

Differential measurements with FCCC at LHCb

Implications of LHCb measurements and future prospects 2023

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on behalf of the LHCb Collaboration



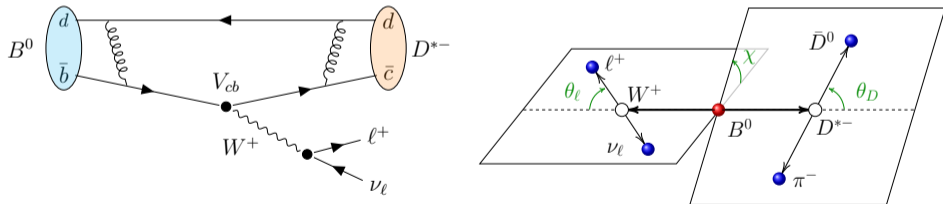
October 25-27 2023, CERN

Angular analyses of semileptonic b -hadron decays

- Differential decay rate can be written as:

$$\frac{d^4\Gamma(B^0 \rightarrow D^{*-} \ell^+ \nu_\ell)}{dq^2 d\cos^2\theta_\ell d\cos\theta_{D^*} d\chi} \propto |V_{cb}|^2 \sum_i \mathcal{H}_i(q^2) f_i(\theta_\ell, \theta_{D^*}, \chi)$$

- q^2 : squared invariant mass of the $\tau\nu_\tau$ system $\equiv (p_B - p_{D^*})^2$
- \mathcal{H}_i : electroweak couplings & QCD form-factors
- f_i : helicity angles distributions, sensitive to New Physics hadronic effects



Angular analyses are good candidates for New Physics search providing complementary information to Lepton Flavour Universality tests (see [Florian's talk](#))

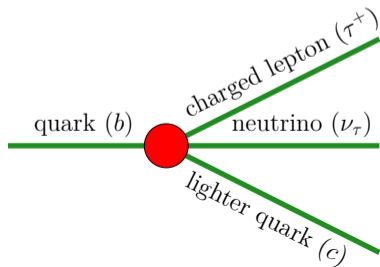
New Physics

- New Physics effects can be described by an effective Hamiltonian consisting in operators with unknown coupling constants

$$\mathcal{H}_{eff} = \frac{G_F}{\sqrt{2}} V_{cb} \sum_i C_i \mathcal{O}_i$$

- \mathcal{O}_i : effective operators (scalar, vector, tensor)
- C_j : Wilson Coefficients (WC) describing the New Physics effects

$$C_j = C_j^{SM} + C_j^{NP}$$



- **Hammer** tool can be used to reweight simulated samples obtaining dynamic templates ([Eur. Phys. J. C80 883 \(2020\)](#))
- Different strategies can be considered:
 - Measure directly WC parameters
 - Measure angular coefficients (amplitudes & q^2 dependence) which relate to the WC

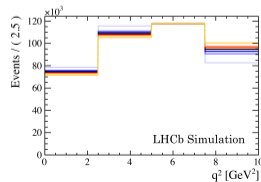
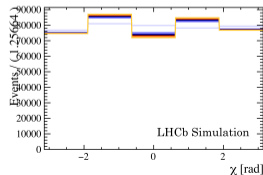
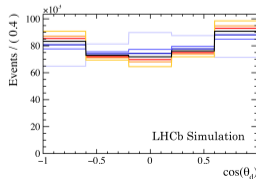
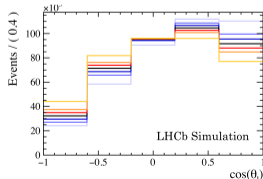
Hammer tool (Eur. Phys. J. C80 883 (2020))

- Vary the amplitude of a simulated sample from the generation model to another one of interest
- A per-event weight is determined according to the FF and WC values of the two models
- New model templates are obtained by contracting the weights tensor

$$\omega_j = \frac{\Gamma_{gen} d^n \Gamma_{new} / dX}{\Gamma_{new} d^n \Gamma_{gen} / dX}$$

- Γ_{gen} : decay rate for the model used to generate the simulated sample
- Γ_{new} : decay rate according to the NP model of interest

- $B^0 \rightarrow D^{*-} \mu^+ \nu_\mu$, varying V_{qRlL} in (-0.5, -0.2, -0.1, 0., 0.1, 0.2, 0.5) [Hammer manual]



$F_L^{D^*}$ measurement
in $B^0 \rightarrow D^{*-} \tau^+ \nu_\tau$ decays

Longitudinal D^* polarization

- Measurement of the longitudinal D^* polarization can provide complementary information to $R(D^*)$, showing NP contribution even if $R(D^*)$ is found compatible with SM expectation
- The differential decay rate can be expressed as 2° polynomial in $\cos \theta_D$:

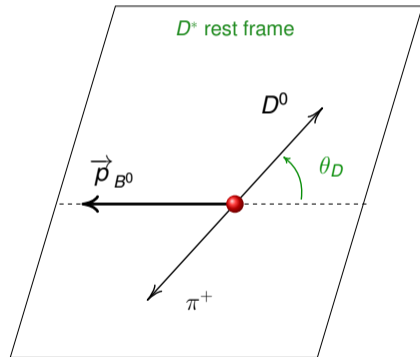
$$\frac{d^2\Gamma}{dq^2 d \cos \theta_D} = a_{\theta_D}(q^2) + c_{\theta_D}(q^2) \cos^2 \theta_D$$

- D^* longitudinal polarization fraction as function of $a_{\theta_D}(q^2)$ and $c_{\theta_D}(q^2)$:

$$F_L^{D^*}(q^2) = \frac{a_{\theta_D}(q^2) + c_{\theta_D}(q^2)}{3a_{\theta_D}(q^2) + c_{\theta_D}(q^2)}$$

- State of art is determined by Belle results:

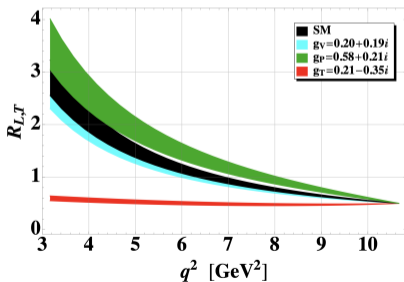
$$F_L^{D^*} = 0.60 \pm 0.08 \pm 0.04$$



arXiv:1903.03102

Expected value of F_L^D

- F_L^{D*} value within the SM scenario has been predicted with different methods
- The most recent theoretical predictions are:
 - 0.441 ± 0.006 [[Phys. Rev. D98, 095018 \(2018\)](#)] Z.-R. Huang, Y. Li, C.-D. Lu, M. A. Paracha, C. Wang
 - 0.457 ± 0.010 [[EPJ C79 268 \(2019\)](#)] S. Bhattacharya, S. Nandi, S. K. Patra
- Predictions for NP scenarios can be found in [arXiv:1907.02257](#)
D. Becirevic, M. Fedele, I. Nisandzic, A. Tayduganovd



[arXiv:1907.02257](#)

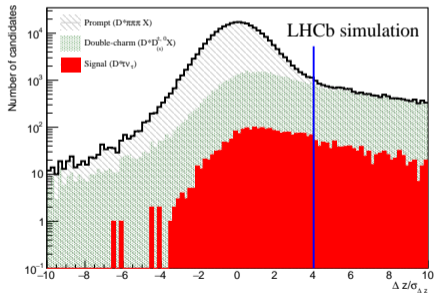
- Expected dependence of $R_{L,T}$ as function of q^2 for three NP models

$$R_{L,T}(q^2) = \frac{d\Gamma_L/dq^2}{d\Gamma_T/dq^2}$$

$$F_L^{D*}(q^2) = \frac{R_{L,T}(q^2)}{1 + R_{L,T}(q^2)}$$

Signal selection

- Analysis performed using hadronic τ decays
- Same strategy used in the $R(D^*)$ hadronic
- Simultaneous fit on 2011-12 (Run 1) and 2015-16 (Run 2) data



[Phys. Rev. D 97, 072013 (2018)]

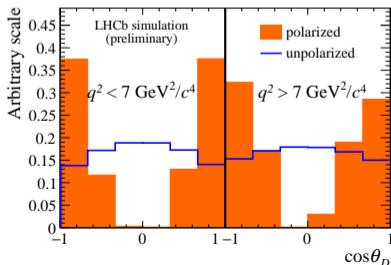
- Advantages of hadronic the decay mode:
 - **only 1 neutrino** in the τ decay:
⇒ event kinematic is properly reconstructed
 - **good purity** ⇒ strong background rejection

Initially dominant backgrounds

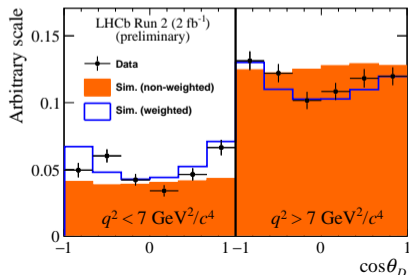
- Prompt decay $B \rightarrow D^{*-} 3\pi^{\pm} X$
 - 3π system from the B meson
 - $\sim 100\times$ signal decays
 - requiring a 3π vertex detached by the B vertex along the beam axis ($\Delta_z/\sigma_{\Delta_z} > 4$)
 - additional BDT in Run 2 to reach Run 1
rejection level: $> 99.9\%$
- Double charm $B \rightarrow D^{*-} D_{(s)}^{+,0} X$ decays
 - signal like topology
 - detached vertex due to non-negligible lifetime
 - rejected through isolation algorithm and dedicated MVA classifiers

Signal & background description [LHCb-PAPER-2023-020] (in preparation)

- $F_L^{D^*}$ determined in two q^2 regions: $\leq 7 \text{ GeV}^2/c^4$
- $F_L^{D^*}$ is extracted from $a_{\theta_D}(q^2)$ and $c_{\theta_D}(q^2)$, determined splitting the signal sample in:
 - **unpolarized** $\implies N_{sig}^{unpol} \propto a_{\theta_D}(q^2)$
 - **polarized** $\implies N_{sig}^{pol} \propto c_{\theta_D}(q^2)$
- $\cos \theta_D$ signal distribution corrected for reconstruction effect



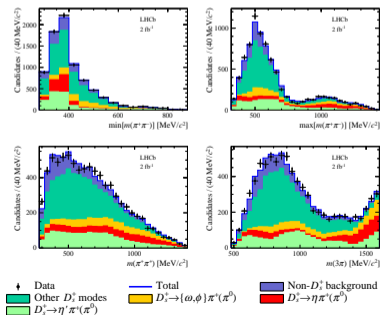
- $D^{*-} DX$ background templates determined from simulation
- Assuming no $F_L^{D^*}$ dependence on the D meson decay mode
- $\cos \theta_D$ distribution corrected through fully reconstructed control samples:
 - $D_s^- \rightarrow 3\pi^\pm$
 - $D^+ \rightarrow K^- 2\pi^+$
 - $D^0 \rightarrow 3\pi^\pm K^-$



Modelling of D_s in $B \rightarrow D^{*-} D_s(X)$ decays

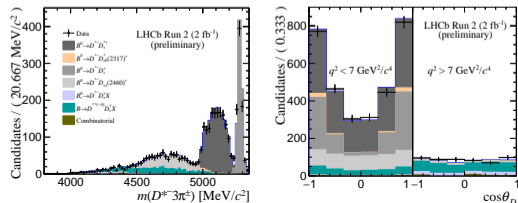
Decays of D_s [Phys. Rev. D108, 012018 (2023)]

- $D_s \rightarrow 3\pi^\pm X$ branching fractions are not well known or correctly simulated
- Data sample selected using D_s BDT output
- Simultaneous fit to: $\min[m(\pi^+\pi^-)]$, $\max[m(\pi^+\pi^-)]$, $m(\pi^+\pi^+)$, $m(3\pi)$
- D_s fractions used to correct simulation



Production of D_s [LHCb-PAPER-2023-020] (in preparation)

- D_s meson arises from $B \rightarrow D^{*-} D_s^{+(*,**)} X$
- Poor knowledge on their relative fractions
- Enriched data sample of $D^{*-} D_s X$ decays reconstructed from $D_s \rightarrow 3\pi^\pm$ mode
- Fraction with respect the $D^* D_s^{*+}$ channel determined through a fit to the $m(D^{*-} 3\pi)$
- Values are used to constrain the various component in the final fit



Extrapolation of F_L^{D*} on simulated sample (III)

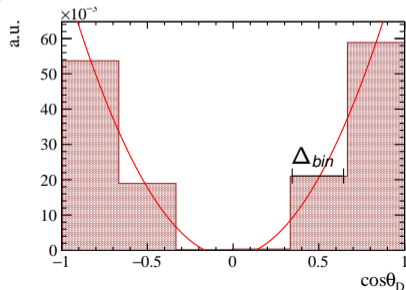
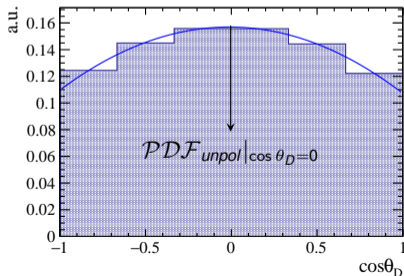
- F_L^{D*} can be determined using the equations in slide 6

$$F_L^{D*}(q^2) = \frac{a_{\theta_D}(q^2) + c_{\theta_D}(q^2)}{3a_{\theta_D}(q^2) + c_{\theta_D}(q^2)}$$

- a_{θ_D} and c_{θ_D} are directly related to the polarized and unpolarized yield

$$a_{\theta_D}(q^2) = N^{unpol} \cdot \mathcal{PDF}_{unpol}|_{\cos\theta_D=0}, \quad c_{\theta_D}(q^2) = \frac{3}{2} N^{pol} \Delta_{bin}$$

signal $\cos\theta_D$ distribution



Signal fit results [LHCb-PAPER-2023-020] (in preparation)

- Signal yields from a 4D-binned template fit:
 - τ^+ lifetime (first row)
 - q^2 & $\cos\theta_D$ (second row)
 - anti- D_s BDT output (third row)
- Fit performed simultaneously on Run 1 and Run 2
- Results are integrated over Run 1 and Run 2

$F_L^{D_s^*}$ value extracted for the 3 q^2 region

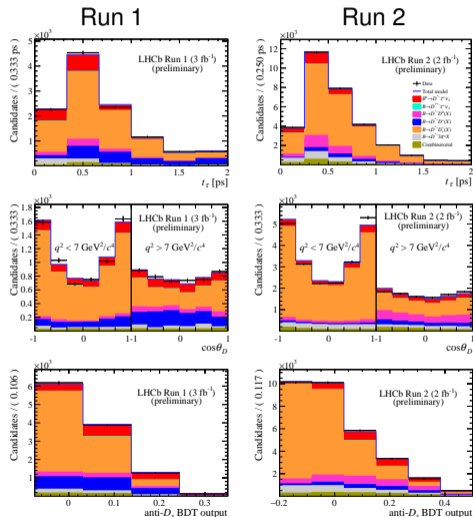
$$q^2 < 7 \text{ GeV}^2/c^4 : \quad 0.51 \pm 0.07(\text{stat}) \pm 0.03(\text{syst})$$

$$q^2 > 7 \text{ GeV}^2/c^4 : \quad 0.35 \pm 0.08(\text{stat}) \pm 0.02(\text{syst})$$

$$q^2 \text{ whole range} : \quad 0.43 \pm 0.06(\text{stat}) \pm 0.03(\text{syst})$$

- All values are found to be compatible with the SM within 1σ
 - expected value in the integrated region ~ 0.44

[Phys. Rev. **D98**, 095018 (2018), EPJ **C79** 268 (2019)]



Systematic uncertainties [LHCb-PAPER-2023-020] (in preparation)

Source	low q^2	high q^2	integrated
Fit validation	0.003	0.002	0.003
FF model	0.007	0.003	0.005
FF parameters	0.013	0.006	0.011
TemplateSize	0.027	0.017	0.019
$\tau^+ \rightarrow 3\pi^\pm \pi^0$ fraction	0.001	0.001	0.001
D^{*+} feed-down	0.001	0.004	0.003
Signal selection	0.005	0.004	0.005
Bin migration	0.008	0.006	0.007
$F_L^{D^{*+}}$ in simulation	0.007	0.003	0.007
D_s decay model	0.008	0.009	0.009
$\cos \theta_D$ $D^{*+} D_s$	0.002	0.001	0.002
$\cos \theta_D$ $D^{*+} D_s^{*+}$	0.007	0.002	0.004
$\cos \theta_D$ $D^{*+} D_s X$	0.007	0.006	0.007
$\cos \theta_D$ $D^{*+} D^+ X$	0.002	0.002	0.003
$\cos \theta_D$ $D^{*+} D^0 X$	0.002	0.002	0.003
$F_L^{D^{*+}}$ integrated	-	-	0.002
Total	0.036	0.023	0.029

Dominant source of systematic are:

- Limited size of the simulation samples
- Form factor parameterization
- Modelling of the D_s
- $\cos \theta_D$ shape in $D^{*+} D_s X$ backgrounds
- Bin migration
- Signal acceptance
- Form factor model

Future updates on the $F_L^{D^*}$ measurement

- $F_L^{D^*}$ measurement performed using Run1 and first part of Run2 data

$$F_L(D^*) = 0.43 \pm 0.06 \text{ (stat)} \pm 0.03 \text{ (syst)}$$

- Plan is to update the $F_L^{D^*}$ value in parallel with the $R(D^*)$ measurement in hadronic τ channel.
- Expected statistical uncertainties for future updates: (optimistic scenario)

	Run1 + (2015+2016)	Run1 + Run2	Run1 + Run2 + Run3
Luminosity	5 fb ⁻¹	9 fb ⁻¹	30 fb ⁻¹
Statistical uncertainty	0.061	~ 0.046	~ 0.025

- Some of the dominant systematic uncertainties should also decrease:
 - $\cos \theta_D$ shape in $D^* D$ backgrounds
 - D_s decay model

$B^0 \rightarrow D^* \mu \nu_\mu$ angular analysis

- Extract directly WC & FF parameters from fit to data
- Shape analysis only
⇒ no attempt to measure $|V_{cb}|$ because of lower sensitivity to yield changes

$$\frac{d^4\Gamma}{dq^2 d \cos \theta_D d \cos \theta_l d \chi} \propto \left[l_{1c} \cos^2 \theta_D + l_{1s} \sin^2 \theta_D \right. \\ + (l_{2c} \cos^2 \theta_D + l_{2s} \sin^2 \theta_D) \cos 2\theta_l \\ + (l_{6c} \cos^2 \theta_D + l_{6s} \sin^2 \theta_D) \cos \theta_l \\ + (l_3 \cos 2\chi + l_9 \sin 2\chi) \sin^2 \theta_l \sin^2 \theta_D \\ + (l_4 \cos \chi + l_8 \sin \chi) \sin 2\theta_l \sin 2\theta_D \\ \left. + (l_5 \cos \chi + l_7 \sin \chi) \sin \theta_l \sin 2\theta_D \right],$$

- Full description using the three helicity angles
- Measure the 12 angular coefficients

Main background contributions

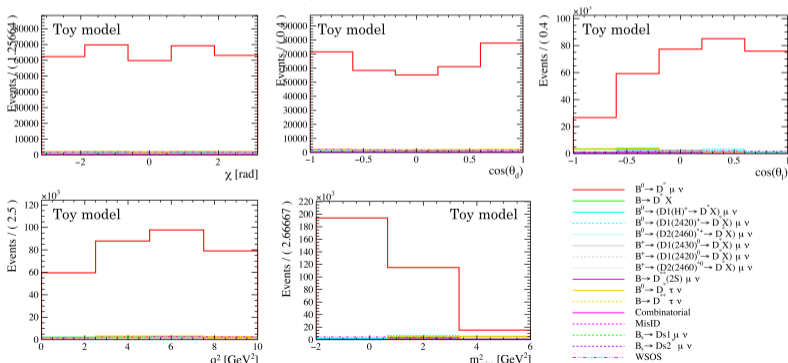
- $B^0 \rightarrow D^* \tau \nu_\tau$: no NP accounted for in the fit
- $B \rightarrow D^{**} \mu \nu_\mu$: modelled with the BLPR FF parameterization [PRD 97 075011 \(2018\)](#)
- Semileptonic decays to heavier charmed hadrons
- MisID and combinatorial (data driven)

Dominant systematic uncertainties

- MisID background shape
- NP contribution from $B^0 \rightarrow D^* \tau \nu_\tau$
- data/simulation differences
- simulated samples size

$B^0 \rightarrow D^* \mu \nu_\mu$ angular analysis (II)

- SM fit using different FF parametrization:
 - BGL [Phys. Rev. Lett. **74** 4603 (1995)]
 - CLN [Nucl. Phys. **B530** 1 (1998)]
 - BLPR [Phys. Rev. **D97** 075011 (2018)]
- Statistical precision comparable (Run1 only) to latest B-factory measurements: (Phys. Rev. **D100** 052007 (2019), Phys. Rev. Lett. **123** 091801 (2019))



$B^0 \rightarrow D^* \mu \nu_\mu$ angular analysis (III)

- Ideally WC determined by means of a fit to data without any assumption about the the NP structure ... (EPJ C80 883 (2020))
- ... but is easier to search for specific NP models (e.g. Bhattacharya et. al. JHEP 05, 191 (2019))
- Studied different NP scenarios (plan to report fit results for each)

Expected (stat - Run1) uncertainty on WC

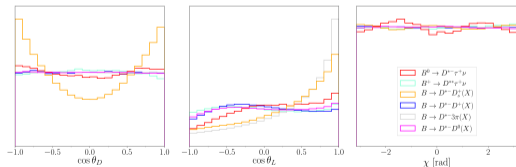
WC floating in fit

	VqRIL	VqLIL	SqRIL (SqLIL)	TqLIL
VqRIL	$\mathcal{I}m \mathcal{O}(10^{-2})$ $\mathcal{R}e \mathcal{O}(10^{-2})$			
VqLIL		$\mathcal{I}m \mathcal{O}(10^{-1})$ $\mathcal{R}e \text{ ---}$		
SqRIL (SqLIL)			$\mathcal{I}m \mathcal{O}(10^{-1})$ $\mathcal{R}e \mathcal{O}(10^{-1})$	
TqLIL				$\mathcal{I}m \mathcal{O}(10^{-3})$ $\mathcal{R}e \mathcal{O}(10^{-3})$
VqRIL+VqLIL+ SqRIL+ TqLIL	$\mathcal{I}m \mathcal{O}(10^{-2})$ $\mathcal{R}e \mathcal{O}(10^{-2})$	$\mathcal{I}m \mathcal{O}(10^0)$ $\mathcal{R}e \text{ ---}$	$\mathcal{I}m \mathcal{O}(10^{-1})$ $\mathcal{R}e \mathcal{O}(10^{-1})$	$\mathcal{I}m \mathcal{O}(10^{-3})$ $\mathcal{R}e \mathcal{O}(10^{-2})$

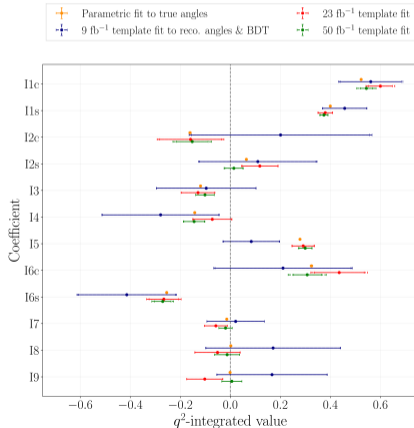
Uncertainties increase, generally within same order of magnitude, fits less stable →

$B^0 \rightarrow D^{(*)} \tau \nu_\tau$ angular analysis

- Ideally shape + rate analysis: $R(D)$ vs $R(D^*)$ determination simultaneous to WC
- Sensitivity studies need to include the full set of backgrounds
- Better angular resolutions with 3-prong hadronic τ decays [JHEP 11, 133 \(2019\)](#)
- Lower statistics than muonic case
- Large backgrounds
- External input for $R(D)$ & $R(D^*)$



[\[JHEP 11, 133 \(2019\)\]](#)

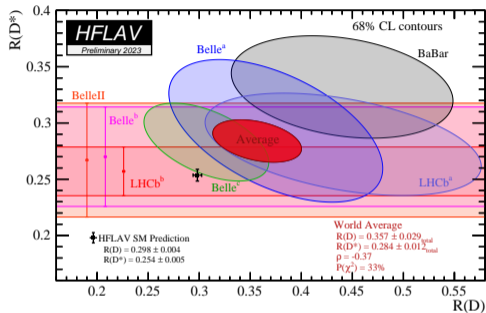


Conclusions

- **First measurement of $F_L^{D^*}$** with hadronic τ decays
 - smallest statistical uncertainty and performed in two q^2 bins
- $F_L^{D^*}$ compatible with the SM expectation
- **Combined $R(D) - R(D^*)$** is still a **3.2σ** tension from the SM
- Full picture to be given by ongoing differential measurements

More to come!

- $R(D^*)$ hadronic & $F_L^{D^*}$ with full Run 1 & Run 2
- Many LFU analysis: $R(D^0)$, $R(D^+)$, $R(D_s)$, $R(D^*)_e$
- Full angular analysis to determine spin and structure of NP
- LHCb Upgrade era has started:
⇒ **exciting time ahead!**



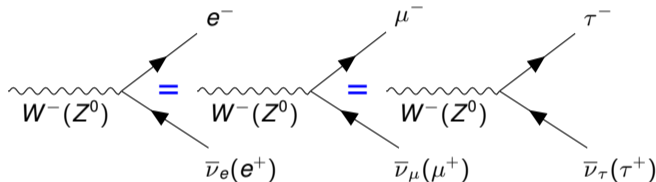
Thank you for your attention!

mail : davide.fazzini@cern.ch

Backup

Lepton Flavour Universality

- Within the Standard Model (SM), the weak interactions towards three generations of leptons are identical



- New physics (NP) may be more sensitive to the 3rd family
- Three typical candidates for NP:
 - **leptoquarks** [PRD 94, 115021, ...](#)
 - **two Higgs doublet models** [PRL 116, 081801, ...](#)
 - **Heavy vector bosons**, e.g. W' [JHEP 07 \(2015\) 142 1506.01705, ...](#)
- Need to cancel for theoretical uncertainties:
 - \implies **measure ratios of \mathcal{B}**

Charged current $b \rightarrow c\ell\nu_\ell$

- Main contribution:
tree level diagrams

$$R(X_c) \equiv \frac{\mathcal{B}(X_b \rightarrow X_c \tau^+ \nu_\tau)}{\mathcal{B}(X_b \rightarrow X_c \ell^+ \nu_\ell)}$$

$$X_b = B_{(s,c)}^{0,+}, \Lambda_b^0, \quad \ell = \mu, e,$$

$$X_c = D_{(s)}^{(*)}, J/\psi, \Lambda_c^+$$

Neutral transition $b \rightarrow s\ell\ell$

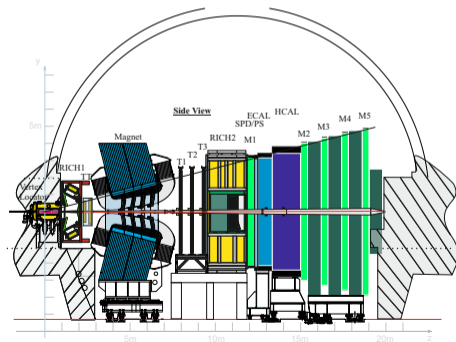
- Main contribution:
penguin or box diagrams

$$R(X_s) \equiv \frac{\mathcal{B}(X_b \rightarrow X_s \mu^+ \mu^-)}{\mathcal{B}(X_b \rightarrow X_s e^+ e^-)}$$

$$(X_b, X_s) = (B^0, K^{*0}) \text{ or } (B^+, K^+)$$

The Large Hadron Collider beauty (LHCb) Experiment

- LHCb detector is a single-arm forward spectrometer optimized for b and c hadron physics
 - pseudorapidity range: $[2,5] \Rightarrow \sim 25\% b\bar{b}$ pairs in LHCb acceptance
- **High precision measurements** in flavour physics (e.g. CKM, beyond SM)
- Collected data:
 - Run1 (2010-2012) $\Rightarrow \approx 3 \text{ fb}^{-1}$
 - Run2 (2015-2018) $\Rightarrow \approx 6 \text{ fb}^{-1}$
- Excellent performances
[Int. J. Mod. Phys. A 30, 1520022 (2015)]:
 - **Momentum resolution:**
 $\frac{\sigma_p}{p} \approx 0.5 - 0.8\%$ ($p < 100 \text{ GeV}/c$)
 - **Impact Parameter (IP) resolution:** $\sigma_{IP} \approx 20 \mu\text{m}$ (at high p_T)
 - **Decay time resolution:**
 $\sigma_t \approx 50 \text{ fs}$
 - **Particle Identification (PID):**
 $\varepsilon(K) \approx 95\%$, π mis-ID $\approx 5\%$ ($p < 100 \text{ GeV}/c$)
 $\varepsilon(\mu) \approx 97\%$, π mis-ID $\approx 1\text{-}3\%$

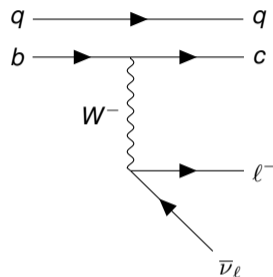


Semileptonic decays at LHCb

- Tree level \implies abundant
- Theoretically clean
 - factorize hadron and lepton current
 - hadron current described by form-factors
- Experimentally tricky due to ≥ 1 missing ν_ℓ

Advantages

- Large H_b production
- Produce B^+ , B^0 , B_s^0 , B_c^+ , Λ_b^0 , etc.
- Boosted H_b
- Efficient tracking \implies isolation

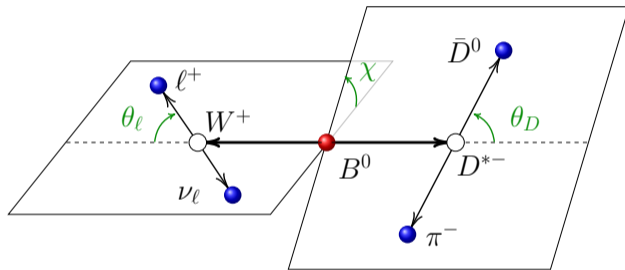


Challenges

- No beam-energy constraint
- Significant backgrounds
(combinatorial + partially reco)
- Reliance on simulation
 - large samples are required

Helicity angles definition

- θ_ℓ , θ_D and χ helicity angles enter directly in decay rate
- They provide (with $q^2 \equiv (p_{H_b} - p_{H_c})^2$) the highest sensitivity to New Physics effects



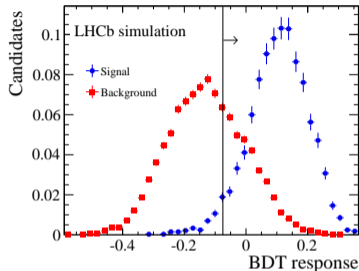
- θ_ℓ : angle between the direction of the lepton and the direction opposite the H_b meson in the virtual W rest frame
- θ_D : angle between the direction of the H_c^0 meson and the direction opposite the H_b meson in the H_c^0 rest frame
- χ : angle between the plane formed by H_c decay and the W decay in the H_b meson rest frame

Description of the anti- D_s BDT

- A dedicated BDT has been trained to suppress the abundant $B \rightarrow D^* D_s X$ background
- Trained performed separately for Run 1 and Run 2 data
- Signal described via simulated $B^0 \rightarrow D^{*-} \tau^+ \nu_\tau$ decays, corrected for data/MC differences
- Background described using simulated sample of $B \rightarrow D^* D_s X$ decays, where the D_s decays in 3π
- Output used in final fit to control $D^* D_s X$ background

Input features

- Output of the isolation algorithms
- Momenta, masses and quality of the reconstruction of the decay chain under the signal and background hypotheses
- Masses of oppositely charged pion pairs
- Energy and the flight distance in the transverse plane of the 3π system
- Mass of the total system

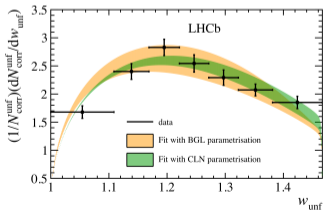


[Phys. Rev. D 97, 072013]

Differential measurement of $B_s \rightarrow D_s^* \mu \nu_\mu$ JHEP12(2020)144

- Measurement of the shape of the $B_s \rightarrow D_s^* \mu \nu_\mu$ decay rate
- Fully reconstruct $D_s^* \rightarrow D_s \gamma$ with $D_s \rightarrow KK\pi$
- Considering two FF schemes: **BGL** (PRL 74 4603 (1995)) and **CLN** (Nucl. Phys. B530 1 (1998))
- Signal yield measured in bins of hadronic recoil parameter

$$W = \nu_{B_s} \cdot \nu_{D_s^*}$$



CLN fit	
Unfolded fit	$\rho^2 = 1.16 \pm 0.05 \pm 0.07$
Unfolded fit with massless leptons	$\rho^2 = 1.17 \pm 0.05 \pm 0.07$
Folded fit	$\rho^2 = 1.14 \pm 0.04 \pm 0.07$
BGL fit	
Unfolded fit	$a_1^f = -0.005 \pm 0.034 \pm 0.046$ $a_2^f = 1.00^{+0.00+0.00}_{-0.19-0.38}$
Folded fit	$a_1^f = 0.039 \pm 0.029 \pm 0.046$ $a_2^f = 1.00^{+0.00+0.00}_{-0.13-0.34}$

