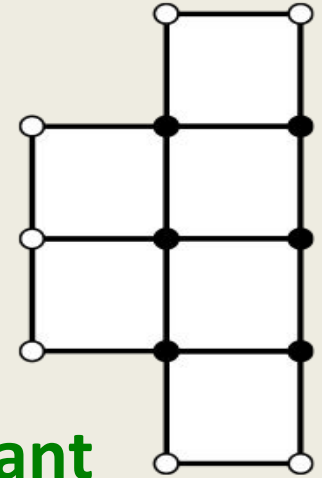
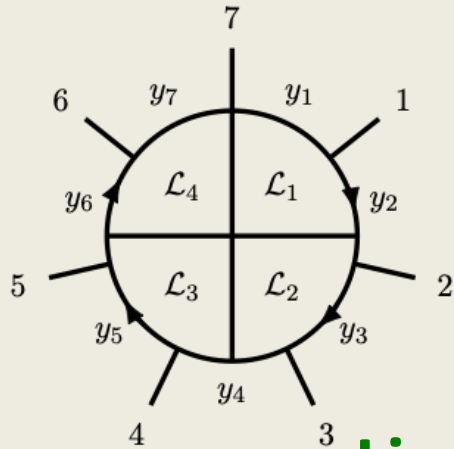


Landau Singularities of Ziggurat Graphs

Anastasia Volovich
Brown University



Lippstreu, Spradlin, Yelleshpur Srikant
2211.16425 + to appear



Prague Spring Amplitudes Workshop, May 2023

Outline

- 1. Some results that I will not talk about:**
 - EFTs with Celestial duals**
 - Coon amplitude**
 - One-loop n-gon in n-dim**
- 2. N=4 Yang-Mills symbol alphabet (review)**
- 3. Symbol alphabet from Landau equations**

1.1. EFTs with Celestial Duals

- **Strominger** showed that the algebra of conformally soft graviton operators in CCFT is w-algebra.
- We computed the modification of the w-algebra due to non-minimal couplings and found that the Jacobi identity is satisfied only when the coupling constants obey the constraints:

$$\begin{aligned}
 (\kappa_{-2,2,2} - \kappa_{0,0,2}) \kappa_{0,2,2} &= 0, & (\kappa_{-2,2,2} - \kappa_{0,0,2}) \kappa_{0,0,2} &= 0 \\
 3\kappa_{0,2,2}^2 &= 10 \kappa_{-2,2,2} \kappa_{2,2,2}.
 \end{aligned}$$

$$\begin{aligned}
 & \begin{array}{c} 2^{++} \\ \diagdown \\ \textcircled{\kappa_{-2,2,2}} \xrightarrow{3^{--}} \\ \diagup \\ 1^{++} \end{array} = \kappa_{-2,2,2} \frac{[12]^6}{[23]^2 [13]^2} \\
 & \begin{array}{c} 2^{++} \\ \diagdown \\ \textcircled{\kappa_{0,2,2}} \xrightarrow{3^\phi} \\ \diagup \\ 1^{++} \end{array} = \kappa_{0,2,2} [12]^4
 \end{aligned}$$

$$\begin{aligned}
 & \begin{array}{c} 2^{++} \\ \diagdown \\ \textcircled{\kappa_{2,2,2}} \xrightarrow{3^{++}} \\ \diagup \\ 1^{++} \end{array} = \kappa_{2,2,2} [12]^2 [23]^2 [13]^2 \\
 & \begin{array}{c} 2^\phi \\ \diagdown \\ \textcircled{\kappa_{0,0,2}} \xrightarrow{3^\phi} \\ \diagup \\ 1^{++} \end{array} = \kappa_{0,0,2} \frac{[12]^2 [13]^2}{[23]^2}
 \end{aligned}$$

1.1. EFTs with Celestial Duals

- These coupling constants constraints can be derived directly from the associativity of the celestial OPEs.
- All four-point amplitudes constructible solely from holomorphic or anti-holomorphic three-point amplitudes vanish on the support of these constraints: $\mathcal{A}_4(1^{++}, 2^{++}, 3^{++}, 4^{++}) = \mathcal{A}_4(1^{++}, 2^{++}, 3^{++}, 4^\phi)$
- All theories with $N \geq 1$ SUSY and stable vacua satisfy these constraints.

Ren, Spradlin, Yellespur Srikant, AV

Ball, Yellespur Srikant, AV, in progress



1.2. On Unitarity of the Coon Amplitude



- **Coon amplitude**

Bhardwaj, De, Spradlin, AV [2212.00764](#)

$$\mathcal{A}^C(s, t) = (1 - q) \exp\left(\frac{\log \sigma \log \tau}{\log q}\right) \prod_{n=0}^{\infty} \frac{\left(1 - \frac{q^n}{\sigma\tau}\right) (1 - q^{n+1})}{\left(1 - \frac{q^n}{\sigma}\right) \left(1 - \frac{q^n}{\tau}\right)},$$

$$\sigma = 1 - (1 - q)(\alpha_0 + s), \quad \tau = 1 - (1 - q)(\alpha_0 + t),$$

Remmen,
Maldacena,
Cheung, Geiser,
Lindwasser,
Figueroa, Tourkine
Chakravarty,
Maity, Mishra,
Jepsen

- **Residue polynomials on each massive pole admits a partial wave expansion in terms of D-dimensional Gegenbauer polynomials**

$$\mathcal{A}^C(s, t) \rightarrow \frac{1}{s - [n]_q} \times R_{[n]_q}^C(t)$$

$$R_{[n]_q}^C(t) = \sum_j B_{n,j}^{(D)}(q) G_j^{(D)}(t)$$

1.2. On Unitarity of the Coon Amplitude



- **Coon amplitude**

Bhardwaj, De, Spradlin, AV [2212.00764](#)

$$\mathcal{A}^C(s, t) = (1 - q) \exp\left(\frac{\log \sigma \log \tau}{\log q}\right) \prod_{n=0}^{\infty} \frac{\left(1 - \frac{q^n}{\sigma\tau}\right) (1 - q^{n+1})}{\left(1 - \frac{q^n}{\sigma}\right) \left(1 - \frac{q^n}{\tau}\right)},$$

$$\sigma = 1 - (1 - q)(\alpha_0 + s), \quad \tau = 1 - (1 - q)(\alpha_0 + t),$$

- We explored unitarity through partial wave expansion using **q-calculus**.
- Manifest positivity on leading ($j=n-1$) and sub-leading Regge ($j=n-2$) trajectories.
- Violation of unitarity for some (q, D) at sub-sub-leading Regge trajectories.

1.3. One-loop n-gon in n-dimensions

- The one loop hexagon integral in n-dim can be interpreted as the volume of a hyperbolic (n-1)-simplex.
- Such simplex can be dissected into *orthoschemes*, a special type of hyperbolic simplex for which the bounding hyperplanes can be ordered in such a way that they are mutually orthogonal.
- **Rudenko [2012.05599](#)**, showed that the volume of the orthoscheme can be written as the polylog function of cross-ratios.

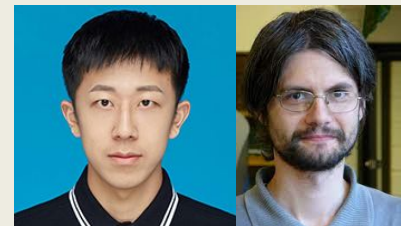
1.3. One-loop n-gon in n-dimensions

- An explicit formula for the volume of a hyperbolic orthoscheme is given in terms of alternating polylogs

$$\text{ALi}_{m_1, \dots, m_k}(\varphi_1, \dots, \varphi_k) := \sum_{\epsilon_1, \dots, \epsilon_k \in \{-1, 1\}} \left(\prod_{i=1}^k \frac{\epsilon_i}{2} \right) \text{Li}_{m_1, \dots, m_k}(\epsilon_1 \sqrt{\varphi_1}, \dots, \epsilon_k \sqrt{\varphi_k})$$

- We can compute the volume of our full simplex by slicing it into orthoschemes recursively.

Ren, Spradlin, Vergu AV to appear



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 - Coon amplitude
 - One-loop n -gon in n -dim
- 2. $N=4$ Yang-Mills symbol alphabet (review)**
- 3. Symbol alphabet from Landau equations**

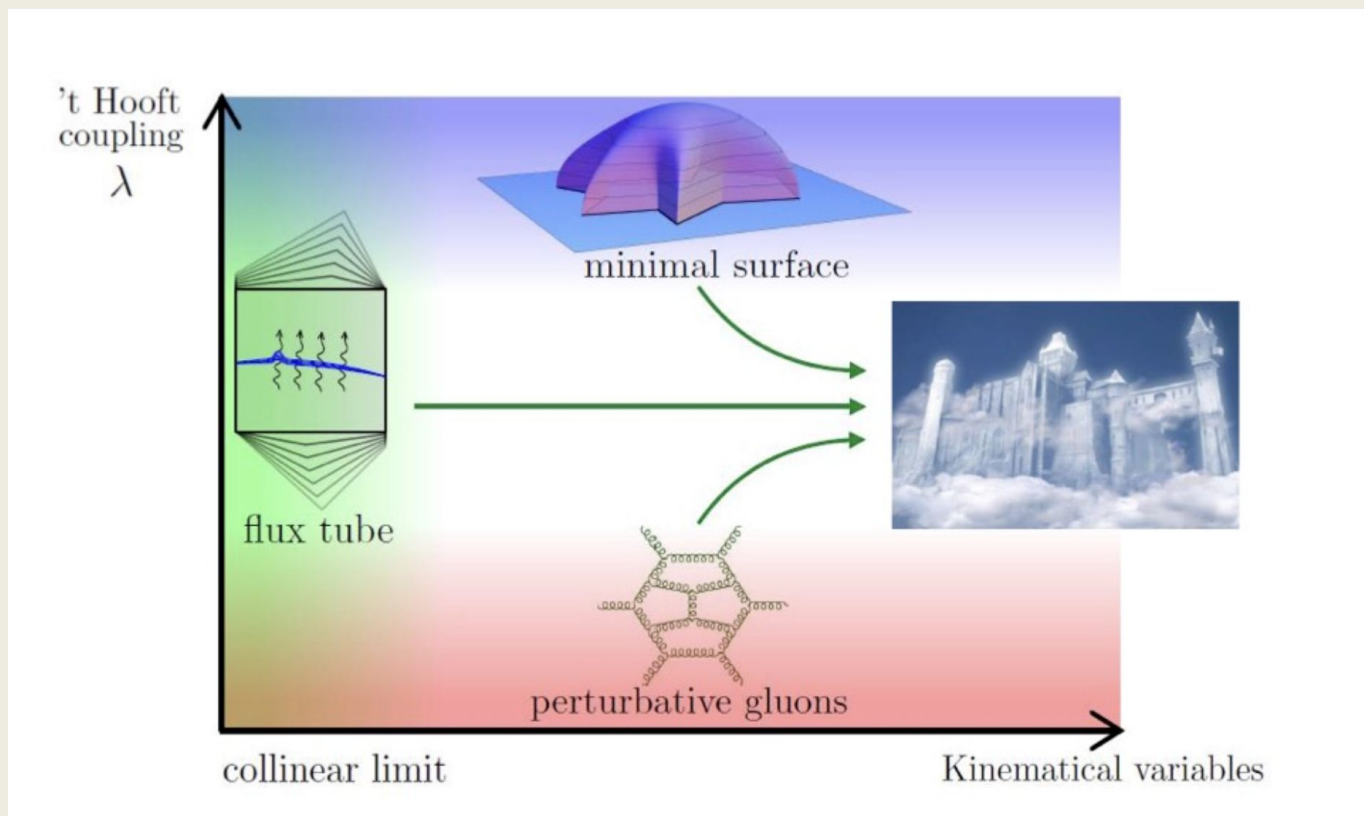
Planar N=4 Yang-Mills Amplitudes

- Planar N=4 Yang-Mills scattering amplitudes have been computed to very high loop order.
- They have many remarkable properties, that have sparked interest from mathematicians working on combinatorics, algebraic geometry and number theory.
- At the same time, several methods that have been developed for N=4 Yang-Mills are directly applicable to, and have greatly aided, computations relevant to actual experiments.

Solving N=4 super-Yang-Mills theory via Scattering Amplitudes: January 8, 2024 – March 8, 2024

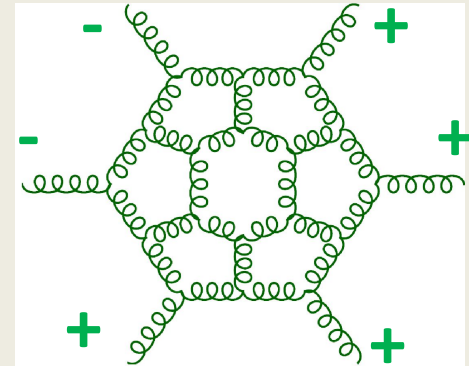
- Benjamin Basso (Ecole Normale Superieure),
- Lance Dixon (SLAC/Stanford U.)
- Jaroslav Trnka (UC Davis)
- Anastasia Volovich (Brown)

This program will bring together physicists and mathematicians with expertise in different facets of particle scattering in planar N=4 super-Yang-Mills theory, in order to try to solve the theory for generic values of the coupling and kinematical variables.



Status: N=4 Yang-Mills amplitudes

n-point L-loop MHV/NMHV



- $n < 6$ all loops [Bern, Dixon, Smirnov 0505205](#)
- $n = 6$ through 7-loops [Caron-Huot, Dixon, Drummond, Dulat, Foster, Gurdogan, von Hippel, McLeod, Papathanasiou, review: 2005.06735](#)
- $n = 7$ through 4-loops [Caron-Huot 1105.5606](#)
- All n MHV through 2-loops [Li, Zhang 2110.00350](#)
- $n = 8$ MHV through 3-loop [He, Li, Zhang 1911.01290](#)
- $n = 8, 9$ NMHV through 2-loops [2009.11471](#)

6 and 7-points: Amplitudes Bootstrap

Caron-Huot, Dixon, Drummond, Dulat, Foster, Gurdogan, von Hippel, Papathanasiou, [2005.06735](#)

Write down the answer as linear combo of functions and then determine the coefficients by solving a system of linear constraints.

Constraint	$L = 1$	$L = 2$	$L = 3$	$L = 4$	$L = 5$	$L = 6$
1. \mathcal{H}_6	6	27	105	372	1214	3692?
2. Symmetry	(2,4)	(7,16)	(22,56)	(66,190)	(197,602)	(567,1795?)
3. Final-entry	(1,1)	(4,3)	(11,6)	(30,16)	(85,39)	(236,102)
4. Collinear	(0,0)	(0,0)	(0*,0*)	(0*,2*)	(1* ³ ,5* ³)	(6* ² ,17* ²)
5. LL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0*,0*)	(1* ² ,2* ²)
6. NLL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0*,0*)	(1*,0* ²)
7. NNLL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0*)
8. N ³ LL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0)
9. Full MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0)
10. T^1 OPE	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0)
11. T^2 OPE	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)

Symbol Alphabet

- One input into the 6 and 7-points amplitudes bootstrap is the symbol alphabet

- **6-points:** 15 letters $\langle a \ a+1 \ b \ c \rangle$

- **7-points:** 49 letters $\langle a \ a+1 \ b \ c \rangle$

$\langle 1(23)(45)(67) \rangle$

$\langle 1(27)(34)(56) \rangle + \text{cyclic}$

Goncharov Spradlin Vergu AV

Golden Goncharov Spradlin Vergu AV

Caron-Huot

8, 9, ..n-points: \bar{Q} equation

$$\bar{Q}_a^A = \sum_{i=1}^n \chi_i^A \frac{\partial}{\partial Z_i^a}$$

$$\bar{Q}_a^A R_{n,k} = \frac{\Gamma_{\text{cusp}}}{4} \text{Res}_{\epsilon=0} \int_{\tau=0}^{\tau=\infty} \left(d^{2|3} \mathcal{Z}_{n+1} \right)_a^A [R_{n+1,k+1} - R_{n,k} R_{n+1,1}^{\text{tree}}] + \text{cyclic}$$

Caron-Huot [1105.5606](#)

He, Li, Zhang [1911.01290](#)

[2009.11471](#)

Li, Zhang [2110.00350](#)

8-point symbol alphabet

180+24 RATIONAL LETTERS

Li, Zhang [2110.00350](#)

- $\binom{8}{4} - 2 = 68$: all $\langle abcd \rangle$ except $\langle 1357 \rangle$ and $\langle 2468 \rangle$;
- 1 cyclic class of $\langle 12(345) \cap (678) \rangle$;
- 7 cyclic classes of $\langle 1(ij)(kl)(mn) \rangle$ with $2 \leq i < j < k < l < m < n \leq 8$;
5 cyclic classes of $\langle 1(28)(kl)(mn) \rangle$ with $2 < k < l < m < n < 8$;
- 5 cyclic classes of $\langle \bar{2} \cap \bar{4} \cap (568) \cap \bar{8} \rangle$, $\langle \bar{2} \cap \bar{4} \cap \bar{6} \cap (681) \rangle$, $\langle (127) \cap (235) \cap \bar{5} \cap \bar{7} \rangle$, $\langle (127) \cap \bar{3} \cap (356) \cap \bar{7} \rangle$, $\langle \bar{2} \cap (278) \cap (346) \cap \bar{6} \rangle$.

$$\bar{a} = (a-1, a, a+1)$$

$$\langle a(bc)(de)(fg) \rangle := \langle abde \rangle \langle acfg \rangle - \langle acde \rangle \langle abfg \rangle,$$

$$\langle ab(cde) \cap (fgh) \rangle := \langle abde \rangle \langle cfgh \rangle + \langle abec \rangle \langle dfgh \rangle + \langle abcd \rangle \langle efgh \rangle,$$

$$\langle (a_1b_1c_1) \cap (a_2b_2c_2) \cap (a_3b_3c_3) \cap (a_4b_4c_4) \rangle := \langle (a_1b_1c_1) \cap (a_2b_2c_2), (a_3b_3c_3) \cap (a_4b_4c_4) \rangle$$

2 x 9 ALGEBRAIC LETTERS (SQUARE ROOTS)

$$\Delta_{1357} = (\langle 1256 \rangle \langle 3478 \rangle - \langle 1278 \rangle \langle 3456 \rangle - \langle 1234 \rangle \langle 5678 \rangle)^2 - 4 \langle 1234 \rangle \langle 3456 \rangle \langle 5678 \rangle \langle 1278 \rangle$$

+ 1 cyclic

9-point symbol alphabet

He, Li, Zhang [2009.11471](#)

59 x 9 RATIONAL LETTERS

- 13 cyclic classes of $\langle 12kl \rangle$ for $3 \leq k < l \leq 8$ but $(k, l) \neq (6, 7), (7, 8)$;
- 7 cyclic classes of $\langle 12(ijk) \cap (lmn) \rangle$ for $3 \leq i < j < k < l < m < n \leq 9$;
- 8 cyclic classes of $\langle \bar{2} \cap (245) \cap \bar{6} \cap (691) \rangle$, $\langle \bar{2} \cap (346) \cap \bar{6} \cap (892) \rangle$, $\langle \bar{2} \cap (346) \cap \bar{2} \cap (782) \rangle$, $\langle \bar{2} \cap (245) \cap \bar{7} \cap (791) \rangle$, $\langle \bar{2} \cap (245) \cap (568) \cap \bar{8} \rangle$, $\langle \bar{2} \cap (245) \cap (569) \cap \bar{9} \rangle$, $\langle \bar{2} \cap (245) \cap (679) \cap \bar{9} \rangle$, $\langle \bar{2} \cap (256) \cap (679) \cap \bar{9} \rangle$;
- 10 cyclic classes of $\langle 1(ii+1)(jj+1)(kk+1) \rangle$ for $2 \leq i, i+1 < j, j+1 < k \leq 8$;
- 6 cyclic classes $\langle 1(2i)(jj+1)(k9) \rangle$ for $3 \leq i < j, j+1 < k \leq 8$, but $(i, k) \neq (3, 8), (4, 7)$;
- 14 cyclic classes of $\langle 1(29)(ij)(kk+1) \rangle$ for $3 < i < j \leq 8, 3 \leq k \leq i-2$ or $j+1 \leq k \leq 7$;
- 1 cyclic class of $\langle 1, (56) \cap \bar{3}, (78) \cap \bar{3}, 9 \rangle$.

11 x 9 ALGEBRAIC LETTERS (SQUARE ROOTS)

$$\Delta_{1357} = (\langle 1256 \rangle \langle 3478 \rangle - \langle 1278 \rangle \langle 3456 \rangle - \langle 1234 \rangle \langle 5678 \rangle)^2 - 4 \langle 1234 \rangle \langle 3456 \rangle \langle 5678 \rangle \langle 1278 \rangle + 8 \text{ cyclic}$$

Symbol Letters from Mathematics

- **Cluster Algebras** Golden, Goncharov, Spradlin, Vergu, AV
- **Tropical Geometry**
- **Dual Polytopes** Drummond, Foster, Gurdogan, Kalousios; Henke, Papathanasiou; Arkani-Hamed, Lam, Spradlin
- **Plabic Graphs** Mago, Schreiber, Spradlin, Srikant, AV; He, Li
- **Tensor Diagrams** Ren, Spradlin, AV
- **Scattering Diagrams** Herderschee
- **Schubert Problem** Yang
- **A-determinant** Dlapa, Helmer, Papathanasiou, Teller

Symbol Letters from Physics

Can we have a general derivation of these letters?

Symbol letters encode information about the location of singularities of an amplitude.

Landau formulated the sufficient conditions for a Feynman integral to have singularities decades ago...

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Landau Singularities

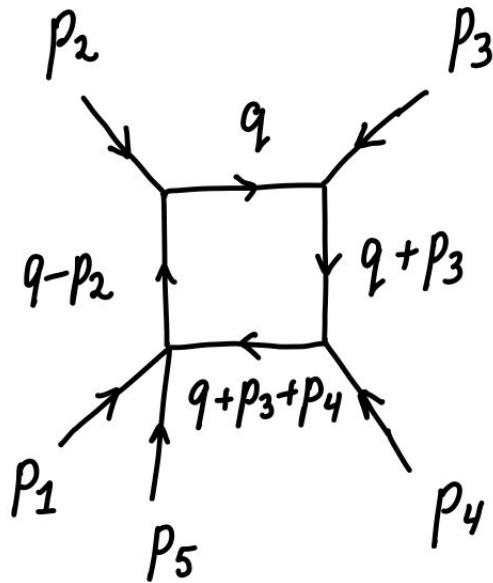
Landau 1959, ELOP

**Landau Equations for a given Feynman integral –
set of kinematic constraints that are necessary for
the appearance of a pole or a branch point in the
integrated function**

$$\sum_{i \in \text{loop}} \alpha_i q_i^\mu = 0 \quad \forall \text{ loops},$$
$$\alpha_i (q_i^2 - m_i^2) = 0 \quad \forall i.$$

**Landau Singularities occur for external momenta
such that the Landau equations have solutions.**

Example



$$\int \frac{d^4 q}{q^2 (q-p_2)^2 (q+p_3)^2 (q+p_3+p_4)^2}$$

$$\begin{cases} (q-p_2)^2 = q^2 = (q+p_3)^2 = (q+p_3+p_4)^2 \\ \alpha_1 (q-p_2) + \alpha_2 q + \alpha_3 (q+p_3) + \alpha_4 (q+p_3+p_4) = 0 \end{cases}$$

$$(p_2+p_3)^2 (p_3+p_4)^2 = 0$$

OR

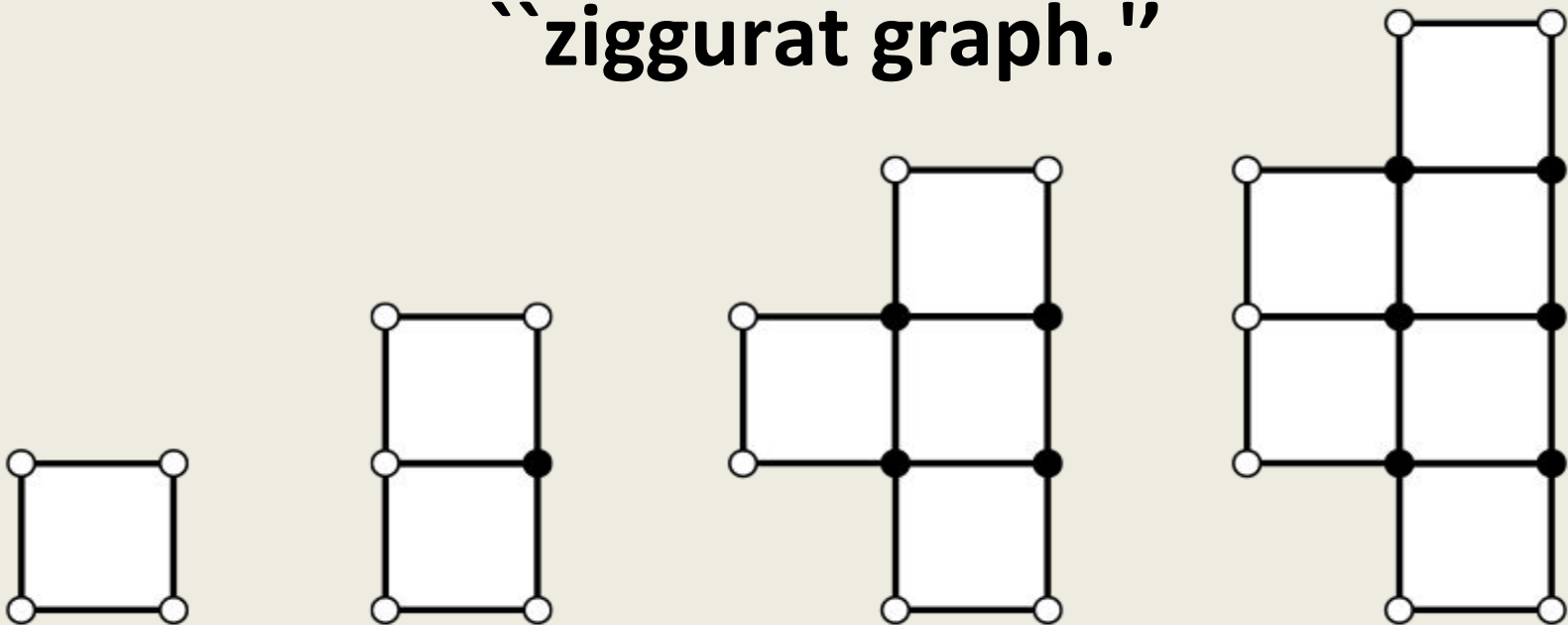
$$\langle 1234 \rangle \langle 2345 \rangle = 0$$



Ziggurat Graphs

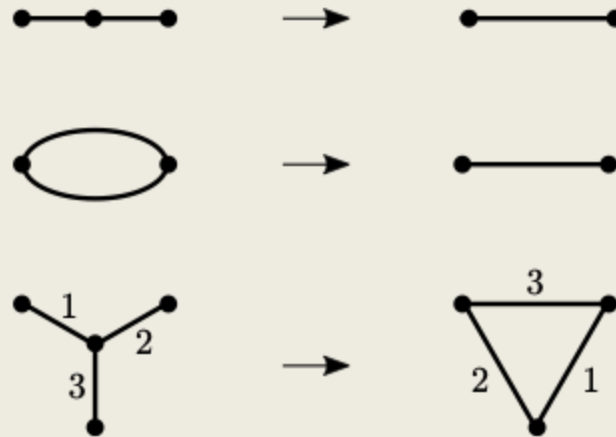
Prlina, Spradlin, Stanojevic

Landau singularities of any n-particle amplitude in any massless, planar theory are a subset of those of a special type of “ziggurat graph.”



Landau singularities from Ziggurats

- Landau equations are invariant under circuit moves



Prlina, Spradlin, Stanojevic

- Any graph can be reduced to ziggurat (or a relaxation) by these moves.

Gitler

Six-point Ziggurat

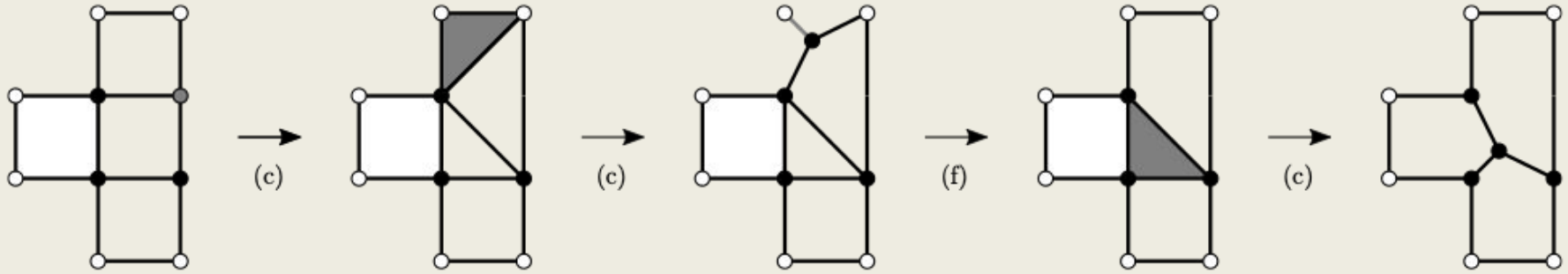
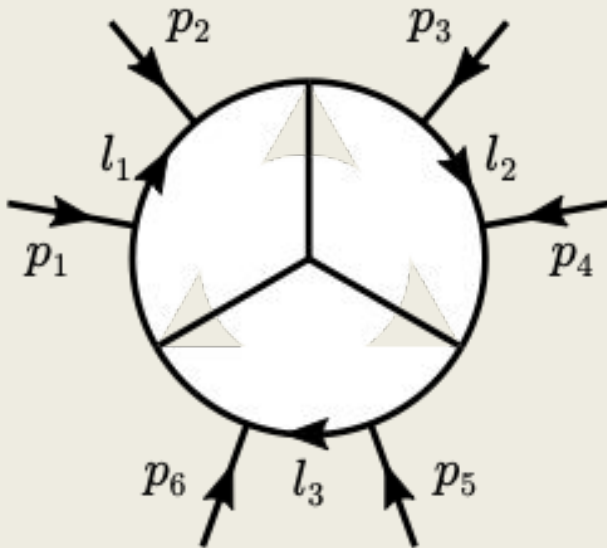


FIG. 3. The six-terminal ziggurat graph can be reduced to a three loop graph by a sequence of three $Y-\Delta$ reductions and one FP assignment. In each case the vertex, edge, or face to be transformed is highlighted in gray.



Matches six-point symbol letters!



Seven-point Ziggurat

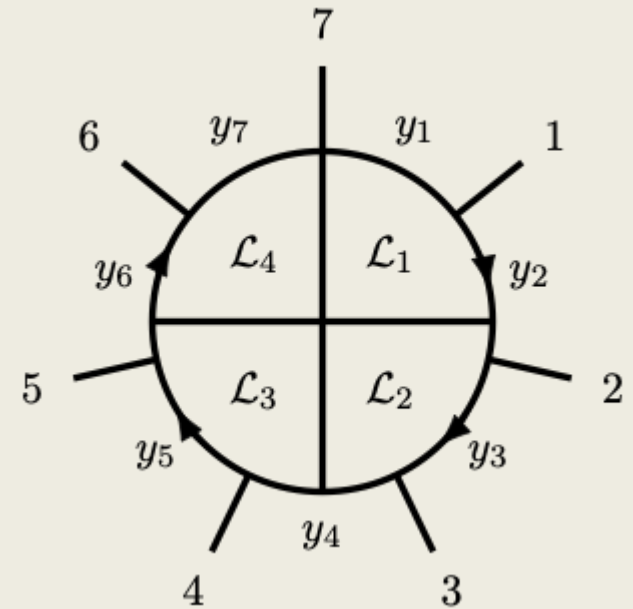
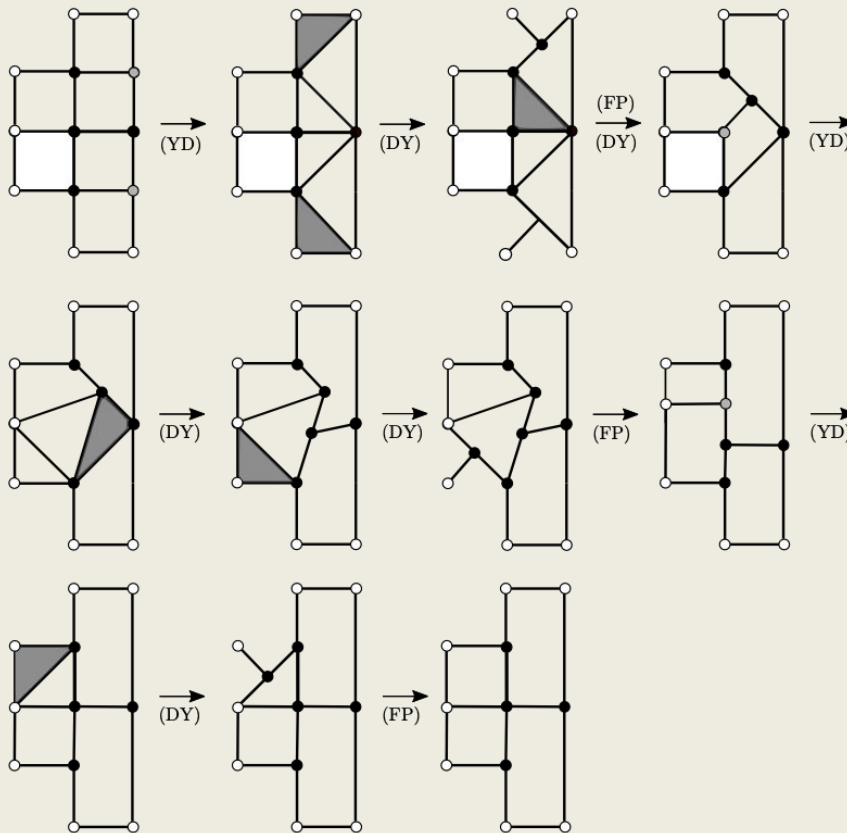


Figure 3: A sequence of graphical moves (see [15]) that transforms the 7-point ziggurat graph Fig. 1(d) into the wheel graph Fig. 2(b). YD indicates wye-delta transformation(s) on the node(s) shaded in grey; DY indicates delta-wye transformations(s) on the triangle(s) shaded in grey, and FP indicates a trivial contraction of external edges.

Seven-point Ziggurat

- We computed leading Landau singularities and found agreement with all but 7 letters.

Lippstreu, Spradlin, AV

- We are now analysing subleading Landau singularities.

Lippstreu, Spradlin, Yelleshpur Srikant, AV, to appear



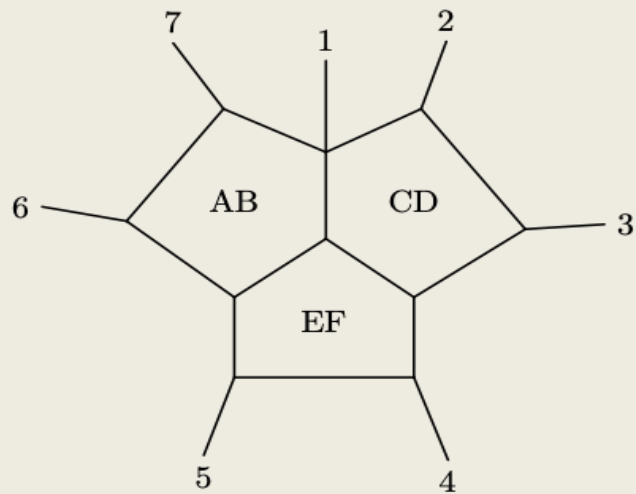


Figure 2: Graph 1. The Landau equations contain the letter $\langle 6(17)(25)(34) \rangle$, which is outside the heptagon symbol alphabet.

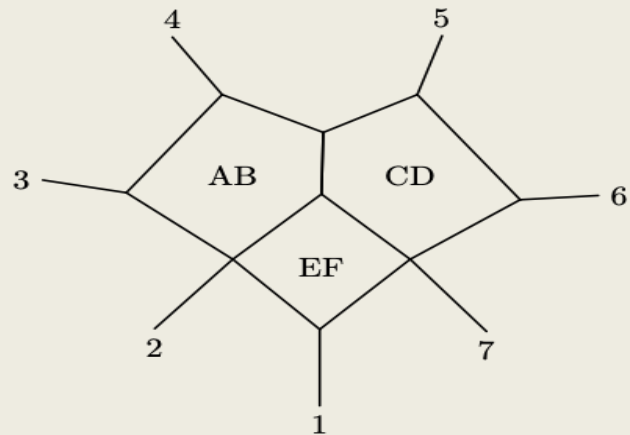
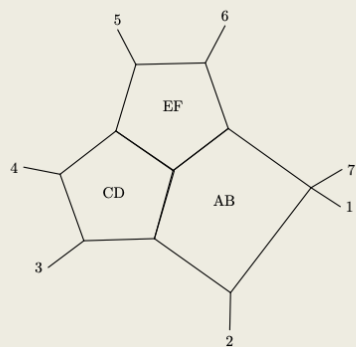
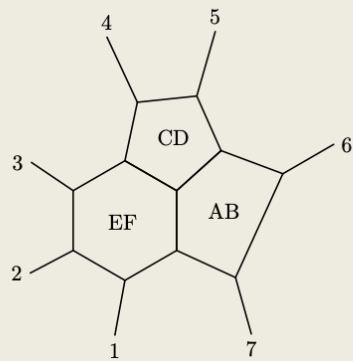


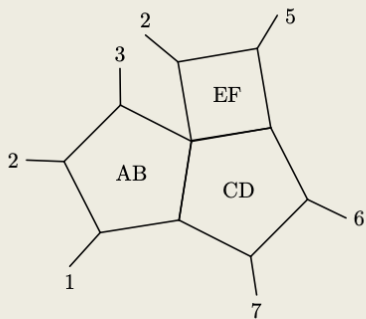
Figure 3: Graph 2. The Landau equations contain the letter $\langle 1(27)(36)(45) \rangle$, which is outside the heptagon symbol alphabet.



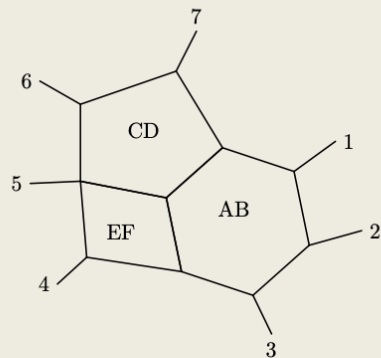
(a) Graph 5



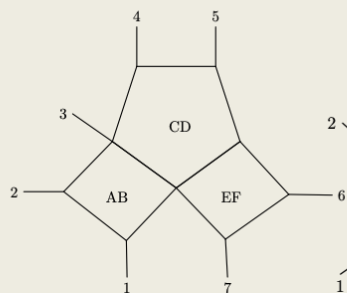
(b) Graph 3



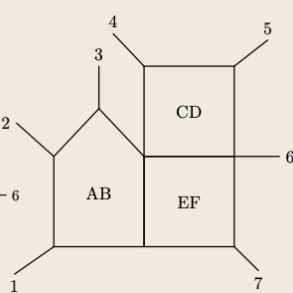
(c) Graph 4



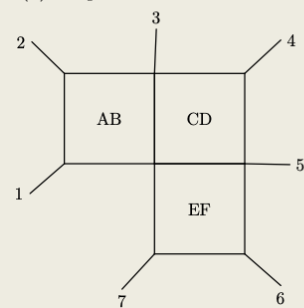
(d) Graph 6



(e) Graph 9



(f) Graph 8



(g) Graph 7

Figure 6: Seven other three-loop graphs we analyzed and did not find any letters beyond the heptagon symbol alphabet.

Seven-point Ziggurat

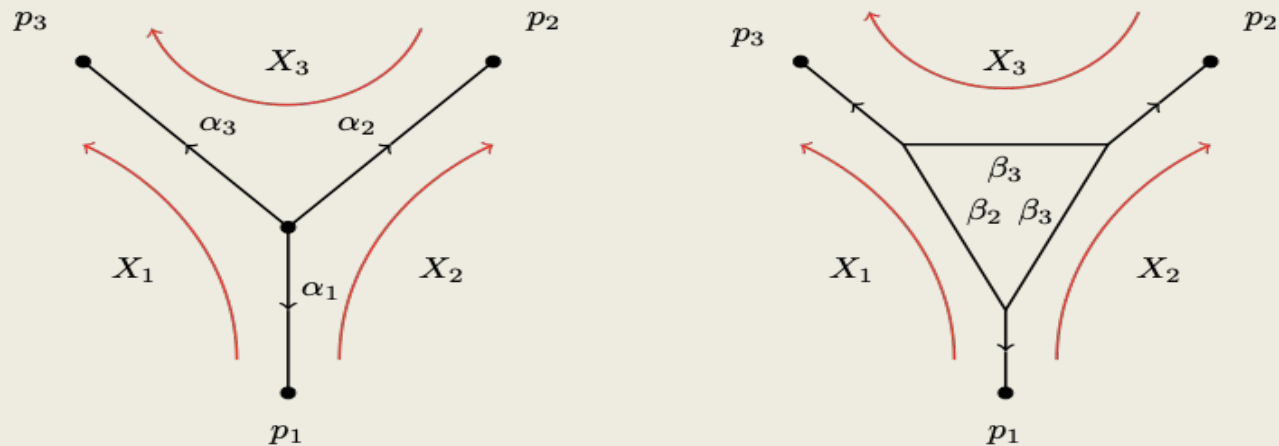
- We computed leading Landau singularities and find agreement with all but 7 letters.
- We are now analysing subleading Landau singularities.
- We found the missing 7, and extra letters $\langle 6(17)(25)(34) \rangle$ and $\langle 1(27)(36)(45) \rangle$
- They appear in certain 3-loop Feynman integrals and it will be interesting to see if they appear in higher loop amplitudes.

Lippstreu, Spradlin, Yellespur Srikant, AV, to appear

Violation of Y->Delta

Lippstreu, Spradlin, Yelleshpur Srikant, AV, to appear

- Along the way, we found that there are some solutions that violate Y-Delta equivalence.



- When $X_1 = X_2 = X_3 = 0$, the Y has extra solutions that are not present in the Delta.

Conclusions

- Planar $N=4$ Yang-Mills amplitudes have been computed to a very high order via bootstrap.
- Their singularity structure is an important input to the bootstrap.
- We were hoping to derive all n singularities from Landau analysis.
- Future: all n , all L , beyond polylogs, beyond planar, beyond $N=4$