

ALPs, the on-shell way

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ALPs = Axion-like particles

- Exact/Approximate shift-symmetry:

$$\phi(x) \rightarrow \phi(x) + \text{cnst} \quad \Rightarrow \quad \text{Physical amplitudes}$$

- Relies on Lagrangian formulation!

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- Relies on Lagrangian formulation!
- Invert the reasoning:

Properties of
amplitudes

$$\stackrel{?}{\Rightarrow} \quad \phi(x) \rightarrow \phi(x) + \text{cnst}$$

On-shell methods

- No fields and Lagrangians, only amplitudes.
- Example for ALP and 2 fermions:

$$\mathcal{A}[\phi\psi_1\bar{\psi}_2] = g_L \langle \mathbf{12} \rangle + g_R [\mathbf{12}]$$

Coefficients

Can depend on
kinematical invariants

Kinematical structures

Fixed by little-group covariance

Some of the progress
in on-shell methods:

- [Parke, Taylor, '86]
- [Bern et al, '95]
- [Britto, Cachazo, Feng, '05]
- [Britto et al, '05]
- [Zwiebel, '12]
- [Evang, Huang, '13]
- [Dixon, '13]
- [Plefka, Henn, '14]
- [Cheung et al, '14]
- [Wilhelm, '15]
- [Arkani-Hamed, Huang, Huang, '17]
- [Durieux et al, '19]
- [Falkowski, Isabella, Machado, '20]
- [Batarella et al, '20]
- [Elias Miró et al, '20]
- [Jiang et al, '20]
- [Balkin et al, '21]
- [Alves, Bertuzzo, GMS, '21]
- [Bonnefoy et al, '21]
- [Travaglini et al, '22]
- [GMS, '22]
- [Machado, Renner, Sutherland, '22]
- [Angelis, '22]

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A set of conditions

- Soft factorization condition:

$$\mathcal{A}_{n+1}[\phi] \xrightarrow{p_\phi \rightarrow 0} \mathcal{A}_n \times \mathcal{S} \text{ is regular}$$



ALP 3-point interactions

- Soft contact condition:

$$\lim_{p_\phi \rightarrow 0} \mathcal{A}_{\text{contact}} = 0$$



Higher-order interactions

- ALP-Higgs obstruction:

$$\lim_{p_\phi, p_H \rightarrow 0} \mathcal{A}[\phi H \dots] = \lim_{p_\phi \rightarrow 0} \lim_{\text{high-energy}} \frac{1}{v} \mathcal{A}[\phi \dots]$$



Connection to SM particle content

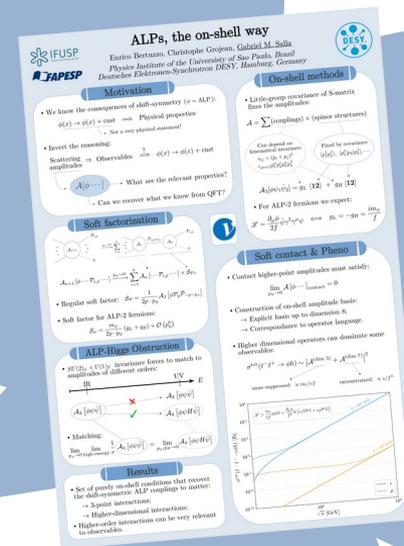
Summary

- Determined set of on-shell and physical conditions to recover shift-symmetry;
- Pheno application of higher-order interactions and more on the poster!

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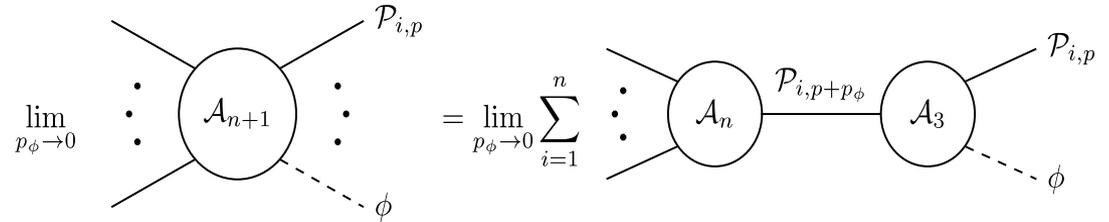
Thank you!



Backup

Soft factorization condition

- Soft limit of ALP in generic amplitude:



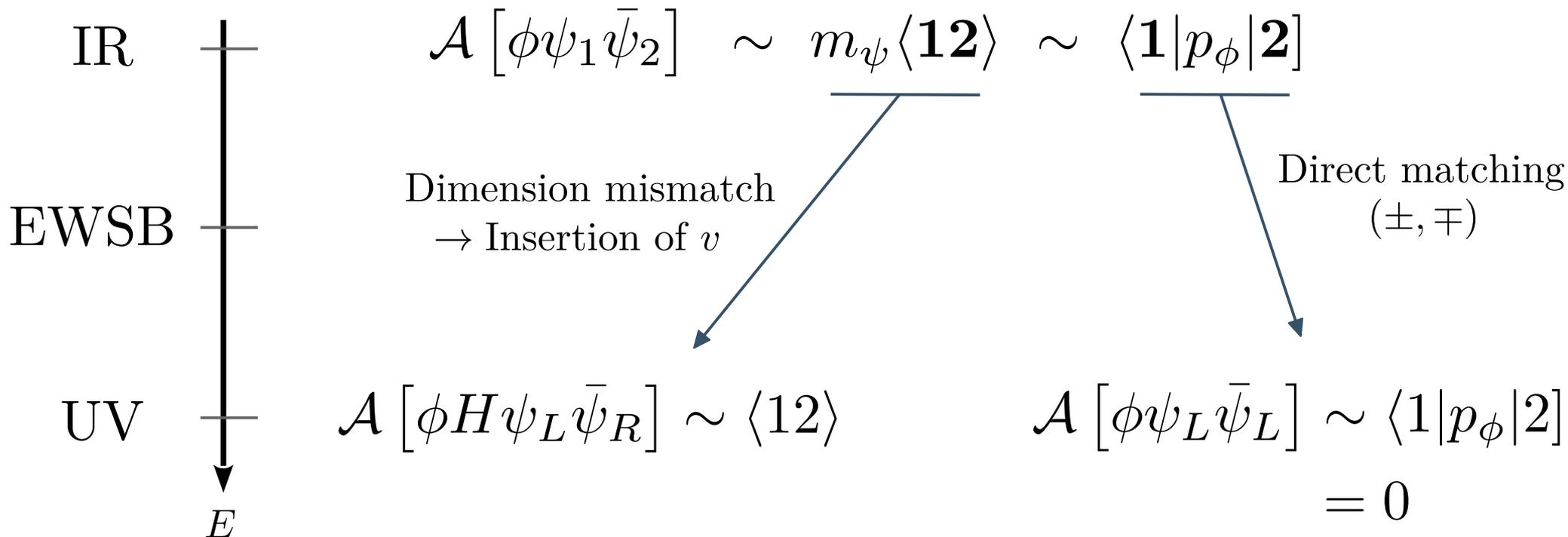
$$\lim_{p_\phi \rightarrow 0} \mathcal{A}_{n+1} \left[\phi \cdots \mathcal{P}_{i,p}^{I_1, \dots, I_{2s}} \cdots \right] = \lim_{p_\phi \rightarrow 0} \sum_{i=1}^n \mathcal{A}_n \left[\cdots \mathcal{P}_{i,p}^{K_1, \dots, K_{2s}} \cdots \right] \times (\mathcal{S}_{\mathcal{P}_i})_{K_1, \dots, K_{2s}}^{I_1, \dots, I_{2s}}$$

- Soft factor imposed to be regular:

$$(\mathcal{S}_{\mathcal{P}_i})_{K_1, \dots, K_{2s}}^{I_1, \dots, I_{2s}} = \frac{\epsilon_{K_1 J_1} \cdots \epsilon_{K_{2s} J_{2s}}}{2p \cdot p_\phi} \mathcal{A}_3 \left[\phi \mathcal{P}_{i,p}^{I_1, \dots, I_{2s}} \bar{\mathcal{P}}_{i,-p-p_\phi}^{J_1, \dots, J_{2s}} \right]$$

ALP-Higgs obstruction

- Soft limit of ALP in generic amplitude:



Determination of higher-order interactions

Dimension 6

Particle content	$\mathcal{A} \times f^2$	$\mathcal{O} \times f^2$
$\phi^2 H \bar{H}$	$p_{\phi_1} \cdot p_{\phi_2}$	$(\partial_\mu \phi)(\partial^\mu \phi) H ^2$

Dimension 7

Particle content	$\mathcal{A} \times f^3$	$\mathcal{O} \times f^3$
$\phi \psi_1 \bar{\psi}_2 H$	$(p_\phi \cdot p_1) \langle 12 \rangle$	$(\partial_\mu \phi) \bar{\psi} H (D^\mu \psi)$
	$(p_\phi \cdot p_2) \langle 12 \rangle$	$(\partial_\mu \phi) (D^\mu \bar{\psi}) H \psi$
$\phi \psi_1 \bar{\psi}_2 V_3$	$\langle 13 \rangle \langle 3 p_\phi 2 \rangle, \langle 23 \rangle \langle 3 p_\phi 1 \rangle$	$(\partial^\mu \phi) \bar{\psi} \gamma^\nu \psi V_{\mu\nu}$
		$(\partial^\mu \phi) \bar{\psi} \gamma^\nu \psi \tilde{V}_{\mu\nu}$
		$(\partial^\mu \phi) \bar{\psi} \gamma^\nu T^A \psi V_{\mu\nu}^A$ $(\partial^\mu \phi) \bar{\psi} \gamma^\nu T^A \psi \tilde{V}_{\mu\nu}^A$
$\phi H_1 \bar{H}_2 V_3$	$\langle 3 p_\phi(p_1 - p_2) 3 \rangle$	$(\partial_\mu \phi) (H^\dagger i \overleftrightarrow{D}_\nu H) V^{\mu\nu}$
		$(\partial_\mu \phi) (H^\dagger i \overleftrightarrow{D}_\nu H) \tilde{V}^{\mu\nu}$
		$(\partial_\mu \phi) (H^\dagger i \overleftrightarrow{D}_\nu^A H) V^{A,\mu\nu}$
		$(\partial_\mu \phi) (H^\dagger i \overleftrightarrow{D}_\nu^A H) \tilde{V}^{A,\mu\nu}$
$\phi H \bar{H} \psi_1 \bar{\psi}_2$	$\langle 1 p_\phi 2 \rangle$	$(\partial_\mu \phi) (\bar{\psi} \gamma^\mu \psi) H ^2$
		$(\partial_\mu \phi) (\bar{\psi} \gamma^\mu T^A \psi) (H^\dagger T^A H)$
$\phi H_1 \bar{H}_2 H_3 \bar{H}_4$	$p_\phi \cdot (p_1 - p_2) + \text{symm.}$	$(\partial^\mu \phi) (H^\dagger i \overleftrightarrow{D}_\mu H) H ^2$

Dimension 8

Particle content	$\mathcal{A} \times f^4$	$\mathcal{O} \times f^4$
ϕ^4	$(p_{\phi_1} \cdot p_{\phi_2})(p_{\phi_3} \cdot p_{\phi_4}) + \text{symm.}$	$(\partial_\mu \phi \partial^\mu \phi)^2$
$\phi^2 H_1 \bar{H}_2$	$(p_{\phi_1} \cdot p_{\phi_2})^2$	$(\partial_\mu \partial_\nu \phi \partial^\mu \partial^\nu \phi) H ^2$
	$(p_{\phi_1} \cdot p_1)(p_{\phi_2} \cdot p_2) + \text{symm.}$	$(\partial^\mu \phi \partial^\nu \phi) (D_\mu H^\dagger D_\nu H)$
$\phi^2 V_1 V_2$	$(p_{\phi_1} \cdot p_{\phi_2}) \langle 12 \rangle^2$	$(\partial_\alpha \phi \partial^\alpha \phi) V_{\mu\nu} V^{\mu\nu}$
		$(\partial_\alpha \phi \partial^\alpha \phi) V_{\mu\nu} \tilde{V}^{\mu\nu}$ $(\partial_\alpha \phi \partial^\alpha \phi) V_{\mu\nu}^A V^{A,\mu\nu}$ $(\partial_\alpha \phi \partial^\alpha \phi) V_{\mu\nu}^A \tilde{V}^{A,\mu\nu}$
$\phi^2 \psi_1 \bar{\psi}_2$	$\langle 1 p_{\phi_1} 2 \rangle \langle 1 p_{\phi_2} 2 \rangle$	$(\partial^\mu \phi \partial_\nu \phi) V_{\mu\alpha} V^{\alpha\nu}$
		$(\partial^\mu \phi \partial_\nu \phi) V_{\mu\alpha}^A V^{A,\alpha\nu}$
$\phi^2 \psi_1 \bar{\psi}_2 H$	$(p_{\phi_1} \cdot p_1) \langle 1 p_{\phi_2} 2 \rangle + \text{symm.}$	$(\partial_\mu \phi \partial_\nu \phi) (\bar{\psi} \gamma^\mu D^\nu \psi)$
$\phi^2 \psi_1 \bar{\psi}_2 H$	$(p_{\phi_1} \cdot p_{\phi_2}) \langle 12 \rangle$	$(\partial_\mu \phi \partial^\mu \phi) \bar{\psi} H \psi$
$\phi^2 H^4$	$p_{\phi_1} \cdot p_{\phi_2}$	$(\partial_\mu \phi \partial^\mu \phi) H ^4$

Phenomenological application

- Process $\ell^- \ell^+ \rightarrow \phi h$ for (not so) high energies
- Cross-section dominated by higher-order interactions

$$\mathcal{L} \supset \frac{C_{\phi\ell^2 D^2 H}^{(1)}}{f^3} \partial_\mu \phi D^\mu \bar{L} H \ell_R + \frac{C_{\phi\ell^2 D^2 H}^{(1)}}{f^3} \partial_\mu \phi \bar{L} H D^\mu \ell_R + h.c.$$

$$\frac{d\sigma^{\text{tot}}(\ell^- \ell^+ \rightarrow \phi h)}{d \cos \theta} \simeq \frac{1}{512\pi f^6} \left[(u^2 + t^2) \left(|C_{\phi\ell^2 HD^2}^{(1)}|^2 + |C_{\phi\ell^2 HD^2}^{(2)}|^2 \right) + 4tu \text{Re} \left(C_{\phi\ell^2 HD^2}^{(1)} C_{\phi\ell^2 HD^2}^{(2)*} \right) \right],$$

