Implications of first LHC results

Large extra dimensions (http://arxiv.org/abs/1101.4919)
 SuperSymmetry (http://arxiv.org/abs/1101.2195)

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The main goal of LHC

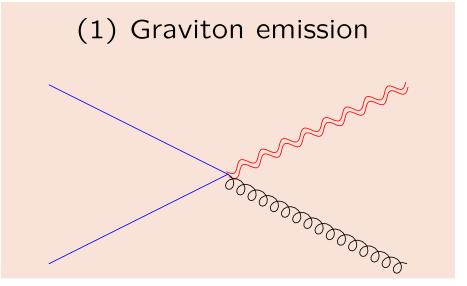
is understanding why the weak scale is small.

Maybe the hierarchy problem was a good guideline. Maybe we lost 30 years.

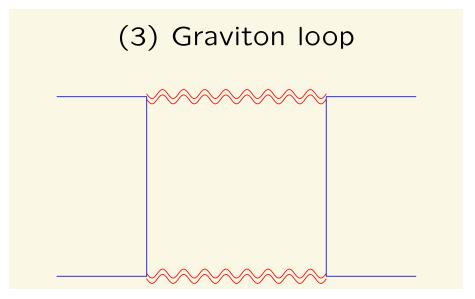
A new symmetry					
Scalar	Vector Fermion				
$H \rightarrow H + \theta$	$A_{\mu} \to A_{\mu} + \partial_{\mu}\theta$	$\Psi ightarrow e^{i heta \gamma_5} \Psi$			
Keeps <i>H</i> massless. Goldstone boson.	Keeps A_{μ} massless In 5 dims: $H = A_5$	Keeps Ψ massless. $H \stackrel{SUSY}{\leftrightarrow} \Psi.$			
Higgs has weak-scale size					
Technicolour	Large extra dims	Warped extra dims			
H bound state like π	H is a string or	Dual to technicolor			

Large extra dimensions

Collider signals



(2) Tree-level graviton exchange





No black holes in first LHC data





"È LA FINE DEL MONDO" La conferma del Cern

Il buco nero creato dal <u>Large Hadron Collider</u> è destinato a inghiottire il nostro pianeta nelle prossime settimane. Secondo <u>Walter Wagner</u> anche l'universo è in pericolo. <u>Scene di panico</u> durante la conferenza stampa del Cern a Parigi. <u>IL VIDEO</u> / <u>LE PRIME IMMAGINI</u>

La Svizzera inghiottita dal buco nero Sarkozy contrario all'evacuazione

"A quoi bon? Tra qualche ora saremo comunque tutti morti". <u>Berlusconi caustico:</u> "risolta la questione Alitalia".

<u>"Pregate, ma non garantisco niente"</u> Lo scetticismo di Papa Benedetto

"Il buco nero è il risultato ultimo di ricerca scientifica dissennata che ha voluto porre l'uomo davanti a Dio. Non resta che la preghiera, forse". Dal suo ritiro in Alto Adige, il Papa gela i pellegrini venuti ad ascoltare un messaggio di speranza. IL VIDEO / LE IMMAGINI / <u>"SANTITA", IL BUCO NERO E' L'INFERNO?"</u> <u>Net Monitor: "Proprio adesso che andavo in ferie"</u> di V. ZAMBARDINO

Hawking deluso: "mi sono sbagliato"

"La mia radiazione avrebbe dovuto eliminare il buco nero, ma evidentemente ho sbagliato qualcosa". Lo scienziato non esclude la possibilità che un nuovo universo possa riformarsi dal caos. "Ci vorranno milioni di anni, ma la prossima volta saremo più prudenti". Software Società Curiosità



Troppa pancia per Babbo Natale è tempo di dieta

Si sono incontrati in 150 per l'annuale congresso in Danimarca. E hanno deciso di perdere qualche chilo. Messa da parte la slitta, per qualche tempo andranno in bici e faranno ginnastica



LE IMMAGINI La dea del metrò colpisce ancora lap dance sui pali della strada

lap dance sui pali della strada



LE IMMAGINI Pechino indossa la sottoveste sono le olimpiadi della lingerie



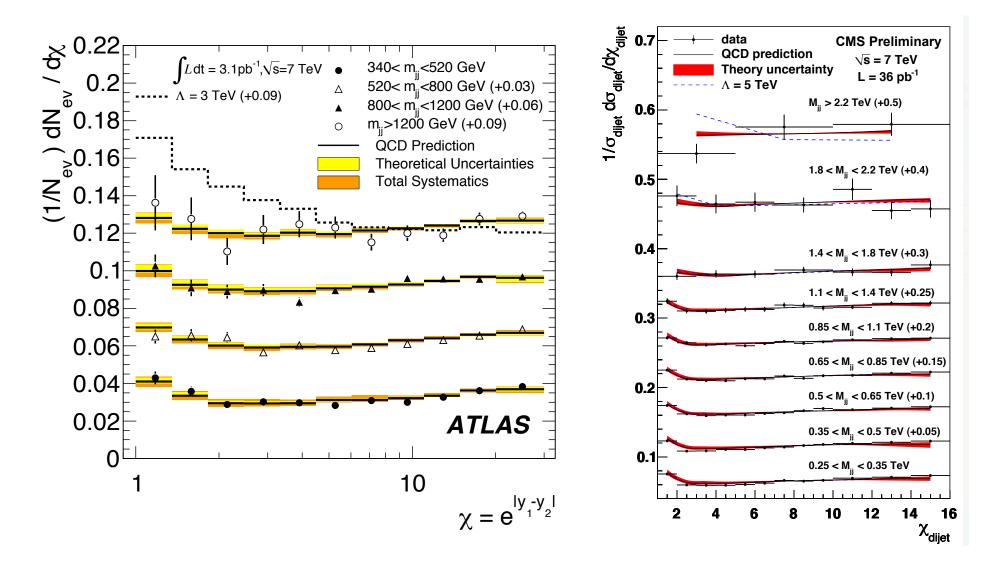
<mark>LE IMMAGINI</mark> Miracolo a Saint Tropez In acqua c'è Bar Rafaeli



LE IMMAGINI Pamela a Central Park spunta l'amico misterioso **First LHC data:** $pp \rightarrow jj$

ATLAS at 3.1/pb

CMS at 36/pb



 χ = angular jj distance. Coulomb-like QCD gives a quasi-flat distribution. New massive particles or effective O give effects at large M_{jj} and small angle χ

Loop level graviton exchange

At low energy is described by the dimension 6 effective operator

$$\mathscr{L} = c_{\Upsilon} \times \Upsilon, \quad \Upsilon = \frac{1}{2} (\sum_{f} \bar{f} \gamma_{\mu} \gamma_{5} f)^{2}$$

95% CL limits on $ c_{\Upsilon}/4\pi ^{-1/2}$ in TeV					
Experiment	Process	+			
LEP combined	$e^+e^- \rightarrow \ell^+\ell^-$	17.2	15.1		
LEP combined	$e^+e^- \rightarrow b\overline{b}$	15.3	11.5		
ZEUS, H1	e^+p and e^-p	4.6	5.3		
DØ	$p\bar{p} \rightarrow e^+e^-$	4.7	5.5		
CDF	$p\bar{p} \rightarrow \ell^+ \ell^-$	4.5	5.6		
CCFR	νN scattering	3.7	5.9		
DØ	$p \overline{p} ightarrow j j$	3.2	3.1		
ATLAS at 7 TeV with 3.1/pb	pp ightarrow jj	5.3	4.2		
CMS at 7 TeV with 36/pb	pp ightarrow jj	11	8.1		
combined			15.7		

LHC improves on TeVatron but not on LEP

Tree level graviton exchange

At low energy is described by the dimension 8 effective operator $\mathscr{L}_{\mathrm{eff}} = c_{\mathcal{T}} \ \mathcal{T}$

$$\mathcal{T} = \frac{1}{2} \left(T_{\mu\nu} T^{\mu\nu} - \frac{\overline{T^{\mu}_{\mu} T^{\nu}_{\nu}}}{\delta + 2} \right)^2 \sim (\bar{\Psi} \partial \Psi + \bar{\Psi} A \Psi + F^2_{\mu\nu})^2$$

Independently produced by brane fluctuations at loop level.

High dimensionality: energy is the key factor and early LHC wins

Coefficient: parameterize as $c_T = 8/M_T^4$ (Hewett normalization).

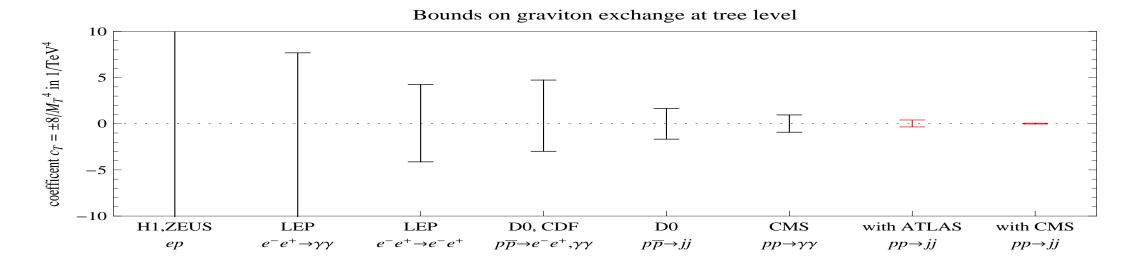
Signals: $pp \rightarrow jj$ is considered dirty by actually is much better at low statistics:

$$\sigma = \left(\frac{2 \,\mathrm{TeV}}{M_{\mathcal{T}}}\right)^8 \times \begin{cases} 12.5 \,\mathrm{pb} & \mathrm{for} \ pp \to jj \\ 10.4 \,\mathrm{fb} & \mathrm{for} \ pp \to \mu^+ \mu^- \\ 21.3 \,\mathrm{fb} & \mathrm{for} \ pp \to \gamma\gamma \end{cases}$$

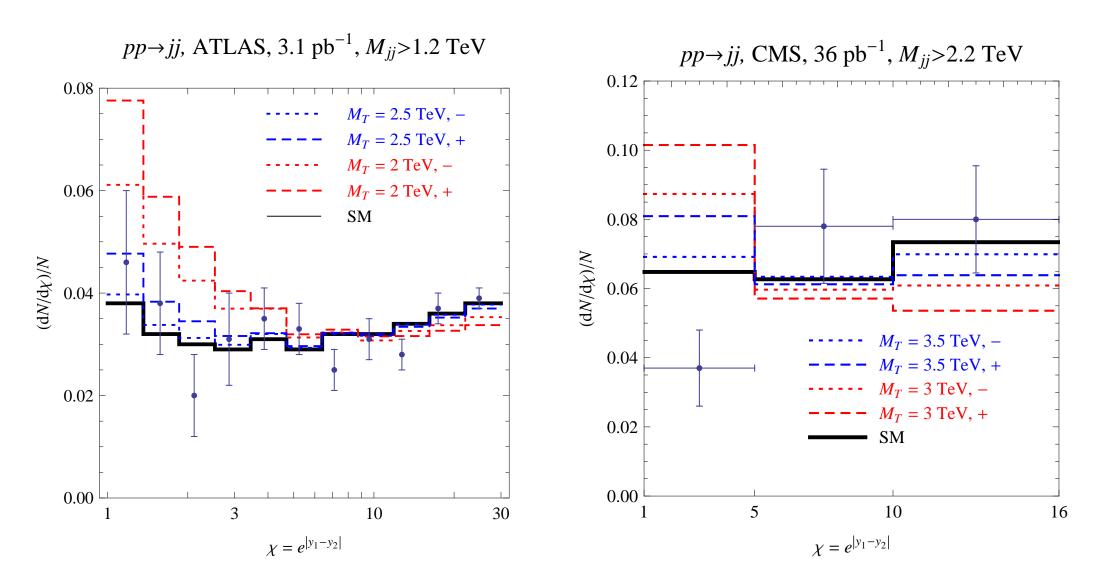
thanks to the energetic $uu \rightarrow uu$. (cuts: $\sqrt{s} = 7 \,\text{TeV}$, $M_{\text{eff}} > 1 \,\text{TeV}$, $\eta < 2.5$)

Bounds on $M_{\mathcal{T}}$

Experiment	Process	+	_
LEP	$e^+e^- \rightarrow \gamma\gamma$	0.93 TeV	1.01 TeV
LEP	$e^+e^- \rightarrow e^+e^-$	1.18 TeV	
H1, ZEUS	e^+p and e^-p	0.74 TeV	
CDF	$p \bar{p} \rightarrow e^+ e^-, \gamma \gamma$	0.99 TeV	0.96 TeV
DØ	$pp \rightarrow e^-e^-, \gamma\gamma$ $p\overline{p} \rightarrow e^+e^-, \gamma\gamma$	1.28 TeV	
	$pp \rightarrow e \cdot e^{-}, \gamma \gamma$ $p\overline{p} \rightarrow jj$		1.48 TeV
CMS at 7 TeV with 34/pb	$pp \rightarrow JJ$ $pp \rightarrow \gamma\gamma$	1.72 TeV	1.40 TeV
ATLAS at 7 TeV with 3.1/pb	pp ightarrow jj $pp ightarrow jj$	2.2 TeV	
CMS at 7 TeV with 36/pb	pp ightarrow jj pp ightarrow jj	4.2 TeV	
	$pp \rightarrow JJ$	T.2 IEV	J.T ICV



The LHC data



Computed implementing in MadGraph and Pythia

Fitting the full amplitude $\mathcal{A} = \mathcal{S}(s)T_{\mu\nu}^2$

Truncate KK tower at $m < \Lambda$ to avoid UV divergence:

$$S(s) = \frac{1}{M_{\text{Pl}}^2} \sum_i \frac{1}{s - m_i^2 + im_i \Gamma_G(m_i)} = \frac{1}{M_D^{2+\delta}} \int_{|q| < \Lambda} \frac{d^{\delta}q}{s - q^2 + i\varepsilon}$$

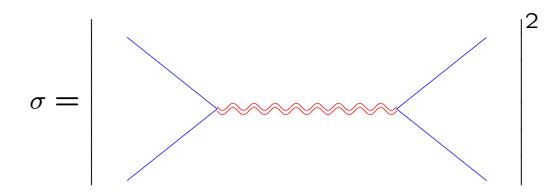
$$c_{\mathcal{T}} = S(s \ll \Lambda^2) = \begin{cases} \frac{\pi^{\delta/2}}{(1 - \delta/2)\Gamma(\delta/2)} \frac{\Lambda^{\delta-2}}{M_D^{\delta+2}} \equiv \frac{8}{M_{\mathcal{T}}^4} & \text{for } \delta > 2\\ \frac{\pi}{M_D^4} \ln \frac{s}{\Lambda^2} & \text{for } \delta = 2\\ \frac{-i\pi}{M_D^3\sqrt{s}} & \text{for } \delta = 1 \end{cases}$$

Subtelty: $S_{\int} = \langle S_{\Sigma} \rangle$ i.e. it is ok to $\Sigma_{KK} \approx \int_q$ and ignore the graviton width

$$\Gamma_G = \frac{283}{960} \frac{m^3}{\pi M_{\rm Pl}^2}$$

 $\langle |\mathcal{S}|^2 \rangle$ vs $|\langle \mathcal{S} \rangle|^2$: subtleties under the carpet Including Γ_G : $\langle |\mathcal{S}|^2 \rangle = (\operatorname{Re} \mathcal{S})^2 + (\operatorname{Im} \mathcal{S})^2 / \epsilon$ with $\epsilon = \pi \Gamma_G / 2\Delta m \sim (s/M_D^2)^{1+\delta/2}$.

Consider just one particle with coupling g ($g \sim E/M_{Pl} \ll 1$ for KK gravitons):



One would guess $\sigma \sim g^4$ but actually $\sigma \sim g^2$, due to $pp \rightarrow$ graviton. Next graviton decays with 100% probability and width $\Gamma \sim mg^2$.

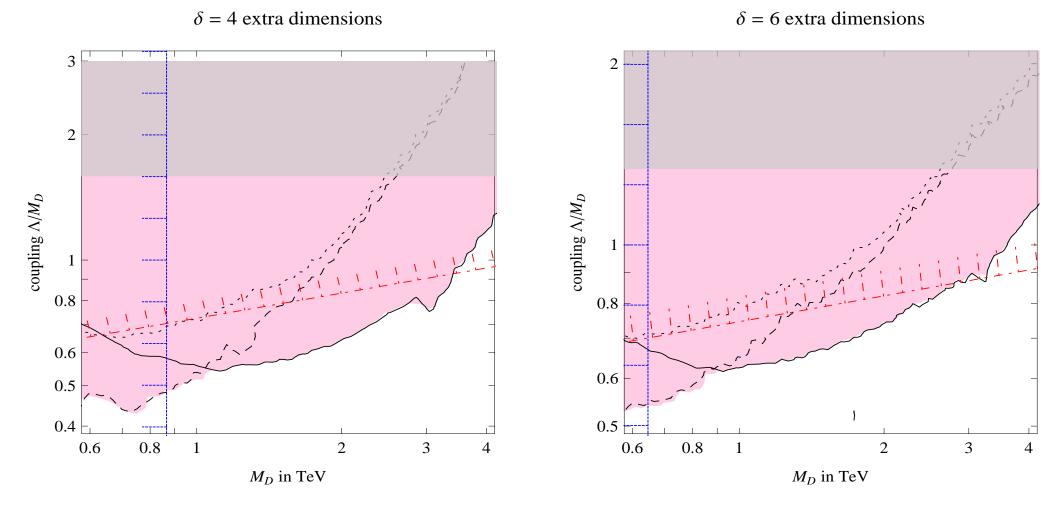
For flat extra dimensions Γ is small: gravitons decay far away from the detector.

Resonant graviton production must be subtracted

We find $\sigma_{\text{subtracted}} \sim -g^4$, up to $\mathcal{O}(g^4)$ terms, such as NLO corrections to Γ . We presume that $\langle |\mathcal{S}|^2 \rangle_{\text{subtracted}} = |\langle \mathcal{S} \rangle|^2$ is the right result

Anyhow, even the unsubtracted $1/g^2$ enhancement (present for $\delta = 1$) is numerically irrelevant for $pp \rightarrow jj$.

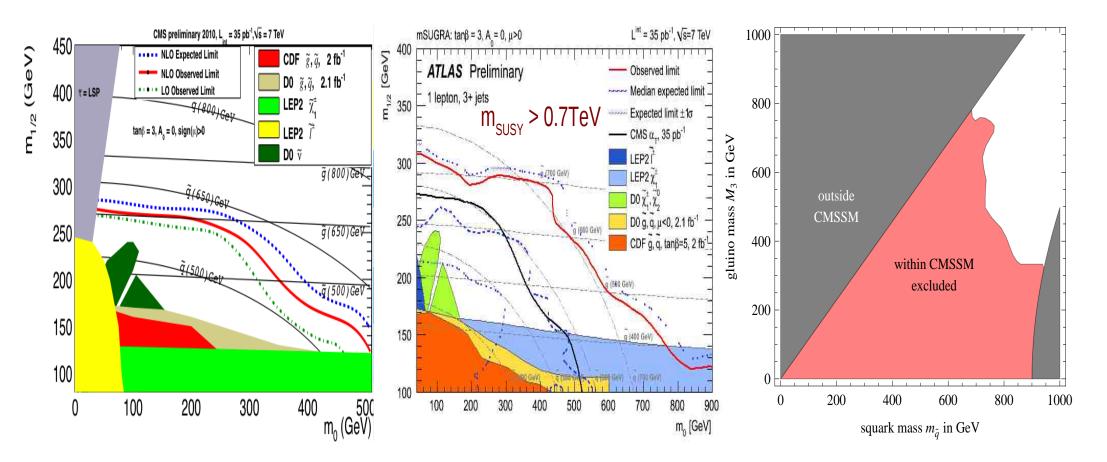
Bound on the full amplitude (M_D, Λ)



- Shaded: LHC (continuous = CMS; dashed = ALTAS; dotted ATLAS F_{χ})
- \bullet Blue: gravition emission, ignoring the dependence on Λ
- Red: NDA estimate of graviton loop
- Gray: non-perturbative quantum gravity



First LHC data



Actually a bound on $m_{\tilde{q}}$ and M_3 , up to 700 GeV in the CMSSM

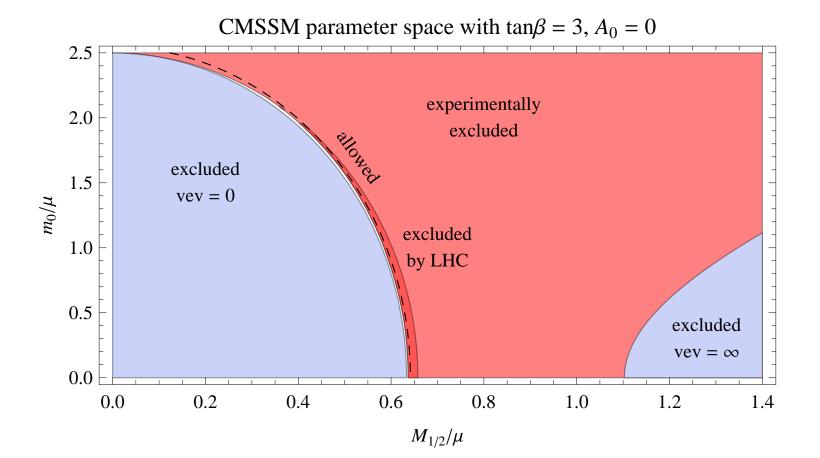
In the following we ignore $g_{\mu} - 2$ or Ω_{DM} ; they point to heavier m_{SUSY} .

The little hierarchy problem

Fix tan $\beta = 3$ and $A_0 = 0$; the overall SUSY mass scale is fixed by

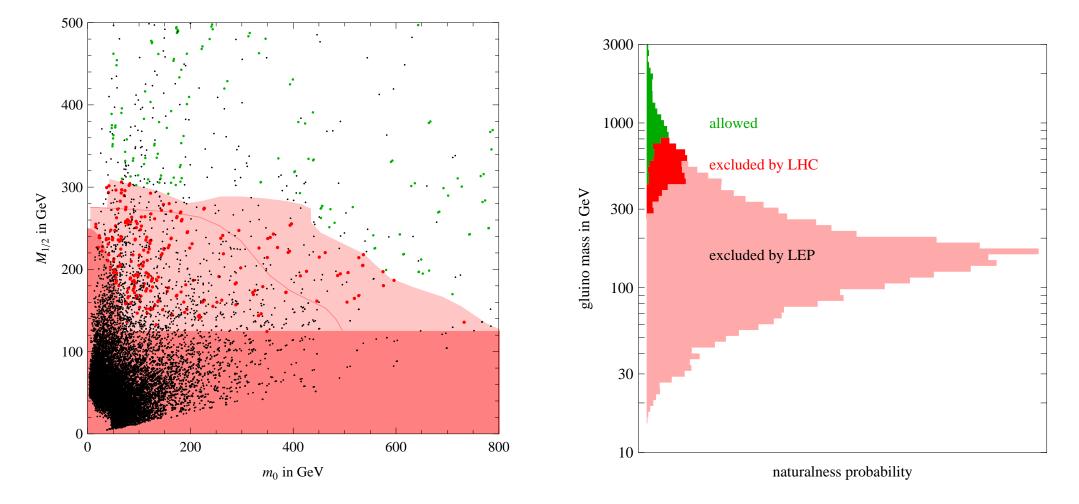
$$m_h^2 \stackrel{\text{tree level}}{\to} M_Z^2 \approx 0.2m_0^2 + 0.7M_3^2 - 2\mu^2 = (91 \,\text{GeV})^2 \times 35(\frac{M_3}{650 \,\text{GeV}})^2 + \cdots$$

Plot this in the plane of the adimensional free parameters $(M_{1/2}, m_0)/\mu$:



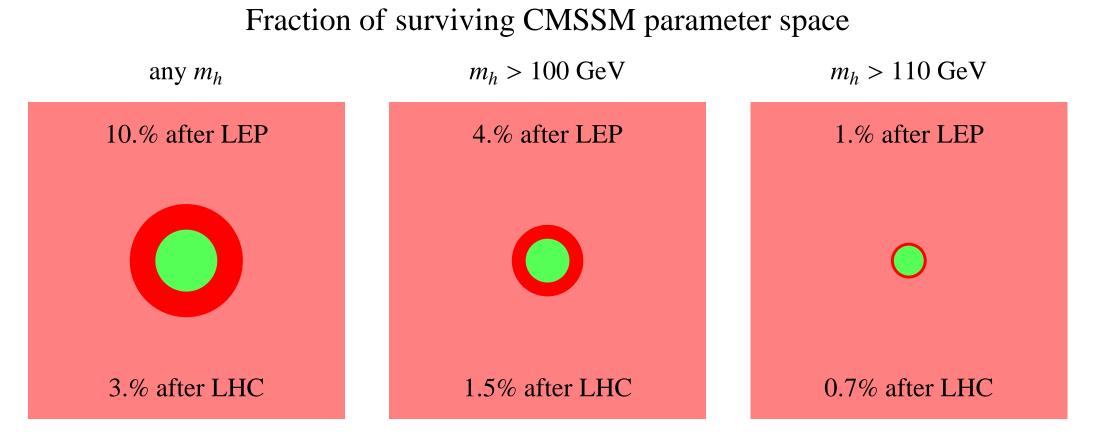
Bayesian MonteCarlo technique

Scan over all adimensional parameters $(m_0/\mu, B_0/m_0, ..., \lambda_t)$ compatible with measured m_t . Compute tan β and m_{SUSY} from V_{MSSM} . Normally $m_{SUSY} \sim M_Z$; rare accidents can make it bigger. All possible fine-tunings are included without using any explicit FT parameter Δ . E.g. focus point is fine-tuning of λ_t .



Black dot = excluded spectrum. Red = excluded by LHC. Green = allowed

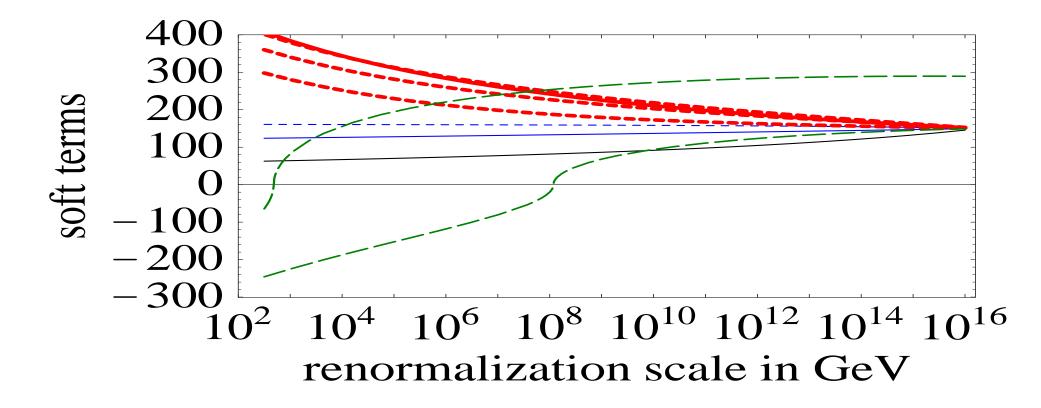
Fraction of alive CMSSM $\approx 1/\Delta$



(The CMSSM prediction for m_h can be circumvented; the theoretical uncertainty in m_h is about 3 GeV)

$M_Z \ll m_{SUSY} \Rightarrow$ late SU(2) breaking

The scale Q_0 at which RGE running makes $m_h^2(Q) < 0$ must be close to m_{SUSY}



SUSY little Higgs as pseudo-Goldstone of some symmetry broken at Q_0 ?

The sliding m_{SUSY} model

Assume that m_{SUSY} is a free parameter determined by minimizing V_{MSSM} (!)

$$\min V_{\text{MSSM}} \sim \begin{cases} 0 & m_{\text{SUSY}} > Q_0 \\ -m_{\text{SUSY}}^4 & m_{\text{SUSY}} < Q_0 \end{cases}$$

Prediction [hep-ph/0005203, BS]: $m_{SUSY} \lesssim Q_0$ and a loop factor above M_Z :

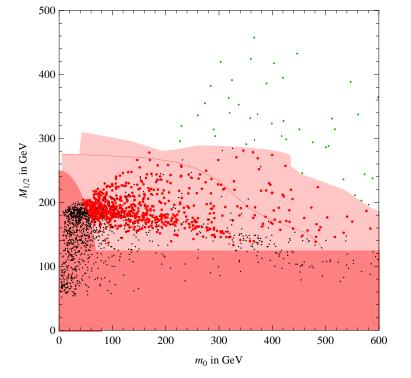
$$\frac{d\mu_{\rm u}^2}{d\ln\mu}\sin^2\beta + \frac{d\mu_{\rm d}^2}{d\ln\mu}\cos^2\beta - 2\frac{d\mu_{\rm ud}^2}{d\ln\mu}\sin\beta\cos\beta = M_Z^2\cos^22\beta \stackrel{\rm one\ loop}{\to} m_h^2$$

RGE loop factors are big: roughly this means

$$m_{\tilde{t}}\approx 4\pi M_Z/\sqrt{12}\approx 400\,{\rm GeV}$$

Predicted: dashed line in the CMSSM plot. Allowed: from 50% to 6% with LHC.

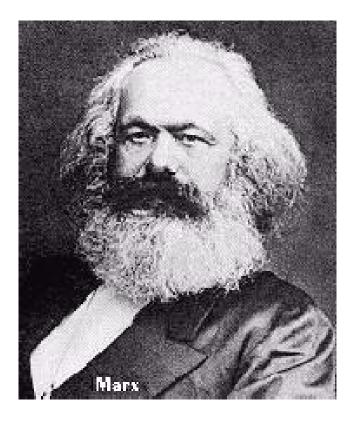
PS: BS hypothesis may be BS: $V \neq V_{MSSM}$. Alternative interpretation in terms of anthropic pressure ($Q_0 \ll m_{SUSY}$): more rare than SM?



Conclusions

These beautiful ideas and the history of the Michelson-Morley experiment teach us that a negative experimental search can have deep theoretical implications.

"History repeats itself, first as tragedy, second as farce".



Conclusions

"Now this is not the end. It is not even the beginning of the end. But it is, perhaps, the end of the beginning".

