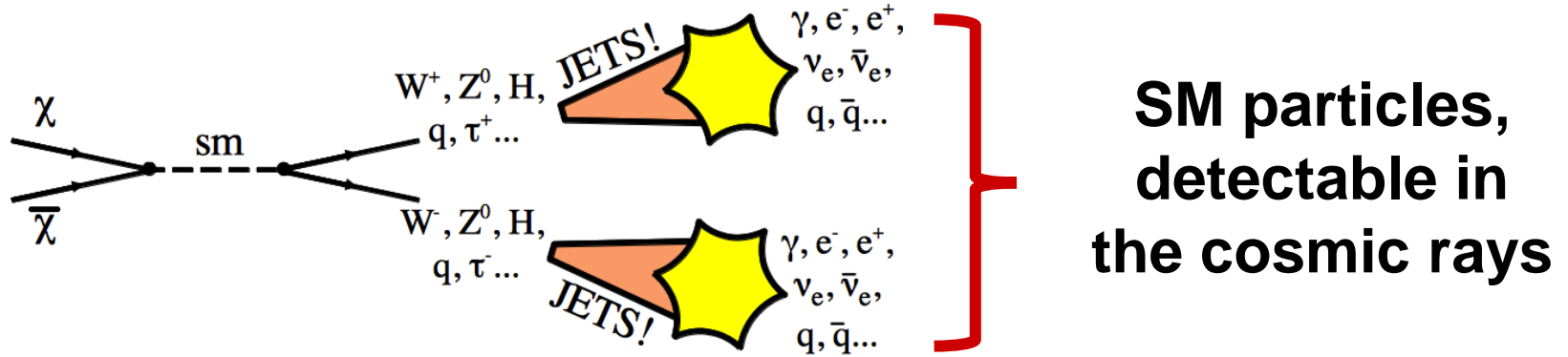


Scuola di Scienze  
Dipartimento di Fisica e Astronomia

# Coalescence parameter study and source characterization in the production of (anti)nuclei in high- energy collisions

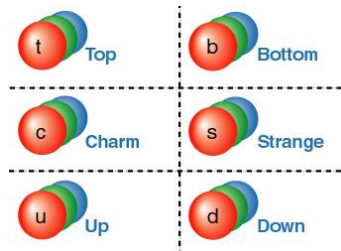
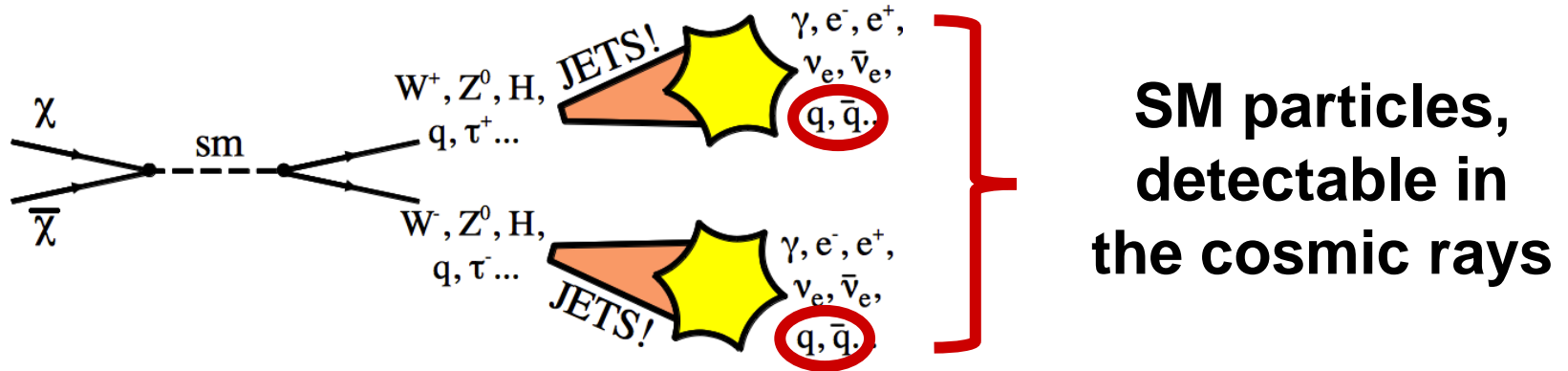
# WIMPs and cosmic (anti)nuclei

WIMPs: Dark matter candidates ( $m_\chi \approx 1 \text{ GeV} \div 1 \text{ TeV}$ )

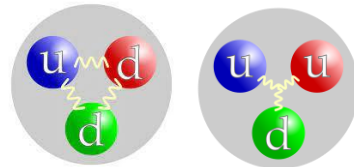


# WIMPs and cosmic (anti)nuclei

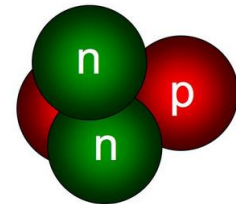
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(Anti)quark



(Anti)nucleons



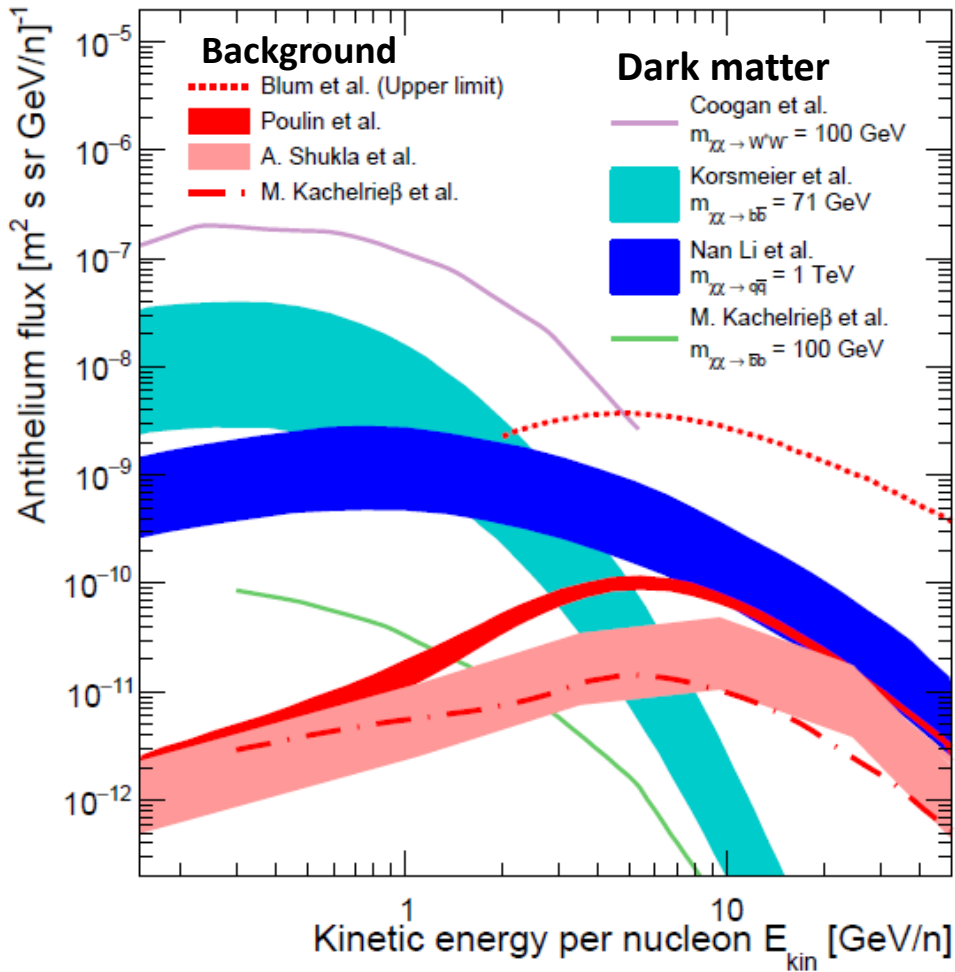
(Anti)nuclei

**Antinuclei** have **low astrophysics «background»** from collisions between cosmic rays and the ISM

**→ smoking guns**

# Cosmic antinuclei flux

*P. von Doetinchem et al., JCAP 08 (2020), 035*



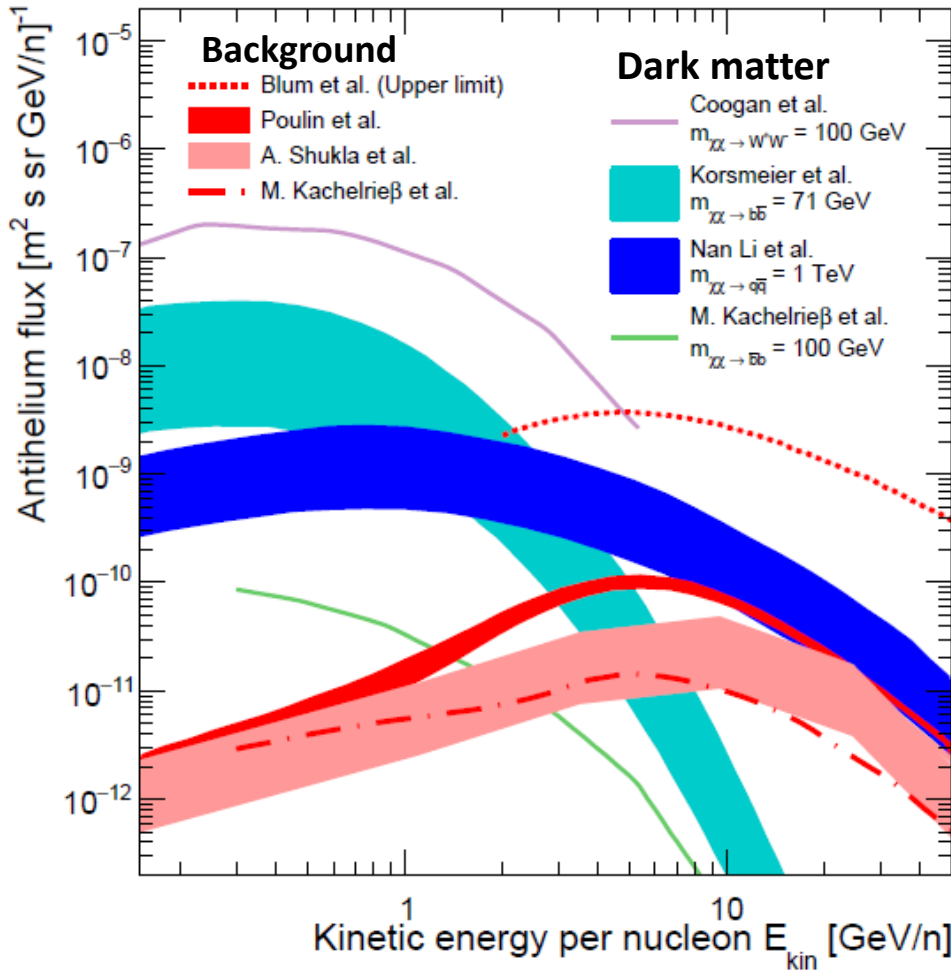
**Predictions**  
(production +  
propagation)

**VS**

**Experimental  
measurements**  
(esperimenti PAMELA,  
BESS, AMS-02...)

# Cosmic antinuclei flux

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**Predictions**  
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**Experimental measurements**  
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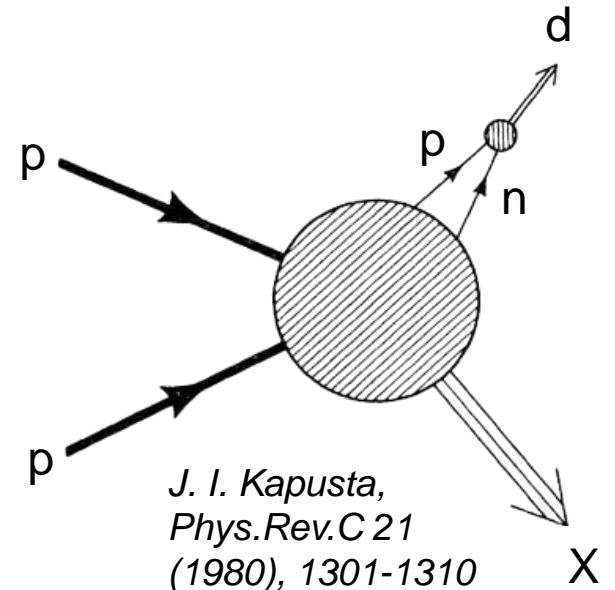
➔ Study of the formation model through LHC pp collisions data, where  $\bar{A}/A \approx 1$

# Coalescence model

Coalescence parameter  $B_A$

Operational definition:

$$E_A \frac{d^3 N}{dp_A^3} = B_A \left( E_{p,n} \frac{d^3 N_{p,n}}{dp_{p,n}^3} \right)^A \quad \vec{p}_p = \vec{p}_n = \frac{\vec{p}_A}{A}$$



Obtained by Bellini, Kalweit (*Phys.Rev.C* 99 (2019) 5, 054905):

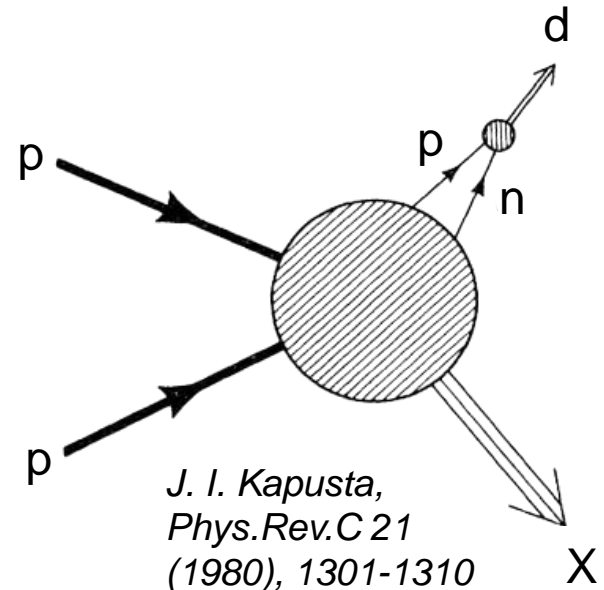
$$B_A = \frac{2J_A + 1}{2^A} \frac{1}{\sqrt{A}} \frac{1}{m_T^{A-1}} \left[ \frac{2\pi}{R^2 + \left(\frac{r_A}{2}\right)^2} \right]^{\frac{3}{2}(A-1)}, \quad m_T = \sqrt{p_T^2 + m_0^2}$$

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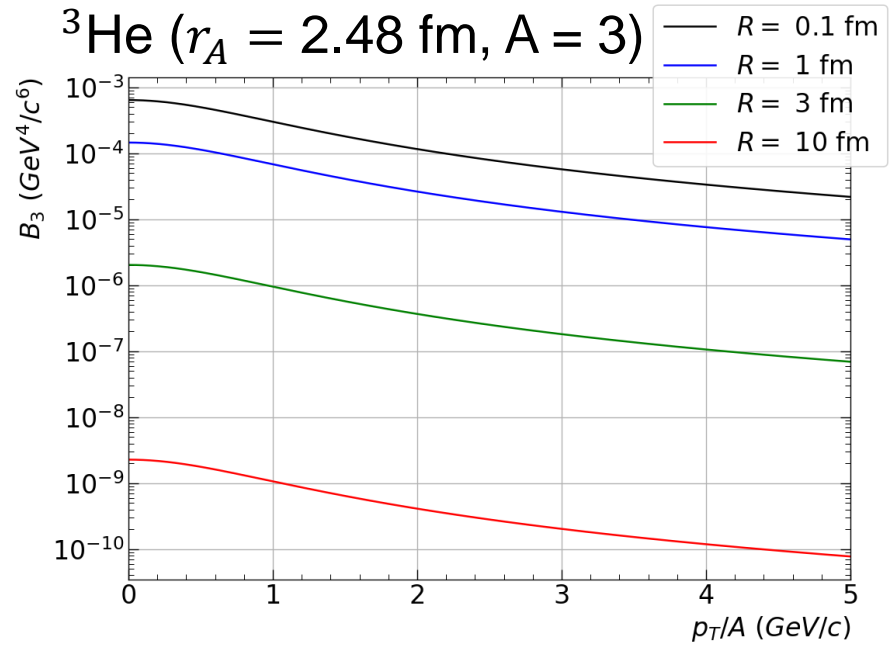
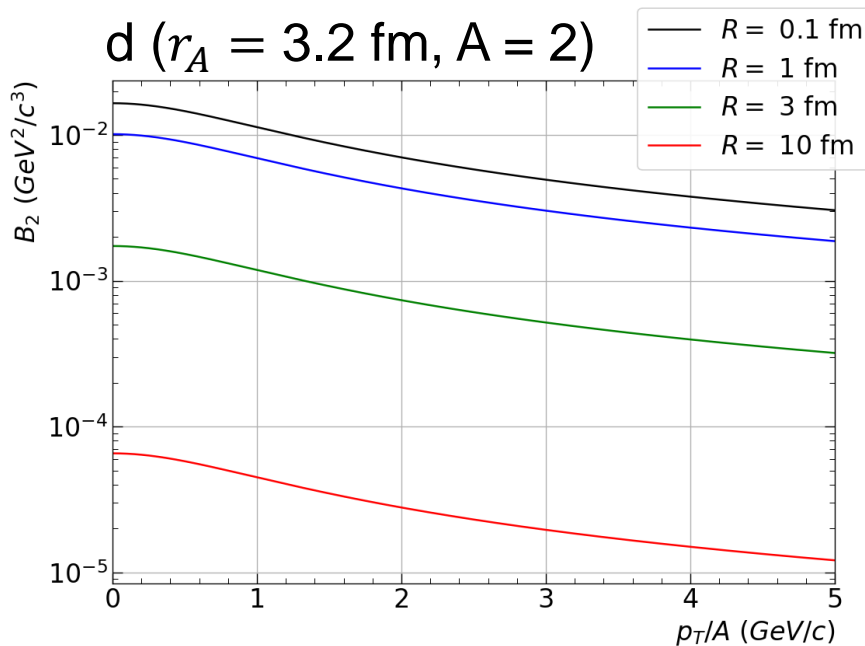
$$B_A = \frac{2J_A + 1}{2^A} \frac{1}{\sqrt{A}} \frac{1}{m_T^{A-1}} \left[ \frac{2\pi}{R^2 + \left(\frac{r_A}{2}\right)^2} \right]^{\frac{3}{2}(A-1)}, \quad m_T = \sqrt{p_T^2 + m_0^2}$$

What is the dependence  $R(p_T)$ ?

# Coalescence parameter

Study of the dependence of  $B_A$  on  $p_T$ , with constant  $R$

$$B_A = \frac{2J_A + 1}{2^A} \frac{1}{\sqrt{A}} \frac{1}{m_T^{A-1}} \left[ \frac{2\pi}{R^2 + \left(\frac{r_A}{2}\right)^2} \right]^{\frac{3}{2}(A-1)}$$



**→  $B_A$  decreases as either  $R$  or  $p_T$  increase**

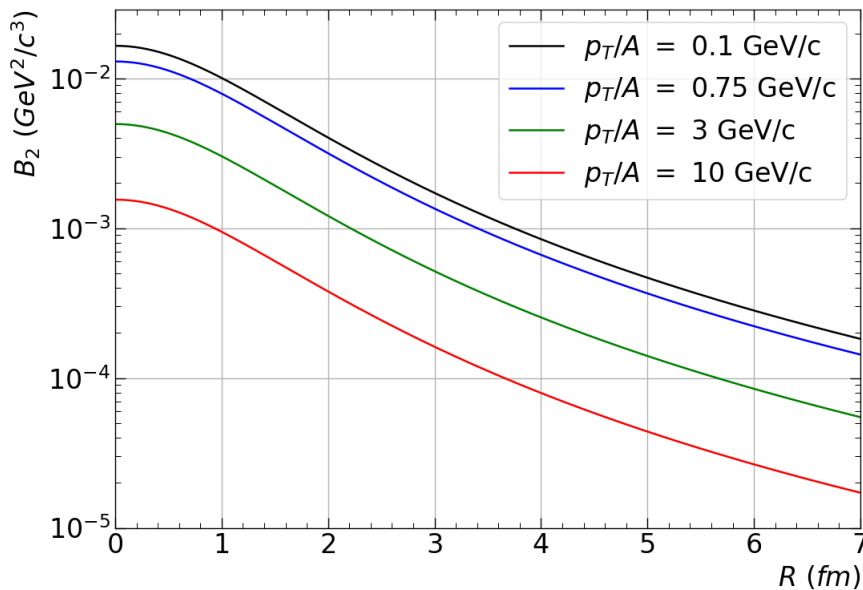


# Coalescence parameter

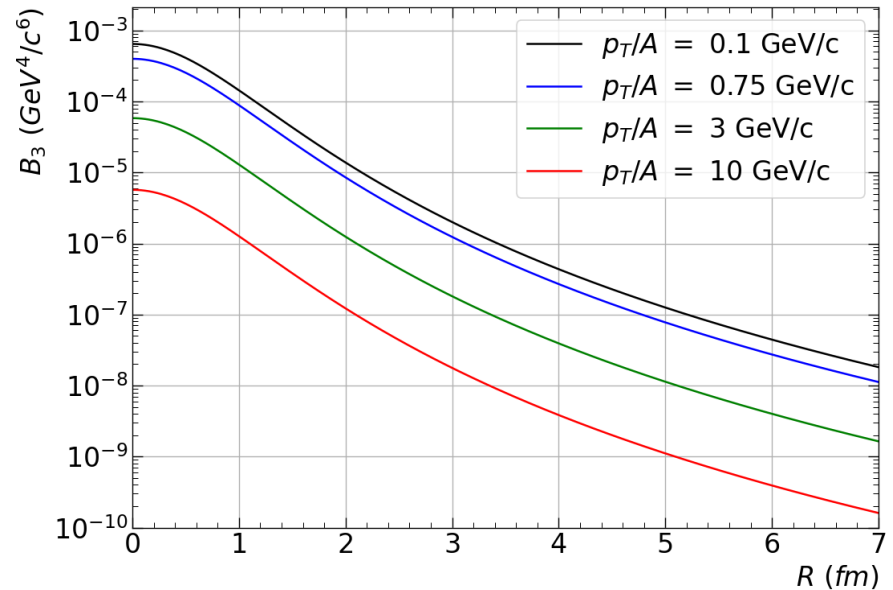
Study of the dependence of  $B_A$  on  $R$  (independent from  $p_T$ ), at  $p_T$  fixed

$$B_A = \frac{2J_A + 1}{2^A} \frac{1}{\sqrt{A}} \frac{1}{m_T^{A-1}} \left[ \frac{2\pi}{R^2 + \left(\frac{r_A}{2}\right)^2} \right]^{\frac{3}{2}(A-1)}$$

d ( $r_A = 3.2$  fm,  $A = 2$ )



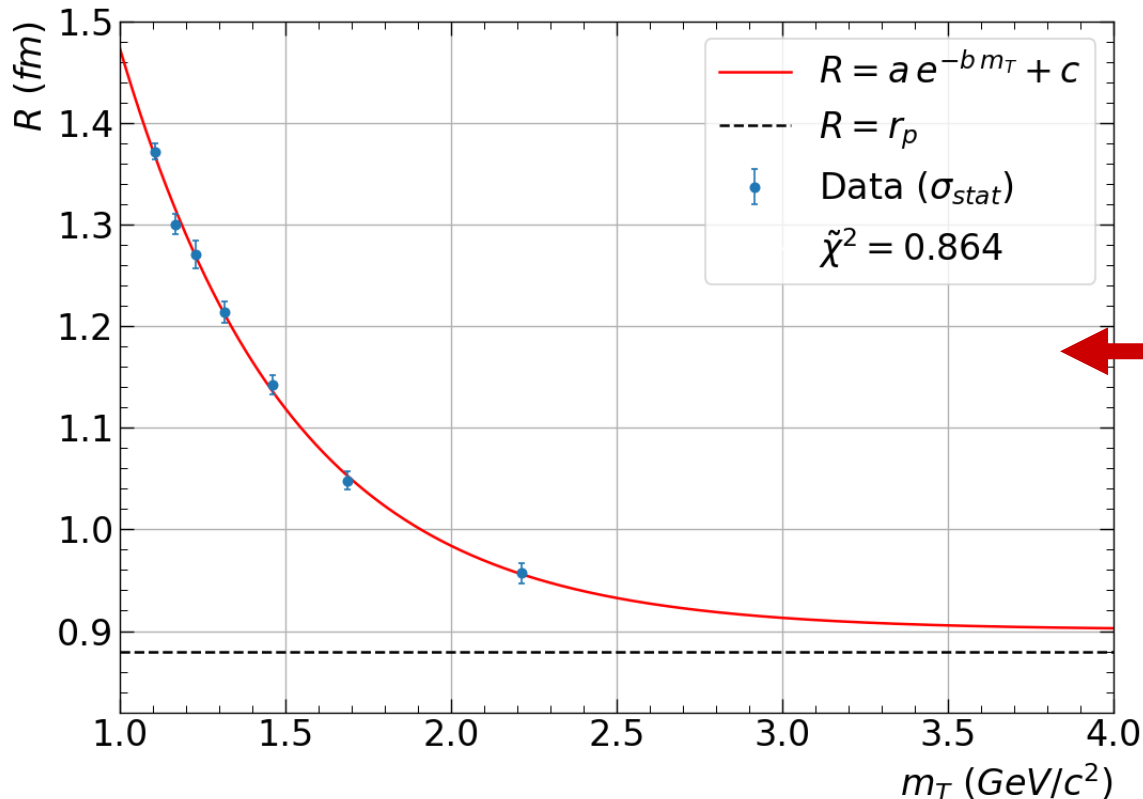
<sup>3</sup>He ( $r_A = 2.48$  fm,  $A = 3$ )



**→  $B_A$  decreases as either  $R$  or  $p_T$  increase**

# Dependence of $R$ on $p_T$

Fit to ALICE pp collisions high multiplicity data,  
at  $\sqrt{s} = 13$  TeV (*Phys.Lett.B* 811 (2020), 135849)



$$R_i(m_T)$$

---

$$R_1 = a e^{-b m_T} + r_p$$

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$$R_2 = a e^{-b m_T} + c$$

---

$$R_3 = a m_T e^{-b m_T} + r_p$$

---

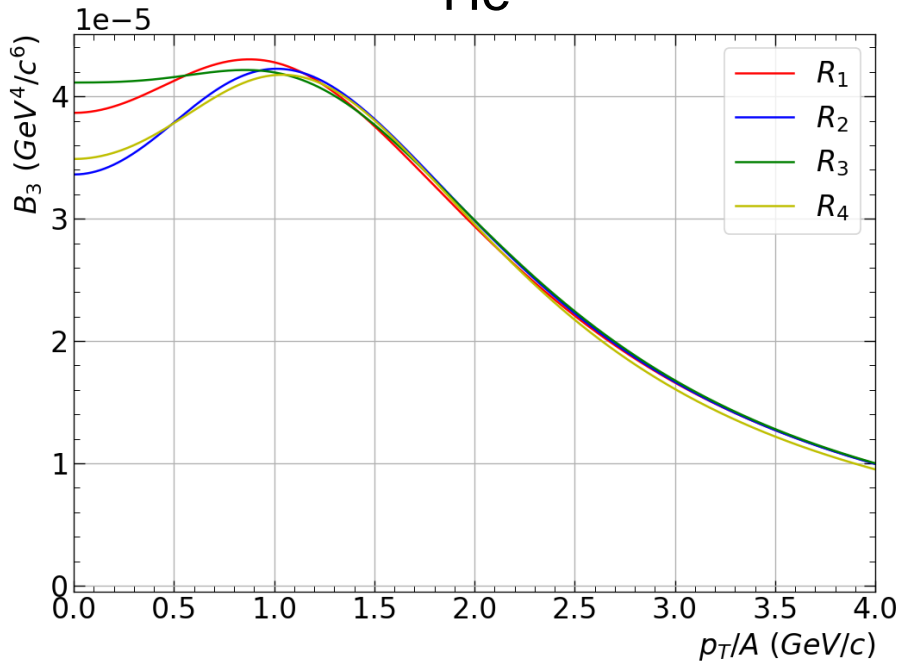
$$R_4 = a m_T e^{-b m_T} + c$$

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# Dependence of $R$ on $p_T$

It is possible to express  $B_A(p_T)$  through the fit functions and to estimate the uncertainty due to the choice of the function.

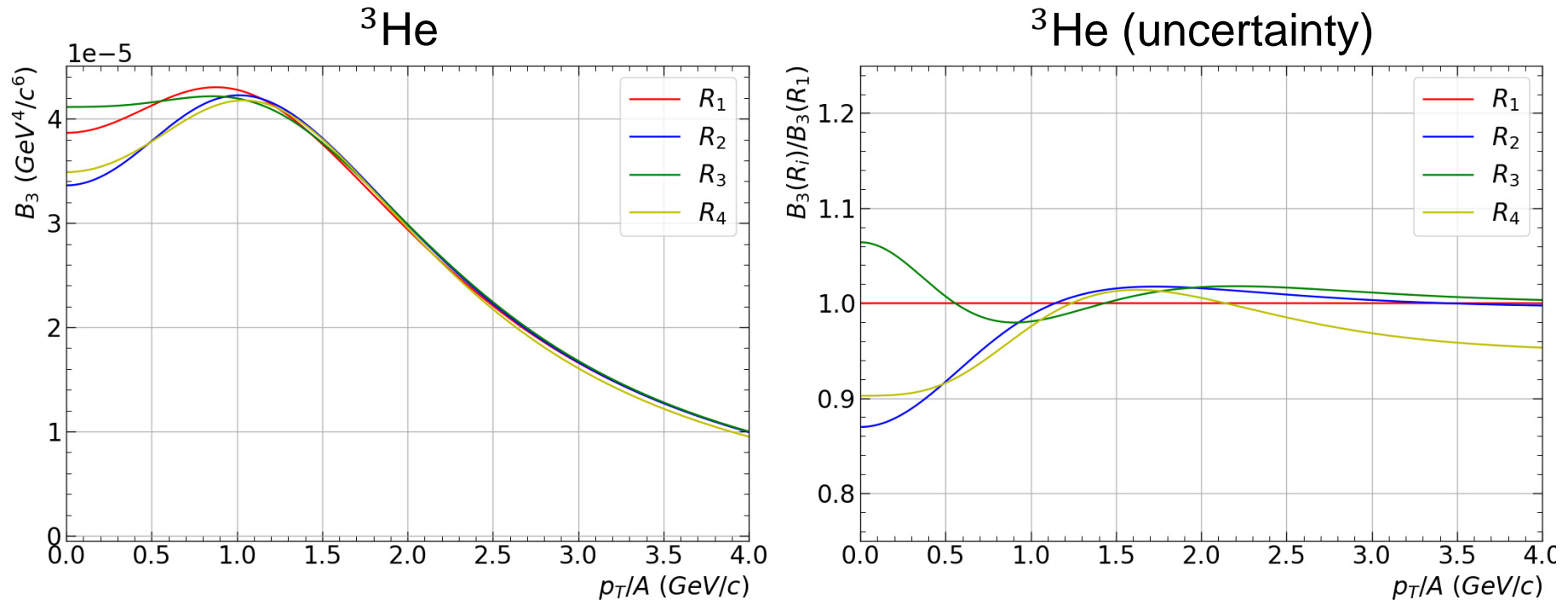
${}^3\text{He}$



→  $B_A$  decreases as  $p_T$  increases

# Dependence of $R$ on $p_T$

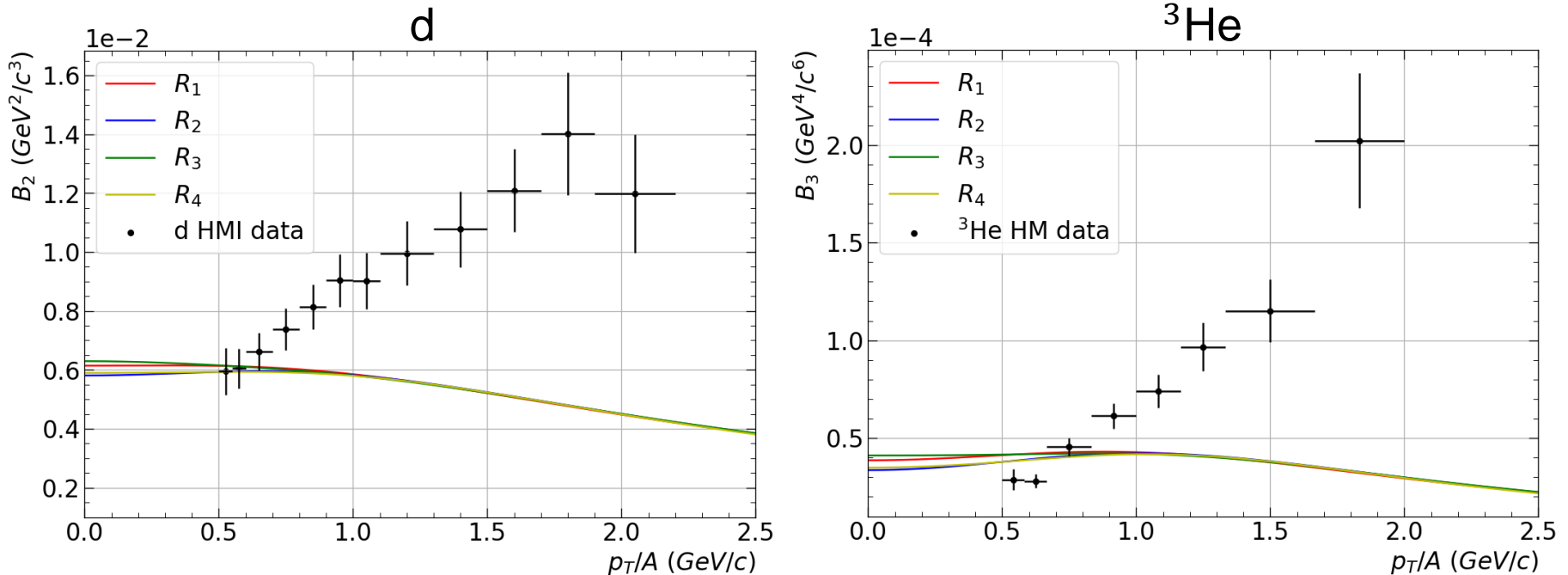
It is possible to express  $B_A(p_T)$  through the fit functions and to estimate the uncertainty due to the choice of the function.



- ➔  $B_A$  decreases as  $p_T$  increases
- ➔ At low momenta, due to the scarcity of data, uncertainties up to **+6%** and **-13%** for  ${}^3\text{He}$  (increasing with  $A$ ).

# Results

Comparison with ALICE data from **d** and  $^3\text{He}$  high multiplicity pp collisions, at  $\sqrt{s} = 13 \text{ TeV}$  (ALICE Collaboration, JHEP 01 (2022), 106)



The curve found **is not able to reproduce** the data, in which  $B_A$  increases as  $p_T$  increases.

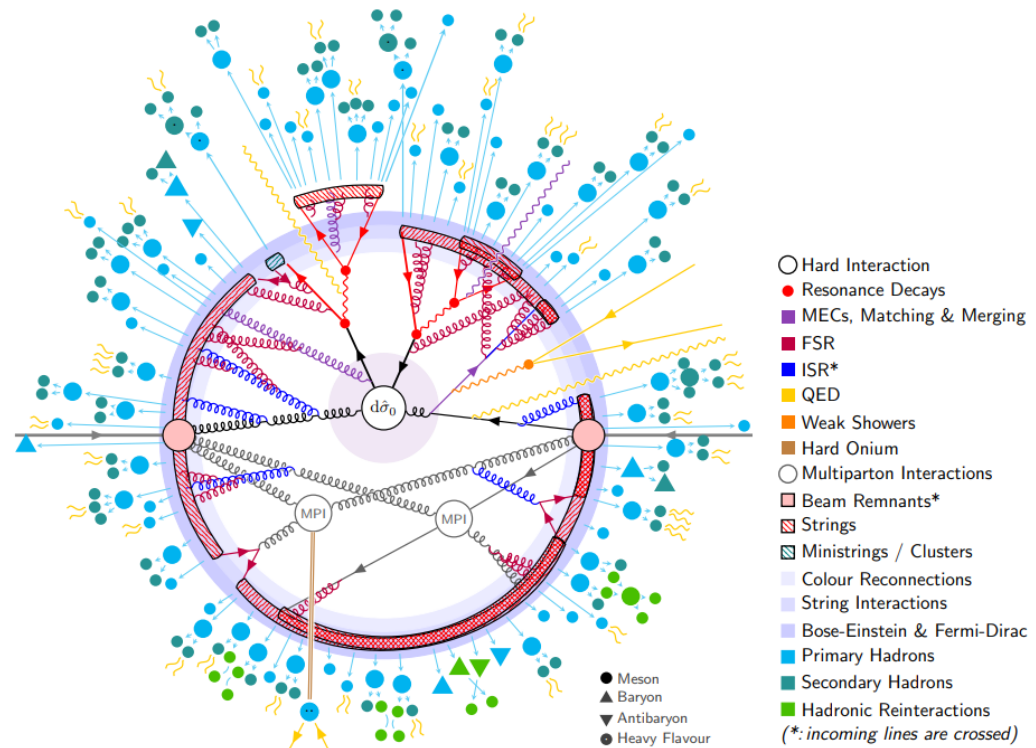
**The possible explanation:** The model used to obtain the  $B_A$  formula was derived for **Pb-Pb collisions**, and thus has to be adapted to **pp collisions**.

# PYTHIA 8.3

A different approach: **event by event simulation** of nuclei formation via coalescence.

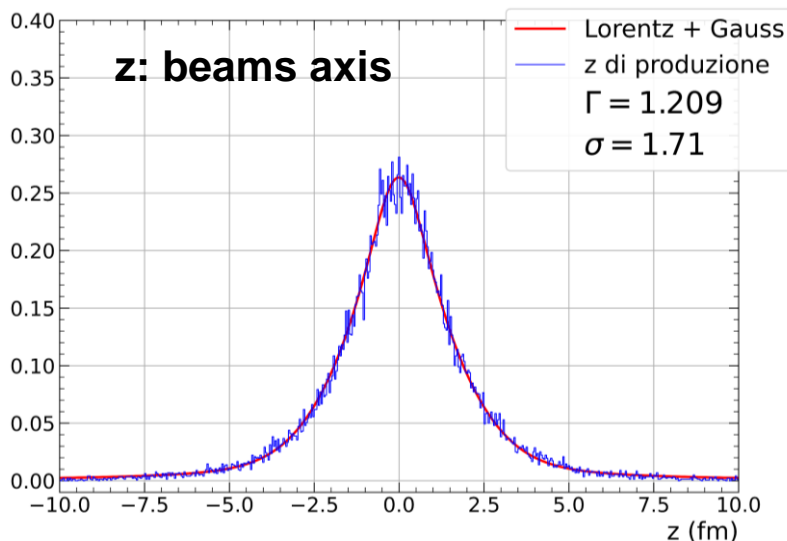
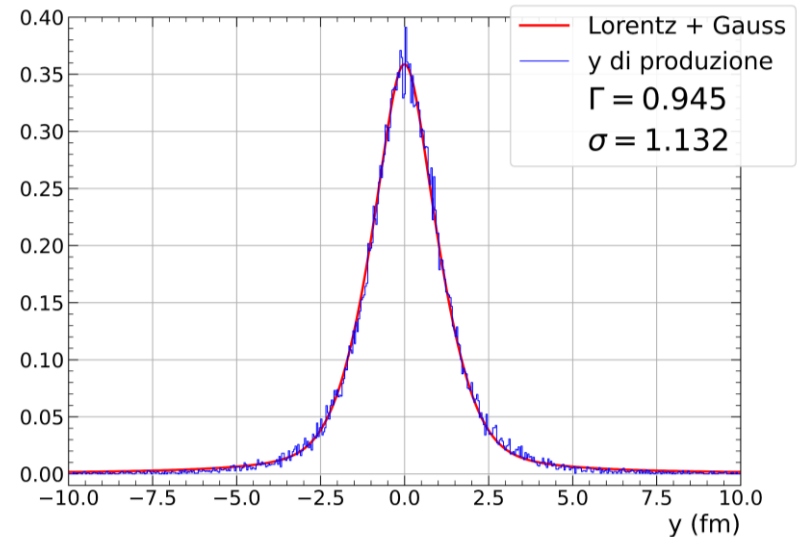
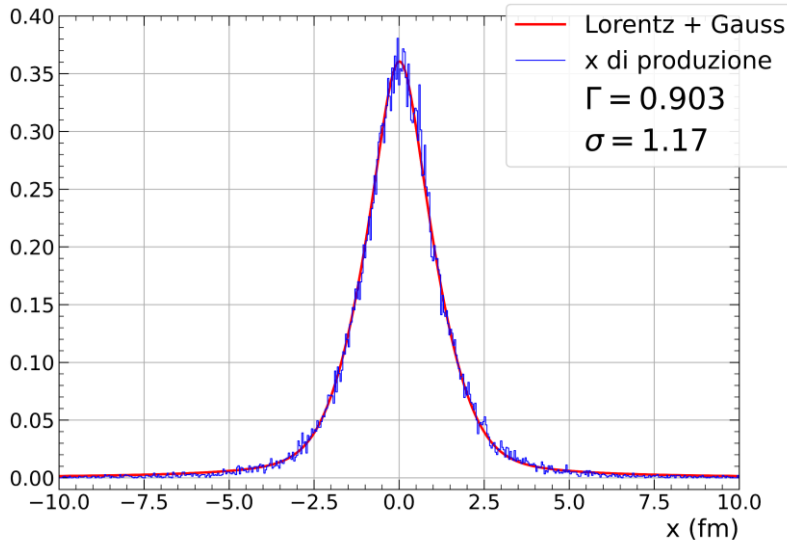
➔ Characterization of the midrapidity ( $|y| < 1$ ) proton source through the Monte Carlo generator **PYTHIA 8.3** with **Monash 2013** configuration.

*A general PYTHIA event scheme.*  
*Christian Bierlich et al.,*  
*ArXiv:2203.11601*



# The source in PYTHIA

Production coordinates distributions (protons):



➔ **Gaussian + lorentzian**

➔ **Spheroidal** source

➔  $\sigma, \Gamma \approx 1$  fm, in agreement with other data

# Conclusions

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## Work

Inclusion of the  $R(p_T)$  dependence in the calculation for the coalescence parameter  $B_A$  via a fit to ALICE data, with applications in the study of cosmic **(anti)nuclei formation**.

## Conclusions

The  $B_A(p_T)$  dependence obtained **does not reproduce the experimental data**. The model at the origin of  $B_A$  must be revised and properly adapted to pp collisions.

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## Work

Study of the **proton source** as generated by PYTHIA 8.3 for future simulations of nuclei formation via coalescence in pp events.

## Conclusions

PYTHIA generates a **spheroidal source**, which is isotropic in the  $xy$  plane. The characteristic dimension is of the order of **1 fm**, comparable with the experimental data.

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**THANK YOU FOR  
YOUR ATTENTION**