Symmetries and Asymmetries with $B \rightarrow D D$ Decays

Jonathan Davies, Martin Jung, Stefan Schacht October 25, 2023


## Some questions

## What?:

- Testing the SM with $B \rightarrow D D$ databranching ratios and CP asymmetries.
- Update of arxiv:1410.8396



## Some questions

## What?:

- Testing the SM with $B \rightarrow D D$ databranching ratios and CP asymmetries.
- Update of arxiv:1410.8396



## Why?:

Any generic BSM theory would have $\mathcal{O}(1)$ weak-phases $\Longrightarrow$ significant CPV expected.

## Some questions

## What?:

- Testing the SM with $B \rightarrow D D$ databranching ratios and CP asymmetries.
- Update of arxiv:1410.8396



## Why?:

Any generic BSM theory would have $\mathcal{O}(1)$ weak-phases $\Longrightarrow$ significant CPV expected.

Who?:
Myself, Martin Jung and Stefan Schacht

## Some questions

## What?:

- Testing the SM with $B \rightarrow D D$ databranching ratios and CP asymmetries.
- Update of arxiv:1410.8396



## Why?:

Any generic BSM theory would have $\mathcal{O}(1)$ weak-phases $\Longrightarrow$ significant CPV expected.

Who?:
Myself, Martin Jung and Stefan Schacht

How?:
Let's see...

Where is this data coming from?



## Which Observables?

Examples include (not an exhaustive list):
$-\frac{f_{s}}{f_{d}} \frac{\mathrm{BR}\left(D_{s}^{+} \rightarrow K^{+} K^{-} \pi^{+}\right)}{\mathrm{BR}\left(D^{+} \rightarrow K^{-} \pi^{+} \pi^{+}\right)} \frac{\operatorname{BR}\left(\bar{B}_{s} \rightarrow D_{s}^{+} D_{s}^{-}\right)}{\operatorname{BR}\left(\bar{B}^{0} \rightarrow D^{+} D_{s}^{-}\right)}$
$-2 f_{+-} \mathrm{BR}\left(B^{-} \rightarrow D^{0} D_{s}^{-}\right) \mathrm{BR}\left(D^{0} \rightarrow K^{-} \pi^{+}\right) \mathrm{BR}\left(D_{s}^{-} \rightarrow \phi \pi^{-}\right)$

- $\operatorname{BR}\left(B_{s}^{0} \rightarrow D_{s}^{+} D_{s}^{-}\right)$
- $\mathrm{BR}\left(D_{s}^{-} \rightarrow K^{+} K^{-} \pi^{+}\right)$
- $A_{C P}\left(B^{-} \rightarrow D_{s}^{-} D^{0}\right)$ including brand new LHCb input! (arxiv:2306.09945)
- $S_{C P}\left(\bar{B}^{0} \rightarrow D^{-} D^{+}\right)$

We take into account correlations and use updated charm BR's to rescale older results

## Adding Theory Arguments to the Mix

Can make approximation

$$
m_{u} \approx m_{d} \approx m_{s} \ll \Lambda_{Q C D}
$$


$u, d$, and $s$ represented by a triplet in an approximate $S U(3)_{F}$ symmetry

## We could do this...

1. Express everything in terms of $S U(3)$ states:

$$
\begin{aligned}
& u=\left|\mathbf{3},\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{3}\right)\right\rangle, d=\left|\mathbf{3},\left(\frac{1}{2},-\frac{1}{2}, \frac{1}{3}\right)\right\rangle, s=\left|\mathbf{3},\left(0,0,-\frac{2}{3}\right)\right\rangle, c=|\mathbf{0}\rangle, b=|\mathbf{0}\rangle \\
& \left|\mathbf{r}_{\mathbf{1}}, \alpha_{1}\right\rangle\left|\mathbf{r}_{\mathbf{2}}, \alpha_{2}\right\rangle=\sum_{i}\left(\mathbf{R}_{\mathbf{i}}, A_{i} \mid \mathbf{r}_{\mathbf{1}}, \alpha_{1}, \mathbf{r}_{\mathbf{2}}, \alpha_{2}\right)\left|\mathbf{R}_{\mathbf{i}}, A_{i}\right\rangle \\
& \mathcal{H}_{u}^{b \rightarrow d} \sim \sqrt{\frac{3}{8}}\left(\mathbf{3}, \frac{1}{2},-\frac{1}{2}, \frac{1}{3}\right)-\frac{1}{2}\left(\overline{\mathbf{6}}, \frac{1}{2},-\frac{1}{2}, \frac{1}{3}\right)+\sqrt{\frac{1}{24}}\left(\mathbf{1 5}, \frac{1}{2},-\frac{1}{2}, \frac{1}{3}\right)+\sqrt{\frac{1}{3}}\left(\mathbf{1 5}, \frac{3}{2},-\frac{1}{2}, \frac{1}{3}\right)
\end{aligned}
$$

2. Parameterise our decays in terms of $S U(3)$ matrix elements:

$$
\left\langle D^{-} D^{0}\right| \mathcal{H}^{b \rightarrow d}\left|B^{-}\right\rangle \quad \rightarrow \quad \sqrt{\frac{3}{8}}\langle 8|(3)|\overline{3}\rangle_{u}-\sqrt{\frac{1}{20}}\langle 8|(\overline{6})|\overline{3}\rangle_{u}+\ldots
$$

... but diagrams are easier.


Linear combinations of $S U(3)$ matrix elements
... but diagrams are easier.

... but diagrams are easier.


## Theoretical Parameterisation

| Mode |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\lambda_{c D} T$ | $\lambda_{c D} A^{c} \lambda_{u D} \tilde{P}_{1}$ | $\lambda_{u D} \tilde{P}_{3}$ | $\lambda_{u D} A_{1}^{u}$ | $\lambda_{u D} A_{2}^{u}$ |  |  |
|  | Spoilers |  |  |  |  |  |  |
| 1 | $B^{-} \rightarrow D^{-} D^{0}$ | 1 | 0 | -1 | 0 | 1 | 0 |
| 2 | $B^{-} \rightarrow D_{s}^{-} D^{0}$ | 1 | 0 | -1 | 0 | 1 | 0 |
| 3 | $\bar{B}^{0} \rightarrow D_{s}^{-} D^{+}$ | 1 | 0 | -1 | 0 | 0 | 0 |
| 4 | $\bar{B}_{s} \rightarrow D^{-} D_{s}^{+}$ | 1 | 0 | -1 | 0 | 0 | 0 |
| 5 | $\bar{B}^{0} \rightarrow D^{-} D^{+}$ | 1 | 1 | -1 | -1 | 0 | 0 |
| 6 | $\bar{B}_{s} \rightarrow D_{s}^{-} D_{s}^{+}$ | 1 | 1 | -1 | -1 | 0 | 0 |
| 7 | $\bar{B}^{0} \rightarrow D_{s}^{-} D_{s}^{+}$ | 0 | 1 | 0 | -1 | 0 | 0 |
| 8 | $\bar{B}_{s} \rightarrow D^{-} D^{+}$ | 0 | 1 | 0 | -1 | 0 | 0 |
| 9 | $\bar{B}^{0} \rightarrow \bar{D}^{0} D^{0}$ | 0 | -1 | 0 | 1 | 0 | -1 |
| 10 | $\bar{B}_{s} \rightarrow \bar{D}^{0} D^{0}$ | 0 | -1 | 0 | 1 | 0 | -1 |

Generate vector of parameterised observables $\underline{\mathrm{P}}(\underline{\mathbf{x}})$

## Predictions

## Construct $\chi^{2}(\underline{\mathbf{x}})$

Scan function

$$
\chi_{\lambda}^{2}(y)=\min \left(\left.\chi^{2}(\underline{\mathbf{x}})\right|_{x_{\lambda}=y}\right)
$$

Find y range where
$\Delta \chi_{\lambda}^{2}<1(4) \Longrightarrow 68(95) \% \mathrm{Cl}$


D branching ratios, fragmentation fractions and efficiencies also floated.

Our main goal is to perform these scans for our physical observables

## Power Counting

A way of managing the relative sizes of the diagrams.

We take $\epsilon \sim 0.3$ and consider the following contributions:

- $\operatorname{SU}(3)$ structure- $\mathcal{O}(\epsilon)$ : For $\mathrm{SU}(3)$-breaking contributions
- CKM suppression- $\mathcal{O}(\epsilon)$ : Where CKM factors cannot be separated from the hadronic matrix elements.
- Colour suppression: Relative counting in $1 /(\#$ colours $) \sim \mathcal{O}(\epsilon)$ for the topologies, based on [Buras, Gèrard, Rückl (1986)]
- Penguin suppression:
- Tree matrix elements of penguin operators- $\mathcal{O}\left(\epsilon^{2}\right)$
- Penguin matrix elements of tree operators- $\mathcal{O}\left(\epsilon^{1 / 2}\right)$
- Annihilation: $\mathcal{O}\left(\epsilon^{1 / 2}\right)$ for annihilation diagram $+\mathcal{O}(\epsilon)$ for $c \bar{c}$ creation.


## Validation

Expected relative scaling from power counting for decay rates and $C P$ asymmetries can be checked with data.

| Relative rate | $b \rightarrow s$ | $b \rightarrow d$ |
| :--- | :---: | :---: |
| Tree-dominated | 1 | $\lambda^{2}$ |
| Annihilation-dominated | $\epsilon^{3}$ | $\bar{\lambda}^{2} \epsilon^{3}$ |


| CP asymmetry | $b \rightarrow s$ | $b \rightarrow d$ |
| :--- | :--- | :---: |
| Tree-dominated | $\lambda^{2} \epsilon^{2.5}$ | $\epsilon^{2.5}$ |
| Annihilation-dominated | $\bar{\lambda}^{2} \epsilon^{2}$ | $\epsilon^{2}$ |

- Measured rates correspond to our scaling usually within $30 \%$, as expected.
- Slightly larger deviations of $\sim 40 \%$ for $\bar{B}^{0} \rightarrow D^{+} D^{-}$and $\bar{B}_{s} \rightarrow D_{s}^{+} D_{s}^{-}$- first sign for negative interference of the sizable $A_{c}$ [arxiv:1410.8396].
- New LHCb measurement for $A_{C P}\left(B^{-} \rightarrow D^{0} D^{-}\right) \sim 2 \%$ is on the lower side of $\epsilon^{2.5} \sim 5 \% \Longrightarrow$ our penguin amplitude scaling is conservative.


## Theoretical Parameterisation

|  | Mode | $\lambda_{c D} T$ | $\lambda_{c D} A^{c}$ | $\lambda_{u D} \tilde{P}_{1}$ | $\lambda_{u D} \tilde{P}_{3}$ | $\lambda_{u D} A_{1}^{u}$ | $\lambda_{u D} A_{2}^{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Counting | 1 | $\varepsilon^{1.5}$ | $\varepsilon^{2.5}$ | $\varepsilon^{3.5}$ | $\varepsilon^{2.5}$ | $\varepsilon^{3.5}$ |
| 1 | $B^{-} \rightarrow D^{-} D^{0}$ | 1 | 0 | -1 | 0 | 1 | 0 |
| 2 | $B^{-} \rightarrow D_{s}^{-} D^{0}$ | 1 | 0 | -1 | 0 | 1 | 0 |
| 3 | $\bar{B}^{0} \rightarrow D_{s}^{-} D^{+}$ | 1 | 0 | -1 | 0 | 0 | 0 |
| 4 | $\bar{B}_{s} \rightarrow D^{-} D_{s}^{+}$ | 1 | 0 | -1 | 0 | 0 | 0 |
| 5 | $\bar{B}^{0} \rightarrow D^{-} D^{+}$ | 1 | 1 | -1 | -1 | 0 | 0 |
| 6 | $\bar{B}_{s} \rightarrow D_{s}^{-} D_{s}^{+}$ | 1 | 1 | -1 | -1 | 0 | 0 |
| 7 | $\bar{B}^{0} \rightarrow D_{s}^{-} D_{s}^{+}$ | 0 | 1 | 0 | -1 | 0 | 0 |
| 8 | $\bar{B}_{s} \rightarrow D^{-} D^{+}$ | 0 | 1 | 0 | -1 | 0 | 0 |
| 9 | $\bar{B}^{0} \rightarrow \bar{D}^{0} D^{0}$ | 0 | -1 | 0 | 1 | 0 | -1 |
| 10 | $\bar{B}_{s} \rightarrow \bar{D}^{0} D^{0}$ | 0 | -1 | 0 | 1 | 0 | -1 |

Power counting enters as hard parameter constraints during $\chi^{2}$ minimisation.

## SU(3) Breaking

But $m_{u} \neq m_{d} \neq m_{s}$ !

Must fix our assumption that $d \equiv s$ :

$$
\mathcal{H}_{S U(3)} \sim \frac{m_{s}-m_{d}}{\Lambda_{Q C D}}(s \bar{s})
$$

Fitting this into diagrammatic language [arxiv:9504326]:


## Breaking table

|  | Mode | $\lambda_{c D} \delta T_{1}$ | $\lambda_{c D} \delta T_{2}$ | $\lambda_{c D} \delta A_{1}^{c}$ | $\lambda_{c D} \delta A_{2}^{c}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | Counting | $\varepsilon^{1(2)}$ | $\varepsilon^{1(2)}$ | $\varepsilon^{2.5}$ | $\varepsilon^{2.5}$ |
| 1 | $B^{-} \rightarrow D^{-} D^{0}$ | 0 | $-\frac{1}{2}$ | 0 | 0 |
| 2 | $B^{-} \rightarrow D_{s}^{-} D^{0}$ | 1 | 0 | 0 | 0 |
| 3 | $\bar{B}^{0} \rightarrow D_{s}^{-} D^{+}$ | 1 | 0 | 0 | 0 |
| 4 | $\bar{B}_{s} \rightarrow D^{-} D_{s}^{+}$ | -1 | $\frac{1}{2}$ | 0 | 0 |
| 5 | $\bar{B}^{0} \rightarrow D^{-} D^{+}$ | 0 | $-\frac{1}{2}$ | $\frac{1}{2}$ | $-\frac{1}{2}$ |
| 6 | $\bar{B}_{s} \rightarrow D_{s}^{-} D_{s}^{+}$ | 0 | 1 | -1 | 1 |
| 7 | $\bar{B}^{0} \rightarrow D_{s}^{-} D_{s}^{+}$ | 0 | 0 | $\frac{1}{2}$ | $\frac{1}{2}$ |
| 8 | $\bar{B}_{s} \rightarrow D^{-} D^{+}$ | 0 | 0 | -1 | 0 |
| 9 | $\bar{B}^{0} \rightarrow \bar{D}^{0} D^{0}$ | 0 | 0 | $-\frac{1}{2}$ | $\frac{1}{2}$ |
| 10 | $\bar{B}_{s} \rightarrow \bar{D}^{0} D^{0}$ | 0 | 0 | 1 | 0 |

## Observables

From amplitudes, it is simple to go to experimental observables:

- Branching Ratios

$$
\mathcal{B}\left(\mathcal{D} \equiv B \rightarrow P_{1} P_{2}\right)=|\mathcal{A}(\mathcal{D})|^{2} \times \text { phase }- \text { space }
$$

- CP Asymmetries

$$
\begin{gathered}
A_{C P}(\mathcal{D})=-\frac{1-|\lambda(\mathcal{D})|^{2}}{1+|\lambda(\mathcal{D})|^{2}}=\frac{|\mathcal{A}(\mathcal{D})|^{2}-|\overline{\mathcal{A}}(\mathcal{D})|^{2}}{|\mathcal{A}(\mathcal{D})|^{2}+|\overline{\mathcal{A}}(\mathcal{D})|^{2}}, \quad S_{C P}(\mathcal{D})=\frac{2 \operatorname{Im}(\lambda(\mathcal{D}))}{1+|\lambda(\mathcal{D})|^{2}} \\
{\left[\lambda(\mathcal{D})=\eta_{C P}^{f} e^{-i \phi_{D}} \frac{\mathcal{A}(\mathcal{D})}{\overline{\mathcal{A}}(\mathcal{D})}\right]}
\end{gathered}
$$

## SM prediction for CP Asymmetries




Predictions extracted both with and without experimental CP information.

## SM prediction for CP Asymmetries



Predictions extracted both with and without experimental CP information.

## We can do 2D scans too!



Blue contours represent $68 \%$ and $95 \% \mathrm{Cl}$ from global fit. Yellow shows experiment-only constraints. Non-trivial correlations, resulting from underlying theory parameterisation, can be seen.

## To conclude...

- Our predictions for CP asymmetries can be used to probe for new physics with future measurements at LHCb
- By working from an assumption of approximate $\operatorname{SU}(3)$ symmetry we can obtain predictions for observables by fitting to experimental data
- Symmetry assumptions found to be valid
- No significant tension with Standard Model found but precision is improved
- We provide predictions for many as-yet unmeasured modes
- Stay tuned for publication soon!


## Time for grilling!

Thanks for your attention.


Hiring?- Find my CV here or email me at jonathan.edward.davies@cern.ch

