

Search for Pentaquarks Decaying to Open Charm Final States

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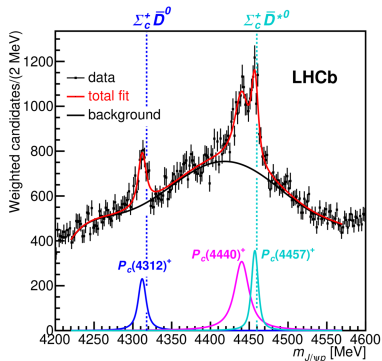
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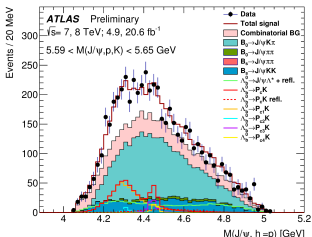
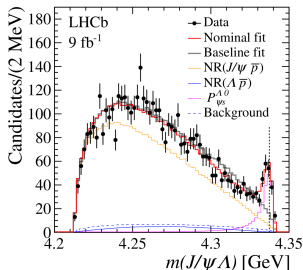
- Search for pentaquarks in wide range of charm hadron combinations (32 in total). - LHCb-PAPER-2023-018 in preparation.
- Observation of $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{(*)0} K^-$ Decay. - LHCb-PAPER-2023-034 in preparation.
- Both in final stages of internal review.
- [LHCb public results](#)
- Prospects in Run 3 and beyond.

Motivation

- Exotic hadrons are ones which have quark content that is not $q\bar{q}$ or qqq .
- Proposed in Gell–Mann’s quark model paper ([Phys. Lett.\(8\) 3 \(1964\)](#)).
- First observation of a pentaquark was by LHCb in 2015 in $J/\psi p$ mass spectrum. ([PRL \(115\), 072001 \(2015\)](#), [PRL \(122\), 222001 \(2019\)](#)).

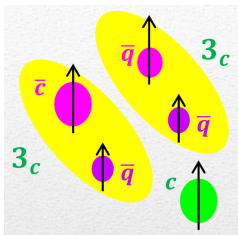


- First pentaquark with strangeness in $J/\psi \Lambda$ mass spectrum (PRL (131) 031901).



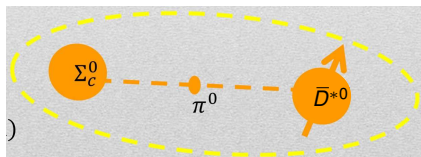
- ATLAS sees states consistent with LHCb pentaquarks (ATLAS-CONF-2019-048).
- Unable to rule out null hypothesis.

What is the structure of the observed pentaquark states?



Compact model:

- Fair understanding of existing spectrum.
- But doesn't explain proximity to threshold.
- Possibly charmonium-like states?



Molecular model:

- Proximity to threshold is natural.
- States are mostly unrelated to each other.
- States would be an order of 10x larger than compact counterparts.

Search for Prompt Production of Pentaquarks in Open-Charm Hadron Final States

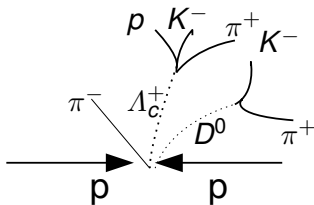
- Observed pentaquarks are close to mass threshold of some charm baryon-meson combinations.
[PRD \(101\), 074030](#)
- Structure is not clear - could they be molecules ([PRD \(101\), 054037](#)) or are they compact ([EPJ A \(56\), 142](#)) states?
- Goal was to search for pentaquark decays into a range of combinations of open-charm Σ_c or Λ_c^+ baryons with D mesons.
- Since no signal was seen, upper limits (ULs) are instead set in each mode (relative to the $\Lambda_c^+ \rightarrow pK^-\pi^+$ normalisation channel).

$$R = \frac{N_P}{N_{\Lambda_c^+}} \times \frac{\epsilon_{\Lambda_c^+}}{\epsilon_P}$$

- ULs set as a function of mass.

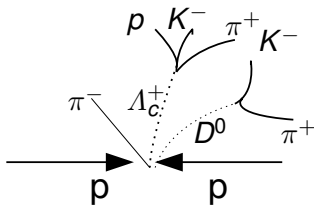
$\Sigma_c^{++} \bar{D}^0$	$\Sigma_c^{++} D^0$	$\Sigma_c^{++} D^-$	$\Sigma_c^{++} D^+$	$\Sigma_c^{++} D^{*-}$	$\Sigma_c^{++} D^{*+}$
$\Sigma_c^0 \bar{D}^0$	$\Sigma_c^0 D^0$	$\Sigma_c^0 D^-$	$\Sigma_c^0 D^+$	$\Sigma_c^0 D^{*-}$	$\Sigma_c^0 D^{*+}$
$\Sigma_c^{*++} \bar{D}^0$	$\Sigma_c^{*++} D^0$	$\Sigma_c^{*++} D^-$	$\Sigma_c^{*++} D^+$	$\Sigma_c^{*++} D^{*-}$	$\Sigma_c^{*++} D^{*+}$
$\Sigma_c^{*0} \bar{D}^0$	$\Sigma_c^{*0} D^0$	$\Sigma_c^{*0} D^-$	$\Sigma_c^{*0} D^+$	$\Sigma_c^{*0} D^{*-}$	$\Sigma_c^{*0} D^{*+}$
$\Lambda_c^+ \bar{D}^0$	$\Lambda_c^+ D^0$	$\Lambda_c^+ D^-$	$\Lambda_c^+ D^+$	$\Lambda_c^+ D^{*-}$	$\Lambda_c^+ D^{*+}$
$\Lambda_c^+ \bar{D}^0 \pi^+$	$\Lambda_c^+ D^0 \pi^+$	$\Lambda_c^+ D^- \pi^+$	$\Lambda_c^+ D^+ \pi^+$	$\Lambda_c^+ D^{*-} \pi^+$	$\Lambda_c^+ D^{*+} \pi^+$
$\Lambda_c^+ \bar{D}^0 \pi^-$	$\Lambda_c^+ D^0 \pi^-$	$\Lambda_c^+ D^- \pi^-$	$\Lambda_c^+ D^+ \pi^-$	$\Lambda_c^+ D^{*-} \pi^-$	$\Lambda_c^+ D^{*+} \pi^-$

- Range of cc ($\Lambda_c^+ D^0$, $\Sigma_c D^+ \dots$) $c\bar{c}$ ($\Lambda_c^+ D^-$, $\Sigma_c \bar{D}^0$) modes, as well as range of total charge.
- ULs are not set in all modes - some statistical limitations.

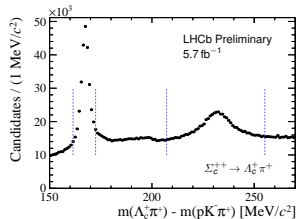
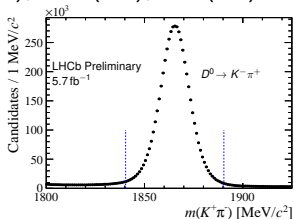
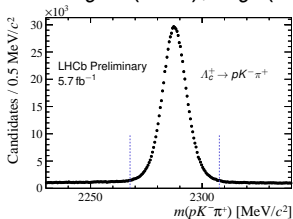


$\Sigma_c^{++} \bar{D}^0$	$\Sigma_c^{++} \bar{D}^0$	$\Sigma_c^{++} D^-$	$\Sigma_c^{++} D^+$	$\Sigma_c^{++} D^{*-}$	$\Sigma_c^{++} D^{*+}$
$\Sigma_c^0 \bar{D}^0$	$\Sigma_c^0 \bar{D}^0$	$\Sigma_c^0 D^-$	$\Sigma_c^0 D^+$	$\Sigma_c^0 D^{*-}$	$\Sigma_c^0 D^{*+}$
$\Sigma_c^{*++} \bar{D}^0$	$\Sigma_c^{*++} D^0$	$\Sigma_c^{*++} D^-$	$\Sigma_c^{*++} D^+$	$\Sigma_c^{*++} D^{*-}$	$\Sigma_c^{*++} D^{*+}$
$\Sigma_c^{*0} \bar{D}^0$	$\Sigma_c^{*0} D^0$	$\Sigma_c^{*0} D^-$	$\Sigma_c^{*0} D^+$	$\Sigma_c^{*0} D^{*-}$	$\Sigma_c^{*0} D^{*+}$
$\Lambda_c^+ \bar{D}^0$	$\Lambda_c^+ D^0$	$\Lambda_c^+ D^-$	$\Lambda_c^+ D^+$	$\Lambda_c^+ D^{*-}$	$\Lambda_c^+ D^{*+}$
$\Lambda_c^+ \bar{D}^0 \pi^+$	$\Lambda_c^+ D^0 \pi^+$	$\Lambda_c^+ D^- \pi^+$	$\Lambda_c^+ D^+ \pi^+$	$\Lambda_c^+ D^{*-} \pi^+$	$\Lambda_c^+ D^{*+} \pi^+$
$\Lambda_c^+ \bar{D}^0 \pi^-$	$\Lambda_c^+ D^0 \pi^-$	$\Lambda_c^+ D^- \pi^-$	$\Lambda_c^+ D^+ \pi^-$	$\Lambda_c^+ D^{*-} \pi^-$	$\Lambda_c^+ D^{*+} \pi^-$

- 10 modes are too statistically limited to set UL.
- Leaves 32.
- *N.B.* Excited doubly charmed baryons could also appear in these spectra.

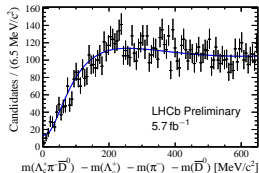
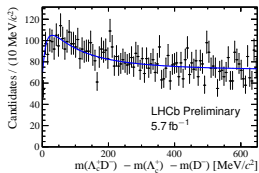
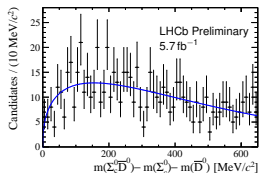


- 2016 - 2018 data set (5.7 fb^{-1}).
- Baryons/mesons are built in trigger, and optimised individually. Optimisation then applied in signal combination.
- Simulation samples used to optimise selection and train multivariate algorithms.
- $\Sigma_c^{++(0)} \rightarrow \Lambda_c^+ \pi^{+(-)}$, $D^{*-} \rightarrow \bar{D}^0 \pi^-$, $D^- \rightarrow K^+ \pi^- \pi^-$
- Σ_c^{++} (uuc), Λ_c^+ (udc), D^- ($d\bar{c}$), D^0 ($c\bar{u}$)



Fitting Procedure

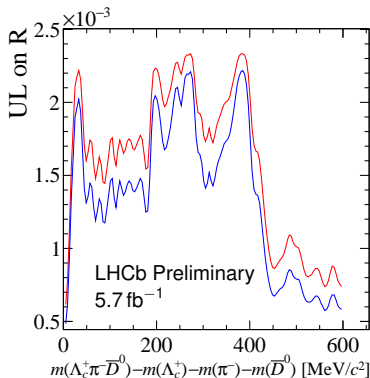
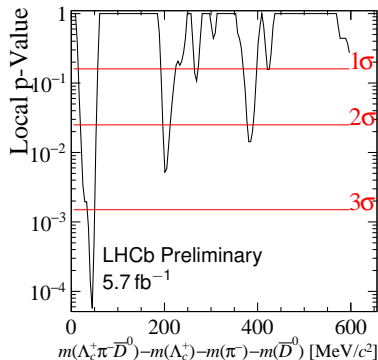
- Fit is done simultaneously to open-charm signal and sideband samples.
- $\Sigma_c^{(*)} D$ modes use threshold function for background.
- $\Lambda_c^+ D$ and $\Lambda_c^+ \pi D$ modes use Chebychev polynomial summed with log normal distribution for background.



- Range of signal models - Gaussian, 5, 10, 15 MeV/c² Voigtians.
- Gives some sensitivity to a state with broader width.

Scanning Method

- In $4 \text{ MeV}/c^2$ steps from threshold to $+600 \text{ MeV}/c^2$ above.
- Fit simultaneously to signal and sideband regions, split by trigger categories, find **local** p -value and UL of fit.
- UL smeared by systematic Gaussian.



Results

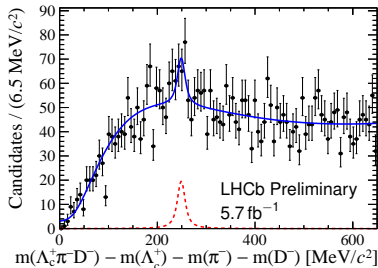
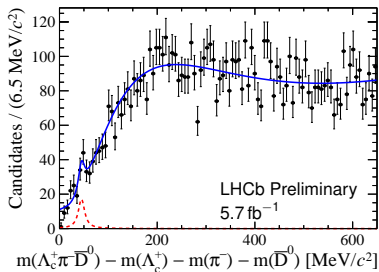
- Result is 4 tables (part of $15 \text{ MeV}/c^2$ shown here), with highest local and global significances with corresponding mass.

Decay Mode	Lowest p -value		Significance (σ)		Corresponding Mass (MeV/c^2)	Signal Yield	Upper Limit ($\times 10^{-3}$)	
	Local	Global	Local	Global			90% CL	95% CL
$\Lambda_c^+ \bar{D}^0$	2.18×10^{-3}	0.16	2.85	1.01	349	46.8 ± 23.4	1.16	1.21
$\Lambda_c^+ D^{*-}$	1.02×10^{-2}	1.00	2.32	0.00	365	15.0 ± 10.3	2.16	2.39
$\Lambda_c^+ \pi^+ D^-$	2.42×10^{-3}	0.16	2.82	0.99	225	68.6 ± 13.3	1.95	2.40
$\Sigma_c^0 \bar{D}^0$	2.85×10^{-2}	1.00	1.90	0.00	65	4.7 ± 4.2	1.02	1.15
$\Lambda_c^+ \pi^- \bar{D}^0$	5.69×10^{-5}	5.23×10^{-3}	3.86	2.56	45	60.1 ± 25.9	1.40	1.70
$\Sigma_c^0 D^-$	2.10×10^{-2}	1.00	2.03	0.00	261	7.0 ± 2.6	0.71	0.89
$\Lambda_c^+ \pi^- D^-$	1.20×10^{-4}	9.49×10^{-3}	3.67	2.35	249	82.8 ± 14.3	2.23	2.67
$\Lambda_c^+ \pi^- D^{*-}$	1.04×10^{-2}	1.00	2.31	0.00	409	23.6 ± 23.0	2.79	3.28
$\Sigma_c^{*++} D^{*-}$	4.10×10^{-2}	1.00	1.74	0.00	453	3.3 ± 2.4	1.24	1.43
$\Sigma_c^{*0} D^-$	3.12×10^{-2}	1.00	1.86	0.00	109	10.7 ± 29.1	1.32	1.59
$\Lambda_c^+ D^+$	5.84×10^{-3}	0.28	2.52	0.59	169	14.9 ± 9.6	1.34	1.50
$\Lambda_c^+ \pi^+ D^0$	6.60×10^{-4}	0.04	3.21	1.72	45	24.8 ± 39.3	0.98	1.18
$\Lambda_c^+ \pi^+ D^{*+}$	3.74×10^{-4}	0.02	3.37	1.99	165	13.8 ± 3.5	0.97	1.22
$\Lambda_c^+ \pi^- D^{*+}$	3.44×10^{-3}	0.28	2.70	0.58	73	5.8 ± 71.3	1.70	1.94
$\Sigma_c^{*++} D^0$	1.76×10^{-2}	1.00	2.11	0.00	113	3.9 ± 2.8	0.87	0.99
$\Sigma_c^{*0} D^+$	1.46×10^{-2}	1.00	2.18	0.00	69	4.7 ± 4.6	1.13	1.32

- Result is 4 tables (part of $15 \text{ MeV}/c^2$ shown here), with highest local and global significances with corresponding mass.

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$\Sigma_c^{*++} D^0$	1.76×10^{-2}	1.00	2.11	0.00	113	3.9 ± 2.8	0.87	0.99
$\Sigma_c^0 D^+$	1.46×10^{-2}	1.00	2.18	0.00	69	4.7 ± 4.6	1.13	1.32

- Highest **global** significances are 2.56σ and 2.35σ in $\Lambda_C^+ \pi^- \bar{D}^0$ and $\Lambda_C^+ \pi^- D^-$ modes.



- Modes would be interesting to look at with larger data set in near future.
- Paper available on arXiv very soon!

- Where kinematically possible, the width and mass of the known pentaquarks can be fitted.

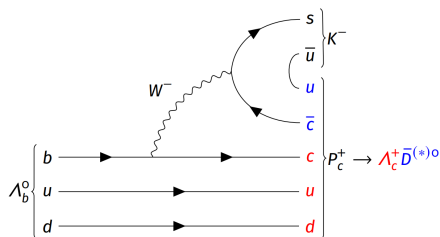
Decay Mode (Threshold)	Pentaquark Hypothesis				
$\Lambda_c^+ \bar{D}^0$ (4151.29 MeV)	$P_c(4312)^+$	$P_c(4312)^+$	$M = 4311.9 \text{ MeV}, \Gamma = 10 \text{ MeV}$		
	$P_c(4440)^+$			$P_c(4440)^+$	$M = 4440 \text{ MeV}, \Gamma = 21 \text{ MeV}$
	$P_c(4457)^+$			$P_c(4457)^+$	$M = 4457.3 \text{ MeV}, \Gamma = 6.4 \text{ MeV}$
$\Sigma_c^0 D^-$ (4323.41 MeV)	$P_c(4440)^+$	$P_c(4457)^+$	$M = 4457.3 \text{ MeV}, \Gamma = 6.4 \text{ MeV}$		
	$P_c(4457)^+$				
$\Lambda_c^+ \pi^+ D^{*-}$ (4436.32 MeV)	$P_c(4440)^+$	$P_c(4457)^+$			
	$P_c(4457)^+$				

- Only consider states with hidden charm content.
- Signal yield consistent with 0 in all cases.
- Full details in paper.

Observation of $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{(*)0} K^-$ Decay



- Comparable to $\Lambda_b^0 \rightarrow J/\psi p K^-$ since P_c production is the same.



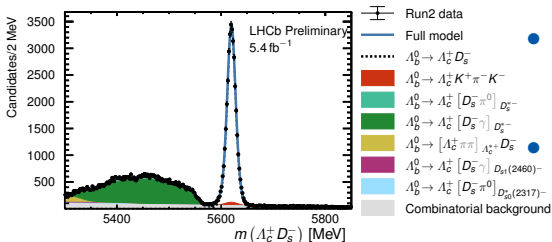
- To probe, we need:

► Fit fraction - $f_{\Lambda_c^+ \bar{D}^{(*)0}}$

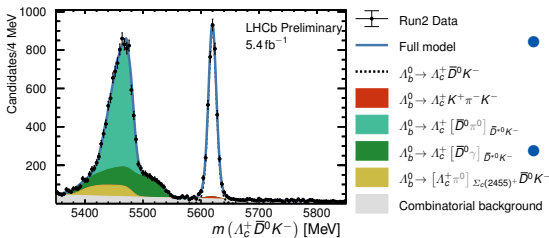
►
$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi p K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{(*)0} K^-)}$$

- BFs relative to $\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-$ found.

- 2016-18 Run 2 dataset (5.4 fb^{-1}).
- Measure $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{(*)0} K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)} = \frac{N_{\Lambda_c^+ \bar{D}^{(*)0} K^-}}{N_{\Lambda_c^+ D_s^-}} \frac{\epsilon_{\Lambda_c^+ \bar{D}^{(*)0} K^-}}{\epsilon_{\Lambda_c^+ D_s^-}} \frac{\mathcal{B}(D_s^- \rightarrow K^- K^+ \pi^-)}{\mathcal{B}(\bar{D}^0 \rightarrow K^+ \pi^-)}$
- Reconstruction and trigger selection identical between $\Lambda_c^+ \bar{D}^{(*)0} K^-$ and $\Lambda_c^+ D_s^-$ decay modes.
- Topological and particle identification requirements different between modes.
- BDT used to clean combinatorial background in $\Lambda_c^+ \rightarrow p K^- \pi^+$ spectrum.
- $\Lambda_c^+ \bar{D}^{*0} K^-$ selected with partially reconstructed $\bar{D}^{*0} \rightarrow \bar{D}^0 \pi^0 (\bar{D}^0 \gamma)$.
- 4 selection strategies, 3 background subtraction methods, 3 weighting methods - validate result.



- $\Lambda_c^+ \bar{D}^{*0} K^-$ selected with partially reconstructed $\bar{D}^{*0} \rightarrow \bar{D}^0 \pi^0 (\bar{D}^0 \gamma)$
- Many other partially reconstructed decays accommodated in the spectra.



- $\Lambda_b^0 \rightarrow \Sigma_c(2455)^+ \bar{D}^0 K^-$ is color suppressed, yet shows significant component.
- Possible enhancement from Ξ_c^{*0} , or $P_c?$

$$N^{\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^{*-}} = 46\,400 \pm 500(\text{stat.}), \quad N^{\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-} = 35\,450_{-210}^{+200}(\text{stat.})$$

$$N^{\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{*0} K^-} = 10\,560_{-290}^{+310}(\text{stat.}), \quad N^{\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0 K^-} = 4\,010 \pm 70(\text{stat.})$$

- Converting to BFs:

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D^0 K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)} = 0.1908_{-0.0034}^{+0.0036}(\text{stat.})_{-0.0018}^{+0.0016}(\text{syst.}) \pm 0.0038(\mathcal{B})$$

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{*0} K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)} = 0.589_{-0.017}^{+0.018}(\text{stat.})_{-0.018}^{+0.017}(\text{syst.}) \pm 0.012(\mathcal{B})$$

- Finally relative to $J/\psi p K^-$ (using [PRL 122, 222001](#)):

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi p K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0 K^-)} = 0.152_{-0.028}^{+0.032} \frac{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi p K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{*0} K^-)} = 0.049_{-0.009}^{+0.011}$$

Reference	$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^{*-})}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)}$
Z. Phys. C 59, 179	0.75
Phys. Rev. D 56, 2799	0.83
Mod. Phys. Lett. A 13, 23	1.54
Phys. Rev. D 58, 014016	1.46
Prog. Theor. Phys. 101, 959	1.84
Phys. Rev. D 99, 054020	0.85
Chin. Phys. C 42, 093101	1.49
Eur. Phys. J. C 78, 528	1.23
Phys. Rev. D 98, 074011	1.70
Eur. Phys. J. C 79, 540	1.51
Phys. Rev. D 100, 034025	1.47
Eur. Phys. J. C 80, 636	1.29
arXiv:2309.12050	2.25

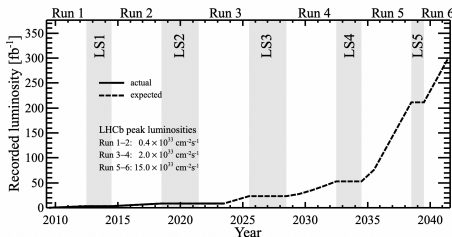
- Also make a measurement of $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^{*-})}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)}$.
- Consistent with several theory predictions.

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^{*-})}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)} = 1.668 \pm 0.022(\text{stat.})_{-0.055}^{+0.061}(\text{syst.})$$

Conclusions

- Search for pentaquarks in wide range of signal channels (LHCb-PAPER-2023-018).
 - ▶ No substantial evidence of signal.
 - ▶ Some hints of interesting areas.
 - ▶ 2.35σ in $\Lambda_c^+ \pi^- D^-$, 2.56σ in $\Lambda_c^+ \pi^- \bar{D}^0$.
 - ▶ With Run 3 data a study of these regions could be interesting.
 - ▶ No signal seen for decays from known pentaquark states.
- Observation of $\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{(*)0} K^-$ decay (LHCb-PAPER-2023-034).
 - ▶ $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi p K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^0 K^-)} = 0.152_{-0.028}^{+0.032}$
 - ▶ $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow J/\psi p K^-)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \bar{D}^{*0} K^-)} = 0.049_{-0.009}^{+0.011}$
 - ▶ $\frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^{*-})}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ D_s^-)} = 1.668 \pm 0.022(stat.)_{-0.055}^{+0.061}(syst.)$ - consistent with several theory predictions.
 - ▶ Future analysis to determine $f_{\Lambda_c^+ \bar{D}^{(*)0}}$ will allow to test model predictions.

- More results still to come from LHCb Run 2 dataset.
 - ▶ Many pentaquark analyses nearing completion.
- LHCb Run 3 has just begun:
 - ▶ New data set is on the horizon - will be $\sim 10x$ (by end of Run 4) larger than the existing dataset.
 - ▶ Improved trigger means cleaner hadronic events.
 - ▶ Statistically limited pentaquark searches will be within our grasp.

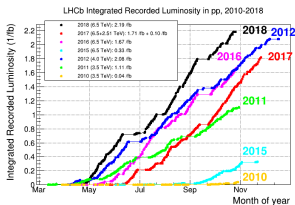
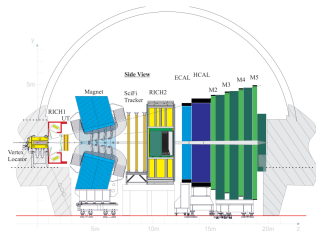


LHCb RUN 3 PROSPECTS

LHCb
LHCb

Backup slides

- Largest heavy-flavour dataset collected – 9 fb^{-1} with Run 1+2.
- Largest production cross-sections of b and c hadrons.
- Specialised trigger for hadronic decays.



- Two Ring Imaging Cherenkov (RICH) detectors allow excellent PID.
- Vertex Locator (VELO) and tracking stations allow precise tracking of particles.

- For $\Sigma_c D$ and $\Sigma_c^* D$ modes the fit model is a threshold function:

$$f(x) = x^\gamma \cdot \exp(-p_1 \cdot x) \quad (1)$$

- $\Lambda_c^+ D$ and $D\Lambda_c^+ \pi$ modes use a Chebychev polynomial summed with a log normal distribution:

$$f(x) = f \cdot \frac{1}{2\pi \cdot \ln k * x} \cdot \exp\left(\frac{-\ln^2\left(\frac{x}{m_0}\right)}{2 \ln^2 k}\right) + (1 - f) \cdot C(x, x_1) \quad (2)$$