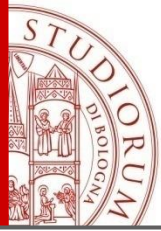


8th BCD ISHEP Cargèse School

NEUTRINO MASS HIERARCHY MEASUREMENT WITH DUNE

In collaboration with: prof. Marumi Kado

Giulia Lupi



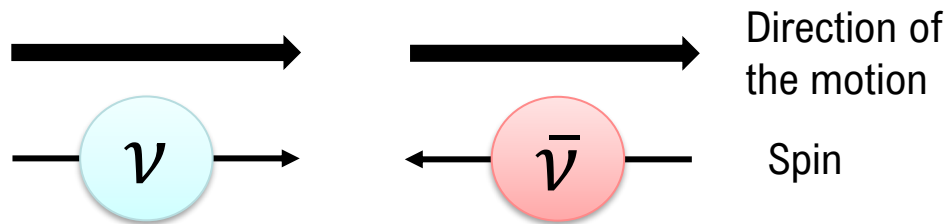
NEUTRINO IN THE STANDARD MODEL

Two notable facts about neutrino

From massless Dirac's equation energy solutions

Neutrino are emitted with negative helicity

Antineutrino are emitted with positive helicity



NEUTRINO \neq ANTINEUTRINO

From the V-A structure of the CC processes

$$J_{CC}^{\mu} = \bar{u}_1 \gamma^{\mu} \frac{1}{2} (1 - \gamma^5) u_2$$

ONLY ν_L AND $\bar{\nu}_R$ ENTER IN WEAK INTERACTIONS

Extract only the *left-handed* component of fermions and *right-handed* component of anti-fermions



DIRAC OR MAJORANA NEUTRINO?

In SM particles can be described either by Dirac or Weyl's equation but massive neutrinos can't:

- Weyl's equation describes massless fermion
- Dirac's equation needs 4 states ($f_L, f_R, \bar{f}_L, \bar{f}_R$) to explain mass of any particles
→ neutrino has only non zero two spinor's components ν_L and $\bar{\nu}_R$

1937 Italian physicist E. Majorana debuted another theory: could Neutrino and Antineutrino be the same particle $\nu_i = \bar{\nu}_i$?



DIRAC FERMIONS

$$\nu_i \neq \bar{\nu}_i$$

same mass but different lepton number $L(\nu_i) = -L(\bar{\nu}_i) = 1$ and the particles will be described as a quadruplet.

MAJORANA FERMIONS

$$\nu_i = \bar{\nu}_i$$

No lepton number conservation and distinction between particle and antiparticle. Neutrino is described by only two spin states.



NEUTRINO OSCILLATIONS

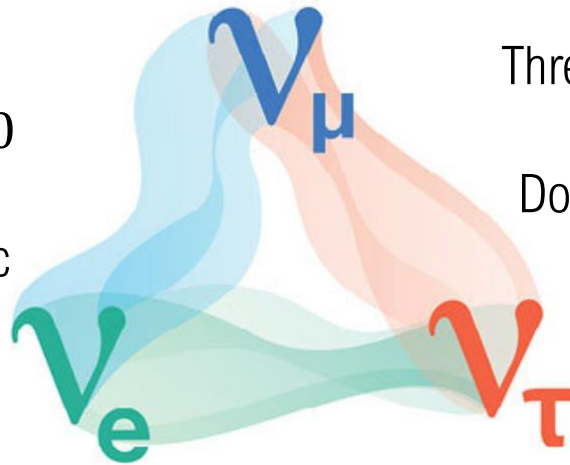
STANDARD MODEL

Three flavour eigenstates

$$\nu_e \quad \nu_\mu \quad \nu_\tau$$

One zero-mass state $m_\nu \simeq 0$

Does conserve family's leptonic number L_α ($\alpha = e, \mu, \tau$)



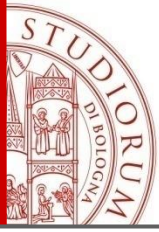
OSCILLATION THEORY

Three flavour eigenstates $\nu_e \quad \nu_\mu \quad \nu_\tau$

Three mass eigenstates $\nu_1 \quad \nu_2 \quad \nu_3$

Does NOT conserve family leptonic number L_α ($\alpha = e, \mu, \tau$)

Experimental observation of flavour changing in propagating neutrino can only be explained by assuming different neutrino mass eigenstates



NEUTRINO OSCILLATIONS

Neutrino with a definite flavour can be expressed as a linear combination of the three mass eigenstate and viceverse

$$|\nu_\alpha\rangle = \sum_i U_{\alpha,i}^* |\nu_i\rangle \quad \alpha = e, \mu, \tau \quad \text{FLAVOUR}$$

$$|\nu_i\rangle = \sum_\alpha U_{\alpha,i} |\nu_\alpha\rangle \quad i = 1, 2, 3 \quad \text{MASS}$$

$U_{\alpha,i}$ PMNS mixing matrix Unitary matrix that assumes 3x3 form for three neutrino families \rightarrow 9 degrees of freedom

DIRAC NEUTRINOS

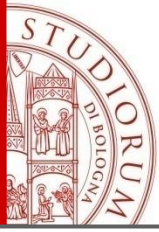
θ_{12} θ_{13} θ_{23} three mixing angles

δ_{CP} one phase related to the CP symmetry violation

MAJORANA NEUTRINOS

θ_{12} θ_{13} θ_{23} three mixing angles

δ_α δ_β δ_γ three phases related to the CP symmetry violation



NEUTRINO OSCILLATIONS

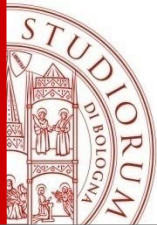
Using Dirac parametrization for the PMNS matrix

- θ_{12} θ_{13} θ_{23} three mixing angles $\theta_{ij} \in [0, \pi]$
- δ_{CP} one phase related to the CP symmetry violation $\in [0, 2\pi]$

(observation of oscillation does not distinguish between Majorana and Dirac neutrino!)

$$U_{\alpha,i} = \underbrace{\begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix}}_{\substack{\text{Atmospheric} \\ \nu_\mu \rightarrow \nu_\tau}} \underbrace{\begin{bmatrix} c_{13} & 0 & s_{13} e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta_{CP}} & 0 & c_{13} \end{bmatrix}}_{\substack{\text{Accelerated} \\ \nu_\mu \rightarrow \nu_e}} \underbrace{\begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{\substack{\text{Solar} \\ \nu_e \rightarrow \nu_\alpha}}$$

With $s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$, $ij = 12, 13, 23$



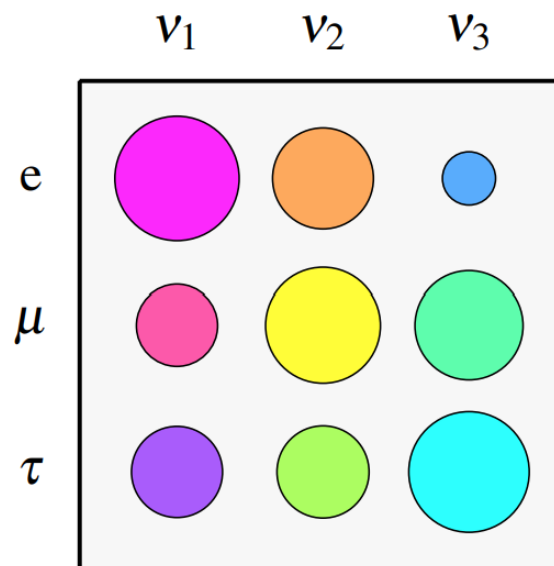
NEUTRINO OSCILLATIONS

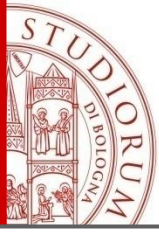
Parameter	Normal ordering	Inverted ordering
$\theta_{12}/^\circ$	$33.45^{+0.77}_{-0.75}$	$33.45^{+0.78}_{-0.75}$
$\theta_{23}/^\circ$	$42.1^{+1.1}_{-0.9}$	$49.0^{+0.9}_{-1.3}$
$\theta_{13}/^\circ$	$8.62^{+0.12}_{-0.12}$	$8.61^{+0.14}_{-0.12}$
$\delta_{CP}/^\circ$	230^{+36}_{-25}	278^{+22}_{-30}

Respect to CKM matrix parameters the mixing is enhanced

The octant for θ_{23} it has not been determined yet

Currently values for δ_{CP} are in favor of the CP symmetry violation





NEUTRINO OSCILLATIONS

Neutrino oscillation occur with a non zero probability that a neutrino with initial flavour α could be detected after a distance L with a different flavour $\beta \neq \alpha$.

- Each neutrino propagates as a plane wave with different phases for each mass state $\lambda = h/p_i$
- Ultra relativistic particle approximation $|p_i| \simeq E - \frac{m_i^2}{2E} \simeq E$

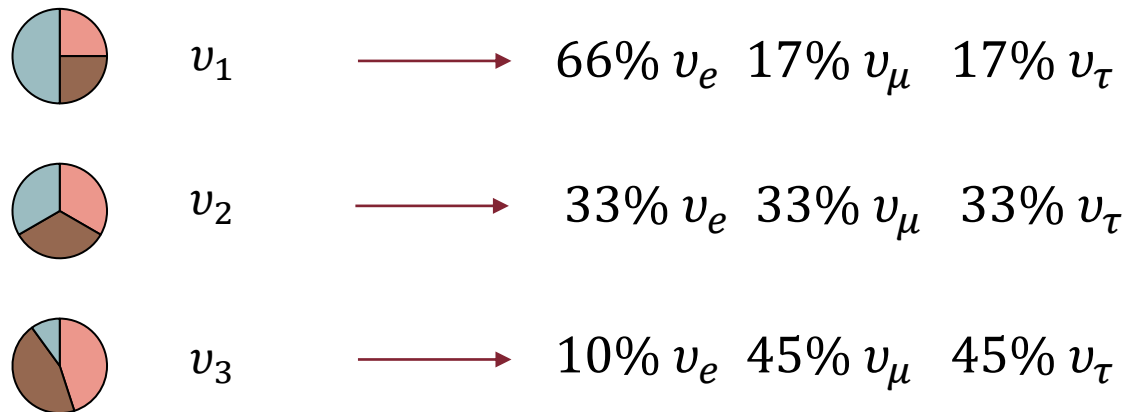
PROBABILITY OF OSCILLATION

$$P_{\alpha \rightarrow \beta} = \sum_i |U_{\alpha,i}|^2 |U_{\beta,i}|^2 + 2\text{Re} \sum_{i>j} U_{\alpha,i}^* U_{\beta,i} U_{\alpha,j} U_{\beta,j}^* e^{-i \frac{\Delta m_{ij}^2 L}{2E}} \quad \Delta m_{ij}^2 = m_i^2 - m_j^2$$

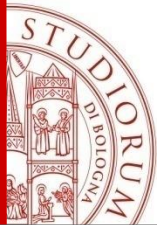
Observation of a flavour's transition implies the existence of different non zero mass states

MASS HIERACHY

- $P_{\alpha \rightarrow \beta} \propto \Delta m_{ij}^2$ **does not give any information about the absolute mass values!**
- Any massive state has different probabilities to interact as a neutrino with a given flavour



- The measurable quantity is $\Delta m_{ij}^2 \rightarrow$ studies on solar and atmospheric neutrinos first measured
- SOLAR (Δm_{12}^2) $\nu_e \rightarrow \nu_\mu \nu_\tau$ Homestake, Kamiokande, SAGE, SNO, GALLEX, Borexino
- ATMOSPHERIC (Δm_{23}^2) $\nu_\mu \rightarrow \nu_\tau$ Super-Kamiokande, K2K, MINOS, OPERa, IceCube



MASS HIERARCHY

Parameter	NO	IO
$ \Delta m_{12}^2 [10^{-5} eV^2]$	7.42 ± 0.21	7.42 ± 0.21
$ \Delta m_{23}^2 [10^{-3} eV^2]$	2.51 ± 0.27	2.49 ± 0.28
$ \Delta m_{31}^2 [eV^2]$	$O(10^{-3})$	$O(10^{-3})$

Eigenvalues m_1^2 and m_2^2 are similar while m_3^2 is more separated from the others

$$|\Delta m_{12}^2| \ll |\Delta m_{23}^2|$$

What is the sign of Δm_{31}^2 ?

Possible configurations that arise from the sign of Δm_{31}^2 are known as *hierarchies* (or ordering)

NORMAL $m_1 < m_2 < m_3$ $\text{sgn}(\Delta m_{31}^2) = +1$

INVERTED $m_3 < m_1 < m_2$ $\text{sgn}(\Delta m_{31}^2) = -1$

MSW EFFECT

- Neutrino interaction with matter can modify the probability of oscillation
- Propagation is altered due to the elastic scattering event with the medium's particles

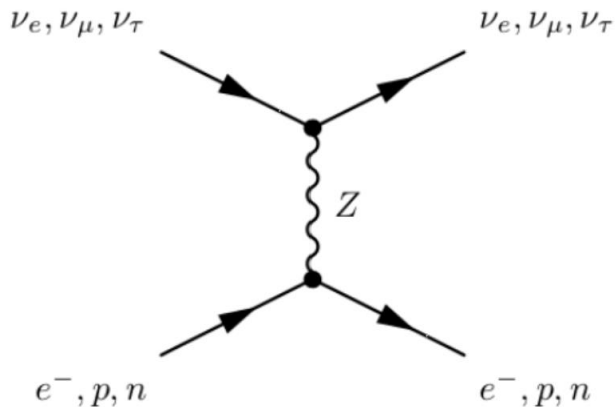
$$H \rightarrow H_{effective} = H + V_{MSW}$$

$$\Delta m^2 \rightarrow \Delta m_M^2 \begin{cases} \text{Increase for } \nu_e \\ \text{Decrease for } \bar{\nu}_e \end{cases}$$

NEUTRAL CURRENT WEAK INTERACTION (NC)

Cross section is independent from the leptonic family, all the three flavours can exchange

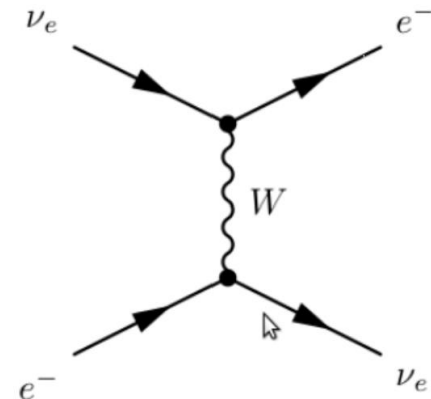
$$Z^0 \rightarrow V = 0$$

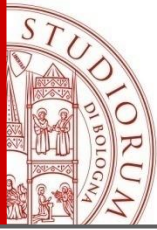


CHARGED CURRENT WEAK INTERACTION (CC)

Only electron neutrinos are involved in these processes, inducing an additional phase in the Hamiltonian

$$W^\pm \rightarrow V_{MSW} = \pm\sqrt{2} G_F N_e$$





MSW EFFECT

- Matter effect modifies the probability of oscillation

$$P_{\alpha \rightarrow \beta}(L, E) = |\Psi_{e,\beta}|^2 = \sin^2 2\theta_M \sin^2 \left(\frac{\Delta m_M^2 L}{4E} \right)$$

$$\Delta m_M^2 = \Delta m^2 \sqrt{(\cos 2\theta - A)^2 + (\sin 2\theta)^2}$$

$$\tan 2\theta_M = \frac{\sin 2\theta}{\cos 2\theta - A}$$

$$\text{where } A = \pm \frac{2\sqrt{2} G_F N_e E}{\Delta m^2}$$

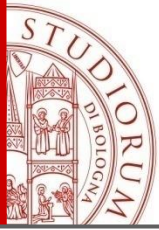
depends on neutrino (+) or antineutrino (-)
so it is sensible to $\text{sgn}(\Delta m_{ij}^2)$ → allows to discriminate the mass hierarchy !

- Becomes more important increasing the ratio

between E and L → *Long-Baseline* $\frac{L}{E} \simeq 10^{-3} \frac{\text{Km}}{\text{GeV}}$

- At fixed parameters results for neutrino and antineutrino are opposite

$$\frac{P_{\nu_\alpha \rightarrow \nu_\beta}}{P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta}} \begin{cases} > 1 & \text{NORMAL (NH)} \\ < 1 & \text{INVERTED (IH)} \end{cases}$$



MSW EFFECT

How to measure $\text{sgn}(\Delta m_{13}^2)$

- Particular values of N_e allows resonant transition between neutrino with different flavour
→ the effect is maximized for the condition:

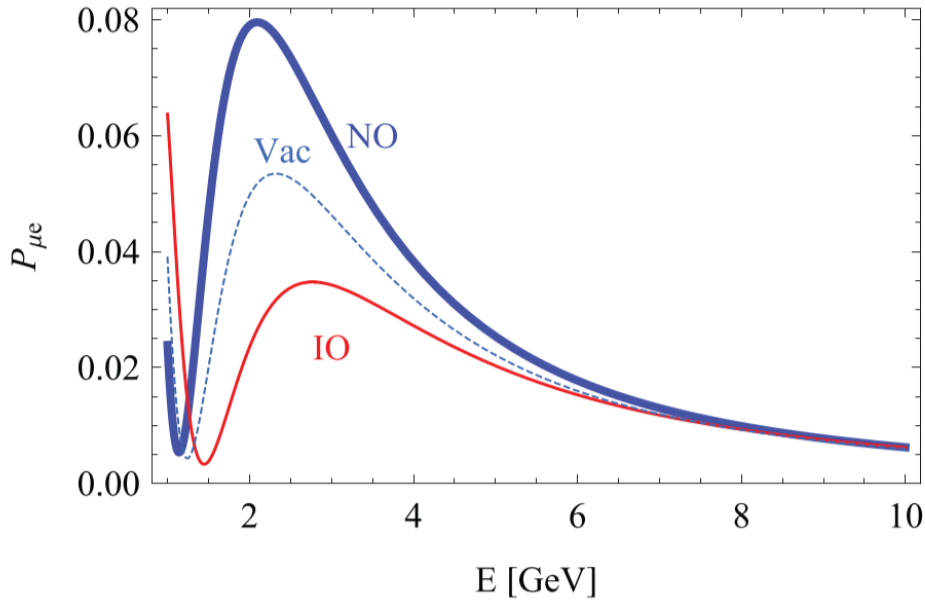
$$N_e^{Resonance} = \Delta m_i^2 \frac{\cos 2\theta_{ij}}{2\sqrt{2} G_F E_\nu} \rightarrow \theta_M = \frac{\pi}{4} \quad \text{MAXIMAL MIXING}$$

- Using different effects between neutrino and neutrino due to the presence of matter is an efficient tool for the ordering discrimination

$$\frac{P_{\nu_\alpha \rightarrow \nu_\beta}}{P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta}} \begin{cases} > 1 & \text{NORMAL (NH)} \\ < 1 & \text{INVERTED (IH)} \end{cases}$$

MSW EFFECT

Example: appearance events for transition $\nu_\mu \rightarrow \nu_e$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ (with neutrino beam)



One of the best ways to use the MSW effect in Earth matter such as in a long baseline neutrino experiment

Observation by MINOS an electron produced by a muon neutrino beam allows the determination of θ_{13}

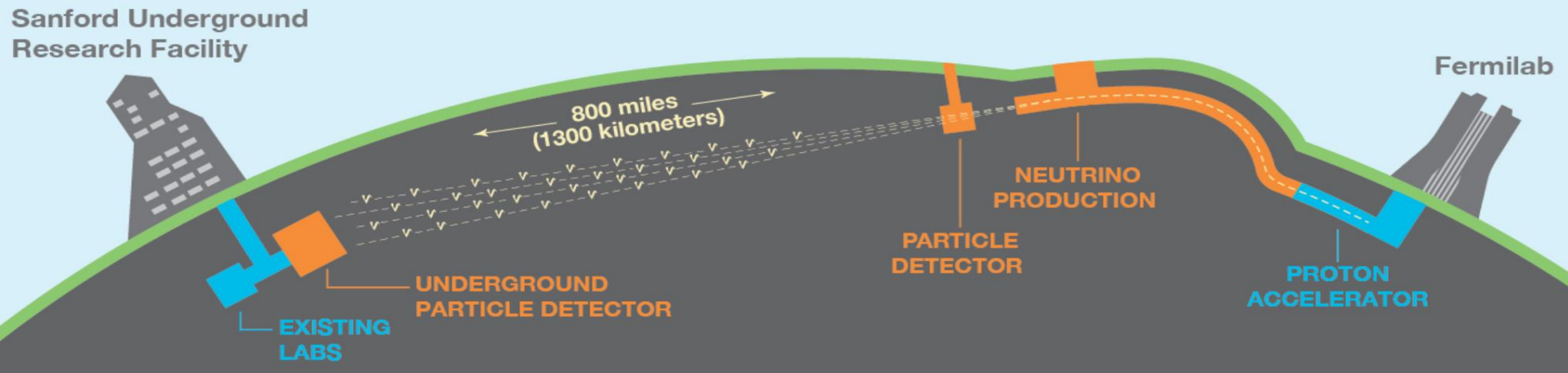
$$P_{\nu_\mu \rightarrow \nu_e} = \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{23}^2 L}{4E_\nu}$$

$P_{\nu_\mu \rightarrow \nu_e}$ $\left\{ \begin{array}{l} \text{Amplified for NH} \\ \text{Suppressed for IH} \end{array} \right.$

$P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e}$ $\left\{ \begin{array}{l} \text{Amplified for IH} \\ \text{Suppressed for NH} \end{array} \right.$

DUNE

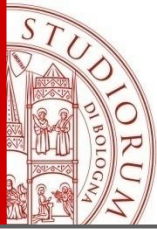
Deep Underground Neutrino Experiment



- Determine the correct mass hierarchy (confidence $>5\sigma$)
- Measuring δ_{CP} to verify CP violation
- Even more precise measurements of the other parameters
- Atmospheric neutrino
- Nuclear decays
- Supernovae neutrinos

PRIMARY GOALS

SECONDARY GOALS



DUNE

Deep Underground Neutrino Experiment

ACCELERATOR An effective way to access neutrino oscillation is by making intense neutrino beams using particle accelerators

Disappearance channels $\longrightarrow |\Delta m_{31}^2|$ and $\sin^2(2\theta_{23})$

Appearance channels $\longrightarrow \Delta m_{31}^2$ θ_{23} θ_{13} δ_{CP}

$$\mathcal{P}(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \left(1.27 \Delta m_{23}^2 \frac{L}{E} \right)$$

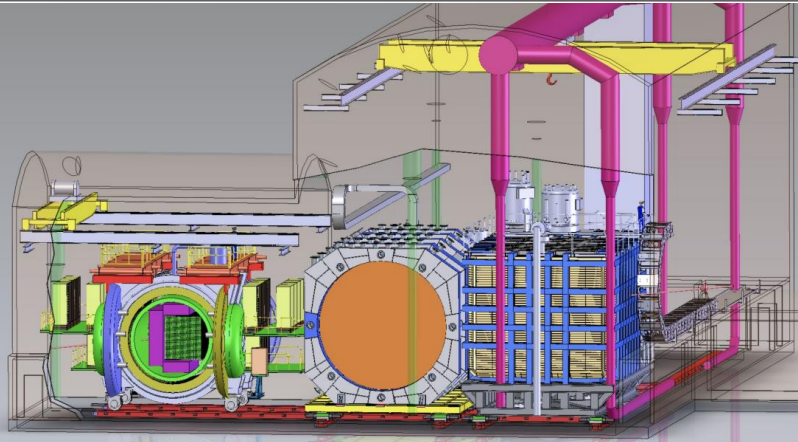


Forward Horn Current (FHC) mode produces a predominantly ν_μ beam

Reverse Horn Current" (RHC) mode produces predominantly $\bar{\nu}_\mu$

DUNE

The Near Detector



SAND **ND-GAr** **ND-LAr**

3 SUBSYSTEM

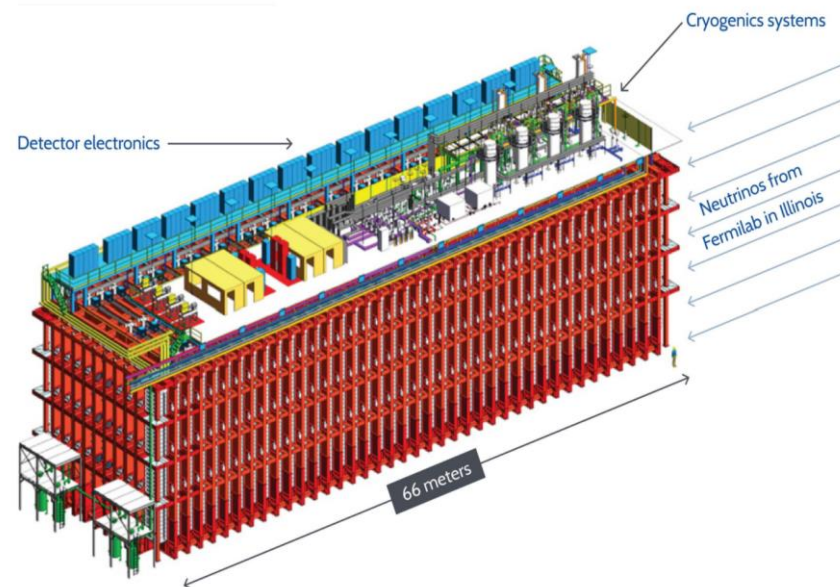
SAND	tracker surrounded by ECAL and magnet On-axis monitor of beam spectrum
ND-GAr	high-pressure GArTPC surrounded by ECAL and magnet muon spectrometer Constraint nuclear Interaction model
ND-LAr	modular pixelated LArTPC, primary target similar to FR

- Higher number of interactions with pure neutrino beam
- Information about the initial beam composition
- Smaller systematic errors on the final measures
- 600 m away from the Fermilab base
- ArgonCube Liquid Argon detector.
- Array of 5x7 ArgonCube modules (5 along and 7 transverse to the beam direction) sharing a common cryostat
- Minimal amount of inactive material

DUNE

Far Detector

- More chances to interact with matter
- Enhancement of oscillation probability
- $\frac{L}{E} \simeq 10^{-3} \frac{Km}{GeV}$
- 1300 km far from Fermilab
- 1.5 km underground
- 4 modules, filled with liquid Argon at -184 each with a total (fiducial) mass of 17 kt (10 kt)



DUNE

Deep Underground Neutrino Experiment

2018
Construction begin



2021
Far-Detector installation



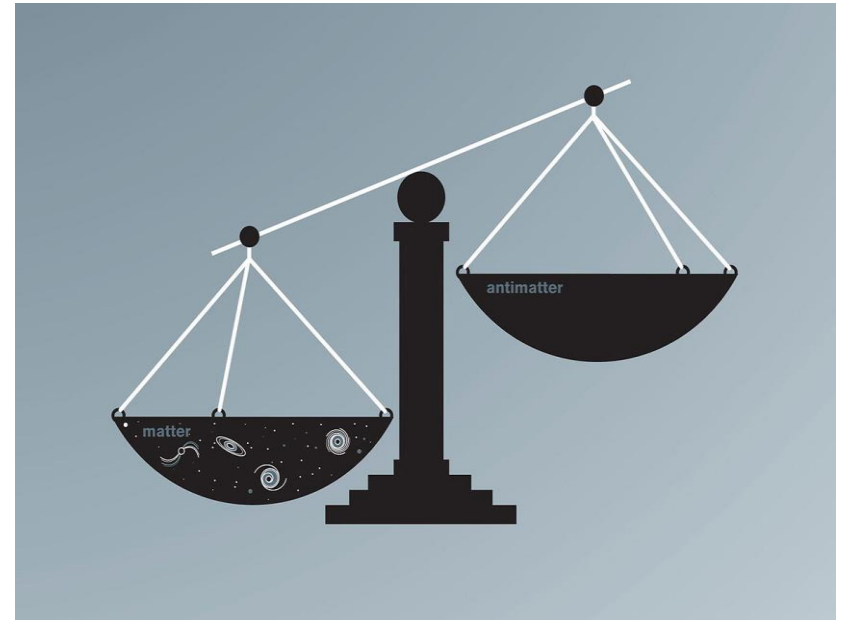
2025
Start ! (maybe..)



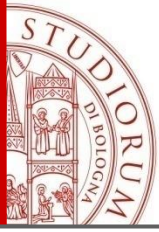
CONCLUSION

Still a lot of questions!

- What is the right mechanism through which neutrinos acquire their mass?
- Are neutrinos Majorana or Dirac fermions?
- Do they violate CP symmetry?



- Could be the explanation of the discrepancies between matter and antimatter in our universe?
- Are new beyond Standard Model theories necessary? (YES!)



8th BCD ISHEP Cargèse School

University of Bologna
a.a 2022-2023

Ajó
THANKS FOR YOUR ATTENTION!

Giulia Lupi