### Searches for long-lived particles at hadron colliders and beyond

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- Mass scale of new physics particles is unknown
- One of the possibilities:  $m_{\rm new \ physics} \ll m_{\rm EW} \Rightarrow$  may be copiously produced at accelerators
- Large couplings ( $\equiv$  small lifetimes) of such particles are often excluded by past experiments  $\Rightarrow$  Long-Lived Particles (LLPs)

LLP	Dark photons $V$	Dark scalars $S$	$\mathbf{HNLs}\ N$	$\mathbf{ALPs} \ a$
$\mathcal{L}$	$-\epsilon F_{\mu\nu}V^{\mu\nu}$	$H^{\dagger}HS$	$Y_l \bar{L} H^c N$	$\frac{a}{f}F_{\mu\nu}\tilde{F}^{\mu\nu}+\ldots$

### Current LHC reach for LLPs I



- FIPs are unstable  $\Rightarrow$  searches for displaced vertices
- Ongoing searches for  $\mathcal{O}(1 \text{ GeV})$  LLPs at the LHC:
  - 1. Tag over production vertex (associated high- $p_T$  particles)
  - 2. Decay vertex within inner trackers

[2201.05578], [2204.11988]

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August 30, 2023 3/27

### Current LHC reach for LLPs II



- FIPs are unstable  $\Rightarrow$  searches for displaced vertices
- Ongoing displaced vertex searches for  $\mathcal{O}(1 \text{ GeV})$  LLPs at the LHC:
  - 1. Tag over production vertex (associated high- $p_T$  particles)
  - 2. Decay vertex within inner trackers

[2201.05578], [2204.11988]

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## Current LHC reach for LLPs III

- Problems of these searches:
  - 1. High- $p_T$  associated primary particle(s): uncommon for light LLPs
  - 2. Small size of inner trackers  $\Rightarrow$  small  $P_{\text{decay}} (c\tau_{\text{LLP}} \propto m_{\text{LLP}}^{-n})$
  - 3. Numerous background  $\Rightarrow$  severe selection



August 30, 2023 5/27

315

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### Current LHC reach for LLPs IV

- Problems of these searches:
  - 1. High- $p_T$  associated primary particle: uncommon for light LLPs
  - 2. Small size of inner trackers  $\Rightarrow$ small  $P_{\text{decay}}$   $(c\tau_{\text{LLP}} \propto m_{\text{LLP}}^{-n})$
  - 3. Numerous background  $\Rightarrow$  severe selection



### We need to go beyond the current searches

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August 30, 2023 6/27

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### Improving the reach of colliders I



- 1. Going beyond inner trackers
- 2. Triggering without tagging over production vertex (ATLAS/CMS: since 2029. LHCb: currently)

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August 30, 2023 7/27

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## Improving the reach of colliders II

### Method 1

Muons from outer space at muon stations

- Decay volume: all space  $|\eta| < 2.5$  between the IP and muon stations
- Limitations: no full kinematics reconstruction, only  $\mu\mu$  final state



1903.11918

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August 30, 2023 8/27

315

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## Improving the reach of colliders III

### Method 2

Non-muonic showers in muon chambers

- Current CMS analysis: triggering over high- $p_T$  prompt leptons Study of the potential of track-triggers: in preparation
- **Limitations**: no full kinematics reconstruction



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Muon detector

August 30, 2023 9/27

Improving the reach of colliders IV

### Method 3

T tracks at LHCb

- Decay volume: LHCb detector  $2 < \eta < 5$  until SciFi
- **Limitations**: low luminosity



### $In\ preparation$

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August 30, 2023 10/27

## Improving the reach of colliders V

- Preliminary studies: good potential to explore light LLPs with  $m_{\rm LLP} \lesssim m_B$ 

In preparation



### Limitations are still there

Dealing with background, finite reconstruction efficiency, ...

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August 30, 2023 11/27

315

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## Going beyond main LHC experiments I

## Workaround

### Going beyond ATLAS/CMS/LHCb

- A displaced decay volume
  - 1. Collider-based
  - 2. Dedicated (beam dump)
- Advantages:
  - Far from IP  $\Rightarrow$  low (zero) background without selection
  - Long decay volume



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August 30, 2023 12/27

Comparison of the BC experiments in nutshell I



[1901.09966]

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August 30, 2023 13/27

## Comparison of the BC experiments in nutshell II



Experiments differ in:

- Placement relative to the beam axis
- Detector equipment
- Decay volume geometry

### Need to choose among them (budget limitations!!)

How?

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August 30, 2023 14/27

Comparison of the BC experiments in nutshell III

### Exclusion potential argument

The event rate should be as large as possible to exclude the parameter space

The event rate:

$$N_{\text{events}} \propto N_{\text{prod}} \times P_{\text{decay}} \propto N_{\text{pp}} \cdot \chi_{\text{mother}}(s_{\text{pp}}) \cdot \epsilon_{\text{geom}} \times \Delta z_{\text{fid}} \langle \gamma_{\text{LLP}}^{-1} \rangle$$
 (1)

 $N_{pp}$ : number of protons.  $\chi_{mother}$ : rate of mother process per pp.  $\epsilon_{geom}$ : fraction of LLPs pointing to the detector.  $\Delta z_{fid}$ : length of the decay volume Two LLP categories (the spread of  $df/d\Omega_{LLP}$ ):

- LLPs produced from heavy flavors (B, D)

- LLPs produced in EM processes, mixing with/decays of light mesons, bremsstrahlung

August 30, 2023 15/27

Comparison of the BC experiments in nutshell IV

### Discovery potential argument

The experiment has to be equipped with a full detector system to reconstruct LLP's properties

- Measuring LLP's mass and spin
- Measuring its decay modes
- Establishing LLP's relation to the resolution of the BSM problems

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August 30, 2023 16/27

Comparison of the BC experiments in nutshell V



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August 30, 2023 17/27

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Comparison of the BC experiments in nutshell VI



### Beam dump:

- SHiP: very large  $\epsilon_{\text{geom}}$ , medium  $\langle \gamma^{-1} \rangle$ , full detector
- Smaller  $\chi_{
  m mother}$  for heavy flavors is compensated by a very large  $N_{
  m PoT}$
- **SHADOWS/HIKE**: smaller  $N_{PoT}$ +geometric limitations

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August 30, 2023 18/27

## Comparison of the BC experiments: a detailed study I

- A more detailed analysis: compare sensitivity curves - regions of exclusion/discovery

### Sensitivities comparison in a perfect world:

- 1. The same LLP phenomenology input Phenomenology is revising continuously
- 2. Control over numerical artifacts Complicated to find errors in pure MC generators
- 3. Transparency and publicity of the sensitivity calculations Sensitivities obtained by collaborations are black-box sourced

August 30, 2023 19/27

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Comparison of the BC experiments: a detailed study II

- Semi-analytic estimates:

$$N_{\rm ev} \approx N_{\rm prod} \times \epsilon_{\rm LLP} \times \langle P_{\rm dec} \rangle \times \epsilon_{\rm dec}$$
 (2)

Each factor may be qualitatively estimated [1902.06240]

– Improved version ( $\boldsymbol{z}$  - LLP long. displacement,  $\boldsymbol{\theta}$  - LLP polar angle,  $\boldsymbol{\phi}$  - az. angle):

$$N_{\rm ev} = \sum_{i} N_{\rm prod}^{(i)} \int dE d\theta dz \ f^{(i)}(\theta, E) \cdot \epsilon_{\rm az}(\theta, z) \cdot \frac{dP_{\rm dec}}{dz} \cdot \epsilon_{\rm dec}(m, \theta, E, z) \cdot \epsilon_{\rm rec} \tag{3}$$

•  $N_{\text{prod}}^{(i)}, f^{(i)}(\theta, E)$  are the total number of produced LLPs and the angle-energy distribution for the given channel i

•  $\epsilon_{\rm az}$  is the azimuthal acceptance for the LLP to decay inside the decay volume

- $\frac{dP_{\text{dec}}}{dz} = \frac{\exp[-z/(\cos(\theta)c\tau\sqrt{\gamma^2-1})]}{\cos(\theta)c\tau\sqrt{\gamma^2-1}}$  is the differential decay probability for the LLP to decay
- $\epsilon_{dec}$  is the decay products acceptance
- $\epsilon_{\rm rec}$  (may be computed externally) is the reconstruction efficiency

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August 30, 2023 20/27

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### Comparison of the BC experiments: a detailed study III





- SensCalc a Mathematica-based sensitivity evaluator
- **Input**: model description (production, decays), experimental setup (geometry, selection cuts), the tabulated distributions of mother particles
- **Output**: tabulated number of events  $N_{\text{events}}(m_{\text{LLP}}, g_{\text{LLP-SM}})$  that may be converted into exclusion/discovery limits

[2305.13383]

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August 30, 2023 21/27

Comparison of the BC experiments: a detailed study IV



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August 30, 2023 22/27

Comparison of the BC experiments: a detailed study V



SHiP: a perfect balance between exclusion and discovery potentials

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August 30, 2023 23/27

## Reconstructing LLPs at SHiP: HNL example I



- SHiP experiment: despite limitations of the ECN3 hall, already has the optimal placement [2305.13383]
- Decision about SHiP: will be made by the end of this year

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### Reconstructing LLPs at SHiP: HNL example II



[1912.05520]

- Standard properties: mass, decay modes, spin
- By measuring the angular distribution of decay products, it is possible to check the nature of HNLs and resolve  $N \bar{N}$  oscillations

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August 30, 2023 25/27

## Reconstructing LLPs at SHiP: HNL example III

 By measuring the decay modes, it is possible to recover the mixing pattern

$$U_e^2 : U_{\mu}^2 : U_{\tau}^2 \tag{4}$$

 This allows us to check whether HNLs are consistent with IH/NH hypothesis



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August 30, 2023 26/27

- We need to go beyond LHC to explore light LLPs
- Comparison of these beyond-collider experiments: unified way was missing, SensCalc provides it
- Detailed study: SHiP is optimal for probing LLPs with  $m \lesssim 5~{\rm GeV}$

# **Backup** slides

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August 30, 2023 1/8

## Validation: dark scalars at MATHUSLA and SHADOWS I



– Setups: taken from the SHADOWS LoI and MATHUSLA Snowmass paper

- Minimal event requirements: scalars must decay inside the decay volume, decay products have to point to the end of the detector
- SensCalc predictions cross-checked with a dedicated simulation under the same input

Image: A (1)

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## Validation: dark scalars at MATHUSLA and SHADOWS II



- The sensitivities obtained by SHADOWS and MATHUSLA people: a huge difference

- **Reason 1**: the setups used in the collab. estimates do not match the setups described publicly:  $\epsilon_{dec} = 1$  for MATHUSLA, a larger decay volume (without clearly studied background status) for SHADOWS
- Reason 2: different description of the scalar production

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## Validation: SHiP sensitivity I



- SensCalc predictions agree with FairShip simulations for the ECN4 setup from [1811.00930], [2011.05115]
- Differences: different phenomenology, simplification for the upper bound calculation in [1811.00930]

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August 30, 2023 4/8

315

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 New physics searches at LHCb using new downstream tracking algorithm (paper in preparation): acceptances perfectly agree with full LHCb simulations

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▲□ ▶ ◆□ ▶ ◆ ■ ▶ ◆ ■ ▶ ● ■ ● つへで August 30, 2023 5/8 LLPs 2022 proceedings, LLPs models BCXX:

- **BC4, BC5** (Higgs-like scalars): some (most) of the experiments use inclusive description of the scalar production, which is wrong at large masses  $m_S \gtrsim 3$  GeV
- **BC10** (ALPs coupled to fermions): most of the experiments use a completely wrong phenomenology (missing important production and decay channels)<sup>1</sup>; some of them include hadronic width, while most of them don't
- BC11 (ALPs coupled to gluons): different definition of the coupling to SM particles is used by different experiments

<sup>&</sup>lt;sup>1</sup>M. Ovchynnikov et al., in progress

## Calculating acceptances I

- 1. Provide geometry input (decay volume, detector)
- 2. Find a grid  $\theta_{\text{LLP}}(z_{\text{LLP}}), \phi_{\text{LLP}}(z_{\text{LLP}})$  for which the LLP is inside the decay volume  $\Rightarrow \epsilon_{\text{az}}(\theta_{\text{LLP}}, z_{\text{LLP}})$ Simple verifications:
  - checking  $\theta_{\min/\max}$  belonging to the decay volume
  - visualization of the points  $\{z_{LLP}, \theta_{LLP}, \phi_{LLP}\}$  – they must belong to the decay volume
  - the integral of ε<sub>az</sub> gives the total volume of the decay volume



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August 30, 2023 7/8

- 3. Consider a grid  $m_{\text{LLP}}, E_{\text{LLP}}, \theta_{\text{LLP}}, z_{\text{LLP}}, \phi_{\text{LLP}}$ , where  $\{\theta, \phi\}$  belong to the decay volume
- 4. Generate phase space of the decay products at rest and boost them given  $E_{\rm LLP}, \theta_{\rm LLP}, \phi_{\rm LLP}$ 
  - Exclusive decay channels (where analytic expression for the matrix element exists): simulating phase space with the weight given by squared matrix element
  - Inclusive decays (into jets/jets+leptons): either simulating decay into jets in Mathematica, or using pre-computed phase space of a typical hadronized final state obtained using pythia
- 5. Require at least two decay products with zero total charge to point to the end of the detector (may be changed to requiring all the decay products to be within the acceptance). Additionally, require some other cuts if needed (the energy cut, the  $p_T$  cut, etc.). This gives  $\epsilon_{\text{dec}}$

Averaging  $\epsilon_{dec}$  over  $\phi$ : reasonable since other quantities (such as  $dP_{dec}/dz$ ,  $f_{LLP}$ ) are  $\phi$ -independent