

# Spontaneous symmetry breaking

오수민

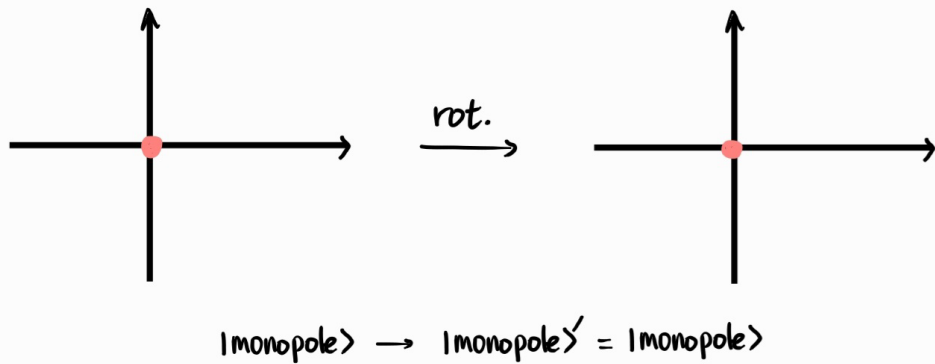
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# Contents

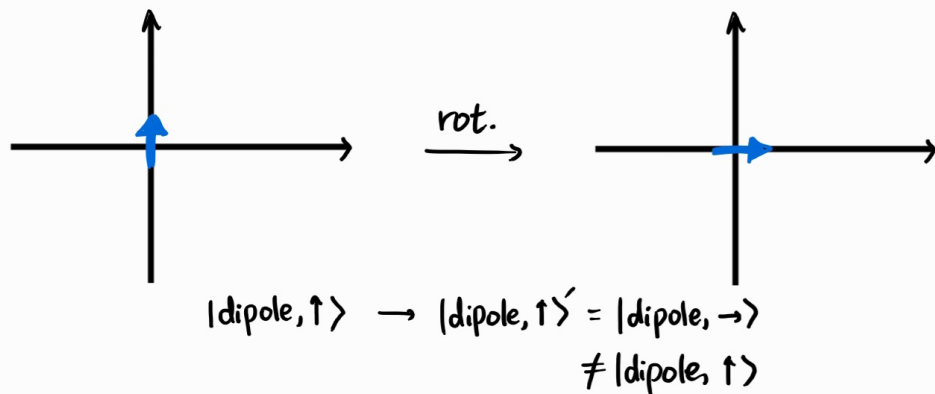
- Symmetry is important!
- Phase transition
- Landau theory
- Higgs mechanism

# Symmetry is important!

- States with different symmetry properties are different
- Ex) monopole vs dipole



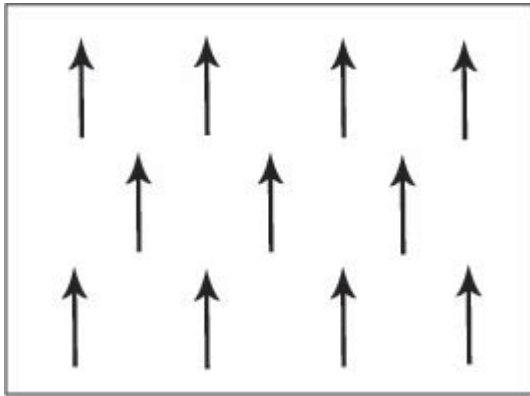
Trivial under rotational symmetry



Non-trivial under rotational symmetry

# Phase transition

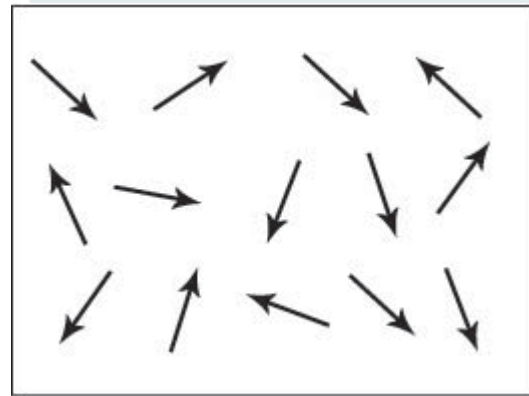
- Physical process of transition between one state and another
- 'Symmetry breaking' can be understood as phase transition
- Ex) ferromagnet vs paramagnet under  $\mathbb{Z}_2 : m \rightarrow -m$



$$T < T_c$$

$$m = \langle \sigma_j \rangle \neq 0$$

symmetry broken



$$T > T_c$$

$$m = \langle \sigma_j \rangle = 0$$

symmetric

# Landau theory

- In equilibrium, systems minimize the free energy!
- Landau's suggestions:
  - 1) Free energy is invariant under symmetry transformations
  - 2) Free energy is analytic

Under these conditions,

one can construct a free energy expression w.r.t. the order parameter

# Landau theory

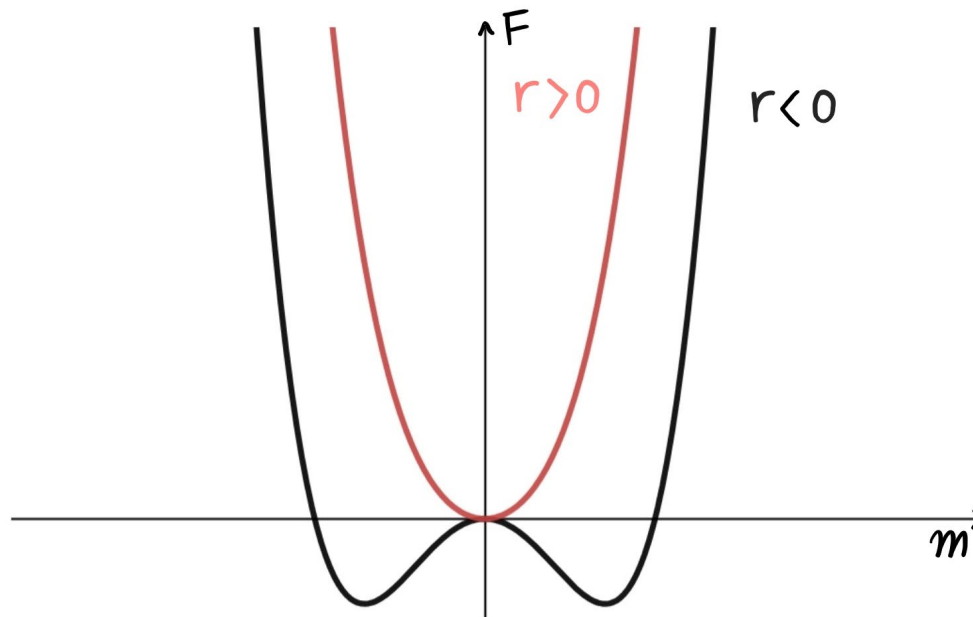
- Consider the previous case  $\mathbb{Z}_2 : m \rightarrow -m$

For the free energy to be invariant under  $\mathbb{Z}_2$  symmetry,

$$F(m) = a_0 + a_2 m^2 + a_4 m^4 + \dots$$

- How can we construct free energy for the ferromagnetic system?

$$F(m) = \frac{r}{2} m^2 + \frac{u}{4} m^4$$



$$\frac{\partial F}{\partial m} = rm + um^3 = 0$$

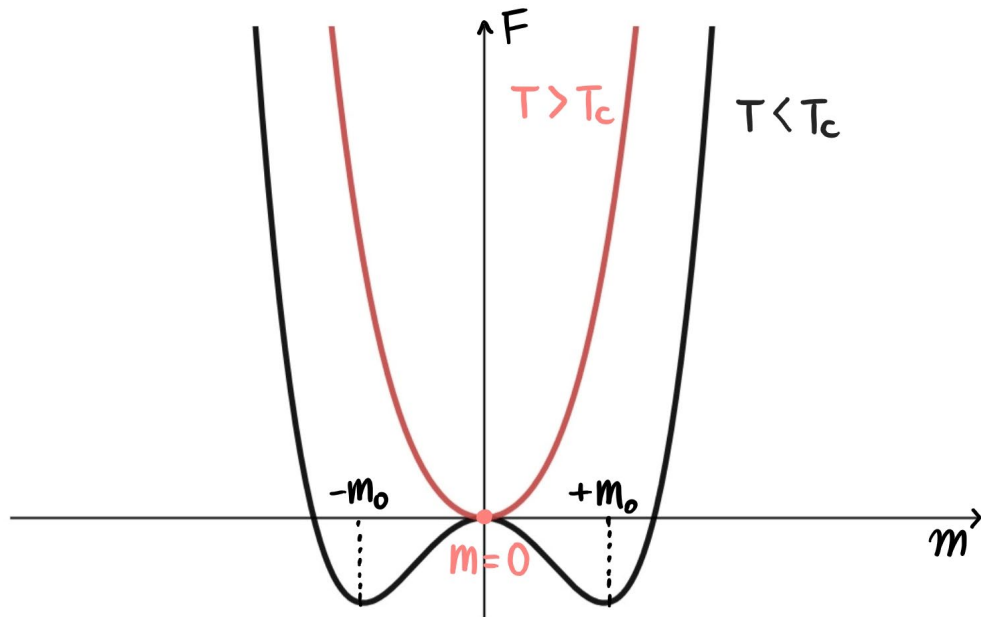
$$r > 0: m = 0$$

$$r < 0: m_0 = \pm \sqrt{\frac{|r|}{u}} \neq 0$$

# Landau theory

- Ferromagnet:  $T < T_C$ ,  $m \neq 0$
- Paramagnet:  $T > T_C$ ,  $m = 0$

$$F(T, m) = \frac{\alpha(T - T_C)}{2} m^2 + \frac{u}{4} m^4$$



As the temperature cools down, the system breaks the  $\mathbb{Z}_2$  symmetry ( $m_0 \neq 0$ ) even though the free energy is still invariant

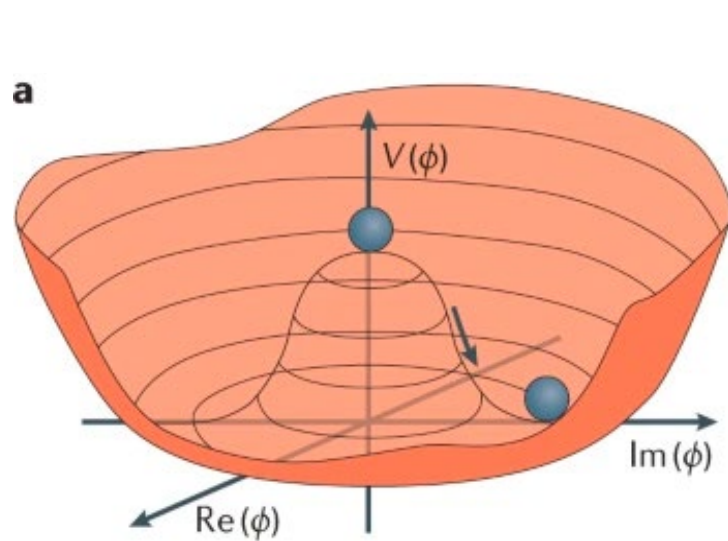
→ **Spontaneous symmetry breaking**

# Higgs mechanism

- The Lagrangian is invariant under electroweak symmetry  $SU(2)_L \times U(1)_Y$

$$\mathcal{L}_{Higgs} = (D^\mu \phi)^\dagger (D_\mu \phi) - V(\phi)$$

$$V(\phi) = -\mu^2 |\phi|^2 + \lambda |\phi|^4$$



Somehow nature chose  $\mu^2 > 0 \rightarrow V(\phi)$  minimum at  $|\phi| = \sqrt{\frac{\mu^2}{2\lambda}} = \frac{v}{\sqrt{2}}$

We may choose a direction:  $\phi_0 = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$

→ Only one component of the doublet gains nonzero value

→ **SU(2) broken**

→ Higgs boson has nonzero hypercharge ( $Y_H = \frac{1}{2}$ ) → **U(1)<sub>Y</sub> broken**

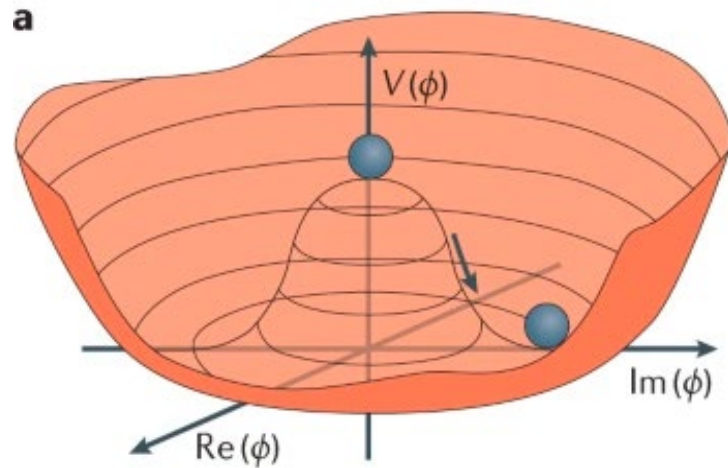


# Higgs mechanism

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$$\text{Electric charge: } Q = T^3 + Y_H = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$U(1)_{EM} : \phi_0 \rightarrow \phi'_0 = e^{iQ\theta} \phi_0 = \phi_0$$

**$U(1)_{EM}$  survives!**

→ **Spontaneous symmetry breaking**  $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$

# Higgs mechanism

- Generating masses to gauge bosons

$$D_\mu \phi = (\partial_\mu + igT^i W_\mu^i + ig'Y_H B_\mu)\phi$$

$$\begin{aligned}(D^\mu \phi_0)^\dagger (D_\mu \phi_0) &= \left| (\partial_\mu + igT^i W_\mu^i + ig'Y_H B_\mu) \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \right|^2 = \dots \\ &= \frac{v^2}{8} \left[ g^2 \left( (W_\mu^1)^2 + (W_\mu^2)^2 \right) + (gW_\mu^3 - g'B_\mu)^2 \right] = \frac{1}{2} \left( \frac{gv}{2} \right)^2 W_\mu^+ W^{-\mu} + \frac{1}{2} \left( \sqrt{g^2 + g'^2} \frac{v}{2} \right)^2 Z_\mu Z^\mu\end{aligned}$$

$$W_\mu^\pm = \frac{1}{\sqrt{2}} (W_\mu^1 \mp iW_\mu^2), \quad m_W = \frac{gv}{2}$$

$$Z_\mu = \frac{1}{\sqrt{g^2 + g'^2}} (gW_\mu^3 - g'B_\mu), \quad m_Z = \sqrt{g^2 + g'^2} \frac{v}{2}$$

$$A_\mu = \frac{1}{\sqrt{g^2 + g'^2}} (g'W_\mu^3 + gB_\mu), \quad m_A = 0$$

# References

- Contents until Landau theory is based on the lectures by E.G.Moon, PH211 2022 Fall, KAIST Physics

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Kien Nguyen (July 2009). *The Higgs Mechanism*, [https://www.theorie.physik.uni-muenchen.de/lsfrey/teaching/archiv/sose\\_09/rng/higgs\\_mechanism.pdf](https://www.theorie.physik.uni-muenchen.de/lsfrey/teaching/archiv/sose_09/rng/higgs_mechanism.pdf)

**THANK YOU**