



Light Thermal Self-Interacting Dark Matter in the Shadow of Non-Standard Cosmology

Shu-Yu HO (KIAS)

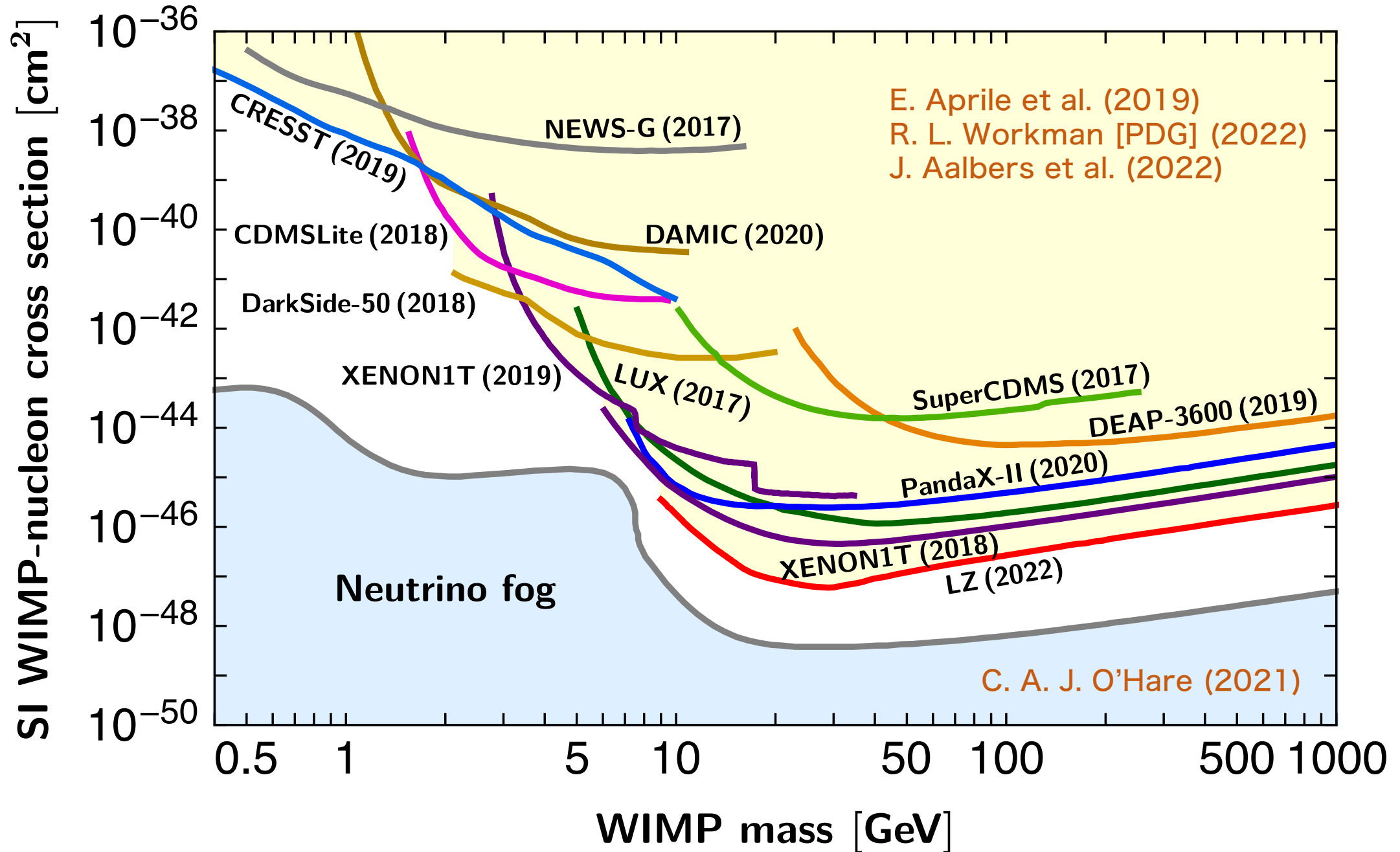
arXiv : 2310.05676

In collaboration with Prof. P. Ko & Dr. N. Dibyendu (KIAS)

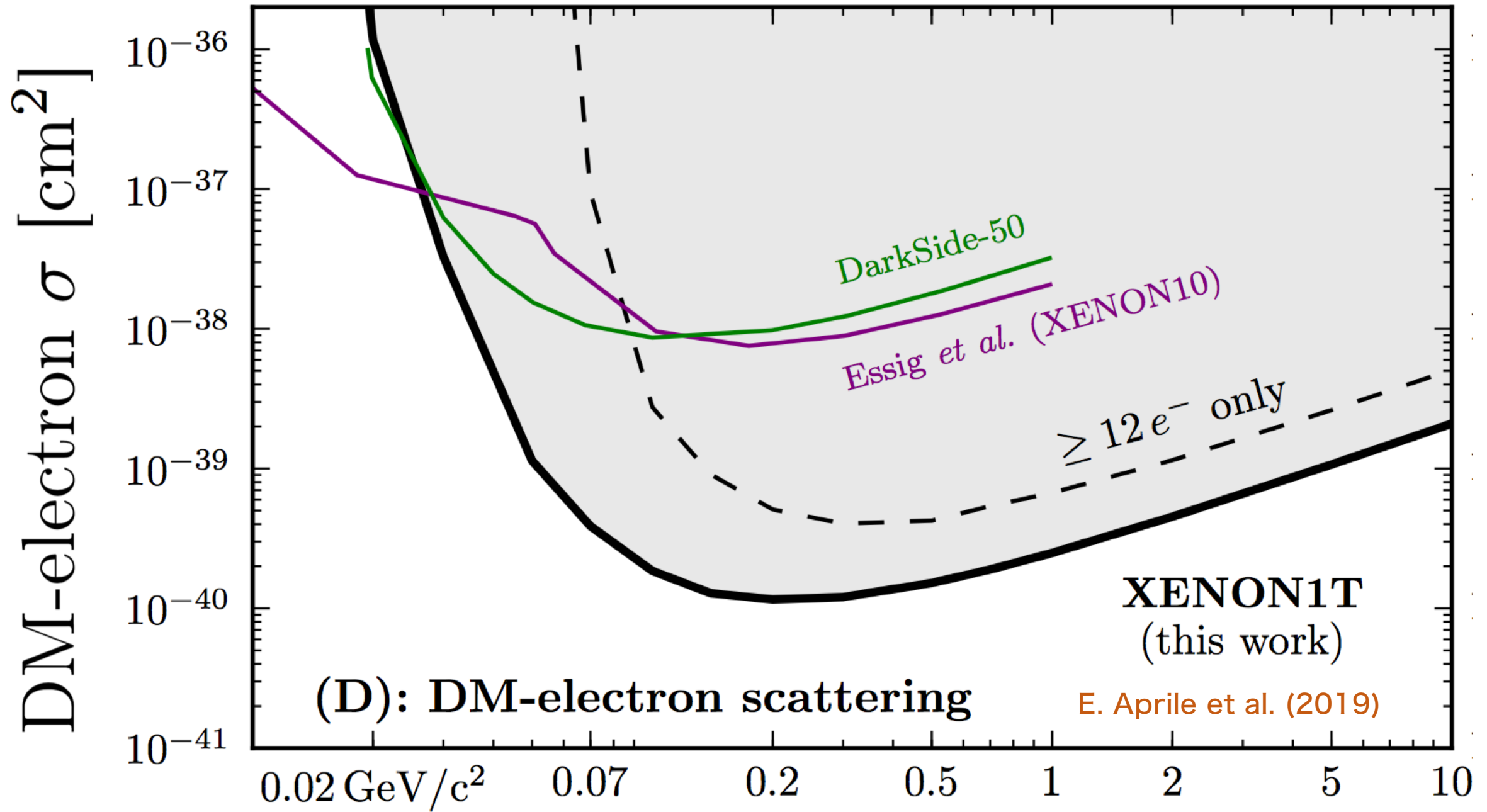
06/Nov/2023

BSM-2023

WIMP Dark Matter (DM) direct searches

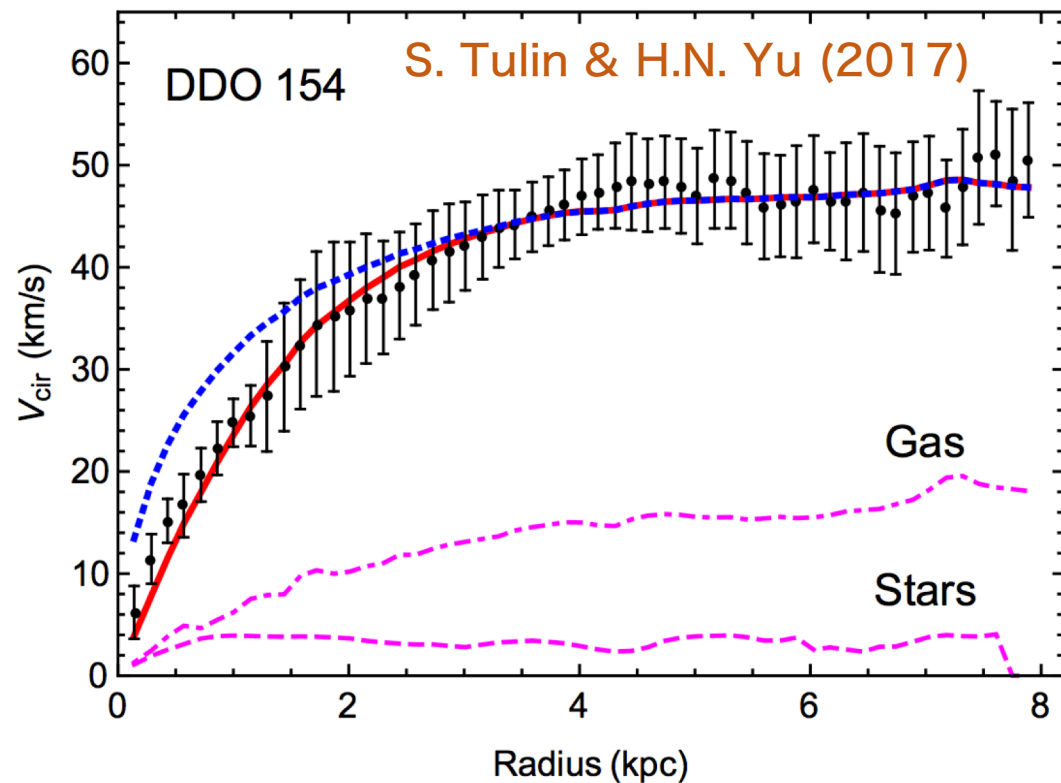


Current experiments of light DM detections

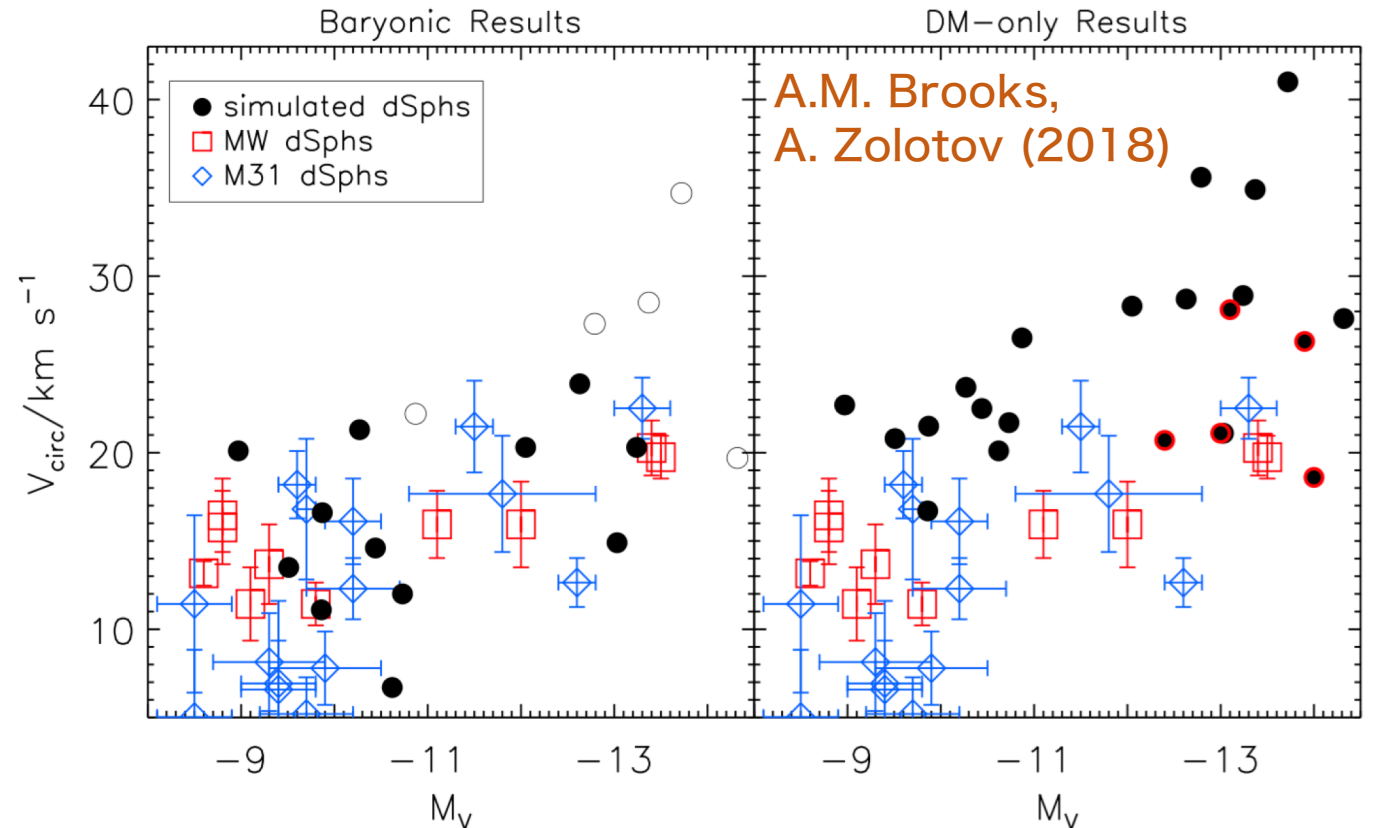


Issues of small scale structures (< 1 Mpc)

- Discrepancy between N-body simulations and observations :



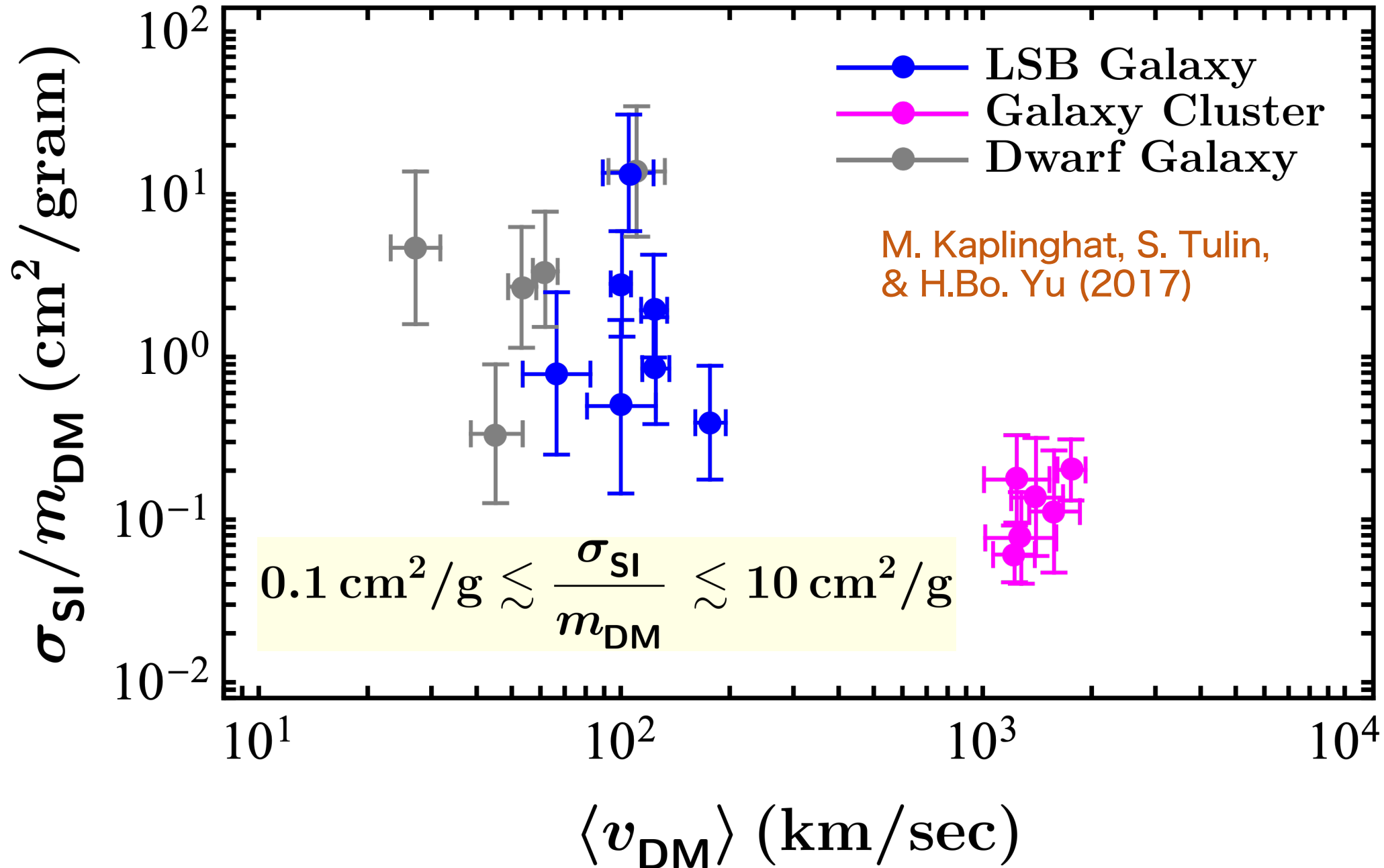
core-vs-cusp problem



too-big-to-fail problem

- DM with a sizable self-interacting (SI) cross-section can resolve these astrophysical problems/issues.

Bounds on DM self-interacting cross-section



**Can we have light thermal
(WIMP) DM with a sizable
self-interacting cross-section?**

WIMP DM

- Relic abundance of WIMP DM

$$\Omega_{\text{WIMP}} h^2 \simeq 0.12 \left(\frac{10^{-8} \text{ GeV}^{-2}}{\langle \sigma v \rangle} \right) \Rightarrow \langle \sigma v \rangle \simeq 10^{-8} \text{ GeV}^{-2}$$

annihilation
cross-section

- Mass scale and coupling strength of WIMP DM

$$\langle \sigma v \rangle = \frac{g^2}{m_{\text{DM}}^2} \Rightarrow g \simeq 10^{-2} \left(\frac{m_{\text{DM}}}{100 \text{ GeV}} \right) \quad (\text{WIMP miracle})$$

g : dimensionless
coupling

$$\simeq 10^{-3} \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right) \quad (\text{Our work})$$

WIMP DM

- SI cross-section via a contact-interaction

$$\left. \frac{\sigma_{\text{SI}}}{m_{\text{DM}}} \right|_{\text{obs}} \simeq 1 \text{ cm}^2/\text{g} \simeq 4.6 \times 10^3 \text{ GeV}^{-3}$$

SIMP, Forbidden DM,...

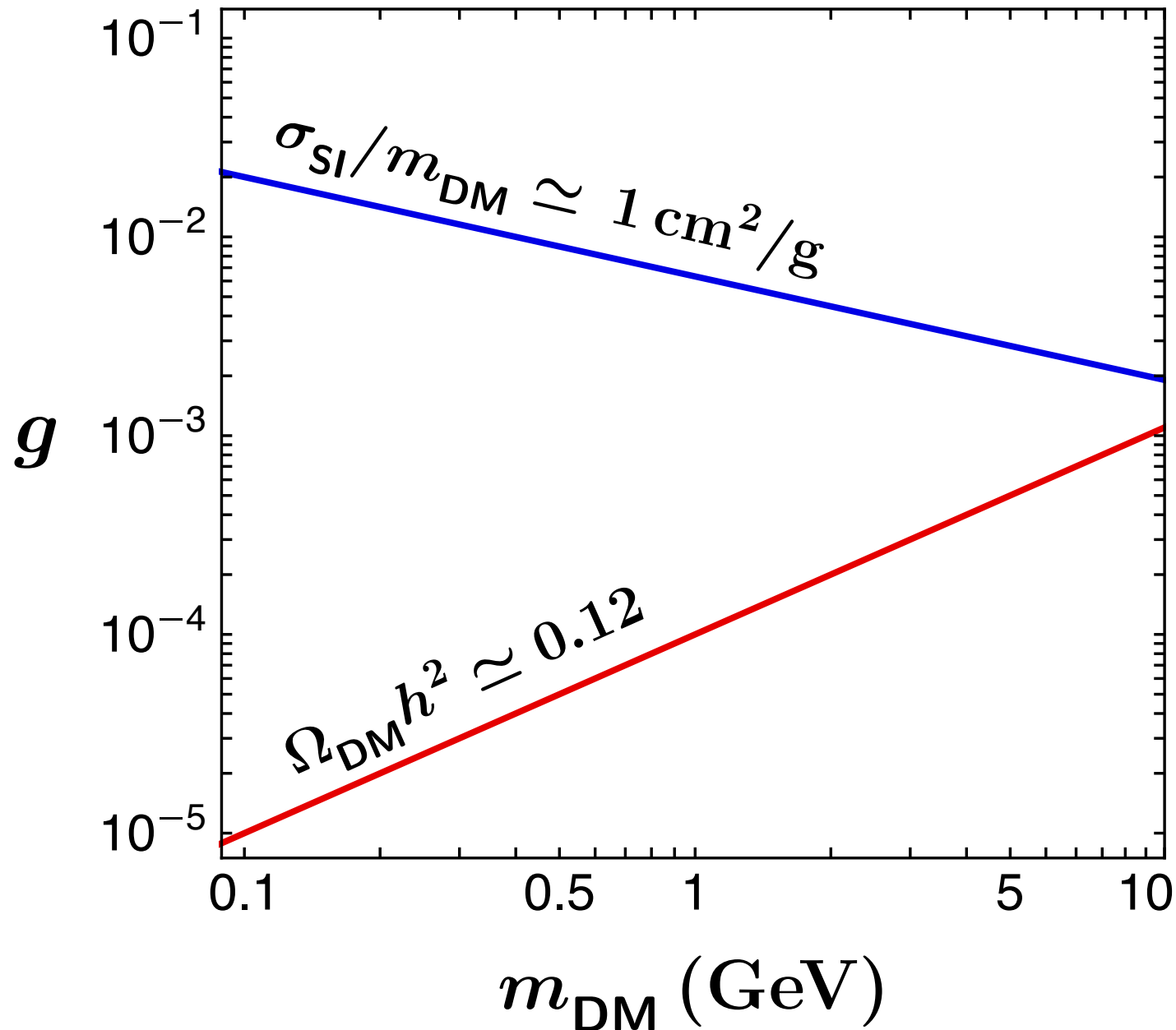
$$\frac{\sigma_{\text{SI}}}{m_{\text{DM}}} = \frac{g^2}{m_{\text{DM}}^3} \Rightarrow g \simeq 2 \times 10^3 \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right)^{3/2} \simeq \mathcal{O}(1) \left(\frac{m_{\text{DM}}}{100 \text{ MeV}} \right)^{3/2}$$

- SI cross-section via a light mediator in the small velocity limit

$$\frac{\sigma_{\text{SI}}}{m_{\text{DM}}} = \frac{g^2}{m_{\text{DM}}^3} \left(\frac{m_{\text{DM}}}{m_{Z'}} \right)^4 \Rightarrow g \simeq 2 \times 10^{-3} \left(\frac{m_{Z'}}{10 \text{ MeV}} \right)^2 \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right)^{-1/2}$$

$\gg 1$

DM mass v.s. coupling



Relic abundance

$$g \simeq 10^{-3} \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right)$$

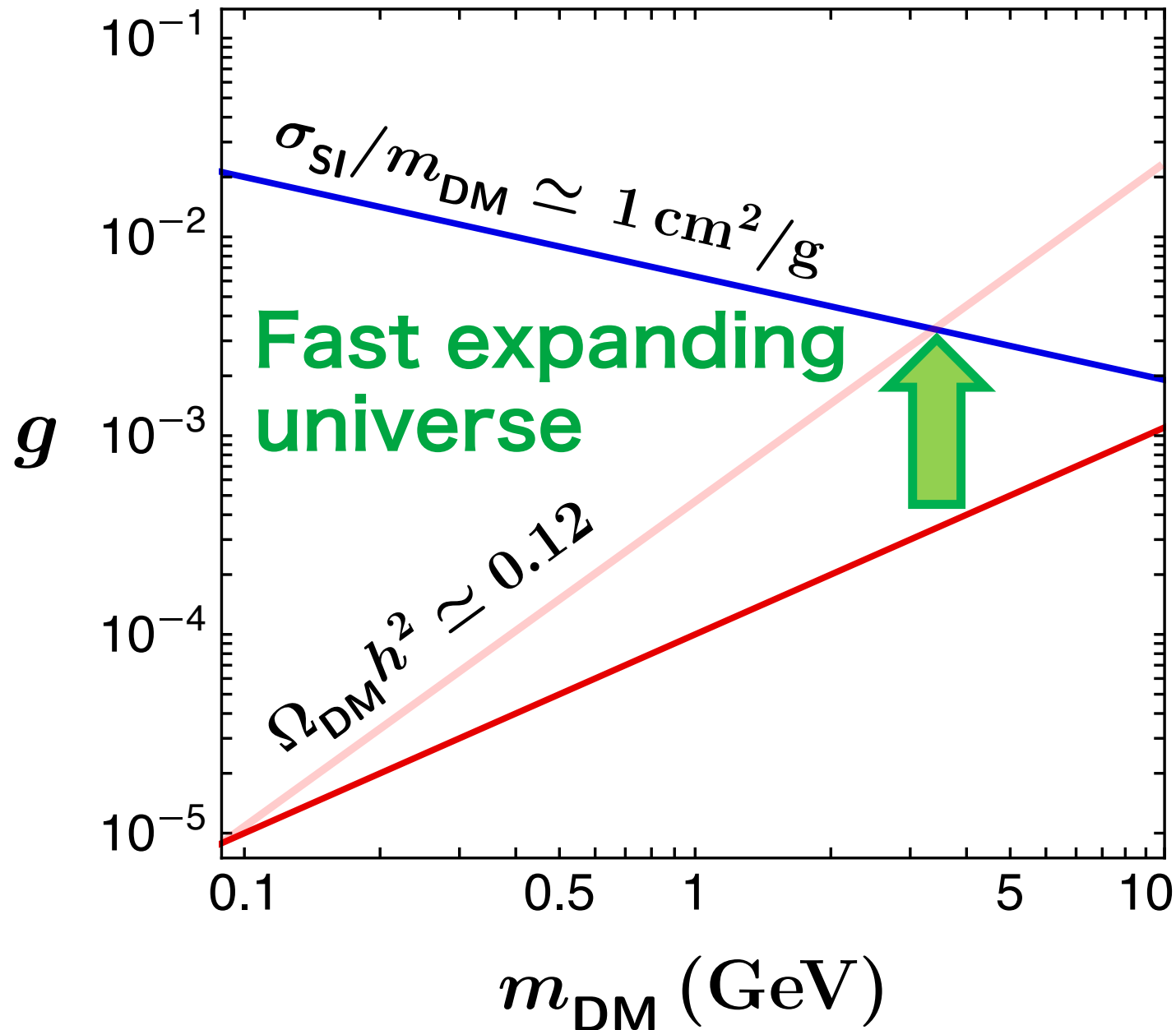
Self-interaction

$$g \simeq 2 \times 10^{-3} \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right)^{-1/2}$$

$$m_{Z'} \sim \mathcal{O}(10) \text{ MeV}$$

DM is under-abundant in low mass regime due to too large annihilation cross section

DM mass v.s. coupling



Relic abundance

$$g \simeq 10^{-3} \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right)$$

Self-interaction

$$g \simeq 2 \times 10^{-3} \left(\frac{m_{\text{DM}}}{10 \text{ GeV}} \right)^{-1/2}$$

$$m_{Z'} \sim \mathcal{O}(10) \text{ MeV}$$

DM is under-abundant in low mass regime due to too large annihilation cross section

Fast expanding universe

D'Eramo, et al (2017)

- Assuming the early universe is dominated by a species ϕ that redshifts faster than radiation :

$$\rho_{\phi}(a) \propto a^{-(4+n)}$$

a : scale factor

$$n > 0$$

- The total energy density :

$$\rho_{\text{tot}}(T) = \rho_{\phi}(T) + \rho_{\gamma}(T) = \rho_{\gamma}(T) \left\{ 1 + \frac{g_{\rho}(T_r)}{g_{\rho}(T)} \left[\frac{g_s(T)}{g_s(T_r)} \right]^{\frac{4+n}{3}} \left(\frac{T}{T_r} \right)^n \right\}$$

$$\mathcal{H}(T) \simeq \sqrt{\frac{\pi^2 g_{\rho}(T)}{90}} \frac{T^2}{m_{\text{Pl}}} \left(\frac{T}{T_r} \right)^{n/2}$$

$$\rho_{\phi}(T_r) = \rho_{\gamma}(T_r)$$

Parameters : (n, T_r)

- $\Delta N_{\nu}(T_{\text{BBN}} \simeq 1 \text{ MeV})$ constraint : $T_r \gtrsim (15.4)^{1/n} \text{ MeV}$

A simple light thermal self-interacting DM model

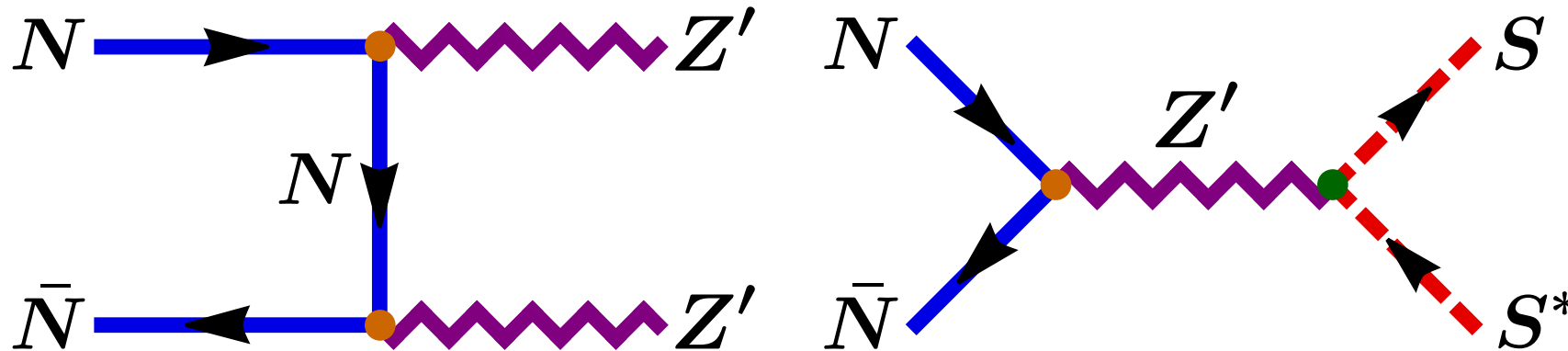
■ Particle content & charge assignment under $G_{\text{SM}} \otimes U(1)_D$

	L	E	H	N	S	Z'
SU(2)	2	1	2	1	1	1
U(1) _Y	-1/2	-1	+1/2	0	0	0
U(1) _D	0	0	0	Q_N	Q_S	0
spin	1/2	1/2	0	1/2	0	1

- N plays the role of fermionic dark matter
- S develops VEV that breaks the **D**ark gauge symmetry
- Z' is a mediator responding the DM self-interaction

Feynman diagrams

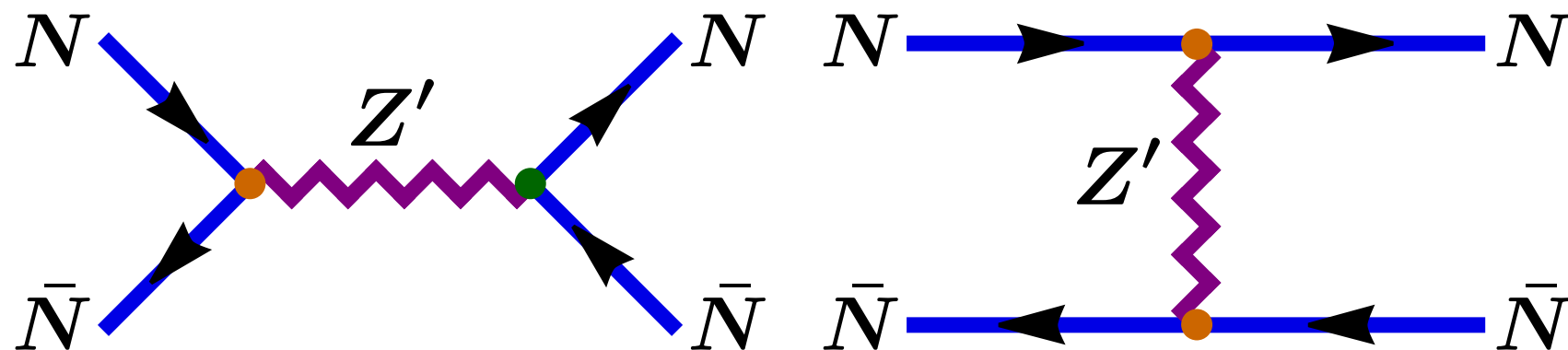
DM annihilation cross-section



$$\langle \sigma v \rangle = \frac{g_D^4}{128\pi m_N^2}$$

(s-wave)

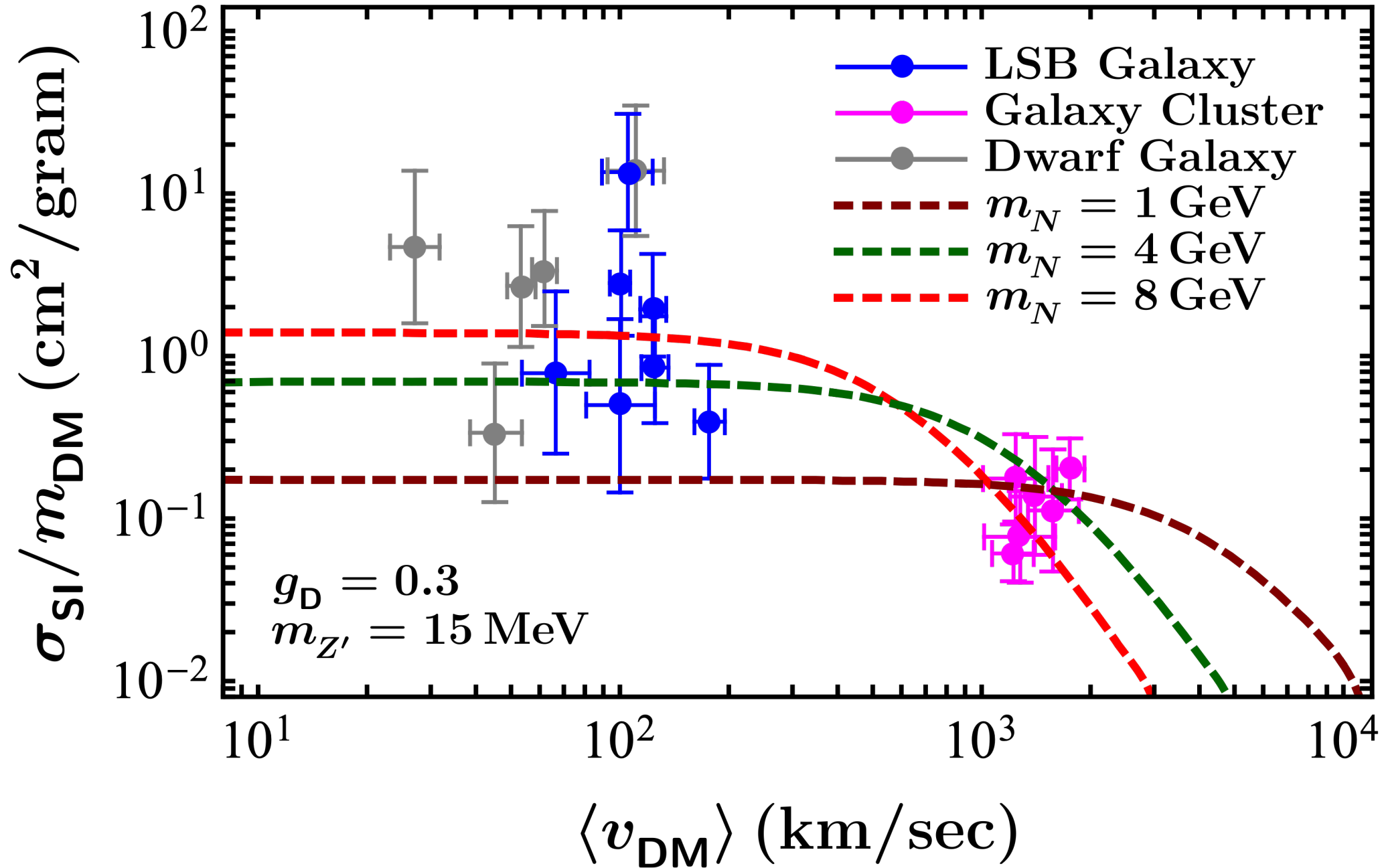
SI cross-section/DM mass



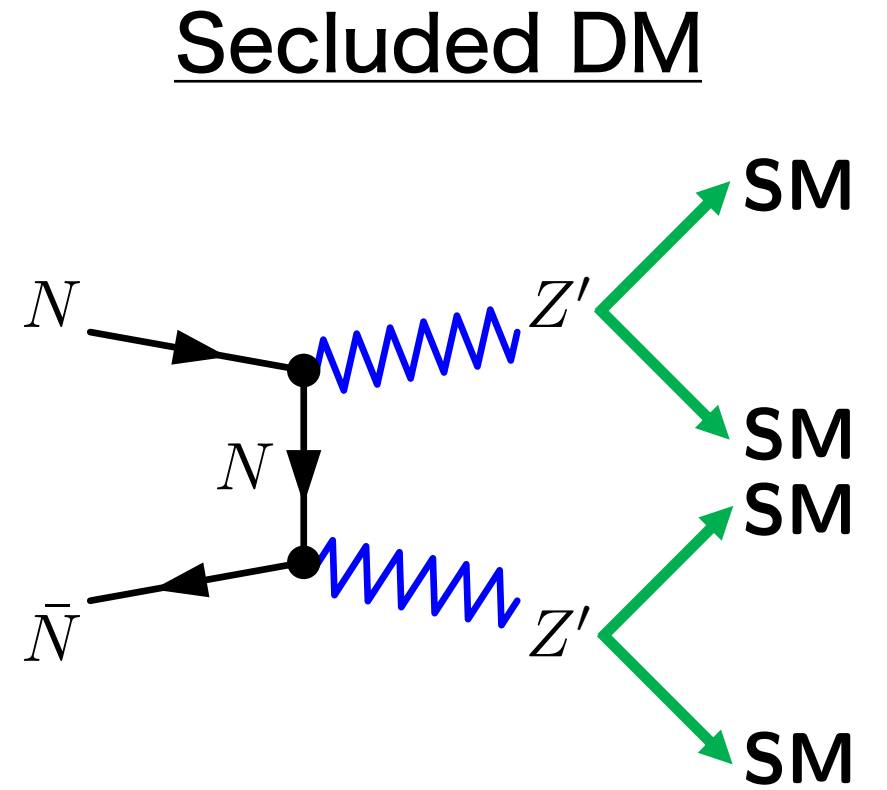
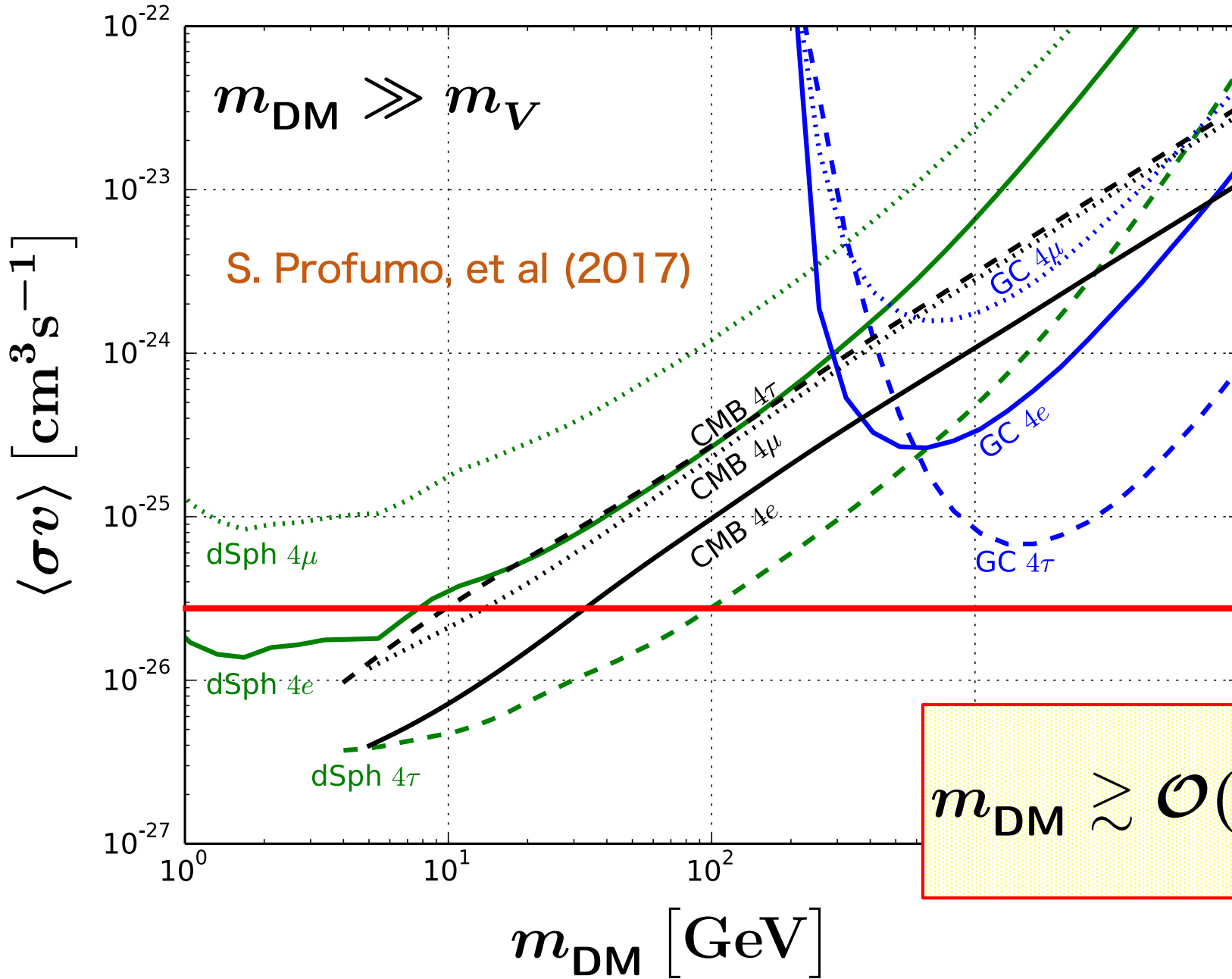
$$\sigma_{\text{SI}} = \frac{\pi}{m_{Z'}^2} f(\beta)$$

$$\beta = \frac{2\alpha_D m_{Z'}}{m_N v_{\text{DM}}^2}$$

Prediction of DM SI cross-section



CMB constraint on light DM mass

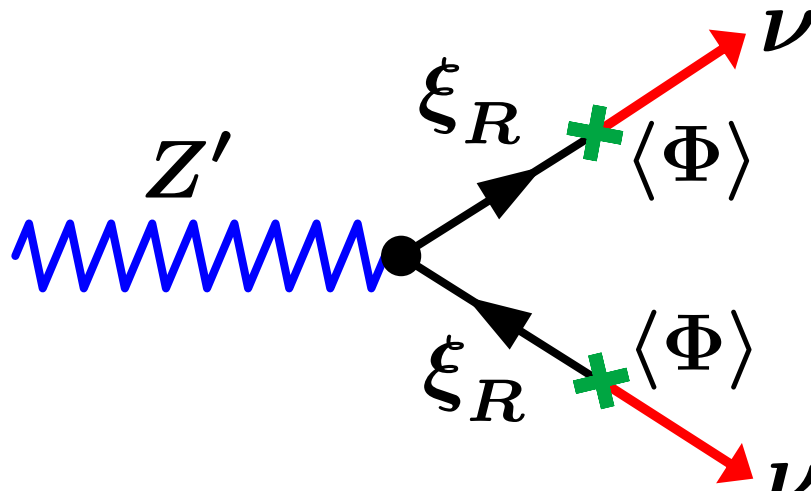


$$m_{\text{DM}} \gtrsim \mathcal{O}(10 \text{ GeV}) \left[\frac{\langle \sigma v \rangle}{10^{-8} \text{ GeV}^{-2}} \right]$$

A viable light thermal self-interacting DM model

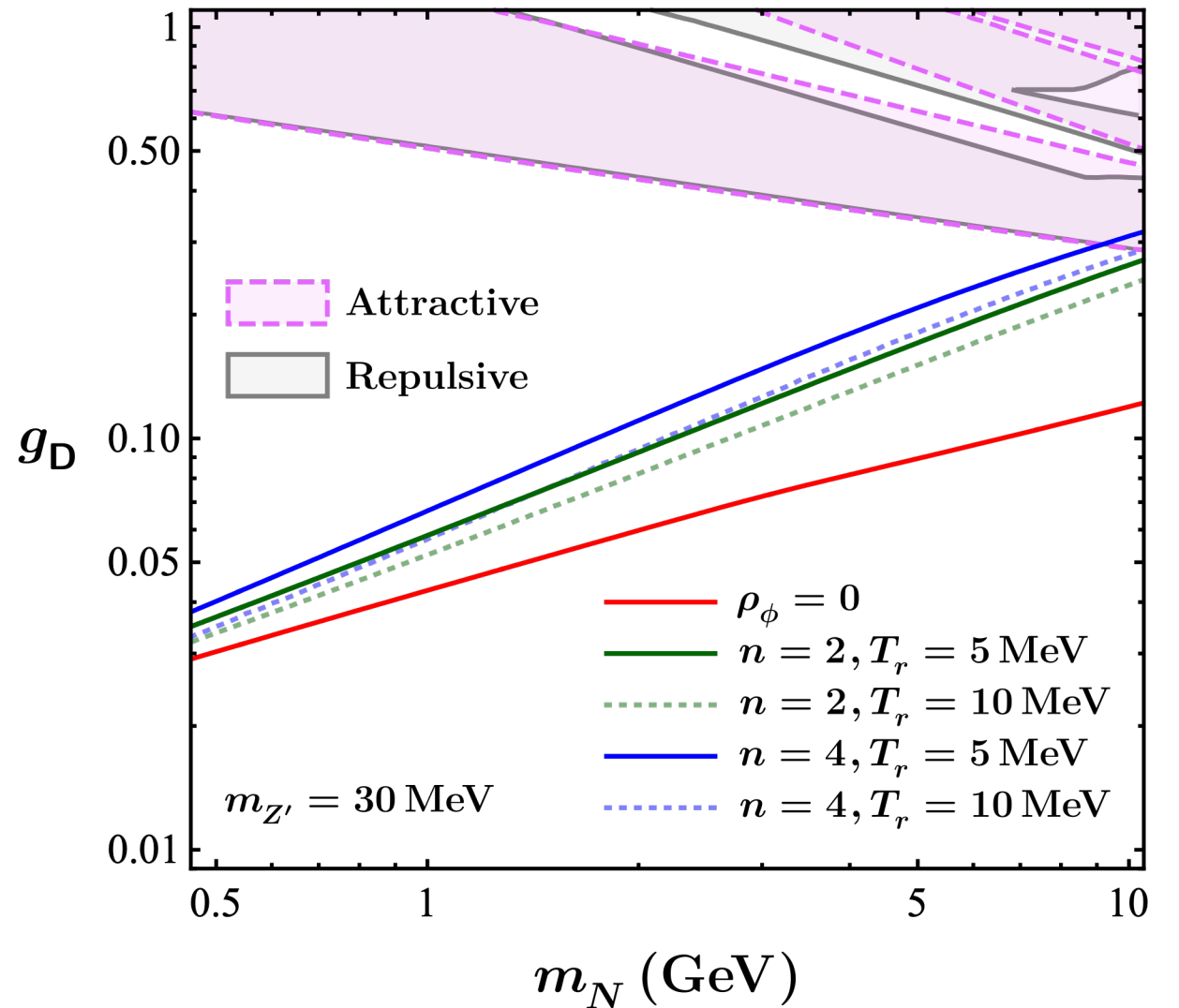
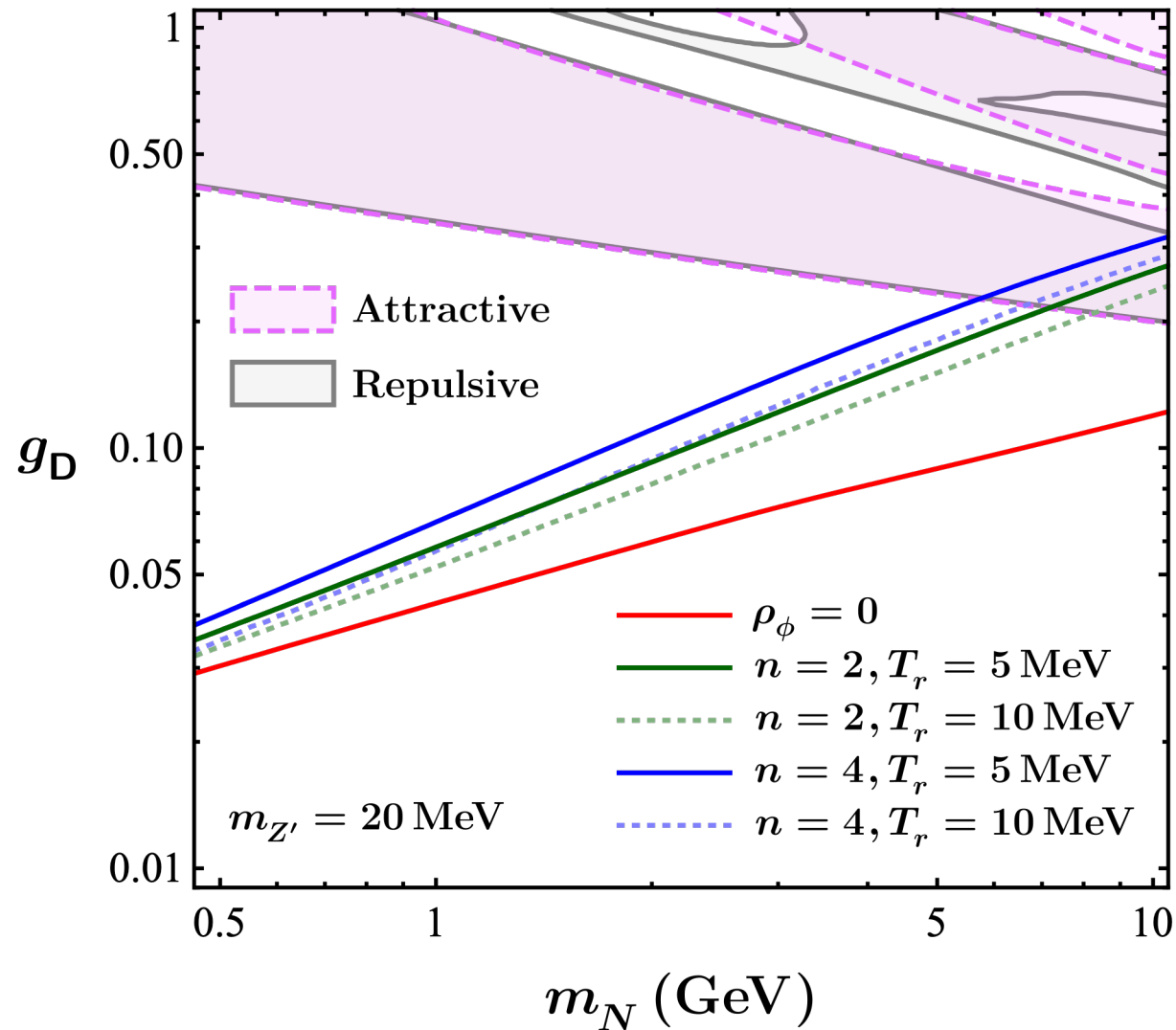
■ Particle content & charge assignment under $G_{SM} \otimes U(1)_D$

	L	E	H	N	ξ_R	χ_L	Φ	S	Z'
SU(2)	2	1	2	1	1	1	2	1	1
U(1) _Y	-1/2	-1	+1/2	0	0	0	+1/2	0	0
U(1) _D	0	0	0	+1/2	+1	+1	+1	+1	0
spin	1/2	1/2	0	1/2	1/2	1/2	0	0	1

● $\mathcal{L} = \mathcal{Y}_\psi \overline{L}_L \tilde{\Phi} \xi_R :$ 

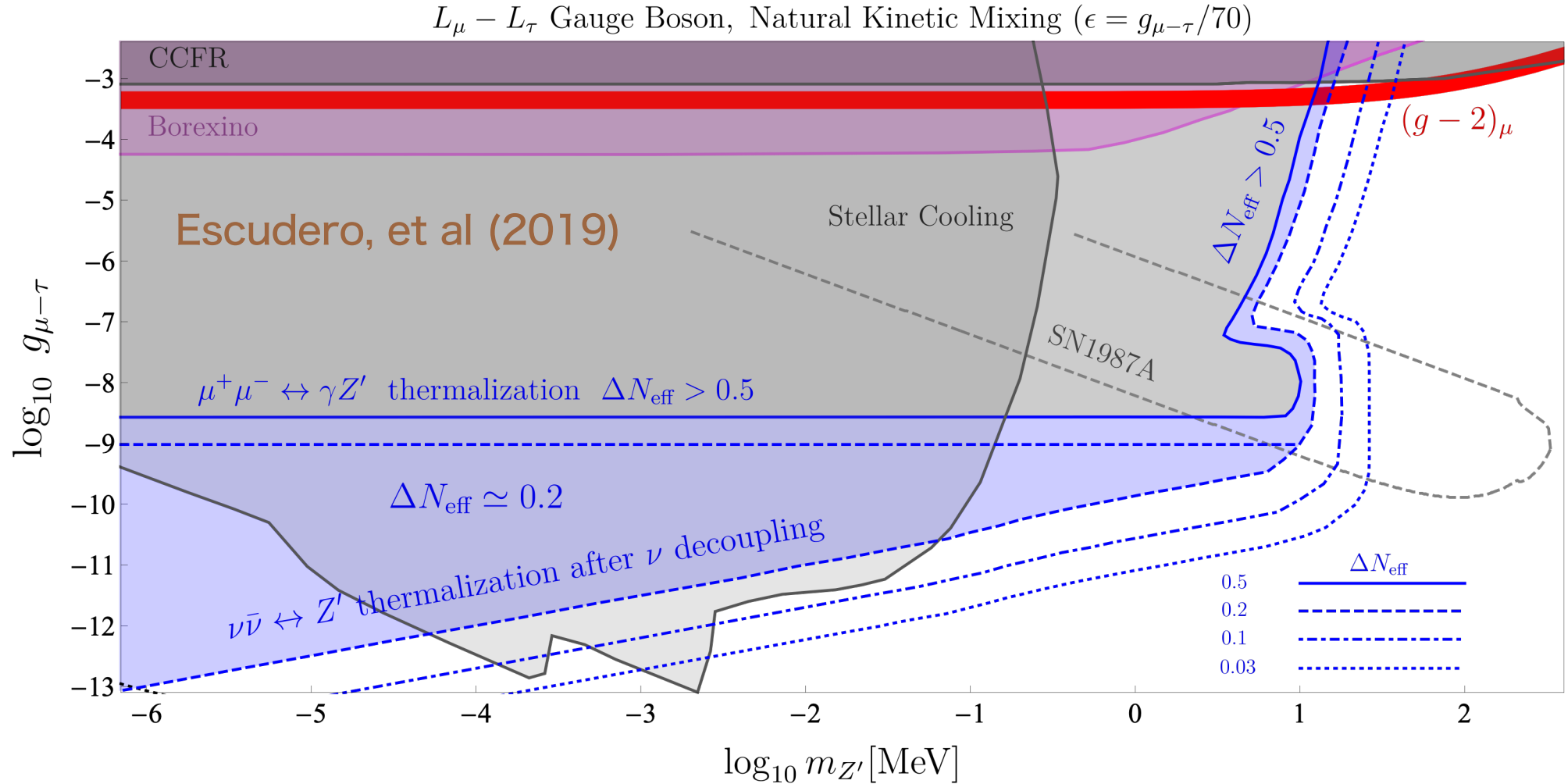
Light mediator
mainly decays
into neutrinos
at CMB epoch

Numerical results



■ Light thermal self-interacting DM can be used to test the non-standard cosmological evolution of the universe.

Backup



- Early Universe Equilibrium:** If $g_{\mu-\tau} \gtrsim 4 \times 10^{-9}$, the Z' population thermalizes with the SM bath at early times and decays into neutrinos when $T \sim m_{Z'}/3$. If these decays occur predominantly after the neutrinos and photons decouple, they contribute to the neutrino energy density and thereby increase the value of N_{eff} . Furthermore, in the presence of non-negligible kinetic mixing with the photon, Z' interactions with charged particles can delay the neutrino-photon decoupling, quantitatively affecting N_{eff} .

