

# Early Universe hypercharge breaking & neutrino mass generation

Álvaro Lozano Onrubia

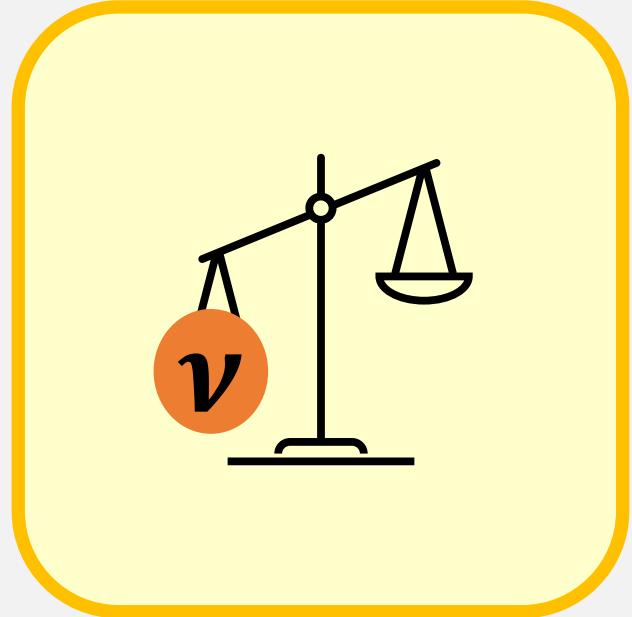
Based on arXiv:2308.09206 [hep-ph]

in collaboration with

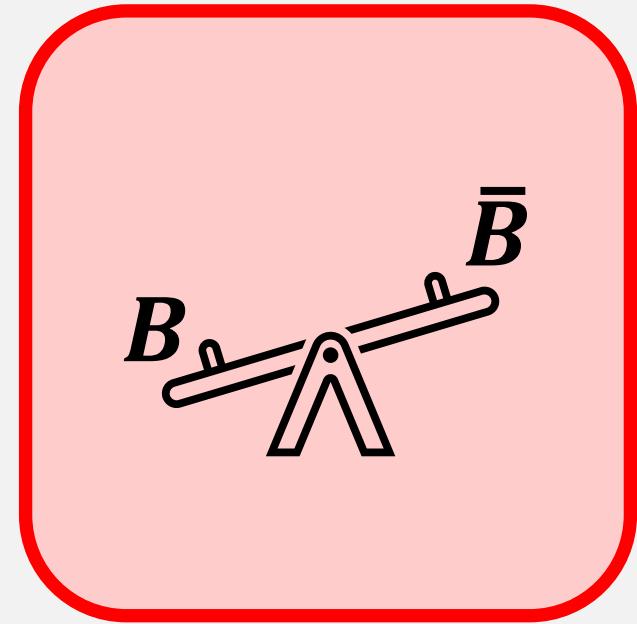
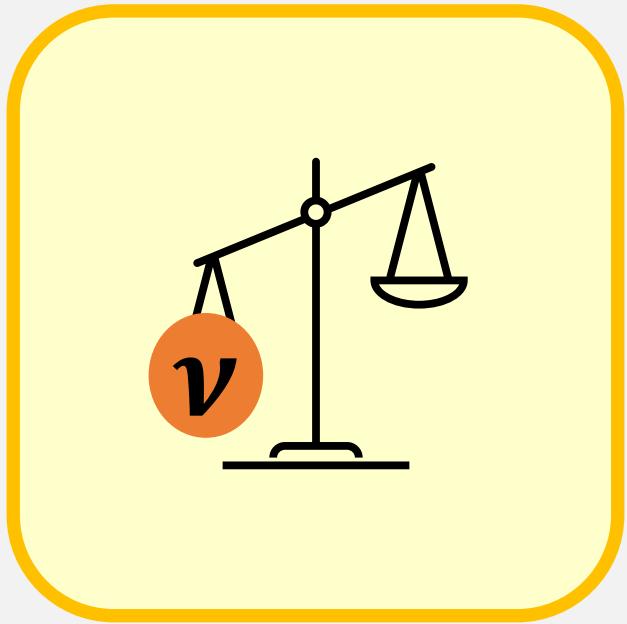
S. López Zurdo, L. Merlo, J. M. No

Invisibles23 Workshop, 29/08/2023

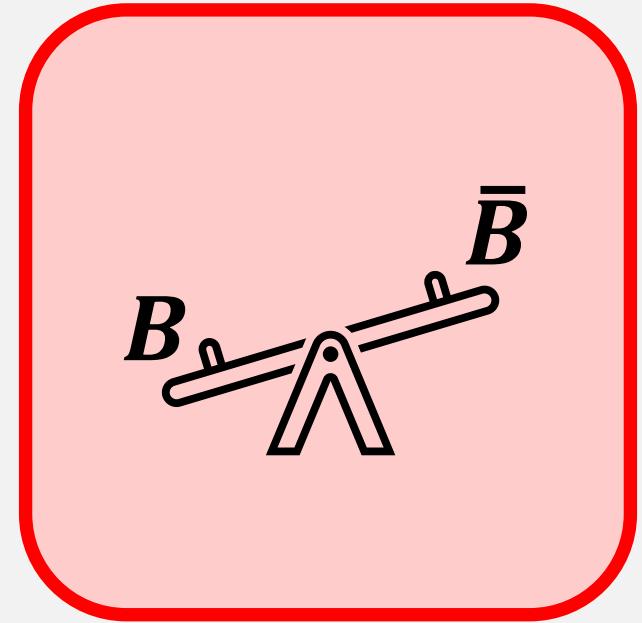
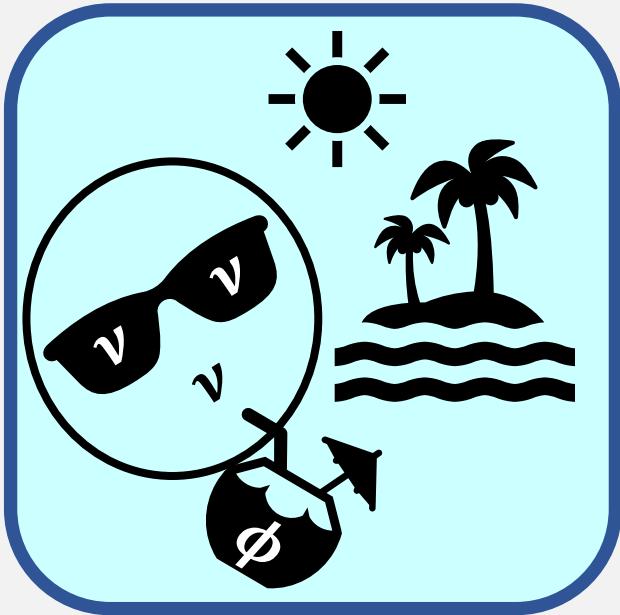
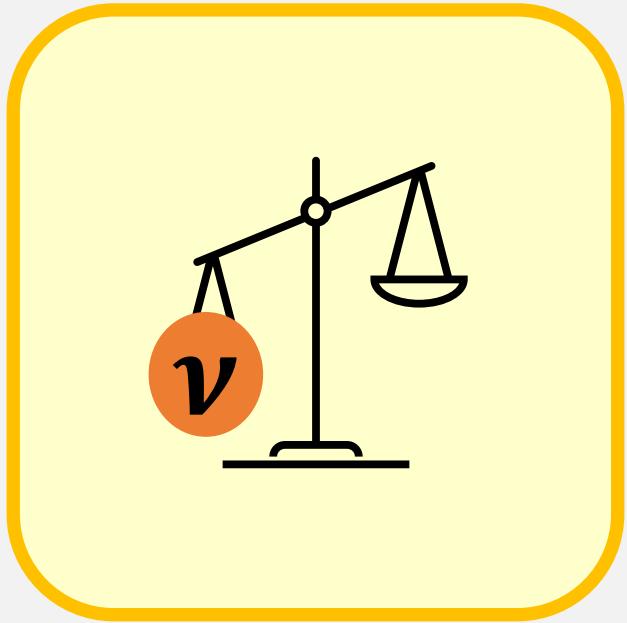




**Neutrino masses**



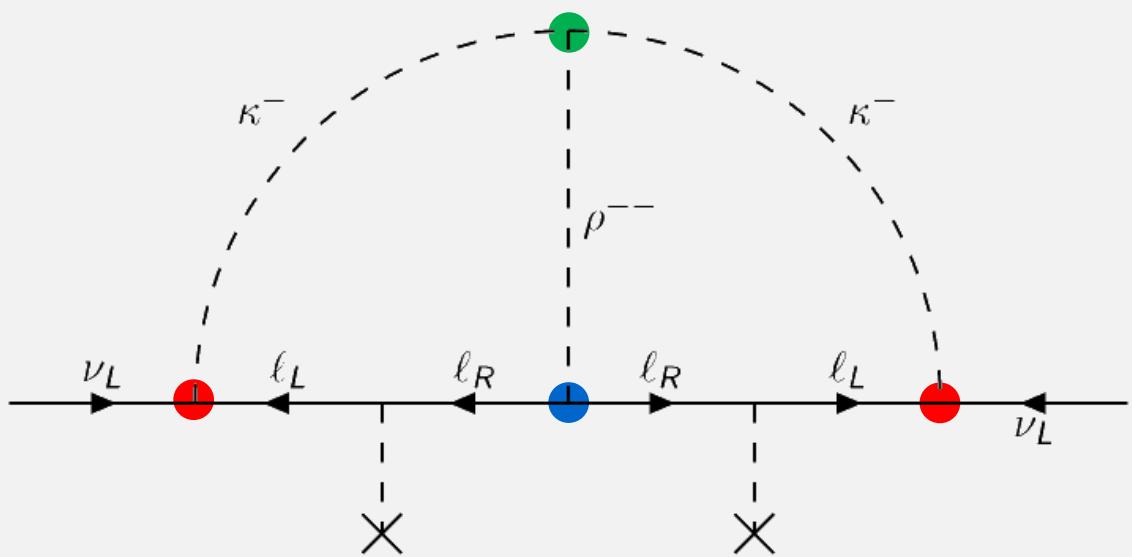
**Baryon asymmetry**



Exotic pheno



# Zee-Babu (ZB) model

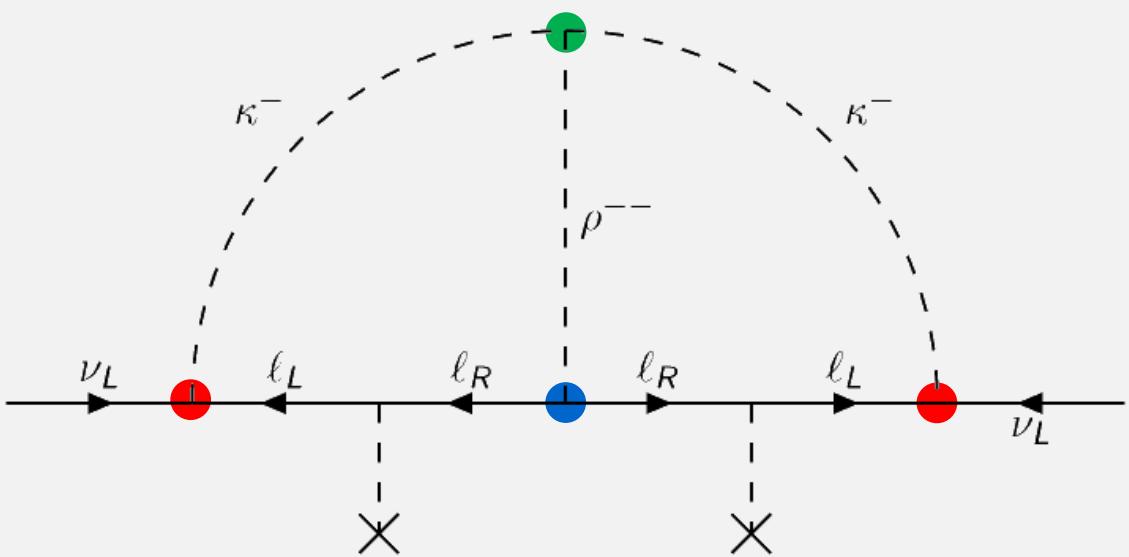


**Fig. 1:** Feynman diagram generating the neutrino masses in the ZB model [1, 2].

- Canonical 2-loop **radiative  $\nu$  mass** model [1, 2]
- Some features:
  - scalars  $\kappa^+ \sim (1_C, 1_L, \mathbf{2}_Y)$  and  $\rho^{++} \sim (1_C, 1_L, \mathbf{4}_Y)$
  - $\mathcal{L}_{ZB} \supset \overline{L}_L \mathbf{f} L_L \kappa^+ + \overline{l}_R^c \mathbf{g} l_R \rho^{++} + \textcolor{teal}{\mu} \rho^{++} \kappa^- \kappa^- + \text{h.c.}$   
with  $\mathbf{f}$  antisymmetric and  $\mathbf{g}$  symmetric.
  - neutrino masses:  $M_\nu \sim \textcolor{teal}{\mu} \mathbf{f} Y_l \mathbf{g}^T Y_l \mathbf{f}^T$



# Zee-Babu (ZB) model

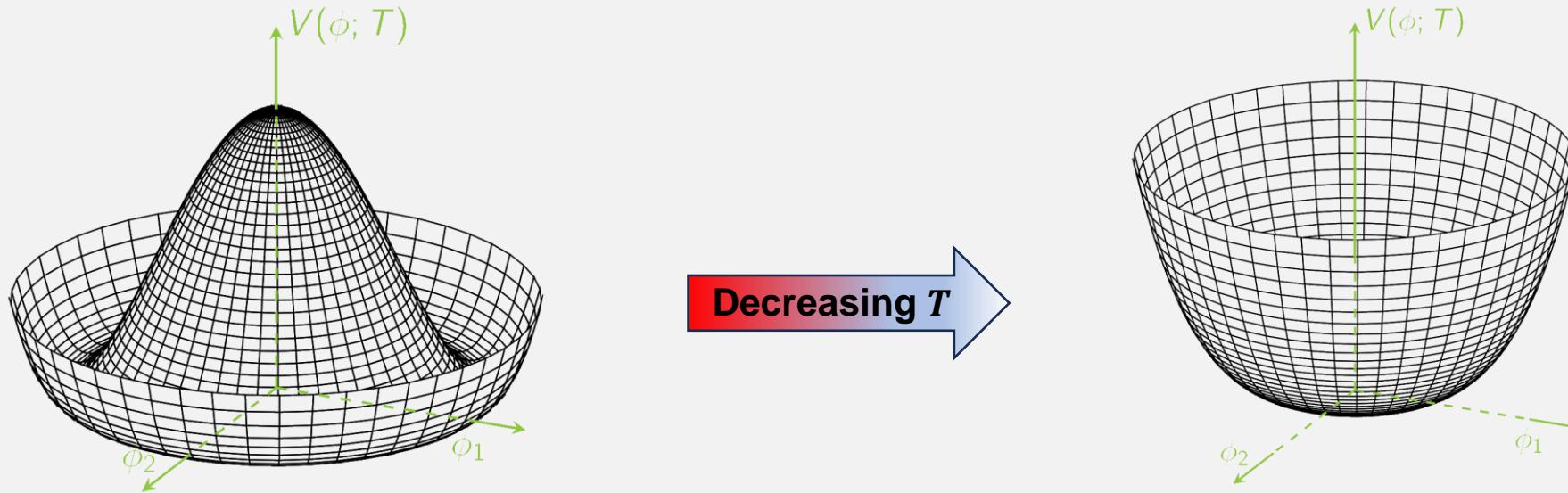


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with  $f$  antisymmetric and  $g$  symmetric.
  - neutrino masses:  $M_\nu \sim \mu f Y_l g^T Y_l f^T$
- **Majorana mass terms** for  $L_L$  and  $l_R$  are interesting.



# Inverse high $T$ symmetry breaking

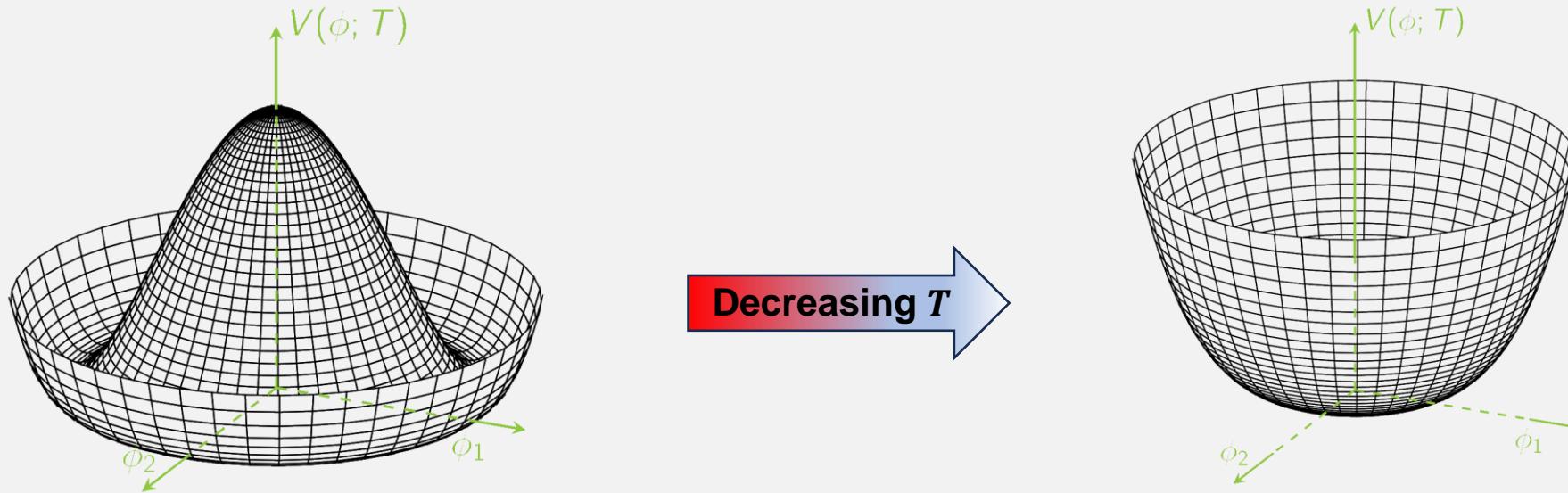


**Fig. 2:** Multi-scalar potential under high- $T$  symmetry breaking and low- $T$  symmetry restoration.

- To leading order in temperature  $T$  and 1-loop:  $V^T_{1l} \sim \sum_{i,j \neq i} [(m_i^2 + \textcolor{red}{C}_{\phi_i} T^2) \phi_i^2 + \lambda_i \phi_i^4 + \lambda_{ij} \phi_i^2 \phi_j^2]$
- For  $\textcolor{red}{C}_{\phi_k} < 0 \rightarrow \langle \phi_k \rangle \neq 0$  at high enough  $T$
- If  $\phi_k$  carries hypercharge: **SSB of  $U(1)_Y$**
- Caveat:  $\textcolor{red}{C}_{\phi_k} < 0$  difficult



# Inverse high $T$ symmetry breaking

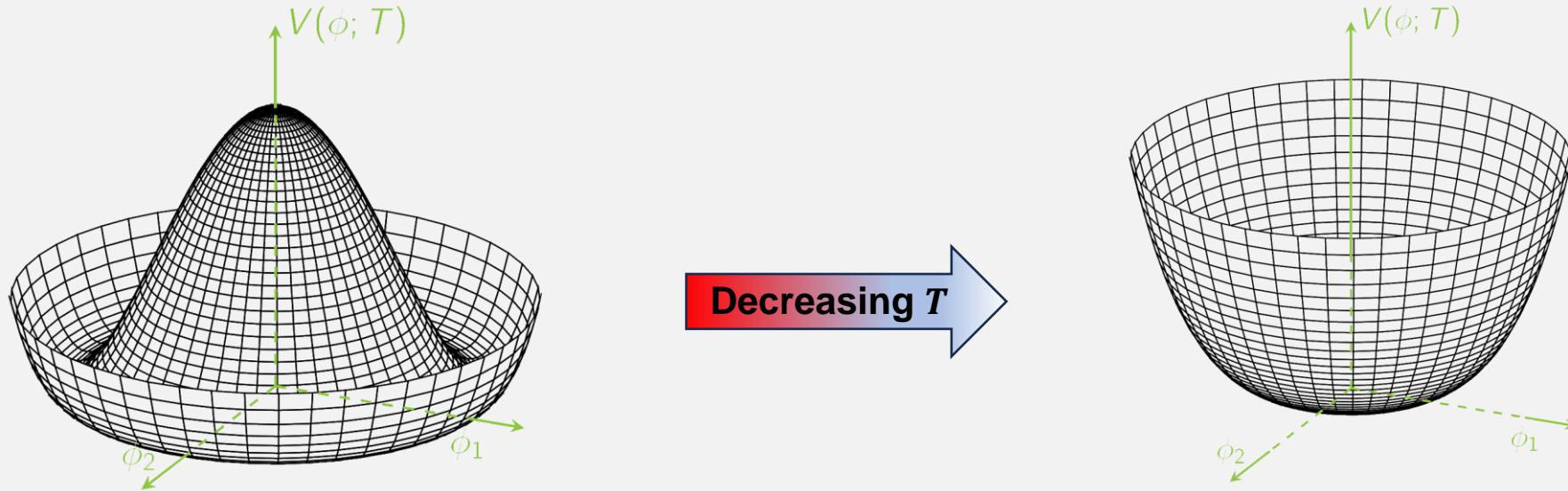


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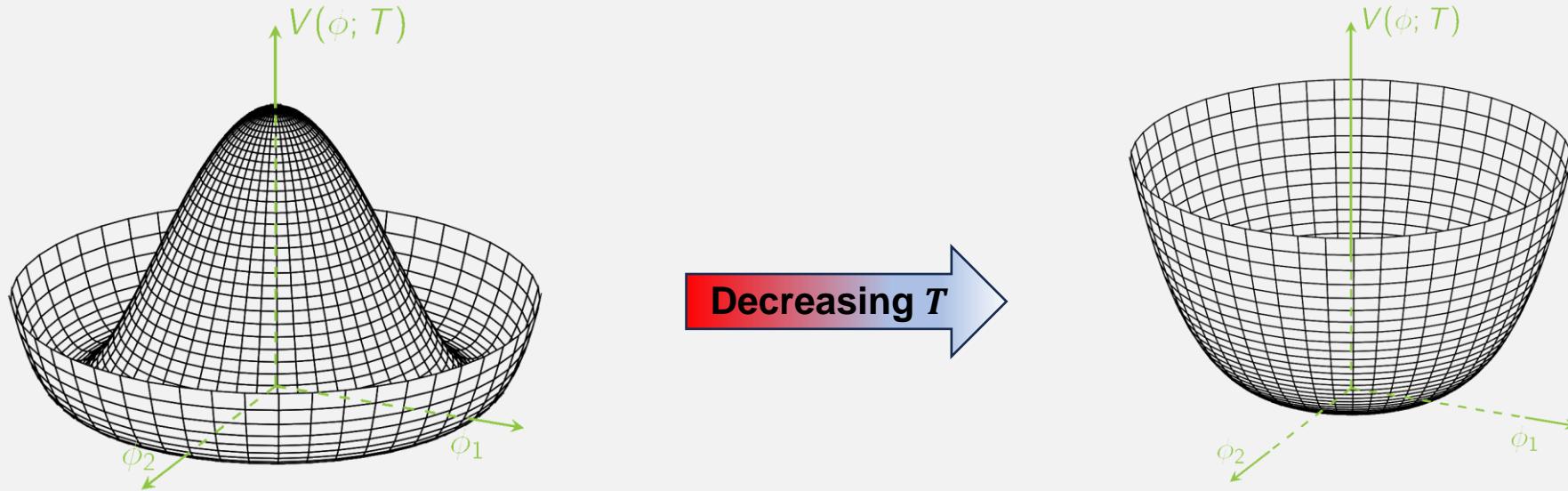


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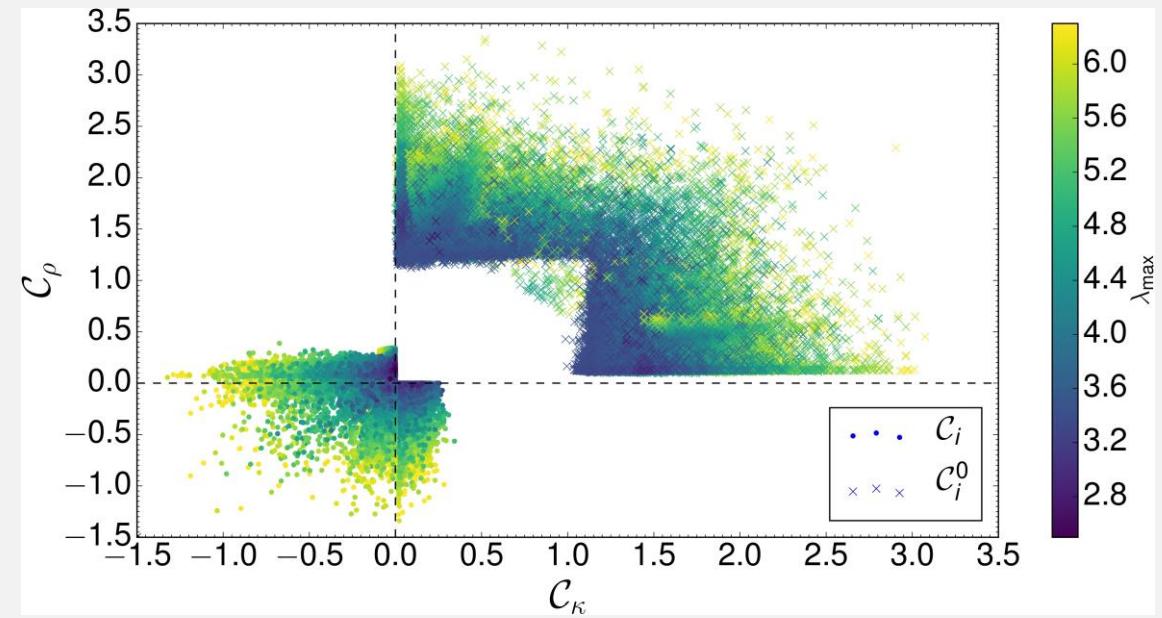


# High $T$ SSB of $U(1)_Y$ in ZB?

$$C_h = \frac{g'^2 + 3g^2}{32} + \frac{y_t^2}{8} + \frac{\lambda_h}{4} + \frac{\lambda_{h\kappa} + \lambda_{h\rho}}{24} + \Delta C_h$$

$$C_\kappa = \frac{g'^2}{4} + \frac{4\lambda_\kappa + 2\lambda_{h\kappa} + \lambda_{\kappa\rho}}{12} + \frac{1}{24} \sum_{j \neq k} |f_{jk}|^2 + \Delta C_\kappa$$

$$C_\rho = g'^2 + \frac{4\lambda_\rho + 2\lambda_{h\rho} + \lambda_{\kappa\rho}}{12} + \frac{1}{24} \sum_{j,k} |g_{jk}|^2 + \Delta C_\rho$$

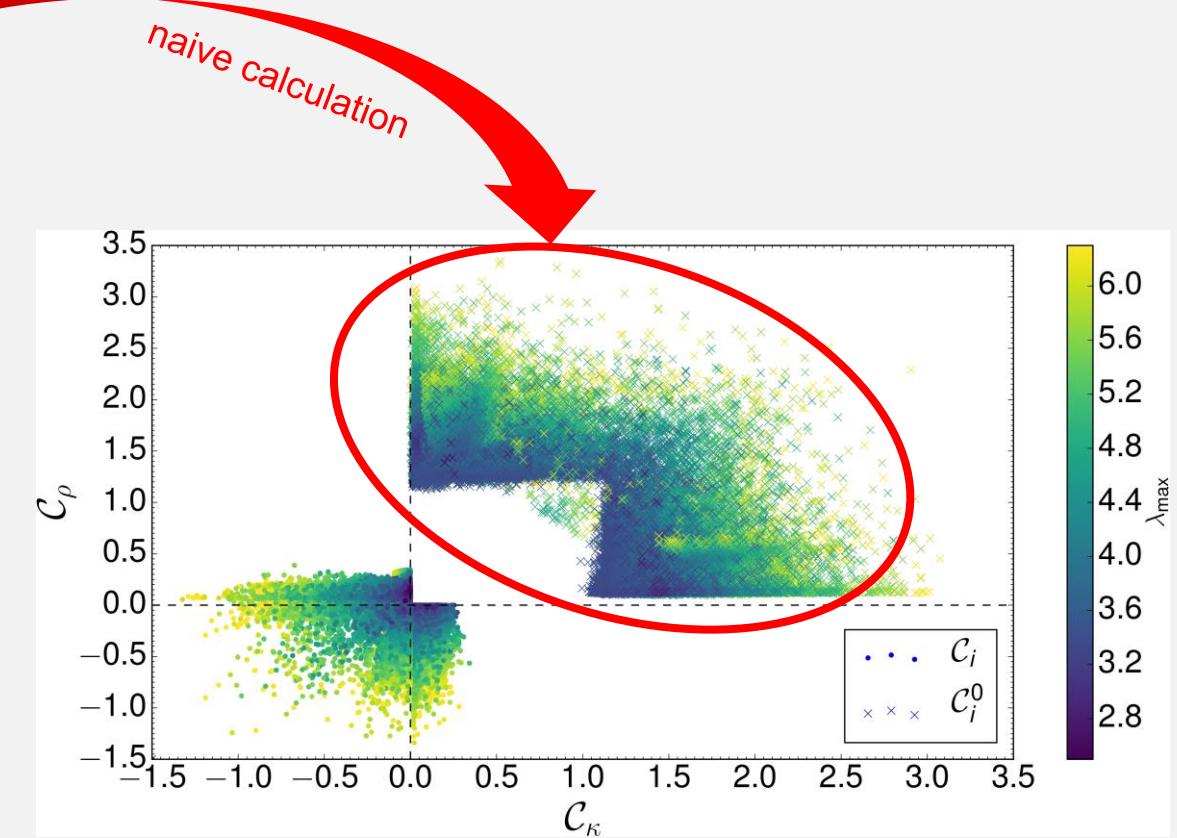


**Fig. 3:** Values of  $C_\kappa$ ,  $C_\rho$  (circles) and  $C_\kappa^0$ ,  $C_\rho^0$  (crosses) for a parameter scan of the quartic couplings in the ZB model. The color scheme encodes  $\lambda_{max} = \max\{\lambda_i\}$ . Both sets correspond to the same  $\lambda_i$ .



# High $T$ SSB of $U(1)_Y$ in ZB?

$$\begin{aligned} C_h &= \frac{g'^2 + 3g^2}{32} + \frac{y_t^2}{8} + \frac{\lambda_h}{4} + \frac{\lambda_{h\kappa} + \lambda_{h\rho}}{24} + \Delta C_h \\ C_\kappa &= \frac{g'^2}{4} + \frac{4\lambda_\kappa + 2\lambda_{h\kappa} + \lambda_{\kappa\rho}}{12} + \frac{1}{24} \sum_{j \neq k} |f_{jk}|^2 + \Delta C_\kappa \\ C_\rho &= g'^2 + \frac{4\lambda_\rho + 2\lambda_{h\rho} + \lambda_{\kappa\rho}}{12} + \frac{1}{24} \sum_{j,k} |g_{jk}|^2 + \Delta C_\rho \end{aligned}$$

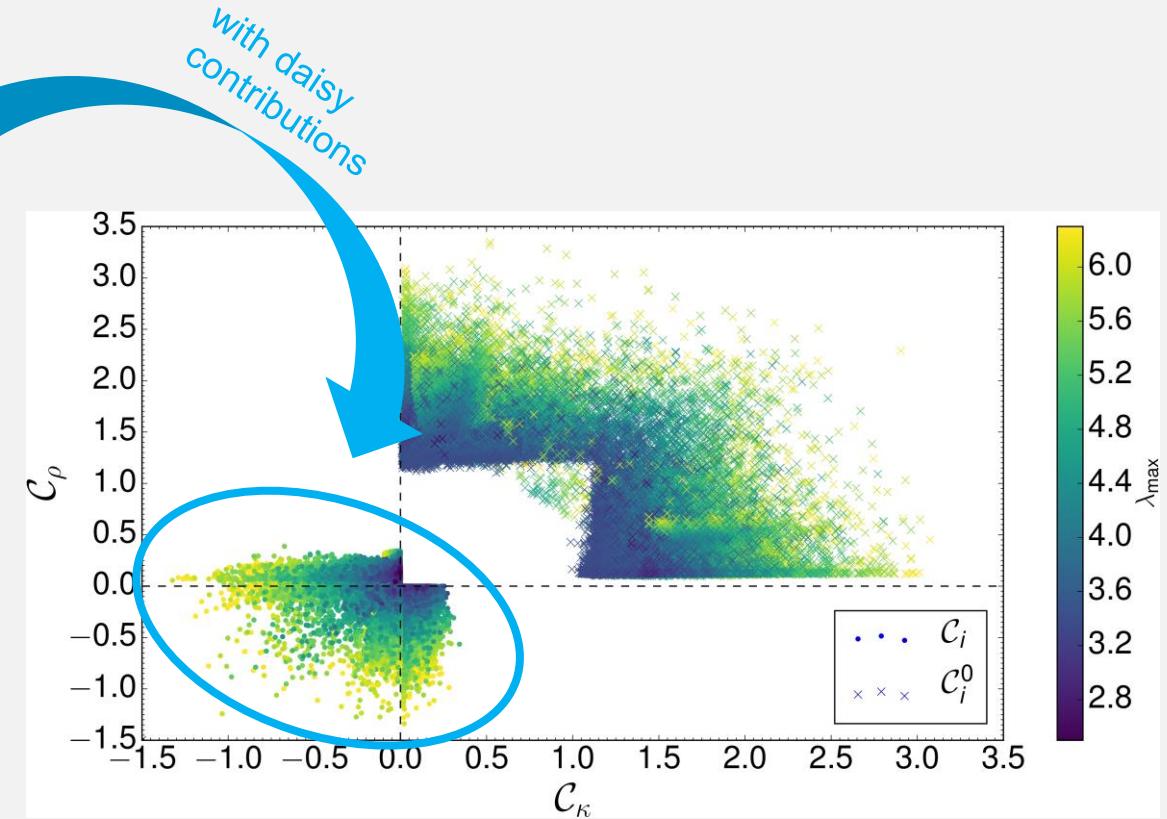


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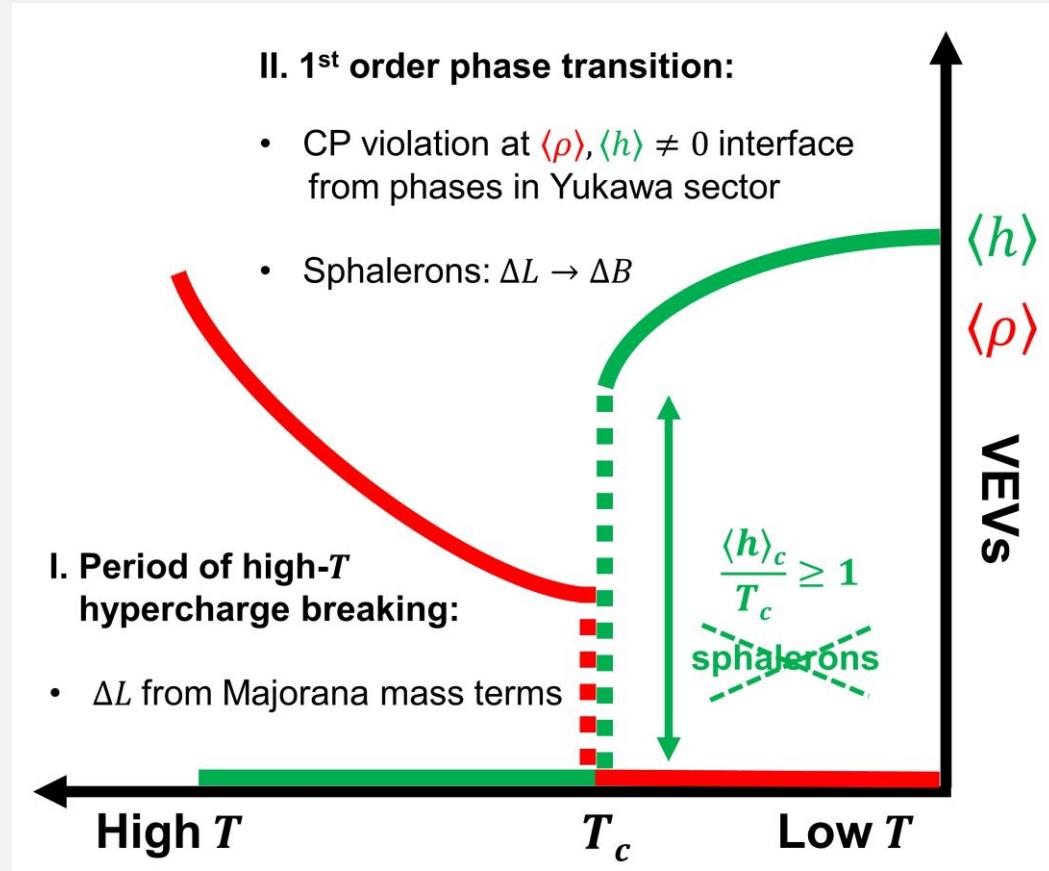
$$\begin{aligned} C_h &= \frac{g'^2 + 3g^2}{32} + \frac{y_t^2}{8} + \frac{\lambda_h}{4} + \frac{\lambda_{h\kappa} + \lambda_{h\rho}}{24} + \Delta C_h \\ C_\kappa &= \frac{g'^2}{4} + \frac{4\lambda_\kappa + 2\lambda_{h\kappa} + \lambda_{\kappa\rho}}{12} + \frac{1}{24} \sum_{j \neq k} |f_{jk}|^2 + \Delta C_\kappa \\ C_\rho &= g'^2 + \frac{4\lambda_\rho + 2\lambda_{h\rho} + \lambda_{\kappa\rho}}{12} + \frac{1}{24} \sum_{j,k} |g_{jk}|^2 + \Delta C_\rho \end{aligned}$$



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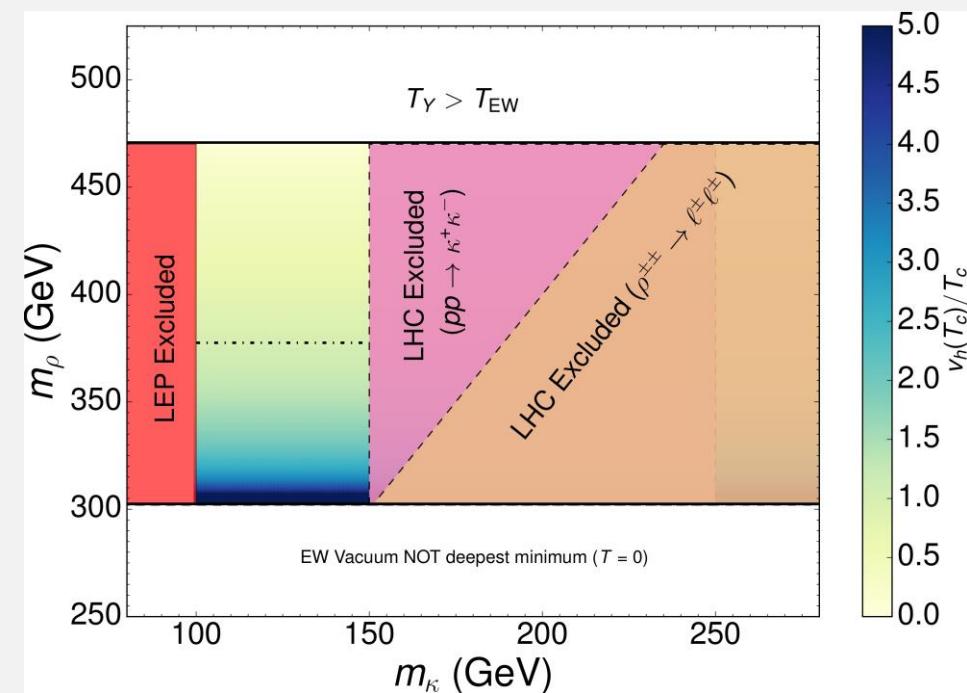


# Baryogenesis



**Table 1: Benchmarks**

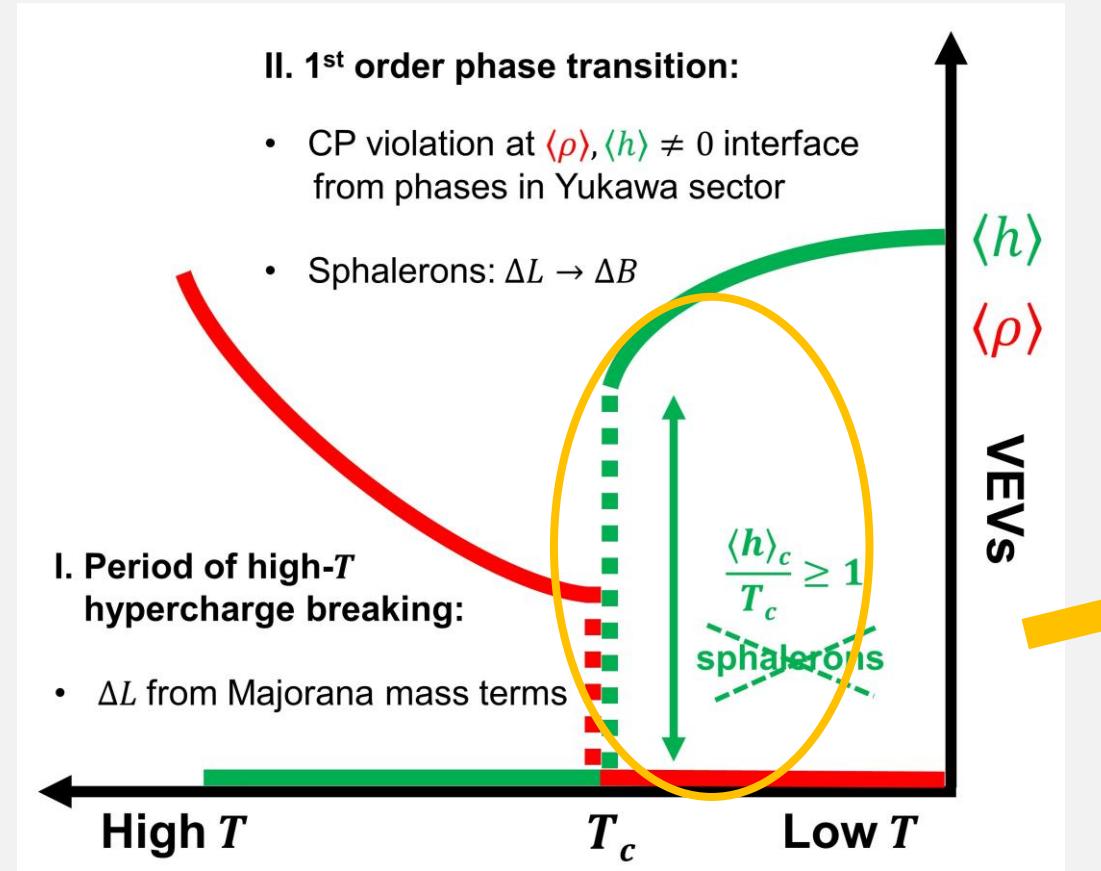
	$\lambda_\kappa$	$\lambda_\rho$	$\lambda_{h\kappa}$	$\lambda_{h\rho}$	$\lambda_{\kappa\rho}$	$C_h$	$C_\kappa$	$C_\rho$
BM1	0.118	5.44	4.70	-0.097	-0.052	0.042	0.048	-0.85
BM2	4.394	0.811	5.134	-0.537	-0.142	0.048	-0.529	0.192



**Fig. 4:** Viable masses ( $m_\kappa, m_\rho$ ) for baryogenesis in BM1.

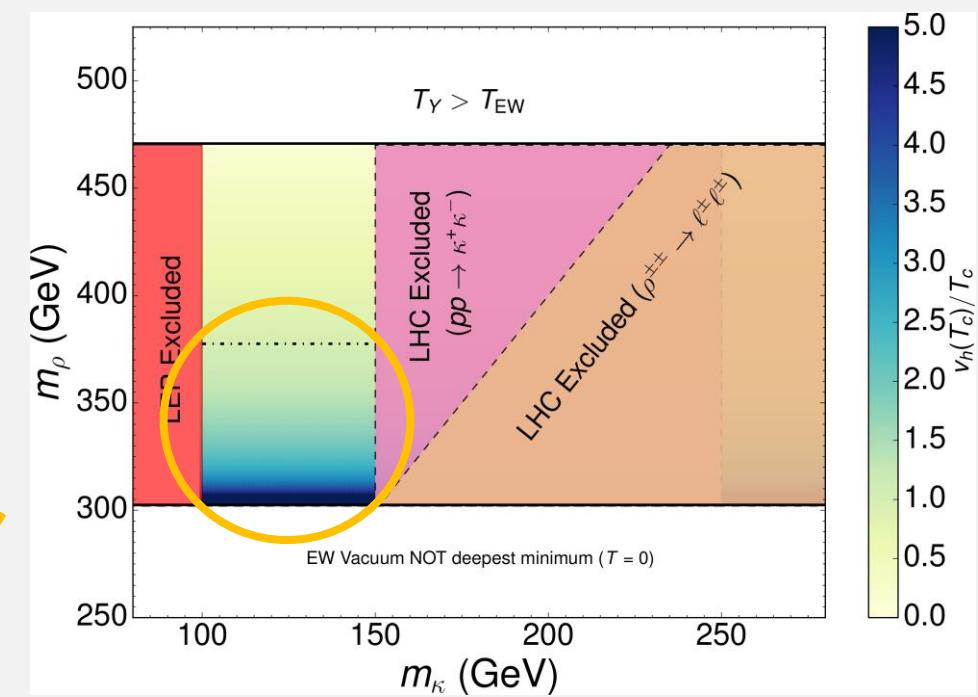


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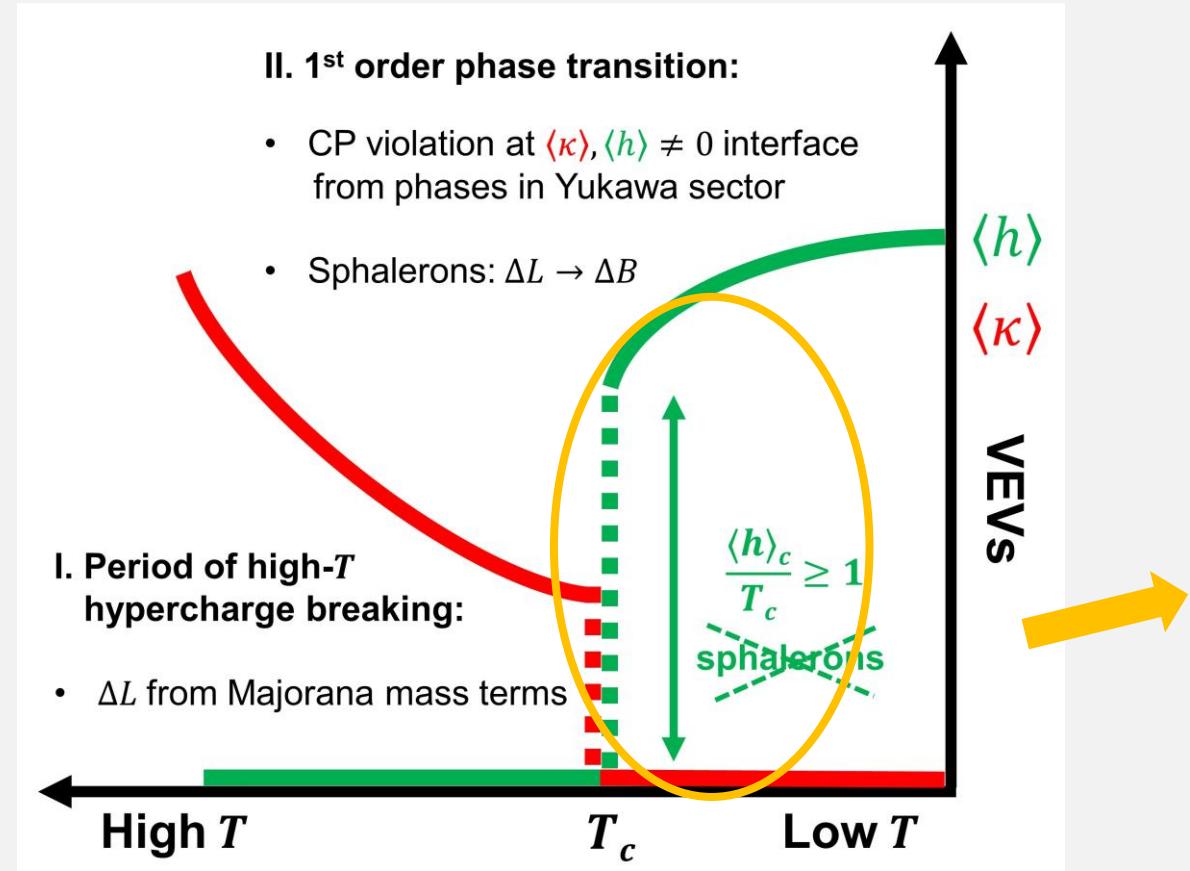
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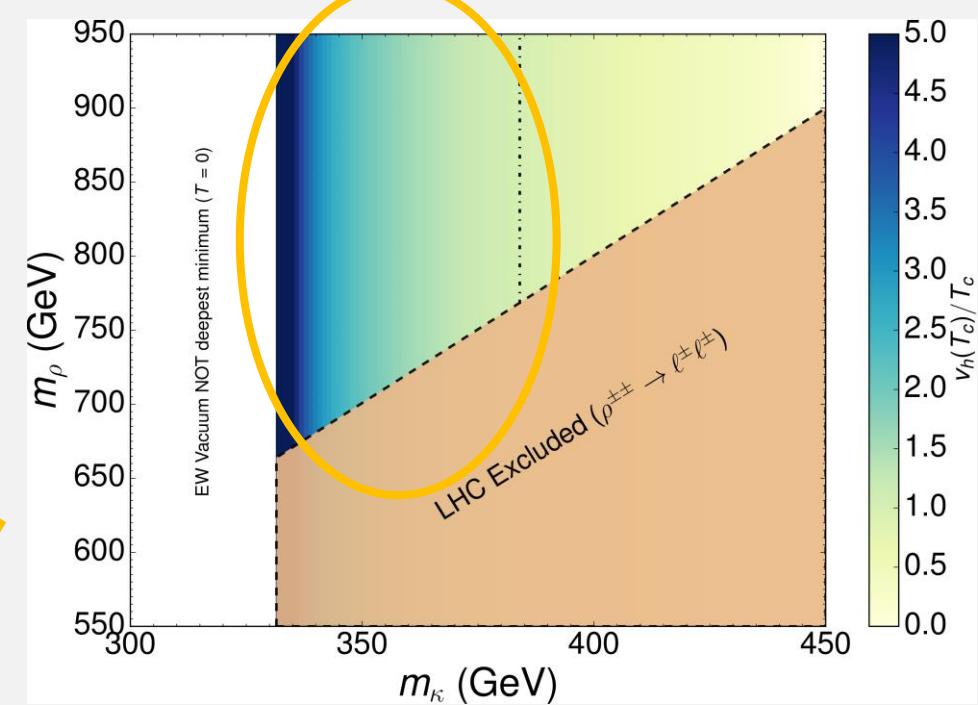


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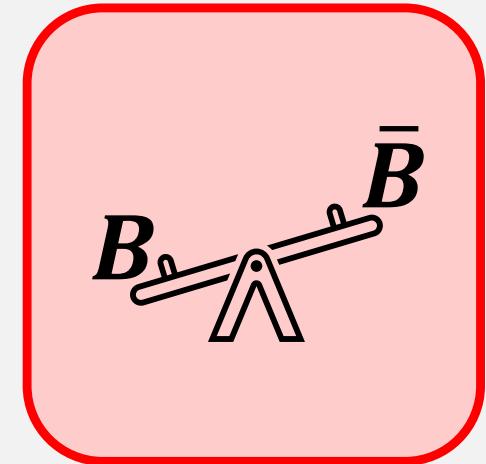
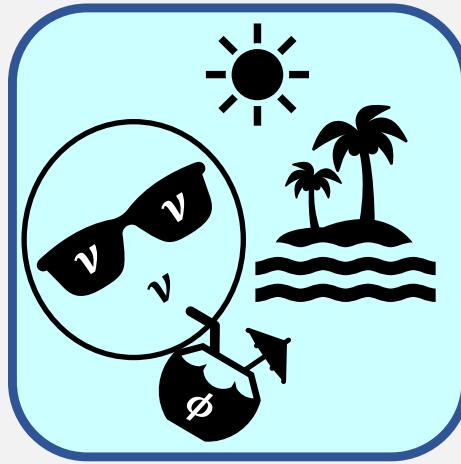
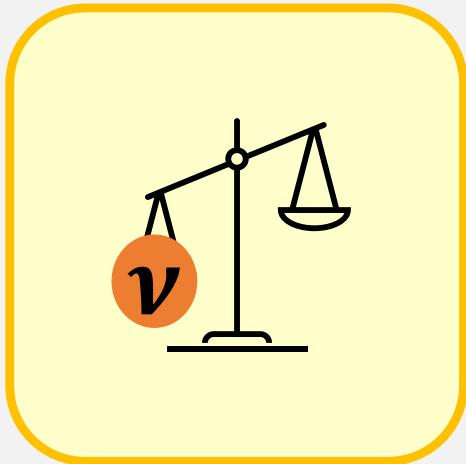
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**Fig. 5:** Viable masses ( $m_\kappa, m_\rho$ ) for baryogenesis in BM2.

# Takeaways



∅

- ZB possesses the **ingredients** for exotic early Universe phenomenology.
- How **generic to radiative models** are these features?

**Stay tuned!**

# Early Universe hypercharge breaking and neutrino mass generation

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 Universidad Autónoma de Madrid, Cantoblanco, 28049, Madrid, Spain

Contact: alvaro.lozano.onrubia@csic.es Based on arXiv:2308.05206 [hep-ph]

## 1. Motivation

- We still do not know how neutrinos acquire mass.
- Radiative models of neutrino mass are very popular.
- They typically rely on extended scalar sectors which we exploit here for exotic phenomenology.

## 2. Zee-Babu model

- The **Zee-Babu** (ZB) model is a canonical and well-studied model that generates neutrino masses at 2-loop [1, 2].

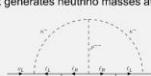


Fig. 1: Feynman diagram generating the neutrino masses in the ZB model.

## 4. Early Universe $U(1)_Y$ breaking

$$C_3 = \frac{g^2 + 3g_Y^2}{32} + \frac{y_\tau^2}{6} - \frac{\lambda_{\text{hyp}}}{4} + \frac{\lambda_{\text{hyp}} + \lambda_{\text{hyp}}}{24} + \Delta C_h$$

$$C_4 = \frac{g^2 + 4\lambda_{\text{hyp}} + 2\lambda_{\text{hyp}} + \lambda_{\text{hyp}}}{4} + \frac{1}{12} \sum_{k=1}^4 |f_k| + \Delta C_h$$

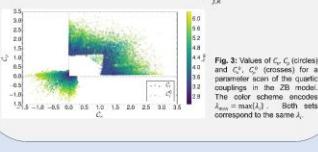
$$C_6 = g^2 + \frac{4\lambda_{\text{hyp}} + 2\lambda_{\text{hyp}} + \lambda_{\text{hyp}}}{12} + \frac{1}{24} \sum_{k=1}^6 |f_k|^2 + \Delta C_h$$


Fig. 3: Values of  $C_3$ ,  $C_4$  (circles), and  $C_6$  (triangles) for a parameter scan of the quartic couplings in the ZB model. The axes are  $\lambda_{\text{hyp}}$  and  $\lambda_{\text{hyp}}$ .  $\lambda_{\text{hyp}} = \max(\lambda_i)$ . Both sets correspond to the same  $\lambda_{\text{hyp}}$ .

## 5. Baryogenesis

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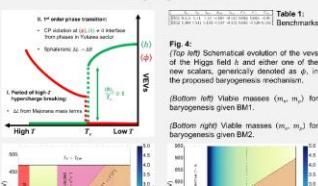


Fig. 4: (left) Schematic evolution of the vevs of the Higgs field  $\langle h \rangle$  and either one of the new scalars, generically denoted as  $\phi$ , in the proposed baryogenesis mechanism. (bottom left) Viable masses ( $m_\phi$ ) for baryogenesis given BHAT. (bottom right) Viable masses ( $m_\phi$ ) for baryogenesis given BLT.

$m_\phi$ (GeV)	$m_h$ (GeV)	$\langle h \rangle$ (GeV)	$\langle \phi \rangle$ (GeV)
100	100	~100	~100
200	200	~200	~200
300	300	~300	~300
400	400	~400	~400

## 3. Inverse symmetry breaking

- Thermal effects significantly modify scalar potentials [3].
- To leading order in temperature and 1-loop:

$$V_{1L} \sim \sum_{i,j \neq i} \left[ (m_i^2 + C_{ij} T^2) \phi_i^2 + \lambda_i \phi_i^4 + \lambda_{ij} \phi_i^2 \phi_j^2 \right]$$

- For  $C_{ij} < 0$  a non-vanishing vacuum expectation value (vev)  $\langle \phi_i \rangle$  is induced at high temperature.
- If  $\phi_i$  carries hypercharge, the vev breaks  $U(1)_Y$  spontaneously
- Caveat:  $C_{ij}$  depend on constrained  $\lambda_i$ ,  $C_{ij} < 0$  difficult.



Fig. 2: Scalar potential under high- $T$  symmetry breaking and low- $T$  symmetry restoration.

## 6. Conclusions

- The ZB model possesses the ingredients for very exotic early Universe phenomenology.
- This looks like a more general feature of radiative neutrino mass models – stay tuned for further work!

## 7. References

- [1] A. Zee, Quantum Numbers of Majorana Neutrino Masses, *Nucl. Phys. B* 264, 99 (1986).
- [2] K. S. Babu, Model of Calculable Majorana Neutrino Masses, *Phys. Lett. B* 203, 132 (1988).
- [3] B. Weinberg, Gauge and Global Symmetries at High Temperature, *Phys. Rev. D* 9, 3357 (1974).

# Thank you!

Contact: alvaro.lozano.onrubia@csic.es

# References

- [1] A. Zee, Quantum Numbers of Majorana Neutrino Masses, Nucl. Phys. B 264, 99 (1986).
- [2] K. S. Babu, Model of 'Calculable' Majorana Neutrino Masses, Phys. Lett. B 203, 132 (1988).