



# Multi-boson production sensitivity to dimension-6 EFT operators at the LHC

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31<sup>st</sup> August

Multi-Boson Interaction 2023, San Diego US

#### **Outline**

- Theoretical introduction: SMEFT
- Motivation for multi-boson measurements
- Experimental results at LHC
- Sensitivity LHE study to constrain dimension-6 SMEFT operators
  - Processes: VBS+WW di-boson and tri-boson
  - Analysis Strategy
  - Results: One and Two-dimensional operator constraints
- Summary and Outlook

#### References



• R.Bellan et al. "A sensitivity study of **VBS** and **diboson WW** to dimension-6 EFT operators at the LHC" -published in <u>JHEP05(2022)039</u>



• R.Bellan et al. "A sensitivity study of **triboson** production processes to dimension-6 EFT operators at the LHC" -published in JHEP08(2023)158



# **Theory introduction: SMEFT approach**

- Multi-boson processes serve as a test of the EW Symmetry Breaking
- **EW sector still unexplored** since **several rare processes** not yet **observed**!
  - Any deviation in kinematic observables could point to New Physics



- SMEFT bottom-up approach:
  - Effective Lagrangian with only light SM particles
  - **BSM effects** incorporated as a **momentum expansion**

$$\mathcal{L} = \mathcal{L}_{ ext{SM}}^{(4)} + \sum_{n,i} rac{1}{\Lambda^{n-4}} c_i^{(n)} Q_i^{(n)}$$
 SU(3)<sub>C</sub>×SU(2)<sub>L</sub>×U(1)<sub>Y</sub> invariant Dim. *n* operator.

CKen Mimasu

# **Motivation for multi-boson measurements**

- These processes probe the non-Abelian nature of SM:
  - direct access to triple/quartic gauge couplings (TGC/QGC)
  - sensitive to couplings between Higgs and gauge bosons
    - complementary to **Higgs** measurements at scales  $> m_{H}$

$$= \frac{1}{2} \left( \frac{1}{2} + \frac{1}{2} +$$

- Portal to BSM physics:
  - **model-independent** via Effective Field Theories (EFTs)
    - 18 bosonic operators in dim-8 EFT tested (aQGCs)
    - 14 fermionic and bosonic dim-6 EFT operators

 $q_{f1}$ 

 $q_{f2}$ 

VBS

Triboson

 $q_{i1}$ 

 $q_{i2}$ 

#### Multi-boson measurements at LHC



- The ATLAS and CMS collaborations have a rich program of multi-boson analyses
- All diboson and VBS processes have been measured in different decay modes
- Many triboson processes have been studied with at least a photon in final state



# **Vector boson scattering measurements**

- Statistically limited, L~137 fb<sup>-1</sup> allows new measurements
- Good agreement with SM: in some VBS VV scattering the EW measurements are ~1σ away from theory
- Accurate modelling of VVjj
   non-VBS contributions crucial
- BSM effects in aQGC (EFT dim-8) dim-6 important and should not be neglected [arXiv:1809.04189]





# **Di-boson measurements**

- Diboson: well-known processes,  ${\color{black}\bullet}$ differential x-sec measurements
- Large cross-section: limited by systematic uncertainties
- In agreement with state-of-the-art **NNLO** theoretical calculations
- BSM effects manifest as **aTGC** (EFT dim-6)





#### **Tri-boson measurements**

- Extremely rare processes:
   σ × BR(to leptons) ~O(1fb)
- Observed three massive gauge boson production and in channel with a single VVy and two photons Vyy
- BSM effects as both aTGC/aQGC: and as anomalous Higgs-gauge coupling



# **Constraints to SMEFT operators**



- Multi-boson processes can be a tool to probe BSM physics at the TeV scale, provided no new light state exists
- Combination of several analyses is key to fit simultaneously all 59 independent
   SMEFT operators to preserve gauge invariance and UV-matching
  - What is the sensitivity reach of a **VBS** and **tri-boson** combination?
  - What is the interplay with **WW** constraints?
- The first step towards **global SMEFT fit** is the **sensitivity study** at parton level of anomalies induced by **dimension-6** EFT operators
- Assess sensitivity interplay between multi-boson analyses at LHC
  - in the future **global EFT fit** will be necessary to provide the most stringent constraints to **SMEFT operators** (top, Higgs, EW, etc)

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#### Motivation for dimension-6 EFT sensitivity study

- Interpretation of VBS/tri-boson results traditionally in terms of dim-8 SMEFT
   operators (aQGCs) at the LHC
- However, dim-6 EFT operators
   can also have an impact in VBS, e.g
   *pp→ZZjj*, [arXiv:1809.04189]
- First LHE sensitivity study of VBS+WW including O(Λ<sup>-4</sup>) dim-6 EFT terms extended to triboson production



• EFT analysis of both **EWK+QCD-induced** processes (**main** background)

## **SMEFT Monte Carlo Generations**



- **Paramerisation** using **15 dim-6 SMEFT operators** from **Warsaw basis** 
  - in triboson study focus on bosonic ops, for Yukawa sector see <u>JHEP04(2021)023</u> Ο
- Generated at LO with <u>SMEFTsim</u> interfaced with MadGraph5\_aMC@NLO (2.6.5)
  - U(3)<sup>5</sup> flavour symmetry
  - $\{m_w, m_7, G_F\}$  input scheme
  - **CP-even**  $\bigcirc$
  - $\Lambda = 1 \text{TeV}$  $\bigcirc$
- **Insertion** of **one operator** per **dia** to generate directly **single compo**

$N \propto  \mathcal{A}_{SM} ^2 + \sum_{\alpha} \frac{c_{\alpha}}{\Lambda^2} \cdot \underbrace{2 \operatorname{Re}(\mathcal{A}_{SM} \mathcal{A}_{Q_{\alpha}}^{\dagger})}_{\text{Lin}}$	$+\frac{c_{\alpha}}{\Lambda^{4}}\cdot\left \mathcal{A}_{Q_{\alpha}}\right ^{2}+\sum_{\alpha,\beta}\frac{c}{\alpha}$	$\frac{\Lambda^{\alpha} \mathbf{C}_{\beta}}{\Lambda^{4}} \cdot \underbrace{\operatorname{Re}(\mathcal{A}_{Q_{\alpha}} \mathcal{A}_{Q_{\beta}}^{\dagger})}_{\mathbf{Mix}} \mathbb{I}^{12}$
SM	Quad	$Q_{ll}^{(1)} = (\bar{l}_p \gamma_\mu l_r) (\bar{l}_r \gamma^\mu l_p)$
generate directly <b>single components</b>	$Q_W = \varepsilon^{ijk} W^{i\nu}_{\mu} W^{j\rho}_{\nu} W^{k\mu}_{\rho}$	$Q_{HB} = (\phi^{\dagger}\phi)B_{\mu\nu}B^{\mu\nu}$
nsertion of one operator per diagram	$Q_{HWB} = (H^{\dagger}\sigma^{i}H)W^{i}_{\mu u}B^{\mu u}$	$Q_{HW} = (H^{\dagger}H)W^{i}_{\mu\nu}W^{i\mu\nu}$
	$Q_{HD} = (H^{\dagger}D_{\mu}H)(H^{\dagger}D^{\mu}H)$	$Q_{H\square} = (H^{\dagger}H)\Box(H^{\dagger}H)$
$\Lambda = 1 \text{TeV}$	$Q_{qq}^{(3)} = (\bar{q}_p \gamma_\mu \sigma^i q_p) (\bar{q}_r \gamma^\mu \sigma^i q_r)$	$Q_{qq}^{(3,1)} = (\bar{q}_p \gamma_\mu \sigma^i q_r) (\bar{q}_r \gamma^\mu \sigma^i q_p)$
CP-even	$Q_{qq}^{(1)} = (ar{q}_p \gamma_\mu q_p)(ar{q}_r \gamma^\mu q_r)$	$Q_{qq}^{(1,1)} = (\bar{q}_p \gamma_\mu q_r)(\bar{q}_r \gamma^\mu q_p)$
<pre>{m<sub>w</sub>,m<sub>z</sub>,G<sub>F</sub>} input scheme</pre>	$Q_{Hq}^{(1)} = (H^{\dagger}i\overleftrightarrow{D_{\mu}}H)(\bar{q}_p\gamma^{\mu}q_p)$	$Q_{Hq}^{(3)} = (H^{\dagger} i \overleftrightarrow{D_{\mu}^{i}} H) (\bar{q}_{p} \sigma^{i} \gamma^{\mu} q_{p})$
U(3) <sup>5</sup> flavour symmetry	$Q_{Hl}^{(1)} = (H^{\dagger}i\overleftrightarrow{D_{\mu}}H)(\bar{l}_{p}\gamma^{\mu}l_{p})$	$Q_{Hl}^{(3)} = (H^{\dagger}i\overleftrightarrow{D_{\mu}^{i}}H)(\bar{l}_{p}\sigma^{i}\gamma^{\mu}l_{p})$

# **Processes of interest: VBS+WW**

**Modelling** of  $2 \rightarrow 6(4)$  processes including non-resonant diagrams  $\circ$  both EWK and QCD-induced contributions for SM and EFT processes

```
• Same-sign W<sup>±</sup>W<sup>±</sup>: pp > e^{\pm}v_e^{\mu^{\pm}}v_{\mu}^{\mu}jj
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• Opposite-sign W<sup>±</sup>W (QCD): pp > e^{\pm}v_e^{\mu}v_u^{\mu}jj
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• W<sup>±</sup>Z+2j(QCD): pp > e^+e^-\mu^{\pm}v_{\mu}^{}jj
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• ZZ+2j(QCD): pp > e<sup>+</sup>e<sup>-</sup>µ<sup>+</sup>µ<sup>-</sup>j j
```

```
• ZV+2j(QCD): pp > ZW<sup>+</sup> (W<sup>-</sup>Z) > I<sup>+</sup>I<sup>-</sup> j j j j
```

• WW: **pp** >  $e^+v_e \mu^-v_\mu$ 



• Integrated luminosity of 100 fb<sup>-1</sup> is assumed for the VBS+WW study

#### **Processes of interest: tri-boson**



- Modelling of  $2 \rightarrow 6(4+\gamma)$  processes including non-resonant diagrams
  - both EWK and QCD-induced contributions for SM and EFT processes



• Integrated luminosity of **300 fb<sup>-1</sup>** is assumed for the triboson study

# **Processes of interest - EFT sensitivity**

- **Summary** of the sensitivity of **each process** to the subset of operators
  - empty cells -> no EFT diagrams

#### **VBS+ WW diboson**

proc / op	Q <sub>HD</sub>	$Q_{H\square}$	Q <sub>HWB</sub>	$Q_{Hq}^{(1)}$	$Q_{Hq}^{(3)}$	Q <sub>HW</sub>	Q <sub>W</sub>	$Q_{Hl}^{(1)}$	Q <sub>Hl</sub> <sup>(3)</sup>	$Q_{ll}^{(1)}$	$Q_{qq}^{(3)}$	$Q_{qq}^{(3,1)}$	$Q_{qq}^{(1,1)}$	$Q_{qq}^{(1)}$	Q <sub>ll</sub>
SSWW-EW	1	1	1	1	1	1	1	(🗸)	1	1	1	1	1	1	(⁄)
OSWW-EW	1	1	1	1	1	1	1	(🗸)	1	1	1	1	1	1	(⁄)
WZ-EW	1	1	1	1	1	1	1	1	1	1	1	1	1	1	(⁄)
ZZ-EW	1	1	1	1	1	1	1	1	1	1	1	1	1	1	(⁄)
ZV-EW	1	1	1	1	1	1	1	1	1	1	1	1	1	1	
ww	1		1	1	1		1	(🗸)	1	1					
ZV-QCD	1		1	1	1		1	1	1	1					
OSWW-QCD	1		1	1	1		1	1	1	1					
WZ-QCD	1		1	1	1		1	1	1	1					(⁄)
ZZ-QCD	1		1	1	1			1	1	1					(⁄)

#### Triboson

$\begin{array}{c} \text{Operators} \rightarrow \\ \bigtriangledown \text{ Processes} \end{array}$	$Q_W$	$Q_{HB}$	$Q_{HW}$	$Q_{HWB}$	$Q_{HD}$	$Q_{H\square}$
$\mathbf{W}\mathbf{Z}\gamma$	1	1	1	1	1	
$\mathbf{Z}\mathbf{Z}\gamma$		1	1	1	1	
$\mathbf{V}\mathbf{Z}\gamma$	1	1	1	1	1	
${f QCD} extsf{-}{f Z}\gamma{f j}{f j}$				~	1	
VZZ	1	(•	1	1	1	(•
QCD-ZZjj				1	1	

• **Brackets** indicate only ops that enter **non-resonant** contributions or negligible

### **Template fit analysis**



- **Dependence** of **EFT-**induced kinematic **anomalies** on **Wilson coefficients**
- Likelihood fit for each variable based on 1σ range (area for 2D fit)



• Optimal variable extracted per operator used in combination

# **Tri-boson fully leptonic: WZγ**



- WZy production extremely rare process ~O(fb) studied in fully leptonic channel
- In the SM, **WZy** depends on **quartic gauge couplings (QGC)** at tree level unlike **ZZy**



 Fox-Wolfram variable provides good discrimination between O(Λ<sup>-4</sup>) dim-6 contribution and SM

$$H_0^T = \sum_{a 
eq b}^{ ext{fin.state}} rac{p_T^a p_T^b}{\left(\sum_i p_T^i
ight)^2}$$



# Tri-boson fully leptonic: ΖΖγ

- Higgs-gauge boson couplings affect the electroweak vertices in ZZy
- **Best variable:** total **p**<sub>T</sub> of same-sign **e**μ related to leptonic final state



# Triboson semi-leptonic: VZZ



- **Study** of inclusive **4ljj** production: **EWK VZZ+QCD-induced** background
- Unique Q<sub>w</sub>-sensitivity of WZZ channel to WWZZ quartic couplings unlike ZZZ channel

Inclusion of the main **background** given by diagrams involving QCD-induced vertices.



$$\sigma_{
m SM} = \sigma_{
m EWK} + \sigma_{
m QCD}$$



# **Triboson semi-leptonic: VZZ/VZγ**



- Impact of QCD background on semi-leptonic channels VZZ/VZy sensitivity
- **Q<sub>HWB</sub>** and **Q<sub>HD</sub>** induce EWK anomalies in QCD diagrams unlike other operators



Included as BSM signal for these operators

• enhances overall EFT-sensitivity despite larger background from QCD

#### **Fit procedure**



• Likelihood scan for each variable varying the Wilson coefficient in a fixed range



- SM expectation for N(c=0)
- Total yield:  $N(\mathbf{c}) = SM + \sum_{\mathbf{c}_{\alpha}} \mathbf{c}_{\alpha} \cdot Lin_{\alpha} + \mathbf{c}_{\alpha}^{2} \cdot Quad_{\alpha} + \sum_{\alpha\beta} \mathbf{c}_{\alpha}\mathbf{c}_{\beta}Mix_{\alpha\beta}$
- Single nuisance: proxy LHC luminosity 2% correlated across all yields & samples (flat prior)
- Sensitivity constraint at 68(95%) CL estimated as -2ΔLL <1(2.30) and -2ΔLL</li>
   <3.84(5.99) for single (pair) Wilson coefficient</li>



• Individual operator constraints with(without) O(Λ<sup>-4</sup>) quadratic terms



- Large impact of quadratic terms on half of the single operator constraints
- Strongest constraints on four-fermion coefficients as expected
- $Q_{HW}$ ,  $Q_{H_{\Box}}$ , ops constrained solely by VBS processes

# Individual constraints - Triboson



• Individual operator constraints with(without) O(Λ<sup>-4</sup>) quadratic terms



- Large impact of quadratic terms on Q<sub>w</sub>, Q<sub>HW</sub>, Q<sub>HB</sub> ops
- For other **operators**, linear interference term dominates

CHD

### **Profiled constraints**



- Performed **global fit** ensuring SMEFT model independence including all **O(Λ<sup>-4</sup>)** terms
  - **single operator** fit with all other **coefficients profiled** (free-floating in fit)



#### VBS+ WW diboson



#### triboson

- Profiled constraints are up to 10x less stringent wrt individual ones
- **Q**<sub>w</sub>-induced anomalies uncorrelated with other operators

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#### 2D constraints - VBS + WW combination

- Impact of **sensitivity** to **operator pairs** using **2D template** fit
  - contours allow assessing interplay of
     VBS and di-boson measurements
- Orthogonal constraints between WW and and VBS for (Q<sub>w</sub>,Q<sub>нwB</sub>)
- 2. **4-quark ops** constrained only by VBS
- 3. flat directions  $Q_{HI}^{(3)} \sim Q_{II}^{(1)}$  resolved thanks to channel combination





#### **2D constraints - Triboson combination**

- In triboson studies, most sensitive process
   WZy fully-leptonic to Q<sub>w</sub>-induced anomalies
- Semi-leptonic VZy leads to strongest constraints for Q<sub>HWB</sub> and Q<sub>HD</sub> operators
- Small mutual interference term for operator pairs Q<sub>w</sub>-Q<sub>нwb</sub>



CHWB



## Summary & Outlook

- Multi-boson measurements powerful tool to explore BSM physics in "UV-agnostic" way : consistency test of EWK sector
- First phenomenological dim-6 study on tri-boson VZZ/VZγ processes at O(Λ<sup>-4</sup>)
  - competitive & complementary constraints w.r.t. combination VBS+di-boson WW
- Expand scope of EFT analysis
  - combination of multi-boson analyses

     (VBS, di-boson, tri-boson) with Higgs
     measurements at reco level to constrain
     both dimension 6 and 8 EFT operators





# **Comparison VBS vs tri-boson combination**

<b>↓</b> Processes	Operators $\rightarrow$	$\mathbf{Q}_W$	$\mathbf{Q}_{HB}$	$\mathbf{Q}_{HW}$	$\mathbf{Q}_{HWB}$	$Q_{HD}$
	Best var.	$p_T^{l1}$	$p_{T(Z)}^{\gamma}$	$p_{T(Z)}^{\gamma}$	$p_{T(Z)}^{\gamma}$	$p_{T(WZ)}^{l1}$
$\mathrm{WZ}\gamma$	68% C.L.	[-0.20,0.21]	[-0.41,0.41]	[-0.44,0.44]	[-0.50, 0.73]	[-1.36,1.79]
	95% C.L.	[-0.31, 0.32]	[-0.60,0.60]	[-0.65, 0.65]	[-0.79,1.04]	[-2.50,11.2]
	Best var.		$p_{T(Z_1)}^{\gamma}$	$m_{\mu\mu}$	$m_{\mu\mu}$	$p_T^{e^+\mu^+}$
$\mathrm{ZZ}\gamma$	68% C.L.	No diagrams	[-0.62,0.61]	[-0.68,0.68]	[-0.81,1.06]	[-1.91,4.55]
	95% C.L.		[-0.90,0.90]	[-0.98,0.99]	[-1.23,1.49]	[-3.27,6.53]
	Best var.	$p_T^{l1}$	$m_{jj}$	$m_{jj}$	$p_{T(\gamma)}^{l1}$	$p_{T(\gamma)}^{l2}$
$\mathrm{VZ}\gamma$	68% C.L.	[-0.26,0.26]	[-0.55,0.54]	[-0.60,0.60]	[-0.11,0.11]	[-0.17,0.17]
	95% C.L.	[-0.37,0.37]	[-0.77,0.76]	[-0.84,0.84]	[-0.22,0.23]	[-0.33,0.34]
	Best var.	$p_T^{l1}$		$p_T^V$	$m_{\mu\mu}$	$p_T^{e^+\mu^+}$
VZZ	68% C.L.	[-0.63, 0.63]	Negligible	[-4.78,4.08]	[-0.80,0.65]	[-2.73,1.82]
	95% C.L.	[-0.97,0.97]		[-6.91,6.17]	[-2.22, 1.20]	[-3.78,2.82]
	68% C.L.	[-0.18,0.19]	[-0.37,0.37]	[-0.40,0.40]	[-0.11,0.11]	[-0.17,0.17]
Combination	95% C.L.	[-0.27,0.28]	[-0.53,0.53]	[-0.57,0.57]	[-0.21,0.21]	[-0.33,0.33]
VBS	95% C.L.	[-0.19,0.18]	-	[-1.02,1.08]	[-1.34, 0.96]	[-1.98,1.74]

WWW fully-leptonic 95% CL bound  $c_{W}$  is [-0.13(-0.22), 0.13(0.22)]

#### VBS, WW inclusive, triboson selection

Process	Variables of interest	Selections	Expected events			
$\mathbf{WW} \\ (pp \to 2l2\nu)$	MET, $m_{ll}$ , $p_{T,l^i}$ , $p_{T,ll}$ , $\eta_{l^i}$	$\begin{array}{l} {\rm MET} > 30 \; {\rm GeV} \\ m_{ll} > 60 \; {\rm GeV} \\ p_{T,l^1} > 25 \; {\rm GeV} \\ p_{T,l^2} > 20 \; {\rm GeV} \\  \eta_{l^i}  < 2.5 \end{array}$	(EW) 30600			
$\begin{array}{l} \mathbf{SSWW+2j} \\ (pp  \rightarrow \\ 2l2\nu jj) \end{array}$	$\begin{split} \text{MET, } m_{jj},  m_{ll},  \phi_{j^i},  p_{T,j^i} \\ p_{T,l^i},  p_{T,ll},  \Delta\eta_{jj},  \Delta\phi_{jj},  \eta_{j^i},  \eta_{l^i} \end{split}$	$\begin{array}{l} \mathrm{MET} > 30 \ \mathrm{GeV} \\ m_{jj} > 500 \ \mathrm{GeV} \\ m_{ll} > 20 \ \mathrm{GeV} \\ p_{T,ll} > 25 \ \mathrm{GeV} \end{array}$	(EW) 197 (EW) 493			
$\begin{array}{c} \mathbf{OSWW+2j} \\ (pp \rightarrow 2l2uij) \end{array}$	MFT m. m. é. n. n.	$p_{T,l^2} > 20 \text{ GeV}$ $p_{T,j^i} > 30 \text{ GeV}$ $\Delta m_{e^{-2}} > 25$	(QCD) 1967			
$\begin{array}{l} \mathbf{WZ+2j} \\ (pp \qquad \rightarrow \\ 2e\mu\nu jj) \end{array}$	$ \begin{array}{l} \text{MD1}, \ m_{jj}, \ m_{ll}, \ \phi_{ji}, \ p_{T,ji}, \ p_{T,li} \\ p_{T,ll}, \ \Delta\eta_{jj}, \ \Delta\phi_{jj}, \ \eta_{ji}, \ \eta_{li}, \ m_{3l} \\ p_{T,3l}, \ m_{WZ}, \ \delta\eta_{WZ}, \ \delta\phi_{WZ}, \ \Phi_{\text{planes}} \\ \theta_{lW}, \ \theta_{lZ}, \ \theta^* \end{array} $	$\begin{aligned} &\Delta \eta_{jj} > 2.5 \\ & \eta_{j^i}  < 5 \\ & \eta_{l^i}  < 2.5 \end{aligned}$	(QCD) 90			
$\begin{array}{cc} \mathbf{ZZ+2j} \\ (pp & \rightarrow \\ 2e2\mu 2j) \end{array}$	$\begin{split} m_{jj},  m_{l^{1}l^{2}},  m_{ll},  m_{4l},  \phi_{j^{i}},  p_{T,j^{i}},  p_{T,l^{i}}, \\ p_{T,l^{1}l^{2}},  p_{T,l\pm l\pm},  p_{T,l\pm l\mp},  p_{T,Z},  \Delta\phi_{jj}, \\ \Delta\eta_{jj},  \eta_{j^{i}},  \eta_{l^{i}} \end{split}$	$\begin{array}{l} m_{jj} > 400  {\rm GeV} \\ 60 & < m_{ll} \\ 120  {\rm GeV} \\ m_{4l} > 180  {\rm GeV} \\ p_{T,l^1} > 20  {\rm GeV} \\ p_{T,l^2} > 10  {\rm GeV} \\ p_{T,l^2} > 10  {\rm GeV} \\ p_{T,l^1} > 5  {\rm GeV} \\ p_{T,j^{1,2}} > 30  {\rm GeV} \\ \Delta \eta_{jj} > 2.4 \\  \eta_{j^i}  < 4.7 \\  \eta_{l^i}  < 2.5 \\ \Delta R(l^i,j^k) > 0.4 \end{array}$	(EW) 11 (QCD) 176			
$\mathbf{ZV+2j}$ $(pp \rightarrow 2ljjjj)$	$ \begin{array}{l} m_{jj}^{\max}, m_{jj}^{nomax}, m_{ll}, \phi_{ji}, p_{T,j^i}, p_{T,l^i} \\ p_{T,ll}, \Delta \eta_{jj}, \Delta \eta_{jj}^{nomax}, \Delta \phi_{jj}^{max} \\ \Delta \phi_{jj}^{nomax}, \eta_{j^i} \eta_{l^i} \end{array} $	$\begin{array}{l} m_{jj} > 1500 \ {\rm GeV} \\ 60 & < \ m_{jj}^V \\ < 110 \ {\rm GeV} \\ 85 < m_{ll} < 95 \ {\rm GeV} \\ p_{T,l^1} > 25 \ {\rm GeV} \\ p_{T,l^2} > 20 \ {\rm GeV} \\ p_{T,j^i} > 100 \ {\rm GeV} \\ \Delta \eta_{jj} > 3.5 \\  \eta_{j^i}  < 5 \\  \eta_{l^i}  < 2.5 \end{array}$	(EW) 142 (QCD) 50			

Process	Variables of interest	Selections
$WZ\gamma$	$ \not\!$	$50 < m_{\mu\nu} < 110  \text{GeV}$
$(pp \rightarrow \mu^{\pm} \overset{(-)}{\nu_{\mu}} 2e \gamma)$	$p_T^{SFOS-ll}, p_T^{l_i}, p_T^{\gamma}, p_T^{e^+\mu^+}, \eta_{l_i}, \eta^{\gamma},$	$60 < m_{ee} < 120  \text{GeV}$
	$\phi^{\gamma}, p_{T(Z\gamma)}^{l_i}, p_{T(Z)}^{l_i}, p_{T(WZ)}^{l_i}, p_{T(W)}^{l_i},$	$p_{T,l^1} > 20 \mathrm{GeV}$
	$p_{T(Z)}^{\gamma}, p_{T(Z)}^{W}, p_{T(Y)}^{W}, p_{T(WZ)}^{\gamma},$	$p_{T,l^2} > 10 \mathrm{GeV}$
14-11 (A)	$H_{\ell}^{T}(ee), H_{\ell}^{T}(3l\nu\gamma)$	$p_{T,l^i} > 5 \text{GeV}   \eta_{l^i}  < 2.5$
Expected	and a contraction of the second	$p_{T,\gamma} > 20 \mathrm{GeV}   \eta_{\gamma}  < 2.5$
events: (EW) 50		$Z_T > 30 \mathrm{GeV} \ \Delta R(l^i,\gamma) > 0.4$
$\mathbf{Z}\mathbf{Z}\gamma$	$m_{SFOS-ll},  m_{4l},  p_T^{Z_i},  p_T^{l_i},$	$60 < m_{SFOS-ll} < 120  \text{GeV}$
$(pp \rightarrow 2e \ 2\mu \ \gamma)$	$p_T^{SFOS-ll}, p_T^{\gamma}, p_T^{e^{\pm}\mu^{\pm}}, \eta_i, \eta^{\gamma},$	$p_{T,l^1} > 20 \mathrm{GeV}$
	$\phi^{\gamma}, p_{T(Z\gamma)}^{l_i}, p_{T(Z_i)}^{l_k}, p_{T(ZZ)}^{l_i}, p_{T(\gamma)}^{l_i},$	$p_{T,l^2} > 10 \mathrm{GeV}$
	$p_{T(Z_i)}^{\gamma}, p_{T(ZZ)}^{\gamma}, H_{\ell}^{T}(SFOS - ll),$	$p_{T,l^i} > 5 \mathrm{GeV}   \eta_{l^i}  < 2.5$
0.0110000000	$H_{\ell}^{T}(4l\gamma)$	$p_{T,\gamma} > 20 \mathrm{GeV}   \eta_{\gamma}  < 2.5$
Expected		$\Delta R(l^i, \gamma) > 0.4$
events: (EW) 22		
$VZ\gamma$	$m_{ll}, m_{jj}, p_T^Z, p_T^V, p_T^{\gamma}, p_T^{l_i}, p_T^{j_i},$	$50 < m_{jj} < 120  \text{GeV}$
$QCD - Z\gamma jj$	$\Delta \eta_{jj}, \ \Delta \phi_{jj}, \ \eta_{ji}, \ \eta_{l_i}, \ \phi_{ji}, \ \eta^{\gamma},$	$60 < m_{ll} < 120  \text{GeV}$
$(pp \rightarrow 2j \ 2l \ \gamma)$	$\phi^{\gamma}, p_{T(Z\gamma)}^{l_i}, p_{T(Z)}^{l_i}, p_{T(VZ)}^{l_i}, p_{T(V)}^{l_i},$	$p_{T,l^1} > 20 \mathrm{GeV}$
and and a second s	$p_{T(\gamma)}^{l_i}, p_{T(\gamma)}^{j_i}, p_{T(Z)}^{\gamma}, p_{T(Z)}^{V}, p_{T(\gamma)}^{V},$	$p_{T,l^2} > 10 \mathrm{GeV}$
	$p_{T(Z_{\tau})}^{V}, p_{T(VZ)}^{\gamma}, H_{\ell}^{T}(jj), H_{\ell}^{T}(ll),$	$ \eta_{l^i}  < 2.5$
Expected	$H_{\ell}^{T}(2l \ 2i\gamma)$	$p_{T,\gamma} > 20 \mathrm{GeV} \qquad  \eta_{\gamma}  < 2.5$
events:		$p_{T,j^{1,2}} > 30 \text{ GeV}  \eta_{j^i}  < 2.5$
(EW) 620		$\Delta R(l^i,\gamma) > 0.4$
(QCD) 31385		$\Delta R(l^i, j^k) > 0.4$
		$\Delta R(\gamma, j^k) > 0.4$
VZZ	$m_{SFOS-ll}, m_{jj}, m_{4l}, m_{4ljj},$	$50 < m_{jj} < 120  \text{GeV}$
QCD – ZZjj	$p_T^Z, p_T^{SFOS-ll}, p_T^{j_i}, p_T^{l_i}, p_T^V,$	$60 < m_{SFOS-ll} < 120  \text{GeV}$
$(pp \rightarrow 2j \ 2e \ 2\mu)$	$p_T^{e^{\pm}\mu^{\pm}}, \ \Delta\eta_{jj}, \ \Delta\phi_{jj}, \ \eta_{j_i}, \ \eta_{l_i}, \ \phi_{j_i},$	$p_{T,l^1} > 20 \mathrm{GeV}$
	$p_{T(ZZ)}^{l_i}, p_{T(ZZ)}^{j_i}, p_{T(Z_i)}^{l_i}, p_{T(Z_i)}^{j_i},$	$p_{T,l^2} > 10 \mathrm{GeV}$
	$p_{T_1(\mathcal{I}_{\epsilon})}^{Z_1}, p_{T_1(\mathcal{I}_{\epsilon})}^V, p_{T_1(\mathcal{I}_{\epsilon})}^V, H_{\epsilon}^T(jj),$	$p_{T,l^i} > 5 \mathrm{GeV}$
Expected	$H^{T}_{\ell}(SFOS-ll), H^{T}_{\ell}(4lij)$	$p_{T,j^{1,2}} > 30 \text{ GeV}$
events:	- e ( , , e ( 5 )	$ \eta_{j^i}  < 2.5   \eta_{l^i}  < 2.5$
(EW) 4		$\Delta R(l^i, j^k) > 0.4$
(QCD) 95		12 19 0 C

#### Impact of QCD bkg on sensitivity -VBS





#### **Best variables**



Study best variables to constrain pairs of coefficients in 1D and 2D fit
 p<sub>T</sub>l<sup>1</sup> for cHw, cw , p<sub>T</sub>j<sup>1,2</sup> for fermionic operators

	$c_{qq}^{(3,1)}$	$c_{qq}^{\left( 1,1 ight) }$	$c_{qq}^{\left(1 ight)}$	$c_W$	$c_{Hl}^{(3)}$	$c_{Hq}^{(3)}$	$c_{ll}^{(1)}$	$c_{Hl}^{\left(1 ight)}$	$c_{HW}$	$c_{Hq}^{(1)}$	$c_{HD}$	$c_{HWB}$	$c_{H\square}$
$c_{qq}^{(3)}$	$p_{T,j^1}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^1}$	$p_{T,j^1}$	$p_{T,j^1}$	$p_{T,j^1}$
$c_{qq}^{(3,1)}$		$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^1}$
$c_{qq}^{(1,1)}$	<u></u>		$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^1}$
$c_{qq}^{(1)}$				$p_{T,l^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^2}$	$p_{T,l^1}$	$p_{T,j^1}$	$p_{T,j^2}$	$p_{T,j^1}$	$p_{T,j^1}$
$c_W$					$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$m_{4l}$
$c_{Hl}^{(3)}$						$p_{T,j^1}$	$m_{jj}$	$m_{jj}$	$p_{T,ee}$	$p_{T,j^1}$	$m_{jj}$	$m_{jj}$	$m_{4l}$
$c_{Hq}^{(3)}$			3. <del></del> -1				$p_{T,j^1}$	$p_{T,j^1}$	$p_{T,l^1}$	$p_{T,j^1}$	$p_{T,l^1}$	$p_{T,l^1}$	$p_{T,l^2}$
$c_{ll}^{(1)}$	<u></u> 21		92 <u> </u>				~ <u> </u>	$m_{jj}$	$p_{T,Z}$	$p_{T,j^1}$	$p_{T,e^-\mu^-}$	$p_{T,Z}$	$p_{T,j^2}$
$c_{Hl}^{(1)}$				—			_	_	$m_{4l}$	$p_{T,j^1}$	$p_{T,e^+\mu^+}$	$p_{T,ee}$	$\Delta \eta_{jj}$
$c_{HW}$			3 <b></b> 31	_	-		<u> </u>	_		$p_{T,l^1}$	$m_{4l}$	$p_{T,Z}$	$m_{4l}$
$c_{Hq}^{(1)}$			_	_	-	-	_				$p_{T,e^-\mu^-}$	$p_{T,e^-\mu^-}$	$p_{T,l^2}$
$c_{HD}$	2-2							-		<u>(</u>	<u></u>	$m_{4l}$	$\Delta \eta_{jj}$
$c_{HWB}$	77-10	_	_	-			_	-	-				$m_{4l}$

# ZZ scattering at the LHC



- **Event topology**:
  - 2 Z-bosons produced centrally
  - 2 energetic tagging jets emitted back-to-back
- Signature based on diboson final states:
  - fully leptonic: 4 e/µ ; 2 jets

- Irreducible **tree-level contributions** to the final state:
  - **EW** =  $O(\alpha_{FW}^6)$  signal component
  - QCD=  $O(\alpha_{ew}^4 \alpha_s^2)$  bkg suppressed in high  $m_{ii} |\Delta \eta_{ii}|$  region
  - **INT = O(α<sup>5</sup><sub>EW</sub>α<sub>c</sub>) term: ~O(%) of** the signal Ο





# Summary of aQGCs limits

						:			
-		Observed	Expected	Observed	Expected	Observed	Expected	Observed	Expected
		$W^{\pm}W^{\pm} + WZ$		Z	Z	$\mathrm{W}\gamma$		$\mathrm{Z}\gamma$	
		(Tel	V <sup>-4</sup> )	(TeV	/_4)	(TeV	$(\text{TeV}^{-4})$		/-4)
	$f_{\rm T0}/\Lambda^4$	[-0.25, 0.28]	[-0.35, 0.37]	[-0.24, 0.22]	[-0.37, 0.35]	[-0.6, 0.6]	[-0.6, 0.6]	[-0.52,0.44]	[-0.64,0.57]
<b>T</b>	$f_{\rm T1}/\Lambda^4$	[-0.12, 0.14]	[-0.16, 0.19]	[-0.31, 0.31]	[-0.49, 0.49]	[-0.4, 0.4]	[-0.3, 0.4]	[-0.65,0.63]	[-0.81,0.90]
Transverse	$f_{\mathrm{T2}}/\Lambda^4$	[-0.35, 0.48]	[-0.49, 0.63]	[-0.63, 0.59]	[-0.98, 0.95]	[-1.0, 1.2]	[-1.0, 1.2]	[-1.36,1.21]	[-1.68, 1.54]
(4 gauge tensors)	$f_{\mathrm{T5}}/\Lambda^4$			2	_	[-0.5, 0.5]	[-0.4, 0.4]	[-0.45,0.52]	[-0.58,0.64]
(194496 (6119019)	$f_{\rm T6}/\Lambda^4$	_	—		_	[-0.4, 0.4]	[-0.3, 0.4]	[-1.02, 1.07]	[-1.30,1.33]
	$f_{\rm T7}/\Lambda^4$		—	—	<u> </u>	[-0.9, 0.9]	[-0.8, 0.9]	[-1.67,1.97]	[-2.15,2.43]
	$f_{\rm T8}/\Lambda^4$	_	—	[-0.43, 0.43]	[-0.68, 0.68]		_	[-0.36,0.36]	[-0.47,0.47]
	$f_{\mathrm{T9}}/\Lambda^4$	_	_	[-0.92, 0.92]	[-1.50, 1.50]	2		[-0.72,0.72]	[-0.91,0.91]
	$f_{\rm M0}/\Lambda^4$	[-27, 29]	[-36,37]			[-81,80]	[-77,76]	[-12.5,12.8]	[-15 8 16 0]
	$f_{ m M1}/\Lambda^4$	[-4.1, 4.2]	[-5.2, 5.5]			[-12, 12]	[-11, 11]	[-28.1,27.0]	[-35.0,34.7]
Mixed	$f_{\mathrm{M2}}/\Lambda^4$			—	—	[-2.8, 2.8]	[-2.7, 2.7]	[-5.21,5.12]	[-6.55,6.49]
(2) Lligge fields	$f_{\mathrm{M3}}/\Lambda^4$			,		[-4.4, 4.4]	[-4.0, 4.1]	[-10.2,10.3]	[-13.0,13.0]
(2 Higgs-fields	$f_{\rm M4}/\Lambda^4$	_		27 <u>-</u>		[-5.0, 5.0]	[-4.7, 4.7]	[-10.2,10.2]	[-13.0,12.7]
2 gauge tensors)	$f_{\rm M5}/\Lambda^4$	-			_	[-8.3, 8.3]	[-7.9, 7.7]	[-17.6,16.8]	[-22.2,21.3]
2 80080 (013013)	$f_{\rm M6}/\Lambda^4$	[-5.4, 5.8]	[-7.2, 7.3]		2. <u></u> 1	[-16, 16]	[-15, 15]		_
	$f_{\rm M7}/\Lambda^4$	[-57,60]	[-78,76]			[-21, 20]	[_10, 10]	[-44 7 45 0]	[-56 6,55 9]
Scalar	$f_{ m S0}/\Lambda^4$	[-5.7, 6.1]	[-5.9, 6.2]	<i>y.</i> y	s			_	_
	$f_{\rm G1}/\Lambda^4$	[-16, 17]	[-18, 18]						

• **Competitive** limits for different final states (**no clipping**)

• Expected/observed limits are in good agreement



*L* = 35.9 fb<sup>-1</sup>

#### **Processes of interest**



- Modelling of VBS ZZ  $2 \rightarrow 6(4)$  processes including non-resonant diagrams
  - both EWK and QCD-induced contributions for SM and EFT processes

Process		Variables of interest	Selections	Expected events
ZZ+2j		$m_{jj}, m_{l^1l^2}, m_{ll}, m_{4l}, \phi_{j^i}, p_{T,j^i}, p_{T,l^i},$	$m_{jj} > 400 \mathrm{GeV}$	(EW) 11
(pp	$\rightarrow$	$p_{T,l^1l^2}, p_{T,l^{\pm}l^{\pm}}, p_{T,l^{\pm}l^{\mp}}, p_{T,Z}, \Delta \phi_{jj},$	$60 < m_{ll} <$	(QCD) 176
$2e2\mu 2j)$		$\Delta\eta_{jj},\eta_{j^i},\eta_{l^i}$	$120{ m GeV}$	
			$m_{4l} > 180 \mathrm{GeV}$	
			$p_{T,l^1} > 20 \mathrm{GeV}$	
			$p_{T,l^2} > 10 \mathrm{GeV}$	
			$p_{T,l^i} > 5 \mathrm{GeV}$	
			$p_{T,j^{1,2}} > 30 \mathrm{GeV}$	
			$\Delta \eta_{jj} > 2.4$	
			$ \eta_{j^i}  < 4.7$	
			$ \eta_{l^i}  < 2.5$	
			$\Delta R(l^i, j^k) > 0.4$	

- Fully-leptonic final state 2e2µ+2j studied with VBS-enriched topology
  - study **observables most sensitive** to EFT-induced anomalies

#### **VBS ZZ NLO correction**

 $10^{-3}$ 

 $10^{-5}$ 

40 20

-20 -40

8 [%]



- VBS approximation less accurate at LO than NLO
  - in full computation **tri-boson processes** contribute where one **V->qq** decays Ο hadronically
  - for real QCD radiation, one of the two tagging jets may arise from V boson Ο