like $O\left(n^{\lceil d/2 \rceil}\right)$, for n points in d dimensions ("curse of dimensionality")
  - *Typical number* of EoS points in modest grid: $O\left(10^5 - 10^7\right)$ in 4D
- **Reverse the curse**: only triangulate the region where interpolation is needed, evaluated at runtime
- How to efficiently find the right region to triangulate?
- Naïve nearest-neighbor look-up may be *very* inefficient

Finding closest simplex efficiently: *k*-d trees



2D example



3D example

A rough algorithm

Compute (T, μ_B, μ_S, μ_Q) distributions on scattered $(e, \rho_B, \rho_S, \rho_Q)$ grids

Identify convex hull inside of which density interpolation is defined

Build k-d trees of density grids

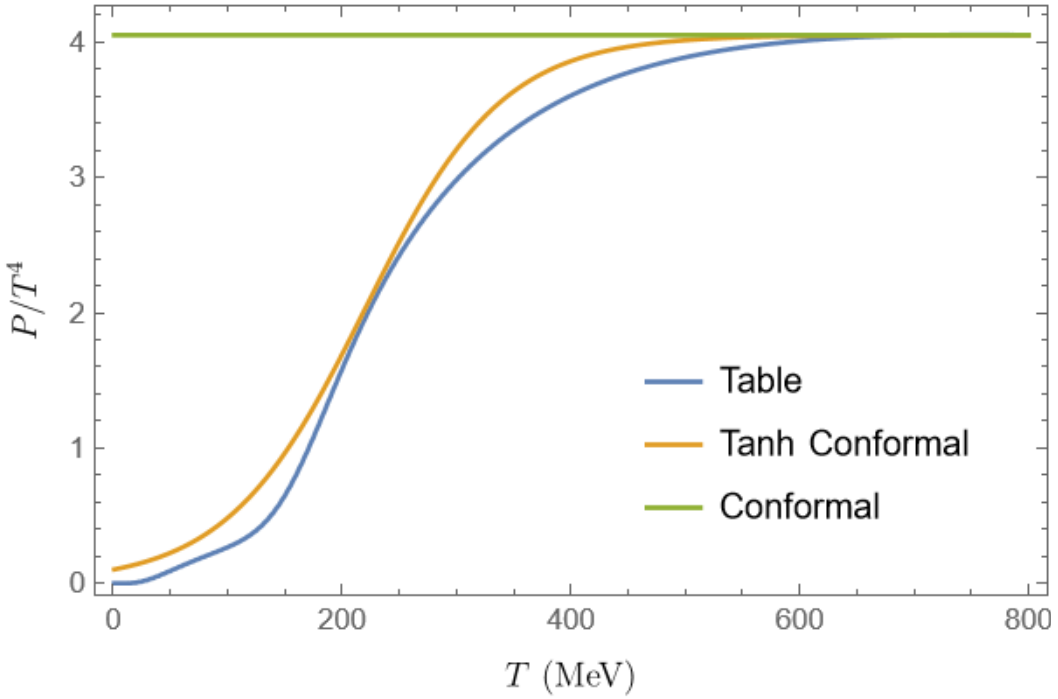
For given densities $(e_0, \rho_{B,0}, \rho_{S,0}, \rho_{Q,0})$:

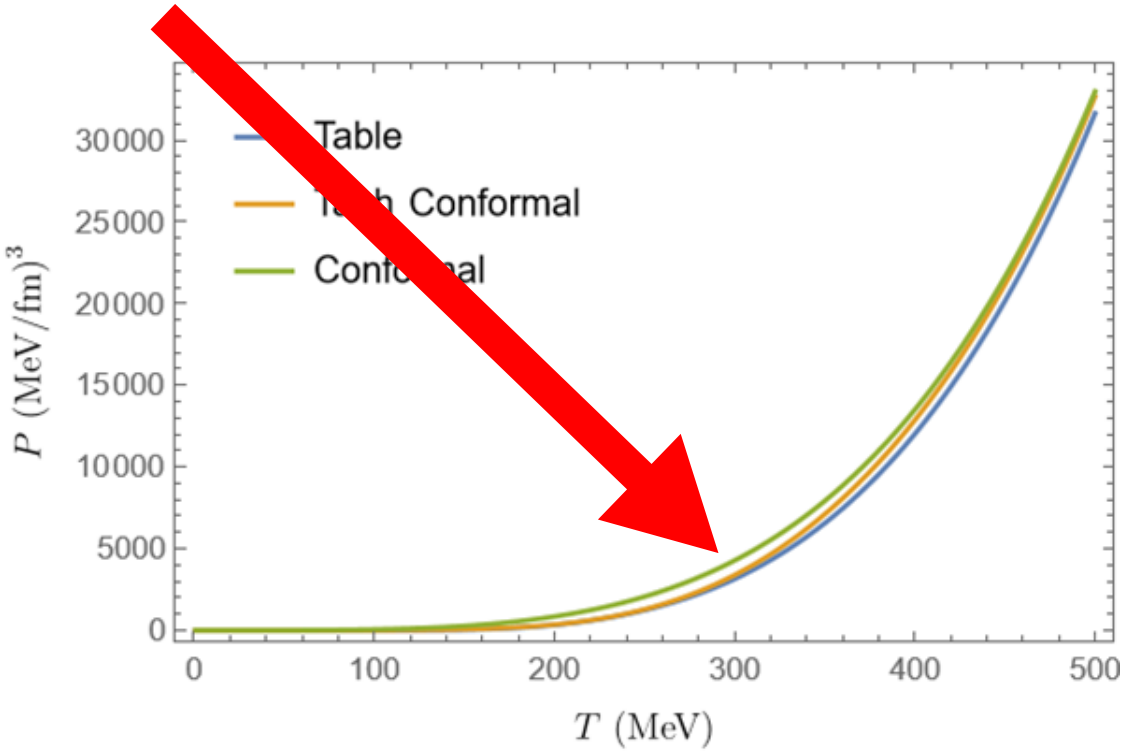
- *locate* containing / neighboring simplices
- *construct* Delaunay triangulation
- *evaluate* unique linear interpolant at input densities

Download: `https://github.com/astrophysicist87/eos_delaunay_demo` (see backup slides)

Comparison: lattice vs. “back-up” EoSs

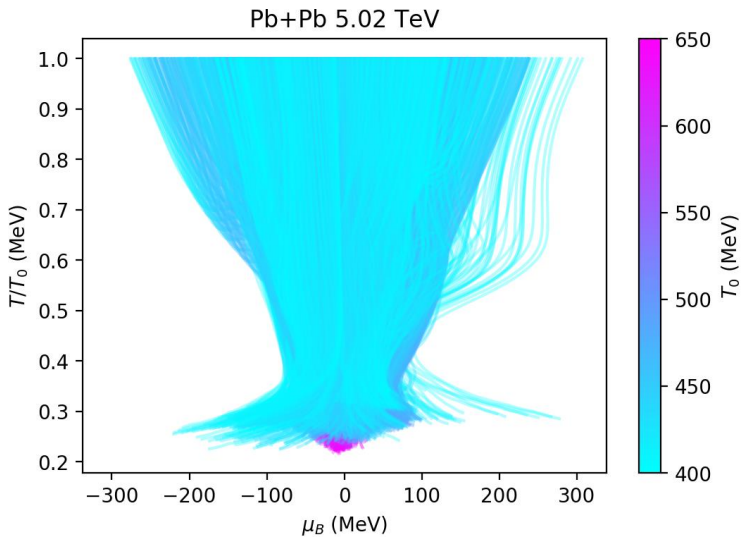
Switching to back-up EoSs produces small violations of energy conservation

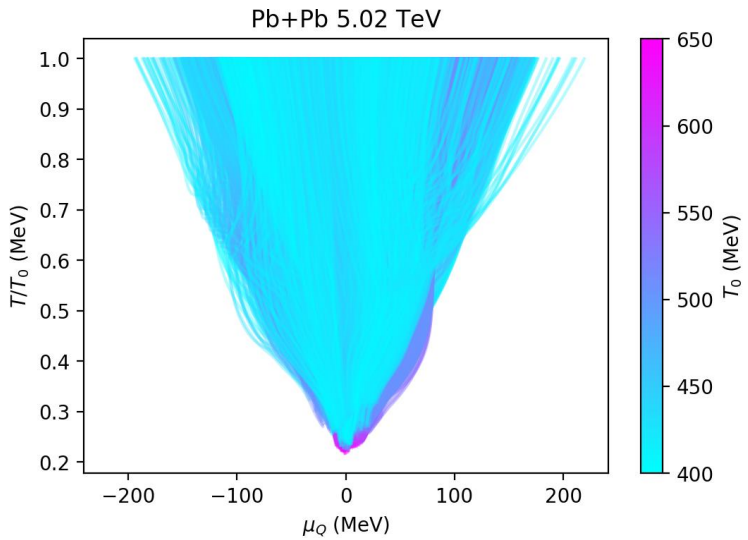


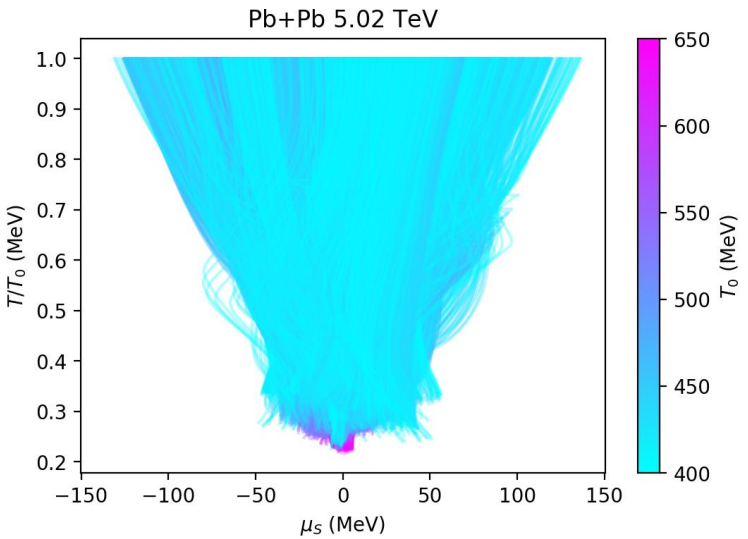


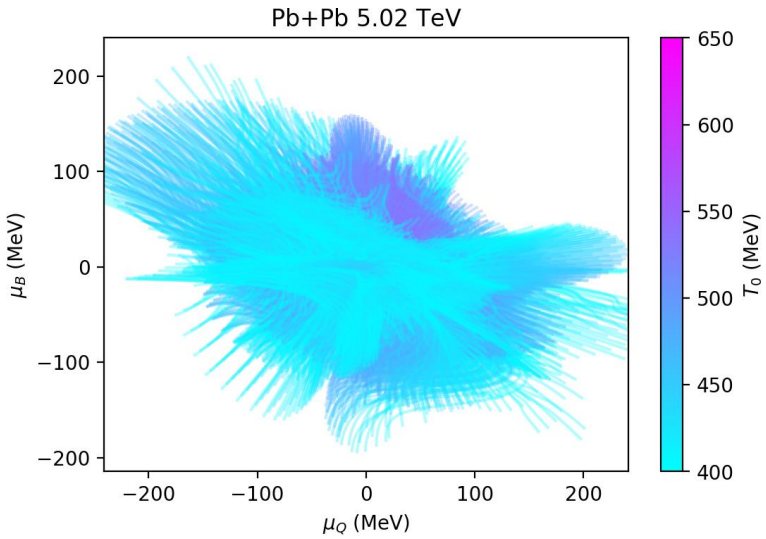
- Conformal-diagonal reduces to Conformal when $\mu_{\mathrm{B,S,Q}} = 0$
- Total energy depends on both energy and pressure
- Total integrated violations below ~0.5%

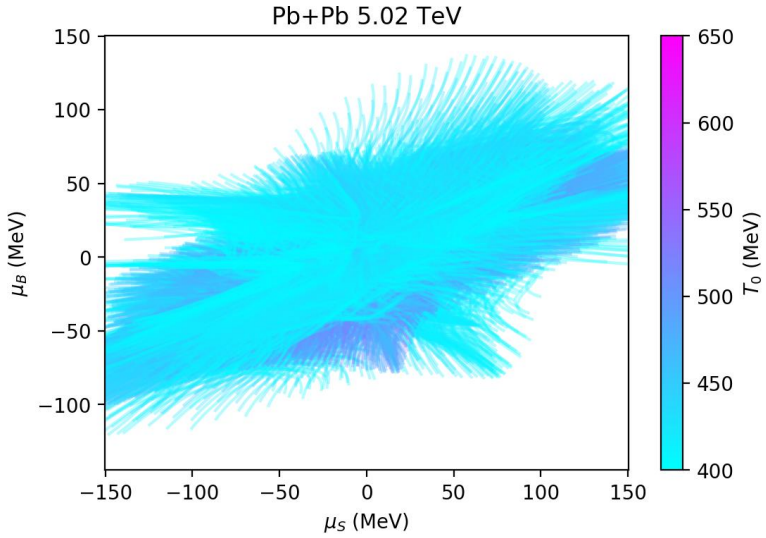
Phase diagram trajectories (0-5%)

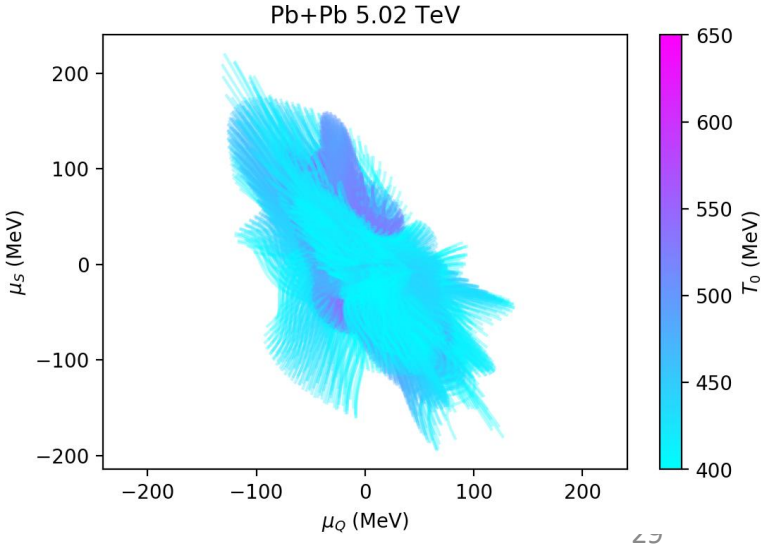




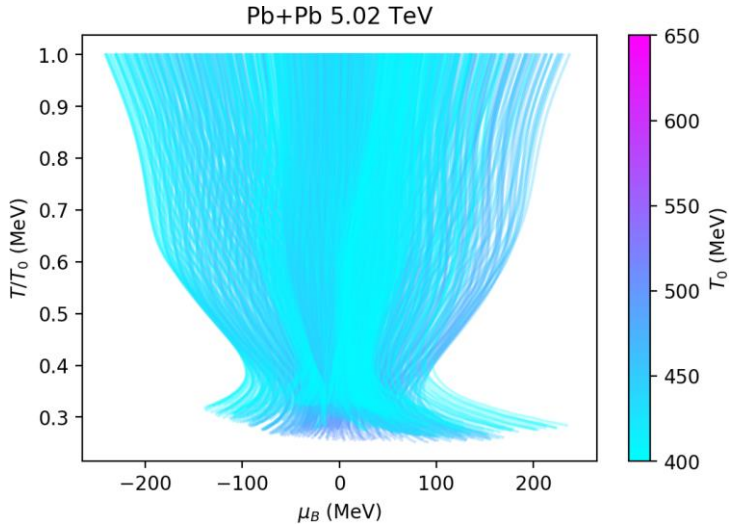


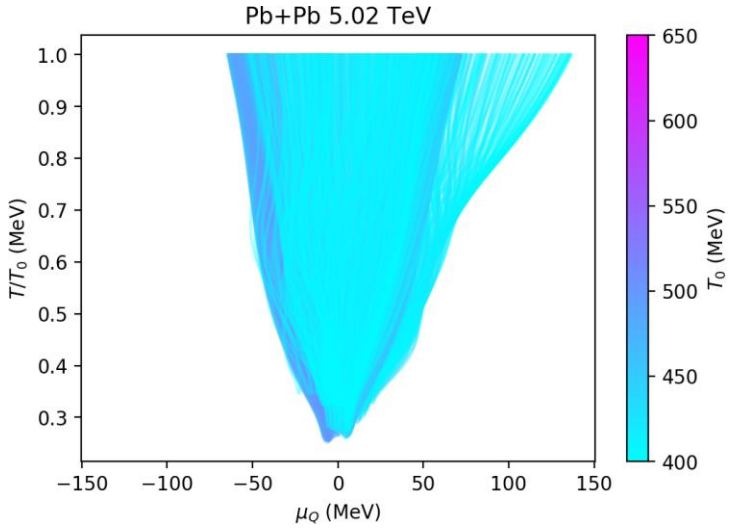


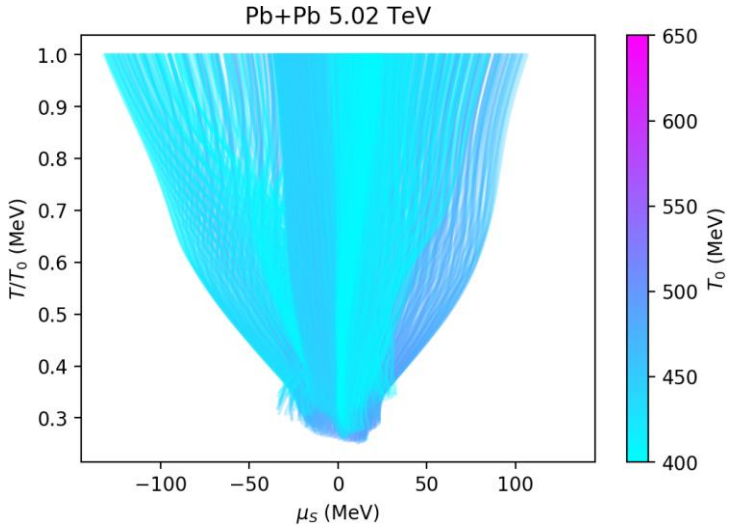


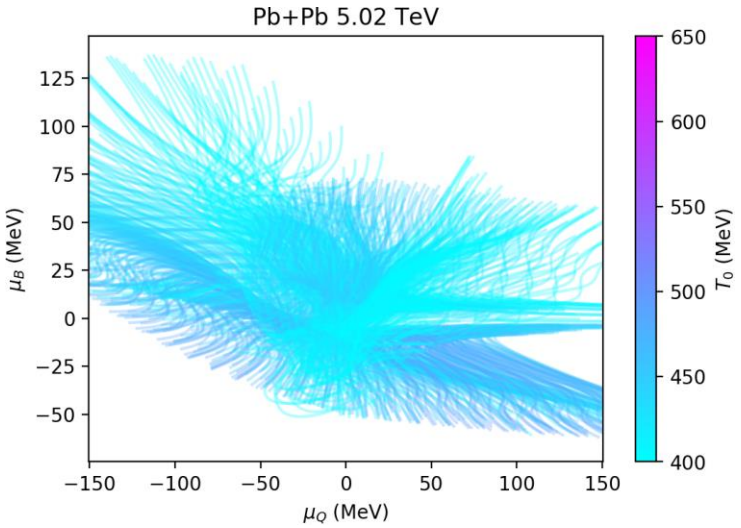


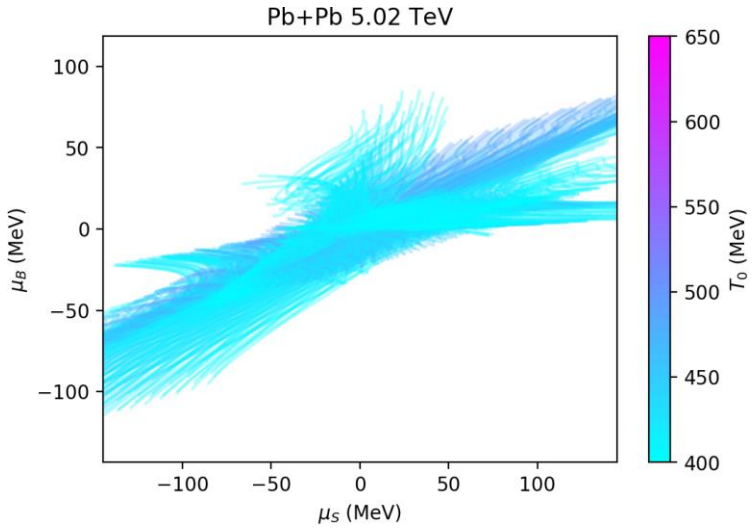
Phase diagram trajectories (20-30%)

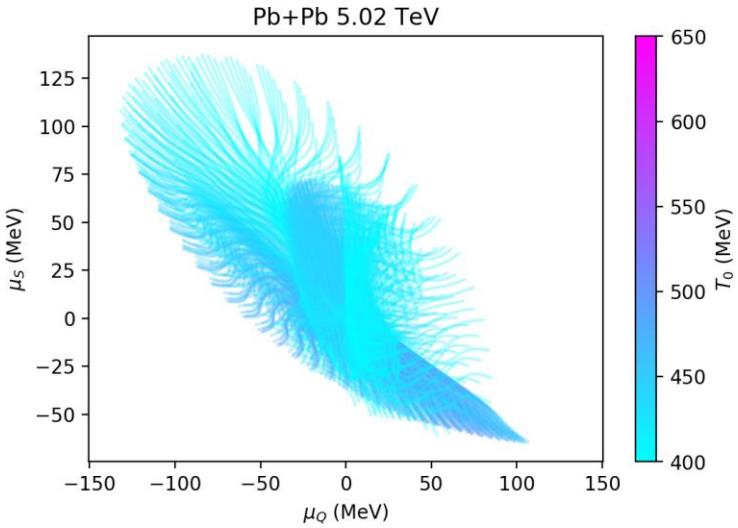












Density distributions (0-5)%