

# Inclusive and diffractive dijet photoproduction in Pb-Pb UPCs at the LHC



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ERC adG YoctoLHC

Based on work in collaboration with **M. Klasen** (U. of Münster, Germany),  
JHEP 04 (2016) 158; EPJ C 76 (2016) 8, 467; PRC 99 (2019) 6, 065202; EPJ C 79 (2019) 5, 396;  
PRD 104 (2021) 11, 114013

## Outline:

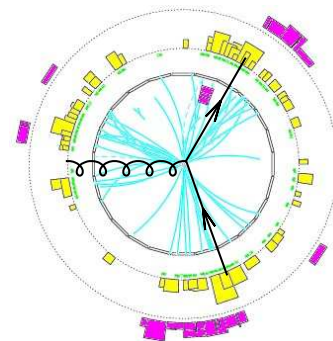
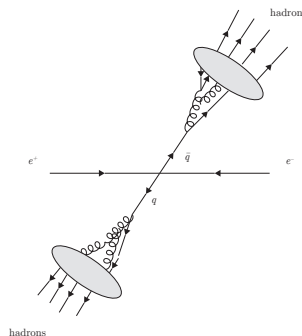
- Jets in QCD and dijet photoproduction in collinear factorization and NLO pQCD
- Inclusive dijet photoproduction in Pb-Pb UPCs at the LHC and nuclear PDFs
- Diffractive dijet photoproduction in Pb-Pb UPCs at the LHC: nuclear diffractive PDFs and QCD factorization breaking

# Jets in QCD

- **Jets** are collimated sprays of hadrons ( $\pi$ , K, p, ...) produced in high-energy  $e^+e^-$ , lepton-hadron, and hadron-hadron collisions, Salam, arXiv:1011.5131; Sapeta, arXiv:1511.09336; Laenen, arXiv:1708.00770

- **Jets** have been instrumental in establishing QCD and concepts of asymptotic freedom (factorization) and confinement (hadronization).

- Classic example: 3-jet events in  $e^+e^-$  annihilation  $\rightarrow$  existence of **gluons**.



- Studies of jet production remain an active field of QCD studies at colliders:

- precision extraction of  $\alpha_s(M_Z)$  from HERA data on ep DIS, H1 Coll. EPJC 75 (2015) 2, 65; ZEUS Coll., NPB 864 (2012) 1

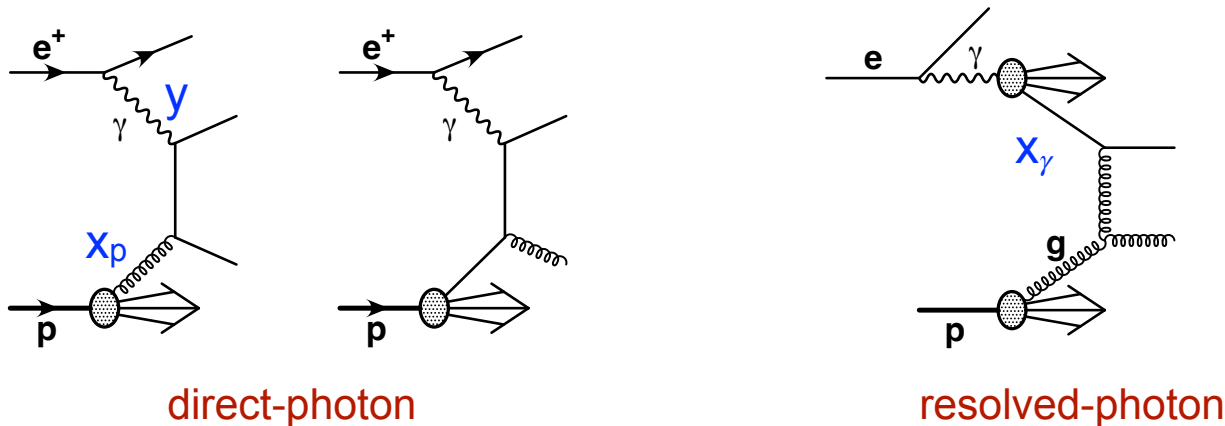
- global QCD fits of parton distribution functions (**PDFs**) of the proton using ep DIS@HERA, anti-pp@Tevatron, and pp@LHC, NNPDF4.0, EPJC 82 (2022) 4, 428; CT18, PRD 103 (2021) 1, 014013, **nuclear PDFs** using dijets in pA@LHC, EPPS21, EPJC 82 (2022) 5, 413, and **photon PDFs**, Slominski, Abramowicz, Levy, EPJC 45 (2006) 633.

- search for small- $x$  BFKL and saturation physics using forward dijet production at LHC, 2022 Snowmass Summer Study, arXiv:2203.08129 and at EIC, Boussarie, Mäntysaari, Salazar, Schenke, JHEP 09 (2021) 178

- Standard Model background for many new physics processes.

# Jet photoproduction in QCD

- All information on jet photoproduction comes from ep scattering at HERA, Newman, Wing, Rev. Mod. Phys. 86 (2014) 3, 1037; Butterworth, Wing, Rept. Prog. Phys. 68 (2005) 2773; Klein, Yoshida, Prog. Part. Nucl. Phys. 61 (2008) 343 + prelim. data on Pb-Pb UPCs@LHC, ATLAS-CONF-2017-011, ATLAS-CONF-2022-021
- Typical leading-order (LO) Feynman graphs: **direct-photon** and **resolved-photon** contributions. The separation is not unique beyond LO, but is still useful.



- **Main interests in studying jet photoproduction:**

- Cross section is sensitive to quark and gluon structure at the same order. When combined DIS cross section, provides of additional constraints on the **gluon PDF**, ZEUS, EPJC 42 (2005) 1
- Test of **QCD factorization** and its violation (in case of diffractive dijet photoproduction)
- Access to photon structure, constraints on the gluon **PDF of the photon**, which are complimentary to those from  $F_2^\gamma(x, Q^2)$  in  $e^+e^-$ , Nisius, Phys. Rept. 332 (2000) 165; Slominski, Abramowicz, Levy, EPJC 45 (2006) 633

# Dijet photoproduction in NLO pQCD

- In photoproduction, jet transverse momentum provides hard scale  $\mu_R = \mu_F = E_T$
- In framework of collinear factorization of perturbative QCD, cross sections are known to next-to-leading order (NLO) accuracy, [Aurenche et al. EPJC 17 \(2000\) 413](#); [Frixione, Ridolfi, NPB 507 \(1997\) 315](#); [Klasen, Kramer, Z. Phys. C 76 \(1997\) 67](#)
- Dijet cross section in NLO pQCD:

$$d\sigma_{ep \rightarrow e + \text{jets} + X} = \sum_{a,b} \int_0^1 dx_\gamma \int_0^1 dx_p f_{\gamma/e}(x_\gamma, \mu_{F\gamma}) f_{b/\gamma}(x_p, \mu_{F\gamma}) f_{a/p}(x_p, \mu_{Fp}) d\hat{\sigma}_{ab \rightarrow cd}(x_\gamma, x_p, \alpha_s(\mu_R), \mu_{F\gamma}, \mu_{Fp}, \mu_R)$$

Photon flux in Weizsäcker-Williams approx.

$$f_{\gamma/e}(y) = \frac{\alpha}{2\pi} \left[ \frac{1 + (1-y)^2}{y} \ln \frac{Q_{\max}^2(1-y)}{m_e^2 y^2} + 2m_e^2 y \left( \frac{1-y}{m_e^2 y^2} - \frac{1}{Q_{\max}^2} \right) \right]$$

Photon PDFs for resolved photon; in  $b=\gamma$  case,

$f_{b/\gamma} = \delta(1-x_\gamma)$  for Born and virtual corr.

Proton PDFs

2→2 and 2→3 hard parton scattering cross section

- This parton-level cross section assumes massless quarks and for comparison with data, needs to be supplemented with [hadronization corrections](#) from Monte Carlo (LO + parton showers), [Helenius, arXiv:1806.07246](#) and [arXiv:1811.10931](#); [Helenius and Rasmusen, EPJC 79 \(2019\) 413](#).

- Parton momentum fractions are determined using their [hadron-level estimates](#) based on measured jet transverse energies  $E_{T1,2}$  and rapidities  $\eta_{1,2}$ .

$$x_\gamma^{\text{obs}} = \frac{p_{T,1} e^{-\eta_1} + p_{T,2} e^{-\eta_2}}{2yE_e}$$

$$x_A^{\text{obs}} = \frac{p_{T,1} e^{\eta_1} + p_{T,2} e^{\eta_2}}{2E_A},$$

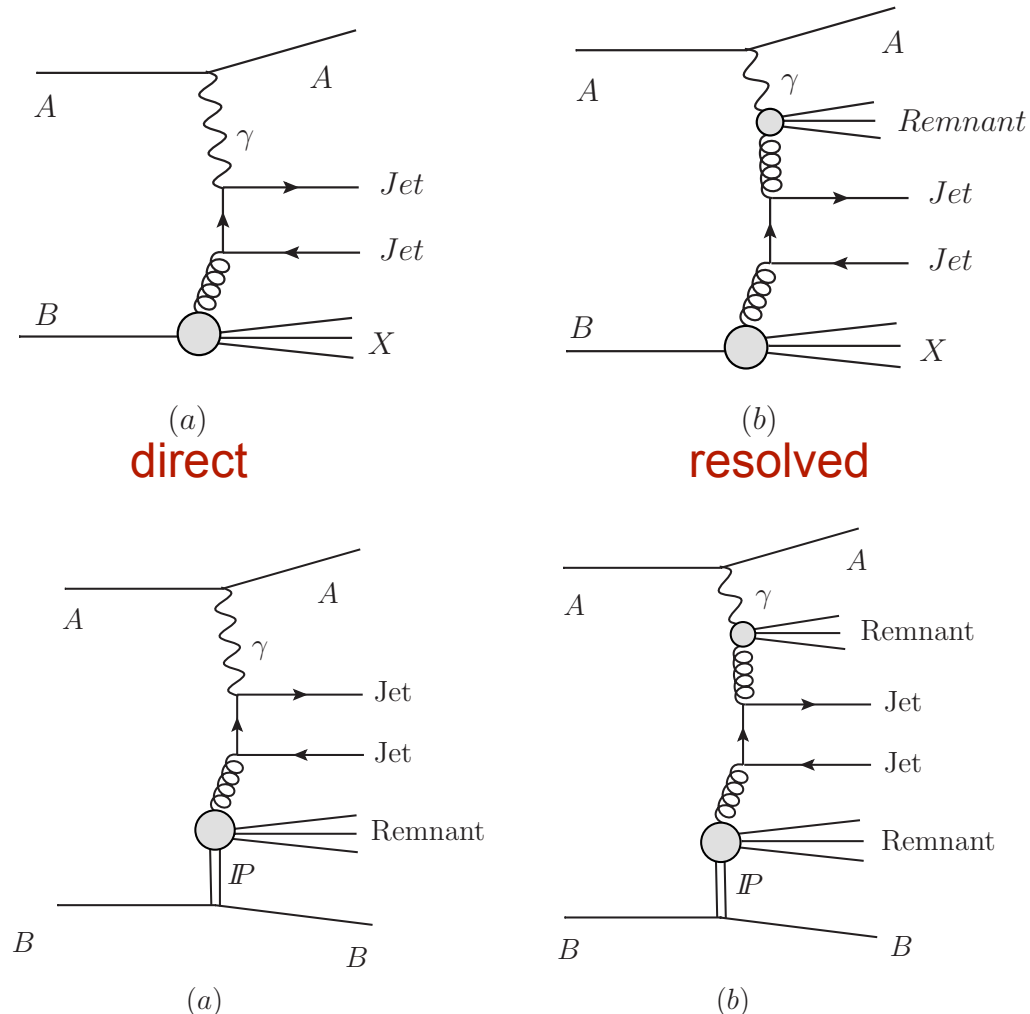
# Dijet photoproduction in UPCs@LHC

- The focus of UPC measurements@LHC has been exclusive (coherent) photoproduction of charmonia ( $J/\psi$ ,  $\psi'$ ) and light vector mesons ( $\rho$ )  $\rightarrow$  new constraints on the gluon density at small  $x$  down to  $x_p \sim 6 \times 10^{-6}$  and  $x_A \sim 10^{-5}$ .

- Nuclear and photon PDFs can also be studied in **inclusive dijet photoproduction** in Pb-Pb UPCs, [ATLAS-CONF-2017-011](#), [ATLAS-CONF-2022-021](#)

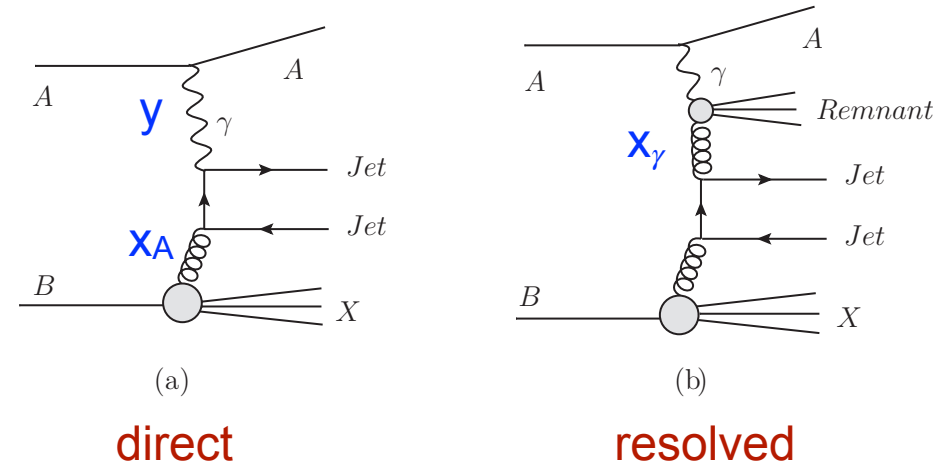
- Requiring intact nuclear target  $\rightarrow$  **diffractive dijet photoproduction** in Pb-Pb UPCs  $\rightarrow$  access to novel nuclear **diffractive PDFs** and mechanism of QCD factorization breaking, [Guzey, Klasen, JHEP 04 \(2016\) 158](#)

- First LO pQCD calculations of heavy-flavor jets in UPCs, [Strikman, Vogt, White, PRL 96 \(2006\) 082001](#)  $\rightarrow$  very large rates for inclusive and diffractive ( $\sim 20\%$ ) dijets.



# Inclusive dijet photoproduction in Pb-Pb UPCs@LHC

- Cross section of dijet photoproduction using collinear factorization and NLO pQCD, which is successful for HERA data on dijet photoproduction in ep scattering, Klasen, Kramer, Z.Phys. C 72 (1996) 107, Z. Phys. C 76 (1997) 67; Klasen, Rev. Mod. Phys. 74 (2002) 1221; Klasen, Kramer, EPJC 71 (2011) 1774



$$d\sigma(AA \rightarrow A + 2\text{jets} + X) =$$

$$\sum_{a,b} \int dy \int dx_\gamma \int dx_A f_{\gamma/A}(y) f_{a/\gamma}(x_\gamma, \mu^2) f_{b/A}(x_A, \mu^2) d\hat{\sigma}_{ab \rightarrow \text{jets}}$$

Photon flux from QED:

- high intensity  $\sim Z^2$
- high photon energy  $\sim \gamma_L$

Photon PDFs (GRVHO)  
(resolved photon), from  
e+e- data

Hard parton  
cross section

$$f_{\gamma/A}(y) = \frac{2\alpha_{\text{e.m.}} Z^2}{\pi} \frac{1}{y} \left[ \zeta K_0(\zeta) K_1(\zeta) - \frac{\zeta^2}{2} (K_1^2(\zeta) - K_0^2(\zeta)) \right]$$

Nuclear PDFs  
(nCTEQ15, EPPS16)

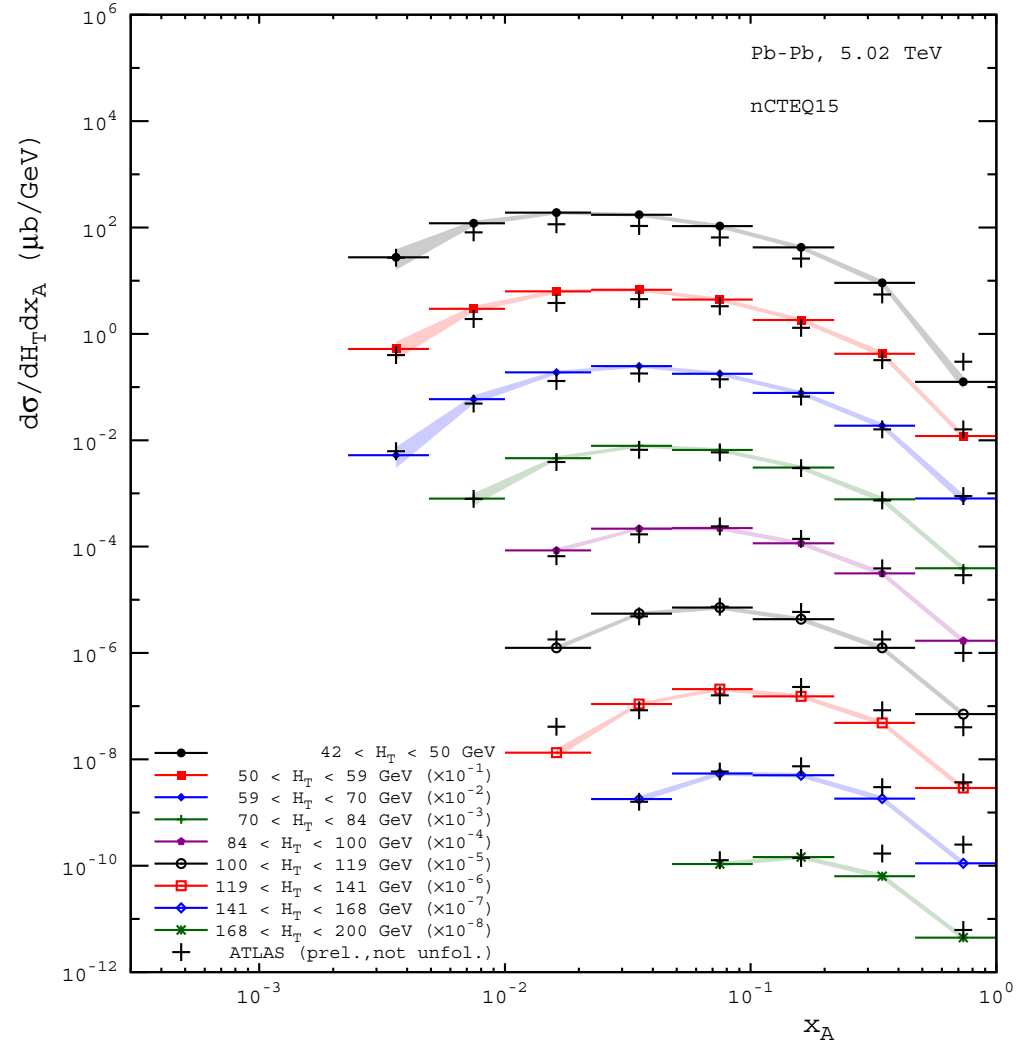
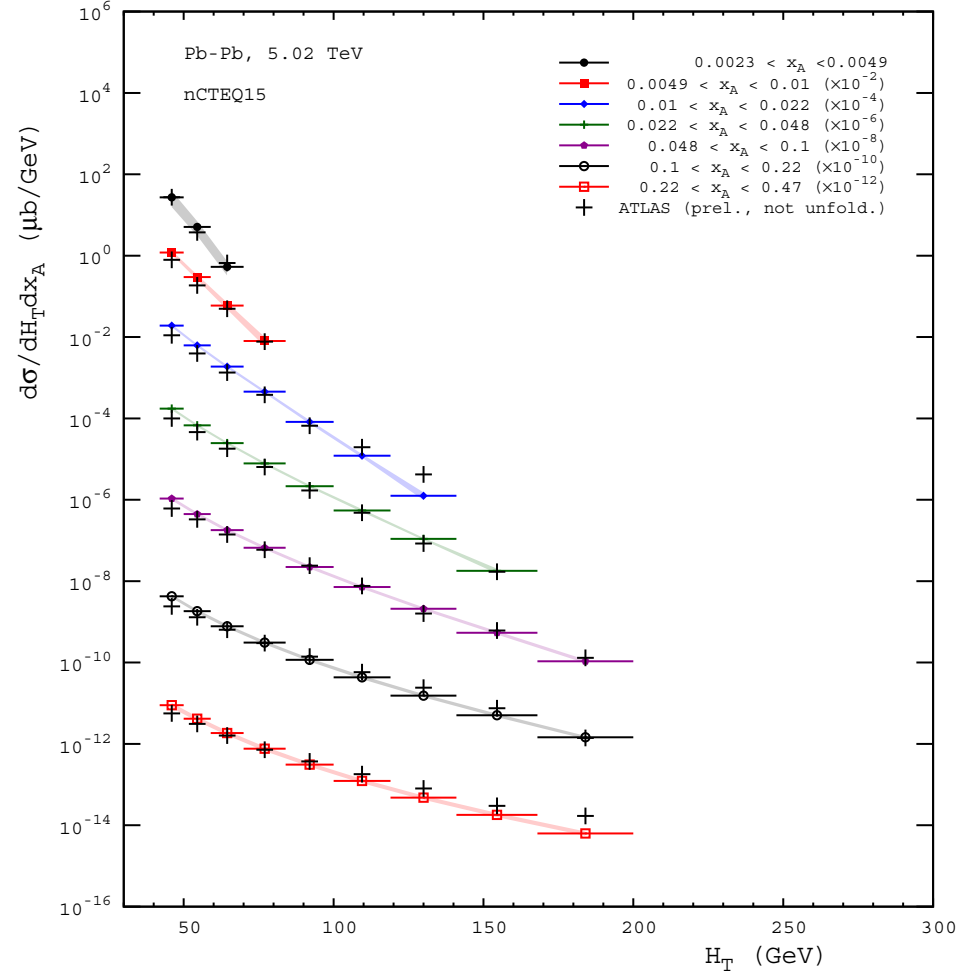
$$\zeta = ym_p b_{\text{min}} \approx ym_p (2R_A)$$

See talk P. Paakinen

# Inclusive dijet photoproduction in Pb-Pb UPCs@LHC (2)

- NLO pQCD vs. ATLAS data as a function of the dijet transverse momentum  $H_T = E_{T,jet1} + E_{T,jet2}$  and nuclear momentum fraction  $x_A = (m_{jets}/\sqrt{s_{NN}})e^{-y_{jets}}$

Guzey, Klasen, PRC 99 (2019) 065202

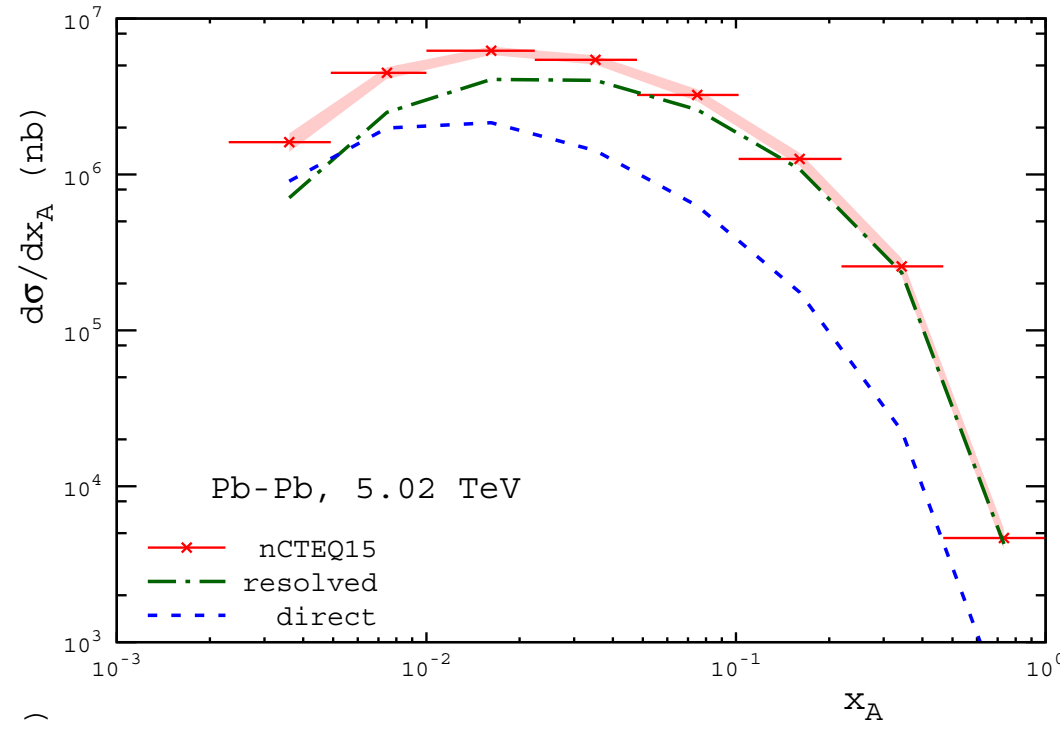


- Shape and normalization of ATLAS data, [ATLAS-CONF-2017-011](#) reproduced well (data not corrected for detector response).
- We haven't compared to the more recent unfolded data, [ATLAS-CONF-2022-021](#).

# Inclusive dijet photoproduction in Pb-Pb UPCs@LHC (3)

- Resolved vs. direct photon contributions:
  - resolved photons dominate for  $x_A > 0.01$
  - resolved and direct are compatible for  $x_A < 0.01$
  - similar trend in leading order (LO) analysis in PYTHIA8 framework, Helenius, Rasmusen, EPJ C 79 (2019) 5, 413.

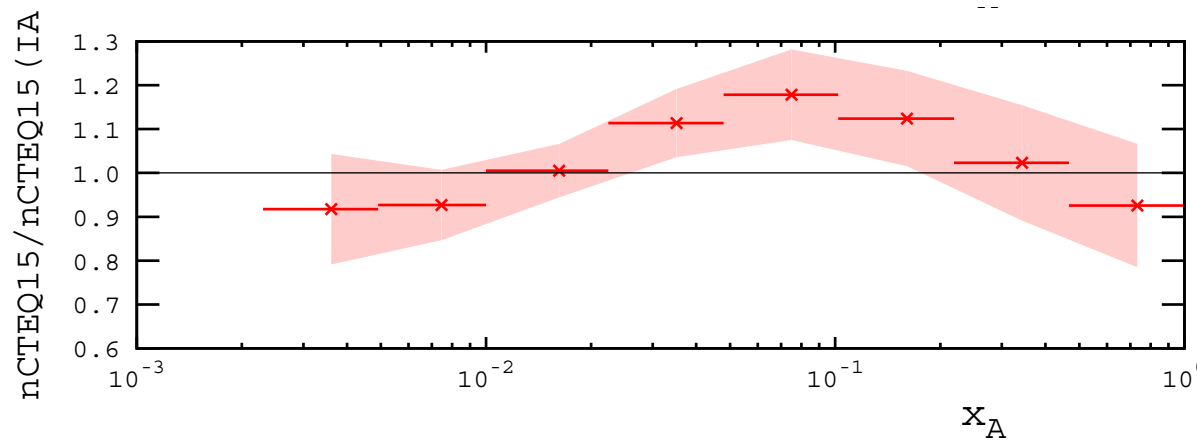
See talk I. Helenius



- Nuclear modifications: shape of repeats that of  $R_g(x) = g_A / Ag_N$ :
  - 10% shadowing for  $x_A < 0.01$
  - 20% antishadowing at  $x_A \sim 0.1$
  - 5-10% EMC effect for large  $x_A$  → can be compared to predictions for EIC, Klasen, Kovarik, PRD 97 (2018) 114013; Guzey, Klasen, PRC 102 (2020) 6, 065201.

$$R = \frac{d\sigma(AA \rightarrow A + 2\text{jets} + X)}{d\sigma^{\text{IA}}(AA \rightarrow A + 2\text{jets} + X)}$$

$$f_{b/A}^{\text{IA}} = Z f_{b/p} + (A - Z) f_{b/n}$$





# Constraints on nPDFs from dijet photoproduction

- Potential of inclusive dijet photoproduction in Pb-Pb UPCs to constrain nPDFs can be assessed using Bayesian reweighting used for pA data, [Armesto et al. JHEP 1311 \(2013\) 015](#); [Paukkunen, Zurita, JHEP 1412 \(2014\) 100](#); [Kusina et al., EPJC 77 \(2017\) 488](#).

- Using error nPDFs, one generates N (N=10,000) replicas:

$$f_{j/A}^k(x, Q^2) = f_{j/A}^0(x, Q^2) + \frac{1}{2} \sum_{i=1}^N \left[ f_{j/A}^{i+}(x, Q^2) - f_{j/A}^{i-}(x, Q^2) \right] R_{ki}$$

↓
↘
↙
↓

central value
error PDFs
random numbers

- Calculate the cross section for each replica:

$$\frac{d\sigma^k}{dx_A} = \sum_{a,b} \int_{y_{\min}}^{y_{\max}} dy \int_0^1 dx_\gamma f_{\gamma/A}(y) f_{a/\gamma}(x_\gamma, \mu^2) f_{b/B}(x_A, \mu^2) d\hat{\sigma}(ab \rightarrow \text{jets})$$

- $\chi^2$  for each k using our calculations with central nPDFs as pseudo-data (need to assign some error **eps=5-15%**)

- statistical weights  $w_k$  for replicas to reproduces the pseudo-data with tolerance T

$$w_k = \frac{e^{-\frac{1}{2}\chi_k^2/T}}{\frac{1}{N_{\text{rep}}} \sum_i^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_i^2/T}}$$

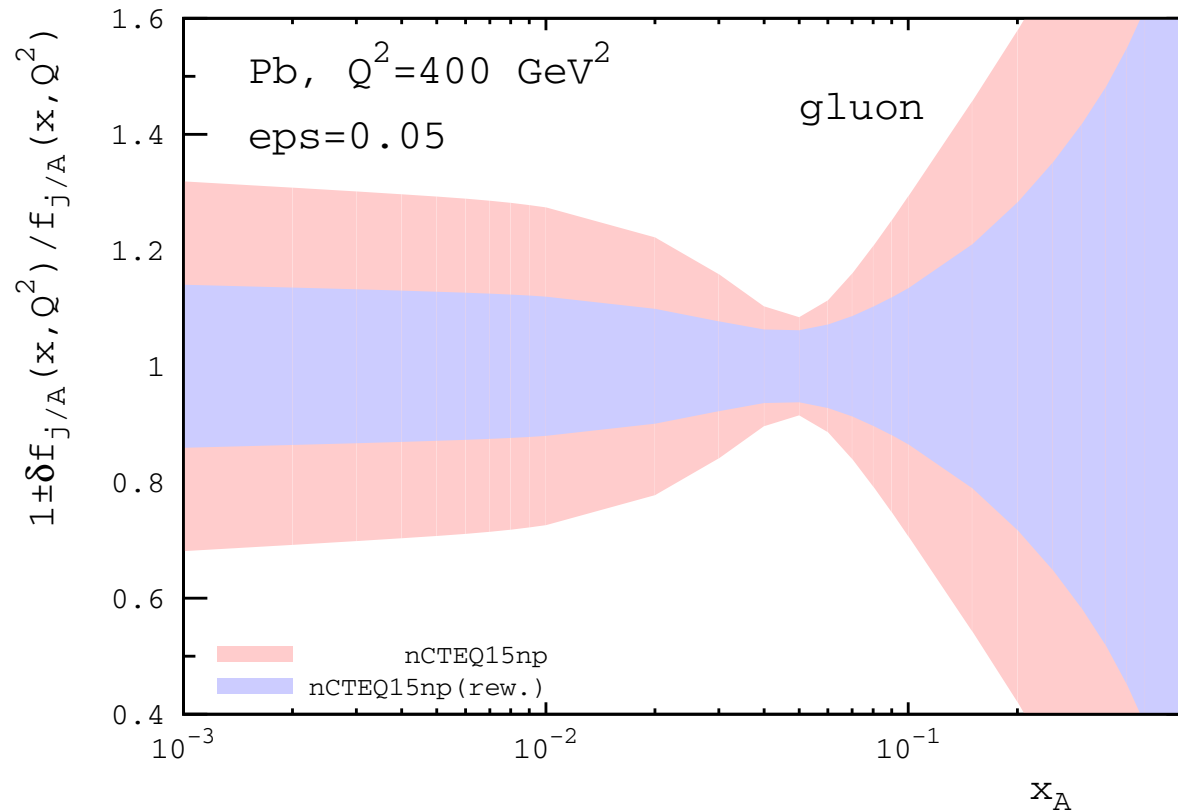
# Constraints on nPDFs from dijet photoproduction (2)

- Re-weighted PDFs and their uncertainties:

$$\langle f_{j/A}(x, Q^2) \rangle_{\text{new}} = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k f_{j/A}^k(x, Q^2),$$

$$\delta \langle f_{j/A}(x, Q^2) \rangle_{\text{new}} = \sqrt{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k \left( f_{j/A}^k - \langle f_{j/A}(x, Q^2) \rangle_{\text{new}} \right)^2}$$

- This quantifies the effect of the pseudo-data on nPDFs and their uncertainties.

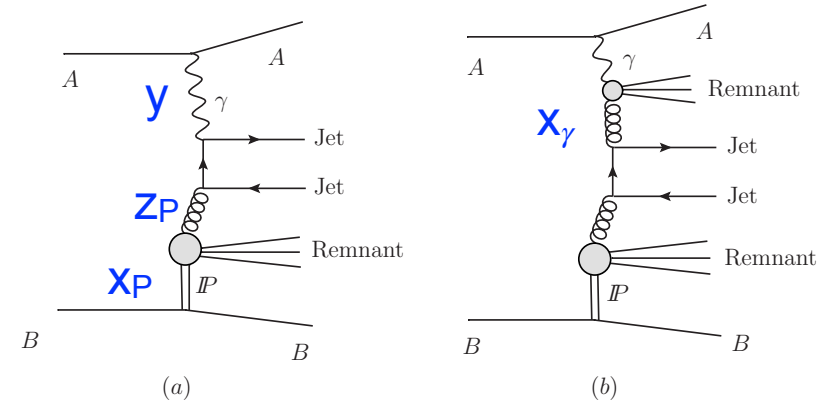


Guzey, Klasen, EPJ C 79 (2019) 5, 396

- Assuming 5% error → reduction of uncertainties by factor 2 at  $x_A = 0.001$ .

# Diffractive dijet photoproduction in Pb-Pb UPCs@LHC

- In framework of collinear factorization & NLO pQCD, it probes novel **nuclear diffractive PDFs**.
- Contribution of right-moving photon source:



$$d\sigma(AA \rightarrow A + 2\text{jets} + X + A)^{(+)} =$$

$$\sum_{a,b} \int dt \int dx_P \int dz_P \int dy \int dx_\gamma f_{\gamma/A}(y) f_{a/\gamma}(x_\gamma, \mu^2) f_{b/A}^{D(4)}(x_P, z_P, t, \mu^2) d\hat{\sigma}_{ab \rightarrow \text{jets}}$$



- Nuclear diffractive PDF  $f_{b/A}^{D(4)}$  = conditional probability to find parton **b** with mom. fraction  $z_P$  with respect to the diffractive exchange (**pomeron**) carrying mom. fraction  $x_P$  provided the nucleus remained intact with mom. transfer **t**.
- $f_{b/A}^{D(4)}$  is subject to nuclear modifications. The leading twist nuclear shadowing model predicts **strong nuclear suppression** (shadowing), Frankfurt, Guzey, Strikman, Phys. Rept. 512 (2012) 255

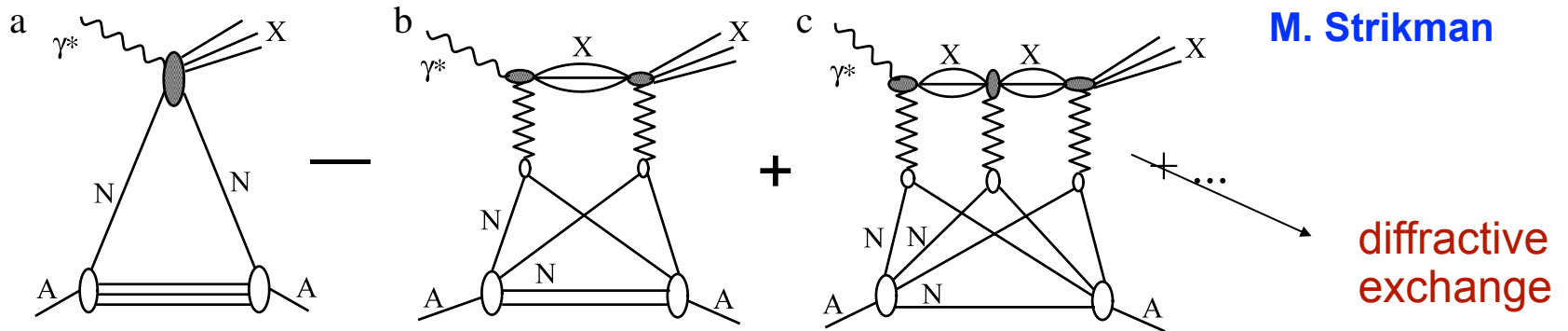
$$f_{b/A}^{D(4)}(x_P, z_P, t, \mu^2) = R_b(x_P, z_P, \mu^2) A^2 F_A^2(t) f_{b/p}^{D(4)}(x_P, z_P, t = 0, \mu^2)$$

$$\approx 0.15 A^2 F_A^2(t) f_{b/p}^{D(4)}(x_P, z_P, t = 0, \mu^2)$$

# Leading twist approach to nuclear shadowing

- Combination of Gribov-Glauber theory with QCD factorization theorems for inclusive and diffractive DIS  $\rightarrow$  prediction for small-x nPDFs at input scale  $Q_0$ , Frankfurt, Guzey, Strikman, Phys. Rept. 512 (2012) 255

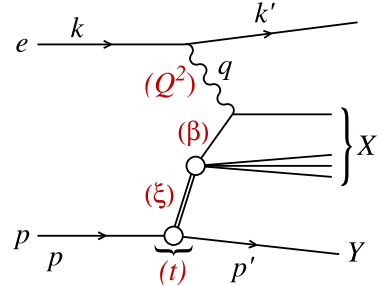
See talks B. Schenke, M. Strikman



$$\beta f_{j/A}^{D(3)}(\beta, Q^2, x_{\mathbb{P}}) = 4\pi A^2 \beta f_{j/N}^{D(4)}(\beta, Q^2, x_{\mathbb{P}}, t_{\min}) \int d^2b \left| \int_{-\infty}^{\infty} dz e^{ix_{\mathbb{P}} m_N z} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q^2)} \int_z^{\infty} dz' \rho_A(b, z') \rho_A(b, z) \right|^2$$

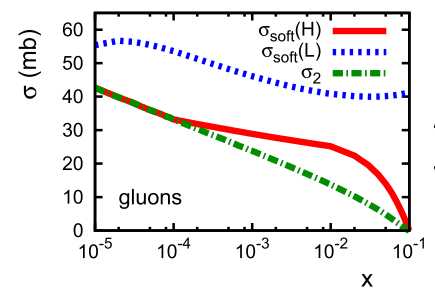
$j=q$  or  $g$ 
 $\downarrow$ 
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N=2 term: proton diffractive PDFs from QCD analysis of HERA data



$N \geq 3$  terms: model-dependent effective cross section

Nuclear density

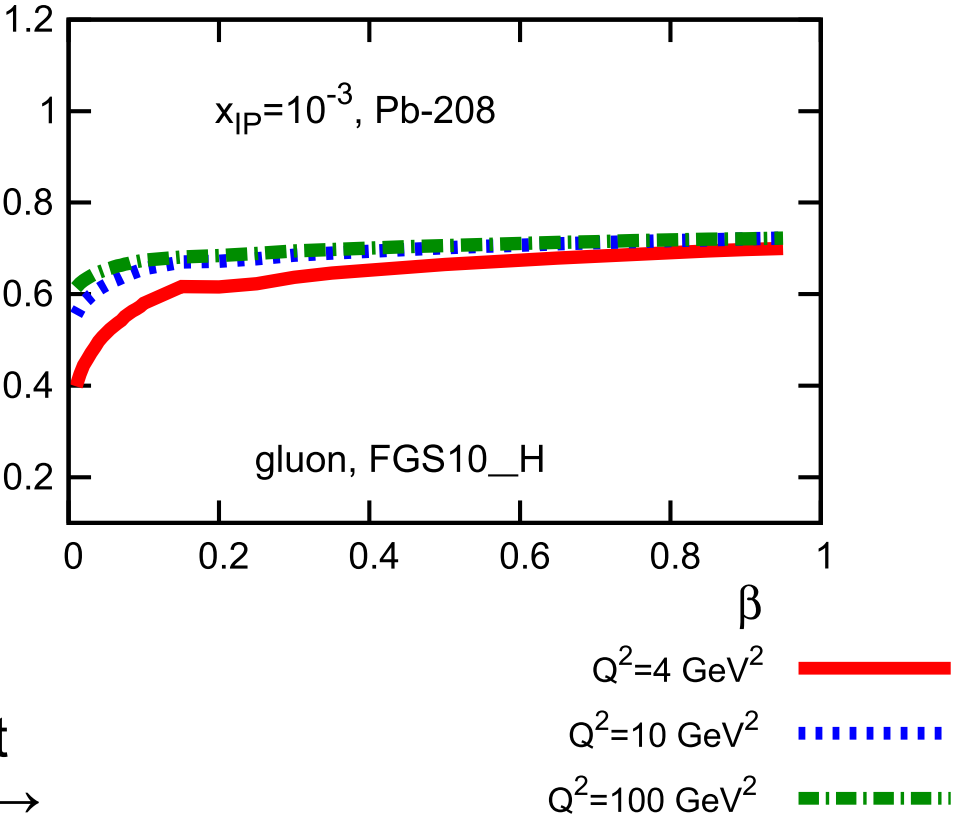
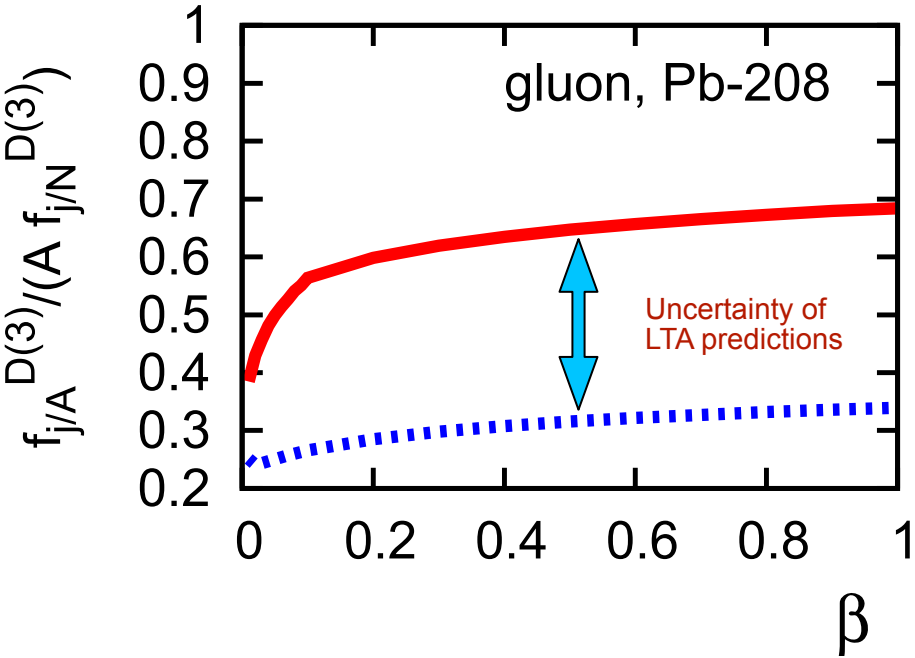


Spread in  $\sigma_{\text{soft}} \rightarrow$  uncertainty of LTA predictions

- Main feature: connects shadowing in eA with diffraction in ep.
- Can also be generalized to incoherent (quasi-elastic diffraction).

# Leading twist approach to nuclear shadowing (2)

- HERA analysis: perturbative Pomeron is made mostly of gluons → LTA model naturally predicts large gluon nuclear shadowing.
- Nucleus/proton ratio of diffractive PDFs is a weak function of  $\beta$ ,  $x_P$  and  $Q^2$ .



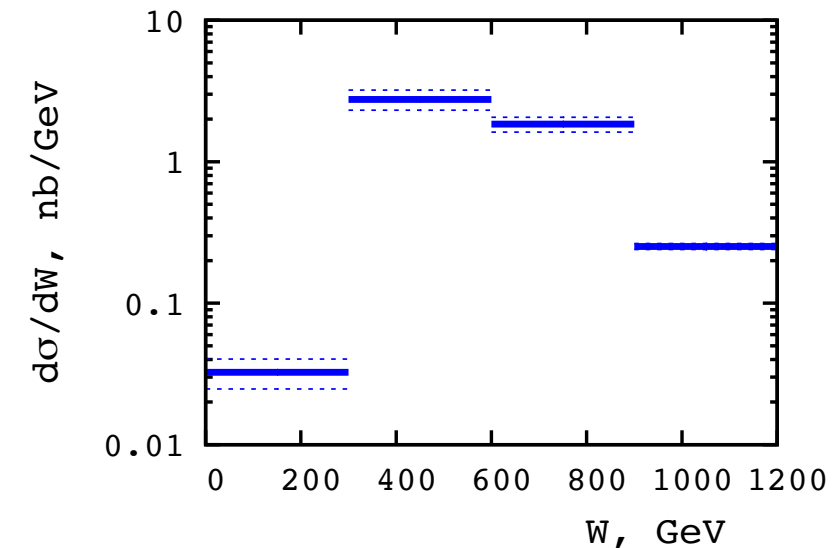
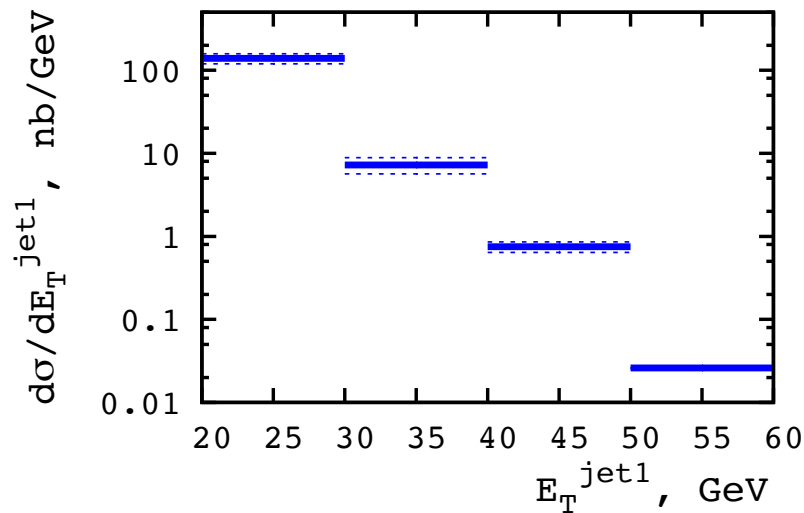
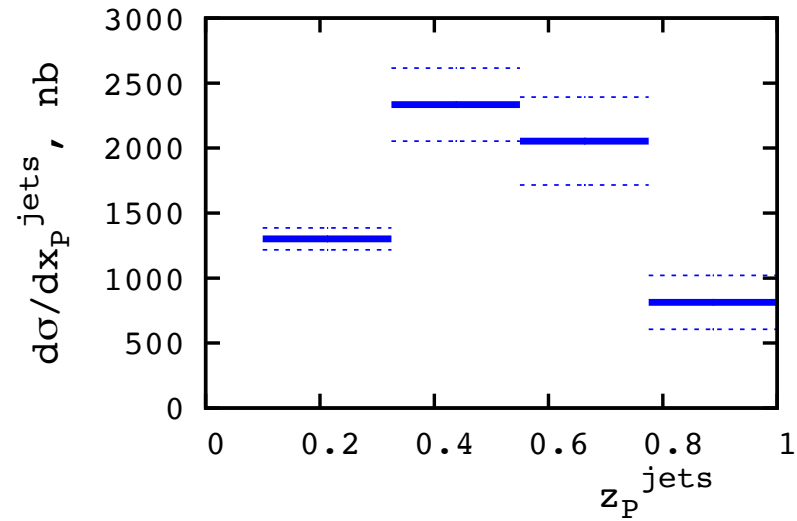
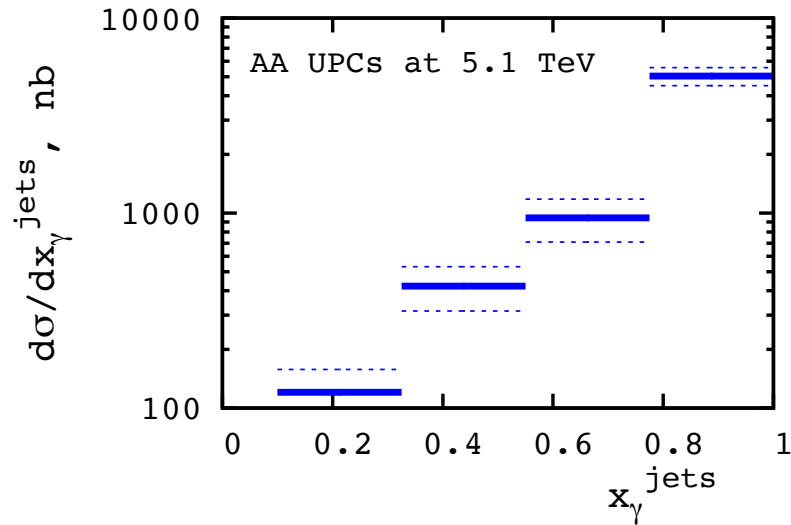
• → we approximate it by a single constant  $R_b=0.15-0.30 < (R_g)^2=(0.6-0.7)^2=0.4-0.5$  → connection to strong shadowing in  $J/\psi$  case,

See talk E. Iancu

# Diffractive dijet photoproduction in Pb-Pb UPCs@LHC (2)

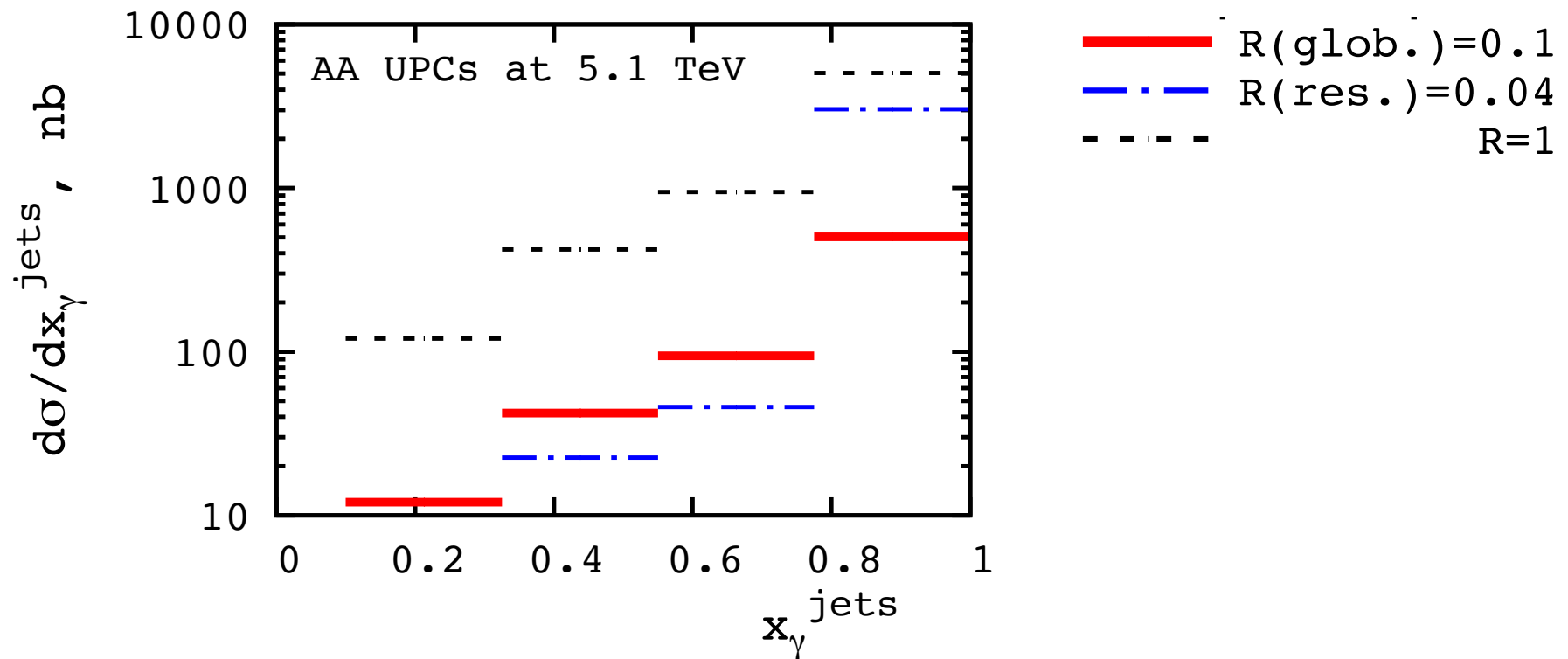
- NLO pQCD predictions as a function of momentum fractions  $x_\gamma$  and  $z_P$ , leading jet transverse momentum  $E_T^{\text{jet1}}$ , and photon-nucleus energy  $W$ .

Guzey, Klasen, JHEP 04 (2016) 158



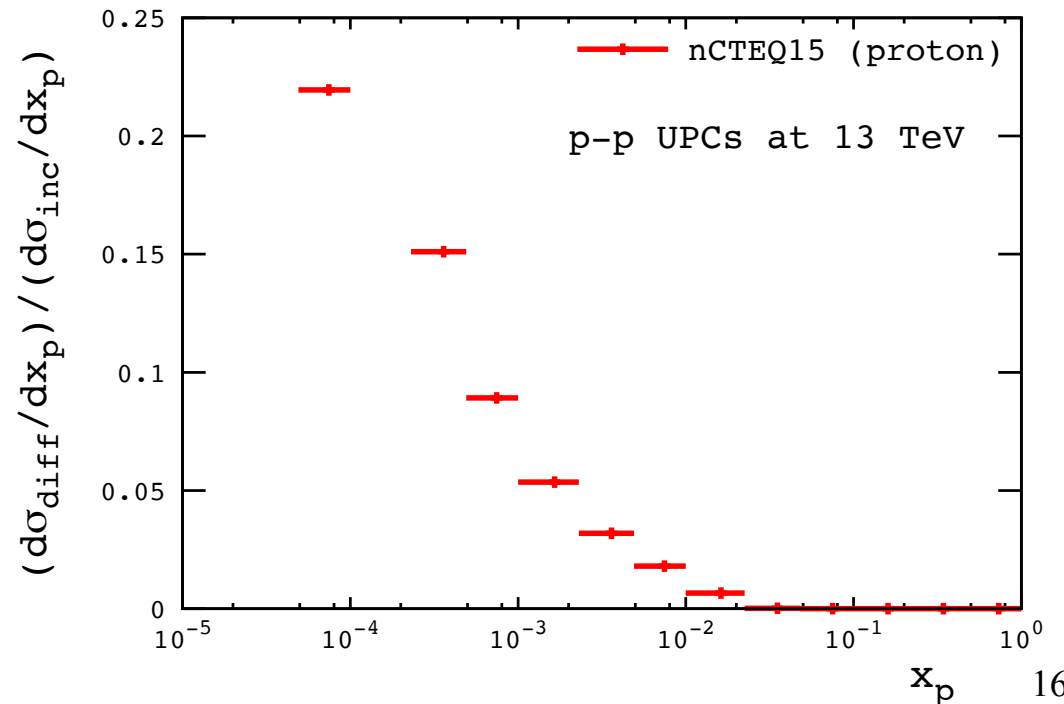
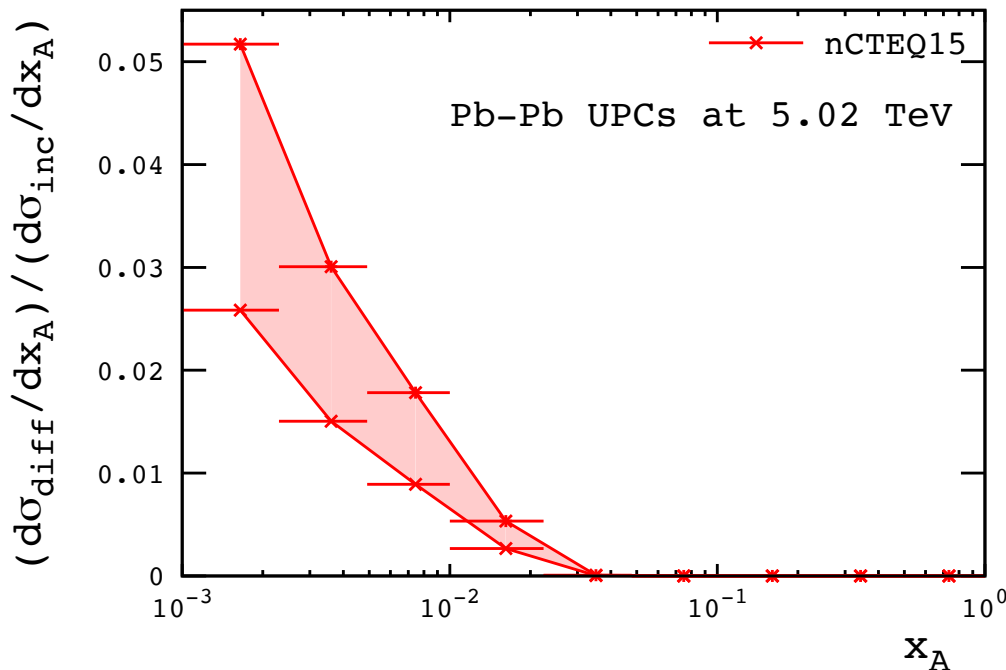
# Diffraction dijet photoproduction in Pb-Pb UPCs@LHC (3)

- Analyses of diffractive dijet photoproduction in ep scattering@HERA  $\rightarrow$  QCD factorization is broken, i.e., NLO calculations overestimate data by factor of  $\sim 2$ , [Klasen, Kramer, EPJ C 38 \(2004\) 93](#); [PRL 93 \(2004\) 232002](#); [JPhys.G 31 \(2005\) 1391](#); [MPLA 23 \(2008\) 1885](#); [EPJ C 70 \(2010\) 91](#); [PLB 508 \(2001\) 259](#); [EPJ C 49 \(2007\) 957](#); [PRD 80 \(2009\) 074006](#); [Guzey, Klasen, EPJ C 76 \(2016\) 8, 467](#)
- The pattern of unknown: either the global suppression factor  $R(\text{glob.})=0.5$  or the resolved-only suppression  $R(\text{res.})=0.34$ , [Kaidalov, Khoze, Martin, Ryskin, EPJ C 66 \(2010\) 373](#), or the flavor-dependent combination of the two, [Guzey, Klasen, EPJ C 76 \(2016\) 8, 467](#)
- One can differentiate between these scenarios by studying  $x_\gamma$  distribution in AA UPCs, [Guzey, Klasen, JHEP 04 \(2016\) 158](#)



# How large is the diffractive contribution?

- Diffractive contribution to inclusive dijet photoproduction in Pb-Pb UPCs in ATLAS kinematics does not exceed **5%** at small  $x_A$ , [Guzey, Klasen, PRD 104\(2021\) 11, 114013](#)
- This is the effect of restricted kinematics,  $p_{T1} > 20 \text{ GeV}$ ,  $x_A > 0.001$  and large shadowing suppression of nuclear diffractive PDFs.
- $\rightarrow$  the diffractive contribution and related ambiguity with the determination of the photon-emitting nucleus (invariant energy  $W$ ) can be safely neglected.
- It is not the case for pp UPCs, where the diffractive contribution can reach **20-25%** at  $x_p \sim 10^{-4}$ .





# Summary and Outlook

- Photoproduction of jets is a standard tool of QCD. Its theory is well-established in NLO pQCD and compares very well to HERA data.
- **Nuclear PDFs** are poorly constrained by available fixed-target and pA LHC data and, hence, there is growing interest in obtaining new constraints on them using hard **photon-nucleus scattering** in heavy ion UPCs at the LHC.
- **Inclusive dijet photoproduction** in Pb-Pb UPCs@LHC probes nPDFs down to  $x_A \sim 0.005$  and can reduce the current small- $x_A$  uncertainties of the gluon distribution by factor of  $\sim 2$ .
- **Diffraction dijet photoproduction** in Pb-Pb UPCs@LHC accesses novel **nuclear diffractive PDFs** and may shed new light on mechanism of QCD factorization breaking in this process.
- Extraction of diffractive events is problematic due to ambiguity in photon source  $\rightarrow$  can be solved by requiring  $M_X \gg M_{\text{dijet}}$ , [Strikman, Vogt, White, PRL 96 \(2006\) 082001](#).
- Both processes can be viewed as precursors of analogous measurements at EIC, [Guzey, Klasen, PRC 102 \(2020\) 6, 065201](#); [JHEP 05 \(2020\) 074](#).