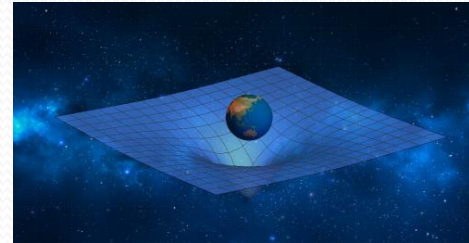
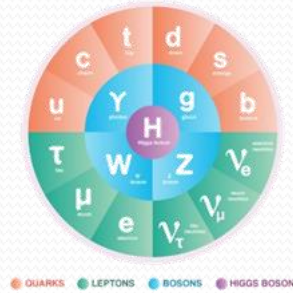
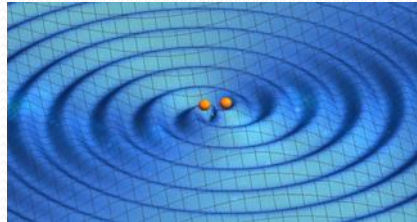
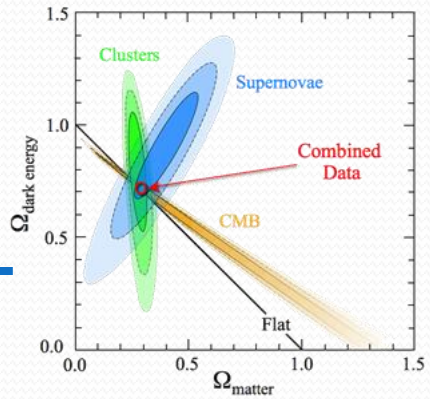
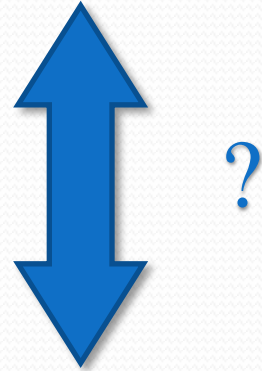
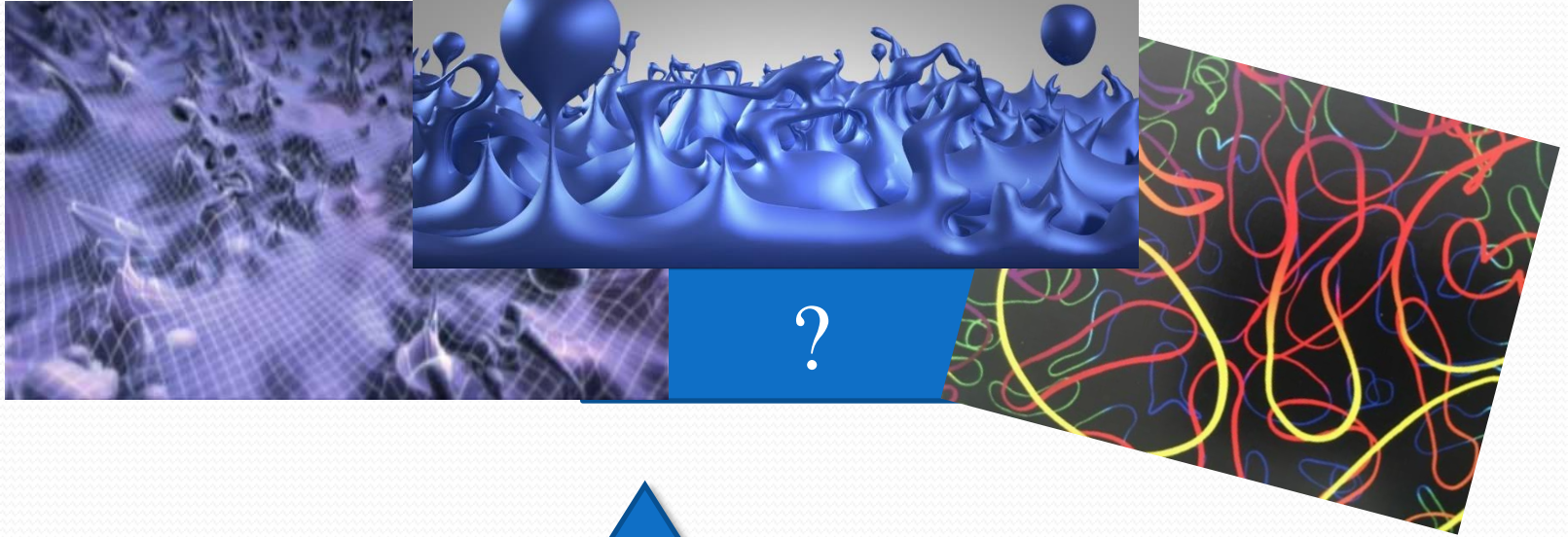
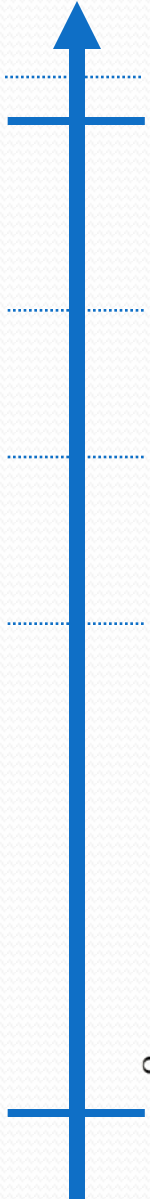




Tools in Effective Field Theories
of (Massive) Gravity

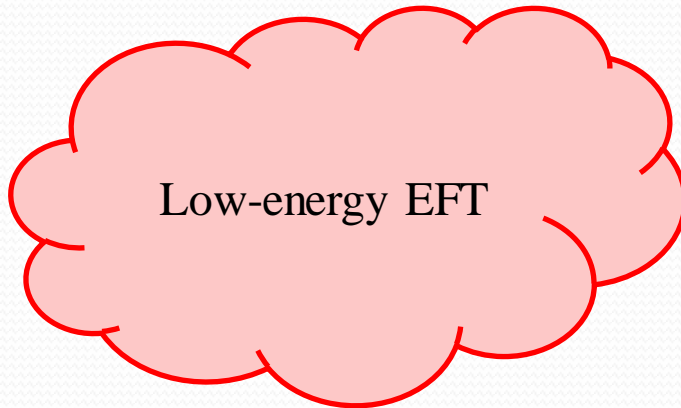
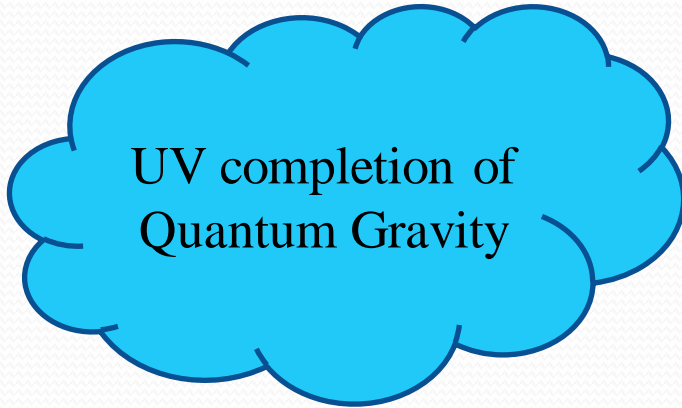
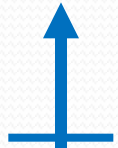


Energy



Effective Field Theory

Energy



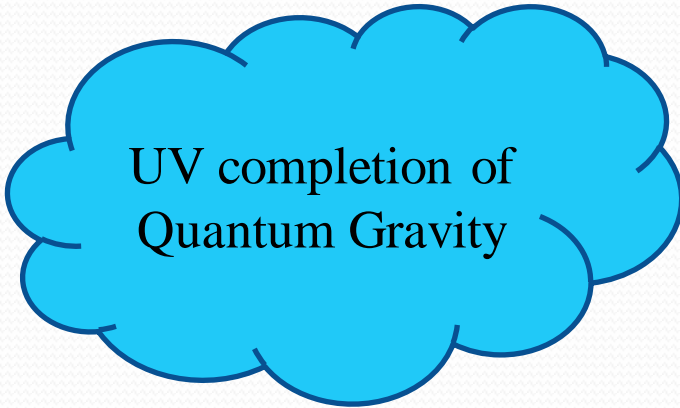
E.g.: GR + irrelevant operators
+ (B)SM fields
+ light degrees of freedom



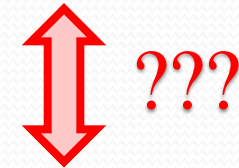
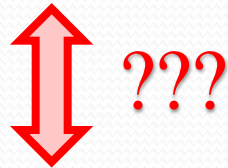
Join Noirmoutier Sandcastle Competition

Effective Field Theory

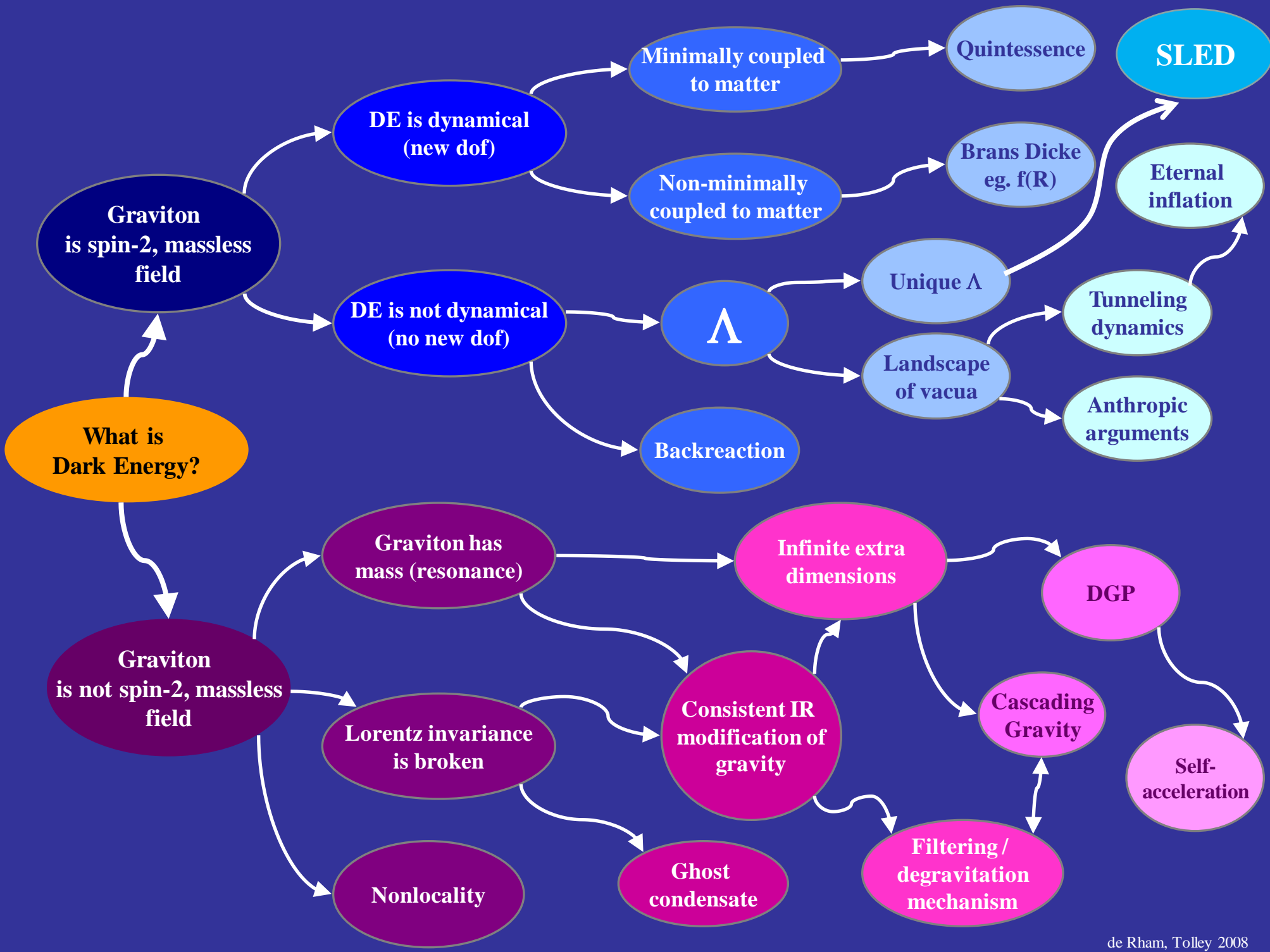
Energy

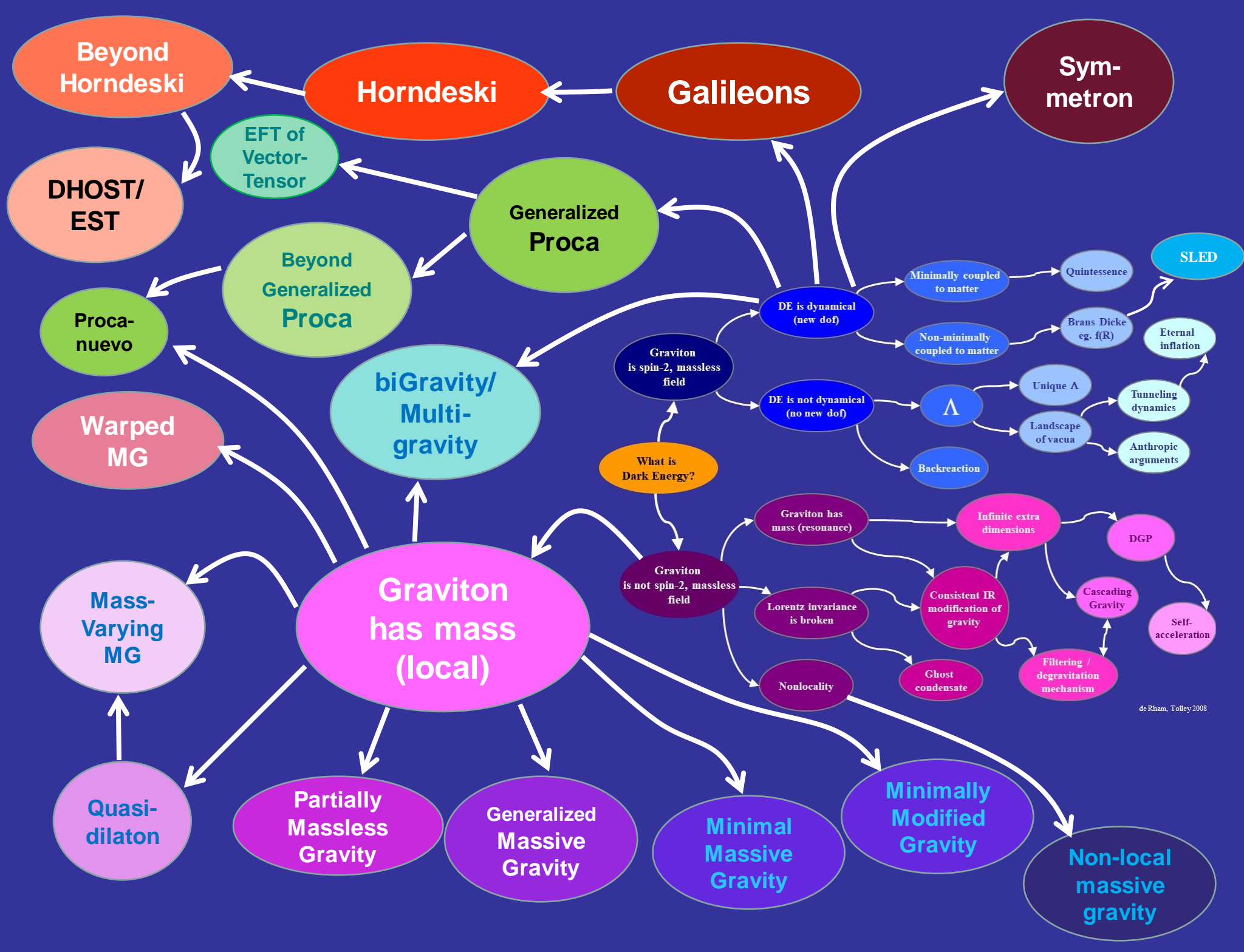


UV completion of
Quantum Gravity



Join
Noirmoutier
Sandcastle
Competition





Dark Matter + SuperCDMS

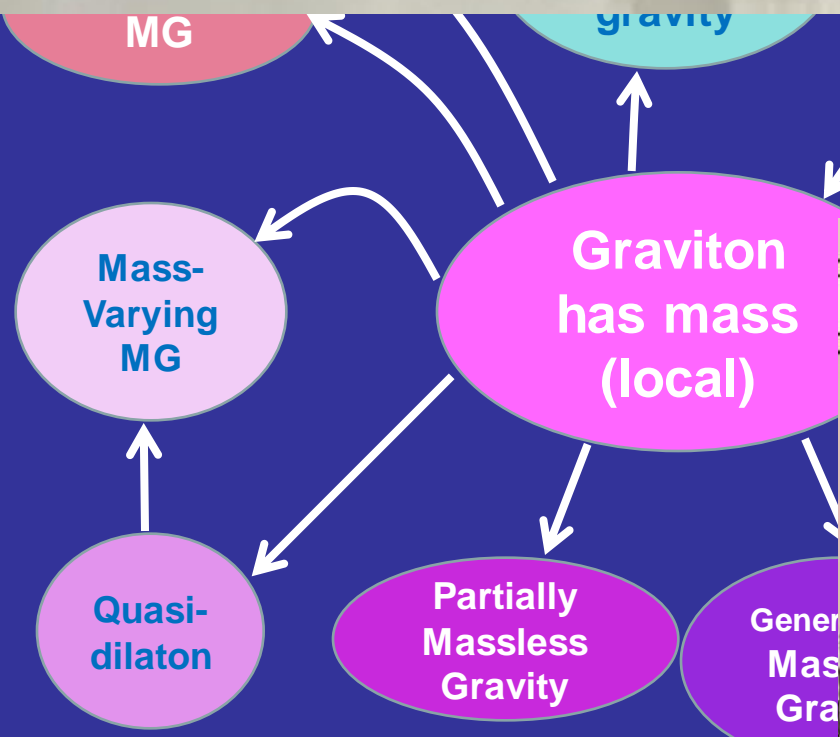
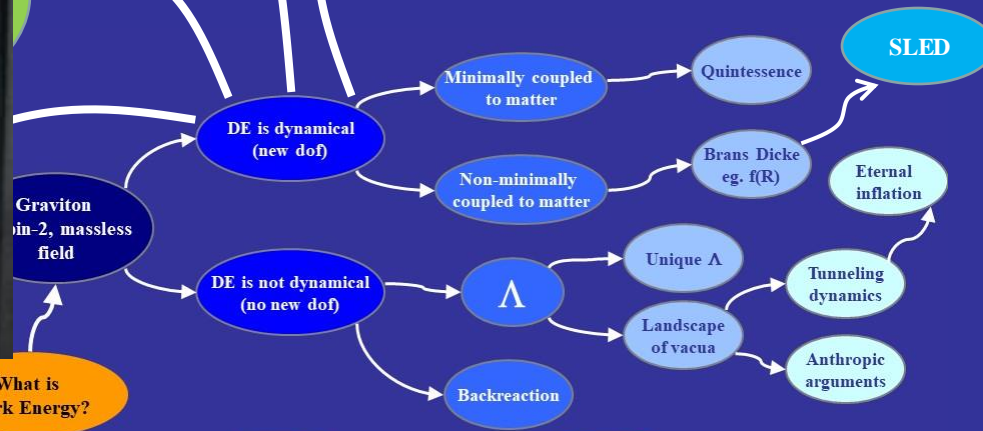
Where and how we're looking



Same story for Dark Matter, eg see Madeleine Zurowski's talk

Galileons

Symmetron



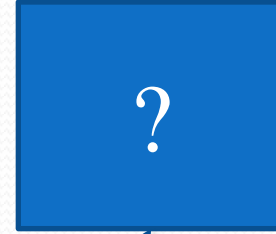
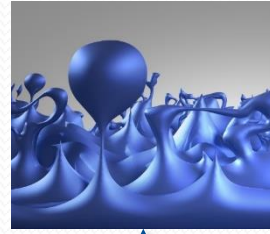
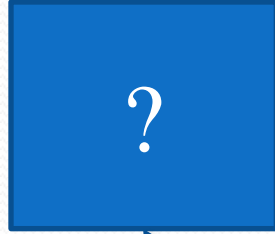
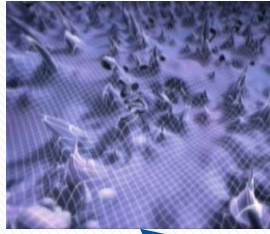
They all ask "What is Dark Energy" and "Where is Dark Energy", but nobody asks "HOW is Dark Energy?"



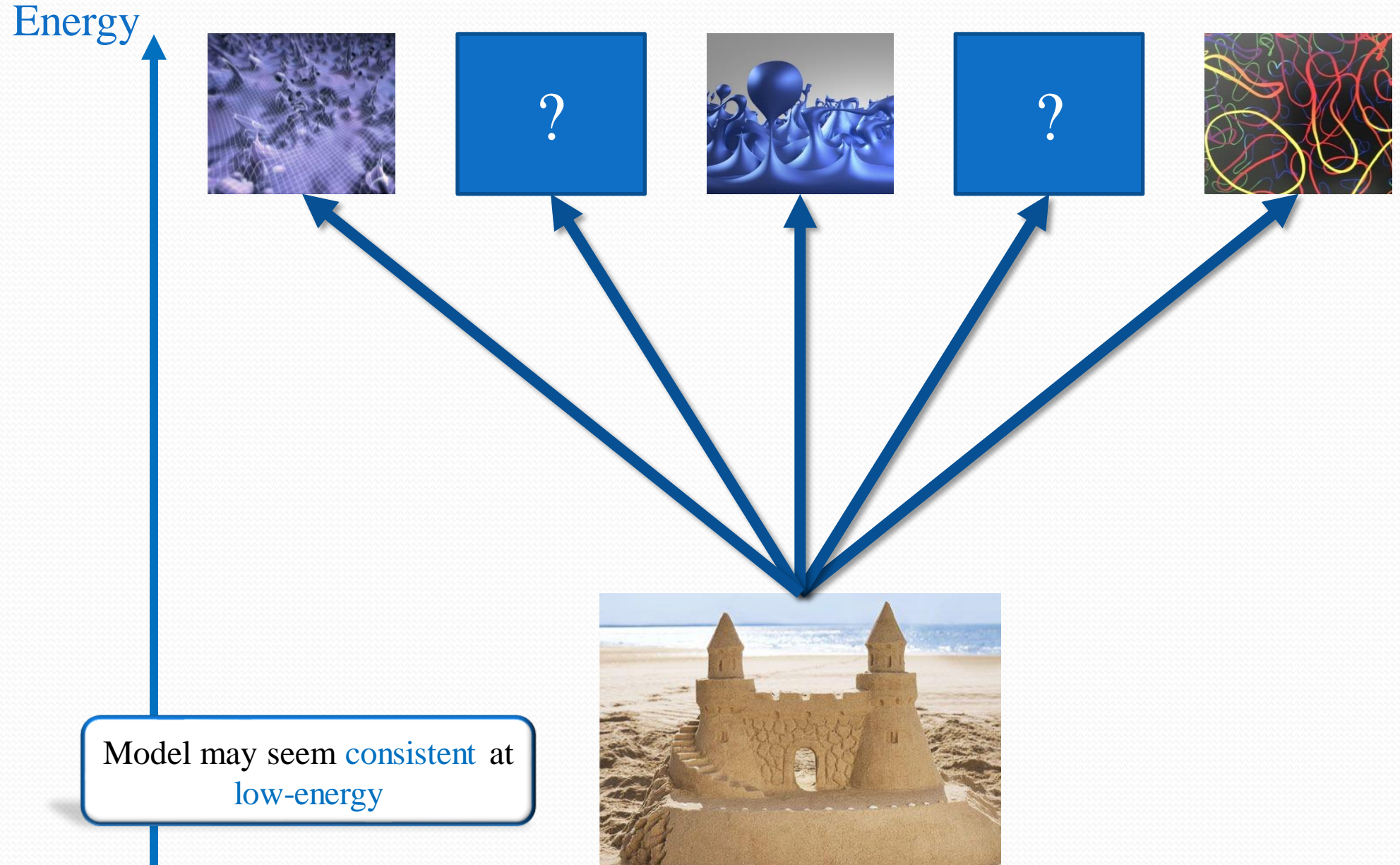
ham, Tolley 2008

High-energy completion

Energy

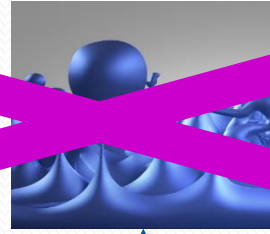
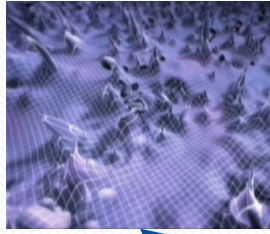


Model may seem *consistent* at
low-energy



High-energy completion

Energy



Yet we know that some models
can **never** enjoy
a **standard UV completion**

60 orders
of magnitude

Model may seem **consistent** at
low-energy



Its **not** that a theory with
parameter $\alpha = 0$ is ruled out
but $|\alpha| > 10^{-60}$ is allowed

Constraints from UV Completion

Starting with the assumption of a healthy Wilsonian UV completion
(string theory or other)

Set up classes of IR constraints

Positivity Bounds

Causality Bounds

Precise statements at the level
of $2 \rightarrow 2$ scattering amplitude
(complement S-matrix bootstrap program)

Unresolvable support outside
front velocity light cone



Alberte



Carrillo-Gonzalez



Chen



Franckfort



Heisenberg



Held



Jaitly



Klapanek



Kozuszek



Margalit



Melville



Momeni



Noller



Pozsgay



Rumbutis



Tokareva



Tolley



Wiseman



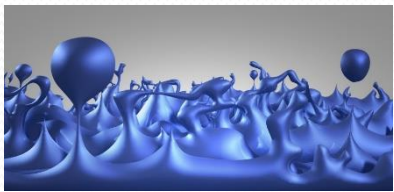
Zhang



Zhou

CoM Energy

Low-energy EFT



- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)
- ✓ CAUSAL (analyticity)
- ✓ Local (Froissart Bound)

A : 2 – 2 elastic amplitude

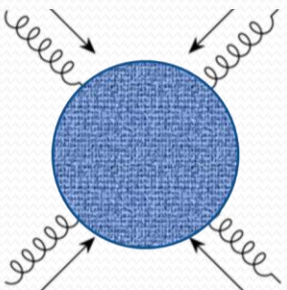
$$2 \text{Im} \left[\text{Diagram 1} \right] = \sum_X \left| \text{Diagram 2} \right|^2 \geq \left| \text{Diagram 3} \right|^2$$



Positivity bounds

(applied to low-energy scattering amplitude)

Low-energy EFT



$$\left. \frac{d^2 \mathcal{A}(s, t)}{ds^2} \right|_{t=0} > 0 \quad s < 4m^2$$

Pham and Truong 1985
 Ananthanarayan, Toublan and Wanders, 1994
 Adams et. al. 2006

Applications to SMEFT

New Physics Beyond the Standard Model

$$\mathcal{L}_{\text{BSM}} = \mathcal{L}_{\text{SM}}^{(4)} + \text{New Physics} \equiv \text{new modes with mass } M \gtrsim \Lambda \gtrsim \text{TeV}$$



$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}}^{(4)} + \frac{1}{\Lambda} \sum_n c_n^{(5)} \mathcal{O}_n^{(5)} + \frac{1}{\Lambda^2} \sum_n c_n^{(6)} \mathcal{O}_n^{(6)} + \dots$$

Quickly leads to a huge parameter space

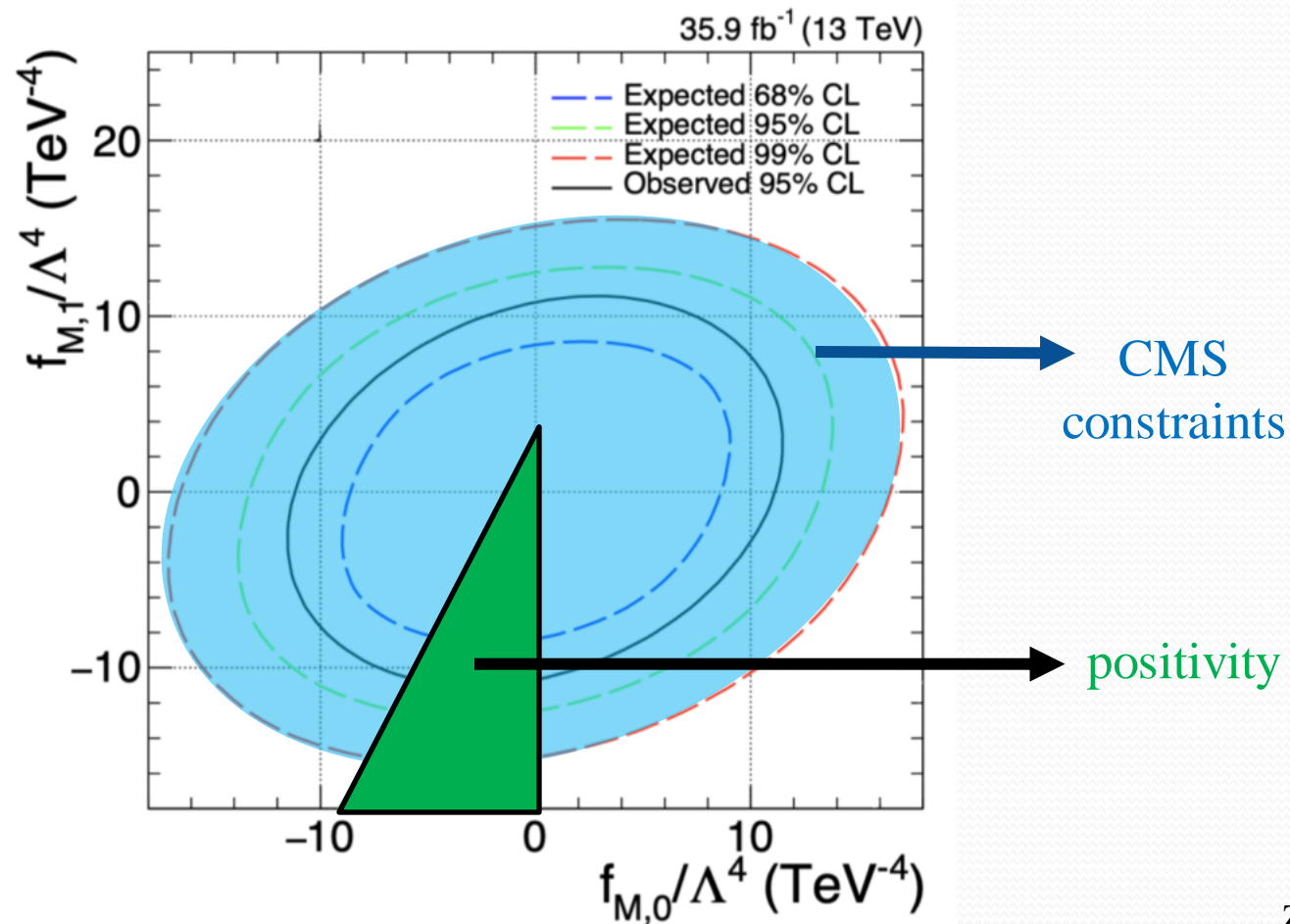
e.g. SMEFT up to dim-6 operators given in 1008.4884 includes 59 operators + B-violating

Becomes even more challenging at higher order
(44,807 dim-8 operators...)

From Grzadkowski, Iskrzyński,
Misiak&Rosiek 1008.4884

SMEFT

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}}^{(4)} + \frac{1}{\Lambda} \sum_n c_n^{(5)} \mathcal{O}_n^{(5)} + \frac{1}{\Lambda^2} \sum_n c_n^{(6)} \mathcal{O}_n^{(6)} + \dots$$



Adding Gravity



Causality \longleftrightarrow analyticity

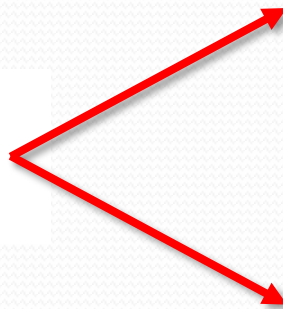
Locality, Froissart Bound with Gravity???

- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)
- ✓ **CAUSAL** (analyticity)
- ✓ Local (Froissart Bound)



positivity bounds

More subtle for
gravitational EFTs



(sub)luminal
sound speed

Applications to EFT of Cosmology/Gravity

Direct application to cosmology
is more subtle due to:

- absence of S-matrix
- lack of Lorentz invariance/crossing symmetry
- ambiguous connection with analyticity
- gravitational exchange pole

Directions in progress:

- Cosmological bootstrap on correlators (only perturbative)

No constraints from UV completion (analyticity, non-perturbative UV unitarity still need to be formulated)

Arkani-Hamed, Baumann, Lee & Pimentel, 1811.00024

Baumann, Duaso Pueyo, Joyce, Lee & Pimentel, 1910.14051, 2005.04234

see Sleight & Taronna for Bootstrapping Inflationary Correlators in Mellin Space, 1907.01143, 2007.09993

Applications to EFT of Cosmology/Gravity

Direct application to cosmology
is more subtle due to:

- absence of S-matrix
- lack of Lorentz invariance/crossing symmetry
- ambiguous connection with analyticity
- gravitational exchange pole

Directions in progress:

- Cosmological bootstrap on correlators (only perturbative)
- Alternative approach is to focus on **approximate S-matrix with broken boosts**
Full formalism still under development

CdR & Melville 1703.00025
Grall & Melville 2102.05683+2022+

See also recent work by Creminelli, Janssen & Senatore, 2207.14224
assuming conformal invariance in 3d

Applications to EFT of Cosmology/Gravity

Direct application to cosmology
is more subtle due to:

- absence of S-matrix
- lack of Lorentz invariance/crossing symmetry
- ambiguous connection with analyticity
- gravitational exchange pole

Directions in progress:

- Cosmological bootstrap on correlators (only perturbative)
- Approximate S-matrix with broken boosts
- Constrain **Wilson coefficients defined around Minkowski vacuum** and translate into constraints around cosmological backgrounds (requires adiabatic assumptions)

1904.05874, 1905.08816, 1906.11840, 1908.08644, 2103.06855, 2103.11195, ...



Will show a quick example

Applications to EFT of Cosmology/Gravity

Direct application to cosmology
is more subtle due to:

- absence of S-matrix
- lack of Lorentz invariance/crossing symmetry
- ambiguous connection with analyticity
- gravitational exchange pole

Directions in progress:

- Cosmological bootstrap on correlators (only perturbative)

- Approximate S-matrix with broken boosts

- Constrain Wilson coefficients defined around **Minkowski** vacuum

quick example

- Use **Infrared causality** as a proxy for positivity

Worked developed with Carrillo Gonzalez, Chen,
Jaitly, Margalit, Pozsgay, Tolley & Zhang



Constraining Models DE/MG

Example of DE/MG: quartic Horndeski with parameters $c_{B,M,T}$

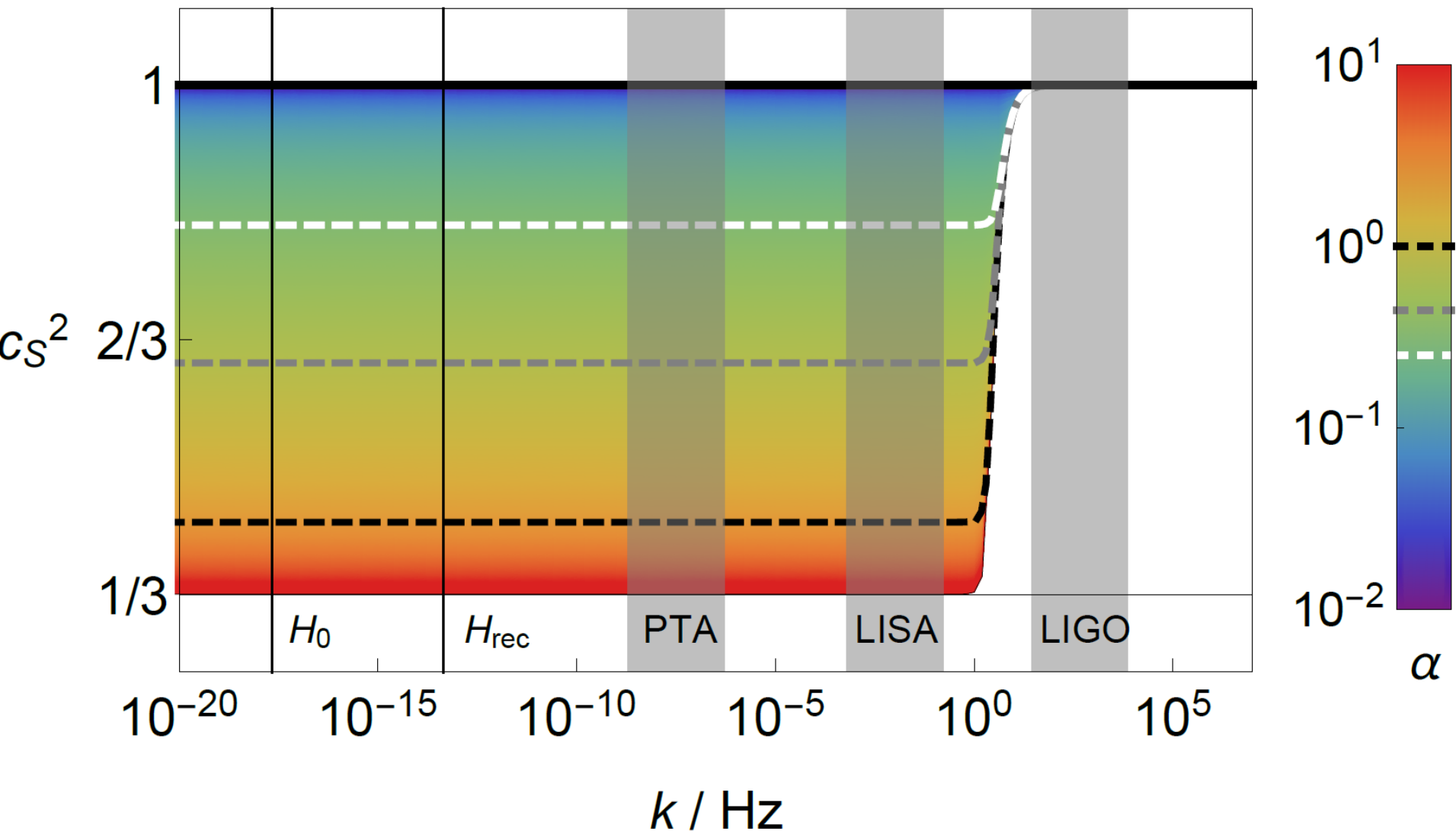
$$\mathcal{L} = M_{\text{Pl}}^2 G_4(X) R + \Lambda^4 G_2(X) + \frac{1}{H_0^2} G_{4,X}(X) \left((\square\phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right) + \mathcal{L}_{\text{matter}}(\psi, g)$$

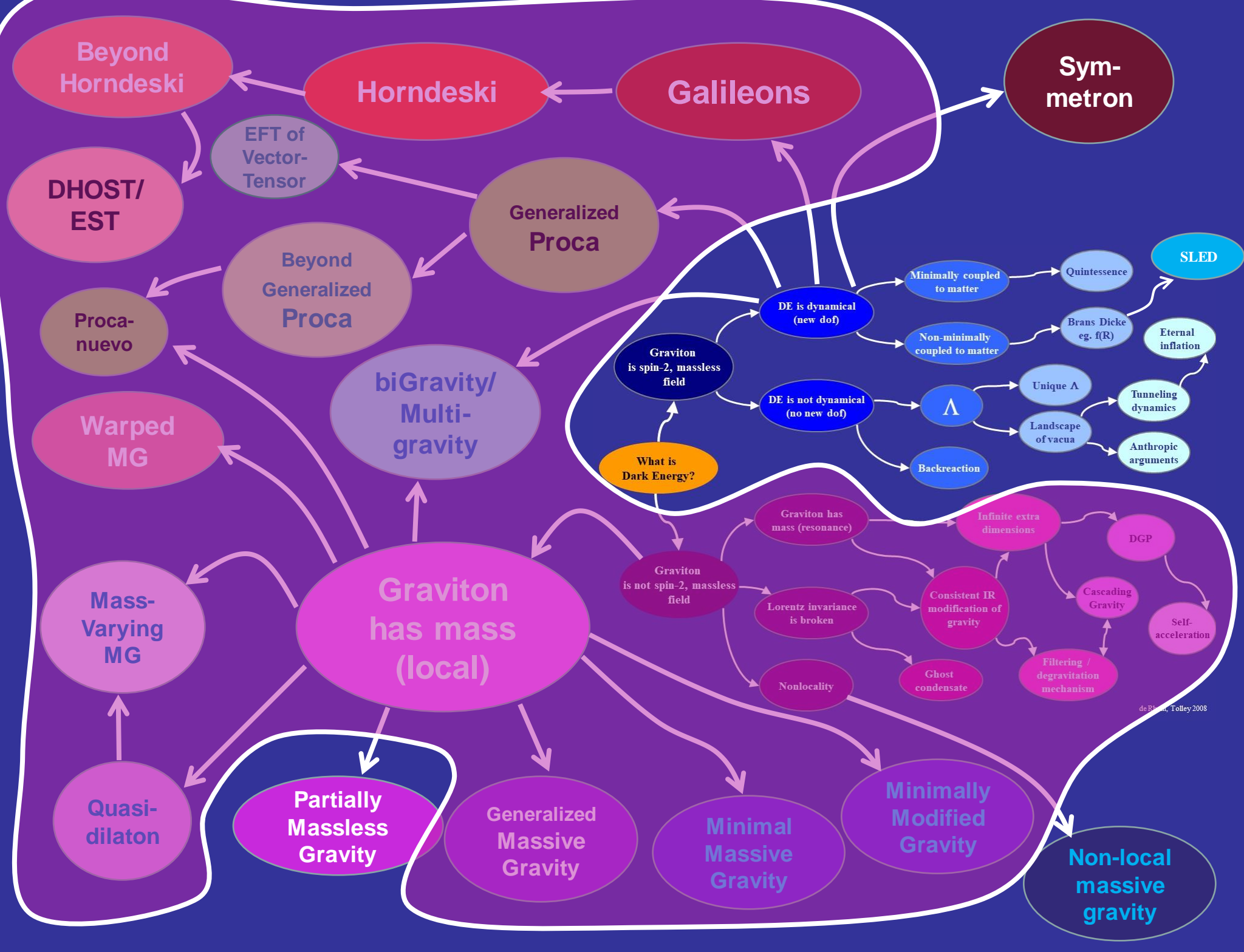
$$c_M = -\frac{2\dot{X}}{HM^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{Effective Planck Mass}$$

$$c_B = \frac{8X}{M^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{“Braiding” (scalar/tensor mixing)}$$

$$c_T = \frac{4X}{M^2\Omega_{\text{DE}}} G_{4,X}, \quad \text{Speed of GWs} - 1$$

$$M^2 = 2(G_4 - 2XG_{4,X})$$





Constraining Models DE/MG

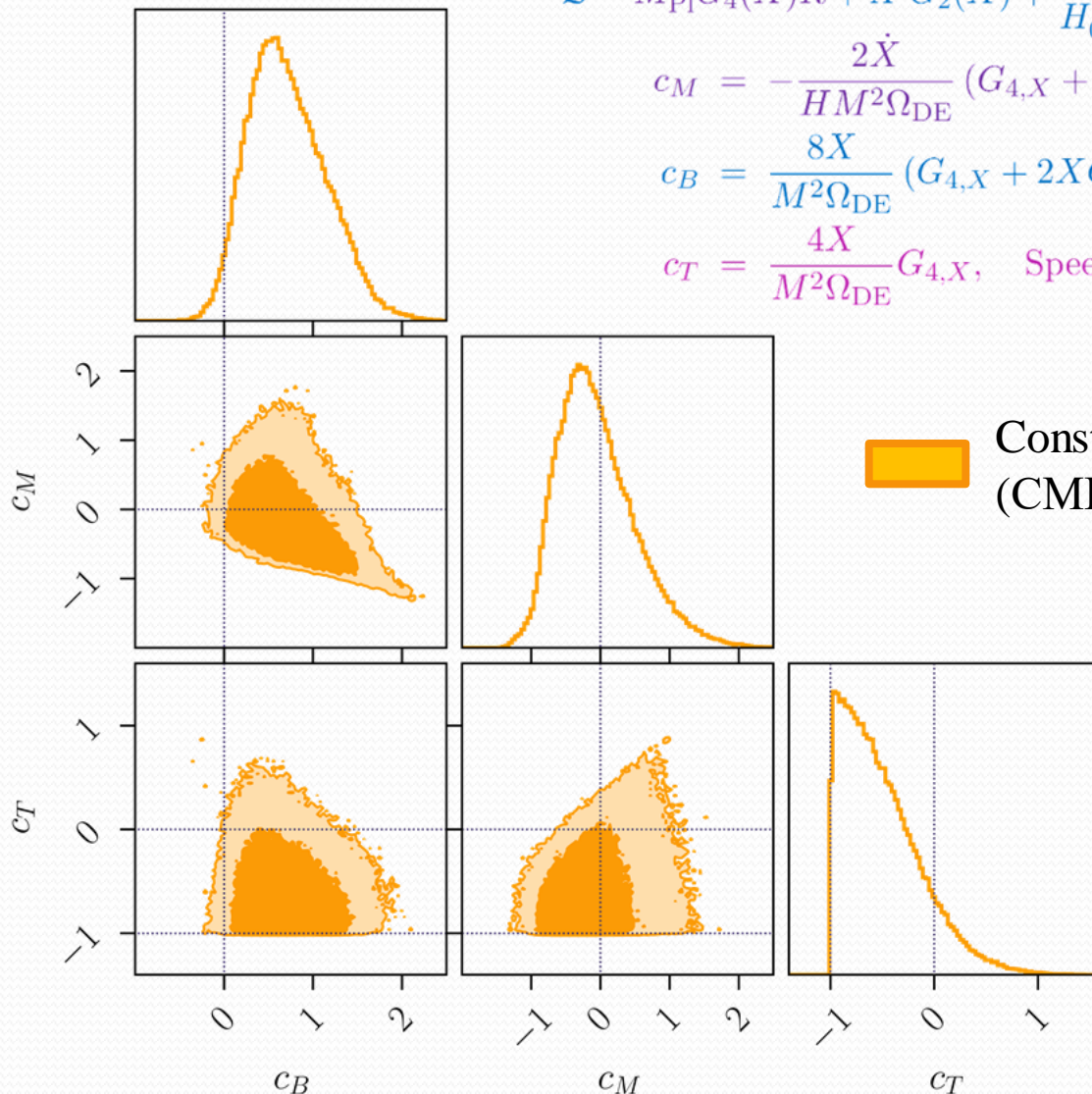
Example of DE/MG: quartic Horndeski with parameters $c_{B,M,T}$


$$\mathcal{L} = M_{\text{Pl}}^2 G_4(X) R + \Lambda^4 G_2(X) + \frac{1}{H_0^2} G_{4,X}(X) \left((\square\phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right) + \mathcal{L}_{\text{matter}}(\psi, g)$$

$$c_M = -\frac{2\dot{X}}{HM^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{Effective Planck Mass}$$

$$c_B = \frac{8X}{M^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{“Braiding” (scalar/tensor mixing)}$$

$$c_T = \frac{4X}{M^2\Omega_{\text{DE}}} G_{4,X}, \quad \text{Speed of GWs} - 1 \qquad M^2 = 2(G_4 - 2XG_{4,X})$$



 Constraints from observations
(CMB, BAO, redshift space distortion, ...)

Constraining Models DE/MG

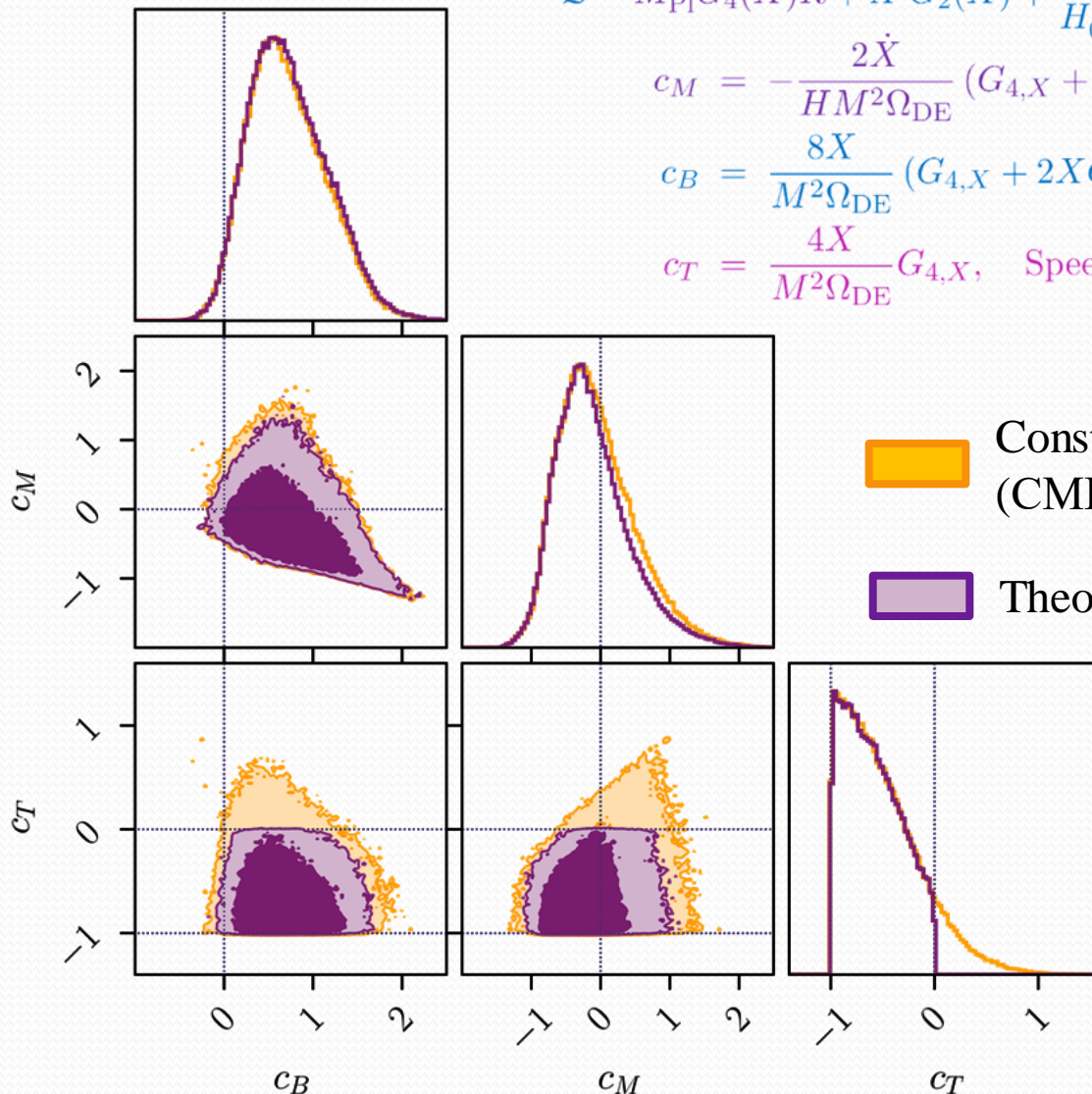
Example of DE/MG: quartic Horndeski with parameters $c_{B,M,T}$

$$\mathcal{L} = M_{\text{Pl}}^2 G_4(X) R + \Lambda^4 G_2(X) + \frac{1}{H_0^2} G_{4,X}(X) \left((\square\phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right) + \mathcal{L}_{\text{matter}}(\psi, g)$$

$$c_M = -\frac{2\dot{X}}{HM^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{Effective Planck Mass}$$

$$c_B = \frac{8X}{M^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{“Braiding” (scalar/tensor mixing)}$$

$$c_T = \frac{4X}{M^2\Omega_{\text{DE}}} G_{4,X}, \quad \text{Speed of GWs} - 1 \quad M^2 = 2(G_4 - 2XG_{4,X})$$

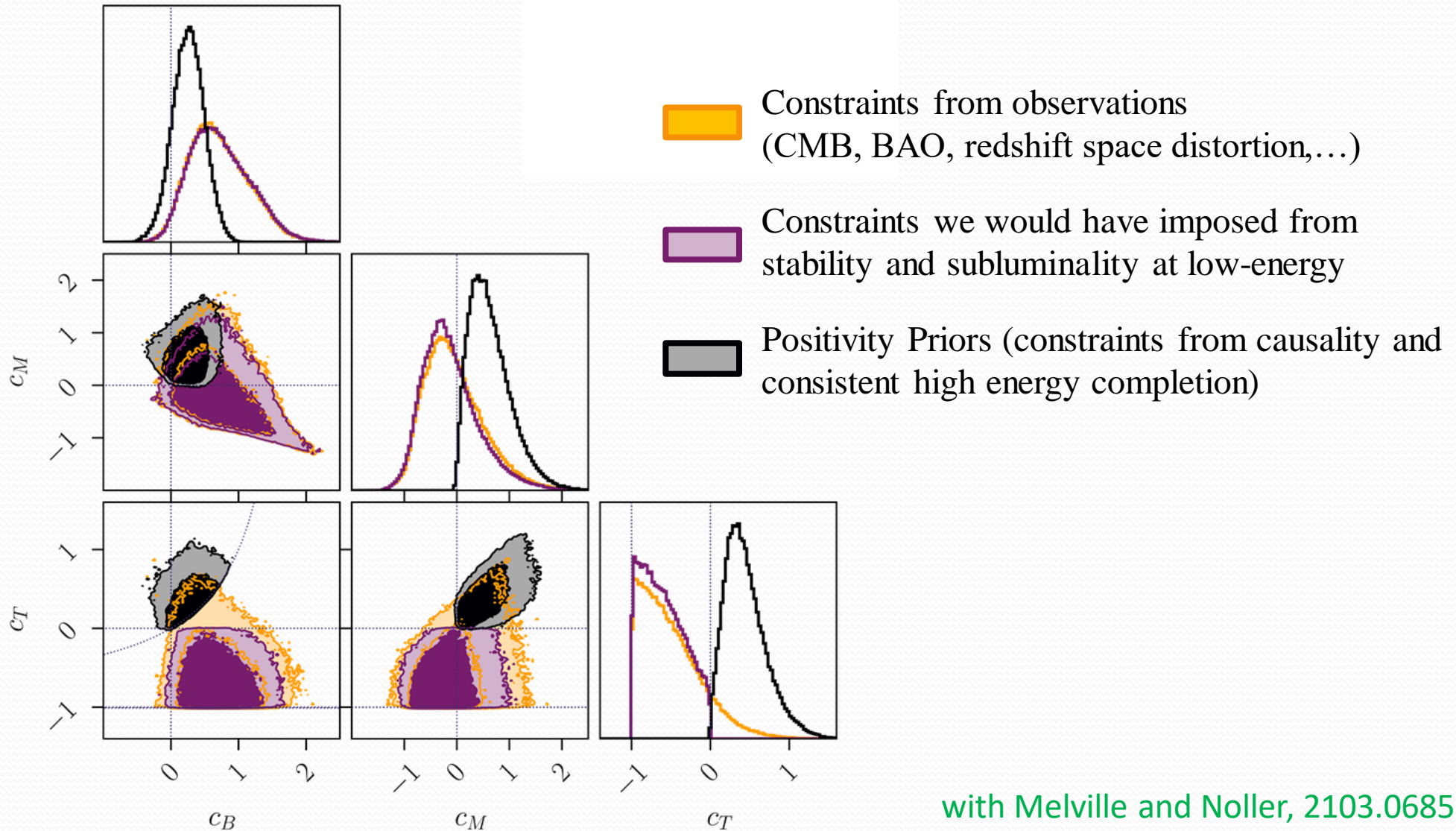


Constraints from observations
 (CMB, BAO, redshift space distortion,...)

Theoretical constraints

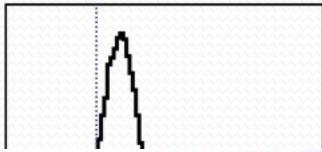
Constraints on Low-energy Models

Example of Dark Energy model (quartic Horndeski)
with parameters $c_{B,M,T}$

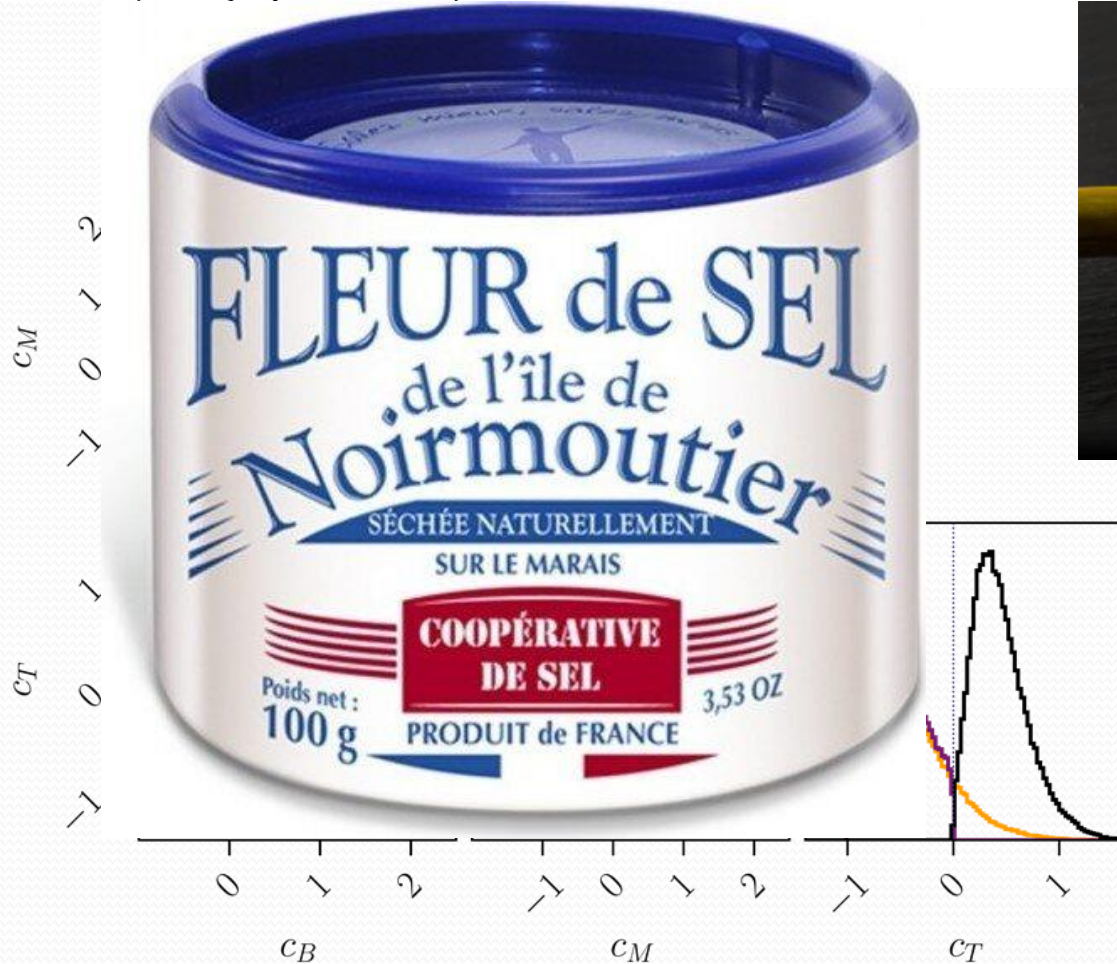


Constraints on Low-energy Models

Example of Dark Energy model (quartic Horndeski)
with parameters $c_{B,M,T}$



Note: To be taken with a pinch of salt



Only makes sense if
assume can connect
back to Minkowski

Applications to EFT of Cosmology

Direct application to cosmology
is more subtle due to:

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- lack of Lorentz invariance/crossing symmetry
- ambiguous connection with analyticity
- gravitational exchange pole

Directions in progress:

- Cosmological bootstrap on correlators (only perturbative)
 - Approximate S-matrix with broken boosts
 - Constrain Wilson coefficients defined around Minkowski vacuum
- Use Infrared causality as a proxy for positivity



Carrillo-Gonzalez



Chen



Jaitly



Margalit



Pozsgay



Tokareva



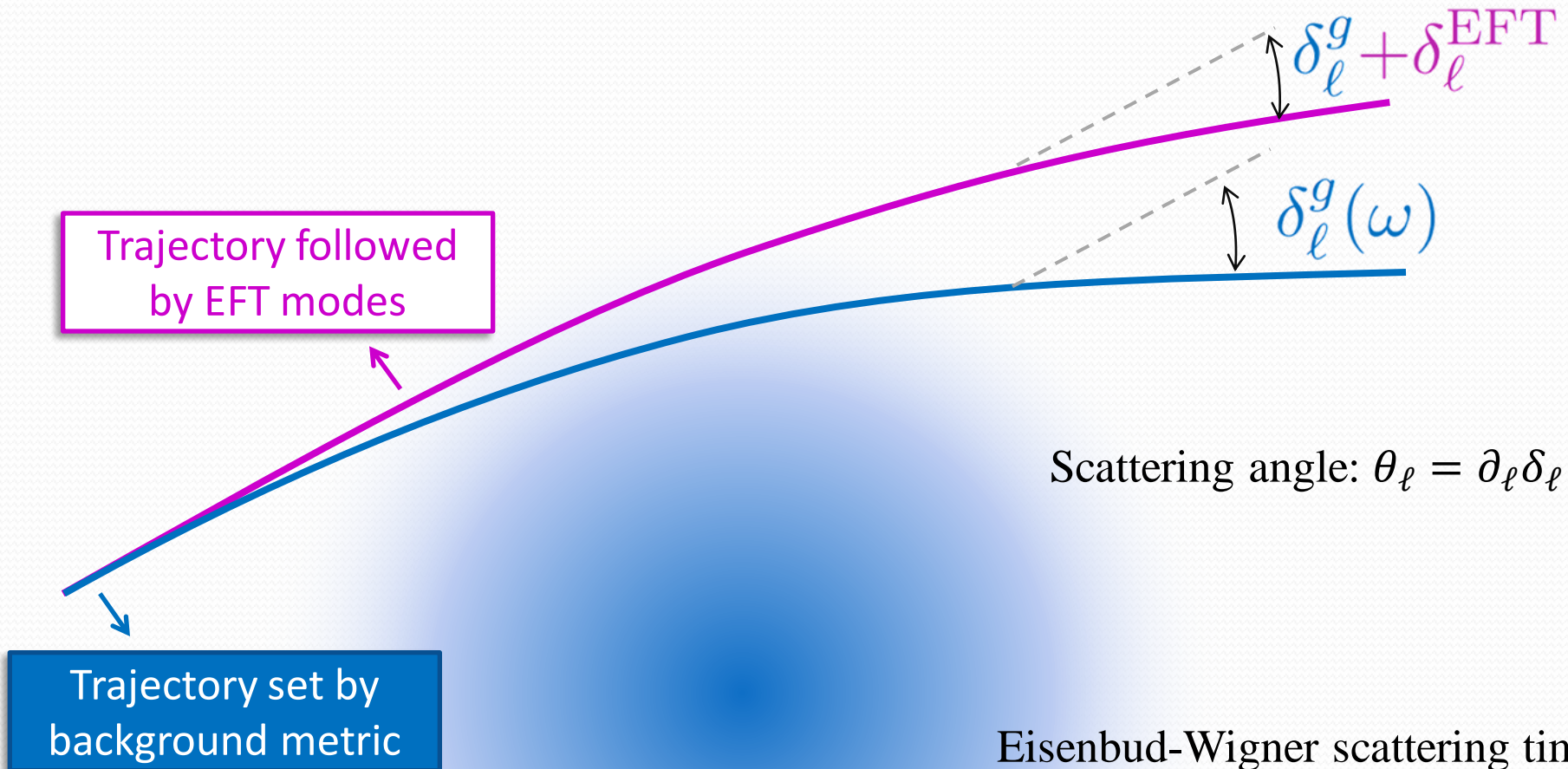
Tolley



Zhang



Scattering



Eisenbud-Wigner scattering time-delay:

$$\Delta T_\ell = 2 \frac{d}{d\omega} \delta_\ell(\omega)$$

$$\Delta T = \Delta T^g + \Delta T^{\text{EFT}}$$

Gravitational EFTs

- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)
- ✓ **CAUSAL** (analyticity)
- ✓ Local (Froissart Bound)



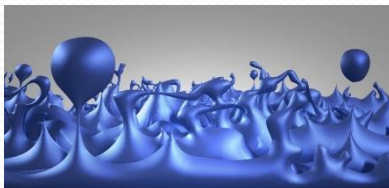
positivity bounds



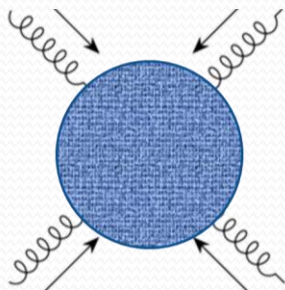
(sub)luminal
sound speed



CoM Energy

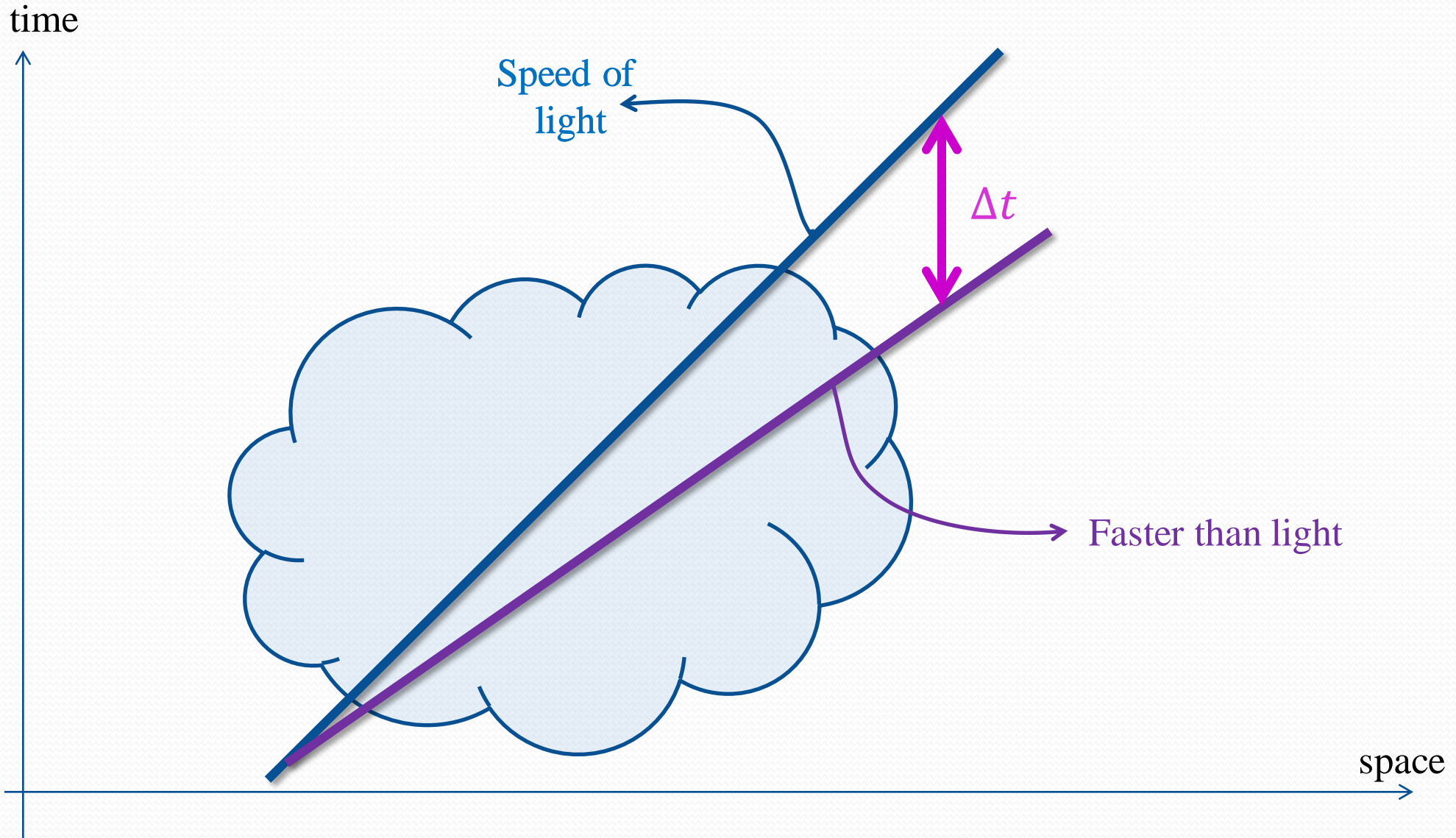


Low-energy
EFT



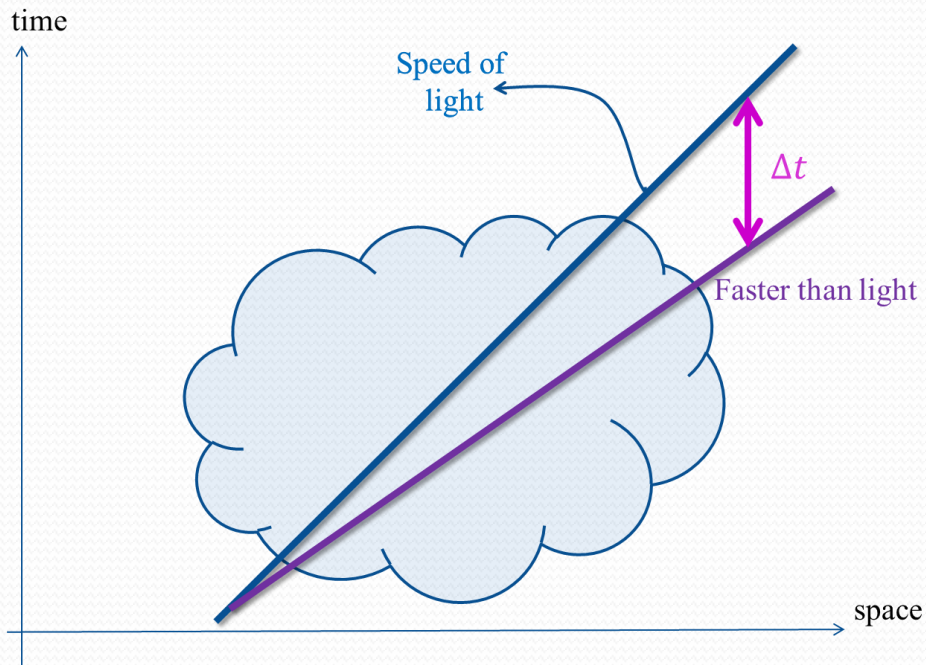
Small Superluminality – without gravity

As soon as a “*substance*” allows the tiniest superluminal speed, nothing prevents us from stacking it so as to end up with a significant/observable “time advance”



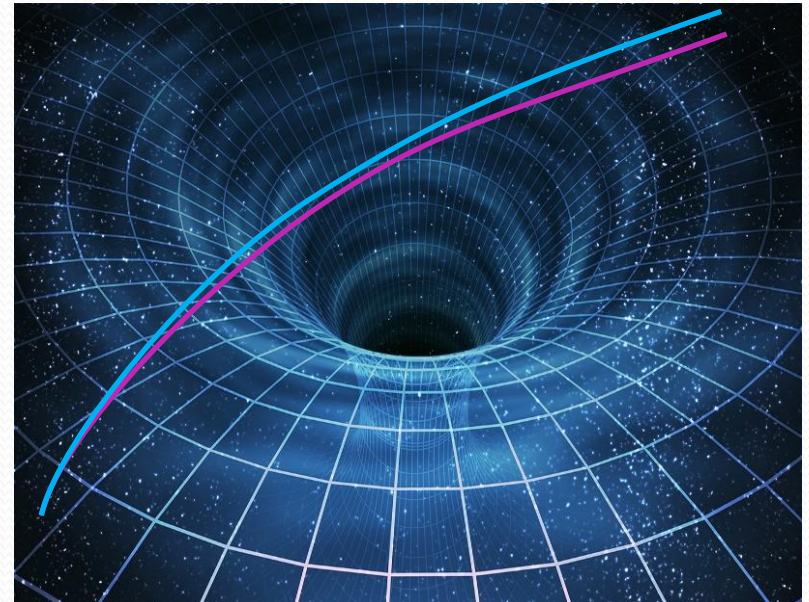
No Gravity

As soon as a “*substance*” allows the tiniest SL, nothing prevents us from stacking it so as to end up with a significant/observable “*time advance*”



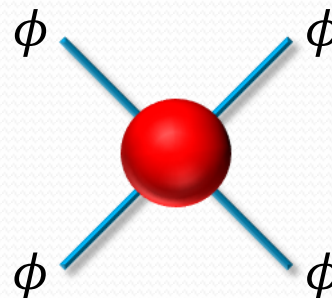
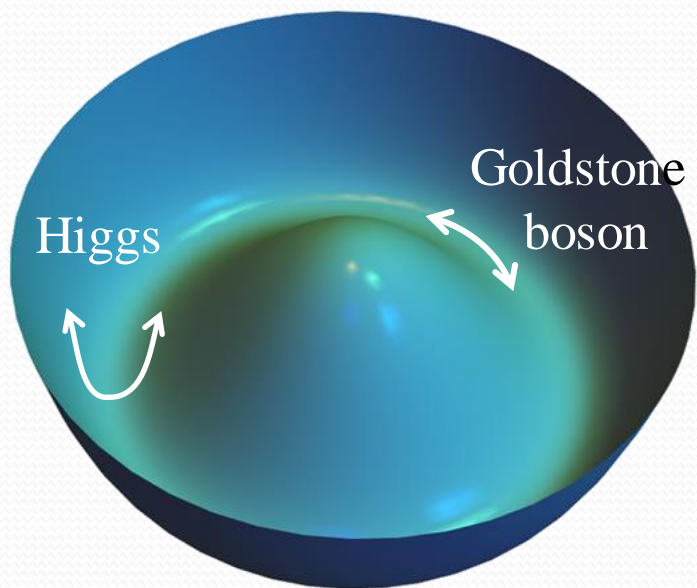
With Gravity

Anything living on the spacetime *inexorably* curves the geometry. There is a limit to “stacking”



Not always meaningful violation of causality

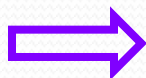
Goldstone Boson without Gravity



Subluminality



Positivity
bounds



$$v^2 \leq c^2$$

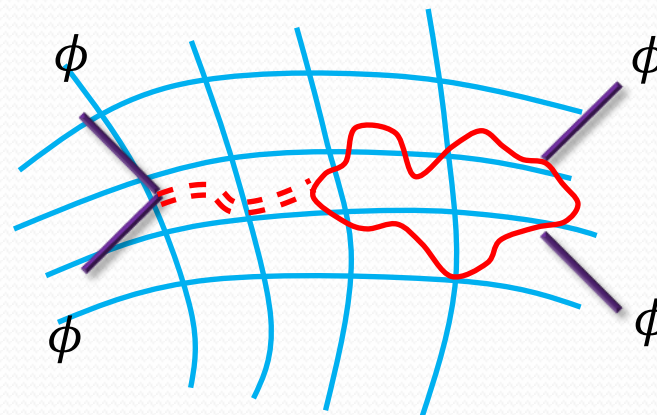
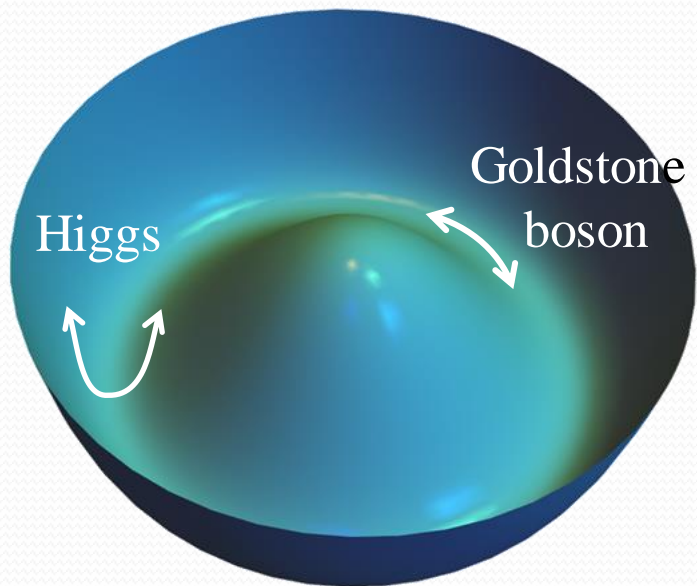


Asymptotic
Causality

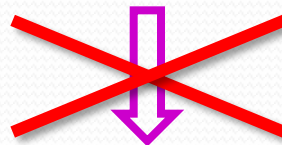


Infrared
Causality

Goldstone Boson with Gravity



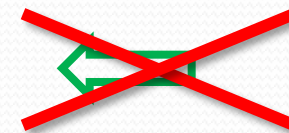
Subluminality
(in Einstein frame)



Positivity
bounds



$$v^2 \leq c^2 + \frac{\Lambda^2}{M_{\text{Pl}}^2}$$



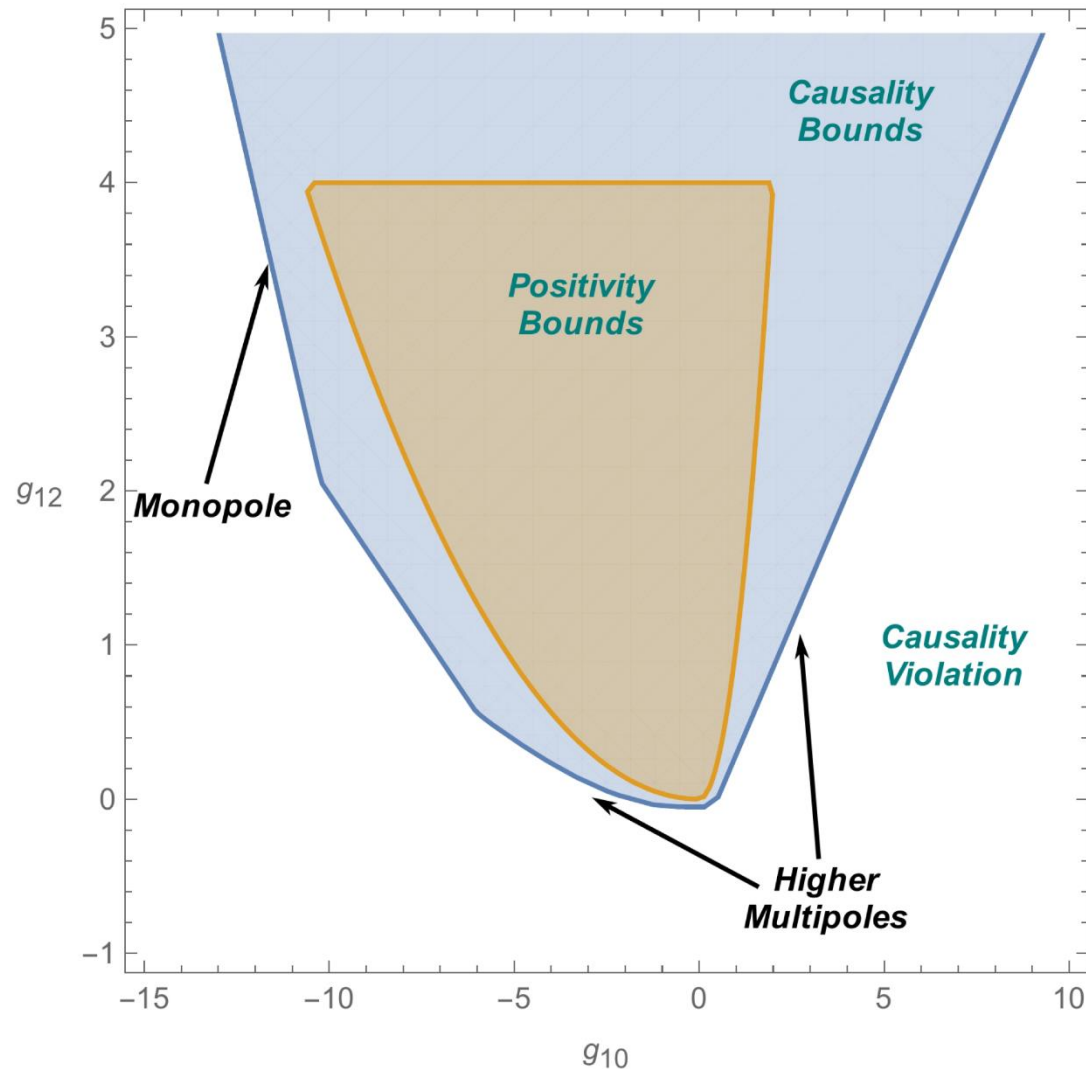
Asymptotic
Causality

$$v^2 \leq c^2 + \frac{\Lambda}{M_{\text{Pl}}}$$

Infrared
Causality

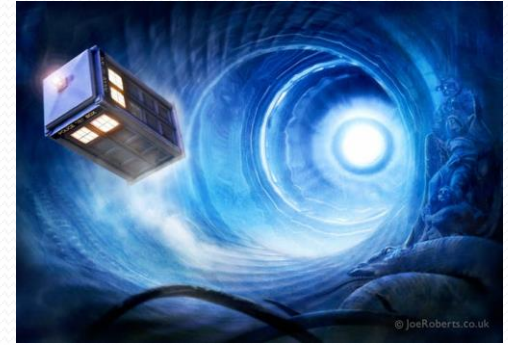
Positivity vs Causality

$$\mathcal{L}_{\text{EFT}} = -\frac{1}{2}(\partial\phi)^2 + \frac{c^2}{\Lambda^4}(\partial\phi)^4 + \frac{g_{10}}{\Lambda^6}(\partial\phi)^2(\phi_{,\mu\nu})^2 + \frac{g_{12}}{\Lambda^8}((\phi_{,\mu\nu})^2)^2 + \dots$$



$$c = 1$$

Tools for EFT of (massive) Gravity

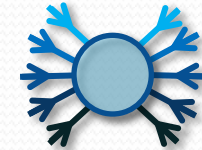
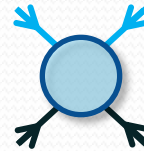


- Positivity bounds applied on low-energy EFTs can probe unknown UV contributions. Connection between UV and IR is made thanks to **analyticity** demanded from causality & **non-perturbative unitarity**.
- Application to SMEFT already proven powerful.
→ Generalization to bounds beyond 2-2 would be instrumental
- Applications to gravitational EFTs & EFT of Cosmology is more subtle
- Causality bounds applied directly at the level of the low-energy EFT can provide powerful complementary bounds
- Applicable to EFT of Gravity, Cosmology and generalizable beyond 2-2

Causality: Towards new Opportunities

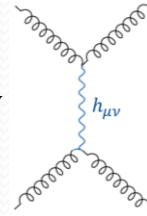
While powerful, current positivity bounds have not yet reached their full potential

1. Positivity **Limited to $2 \rightarrow 2$** scattering statements



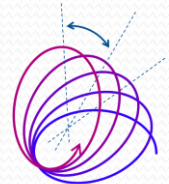
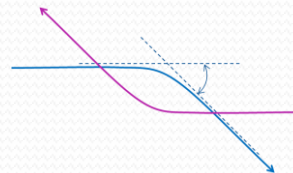
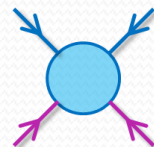
$$\partial^m \phi^n$$

2. **Gravitational pole spoils** full applicability

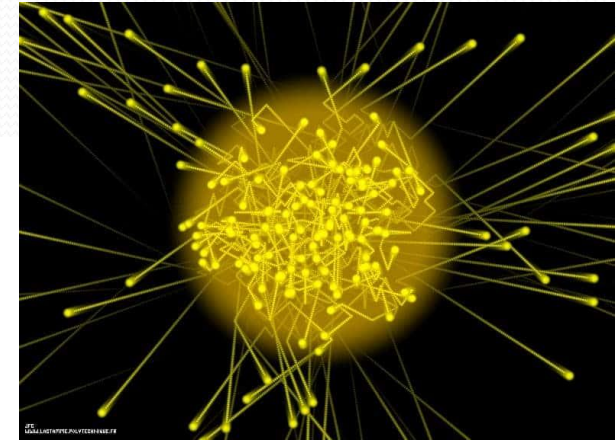


Current causality
can go beyond

3. **Need of more precise dictionary** with actual (gravitational) observables
real world, not *empty* & *Minkowski as required by S-matrix bootstrap bounds*



EFT of photons



$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \\
 & + \frac{c_1}{\Lambda^4}F^{\mu\nu}F_{\mu\nu}F^{\alpha\beta}F_{\alpha\beta} + \frac{c_2}{\Lambda^4}F^{\mu\nu}F^{\alpha\beta}F_{\mu\alpha}F_{\nu\beta} \\
 & + \frac{c_3}{\Lambda^6}F^{\alpha\mu}F^{\nu\beta}\partial_\mu F_{\beta\gamma}\partial_\nu F_{\alpha\gamma} + \frac{c_4}{\Lambda^6}F^{\alpha\mu}F^{\nu\beta}\partial_\beta F_{\mu\gamma}\partial^\gamma F_{\alpha\nu} + \frac{c_5}{\Lambda^6}F^{\alpha\mu}F^{\nu\beta}\partial_\beta F_{\mu\alpha}\partial_\nu F_{\beta\gamma} \\
 & + \frac{c_6}{\Lambda^8}F^{\mu\nu}F^{\alpha\beta}\partial_\mu\partial_\nu F_{\alpha\beta} + \frac{c_7}{\Lambda^8}F^{\mu\nu}F^{\alpha\beta}\partial_\mu\partial_\nu F_{\alpha\beta} + \frac{c_8}{\Lambda^8}F^{\mu\nu}F^{\alpha\beta}\partial_\mu\partial_\nu F_{\alpha\beta} + \dots \\
 & + \text{many higher order operators}
 \end{aligned}$$

From integrating out heavy charged fields (electron,...)

4-point tree-level scattering amplitude arising from this theory can be parameterized as

$$A_{++++} = f_2 (s^2 + t^2 + u^2) + f_3 stu + f_4 (s^2 + t^2 + u^2)^2$$

$$A_{++--} = g_2 s^2 + g_3 s^3 + g_4 s^4 + g'_4 s^2 tu$$

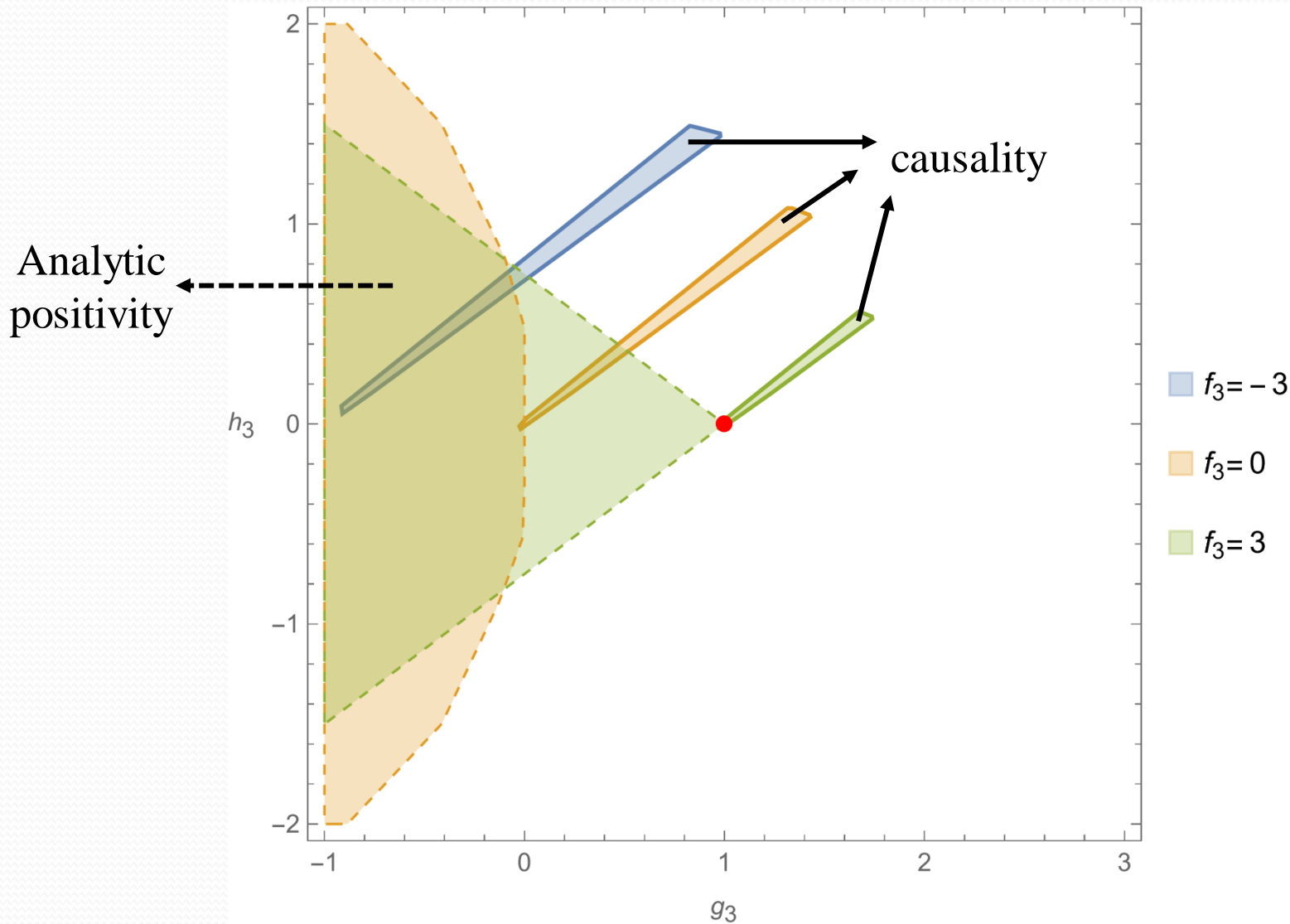
$$A_{++++-} = h_3 stu$$

$$f_2 = 2(4c_1 + c_2), \quad g_2 = 2(4c_1 + 3c_2)$$

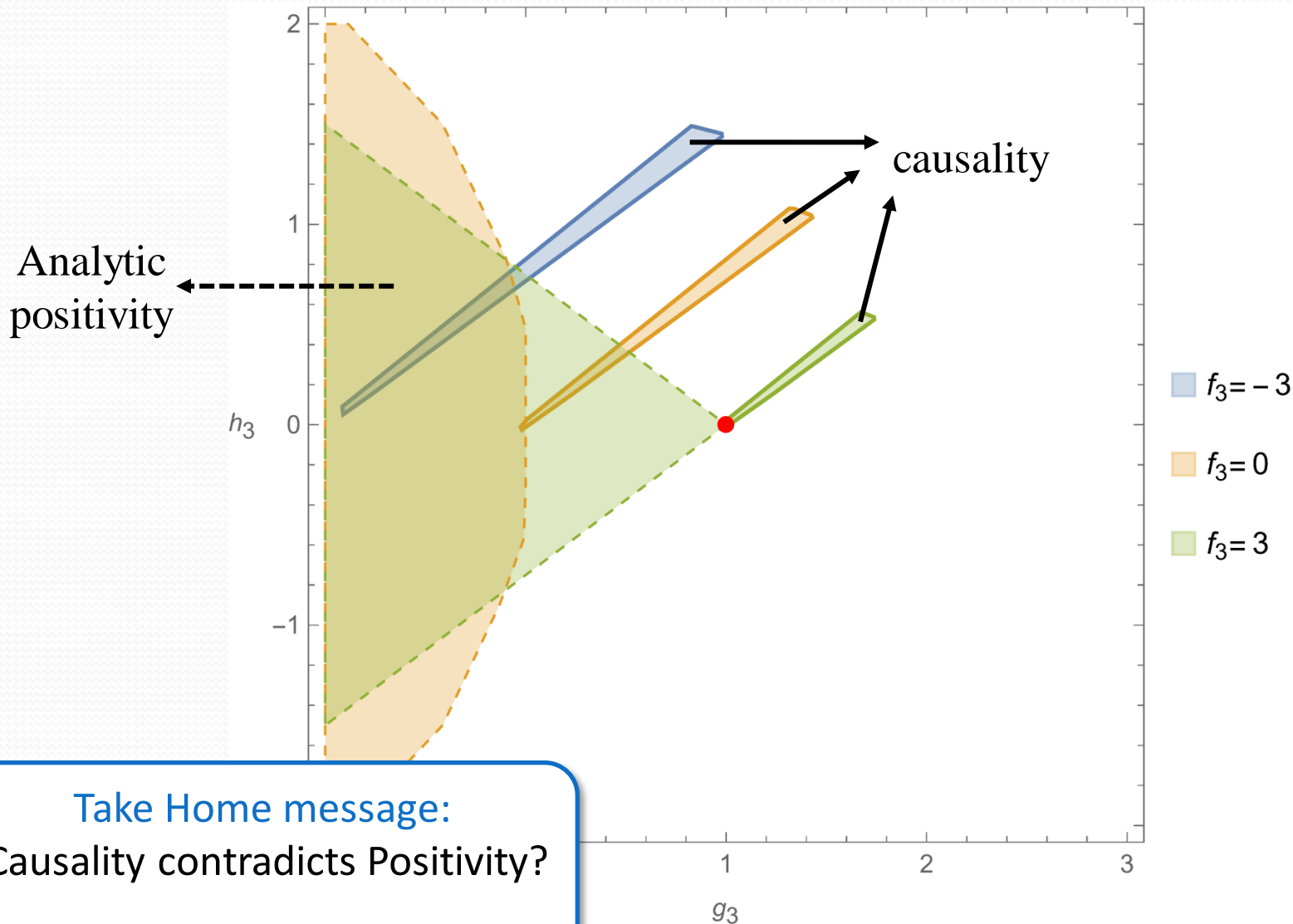
$$f_3 = -3(c_3 + c_4 + c_5), \quad g_3 = -c_5, \quad h_3 = -\frac{3}{2}c_3,$$

$$f_4 = \frac{1}{4}c_6, \quad g_4 = \frac{1}{2}(c_6 - c_8) + c_7, \quad g'_4 = -\frac{1}{2}(c_7 + c_8).$$

Positivity vs Causality in EFT of photons

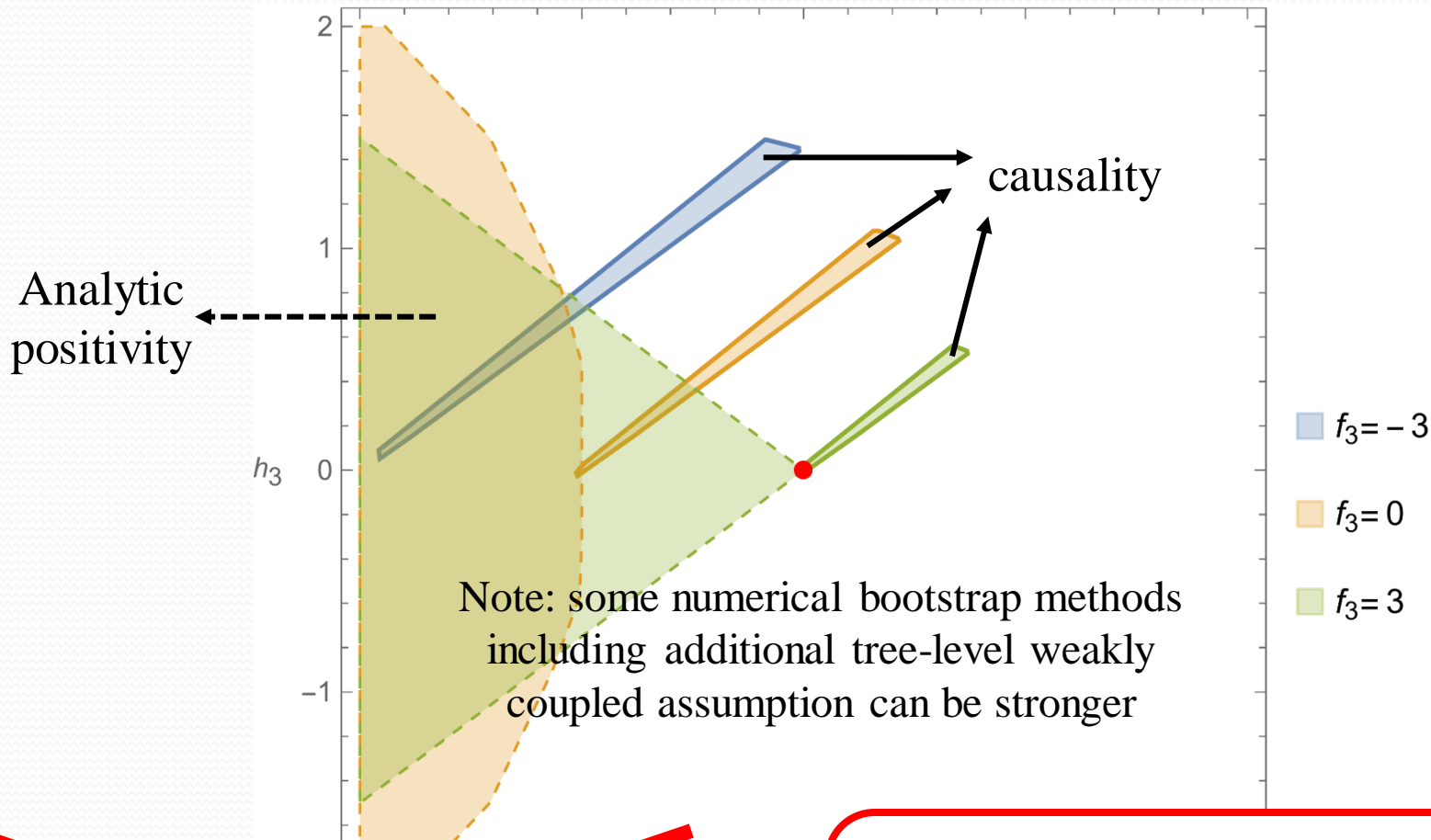


Positivity vs Causality in EFT of photons



Take Home message:
Causality contradicts Positivity?

Positivity vs Causality in EFT of photons



~~Take Home message:
Causality contradicts Positivity?
(NO!!!)~~

Take Home message:
Causality can probe uncharted territories,
not yet constrained by analytic positivity!
Combining both is powerful!

Constraints on Dark Energy

Example of Dark Energy model (quartic Horndeski)

$$\mathcal{L} = M_{\text{Pl}}^2 G_4(X) R + \Lambda^4 G_2(X) + \frac{1}{H_0^2} G_{4,X}(X) \left((\square\phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right) + \mathcal{L}_{\text{matter}}(\psi, g)$$

$$X = -\frac{1}{2\Lambda^4} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi$$

$$\Lambda^2 = H_0^2 M_{\text{Pl}}^2 \sim 10^{-3} \text{eV}$$

$$c_M = -\frac{2\dot{X}}{HM^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{Effective Planck Mass}$$

$$c_B = \frac{8X}{M^2\Omega_{\text{DE}}} (G_{4,X} + 2XG_{4,XX}), \quad \text{“Braiding” (scalar/tensor mixing)}$$

$$c_T = \frac{4X}{M^2\Omega_{\text{DE}}} G_{4,X}, \quad \text{Speed of GWs} - 1$$

$$M^2 = 2(G_4 - 2XG_{4,X})$$

Constraints on Dark Energy

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Assume a Λ CDM background.

Linear perturbations controlled by 4 background functions
(including kineticity related to G_2)

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Constrained performing a Markov chain Monte Carlo analysis, using:

- Planck 2015,
- CMB lensing and low- ℓ polarisation data,
- SDSS/BOSS BAO
- SDSS DR4 LRG matter power spectrum shape
- redshift space distortion constraints from BOSS and 6dF

Constraints on Dark Energy

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In addition, these parameters can be bounded using positivity bounds

Constraints on Dark Energy

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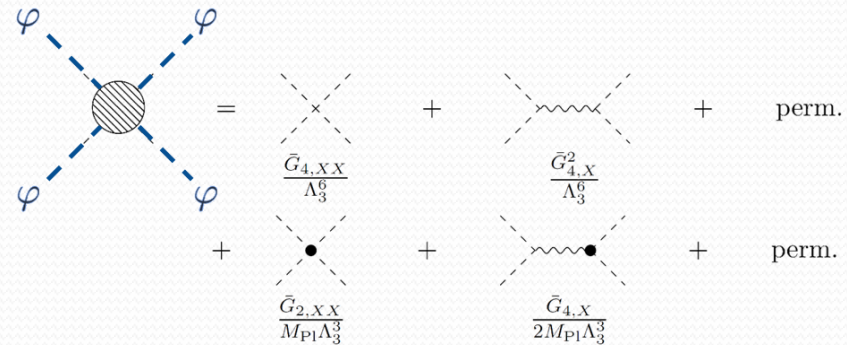
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Positivity Bounds from

$$\varphi\varphi \rightarrow \varphi\varphi$$



$$\mathcal{A}_{\varphi\varphi \rightarrow \varphi\varphi} = \frac{\bar{G}_{2,XX}}{\Lambda^4} (s^2 + t^2 + u^2) - \frac{\bar{G}_{4,X}}{\bar{G}_4 \Lambda^4} (su + st + ut) + 3! \frac{\bar{G}_{4,XX} + \bar{G}_{4,X}^2 / \bar{G}_4}{H_0^2 \Lambda^4} stu$$

$$\bar{G}_{2,XX} \geq -\bar{G}_{4,X} \frac{\bar{G}_{2,X}}{\bar{G}_4}, \quad \bar{G}_{4,XX} + \frac{\bar{G}_{4,X}^2}{\bar{G}_4} \leq 0 \quad \varphi\varphi \rightarrow \varphi\varphi$$

Constraints on Dark Energy

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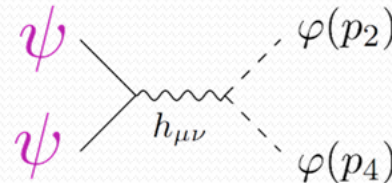
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Standard (minimal)
coupling to
matter fields ψ

Positivity Bounds from

$$\varphi\psi \rightarrow \varphi\psi$$



$$= \frac{\bar{G}_{4,X}}{2\Lambda_2^4} (s^2 + u^2 - t^2)$$



$$\bar{G}_{4,X} \geq 0 \quad \varphi\psi \rightarrow \varphi\psi$$

$$\bar{G}_{2,XX} \geq -\bar{G}_{4,X} \frac{\bar{G}_{2,X}}{\bar{G}_4}, \quad \bar{G}_{4,XX} + \frac{\bar{G}_{4,X}^2}{\bar{G}_4} \leq 0 \quad \varphi\varphi \rightarrow \varphi\varphi$$