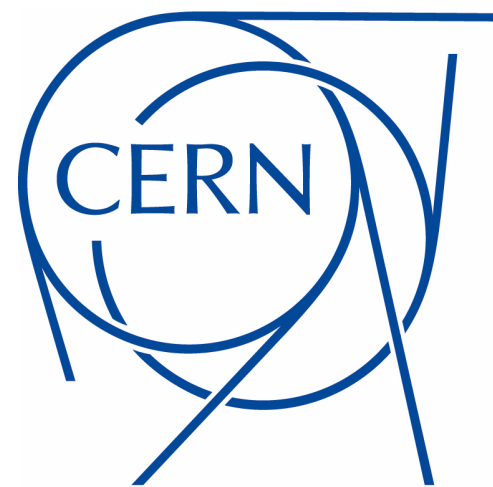
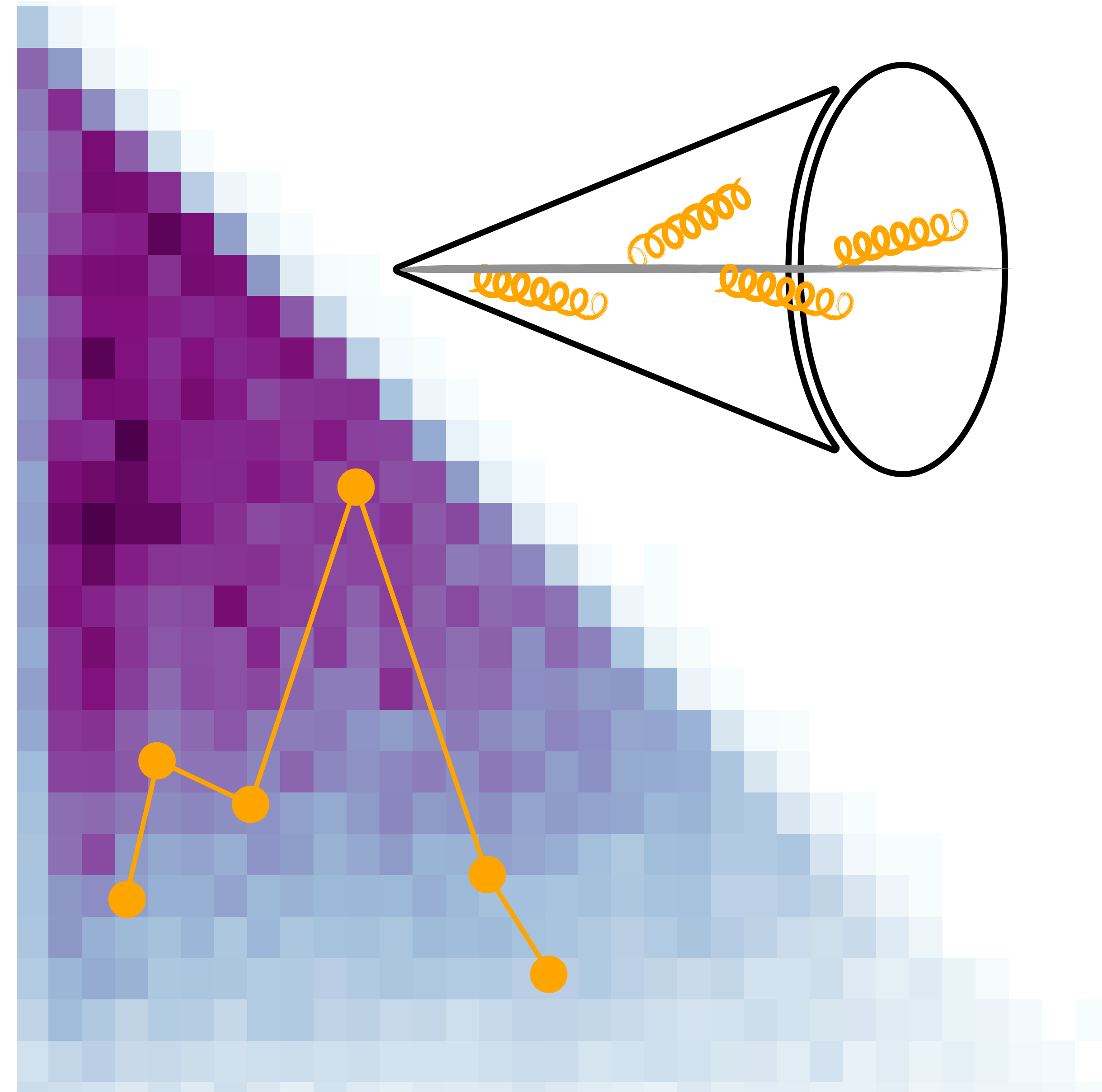


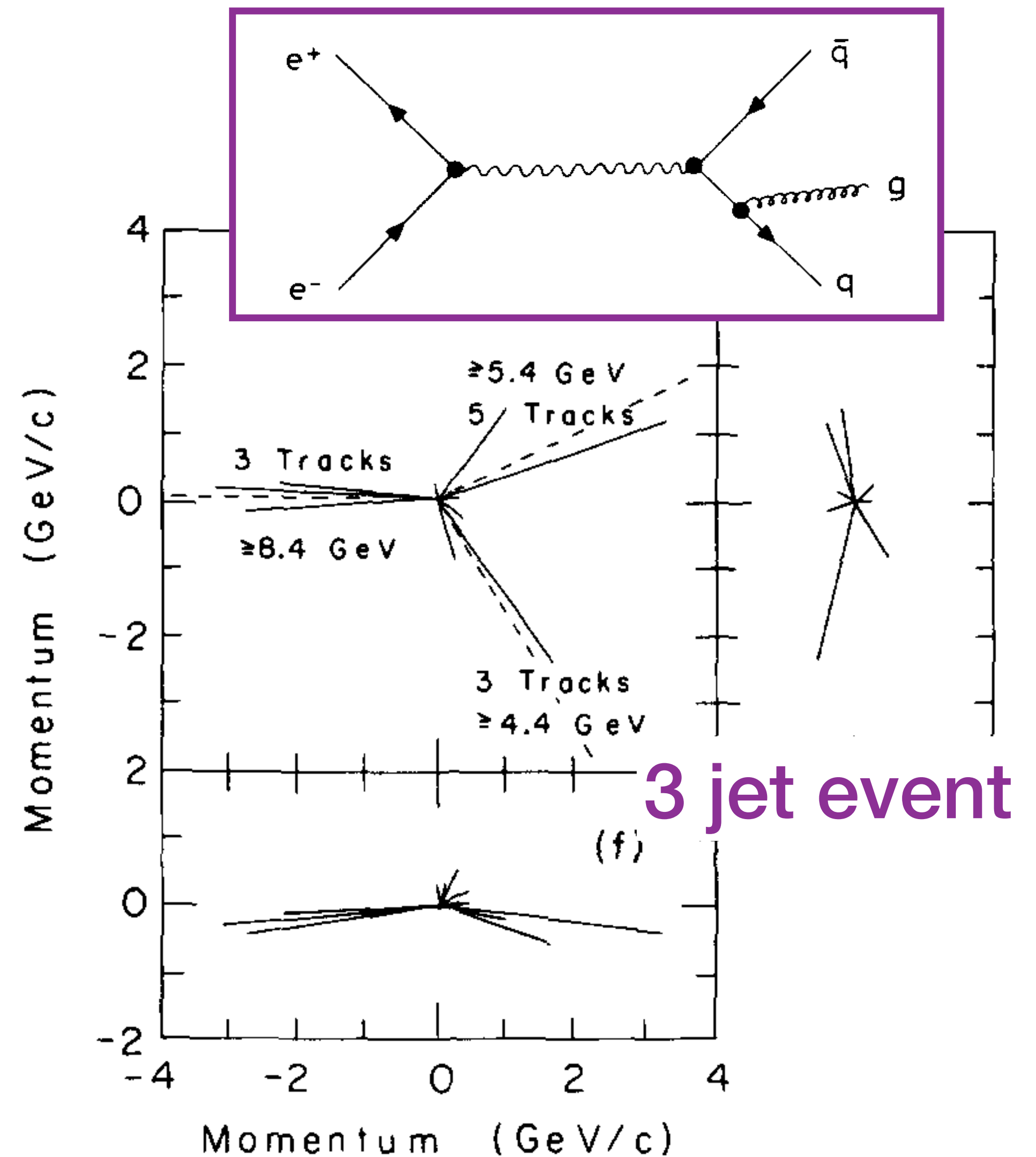
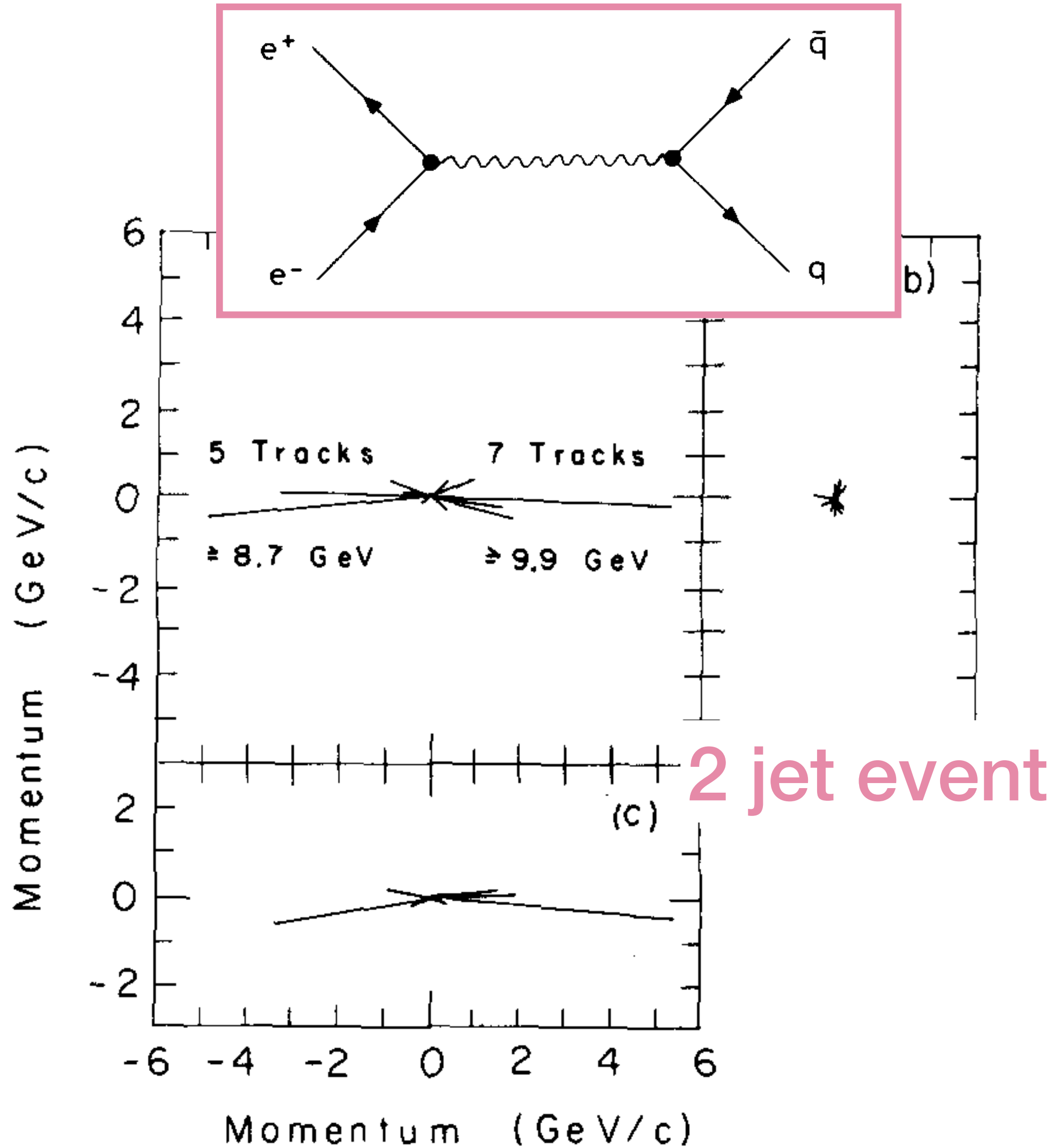
Counting jets in the Lund plane



Alba Soto-Ontoso
1st Lund jet plane Institute
CERN, 3rd July, 2023

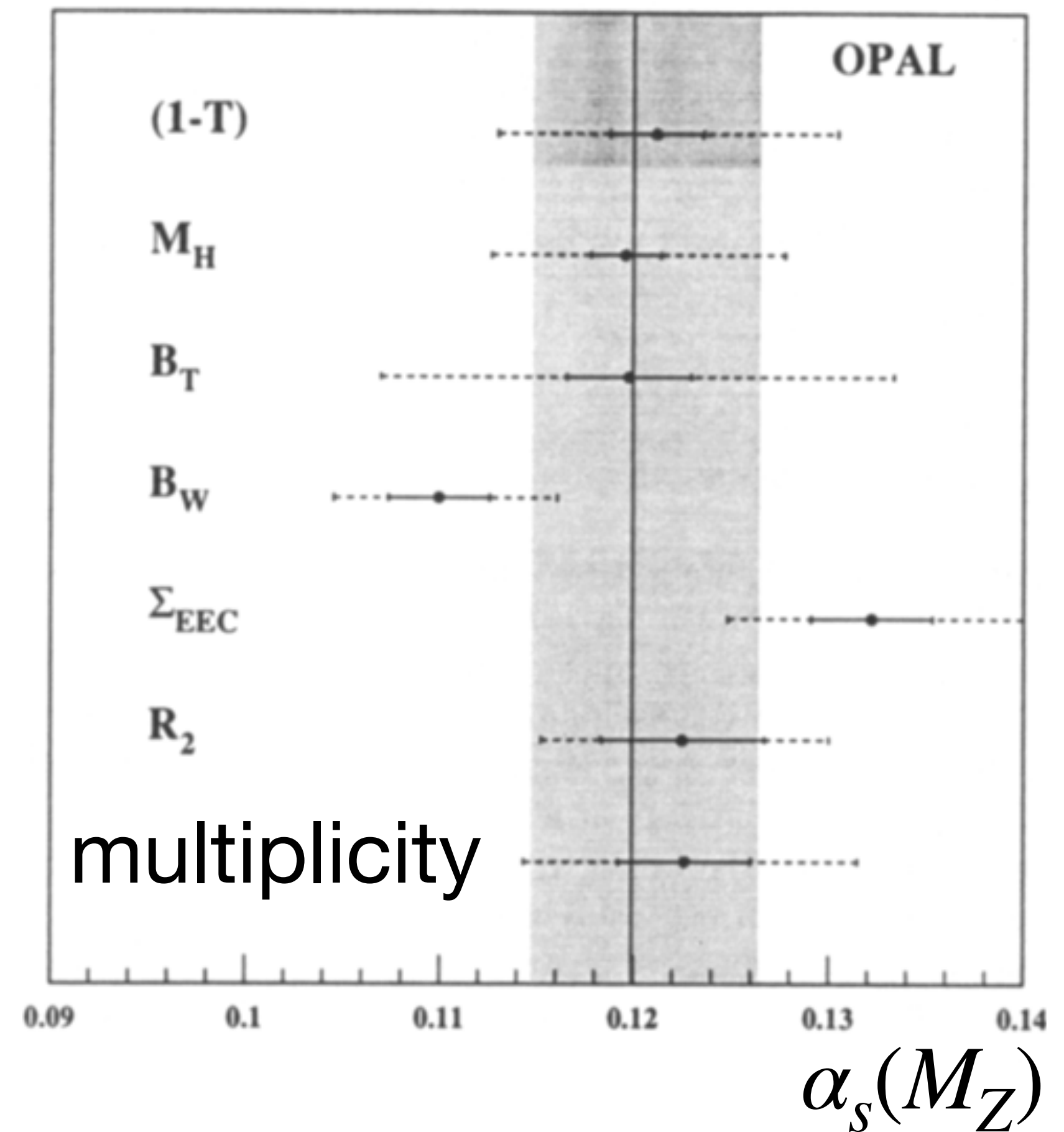
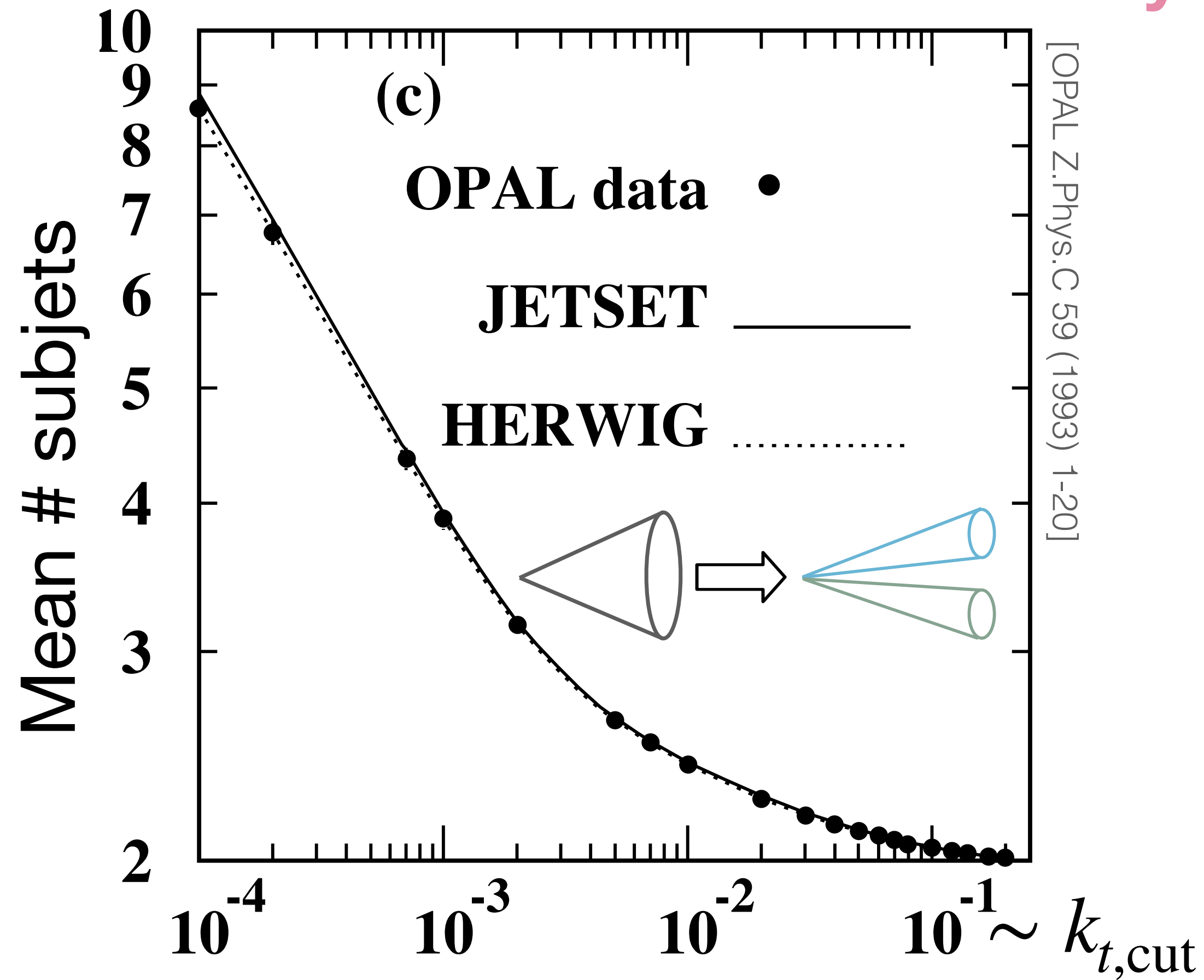


Why counting jets? e.g. to discover the gluon



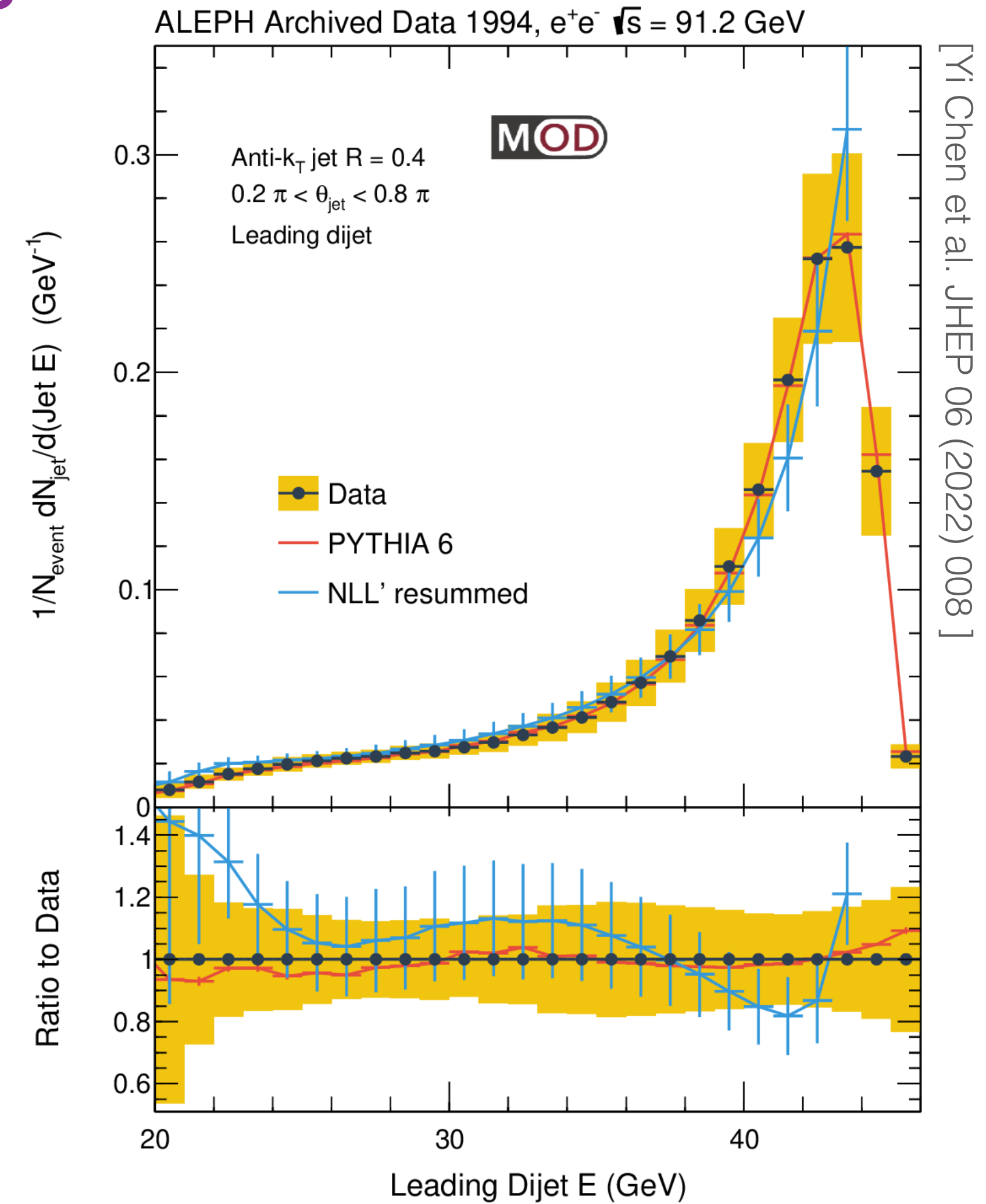
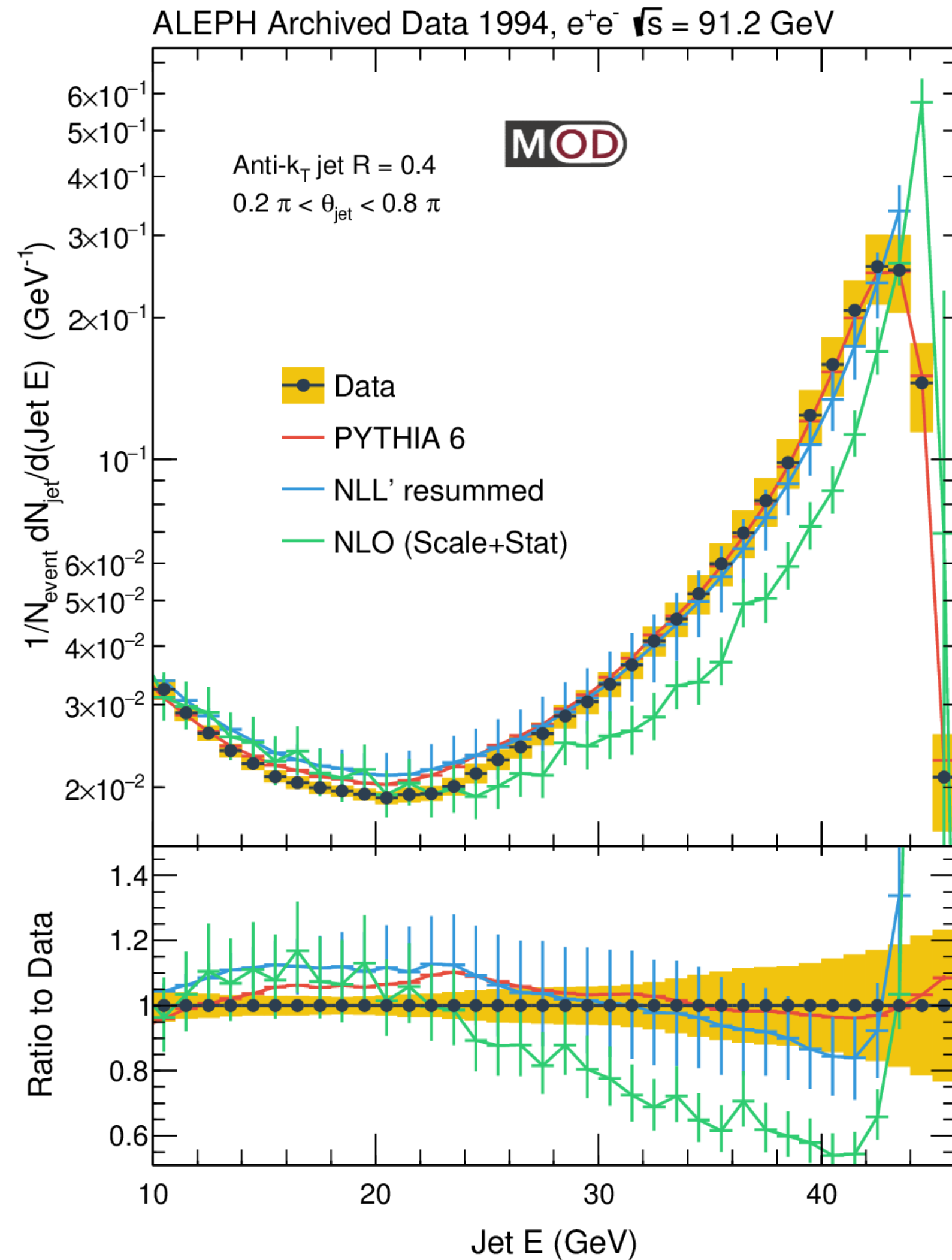
Why counting jets? e.g. to extract α_s in e^+e^- @ LEP

early 90's

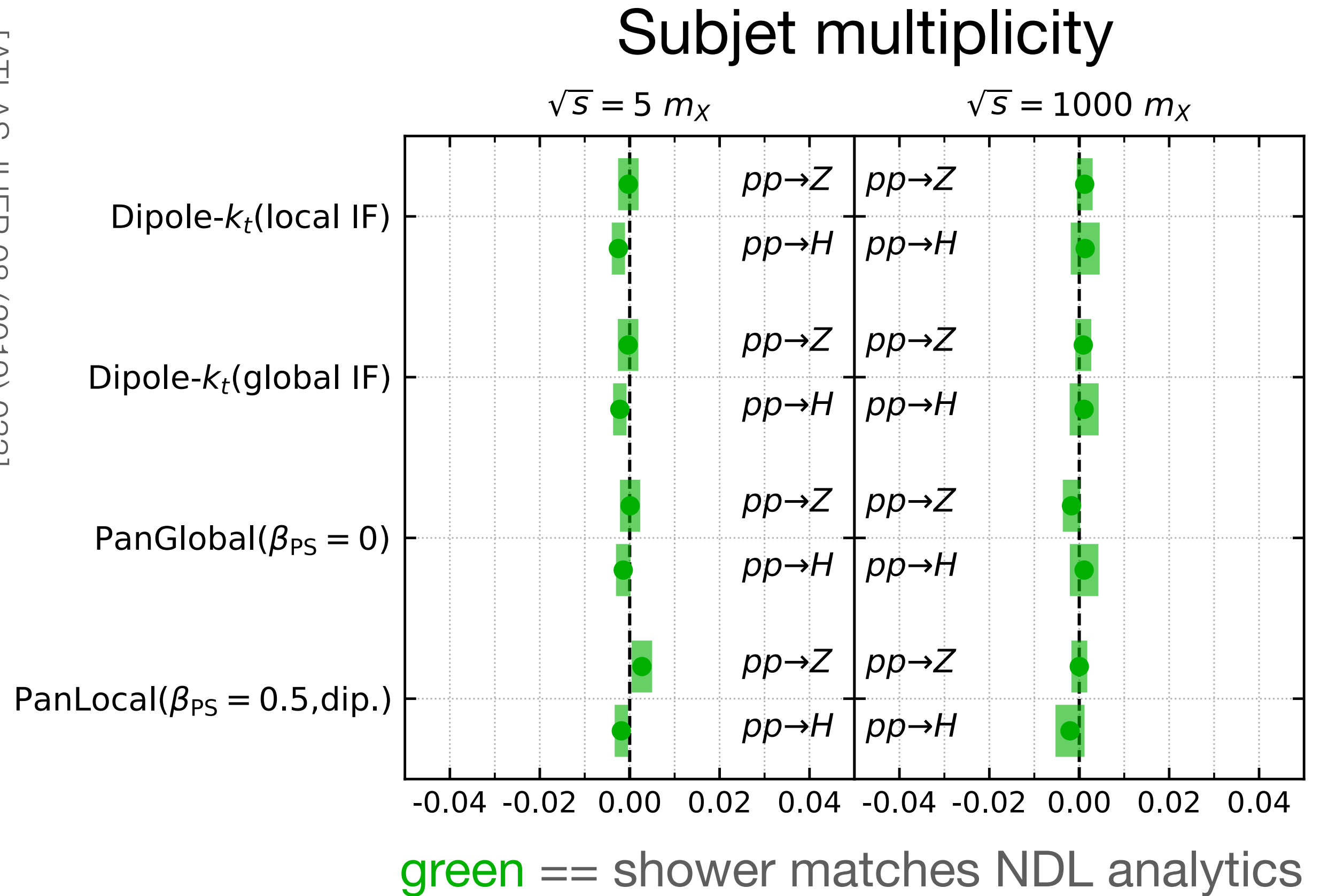
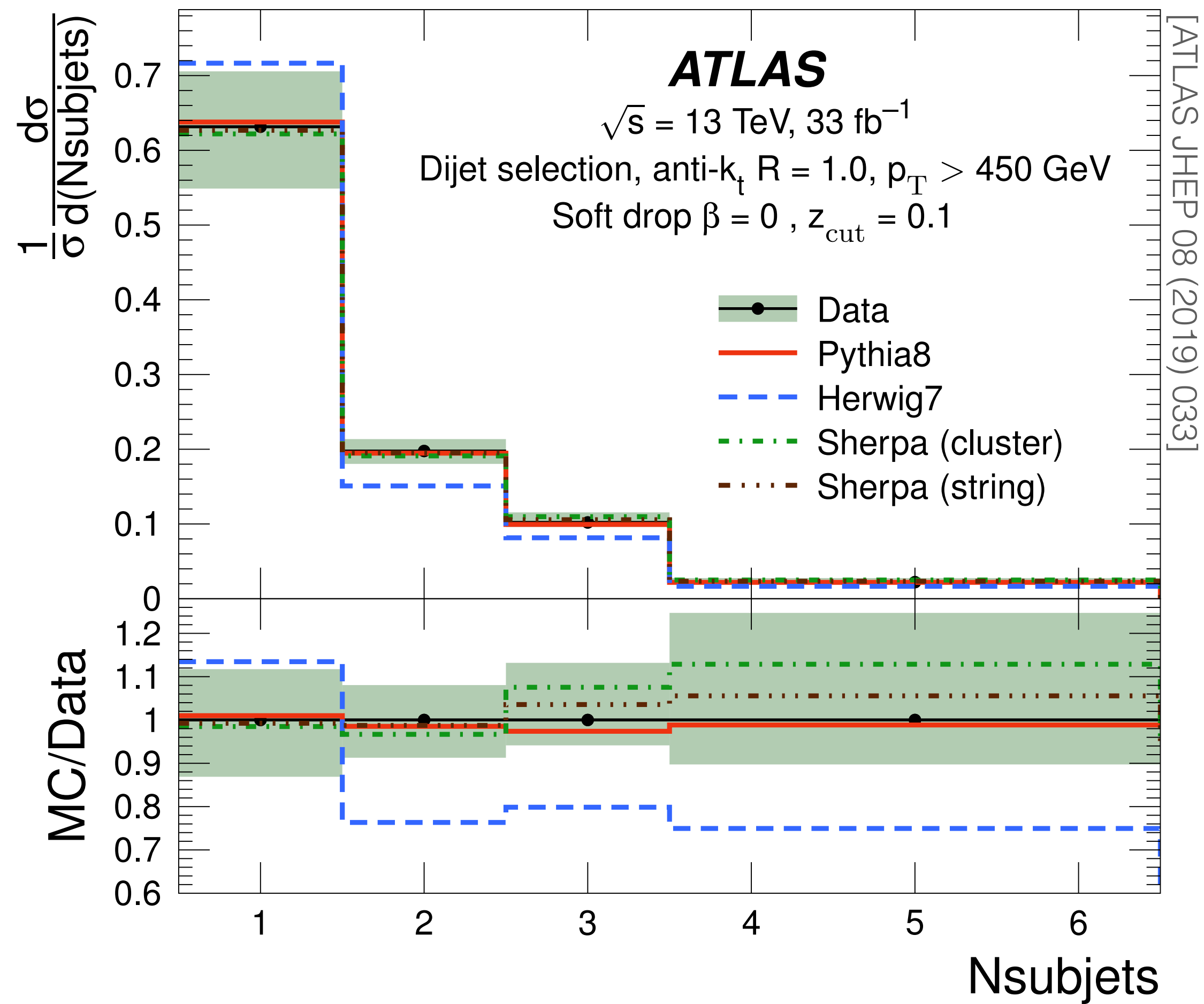


Why counting jets? e.g. to extract α_s in e^+e^- @ LEP

early 2020's

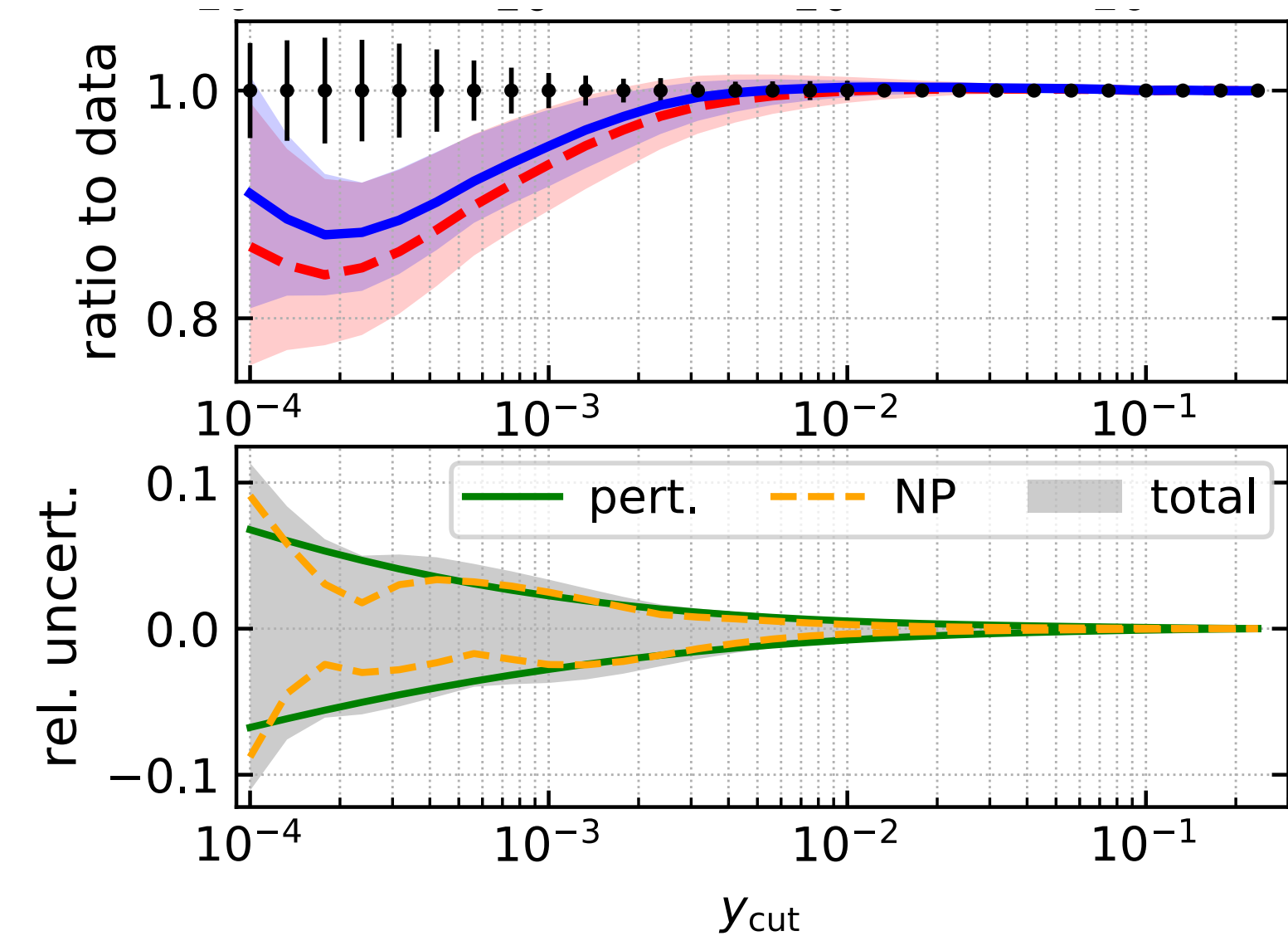
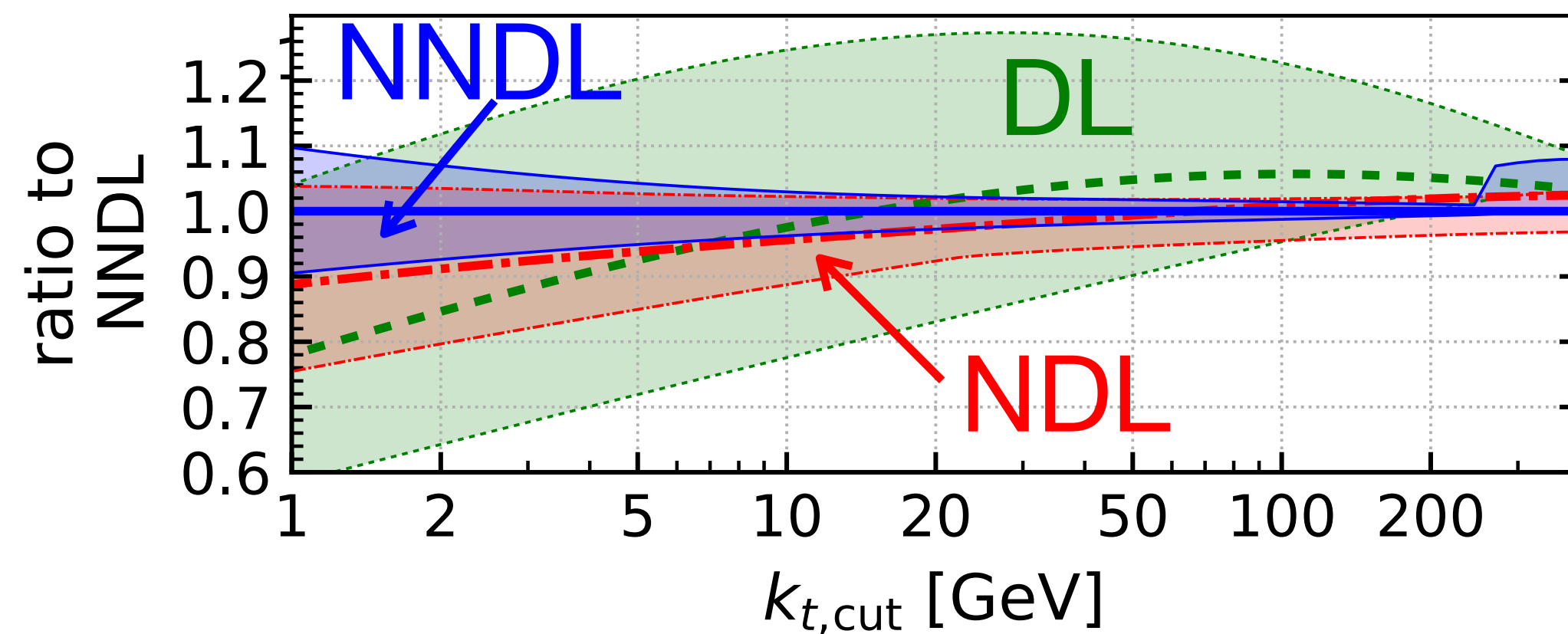


Why counting jets? e.g. to test parton showers

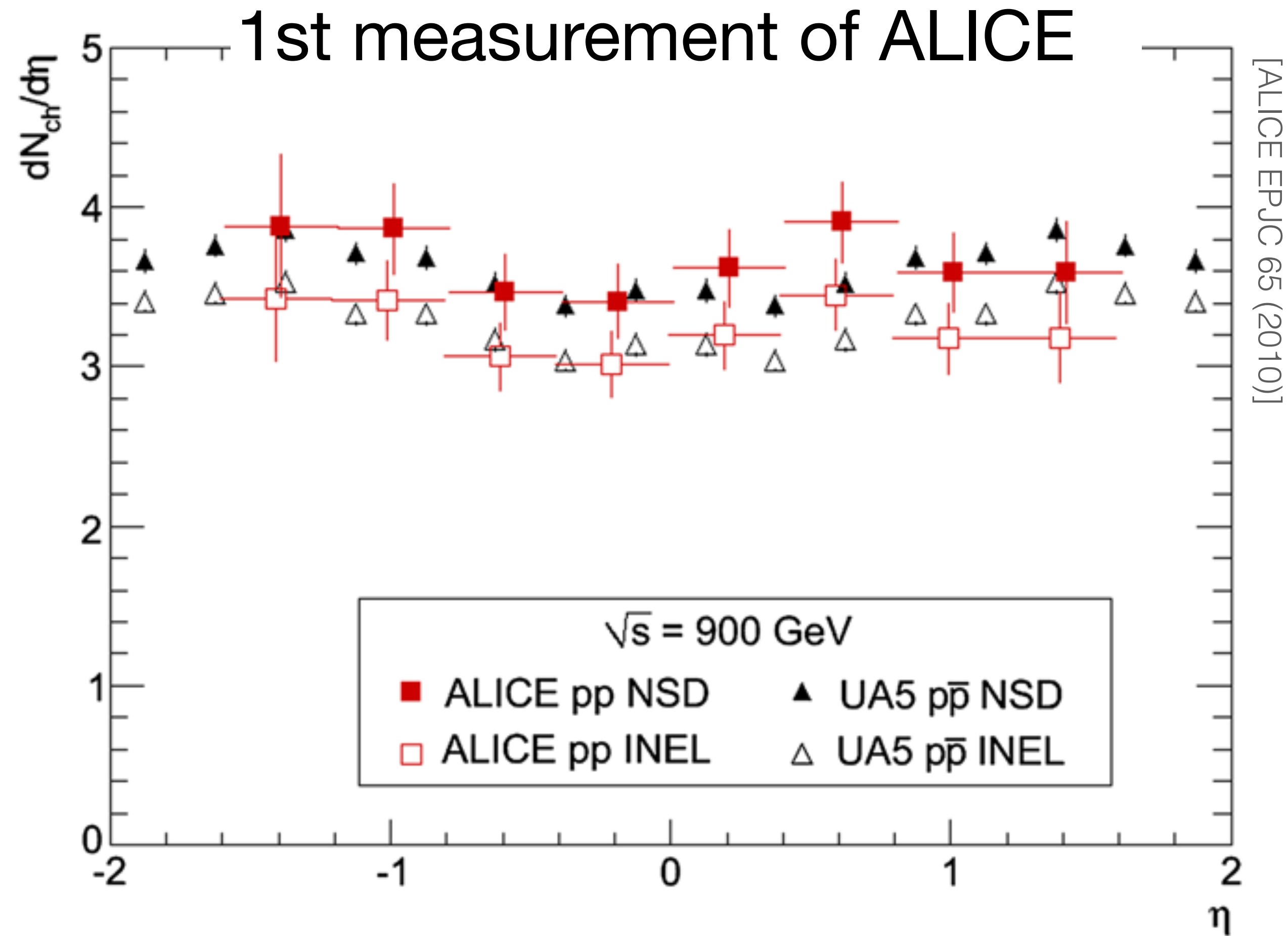


Outline of this talk

- 1 Multiplicity definition: charged particle, jets, subjets?
- 2 Lund multiplicity: resummation structure in both e^+e^- and pp
- 3 Phenomenological studies: NLO+NNDL+NP



Why counting jets and not particles?



Particle multiplicities are non-perturbative objects

Jet algorithms multiplicity

1

Choose a jet algorithm: k_t -like or angular ordered

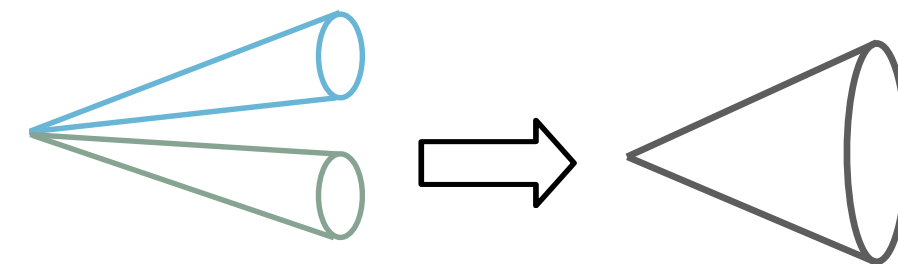
2

Choose a resolution parameter: e.g. $k_{t,\text{cut}}$

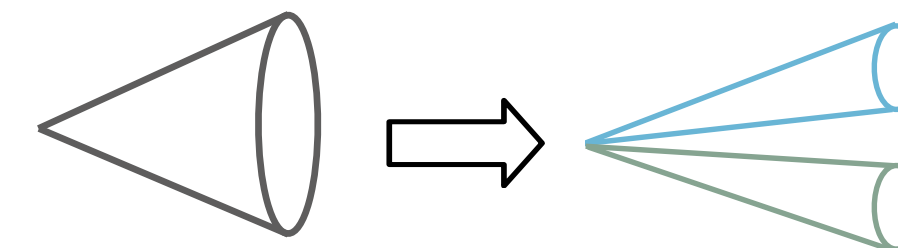
3

Count how many pairs of subjects satisfy $k_t > k_{t,\text{cut}}$

- either during the clustering



- or during the de-clustering



Choices impact calculability of the observable

Analytic structure of the average subjet multiplicity

The perturbative expansion of the average subjet multiplicity reads

$$\langle N(\alpha_s; L) \rangle = \left[\underbrace{h_1(\alpha_s L^2)}_{\substack{\alpha_s \ll 1 \\ \alpha_s L^2 \sim 1 \\ \text{DL}}} + \underbrace{\sqrt{\alpha_s} h_2(\alpha_s L^2)}_{\text{NDL}} + \underbrace{\alpha_s h_3(\alpha_s L^2)}_{\text{NNDL} \sim 10\%} + \dots \right] \quad L = \ln \frac{Q}{k_{t,\text{cut}}}$$

where N^kDL accuracy implies control over $\alpha_s^n L^{2n-k}$ terms with $0 < n < \infty$

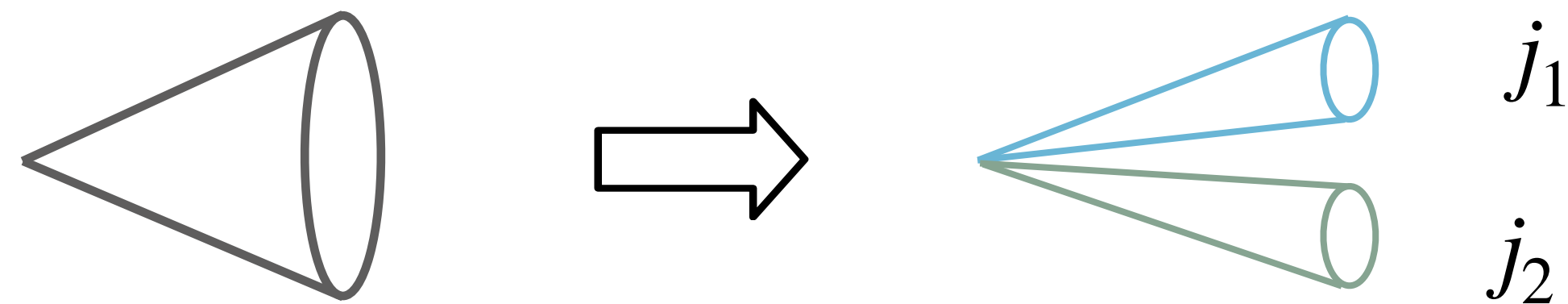
Jet algorithm and/or exact definition of $k_t > k_{t,\text{cut}}$ matter at **NNDL**

Lund-based definition of subjet multiplicity

[Medves, ASO, Soyez, JHEP 10 (2022) 156, JHEP 04 (2023) 104]

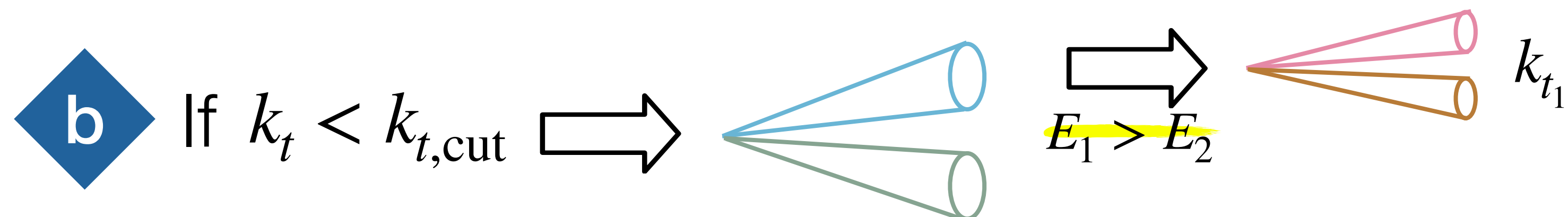
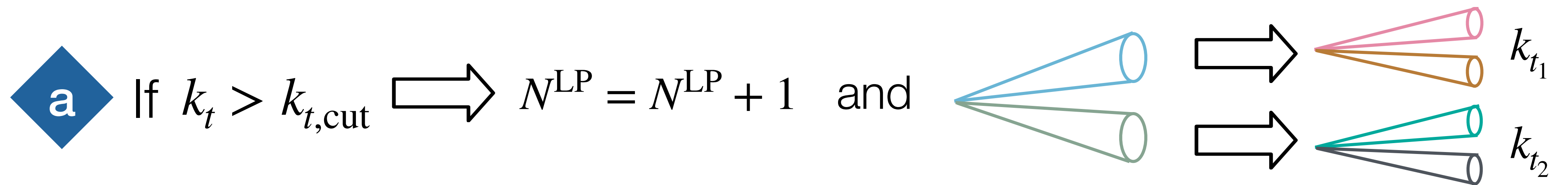
1 Recluster a jet with the Cambridge/Aachen algorithm

2 Traverse backwards the angular ordered sequence



$$k_t^{e^+e^-} = \min(E_1, E_2) \sin \theta_{12}$$

$$k_t^{pp} = \min(p_{t1}, p_{t2}) \Delta R_{12}$$

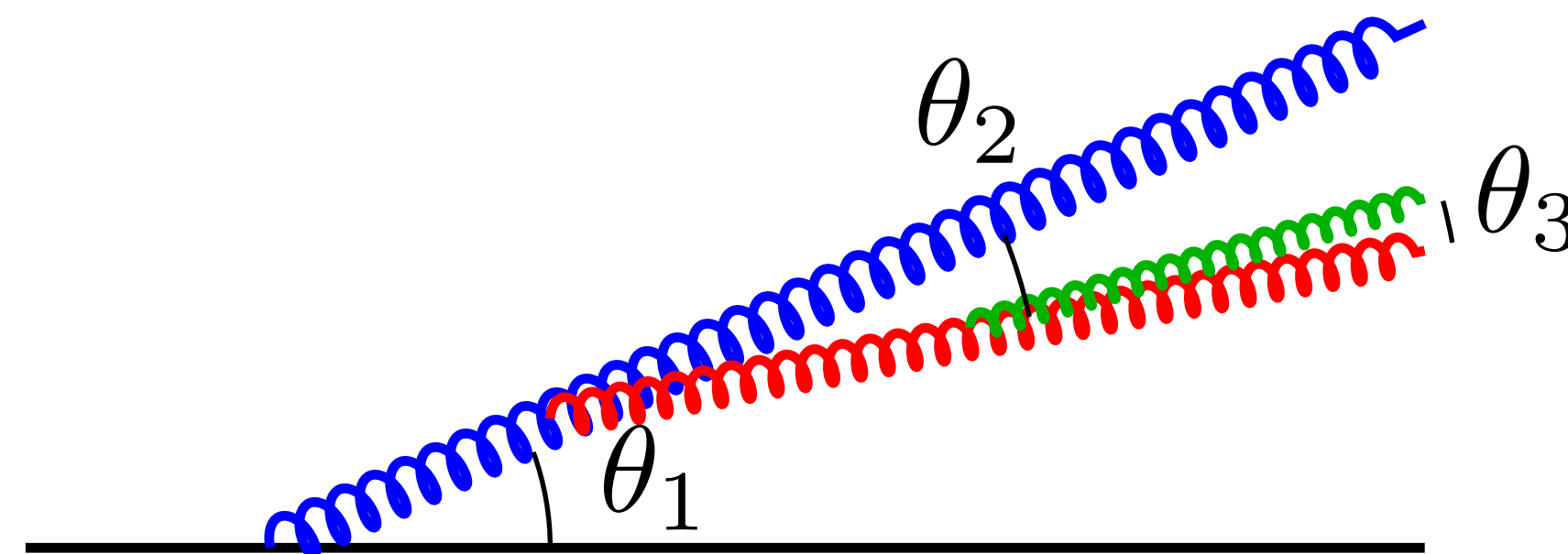


Lund multiplicity: all orders (DL) $(\alpha_s L^2)^n$

$$\langle N \rangle_{\text{DL}} = 1 + \frac{C_i}{C_A} \sum_{n=1}^{\infty} \bar{\alpha}^n \underbrace{\int_0^{\infty} d\eta_1 \int_{\eta_1}^{\infty} d\eta_2 \dots \int_{\eta_{n-1}}^{\infty} d\eta_n}_{\text{angular-ordering}} \underbrace{\int_0^1 \frac{dx_1}{x_1} \int_0^{x_1} \frac{dx_2}{x_2} \dots \int_0^{x_{n-1}} \frac{dx_n}{x_n}}_{\text{energy-ordering}} \underbrace{\Theta(x_n e^{-\eta_n} > e^{-L})}_{k_t > k_{t,\text{cut}}}$$

$$Q \gg E_1 \gg E_2 \gg E_3 \gg k_{t,\text{cut}}$$

$$1 \gg \theta_1 \gg \theta_2 \gg \theta_3$$

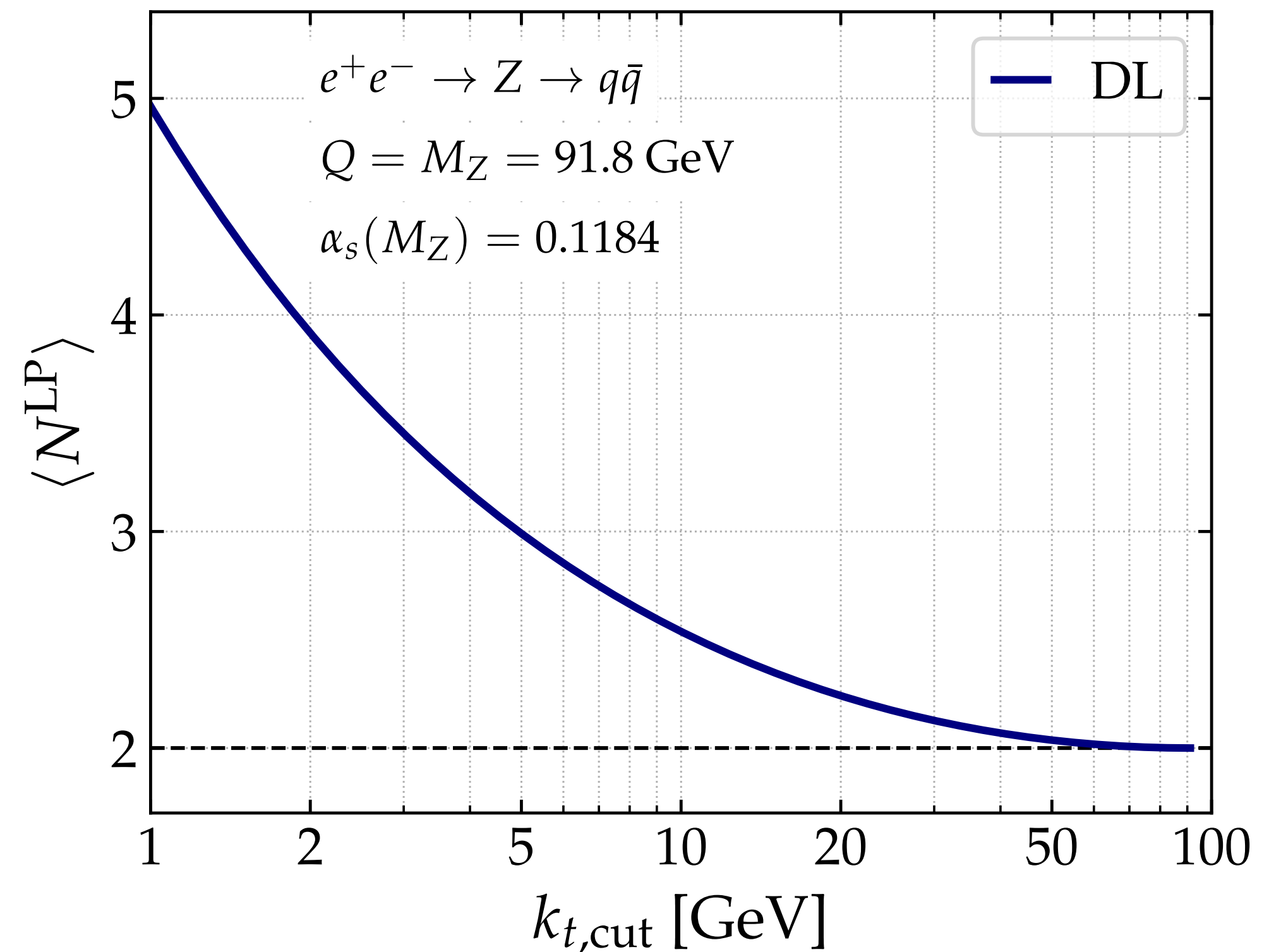


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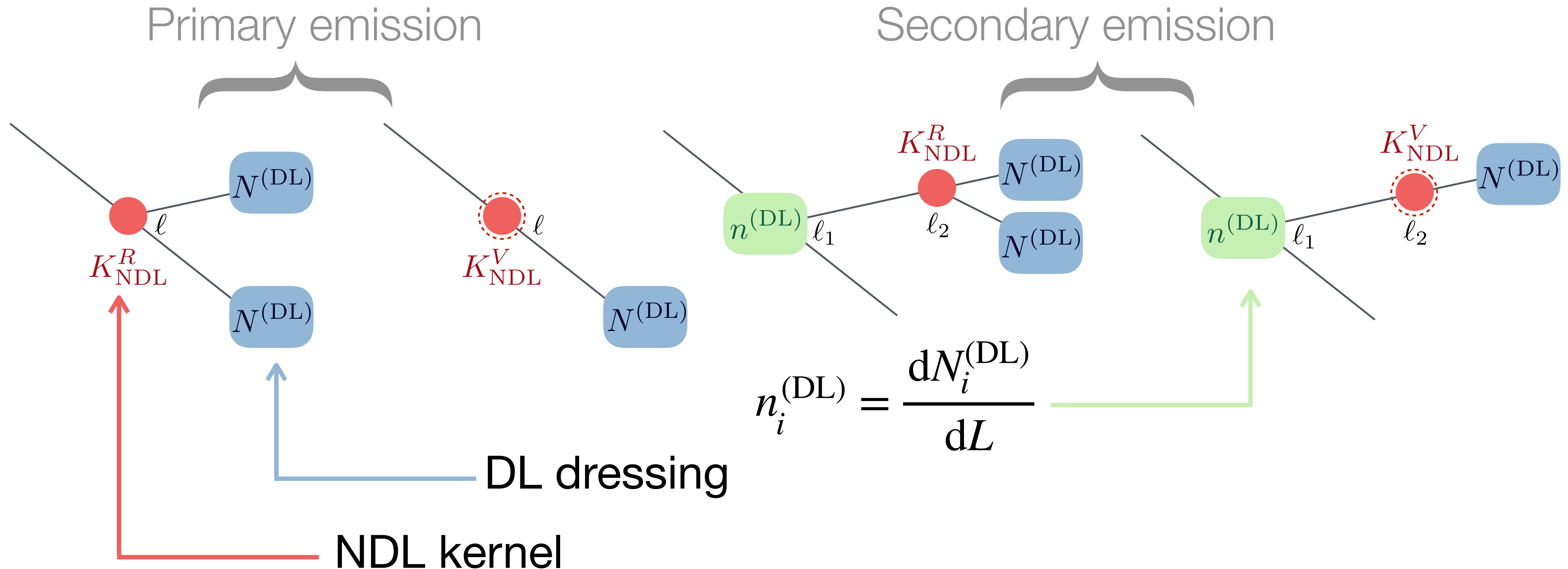
$$\langle N \rangle_{\text{DL}} = 1 + \frac{C_i}{C_A} [\cosh \nu - 1]$$

$$\nu = \sqrt{2\alpha_s C_A L^2 / \pi}$$



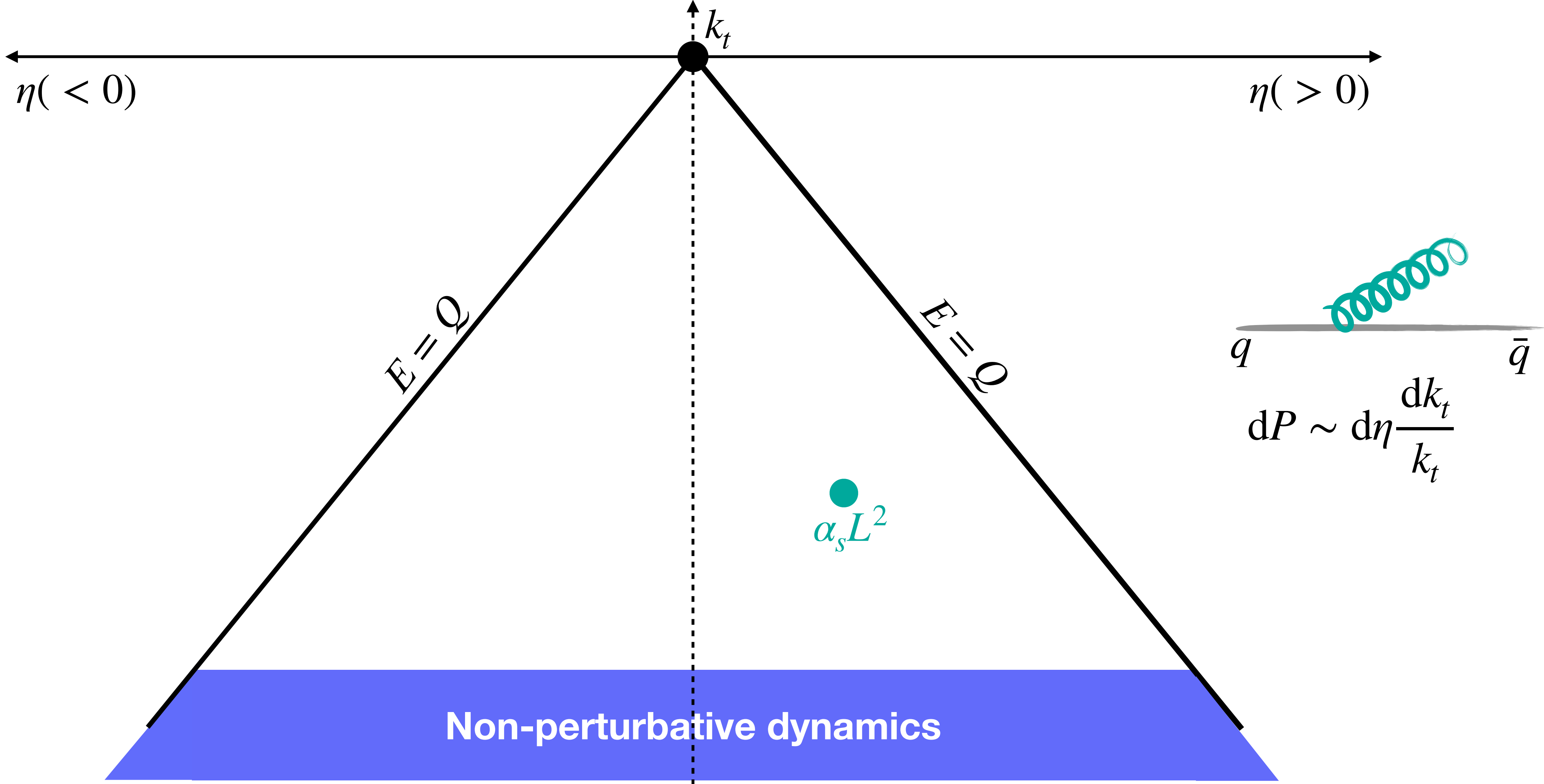
Subject multiplicity: NDL resummation strategy

NDL == one NDL-like emission in the chain at a momentum scale ℓ

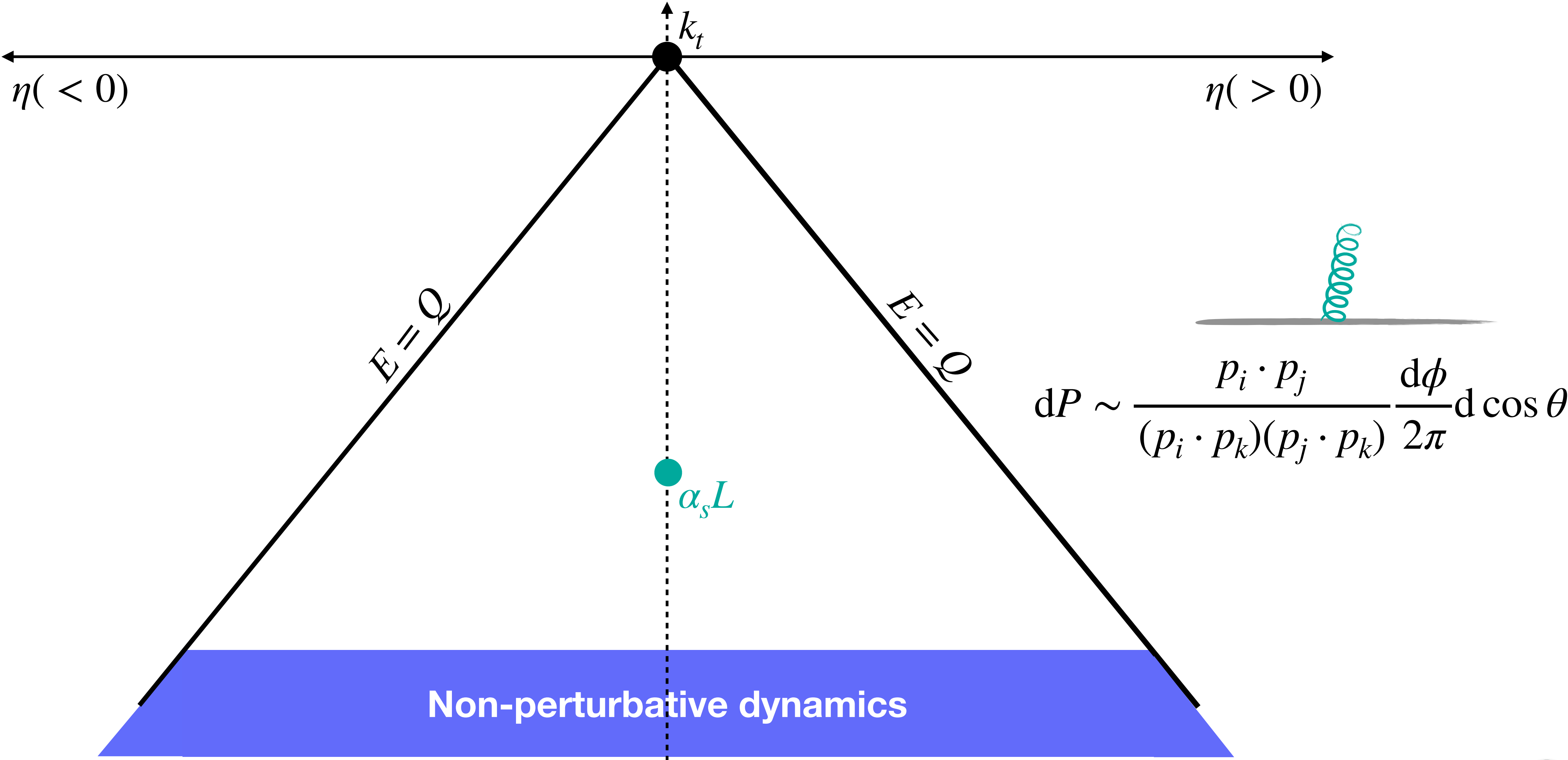


Idea: use the Lund plane to systematically identify all NDL terms

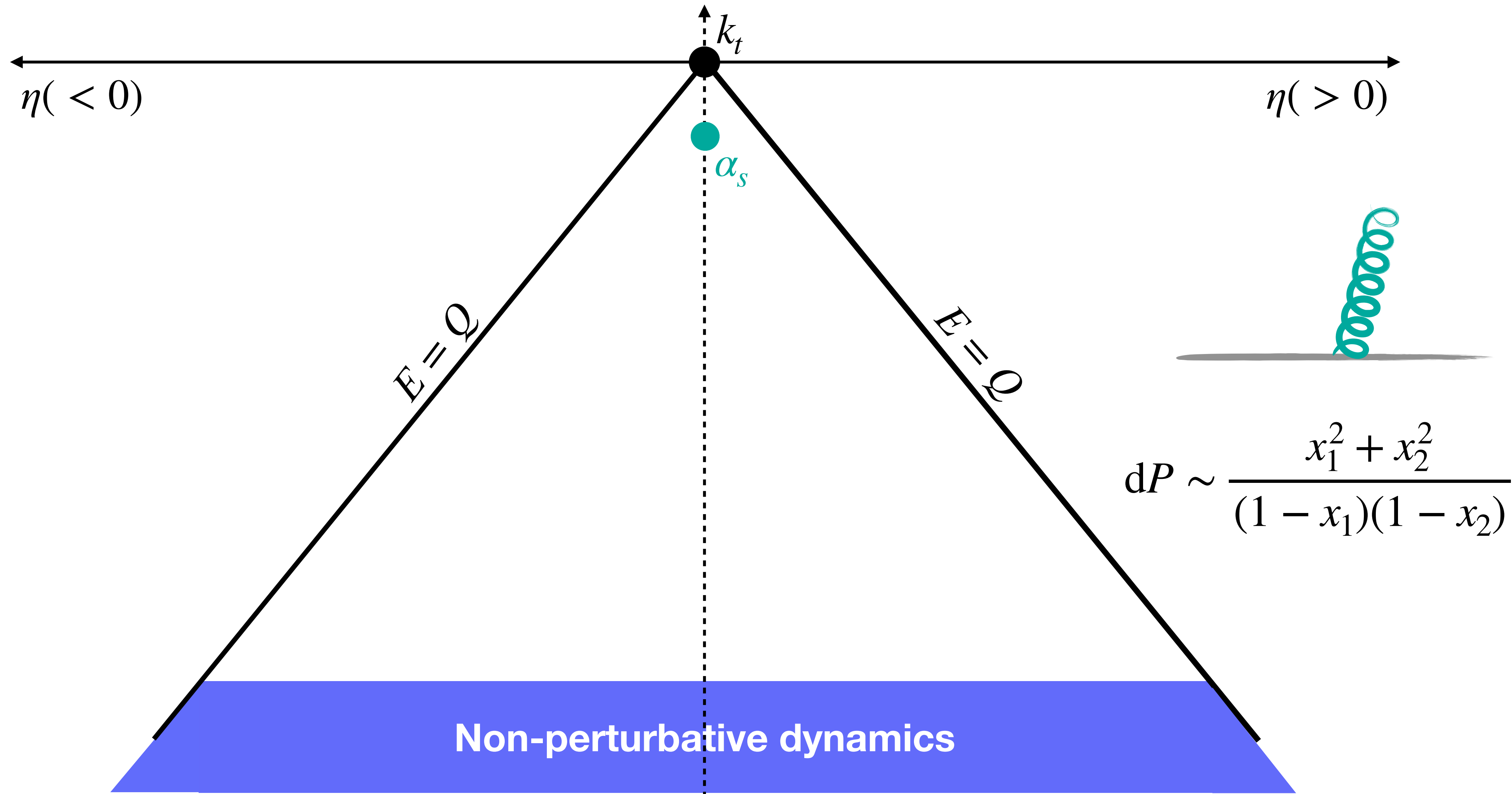
Interlude: counting logs in the Lund jet plane [Dreyer, Salam, Soyez JHEP 12 (2018) 064]



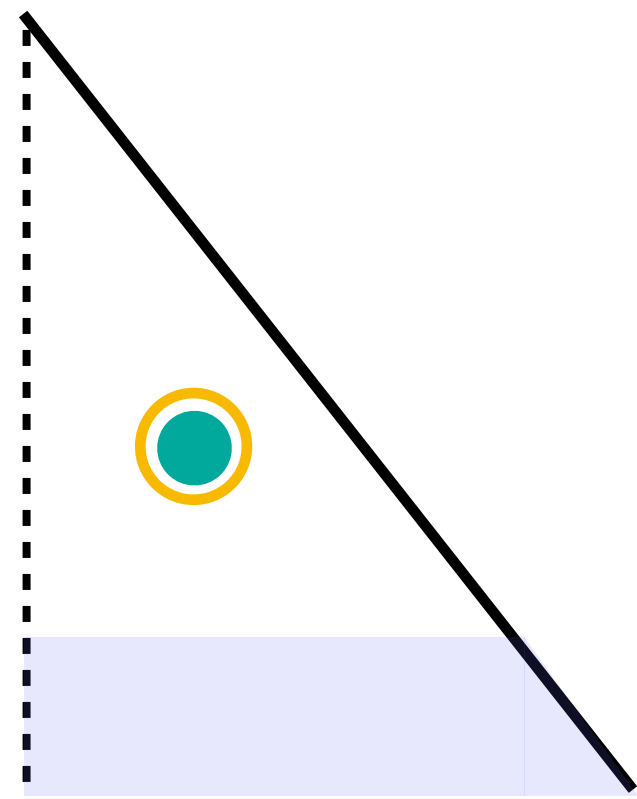
Interlude: counting logs in the Lund jet plane [Dreyer, Salam, Soyez JHEP 12 (2018) 064]



Interlude: counting logs in the Lund jet plane [Dreyer, Salam, Soyez JHEP 12 (2018) 064]



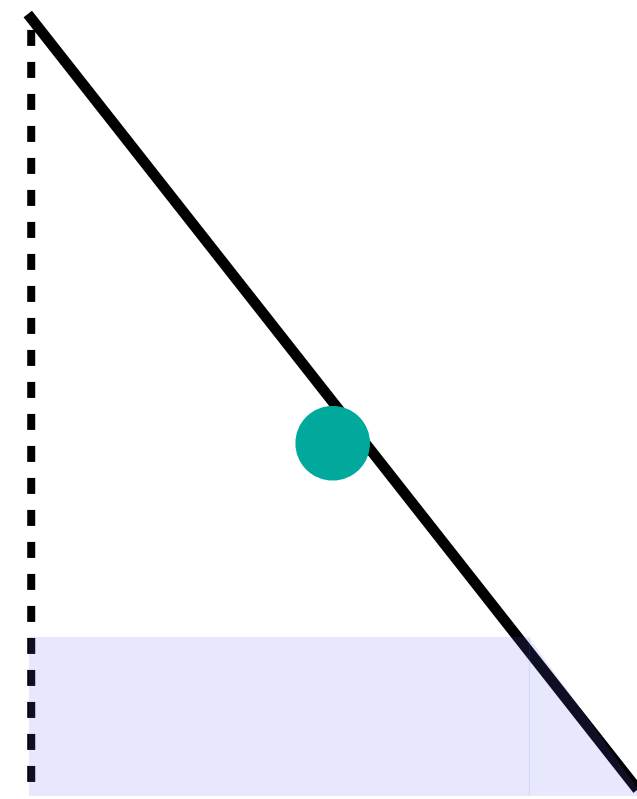
Lund multiplicity: NDL resummation $\alpha_s L (\alpha_s L^2)^n$



Running coupling

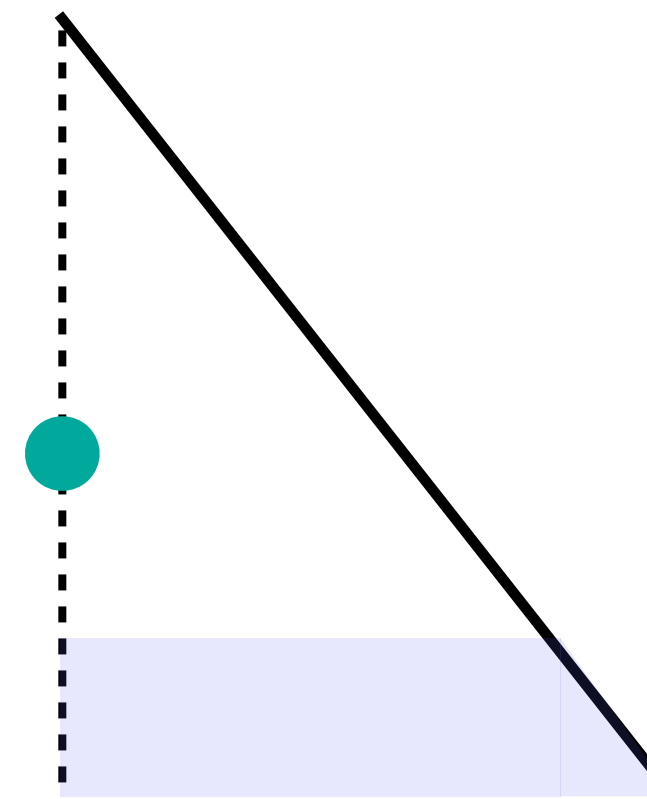
$$\alpha_s \rightarrow \alpha_s - 2\alpha_s^2 \beta_0 \ell + \mathcal{O}(\alpha_s^3)$$

with $\ell \equiv \ln(k_t/Q)$



Hard-collinear

$$\frac{1}{z} \rightarrow C_F \left(\frac{1-z}{z} + \frac{z}{2} \right)$$

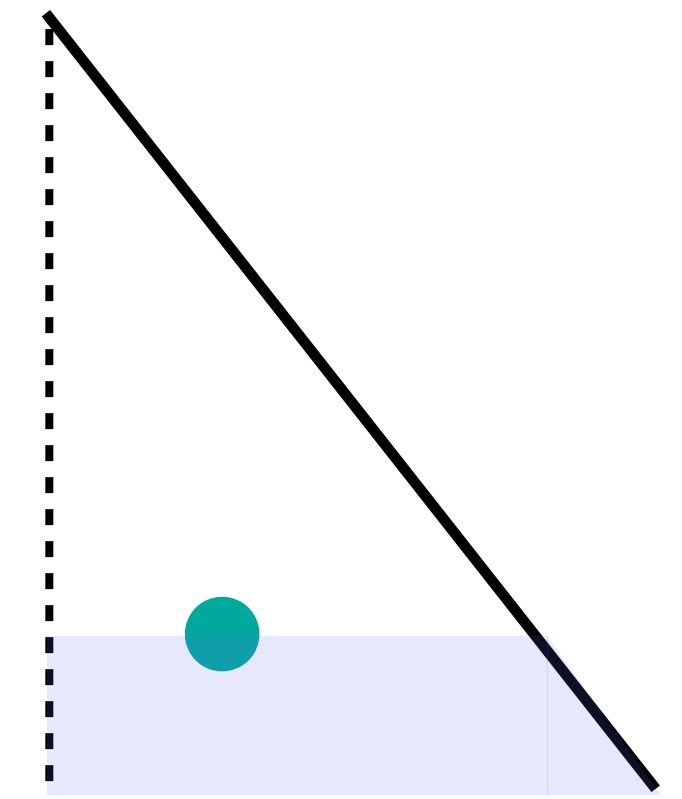


Large-angle

$$\frac{dz}{z} d\eta$$

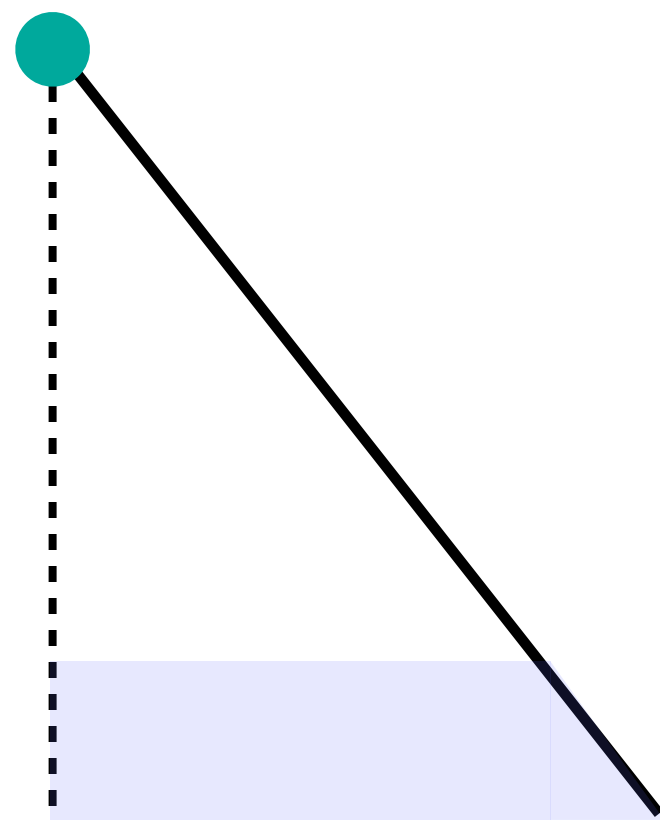
↓

$$\frac{p_i \cdot p_j}{p_i \cdot p_k p_j \cdot p_k} \frac{d\phi}{2\pi} d \cos \theta$$

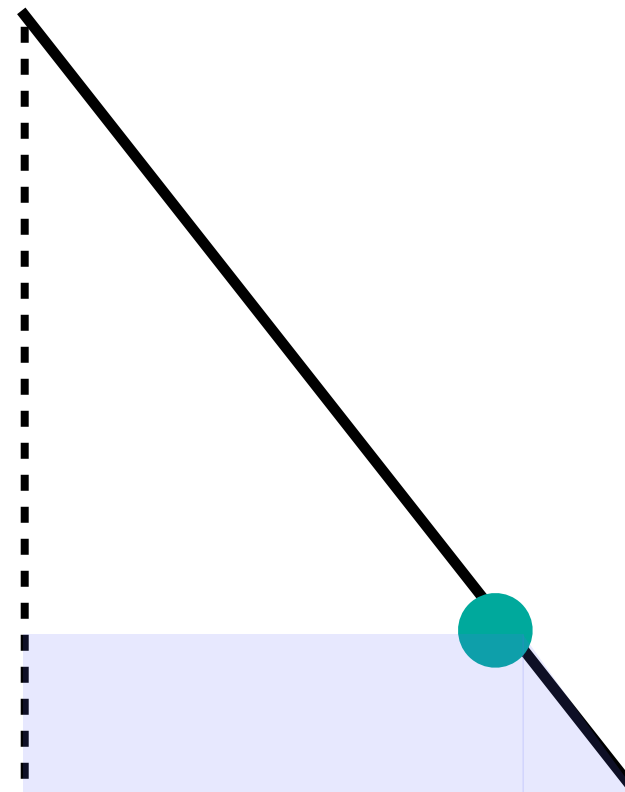


$k_t \sim k_{t,\text{cut}}$

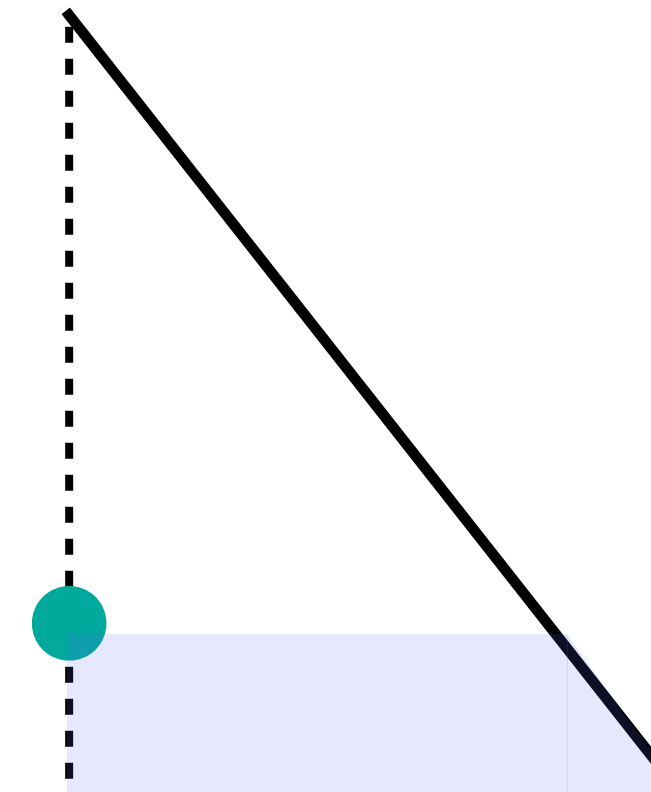
Lund multiplicity: NNLL resummation $\alpha_s(\alpha_s L^2)^n$



Hard ME

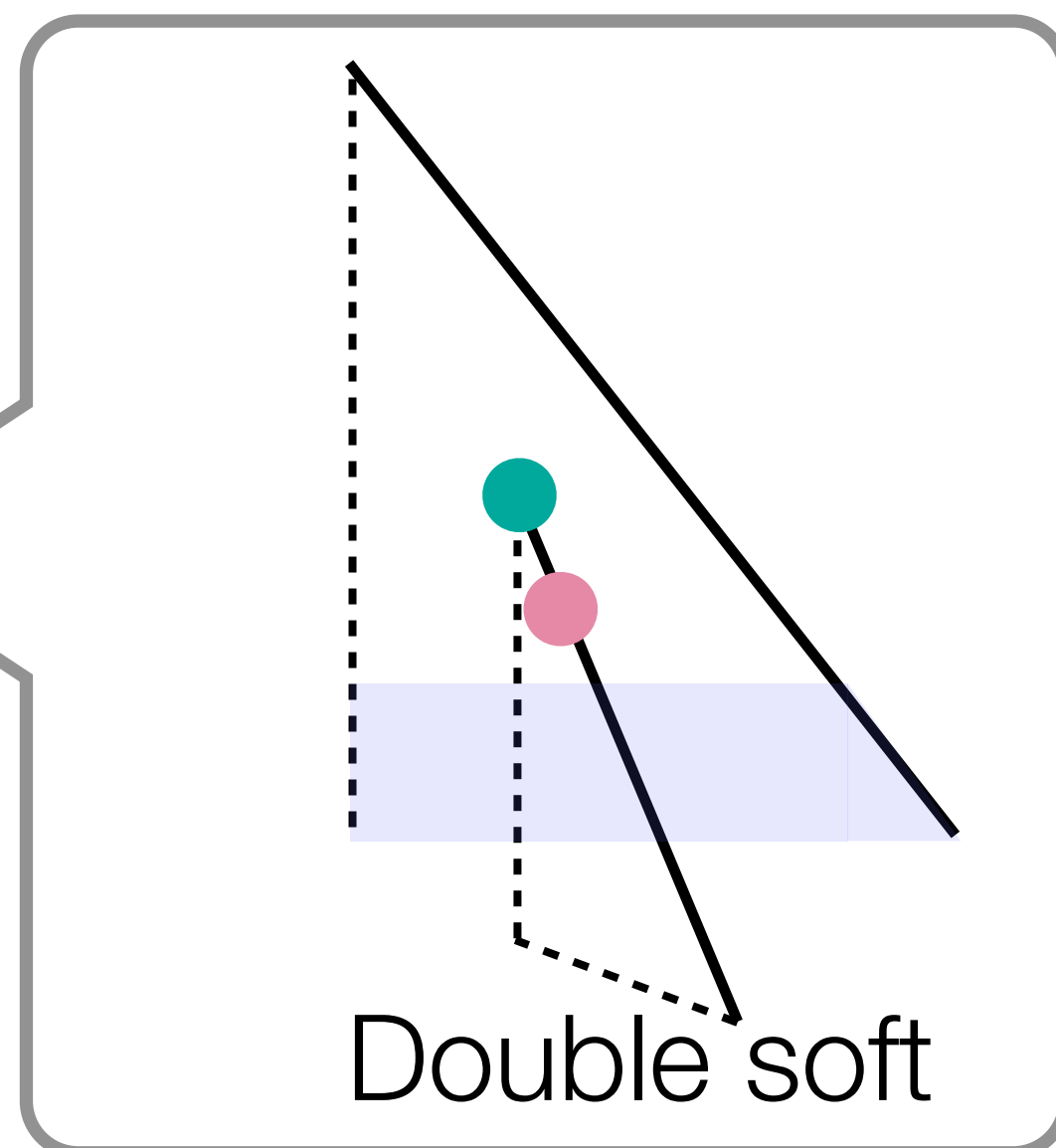


Collinear end-point

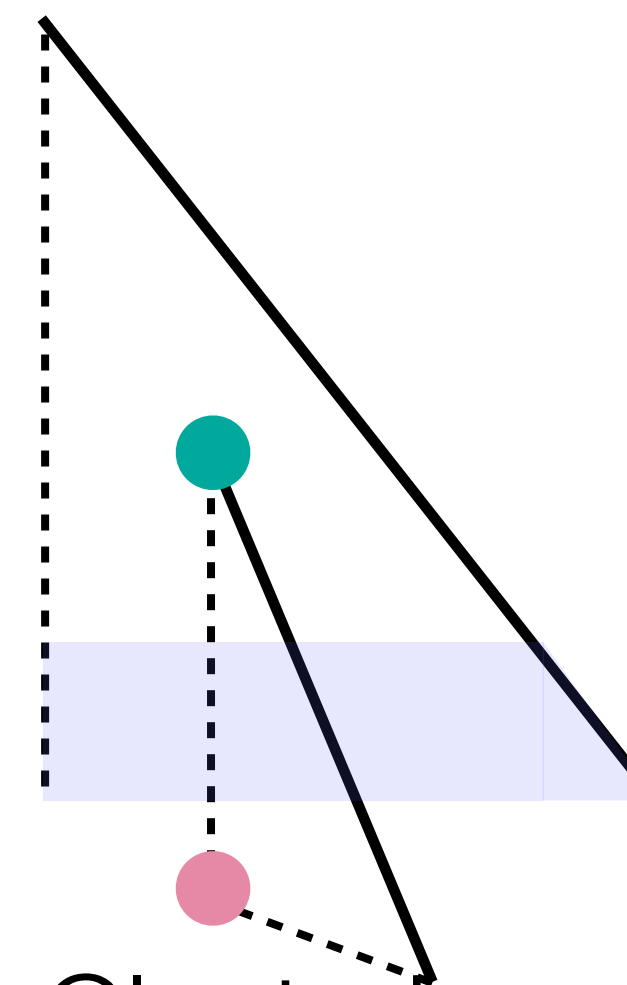


Large angle
+ $k_t \sim k_{t,\text{cut}}$

Crucial for
NNLL showers

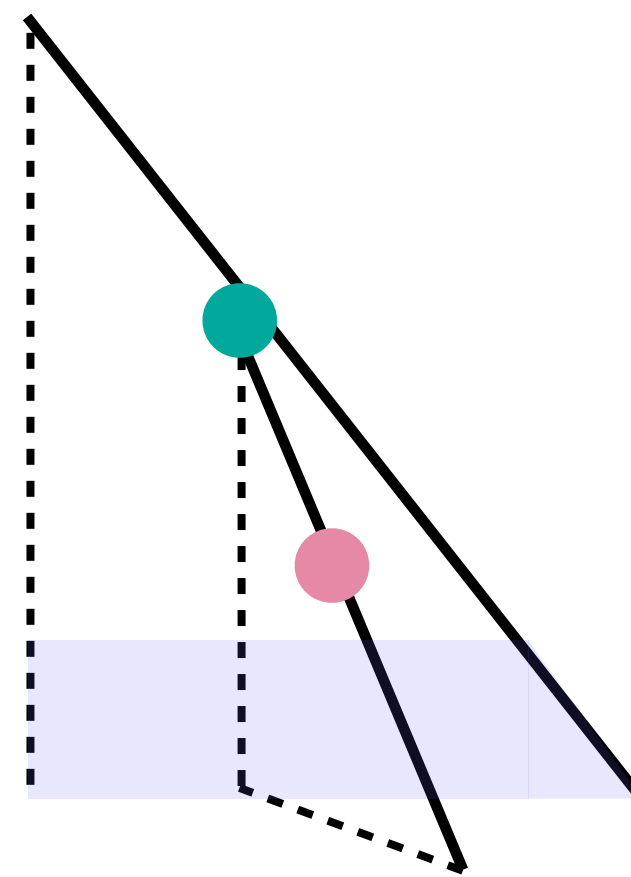


Double soft

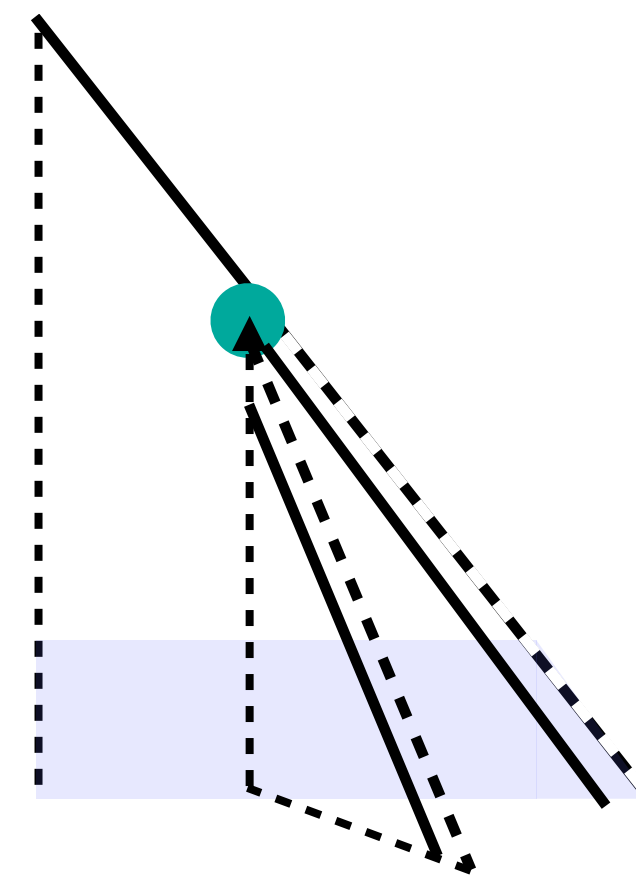


Clustering

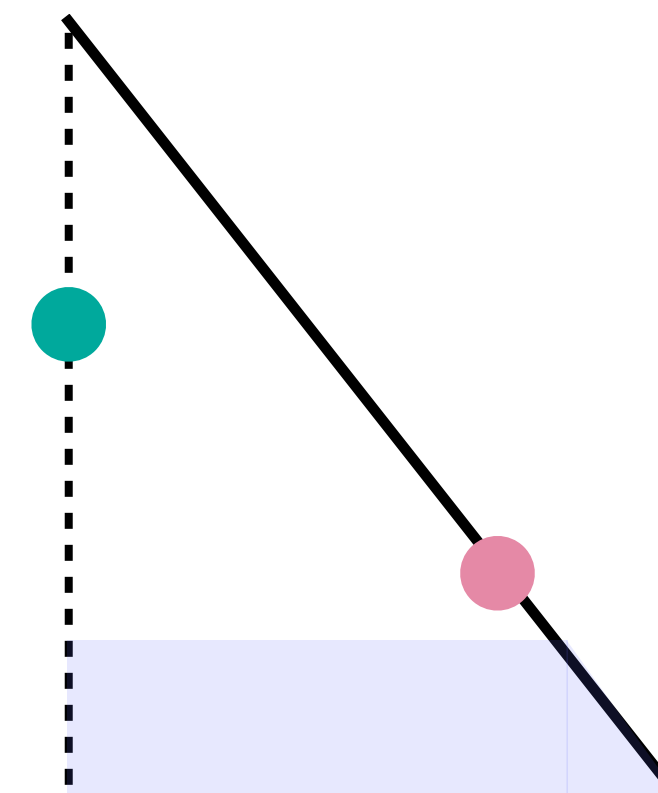
Lund multiplicity: NN DL resummation $(\alpha_s L)(\alpha_s L)(\alpha_s L^2)^n \sim (\text{NDL})^2$



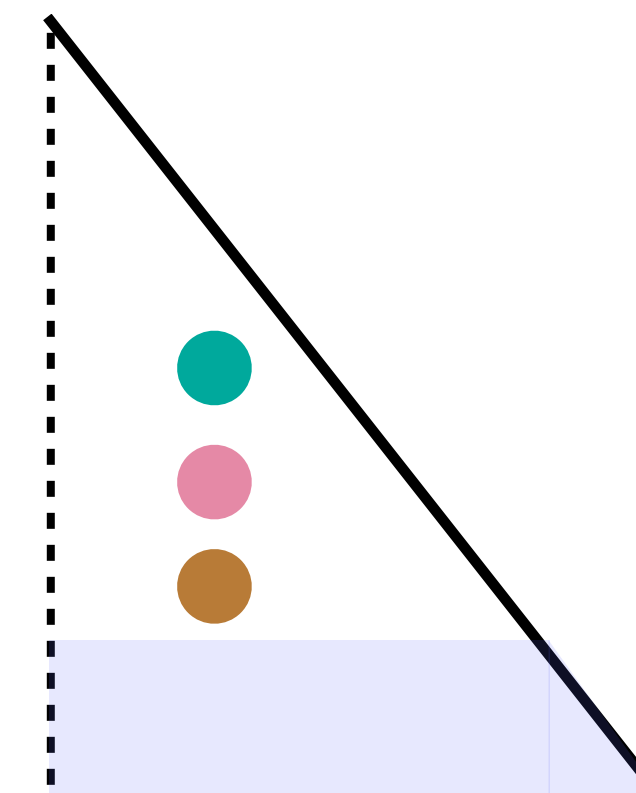
2 hard collinear



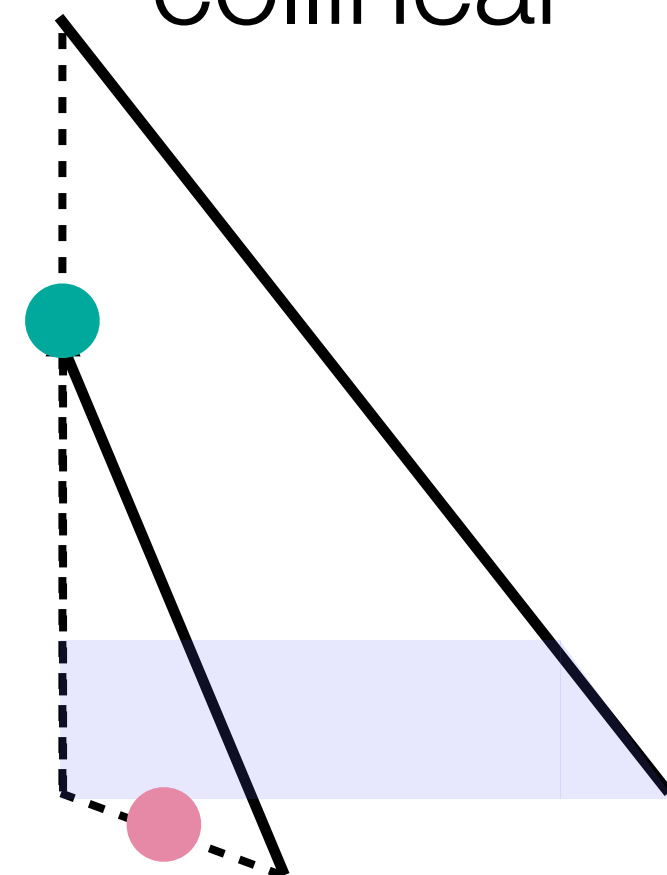
Energy loss



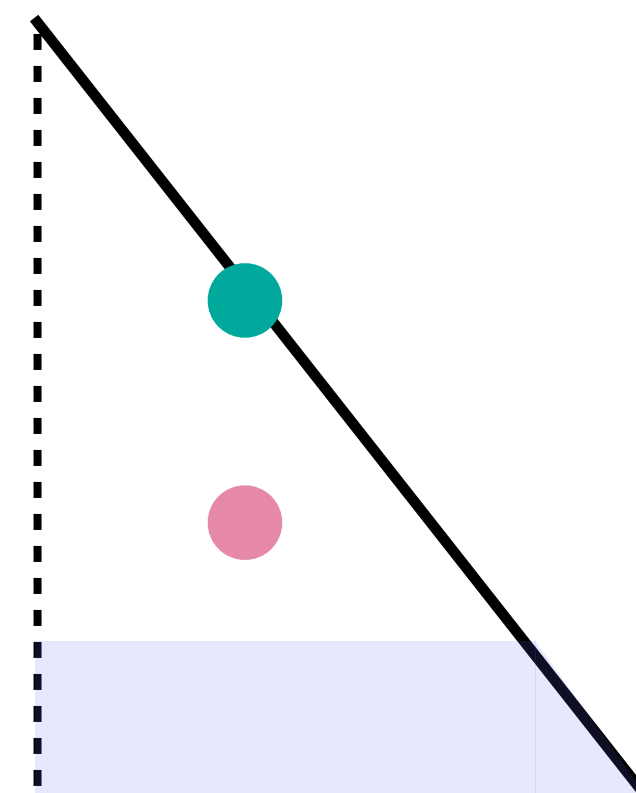
Large angle + hard-collinear



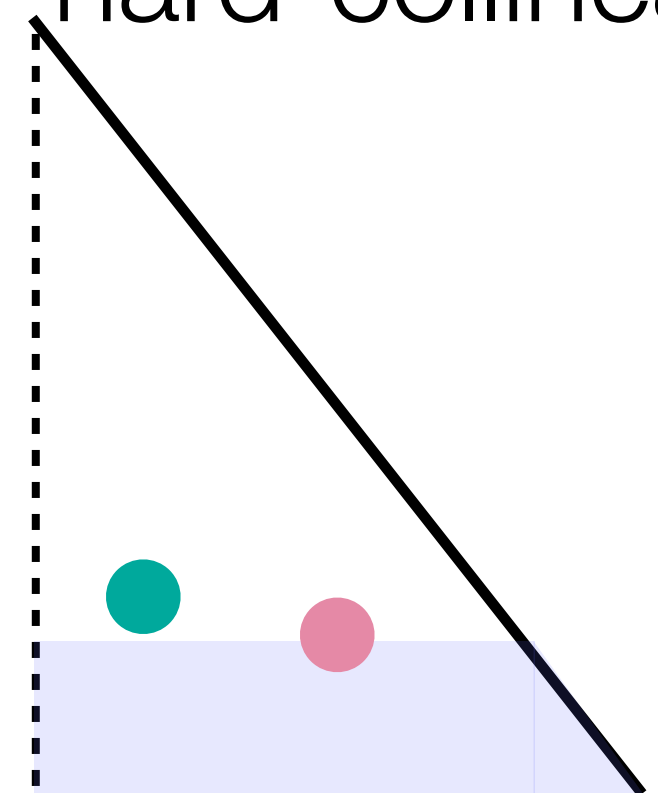
3 commensurate angles



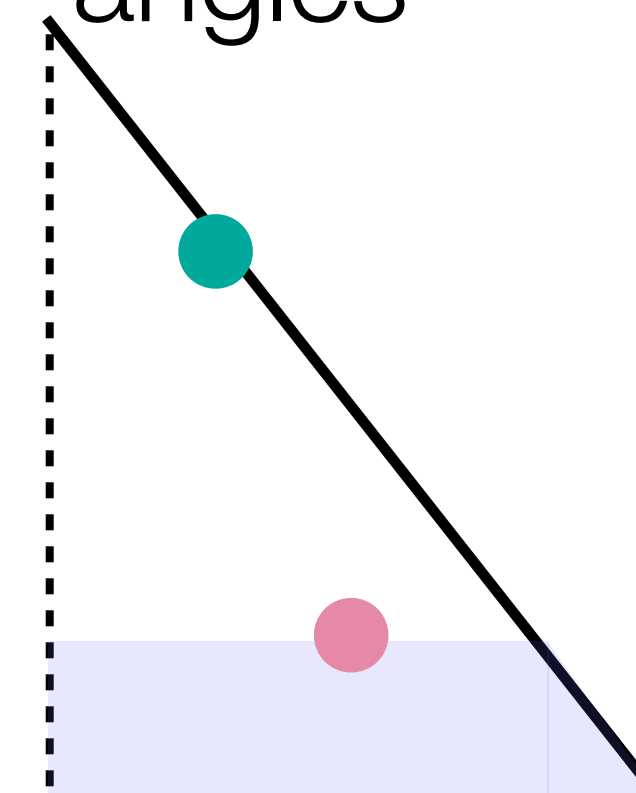
Large angle + $k_t \sim k_{t,\text{cut}}$



Hard coll. + commensurate η

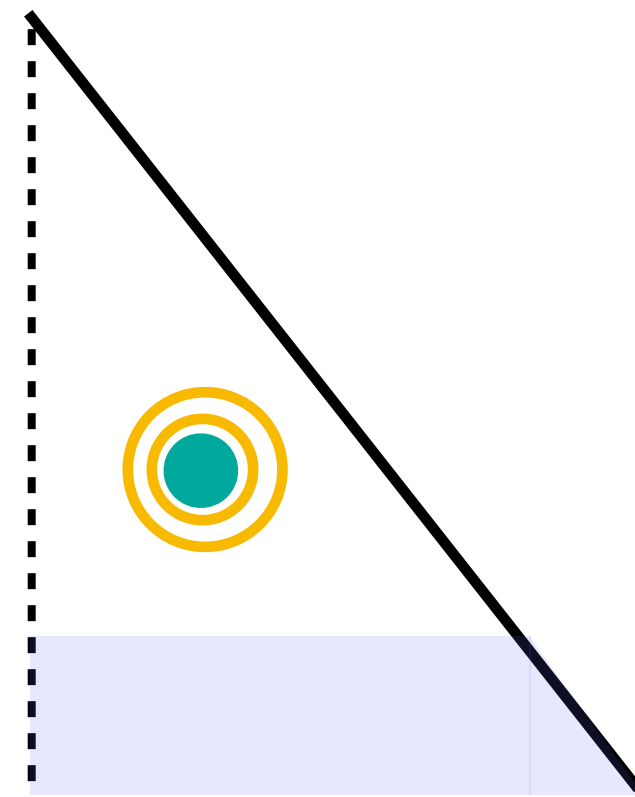


2 $k_t \sim k_{t,\text{cut}}$

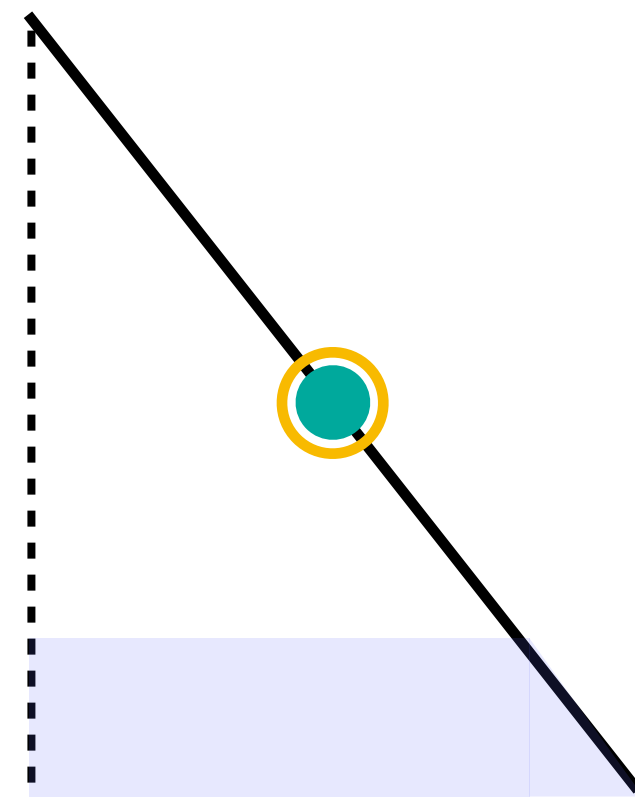


Hard coll. + $k_t \sim k_{t,\text{cut}}$

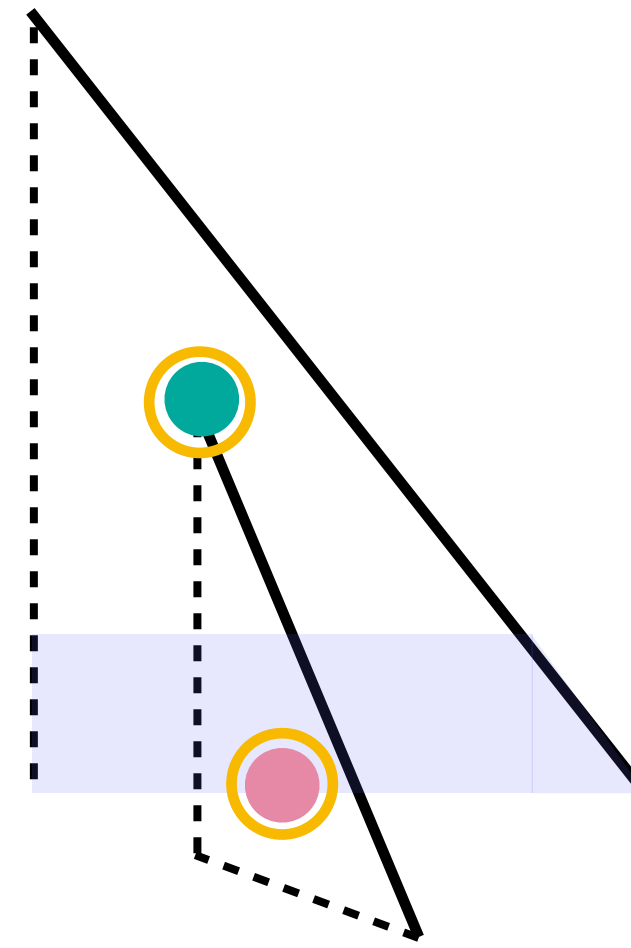
Lund multiplicity: NNDL resummation (running coupling)



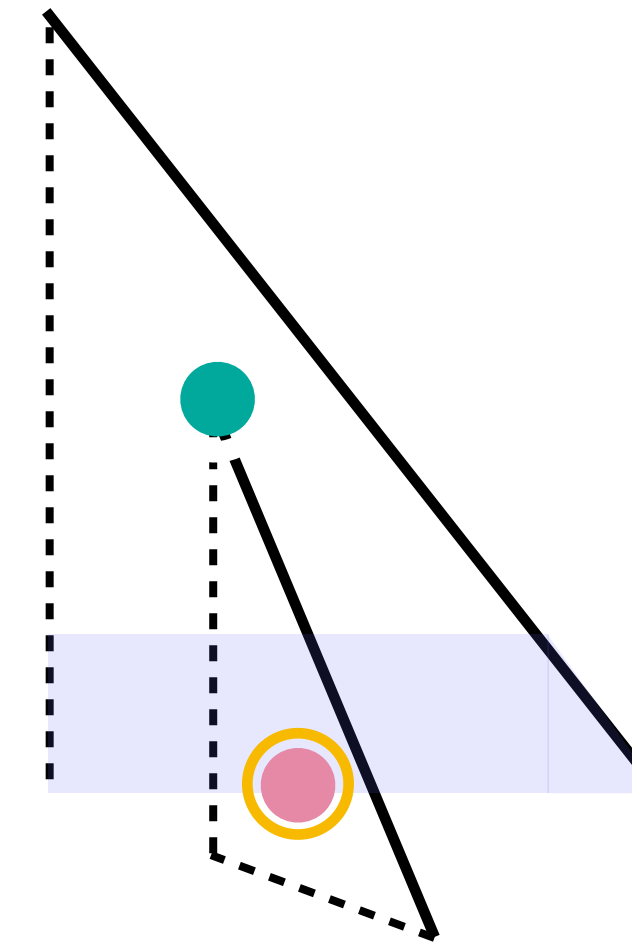
β_0^2



$\beta_0 + \text{hard-collinear}$



$\beta_0 \times \beta_0$



$\beta_0 \times \text{hard-collinear}$

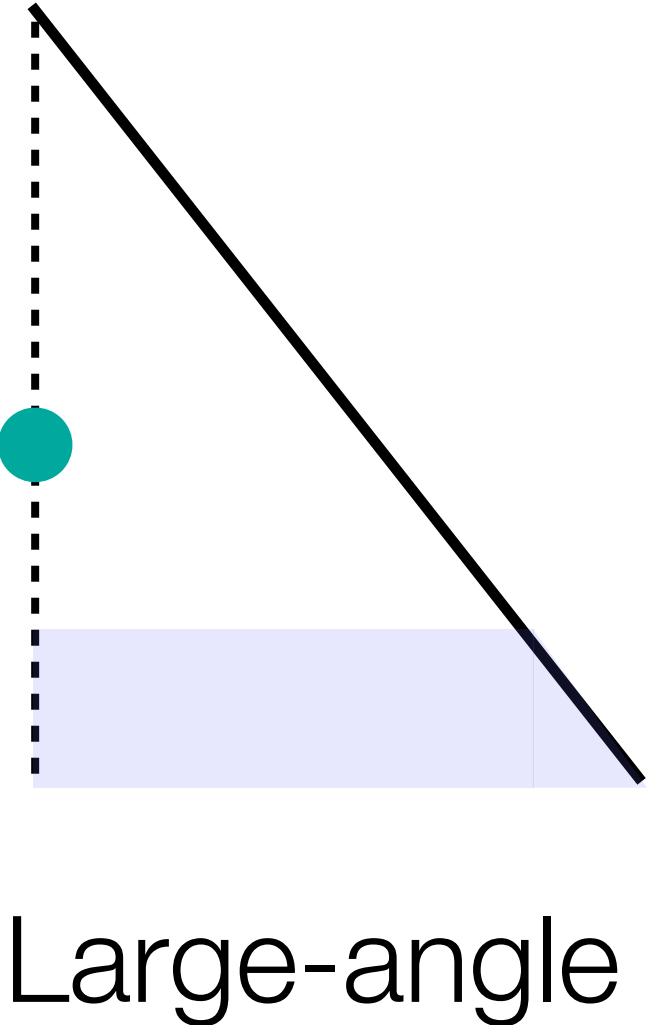
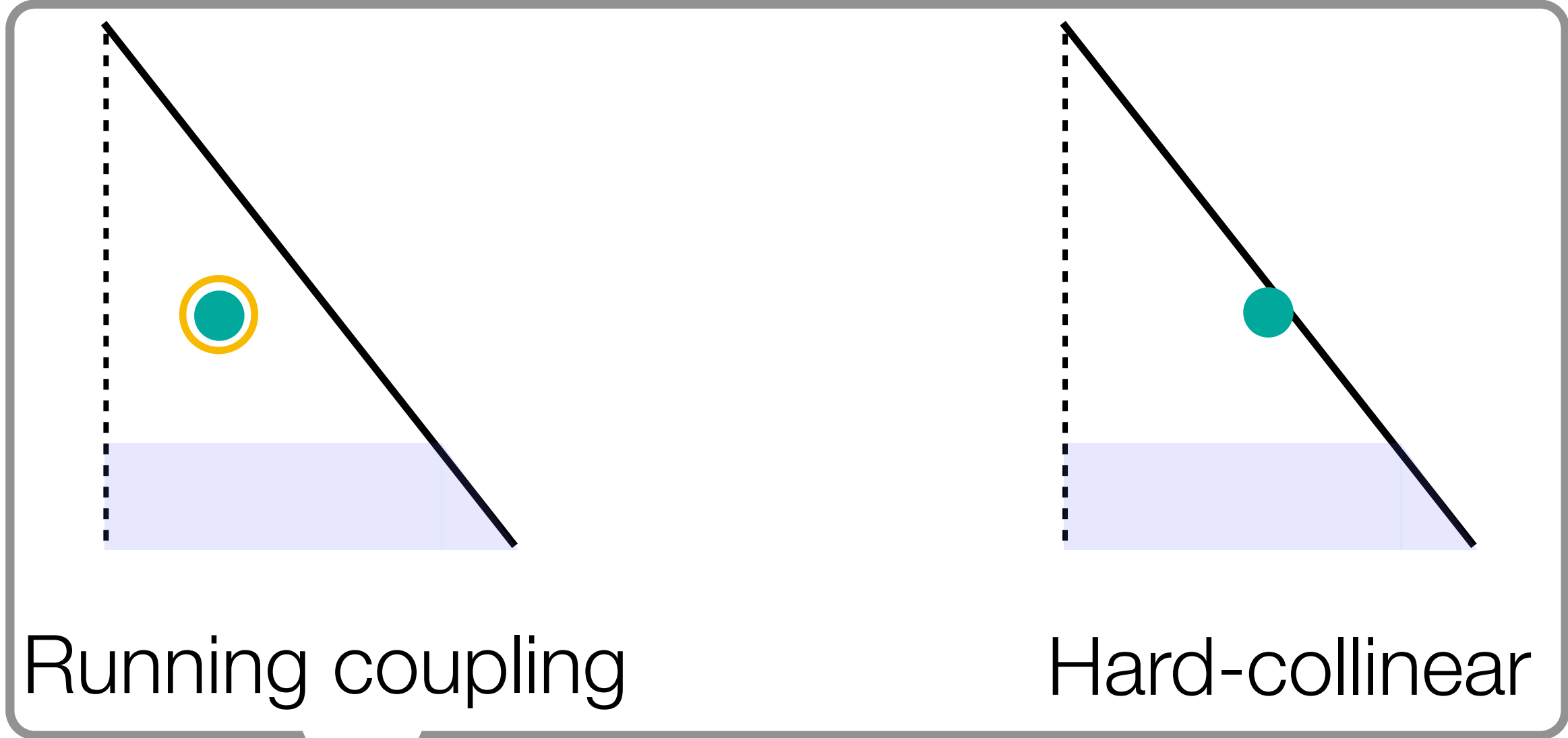
Lund multiplicity: NN DL result

$$\begin{aligned}
 \text{NN DL: } 2\pi h_3^{(q)} = & D_{\text{end}}^{q \rightarrow qg} + \left(D_{\text{end}}^{g \rightarrow gg} + D_{\text{end}}^{g \rightarrow q\bar{q}} \right) \frac{C_F}{C_A} (\cosh \nu - 1) + D_{\text{hme}}^{qqg} \cosh \nu \\
 & + \frac{C_F}{C_A} \left[(1 - c_\delta) D_{\text{pair}}^{q\bar{q}} (\cosh \nu - 1) + \left(K + D_{\text{pair}}^{gg} + c_\delta D_{\text{pair}}^{q\bar{q}} \right) \frac{\nu}{2} \sinh \nu \right] \\
 & + C_F \left[\left(\cosh \nu - 1 - \frac{1 - c_\delta}{4} \nu^2 \right) D_{\text{clust}}^{(\text{prim})} + (\cosh \nu - 1) D_{\text{clust}}^{(\text{sec})} \right] \\
 & + \frac{C_F}{C_A} \left[D_{\text{e-loss}}^g \frac{\nu}{2} \sinh \nu + (D_{\text{e-loss}}^q - D_{\text{e-loss}}^g) (\cosh \nu - 1) \right] \\
 & + \frac{C_F}{2} \left\{ (B_{gg} + c_\delta B_{gq})^2 \nu^2 \cosh \nu + 8 \left[2c_\delta B_{gg} - 2c_\delta B_q - (1 - 3c_\delta^2) B_{gq} \right] B_{gq} \cosh \nu \right. \\
 & \quad + [4B_q (B_{gg} + (2c_\delta + 1) B_{gq}) - (B_{gg} + c_\delta B_{gq}) (B_{gg} + 9c_\delta B_{gq})] \nu \sinh \nu \\
 & \quad \left. + 4(1 - c_\delta^2) B_{gq}^2 \nu^2 + 8 \left[2c_\delta B_q - 2c_\delta B_{gg} + (1 - 3c_\delta^2) B_{gq} \right] B_{gq} \right\} \\
 & + \frac{C_F}{C_A} \frac{\pi \beta_0}{2} \left\{ (B_{gg} + c_\delta B_{gq}) \nu^3 \sinh \nu + [2B_q - 2B_{gg} + (6 - 8c_\delta) B_{gq}] \nu \sinh \nu \right. \\
 & \quad \left. + 2(B_q + B_{gg} + B_{gq}) \nu^2 \cosh \nu - 4(1 - c_\delta) B_{gq} (2 \cosh \nu - 2 + \nu^2) \right\} \\
 & + \frac{C_F}{C_A} \frac{\pi^2 \beta_0^2}{8C_A} [3\nu(2\nu^2 - 1) \sinh \nu + (\nu^4 + 3\nu^2) \cosh \nu] \quad @ \text{ full colour}
 \end{aligned}$$

Extension to QCD jets: DL and NDL

DL: $h_{1,pp}(\xi) = f_q h_{1,e^+e^-}^{(q)}(\xi) + f_g h_{1,e^+e^-}^{(g)}(\xi)$

NDL:

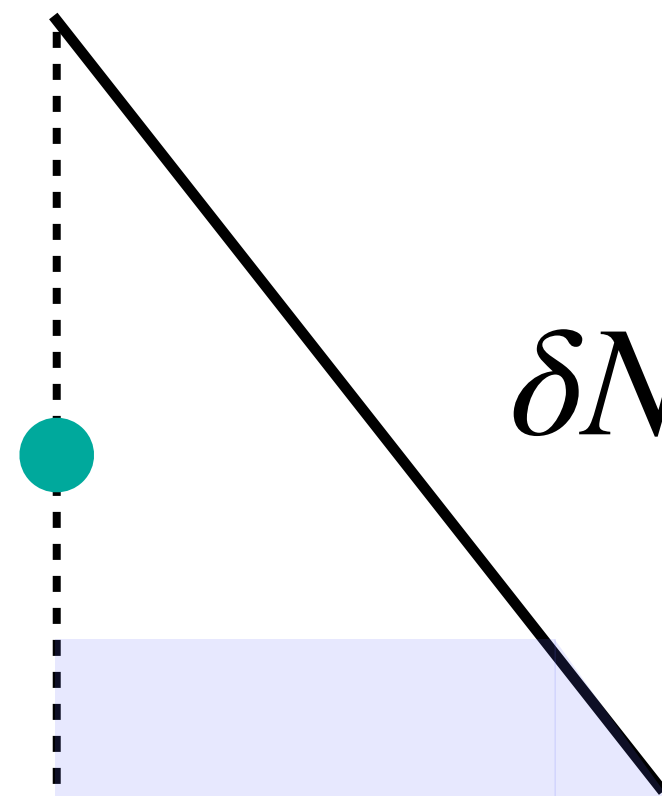


Collinear physics is universal: recompute only wide-angle diagrams

Extension to QCD jets: DL and NDL

DL: $h_{1,pp}(\xi) = f_q h_{1,e^+e^-}^{(q)}(\xi) + f_g h_{1,e^+e^-}^{(g)}(\xi)$

NDL:



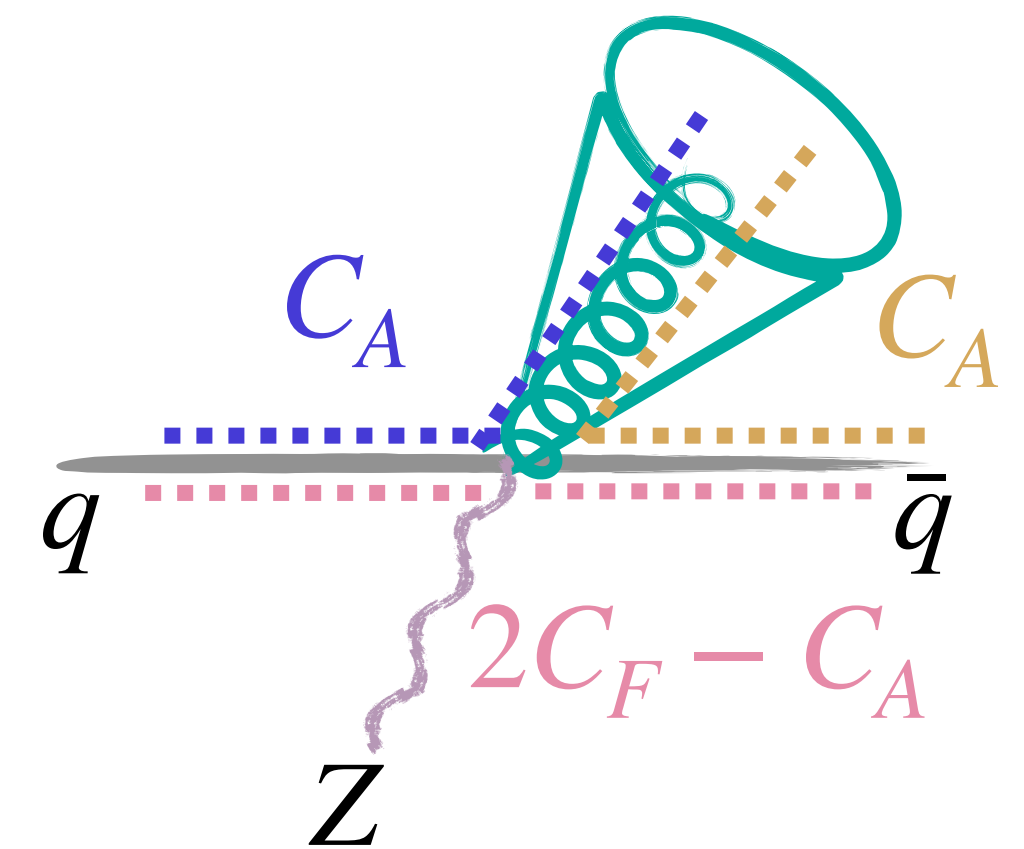
Large-angle

$$\delta N_{\mathcal{C},(ab),\mathcal{O}(\alpha_s)}^{(\text{NDL})} = \frac{\alpha_s}{\pi} L \omega_{ab}^{\mathcal{C}} D_{ab}^{\text{la}}$$

$$= \frac{\alpha_s}{\pi} L \omega_{ab}^{\mathcal{C}} D_{ab}^{\text{la}}$$

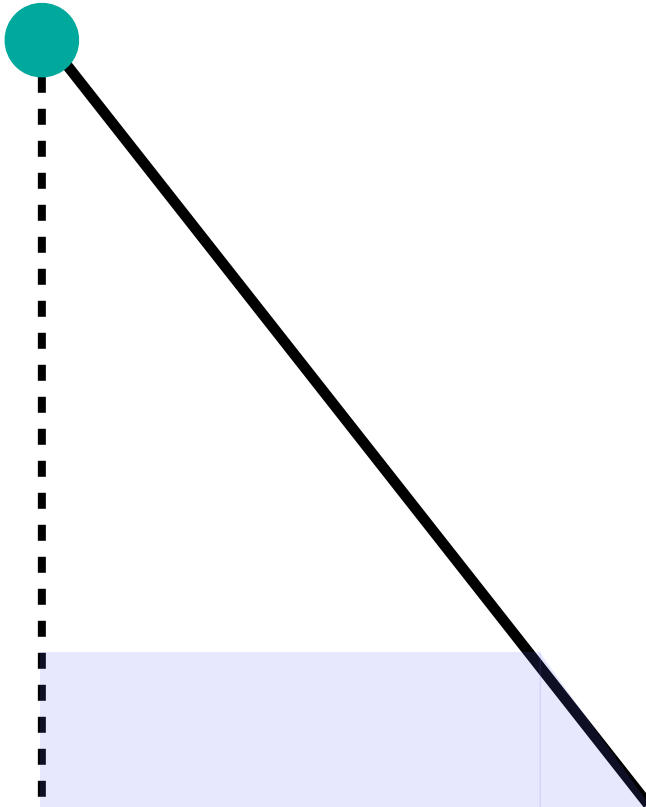
e.g. $\omega_{ab}^{\mathcal{C}} = (-2\mathbf{T}_a \cdot \mathbf{T}_b)$

$$\int [dk] \frac{P_a P_b}{(p_a p_k)(p_b p_k)} \Theta(\Delta_k < R)$$



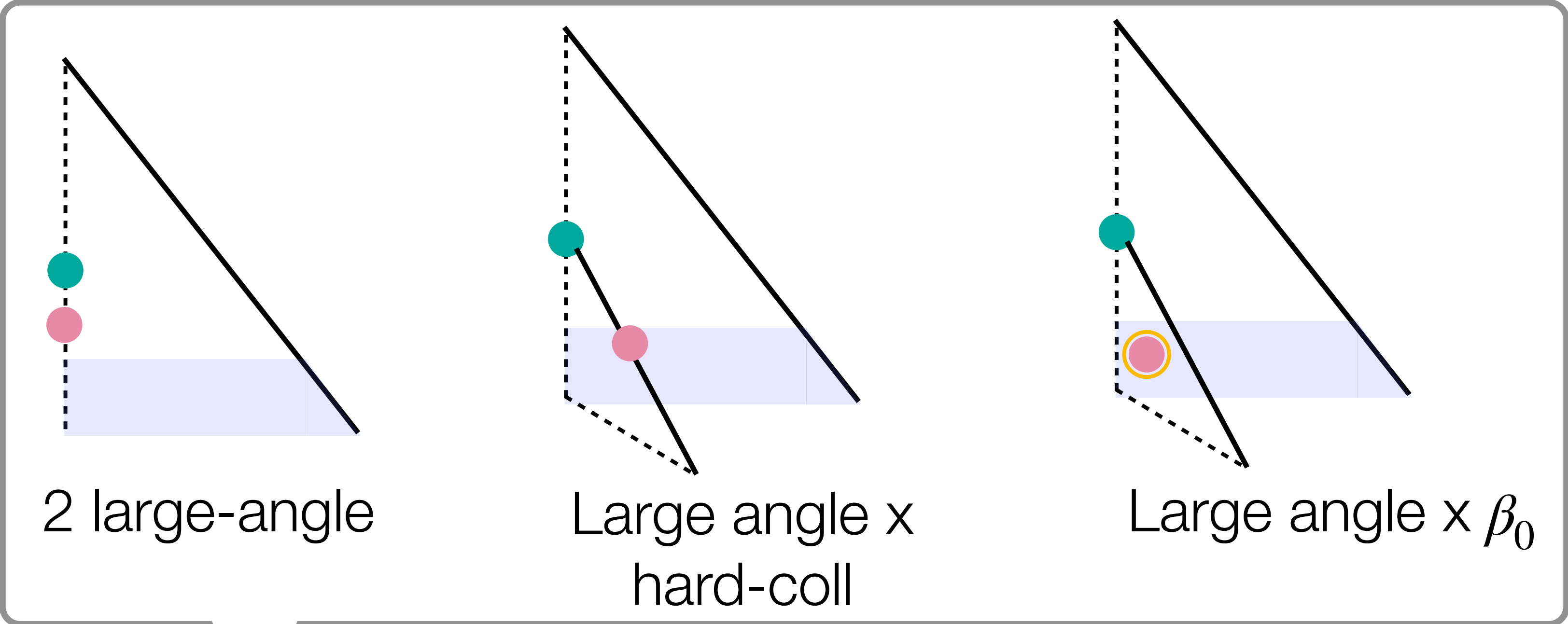
Extension to QCD jets: NNDL

NNDL:



Hard ME

$$\alpha_s (\alpha_s L^2)^n$$



2 large-angle

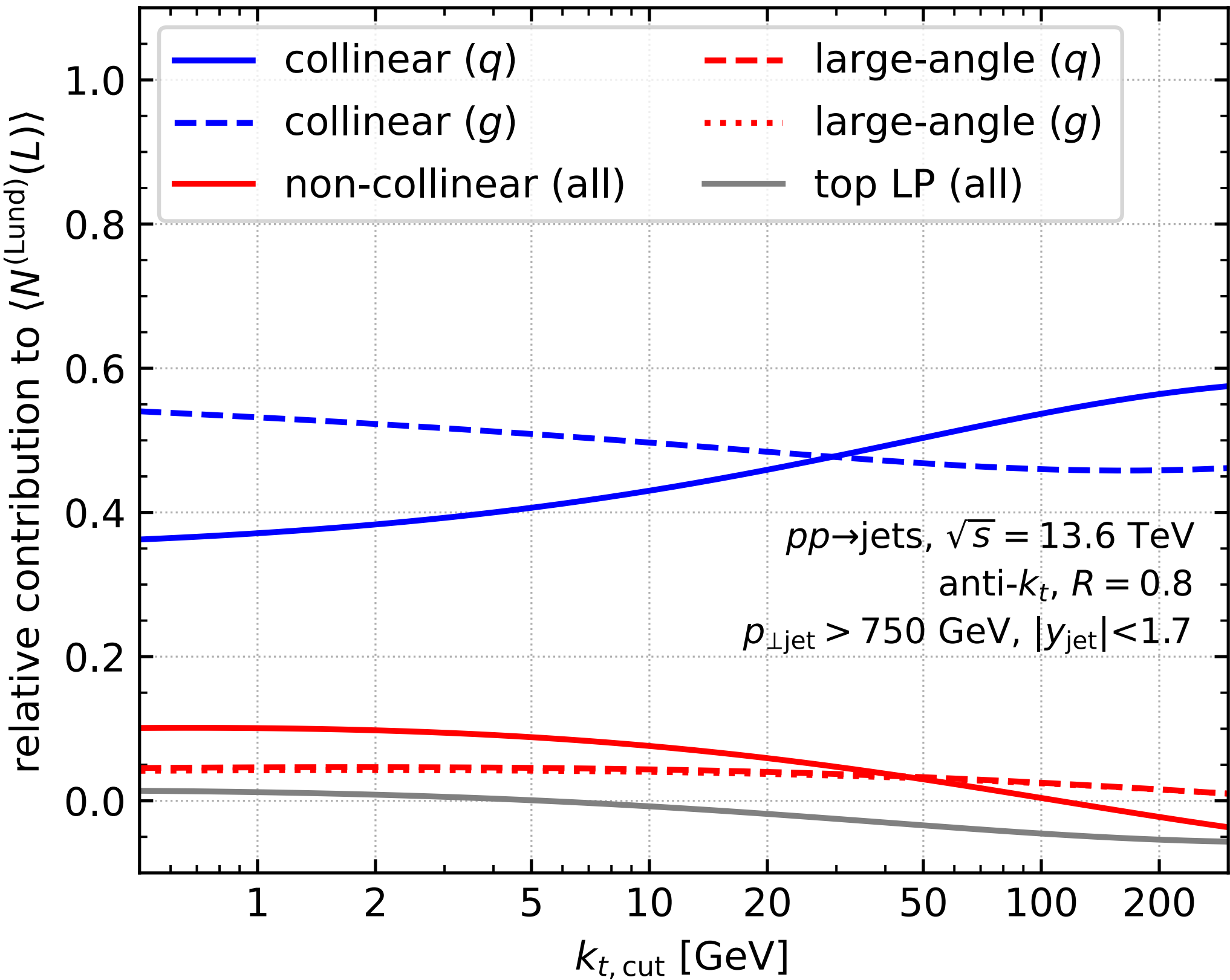
Large angle x
hard-coll

Large angle x β_0

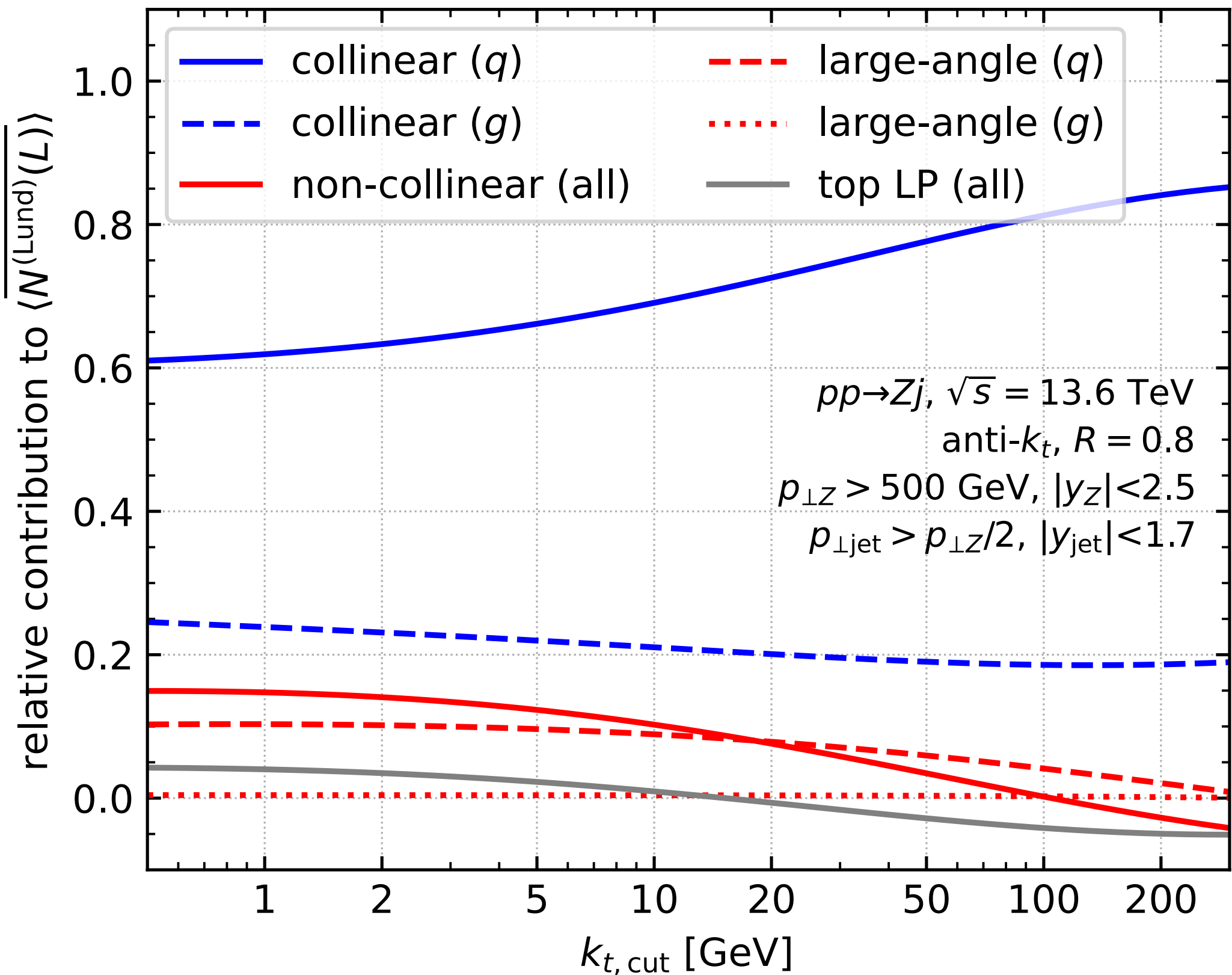
$$(\text{NDL})^2 = (\alpha_s L)(\alpha_s L)(\alpha_s L^2)^n$$

Universal vs large-angle contributions to multiplicity

inclusive jets

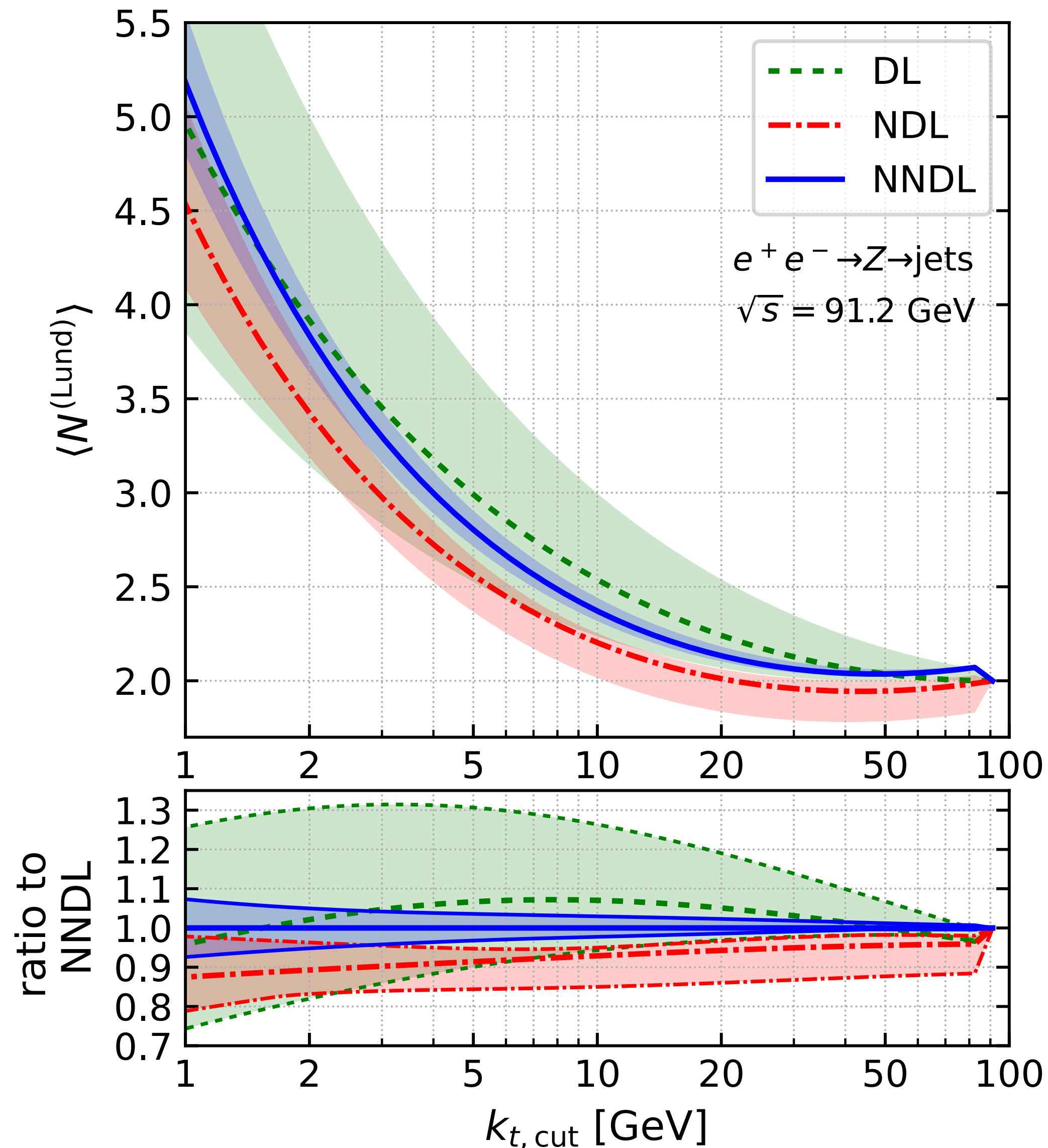


Z+jet events



Almost 90% of the Lund multiplicity has a collinear origin

The importance of higher logarithmic accuracy: e^+e^-

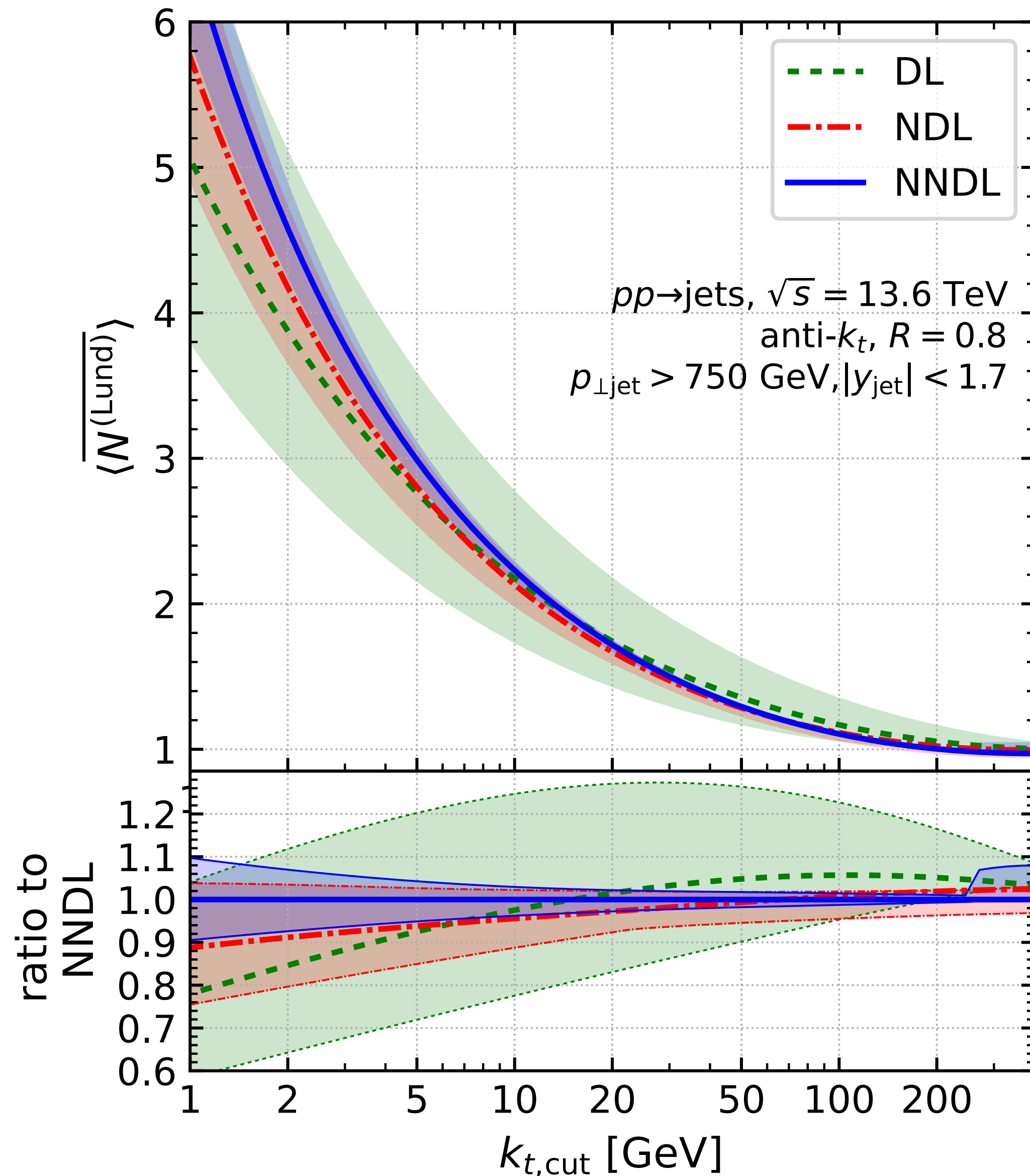


$$\text{Band} = \alpha_s(x_R Q), \ln(x_L Q / k_{t,\text{cut}})$$

The uncertainty of the theoretical prediction at $k_{t,\text{cut}} = 5 \text{ GeV}$ is

DL: 20%, NDL: 6%, NNDL: 3%

The importance of higher logarithmic accuracy: pp

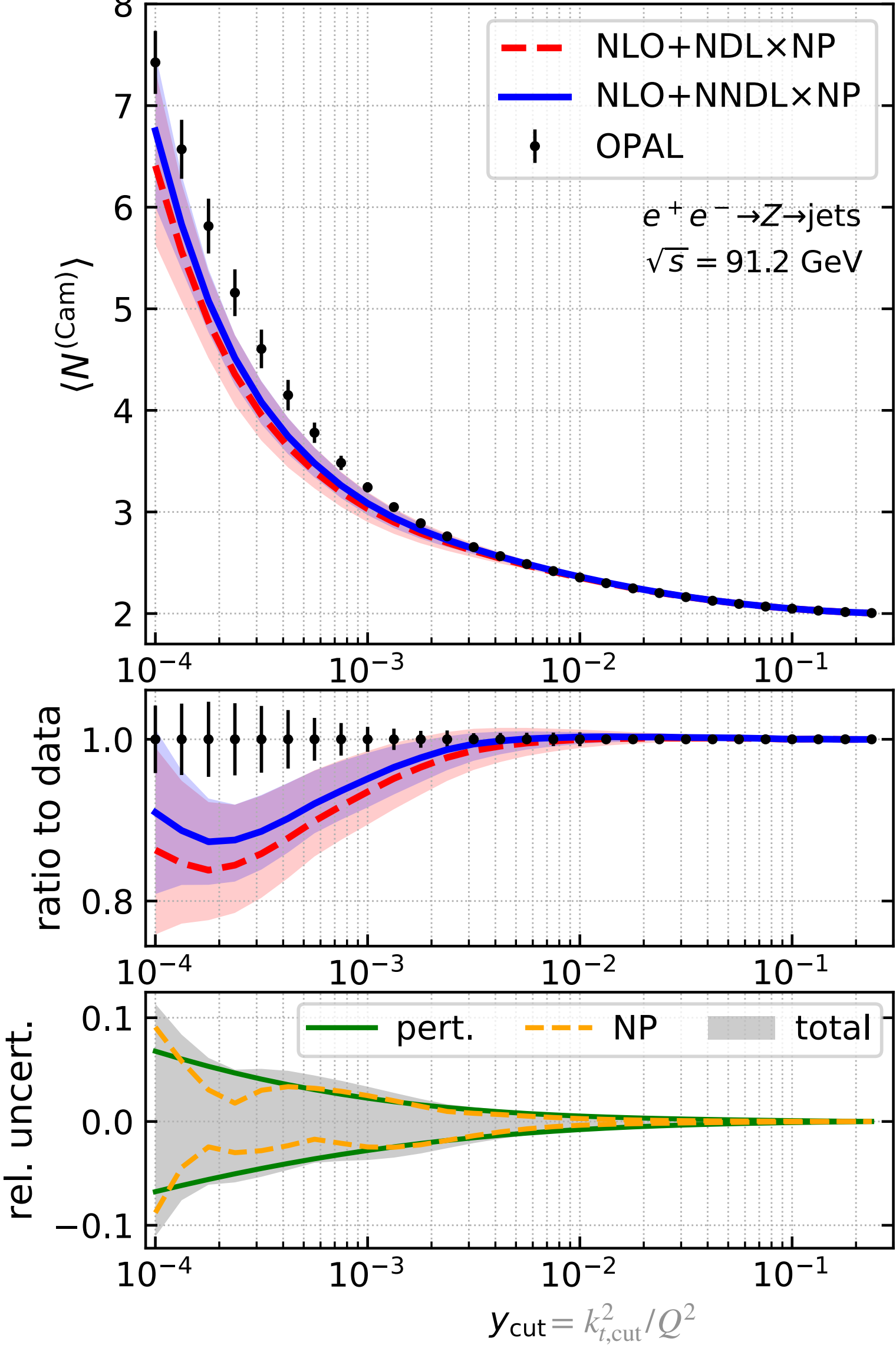


The uncertainty of the theoretical prediction at $k_{t,\text{cut}} = 5 \text{ GeV}$ is

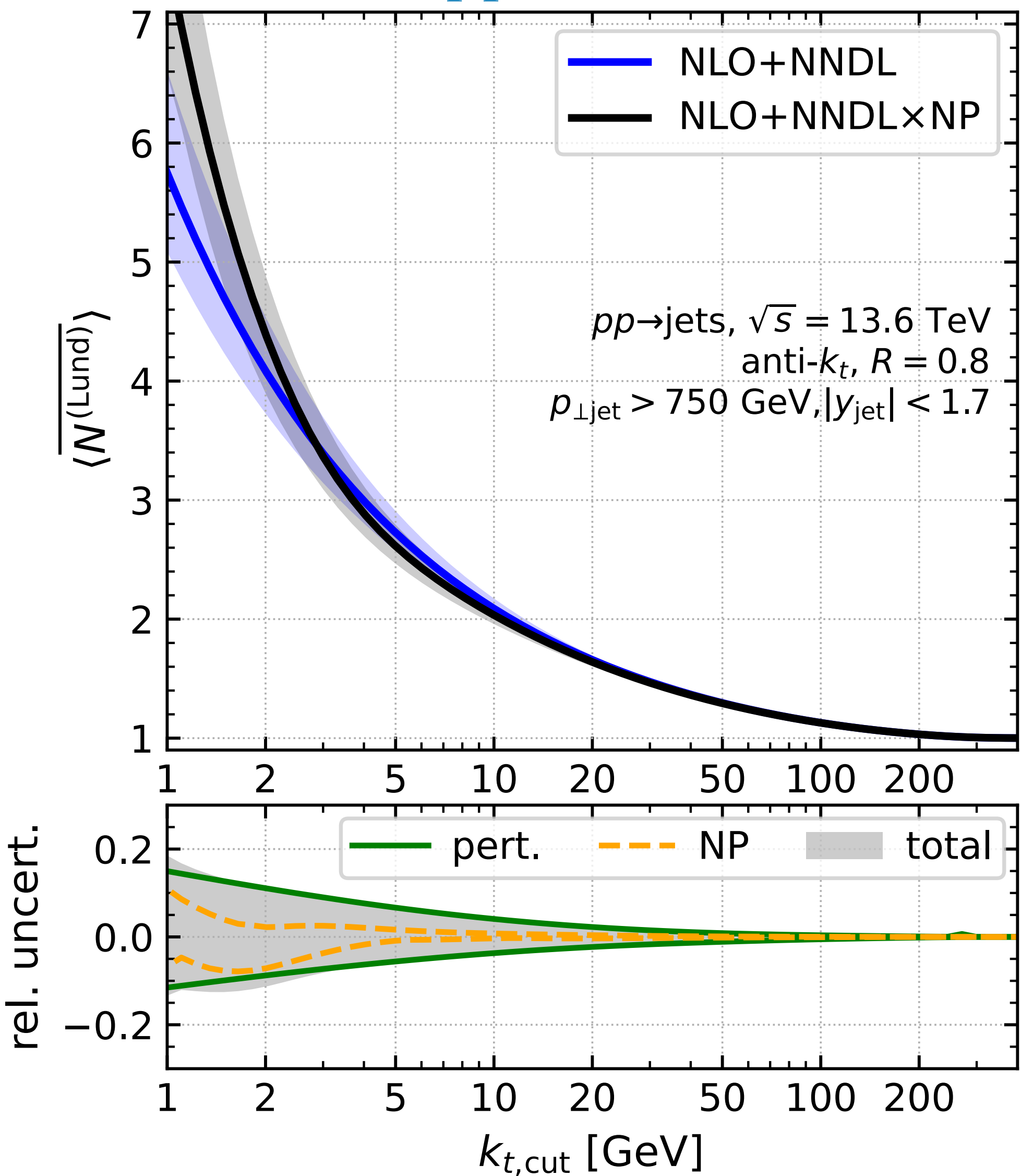
DL: 28 % , **NDL:** 10 % , **NNDL:** 5 %

Phenomenological studies

$e^+e^- @ \text{LEP}$



$pp @ \text{LHC}$



Wrap up

- Jet multiplicity has played a historic role in the development of pQCD
- No progress in the analytic prediction since 1992
- New, Lund-based definition is amenable to high-precision calculations
- Predictions status: NLO ($\mathcal{O}(\alpha_s^2)$) exact + MC non-perturbative corrections +
 - e⁺e⁻: NNDL massless quarks, NDL massive quarks
 - pp: NDL in DY production, NNDL for jets

Stay tuned for new experimental data and parton shower tests