



PennState

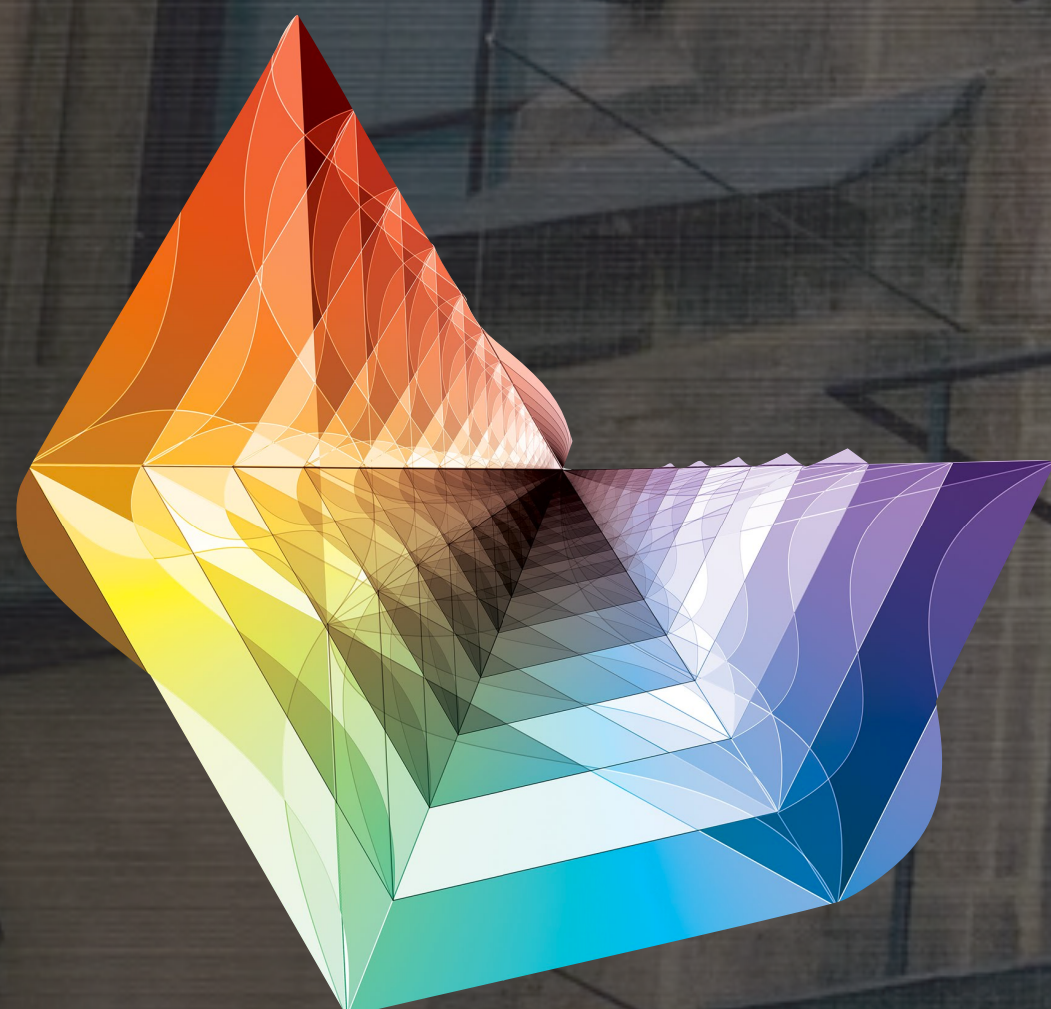
Adventures in Perturbation Theory

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Roadmap



◆ *Spiritus Movens*

- ▶ observables in QFT *surprisingly* simple!
- ▶ what do amplitudes *look like*—functionally?
- ▶ why is perturbation theory *so hard*?
—and how can we make it *easier*?

◆ Generalized Unitarity: *a modern perspective*

- ▶ *stratifying* theories and stratified Feynman integrand bases
- ▶ what makes for a *good* Feynman integral (basis)?

◆ Prescriptivity, Purity, and Polylogarithmicity

- ▶ Impurities, Calabi-Yau Geometries,...; tensions and resolutions



What Form do Observables Take?

- In a general (say, 4d) QFT, it would have *recently been* expected by most that observables took the following general form:

$$\mathcal{A} = \mathcal{A}^{\text{tree}} + \hbar \mathcal{A}^{(L=1)} + \hbar^2 \mathcal{A}^{(L=2)} + \dots + \hbar^L \mathcal{A}^{(L)} + \dots$$

(general dimension d : $2 \mapsto \lfloor d/2 \rfloor$)

rational	+	$\left(\begin{array}{c} \text{weight-2} \\ \text{polylogs} \end{array} \right)$	+	$\left(\begin{array}{c} \text{weight-4} \\ \text{polylogs} \end{array} \right)$	+	\dots	+	$\left(\begin{array}{c} \text{weight-2L} \\ \text{polylogs} \end{array} \right)$	+	\dots
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+	$\left(\begin{array}{c} \text{weight-1} \\ \text{polylogs} \end{array} \right)$	+	\vdots	+	\dots	+	\vdots	planar $\mathcal{N}=4$?
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+	rational	+	$\left(\begin{array}{c} \text{weight-1} \\ \text{polylogs} \end{array} \right)$	+	\dots	+	\vdots
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coefficients: weight-0, + rational

leading singularities (gauge invariant, etc.)



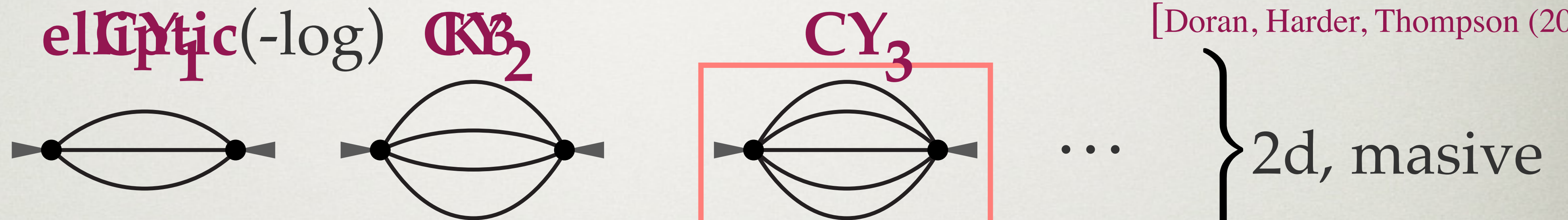
What Form do Observables Take?

◆ Unfortunately, many pesky counterexamples were to be found:

[Bloch, Kerr, Vanhove; Broadhurst;...]

[Doran, Harder, Thompson (2019)]

sunrises:



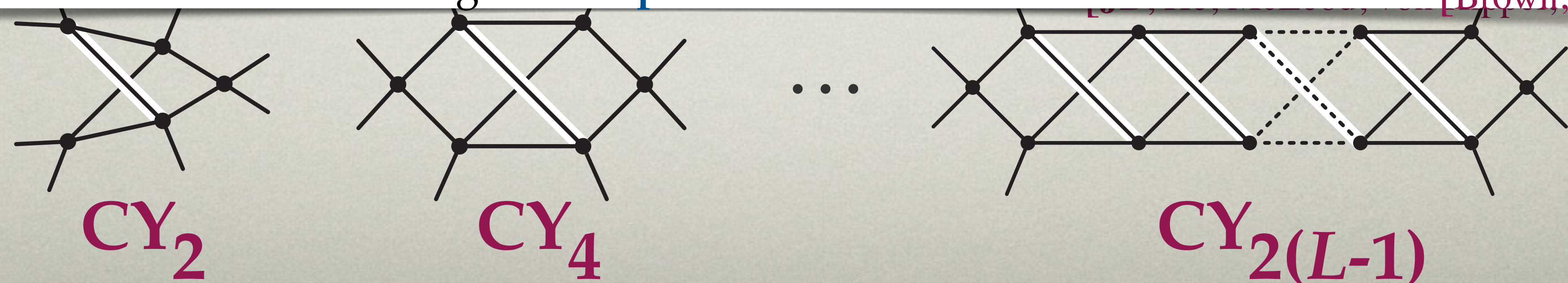
◆ In every instance known, the Calabi-Yau itself is **very simple**:

▶ degree- $(d+1)$ hypersurface in \mathbb{P}^d (or multiple-cover thereof)

◆ kinematic data (momenta/masses) control the moduli

▶ often **extremely singular**
an 8-loop vacuum graph evaluating to a **K3 period**

tardigrades:



SS

(2018)

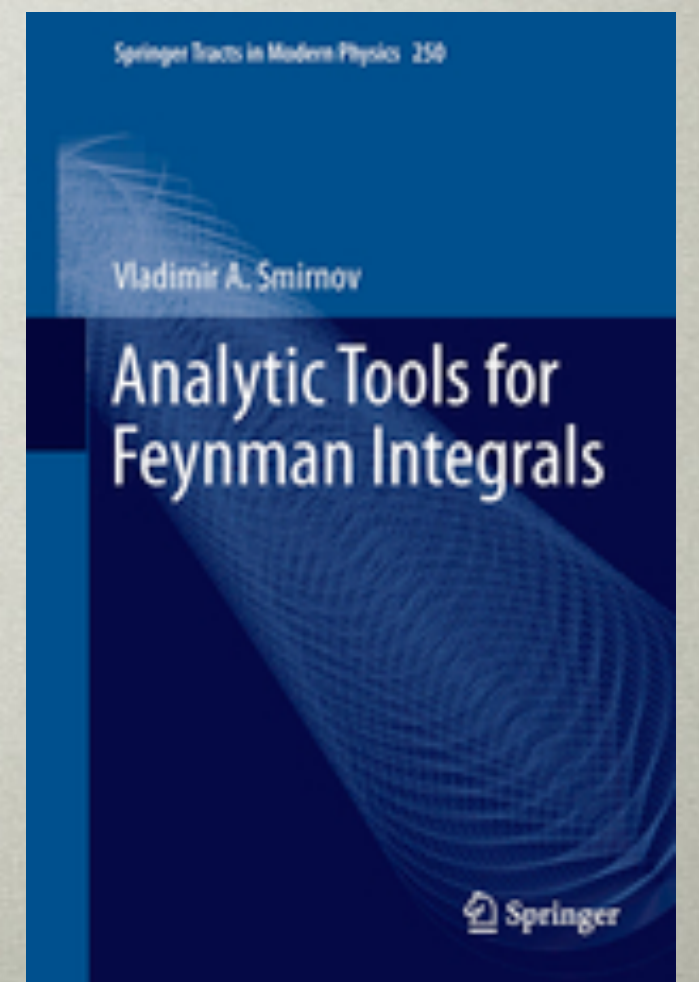
[Brown, Schnetz (2018)]

Why is Perturbation Theory so Hard?



- ◆ **Feynman integrals** (esp. with *scalar* numerators) are *horrible*
 - ▶ difficult to integrate, explosive in number, non-physical,...
- ◆ **Regularization** obscures symmetries (+is technically difficult)
- ◆ **Most** familiar **master integrand bases** are the *unnecessarily bad*:
 - ▶ don't satisfy nice / canonical differential equations
 - ▶ contain *multiple elliptic*(+worse(!)) geometries,
 - ▶ ...

$$I_{a_1 a_2 a_3 a_4 a_5 a_6 a_7 a_8 a_9} = e^{2\gamma_E \epsilon} \iint \frac{d^D k_1 d^D k_2}{(i\pi^{d/2})^2} \frac{P_8^{-a_8} P_9^{-a_9}}{P_1^{a_1} P_2^{a_2} P_3^{a_3} P_4^{a_4} P_5^{a_5} P_6^{a_6} P_7^{a_7}}$$





How can We Make it Easier?

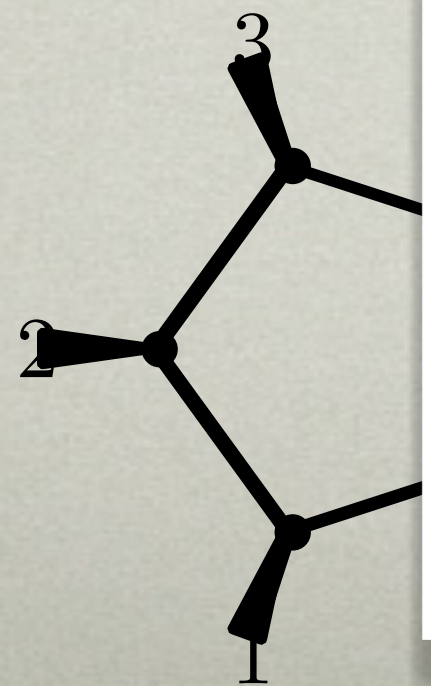
- ◆ **Use unitarity** to choose the *niciest/easiest* integrals to integrate (of course, integration “ease” changes with time and new methods)
 - ▶ search for as many *pure* integrals as you can —those which satisfy nice (canonical) differential equations

Definition: a function $f(s)$ is called *pure* if:

- ▶ there exists a grading of functions by “transcendental” *weight*
- ▶ any **derivative** of $f(s)$ is *strictly lower in weight*

e.g. $g(s) \log(f(s))$ would be *impure*

$$\frac{\partial}{\partial s} [g(s) \log(f(s))] = g'(s) \log(f(s)) + g(s) f'(s) / f(s)$$



(ANSWER. SOMETIMES)

er, Patatoukos (2021); ...]

a sum of
arithms &

elliptic curves

Unitarity-Based Strategies
& the stratification of loop integrands

Generalized Unitarity (in brief)



- ◆ The basic idea behind **unitarity**-based methods is that any *Feynman integrand* is a *rational differential form on loop momenta*
 - ▶ as such, it can be expanded into a **basis** \mathfrak{B} of such forms:

$$A = \sum_{\mathfrak{b}^i \in \mathfrak{B}} a_i \mathfrak{b}^i$$

- ◆ For any fixed QFT (spacetime dimension, particle content), the space of **all amplitude integrands** is **finite-dimensional**
 - ▶ *all-multiplicity amplitudes* can be expressed in a *finite basis*!
- ◆ **Key observation:** viewed as a potential element of *some* basis, *every* Feynman integrand can be interesting!
 - ▶ Why not try to find the *best/easiest* integrands—and use these?



Stratifying Integrand Bases

- ◆ Suppose that a basis could be carved up into **subspaces** (by any, arbitrary means):

$$A = \sum_{\mathfrak{b}^i \in \mathfrak{B}} a_i \mathfrak{b}^i$$
$$\mathfrak{B} =: \bigoplus_{p=0}^{\infty} \mathfrak{B}_p \quad A = \sum_p \mathcal{A}_p \quad \text{with} \quad \mathcal{A}_p := A \cap \mathfrak{B}_p := \sum_{\mathfrak{b}_p^i \in \mathfrak{B}_p} a_i \mathfrak{b}_p^i$$

- ◆ Such a stratification could be given by “*power-counting*” (some proxy for) ultraviolet behavior
 - ▶ recently, we gave an **intrinsically graph-theoretic** definition of power-counting for ***non-planar*** integrand bases

[**JB**, Herrmann, Langer, Trnka (2020)]



Stratifying Integrands Bases

- ◆ Suppose that a basis could be carved up into subspaces (by any arbitrary means):

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- ◆ ¿Is it possible to stratify integrand bases by *physical structure*?

(max-weight) \oplus (next-to-max-weight) $\oplus \dots \oplus$ (rational)

(polylogs) \oplus (elliptic-logs) \oplus (K3's) $\oplus \dots$

(finite) \oplus (\mathbb{H} -divergent) \oplus (\mathbb{R}^2 -divergent) $\oplus \dots$

\oplus ($\log(m)$ -divergent) \oplus ($\log(m)^2$ -divergent) $\oplus \dots$

[JB, Kalyanapuram]

[JB, Langer, Zhang]

[JB, Herrmann, Langer, Patatoukos, *et al*]



Prescriptive Integrand Bases

- How *generalized unitarity* has been used to match amplitudes:

$$A = \sum_i a_i \mathcal{I}_i^0$$

with coefficients a_i determined by **cuts**: a spanning set of cycles $\{\Omega_j\}$

$$\mathcal{I}_j := \sum_i \left\{ \mathcal{I}_i^0 (\mathbf{M}^{-1})_{i,j} \right\} \int_{\Omega_j} \left\{ \mathcal{I}_i^0 (\mathbf{M}^{-1})_{i,j} \right\} \delta_{i,j}$$

$$a_j := \int_{\Omega_j} A \left\{ \sum_i C_i \mathcal{I}_i^0 \right\} \int_{\Omega_j} \left\{ \mathcal{I}_i^0 (\mathbf{M}^{-1})_{i,j} \right\} C_j = \sum_i \left\{ \mathcal{I}_i^0 (\mathbf{M}^{-1})_{i,j} \right\} a_i$$

- A basis is called *prescriptive* if it is the **cohomological dual** of a spanning set of **cycles** $\{\Omega_j\}$

Strategies for Building Bases



- ◆ Given *some* integrand basis (or strata thereof), one should *diagonalize* the space of integrands according to a

homological/cohomological pairing:

- ▶ choose a *spanning-set* of **compact, max-dimensional** contours Ω_j
- ▶ normalize and diagonalize the basis by the requirement

$$\int_{\Omega_j} \mathbf{b}_i = \delta_{ij}$$

Ω_j : $4L$ -dimensional compact contours

{	“residues”
	elliptic periods
	K3 periods, etc.

- ◆ This *trivializes* the representation of amplitudes:
 - ▶ the **coefficient** of any amplitude in this basis will simply be the *on-shell function* evaluated on the contour (a **leading singularity**)
- ◆ Choosing a **maximal** set of IR/UV-**divergence-probing contours** ensures(?) that the basis is split into finite / divergent subspaces



Prescriptivity and Purity

◆ Prescriptive integrand bases are naturally *pure*

$$\mathcal{I}(N) \Leftrightarrow \begin{array}{c} \text{Diagram with vertices } a_1, a_2, a_3, a_4 \text{ and } b_1, b_2 \\ \text{Regions } l_1, l_2 \end{array} = \left(\text{pure polylogarithms} \right) + f(\vec{p}) \times \left(\text{other pure polylogarithms} \right)$$

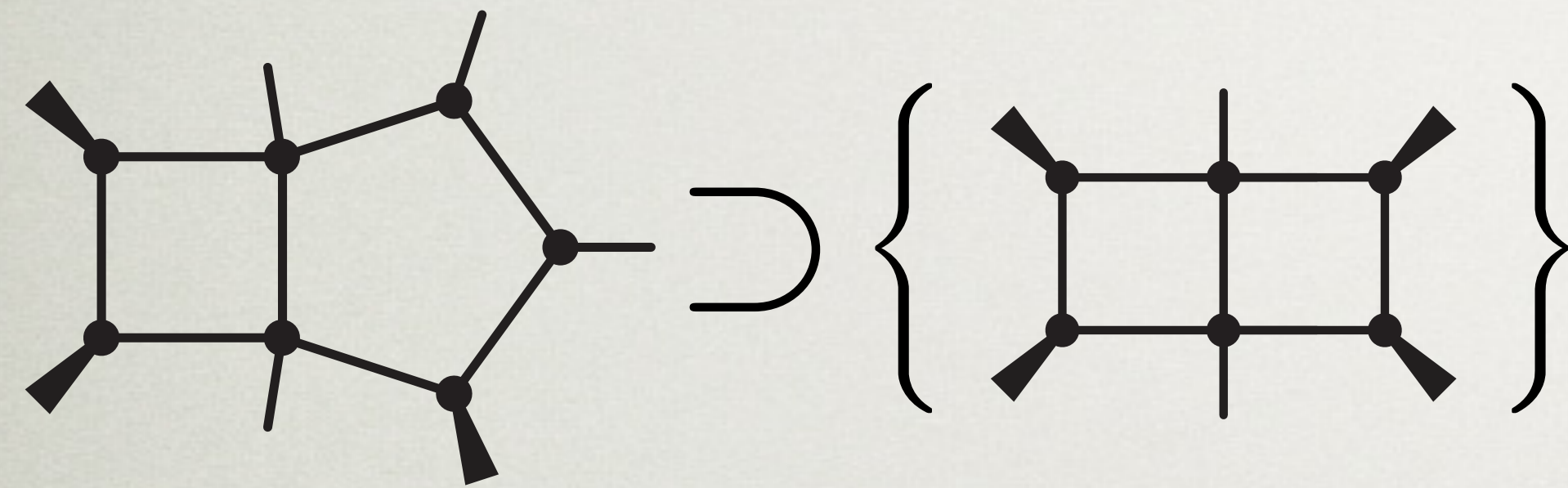
$$\int_{l_i \in \mathbb{R}^4} \begin{array}{c} \text{Diagram with vertices } a_1, a_2, a_3, a_4 \text{ and } b_1, b_2 \\ \text{Regions } l_1, l_2 \end{array} \propto \left(\text{other pure polylogarithms} \right)$$

$$\{\Omega_1, \dots, \Omega_6\} := \left\{ \begin{array}{l} n(a_3) \Omega_1, \Omega_2, \Omega_3, \Omega_4, \Omega_5, \Omega_6 \\ \left(\begin{array}{c|c} l_1 & N_1^0 \\ \hline l_1 & N_2^0 \end{array} \right) \begin{array}{c} a_4 \\ b_1 \\ b_2 \\ a_1 \end{array}, \left(\begin{array}{c|c} l_1 & N_1^0 \\ \hline l_1 & N_2^0 \end{array} \right) \begin{array}{c} a_4 \\ b_1 \\ b_2 \\ a_1 \end{array}, \left(\begin{array}{c|c} l_1 & N_1^0 \\ \hline l_1 & N_2^0 \end{array} \right) \begin{array}{c} a_4 \\ b_1 \\ b_2 \\ a_1 \end{array}, \left(\begin{array}{c|c} l_1 & N_1^0 \\ \hline l_1 & N_2^0 \end{array} \right) \begin{array}{c} a_4 \\ b_1 \\ b_2 \\ a_1 \end{array}, \left(\begin{array}{c|c} l_1 & N_1^0 \\ \hline l_1 & N_2^0 \end{array} \right) \begin{array}{c} a_4 \\ b_1 \\ b_2 \\ a_1 \end{array}, \left(\begin{array}{c|c} l_1 & N_1^0 \\ \hline l_1 & N_2^0 \end{array} \right) \begin{array}{c} a_4 \\ b_1 \\ b_2 \\ a_1 \end{array} \end{array} \right\} =: M^0$$

Stratifying Rigidity



♦ Is it possible to *stratify* integrands according to *rigidity*?



$n(\ell_1)$	Ω_a	Ω_b	Ω_{poles}^e
$(\ell_1 N_1)$	J_1^a	J_1^b	0
$(\ell_1 N_2)$	J_2^a	J_2^b	0
$(\ell_1 N_3)$	J_3^a	J_3^b	1

[JB, Kalyanapuram]

$=: \mathbf{M}$

$$\overline{|N_1)} := \frac{1}{J_1^a \det(\mathbf{M})} \left[\frac{J_2^b |N_1) - J_1^b |N_2)}{\pi y_4 \sqrt{r_{14} r_{23}}} \right] K[\varphi]$$

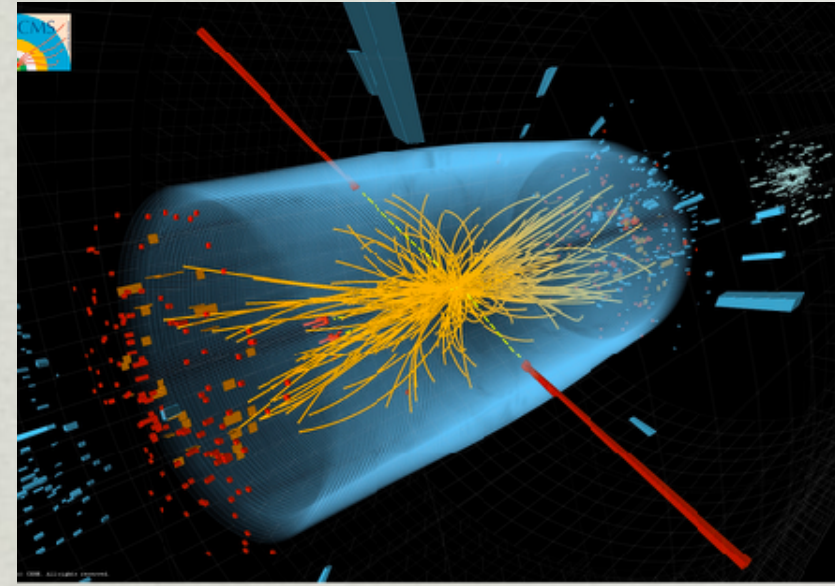
$$\overline{|N_2)} := \frac{\text{a-cycle}}{J_2^a \det(\mathbf{M})} \left[-J_2^a |N_1) + J_1^a |N_2) \right] \frac{y(q_+)}{\pi y_4 \sqrt{r_{14} r_{23}} q_4^+} \left(K[\varphi] + \frac{r_{24}}{q_2^+} \Pi \left[\frac{q_4^+ r_{12}}{q_2^+ r_{14}}, \varphi \right] \right) - (q_+ \leftrightarrow q_-)$$

$$\overline{|N_3)} := \frac{1}{\det(\mathbf{M})} \left[(J_2^a J_3^b - J_2^b J_3^a) |N_1) + (J_1^b J_3^a - J_1^a J_3^b) |N_2) \right] + |N_3)$$

$$J_3^a := \frac{\text{a-cycle}}{\det(\mathbf{M})} \frac{y(q_+)}{\pi y_4 \sqrt{r_{14} r_{23}} q_4^+} \left(K[\varphi] + \frac{r_{24}}{q_2^+} \Pi \left[\frac{q_4^+ r_{12}}{q_2^+ r_{14}}, \varphi \right] \right) + (q_+ \leftrightarrow q_-)$$

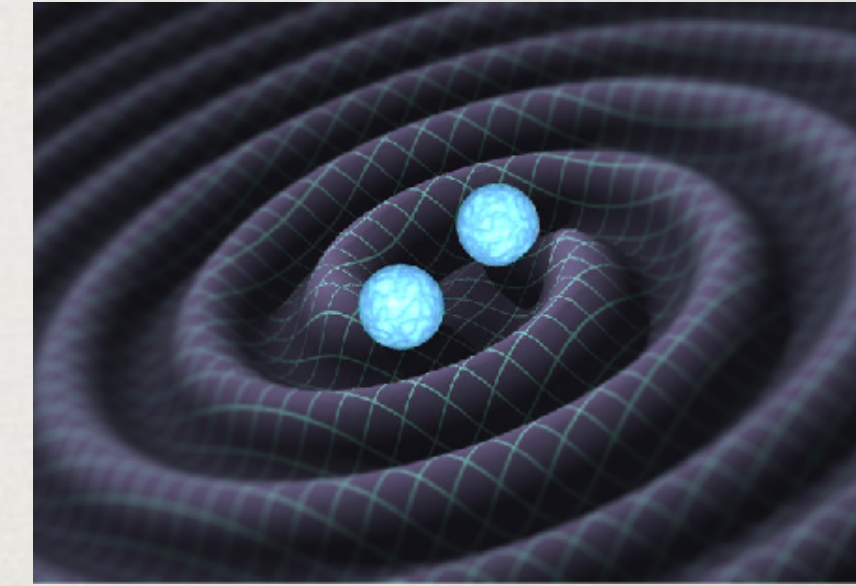
$$\det(\mathbf{M}) = \frac{y(q_+)}{\pi^2 y_4^2 q_3^+ q_4^+ r_{14} r_{23}} K[\varphi] \Pi \left[\frac{q_4^+ r_{13}}{q_3^+ r_{14}}, 1 - \varphi \right] - (q_+ \leftrightarrow q_-) - (r_2 \leftrightarrow r_3)$$

Amplitudes: a Virtuous Cycle



Compute Something

beyond the reach of
recent imagination



Exploit Simplicity

to build more powerful
computational technology

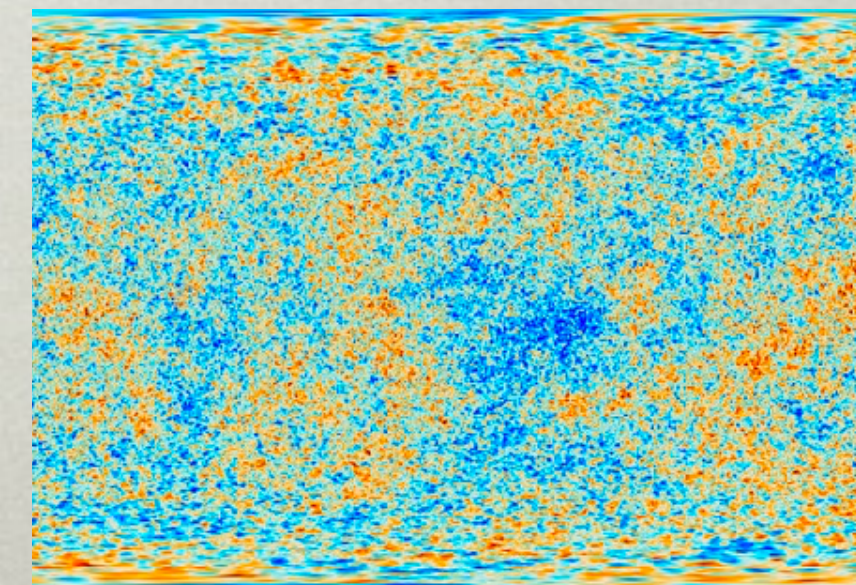
Discover Simplicity

beyond expectations



Understand Why

study it, understand it,
& explore consequences



Thank you!