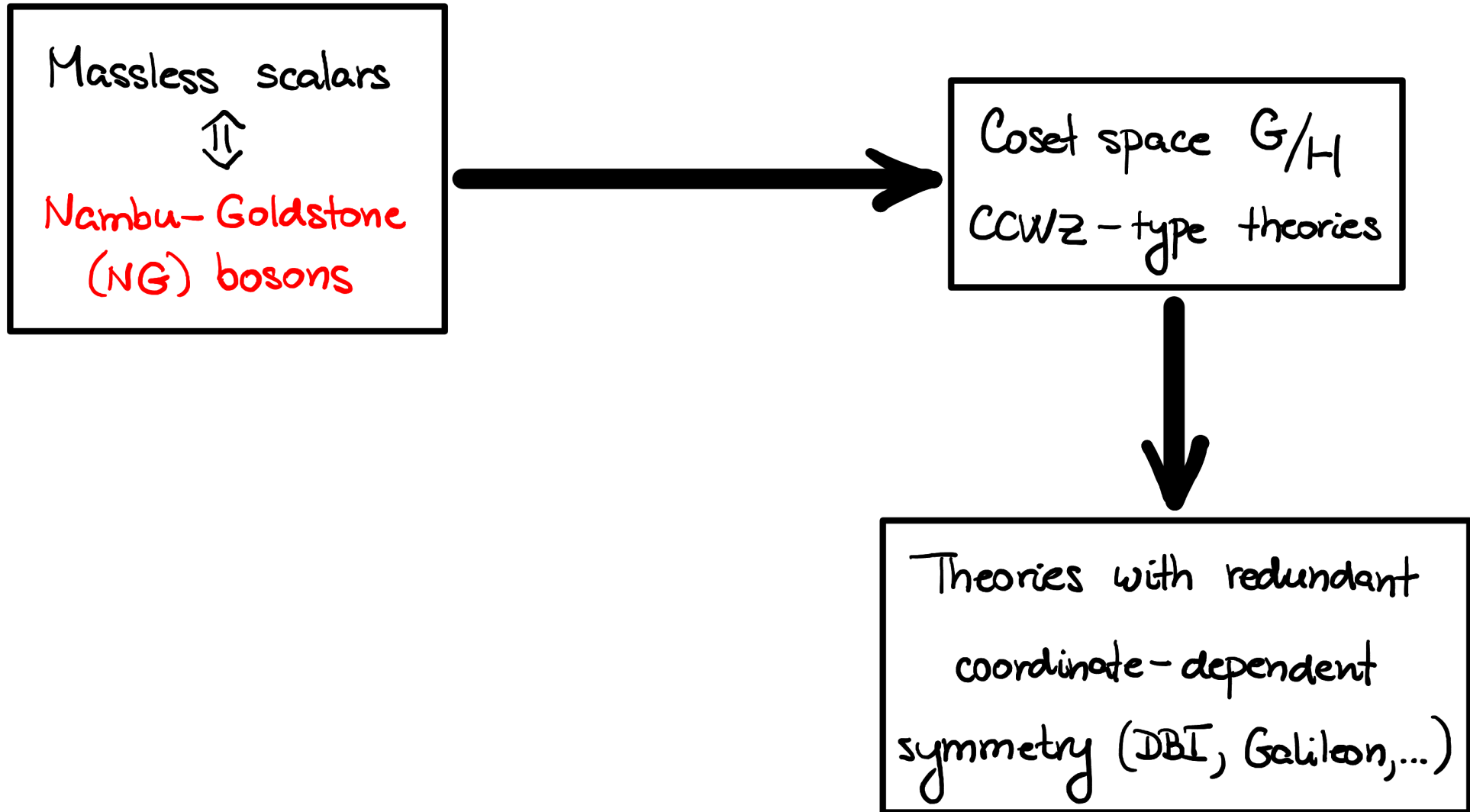


SCATTERING AMPLITUDES IN NONRELATIVISTIC EFT

PROLOGUE

RELATIVISTIC EFTS



SOFT LIMIT OF SCATTERING AMPLITUDES

$$\mathcal{A}_{\alpha \rightarrow \beta + \pi(p)} = \text{Diagram}$$

Asymptotic behavior of
the amplitude :

$$\mathcal{A}_{\alpha \rightarrow \beta + \pi(p)} = \# p^\sigma + \dots$$

soft scaling parameter

- $\sigma = 0$: unexpected exception?
- $\sigma = 1$: Adler zero (default)
- $\sigma \geq 2$: enhanced soft limit

UNEXPECTED EXCEPTIONS?

CCWZ-type EFTs : $\mathcal{L}_{\text{eff}} = \frac{1}{2} g_{ab}(\pi) \partial_\mu \pi^a \partial^\mu \pi^b$

NG fields = coordinates on G/H

G -invariant metric on G/H

$$\lim_{p \rightarrow 0} \sum_b F_a^b \mathcal{R}_{\alpha \rightarrow \beta + \pi^b(p)} = \frac{1}{2} \sum_{b \in \alpha \cup \beta} \sum_c (f_{cab} - f_{bac} - f_{abc}) \mathcal{R}_{\alpha \rightarrow \beta}^{b \rightarrow c}$$

$\langle 0 | J_a^\mu(0) | \pi^b(\vec{p}) \rangle = -i p^\mu F_a^b$
structure constants of G

Main conclusion :

Adler zero $\Leftrightarrow G/H$ is symmetric

Kampf, Novotný, Shifman, Trnka (2020)
 Cheung, Helset, Parra-Martinez (2022)

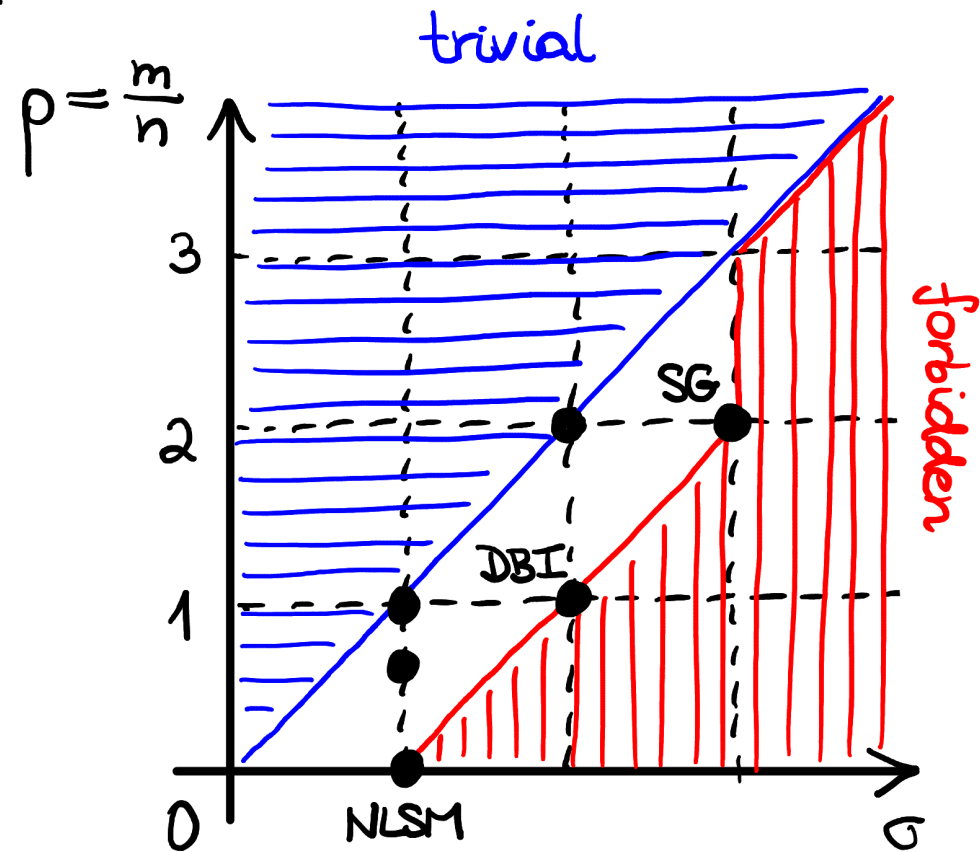
BEYOND ADLER ZERO

EFTs of a single NG boson: $\mathcal{L}_{\text{eff}} \sim (\partial\pi)^2 \sum_{m,n=0}^{\infty} c_{m,n} \partial^m \pi^n$

Not all values of $\underline{\sigma}$ are allowed.

Exceptional theories: maximize $\underline{\sigma}$ for given \underline{p} .

- NLSM : $(\underline{p}, \underline{\sigma}) = (0, 1)$
- DBI : $(\underline{p}, \underline{\sigma}) = (1, 2)$
- SG : $(\underline{p}, \underline{\sigma}) = (2, 3)$



Cheung et al. (2015...)

FINDING NEW THEORIES

Bottom-up
soft bootstrap

Cheung et al. (2015...)
Elvang et al. (2018...)

Top-down
symmetry classification

Bogers, TB (2018)
Roest, Stefanyshyn, Werkman (2019)

Importance of coordinate-dependent symmetries : invariance under

$$\delta_\epsilon \pi(x) = \epsilon_{\mu_1 \dots \mu_n} \left[x^{\mu_1} \dots x^{\mu_n} + \int_A^{\mu_1 \dots \mu_n}(x) \partial^A [\pi](x) \right] \text{ implies } \sigma = n+1$$

Cheung et al. (2017)

NONRELATIVISTIC NG BOSONS

MOTIVATION

- Some kind of NG boson exists in all phases of matter.
- Scattering (amplitudes) also important in a medium.
- Nonrelativistic gravity / string theory?

CLASSIFICATION OF NG BOSONS

Type	Dispersion relation	Propagator
A_m	$\omega \sim \vec{p} ^m$	$\frac{1}{\omega^2 - \vec{p}^{2m}}$
B_{2m}	$\omega \sim \vec{p} ^{2m}$	$\frac{1}{\omega - \vec{p}^{2m}}$

TB, Watanabe, Murayama, Hidaka (2011-2012)

Griffin, Grosvenor, Horava, Yan (2015)

CONSTRAINTS ON THE SPECTRUM

Type	Dispersion relation	Allowed dimension ($T=0$)	Allowed dimension ($T>0$)
A_m	$\omega \sim \vec{p} ^m$	$D \geq m+2$	$D \geq 2m+2$
B_{2m}	$\omega \sim \vec{p} ^{2m}$	any $D \geq 2$	$D \geq 2m+2$

Mermin, Wagner (1966)
 Hohenberg (1967)
 Coleman (1973)

Griffin, Grosvenor, Horava, Yan (2013-2015)
 Watanabe, Murayama (2014)
 Argurio, Naegels, Pasternak (2019)

MASSIVE NG BOSONS

System with chemical potential μ : excitation energy = eigenvalue of $\tilde{H} = H - \mu Q$.

$$[Q, Q_i^\pm] = \pm q_i Q_i^\pm$$

& Q_i^\pm spontaneously broken



state in the spectrum
with $\lim_{\vec{p} \rightarrow \vec{0}} \omega(\vec{p}) = \mu q_i$

- Exact gap protected by symmetry.
- Interactions dictated by symmetry.
- Scattering amplitudes with Adler zero.

Nicolis, Piazza (2012)

Watanabe, TB, Murayama (2013)

TB, Jakobsen (2018)

ADLER ZERO FOR NG BOSONS?

Generic shift-invariant EFT for a single type- A_1 NG boson :

$$\mathcal{L}_{\text{eff}} = \frac{1}{2} (\partial_\mu \pi)^2 + c_1 (\partial_0 \pi)^3 + c_2 \partial_0 \pi (\vec{\nabla} \pi)^2 + \dots$$

- Such cubic vertices cannot be removed by field redefinition!
- Violates Adler zero unless $c_2 = -c_1$:

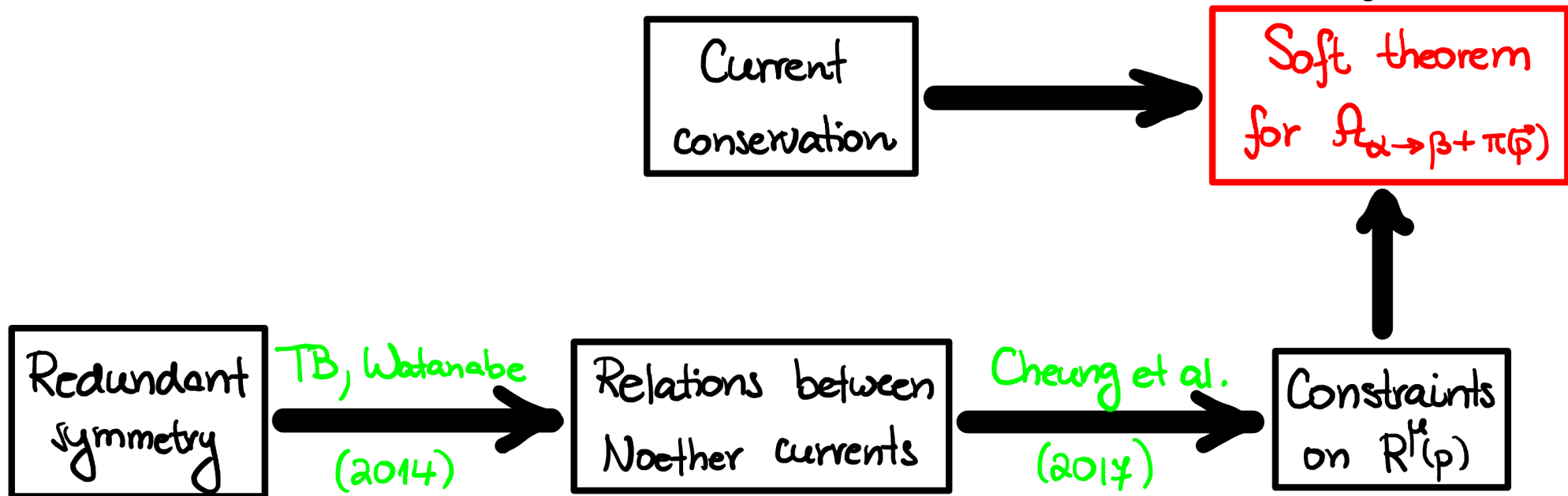
$$\mathcal{L}_{\text{eff}} \longrightarrow \frac{1}{2} (\partial_\mu \pi)^2 + c_1 \partial_0 \pi (\partial_\mu \pi)^2 + \dots$$

General criterion for Adler zero in terms of symmetry not available yet.

ENHANCED SOFT LIMIT

General approach to Adler zero and its enhancement :

$$\langle \beta | J^\mu(0) | \alpha \rangle = \langle 0 | J^\mu(0) | \pi(\vec{p}) \rangle \frac{i}{p^0 - \omega(|\vec{p}|)} \mathcal{A}_{\alpha \rightarrow \beta + \pi(\vec{p})} + R^\mu(p)$$



ENHANCED SOFT LIMIT

Invariance under spatial polynomial shift symmetry:

$$\delta_{\epsilon} \pi(x) = \epsilon_{i_1 \dots i_n} \left[x^{i_1} \dots x^{i_n} + \int_A^{i_1 \dots i_n}(x) \mathcal{O}^A[\pi](x) \right]$$

type- A_m NG bosons

$$\sigma = \min(m, n+1)$$

type- B_{2m} NG bosons

$$\sigma = \min(2m, n+1)$$

Mojahed, TB (2022)

SOFT BOOTSTRAP

Mojahed, TB (2021-2022)

NAIVE LAGRANGIAN SCAN - TYPE A₁

Theories with $p=1$ (one derivative per field):

two-parametric family of theories with $\sigma=2$

$$\mathcal{L}_{\text{eff}} \sim \sqrt{(1 - c_1 \partial_0 \pi)^2 - c_2 [(\partial_0 \pi)^2 - (\vec{\nabla} \pi)^2]}$$

Special cases

● $c_1 = 0$: relativistic DBI

● $c_1 = c_2 = 1$: $\mathcal{L}_{\text{eff}} \sim \sqrt{1 - 2\partial_0 \pi + (\vec{\nabla} \pi)^2}$

Galilei-invariant superfluid

Tree-level amplitudes independent of c_1
and equal to those of the
relativistic DBI theory.

Emergent Lorentz invariance!

SOFT RECURSION - TYPE A₁

Seed 4-point amplitudes: polynomials in (s, t, u) (ω_a) ($a=1, \dots, 4$)
relativistic Mandelstams energies

Only three amplitudes consistent

with recursion:

- $s^2 + t^2 + u^2$... relativistic DBI ($\sigma=2$)
- $s^3 + t^3 + u^3$... relativistic special Galileon ($\sigma=3$)
- $s^2(\omega_1\omega_2 + \omega_3\omega_4) + t^2(\omega_1\omega_3 + \omega_2\omega_4) + u^2(\omega_1\omega_4 + \omega_2\omega_3)$
... "spatial Galileon" ($\sigma=2$)

Exceptional theories are automatically Lorentz-invariant: DBI & special Galileon.

There are no exceptional type-A₁ theories with $\sigma > 3$.

NAIVE LAGRANGIAN SCAN - TYPE B₂

Included $p=1$ theories of the type $\mathcal{L}_{\text{eff}} = \psi^\dagger (i\partial_0 + \vec{\nabla}^2) \psi + \mathcal{L}_{\text{int}}(\psi, \psi^\dagger)$.

The only theory in this class with $\sigma=2$ is the Schrödinger-DBI theory:

$$\mathcal{L}_{\text{eff}} = \psi^\dagger i\partial_0 \psi + s \sqrt{1 - 2s \vec{\nabla} \psi^\dagger \cdot \vec{\nabla} \psi + (\vec{\nabla} \psi^\dagger \cdot \vec{\nabla} \psi)^2 - |\vec{\nabla} \psi \cdot \vec{\nabla} \psi|^2}$$

- $s = +1$
ISO($d+2$) spatial symmetry

- $s = -1$
ISO($d,2$) spatial symmetry

} ψ, ψ^\dagger are fluctuations of a d -dimensional brane embedded in a $(d+2)$ -dimensional pseudo-Euclidean space

SOFT RECURSION - TYPE B_2

Seed 4-point amplitudes : polynomials in nonrelativistic Mandelstams $S_{ab} = \vec{p}_a \cdot \vec{p}_b$

Only three amplitudes consistent

- with recursion :
- $S_{13} + S_{24}$... CP^1 NLSM ($\sigma=1$), ferromagnetic spin waves
 - $(S_{13} + S_{24})^2$... Schrödinger-DBI ($\sigma=2$)
 - lengthy $O(S_{ab}^3)$ expression ... "Schrödinger-Galileon" ($\sigma=2$)
- } exceptional

There are no exceptional type- B_2 theories with $\sigma > 2$.

EPILOGUE

SOME MORALS

- Adler zero and its enhancements, soft recursion & bootstrap are straightforward to generalize to nonrelativistic EFTs.

- Top-down symmetry approach & bottom-up bootstrap complement each other.

easy generation of concrete theories,
see Mojahed, TB (2022) for a technically
natural type- B_4 theory

easy elimination of large parts
of EFT parameter space

- Some surprises in nonrelativistic EFTs, and even more in theories of nonscalar particles: fractional scaling of amplitudes

TB, Esposito, Penco (2022)