

# Introduction to Particle Physics

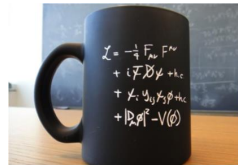
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SWEDISH PHYSICS TEACHERS 2023

# Handout 2

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- Symmetry, invariance and conservation laws
- Toward the SM Lagrangian
- Higgs field
- the SM Lagrangian



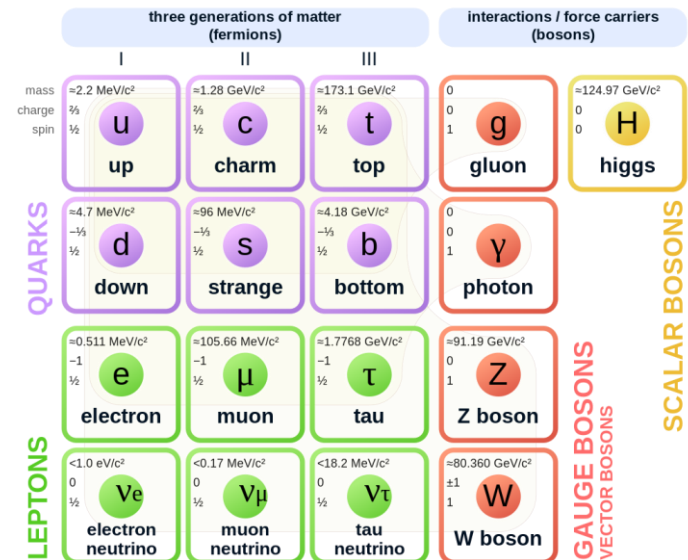
Credits to:

C. Grosjean and D. Tong : <https://indico.cern.ch/event/1254879/timetable/>

# Summary of Fundamental Interactions

PROPERTIES OF THE INTERACTIONS					
Property \ Interaction	Gravitational	Weak	Electromagnetic	Strong	
		Unification		Fundamental	Residual
Acts on:	Mass - Energy	Flavor	Electric Charge	Color Charge	See Residual Strong Interaction Note
Particles experiencing:	All	Quarks, Leptons	Electrically charged	Quarks, Gluons	Hadrons
Particles mediating:	Graviton (not yet observed)	$W^+ W^- Z^0$	$\gamma$	Gluons	Mesons
Strength relative to electromag for two u quarks at:	$10^{-41}$	0.8	1	25	Not applicable to quarks
for two protons in nucleus	$10^{-41}$	$10^{-4}$	1	60	
	$10^{-36}$	$10^{-7}$	1	Not applicable to hadrons	20

## Standard Model of Fundamental Particles



# Invariance of the physics laws

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The laws of nature should be **invariant** under certain transformations since “nature does not know” how we observe it!

For example, the physics laws should be invariant w.r.t. the changing of reference systems 😊 (Special Relativity)

**Abandon the anthropocentric view and move towards a more abstract mathematical view, universally valid!**

# Symmetry, invariance and conservation laws

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1. System is **invariant** under a transformation  $\leftrightarrow$  it **posses a symmetry**;

2. System is **invariant** under a transformation  $\leftrightarrow$  **conserved quantities**  
(A.E. Nother Theorem)



# Lagrangian in the classical Mechanics

The Newton law of classical mechanics

$$\vec{F} = m\vec{a} \quad \text{or} \quad V'(x) = -m\ddot{x}$$

can be obtained by requiring the least action principle

$$\delta S = 0$$

where

the action:  $S = \int_{t_1}^{t_2} dt \mathcal{L}(x, \dot{x})$  with the (classical) Lagrangian:  $\mathcal{L}(x, \dot{x}) = \frac{1}{2}m\dot{x}^2 - V(x)$

(Hamiltonian/energy:  $\mathcal{H} = \dot{x} \frac{\delta \mathcal{L}}{\delta \dot{x}} - \mathcal{L} = \frac{1}{2}m\dot{x}^2 + V(x)$ )

Kinetic term

Potential term

Euler-Lagrange equations

$$\delta S = \int_{t_1}^{t_2} dt \left( \frac{\delta \mathcal{L}}{\delta x} - \frac{d}{dt} \frac{\delta \mathcal{L}}{\delta \dot{x}} \right) \delta x + \text{boundary terms} = 0 \quad \Rightarrow \quad \frac{\delta \mathcal{L}}{\delta x} = \frac{d}{dt} \frac{\delta \mathcal{L}}{\delta \dot{x}}$$

For the classical Lagrangian:  $-V'(x) = m\ddot{x}$

# Classical invariance and conservation laws

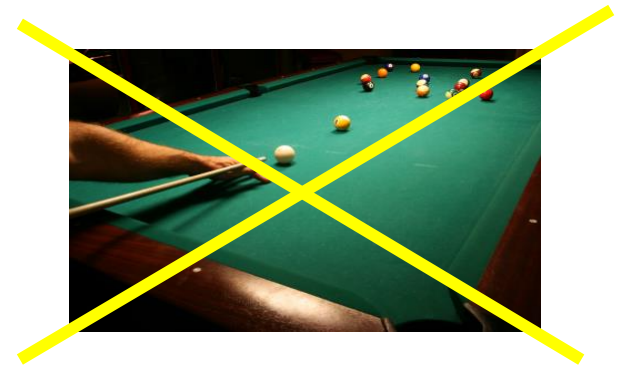
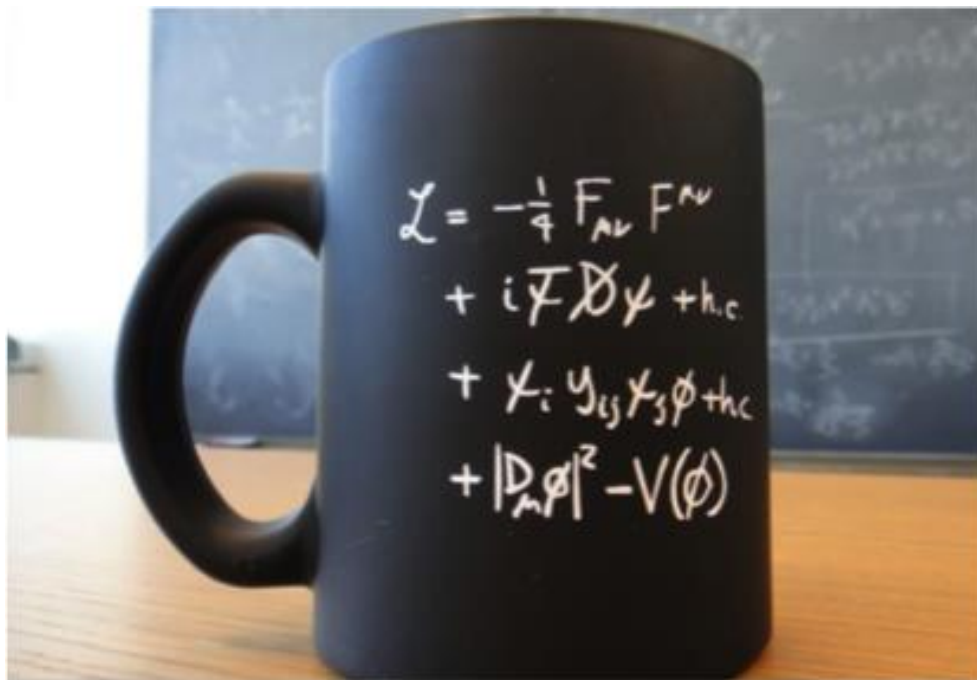
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The evolution of a physical system of **interacting particles** can be fully described by **equations** in the **Lagrangian (or Hamiltonian) formulation**; (Eulero-Lagrange eq.  $\rightarrow x=x(t)$ )

1. If **L** is **invariant** under time translation ( $t \rightarrow t+t_0$ )  $\rightarrow$   **$E_{\text{tot}}$**  is conserved;
2. If **L** is **invariant** under space translation ( $x \rightarrow x+x_0$ )  $\rightarrow$  the **momentum  $\mathbf{p}$**  is conserved;
3. If **L** is **invariant** under axes rotation ( $\theta, \varphi$ )  $\rightarrow$  the **angular momentum  $\mathbf{J}$**  is conserved.

# Toward the “Lagrangian of the SM

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REM:

- particle  $\leftrightarrow$  quantum field
- Complementary principle;
- Heisenberg indetermination principle.



# Lagrangians for Particle Physics

Equations of motion, like  $\vec{F} = m\vec{a}$ , are **covariant** under the action of a symmetry.

(Under the Transformation of Galileo between two reference systems) gdc

Lagrangians are **invariant**.

That makes identifying the symmetries of Nature much easier.

## —Particle Physics—

particles  $\leftrightarrow$  fields with specific transformation properties under some fundamental symmetries

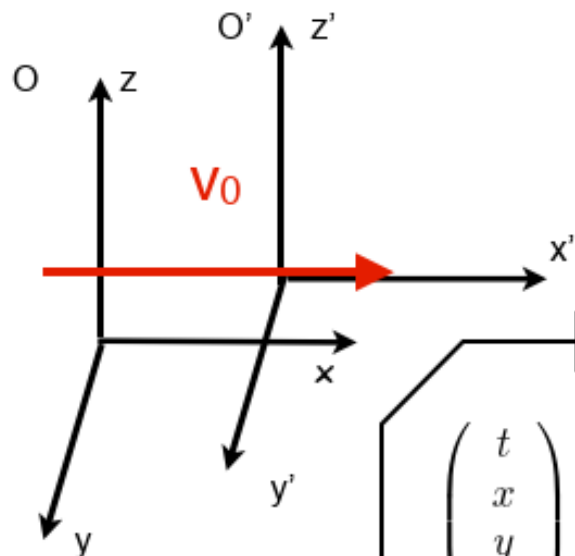
build a Lagrangian (i.e. a function of the these fields and their space-time derivatives) that remains invariant under the action of the symmetry transformations.

- 1) L is invariant under the Special Relativity space-time transform (Lorentz invariant)
- 2) L is invariant under (local)transformations of the particle fields and of the potentials (gauge invariant)

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# Galilean and Lorentz Transformations



Consider two observers

in relative motion with a constant speed  $v_0$  along the x-axis  
they use their own systems of coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$

## Galilean transformations

$$\begin{pmatrix} t \\ x \\ y \\ z \end{pmatrix} \rightarrow \begin{pmatrix} t' = t \\ x' = -\beta_0 ct + x \\ y' = y \\ z' = z \end{pmatrix} \quad \text{with} \quad \beta_0 = \frac{v_0}{c}$$

in particular

$$v' = v - v_0$$

the speed can be arbitrarily large.

## Lorentz transformations

$$\begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix} \rightarrow \begin{pmatrix} ct' = \gamma_0 (ct - \beta_0 x) \\ x' = \gamma_0 (-\beta_0 ct + x) \\ y' = y \\ z' = z \end{pmatrix} \quad \text{with} \quad \beta_0 = \frac{v_0}{c}$$

$$\gamma_0 = \frac{1}{\sqrt{1 - \beta_0^2}}$$

in particular

$$v' = \frac{v - v_0}{1 - v \cdot v_0 / c^2}$$

The speed of light is the same for all observers:

if  $v=c$  then  $v'=c$  too

Note:  $\Delta^2 \equiv (ct)^2 - x^2 - y^2 - z^2 = (ct')^2 - x'^2 - y'^2 - z'^2 \equiv \Delta'^2$

# Summary on Simmetries

SYMMETRY name	INVARIANCE	CONSERVED QUANTITY
Classical Space and time translation symmetries; axes rotation symmetry.  Space and time symmetry (Galilean transformation)	System invariance under: <ul style="list-style-type: none"> <li>• space and time <b>translations</b>;</li> <li>• Axes rotation</li> <li>• Change of reference sys.</li> </ul>	Momentum $\mathbf{p}$ , Energy $E$ and angular momentum $\mathbf{L}$ , respectively
SPACE-TIME symmetries	Lorentz invariance	Square modules of four-vectors: $\mathbf{x}$ , $E-\mathbf{p}$
Internal local symmetry of Lagrangian interaction: <b>GAUGE symmetry</b> (classic ED, QED, QCD, EW)	System Invariant under GAUGE transform	Strength of the interaction

# The SM ingredients

Feynman diagram of the muon beta decay

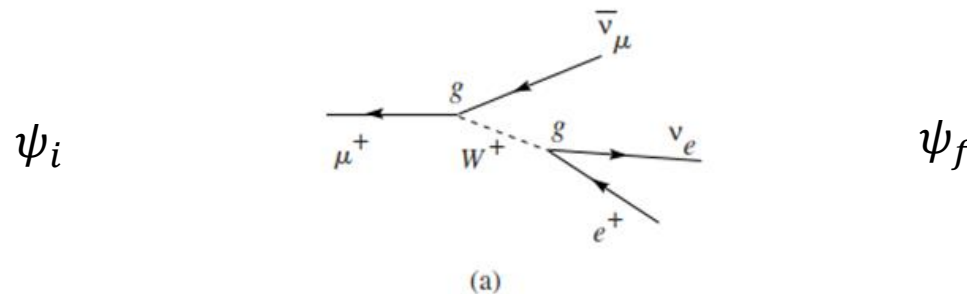


Fig. 7.4. Muon beta decay.

The lifetime  $\tau$  of the  $\mu$  depends on the transition amplitude from  $\psi_i$  to  $\psi_f$ .

It is given by the “matrix element” of the interaction Lagrangian between the final  $|f\rangle$  and the initial  $|i\rangle$  states. With QM symbols:

$$M_{fi} = \langle \psi_f | L_{int} | \psi_i \rangle$$

$$|M_{fi}|^2 \rightarrow \tau_\mu = 10^{-6} s \text{ that is the experimental value!}$$

# L properties

## Fermion Lagrangian

$$\mathcal{L} = \psi^\dagger \gamma^0 (i\gamma^\mu \partial_\mu - m) \psi$$

$\psi$  4-component Dirac spinor  
describes a spin-1/2 particle  
when quantised

$\gamma^\mu$  ( $\mu = 0, 1, 2, 3$ ) are four 4x4 matrices

Kinetic term

Potential term

- Equation of motion:

$$0 = \delta\mathcal{L} = \psi^\dagger \gamma^0 (i\gamma^\mu \partial_\mu - m) \delta\psi \quad \longrightarrow \quad \text{Dirac equation} \quad (i\gamma^\mu \partial_\mu - m) \psi = 0$$

# Dirac Equation

Equation of motion of a fermion :

<b>Dirac Equation (1928):</b>	$\left( i\gamma^\mu \partial_\mu - \frac{mc}{\hbar} \right) \Psi = 0$
$E = \begin{cases} +\sqrt{p^2c^2 + m^2c^4} & \text{matter} \\ -\sqrt{p^2c^2 + m^2c^4} & \text{antimatter} \end{cases}$	$E = \vec{\alpha}\vec{p}c + \beta mc^2$ $\gamma^0 = \beta, \quad \gamma^i = \beta\alpha^i, \quad \{\gamma^\mu, \gamma^\nu\} = 2\eta^{\mu\nu}$

positron (e<sup>+</sup>) discovered by C. Anderson in 1932

In this equation,  $\psi$  is the Dirac bi-spinor

$$\psi(x) = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix} = \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad \varphi = \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix} \quad \chi = \begin{pmatrix} \chi_1 \\ \chi_2 \end{pmatrix}. \quad (2.28)$$

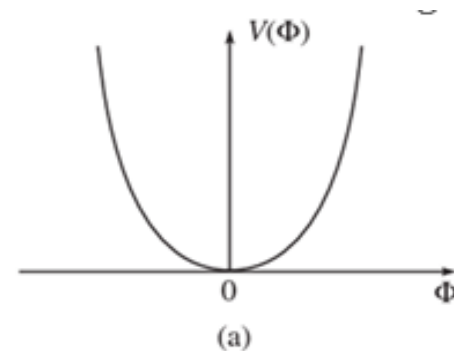
# The Higgs mechanism

The **Higgs mechanism** was proposed (1960) to provide mass to the weak gauge bosons  $W^\pm$  and  $Z^0$ , the interaction carriers, **without destroying the gauge invariance of the EW Lagrangian** !

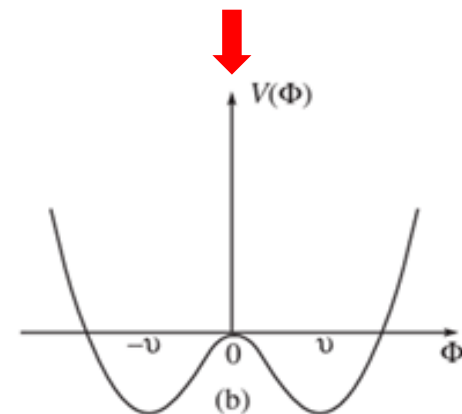
this was to unify the Weak and EM interactions in a gauge QFT (Electroweak interaction);

After about 50 years, The Higgs boson was observed at the LHC and presented on the 4<sup>th</sup> July 2012 at CERN;

The Higgs field fills the Universe with a non zero value and provides mass to the fundamental particles;



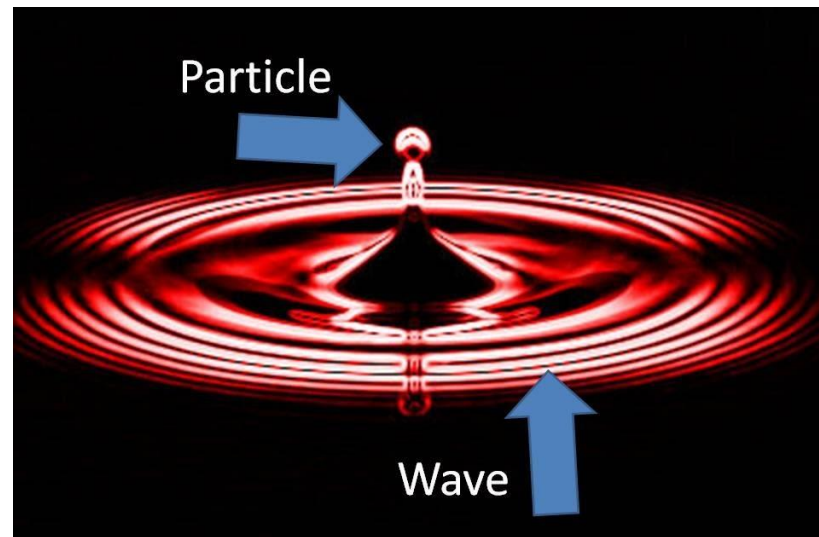
$\sim 10^{-12}$  s after the Big-Bang



# Observation of the Higgs boson

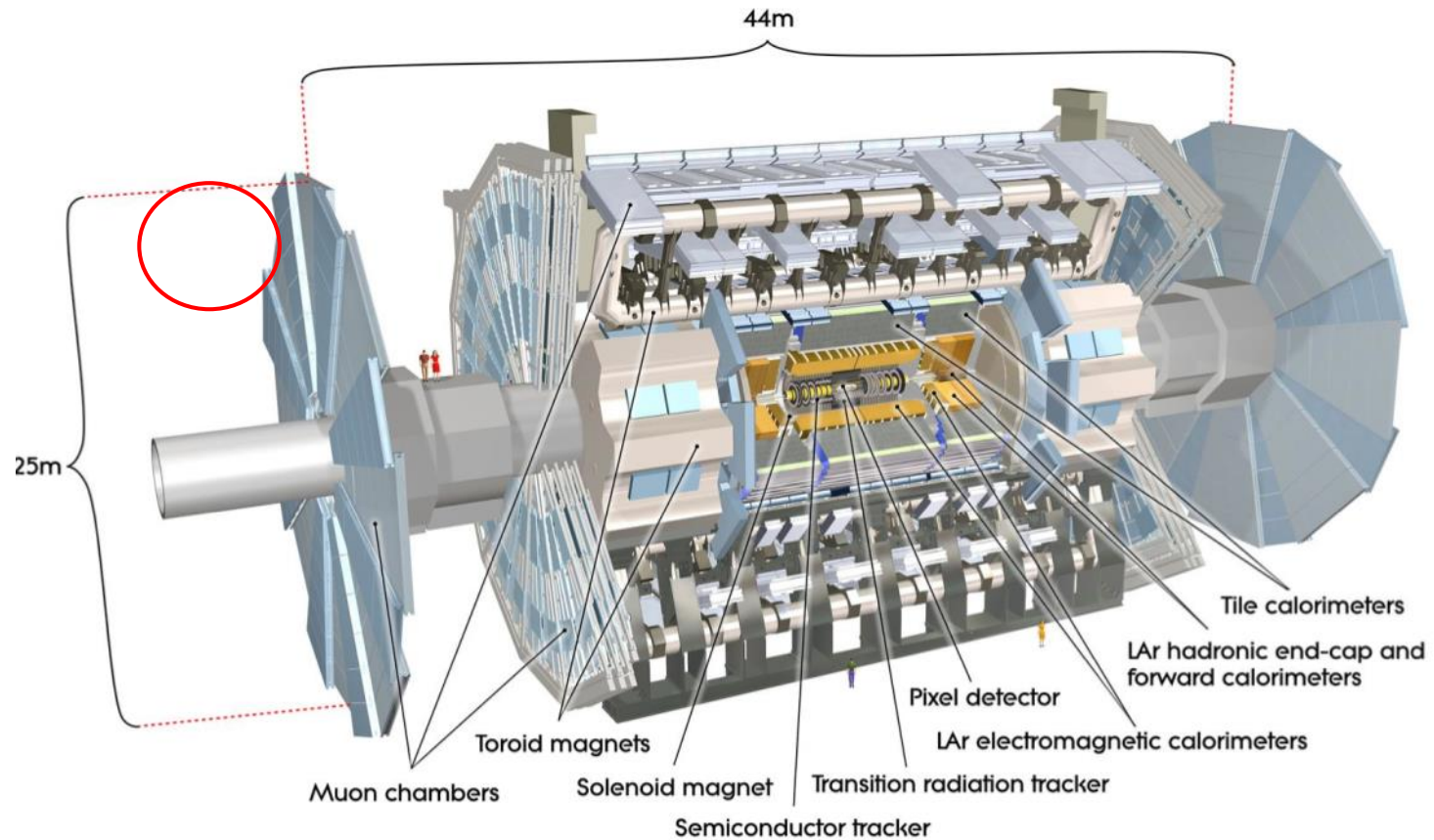
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Two protons interacting in the LHC excite the Higgs field that manifests itself with the emission of the boson with a mass of 126 GeV.

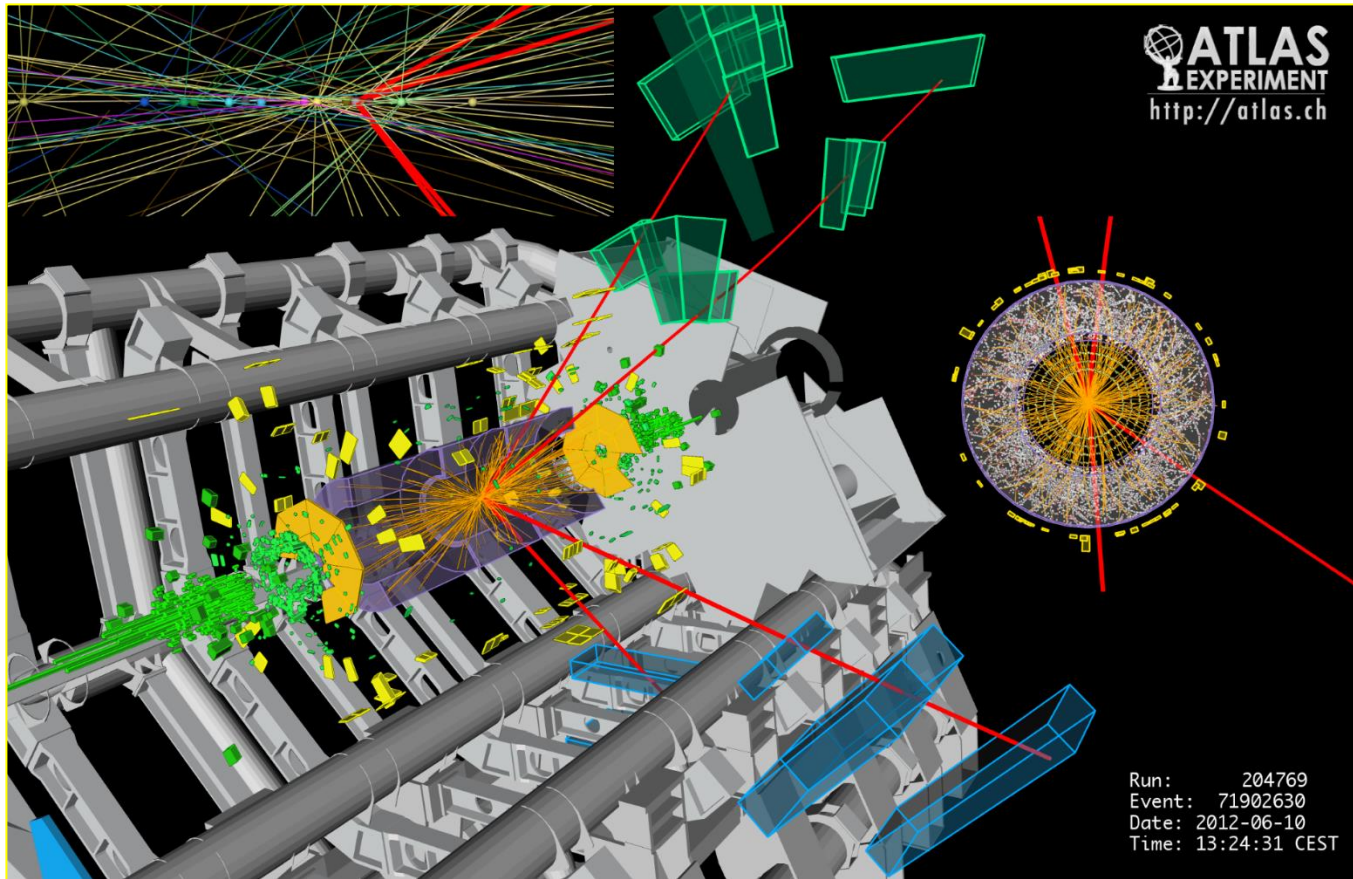




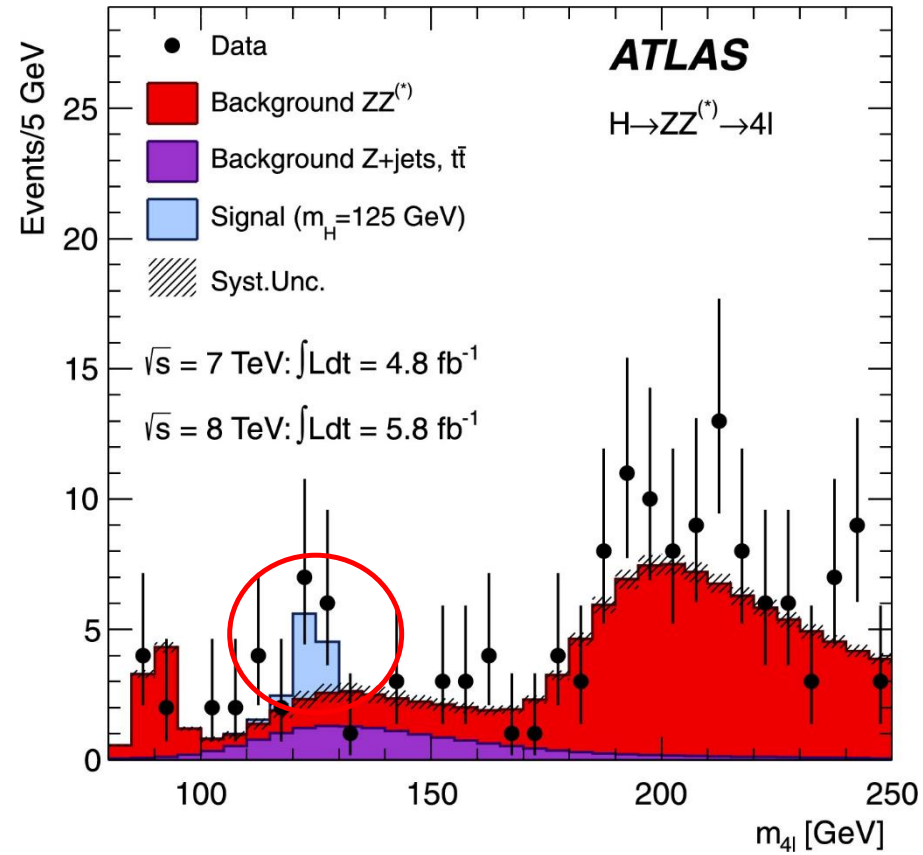
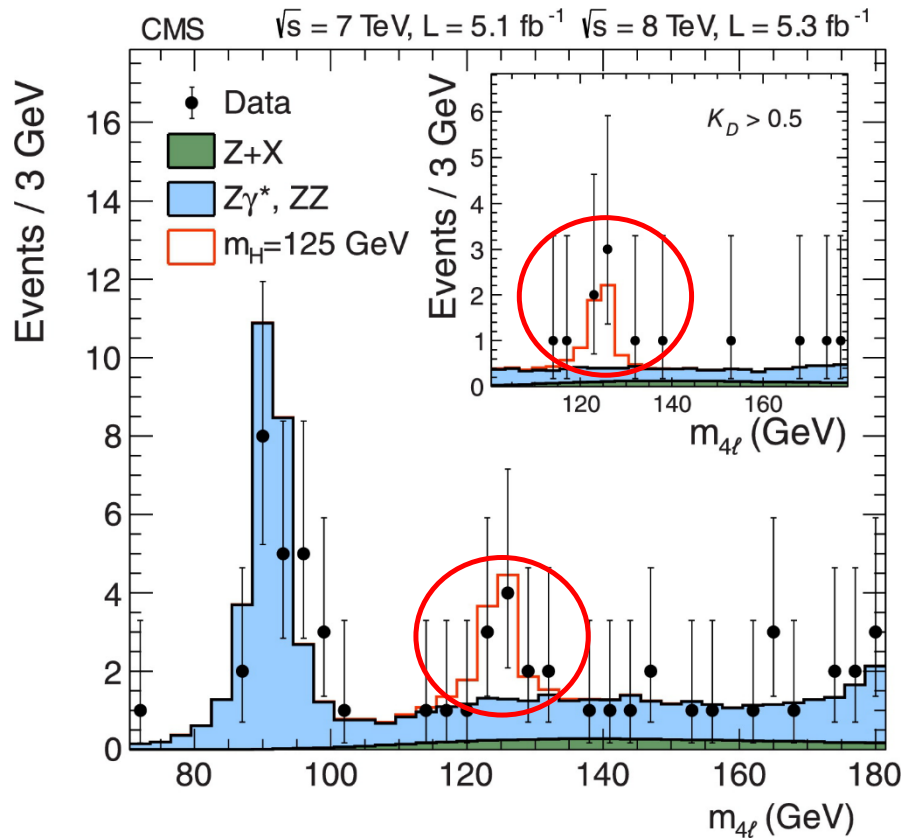
# The “giant” of LHC: ATLAS



$H \rightarrow Z^0 Z^0 \rightarrow 4 \text{ leptons}$

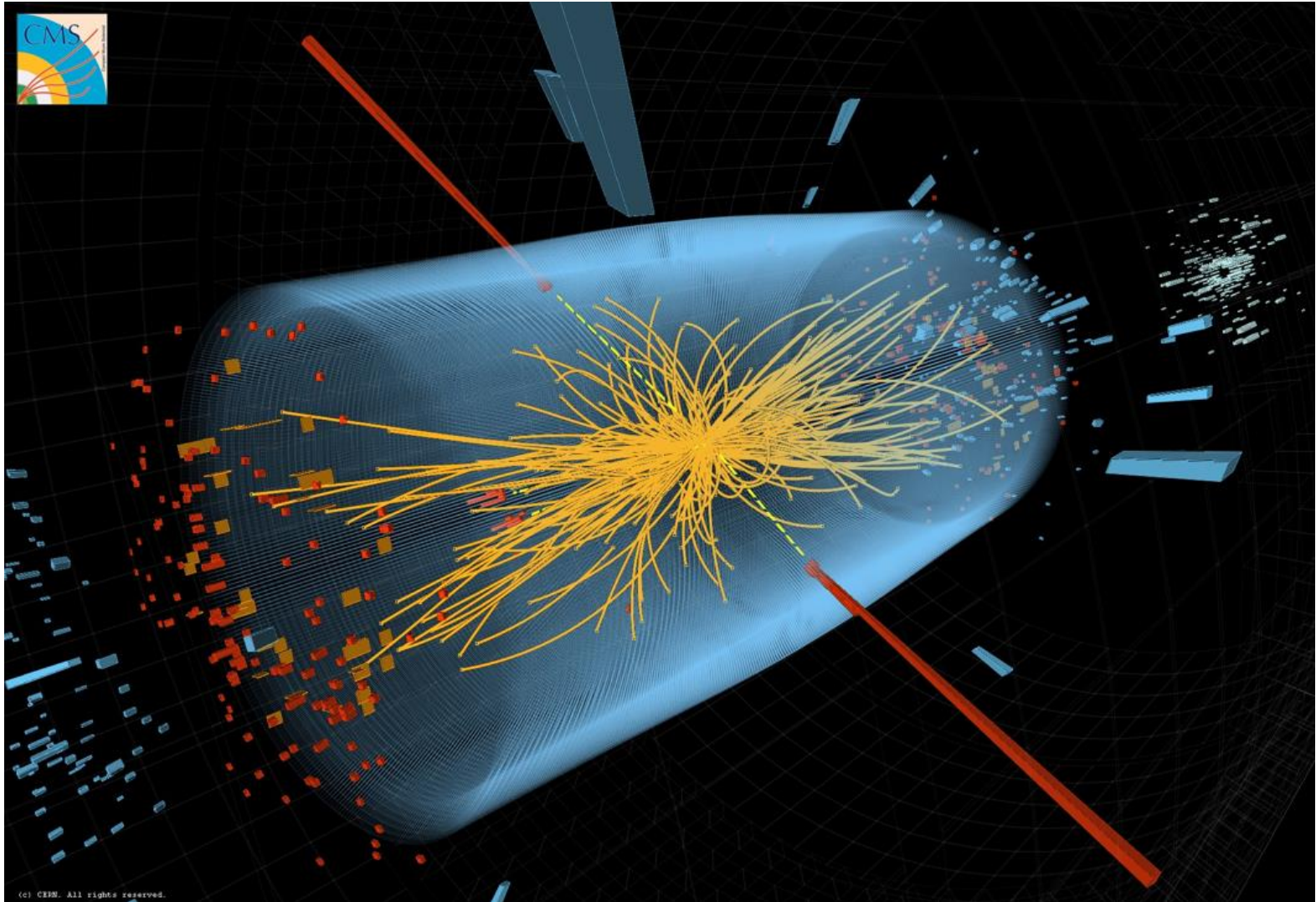


# Invariant mass plot $H \rightarrow ZZ \rightarrow 4 \text{ leptons}$



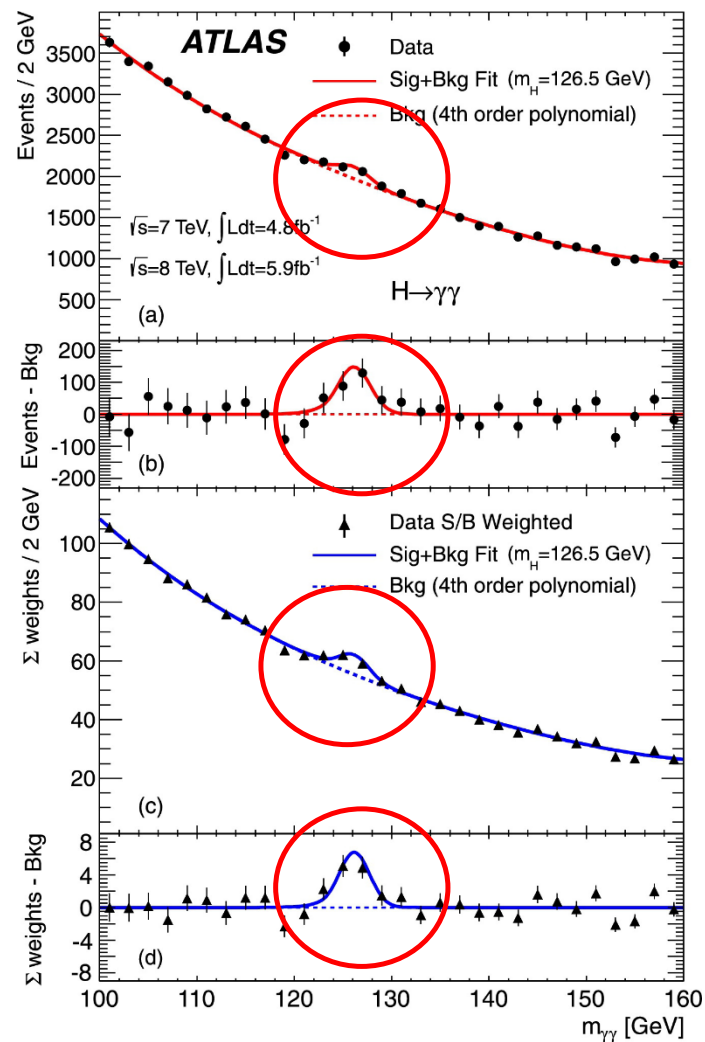
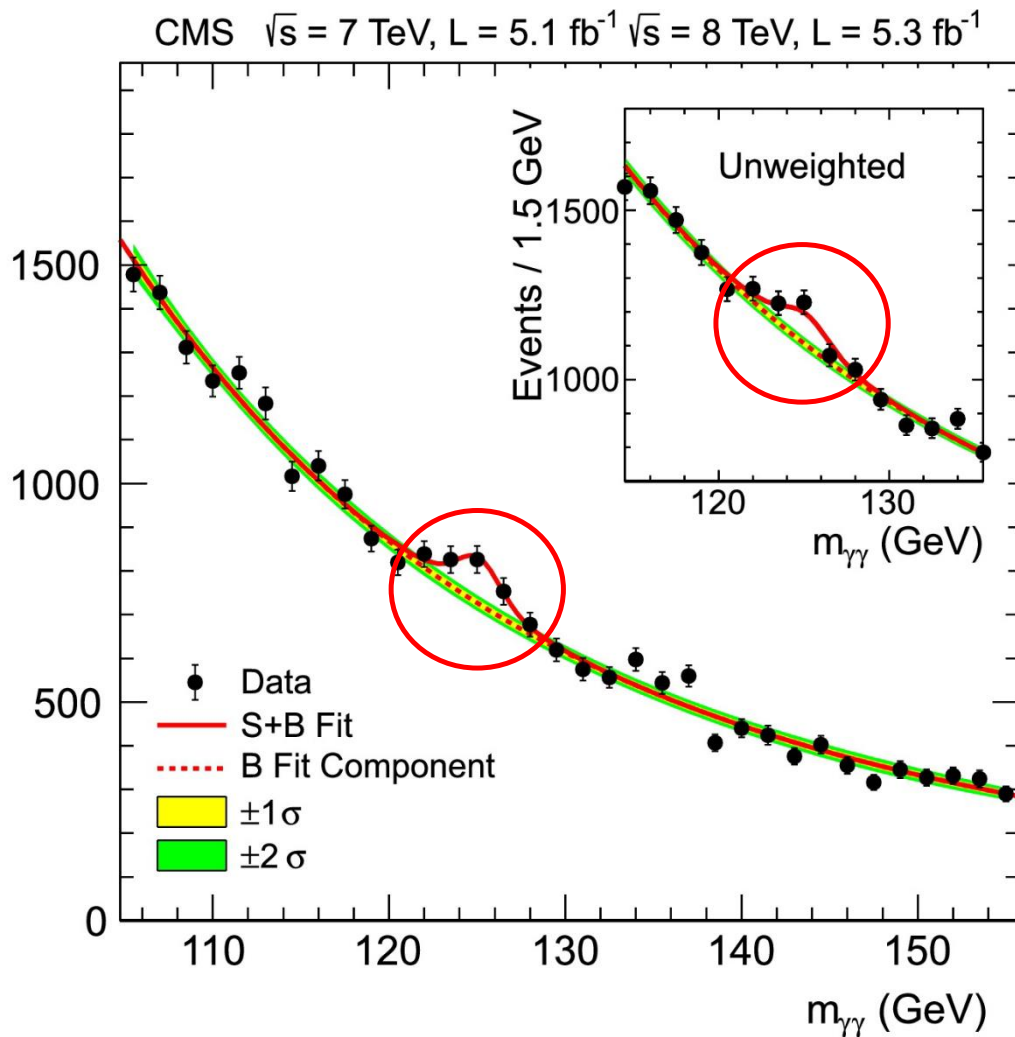


$$H \rightarrow \gamma\gamma$$

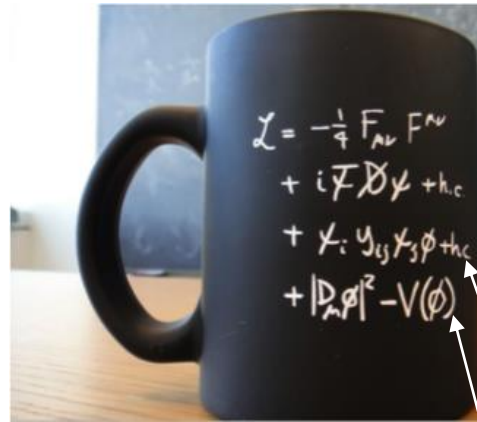


# H → $\gamma\gamma$

S/(S+B) Weighted Events / 1.5 GeV



# The SM ingredients



Kinetic term of all gauge bosons fields  
(W's, Z<sub>0</sub>, photon and gluons)

Kinetic term of fermion fields  
(quarks and leptons)

Interaction with H field giving mass to  
the q's and leptons

kinetic term for Higgs boson and H potential



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# Summary

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- From the classical Lagrangian to the particle physics Lagrangian
- Dirac equation and the anti-matter
- Feynman diagram and the computation of the muon the life-time
- Brief description of the SM Lagrangian .