

New results for thermal axion production in the early Universe

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SUBATECH

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Based on *K. Bouzoud, J. Ghiglieri - Thermal axion production at hard and soft momenta (2404.06113)*

Axion:

- Good dark matter candidate ($m_a \lesssim 10^{-1}$ eV from experimental bounds)
- Solves the “strong CP problem”

Kim-Shifman-Vainshtein-Zakharov (KSVZ) model

$$\mathcal{L}_{\text{int}} = \frac{g_3^2}{32\pi^2} \frac{a}{f_{\text{PQ}}} G \tilde{G} \quad (1)$$

f_{PQ} : energy scale below which the axion exists ($\geq 4 \times 10^8$ GeV from experimental bounds)

The axion only interacts with gluons

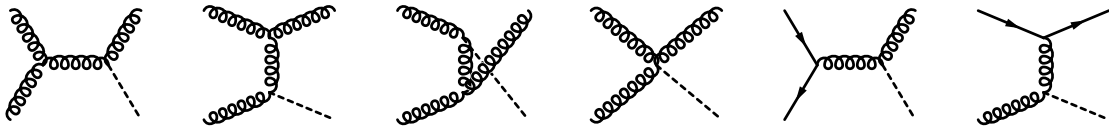


Figure: All tree-level processes in the KSVZ model producing **one** axion in the final state

If \mathcal{L}_{int} is realized in nature:

- Bose-Einstein condensate axion population \equiv dark matter
- Ultra-relativistic (“hot”) axion population \equiv *dark radiation* \rightarrow focus of this talk

Extra contribution to the effective number of neutrinos

Measure of the contribution of the hot axion population to the energy density of the Universe. Defined as:

$$\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{e_a}{e_\gamma} \Big|_{\text{CMB}} \quad (2)$$

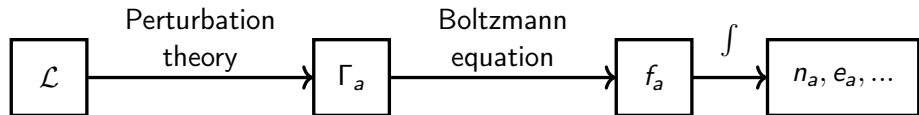


Figure: Phenomenological computation flowchart

Current results

From **Planck** at 2σ : $\Delta N_{\text{eff}} < 0.3$

Future experiments

Expected constraint (at 2σ):

- From **Simons Observatory**: $\Delta N_{\text{eff}} < 0.1$
- From **CMB-S4**: $\Delta N_{\text{eff}} < 0.06$
- From **CMB-HD**: $\Delta N_{\text{eff}} < 0.028$

→ More precise future experiments motivate more precise theoretical computations

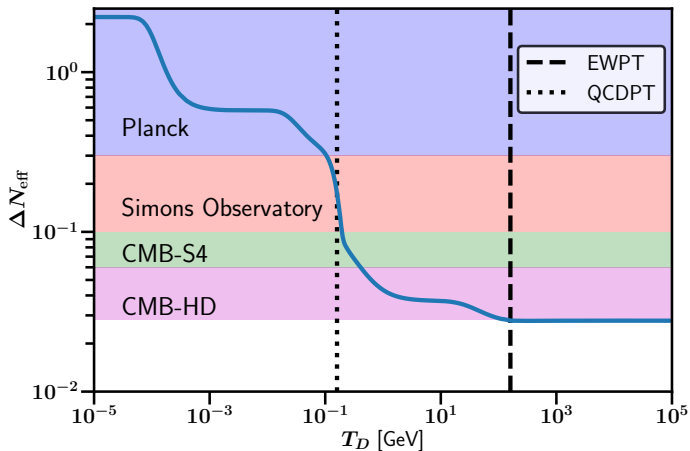


Figure: ΔN_{eff} as a function of the decoupling temperature for a BSM scalar

Phase space integral

 $2 \rightarrow 2$ scatterings

$$\Gamma_a(k) = \frac{1}{4k} \int d\Omega_{2 \rightarrow 2} \sum_{A,B,C} |\mathcal{M}_{A+B \rightarrow C+a}|^2 \frac{f_A(p_1) f_B(p_2) [1 \pm f_C(k_1)]}{n_B(k)} \quad (3)$$

Thermal distributions

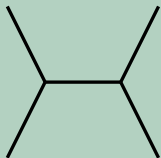
 $f_i \equiv n_B$ or n_F 

Figure: s channel

$|\mathcal{M}|^2 \propto s^{-1}$. $s \rightarrow 0$ limit is not problematic

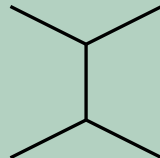


Figure: t channel

$|\mathcal{M}|^2 \propto t^{-1}$. Sensitivity to collective modes in $t \rightarrow 0$ limit (bc. $\lambda \rightarrow \infty$)

Different methods exist for implementing the collective effects. They converge for small values of g_3 (QCD coupling constants) but become extrapolations at larger g_3 .

The “spread” between those methods can be used to quantify the theoretical uncertainty.

We implemented three different methods:

“Strict LO”

Already implemented for the axion case in Peter Graf, Frank Daniel Steffen, *Thermal axion production in the primordial quark-gluon plasma* (1008.4528)

- Leads to large negative values

“Subtraction”

Previously implemented for the neutrino case in J. Ghiglieri, M. Laine, *Neutrino dynamics below the electroweak crossover* (1605.07720)

- Reduces the negativity but does not fix it entirely

“Tuned”

Previously implemented for YM dynamics in M. C. A. York, A. Kurkela, E. Lu, G. D. Moore, *UV Cascade in Classical Yang-Mills via Kinetic Theory* (1401.3751)

- No negative values

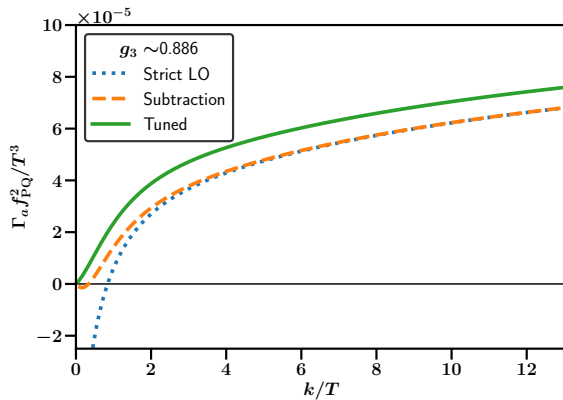


Figure: Axion production rate at $T = 10^4$ GeV

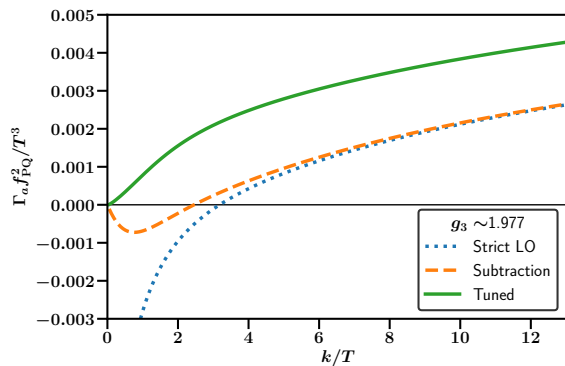
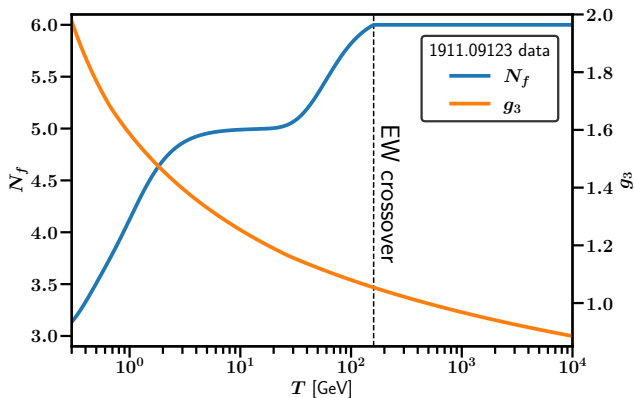


Figure: Axion production rate at $T = 0.3$ GeV

Time/temperature evolution

$$\partial_t f_a = \Gamma_a [n_B - f_a] \quad (4)$$

Creation/destruction



- Use $\Gamma_{\text{used}} \equiv \max(0, \Gamma_{\text{computed}})$.
- Initial condition:
 $f_a(T_{\text{max}}) = n_B(T_{\text{max}})$.

Beyond momentum-independent approximation

Solve for multiple momenta **instead of** using integrated *approximate* form giving energy density

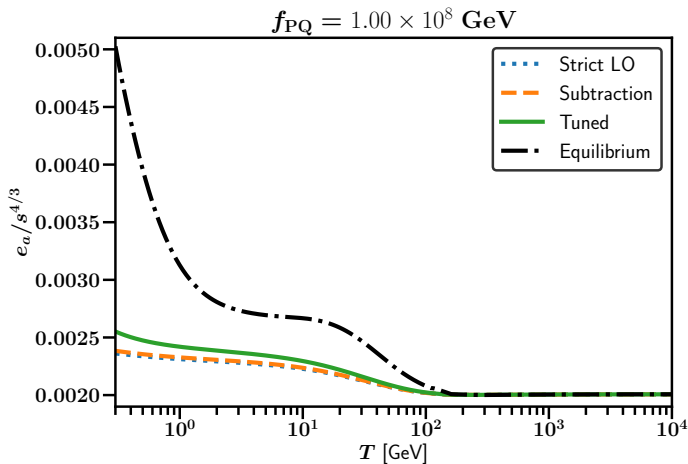
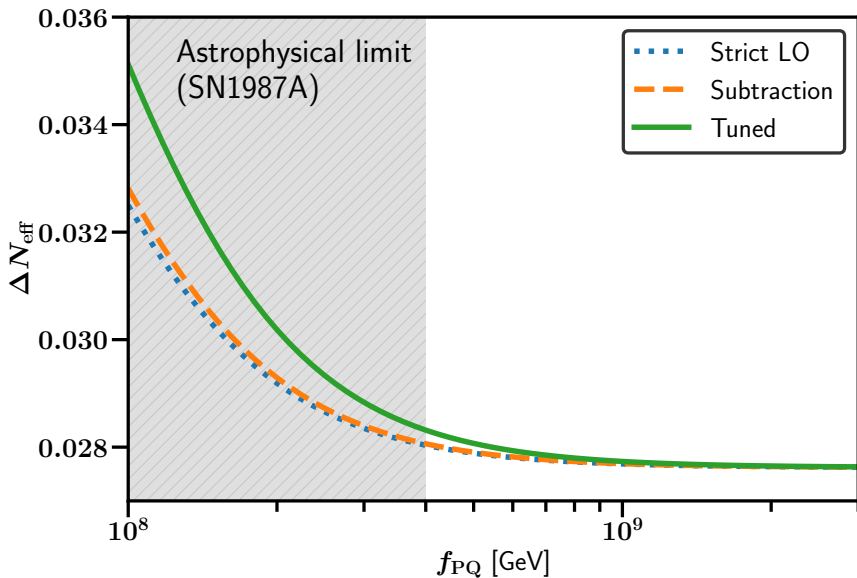


Figure: (Normalized) axion energy density

“Equilibrium” \equiv energy density for $f_a = n_B$



- Main source of theoretical uncertainty is the choice of production rate computation method (~ 0.002 at most)
- The error due to the momentum-independent approximation is smaller than that (~ 0.0005 at most)
- These conclusions might change if delayed production (at QCD transition or after) happens
- CMB-HD would be more constraining than the current astrophysical limit

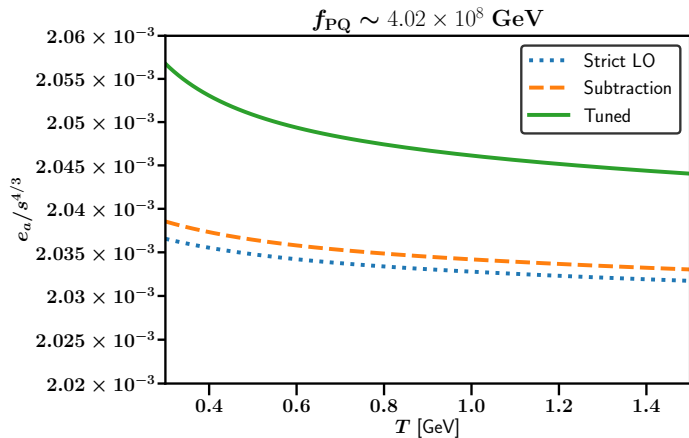


Figure: Normalized axion energy density for f_{PQ} close to the astrophysical limit

Inclusion of QCD transition effects

- Use effective number of colors $N_{C,\text{eff}}$
- Do the same analysis below QCDPT using the χPT Lagrangian and merge the two rates (similar to d'Eramo *et al.* in 2108.05371)

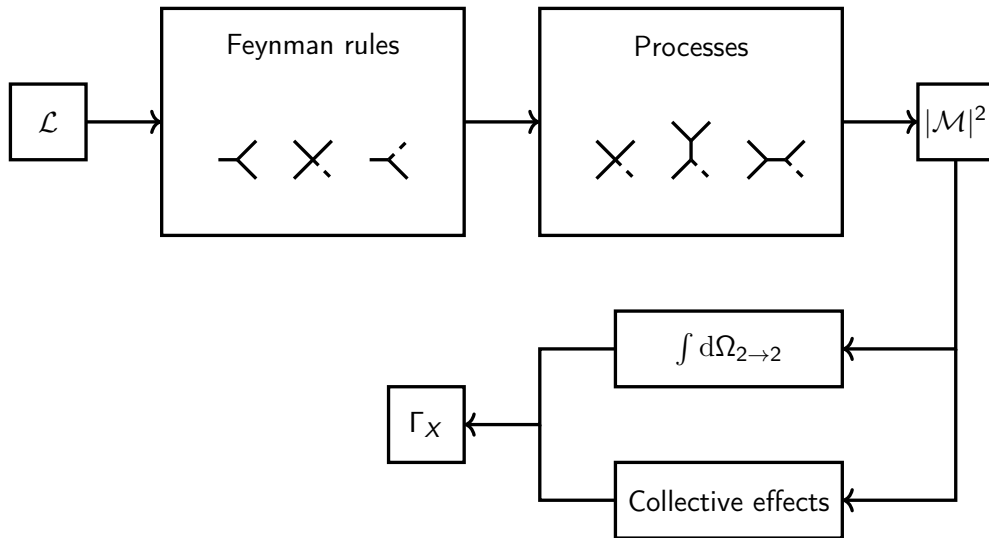


Figure: Flowchart of the production rate computation process

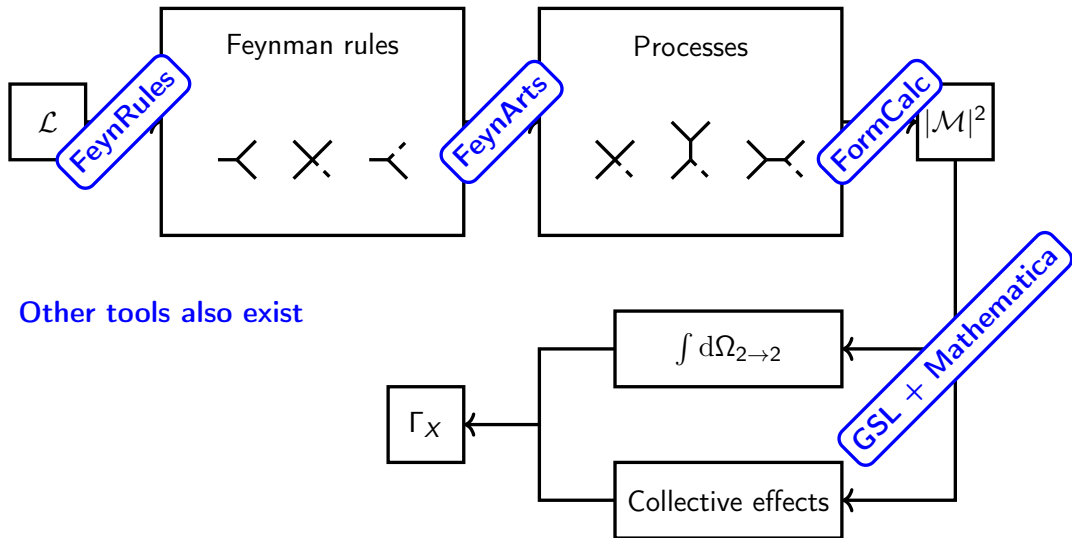
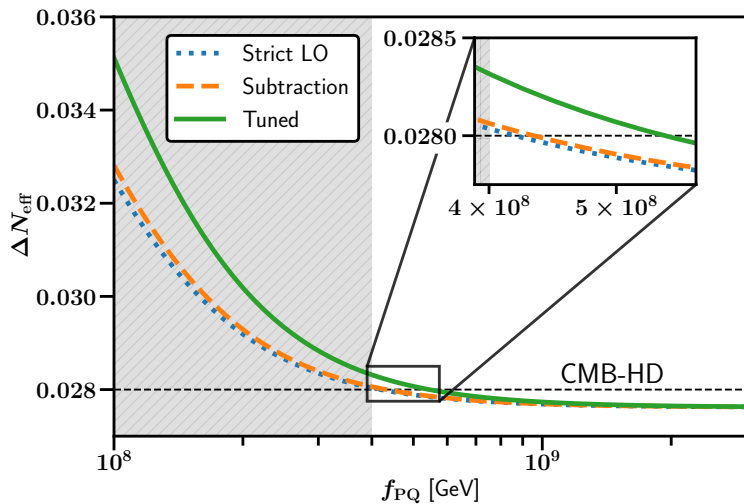
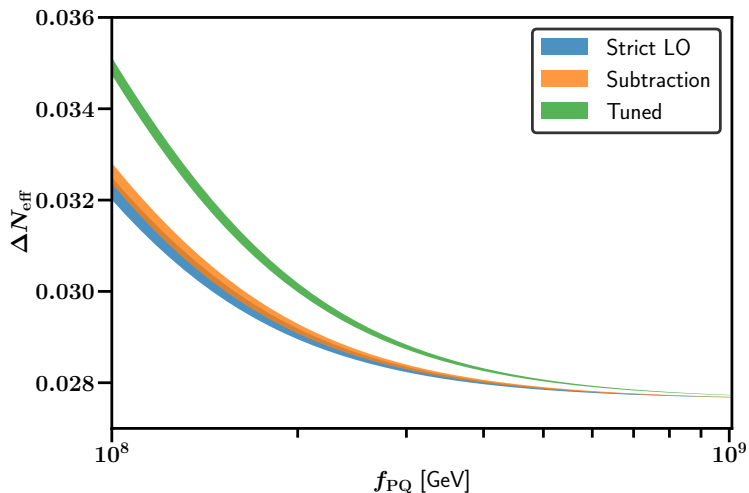
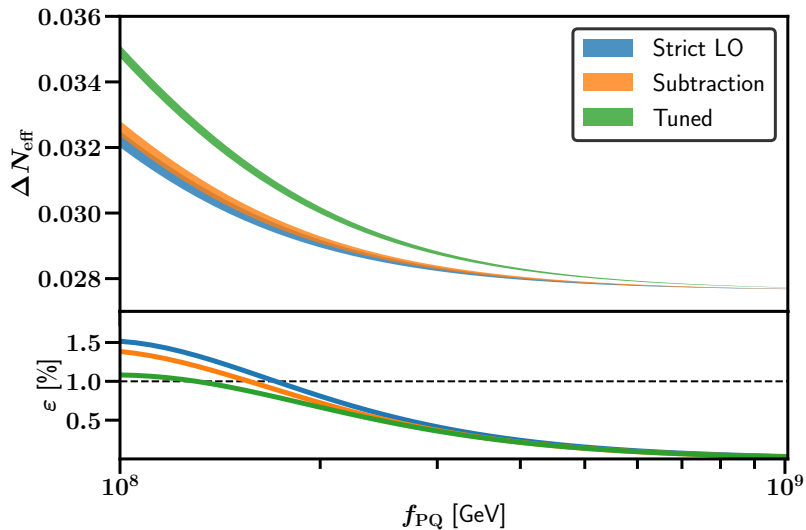


Figure: Flowchart of the production rate computation process



Momentum-dependent method vs. momentum-averaged approximation





Relative error

$$\epsilon = \frac{|\text{dep} - \text{avg}|}{\text{dep}} \quad (5)$$

- Planck Collaboration, *Planck 2018 results. VI. Cosmological Parameters* (1807.06209)
- Simons Observatory Collaboration, *The Simons Observatory: Science goals and forecasts* (1808.07445)
- CMB-S4 Collaboration, *Snowmass 2021 CMB-S4 White Paper* (2203.08024)
- CMB-HD Collaboration, *Snowmass 2021 CMB-HD White Paper* (2203.05728)
- Jihn E. Kim, *Weak-Interaction Singlet and Strong CP Invariance* (Phys. Rev. Lett. 43, 103)
- Mikhail A. Shifman, A.I. Vainshtein, Valentin I. Zakharov, *Can Confinement Ensure Natural CP Invariance of Strong Interactions* (Nucl. Phys. B 166)
- Pierluca Carenza, Tobias Fischer, Maurizio Giannotti, Gang Duo, Gabriel Matrinez-Pinedo, Alessandro Mirizzi, *Improved axion emissivity from a supernova via nucleon-nucleon bremsstrahlung* (1906.11844)