

Analysis, simulations and measurements for the operation of accelerator magnets

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Career history



Magnetic field prediction for operation

Doctoral Student in TE-ABT-PPE (2016-2019)

Optimization of the beam screen for the FCC injection kicker magnets, Ph.D. at EPFL (2019)

- Contribution to the **upgrade** and **commissioning** of the LHC Injection Kicker Magnet (LHC MKI)
- Development of the novel concept of a helical beam screen
 - Improved high-voltage performance
 - Reduced beam-coupling impedance



Beam screen for the kicker magnets

Introduction



What do we mean by **injection**?

A particle beam must be injected into a circular accelerator or accumulator **at the right time** while:

- placing the newly injected particles onto the correct trajectory,
- minimizing beam loss.

Machine protection related issues:

high beam energy (up to 2.4 MJ for each LHC injection and twice more for FCC-hh).

Technological challenges:

- fast, high voltage pulse generators,
- special types of magnets (kickers and septa).

Injection kicker magnet (I)



- Kicker magnet is a pulsed dipole magnet with very fast rise and/or fall time (from tens to hundreds of nanoseconds)
- It "sees" the circulating beam!



• **Ni-Zn ferrite yoke** is always used for fast applications (e.g. 8C11, CMD5005)

Injection kicker magnet (II)

Transmission line ferrite-loaded kicker magnet is the best option to achieve fast field rise and fall times with low pulse ripple.



- Complicated design
- Field rise time (τ_r) depends on pulse propagation time through the magnet:

$$\tau_m = n \sqrt{L_{\text{cell}} C_{\text{cell}}} = n \frac{L_{\text{cell}}}{Z}$$

• Fast: field rise times $\ll 1 \ \mu s$



• Magnetic flux:

$$\Phi = \int (V_{\rm in} - V_{\rm out}) dt$$

• Field rise time:

$$\tau_r \sim \tau_m + \tau_c$$

 τ_c – rise time of the voltage pulse

Example: LHC injection kicker magnet (MKI)

| 1.º | The second se | Parameter | Value | Unit |
|-----------------------------|---|-------------------------------------|------------|------|
| | | Number of magnets per system | 4 | - |
| • | | Kicker strength per magnet | 0.3 | T·m |
| • | | Number of cells per magnet | 33 | - |
| | | Characteristic impedance | 5 | Ω |
| · | | Designed magnet voltage | 27 | kV |
| | | Aperture height/width | 54/54 | mm |
| | | Magnetic field flat-top duration | up to 7.86 | μs |
| | | Designed field rise time 0.5%-99.5% | 0.9 | μs |
| | | Designed field fall time 99.5%-0.5% | 3 | μs |
| | | Beam screen inside diameter | 38 | mm |
| | | Magnet length | 2.7 | m |
| | | | | |
| Counter rotating | | | | |
| Injected + circulating beam | | | | |
| beam Beam screen | | | | |

Major problem: beam-induced heating due to the real component of the longitudinal beam-coupling impedance \rightarrow a beam screen is placed in the aperture of each LHC MKI.

Beam coupling impedance

Reference system:

The wake field generated by the source particle affects the witness particle.



The net force acting on the witness charge:

$$\begin{split} \vec{F}(\vec{r}_1, s_1, \vec{r}, s, t) &= q \Big[\vec{E}(\vec{r}_1, s_1, \vec{r}, s, t) + \vec{v} \times \vec{B}(\vec{r}_1, s_1, \vec{r}, s, t) \Big] \\ \vec{F}(\vec{r}_1, s_1, \vec{r}, s, t) &= \vec{F}_{\parallel} + \vec{F}_{\perp} \end{split}$$

Longitudinal wake function:

$$W_{\parallel}(z) = -\frac{1}{q_1 q} \int_0^L \vec{F}_{\parallel} \left(\vec{r}_1, s_1, \vec{r}, s, t \right)_{|t = \frac{S-Z}{\beta c}, s_1 = s-z} ds.$$
$$W_{\parallel}(z) = -\frac{\Delta E(\vec{r}_1, \vec{r}, z)}{q_1 q}.$$

Longitudinal beam coupling impedance:

$$Z_{\parallel}(\omega) = \frac{1}{\beta c} \int_{-\infty}^{+\infty} W_{\parallel}(z) e^{j\omega \frac{z}{\beta c}} dz.$$

What happens to the energy lost by the source?

- Partially dissipates in the walls or HOM absorbers:
 → Heat load
- Partially transferred to the following particles
 → Beam instabilities

Beam-induced heating

Power loss (multipole pass):

$$P = -2(qN_pN_bf_0)^2 \left[\left(\sum_{p=1}^{+\infty} \left| \hat{\lambda}_{\text{beam}}(p\omega_0) \right|^2 \operatorname{Re}\left[Z_{\parallel}(p\omega_0) \right] \right) + \frac{1}{2} \left| \hat{\lambda}_{\text{beam}}(0) \right|^2 \operatorname{Re}\left[Z_{\parallel}(0) \right] \right]$$

 N_b – number of bunches, N_p – bunch intensity, f_0 – revolution frequency, $\hat{\lambda}_{\text{beam}}$ – beam spectrum.

Beam spectrum 100 Beam spectrum: 10400b 0.8 Bunch spectrum 80 $\hat{\lambda}_{beam}^{}$ [a.u.] $\text{Re}(Z_{||})$ [Ω] 60 40 0.2 20 0 500 750 1000 1250 1500 250 200 500 1000 1500 Freq. [MHz] Freq. [MHz]

25 ns bunch spacing \rightarrow 40 MHz beam harmonics





Shift the peaks towards higher frequencies!

Consequences:

For high-intensity beams, heating of the ferrite yoke can result in:

- \rightarrow local vacuum degradation
- \rightarrow loss of the magnetic properties (Curie point exceeded)



Beam screening: Conventional design (I)



More than 15 years of development...

General features:

- 24 (NiCr) straight screen conductors
 → reduced beam coupling impedance
- Capacitive coupling to the beam pipe at the upstream end and connection to the beam pipe at the downstream end
 - \rightarrow fast magnetic field rise time
- Overlap length determines resonances in the longitudinal impedance spectrum
 - → possibility to shift resonant modes to higher frequencies:

$$f_{\rm conv}^{(n)} = \frac{nc}{2\sqrt{\varepsilon_{r,\rm eff}} + (L_{\rm overlap} + \delta_{\rm fringe})}$$

Ferrite rings at both ends of the beam screen
 → power loss reduction in the ferrite yoke

Beam screening: Conventional design (II)

During magnetic field rise and fall, a significant voltage is induced on the screen conductors. → Possibility of HV breakdowns



Induced voltage:

$$V_i = \frac{d\Phi}{dt} = \frac{dB}{dt}(d + R - y_i)L$$

Magnet voltage and voltage of screen conductor:



Maximum conductor and inter-conductor voltages (for 60 kV PFN voltage):



- 1. Analysis of the LHC MKI beam screen (not discussed here)
- 2. Design of the conventional beam screen for FCC-hh (not discussed here)
- 3. Development of a novel helical beam screen for FCC-hh (next slides)

Design of a kicker magnet module for FCC-hh



| Baseline parameters: | | | | |
|------------------------------|-------|--|--|--|
| Injecton energy [TeV] | 3.3 | | | |
| Angle [mrad] | 0.18 | | | |
| Pulse duration [µs] | 2 | | | |
| Current [kA] | 2.4 | | | |
| Voltage [kV] | 15 | | | |
| System impedance [Ω] | 6.25 | | | |
| System length [m] | 40 | | | |
| Field rise time [µs] | 0.43 | | | |
| Aperture dimensions [mm] | 48/48 | | | |
| Magnet fill time [µs] | 0.355 | | | |
| Good field region (h/v) [mm] | 8/8 | | | |
| Flat-top tolerance [%] | ±0.5 | | | |

18 magnets per injection system



New PSpice model and CST model

Power deposition estimates



LHC MKI limit: 160 W/m.

Power loss is at unacceptable level.

Kicker magnet with the conventional beam screen:



Worst case scenario: 30.2 W/33.4 W/29.3 W

- → significant reduction in comparison with unshielded kicker (P~355 W/m),
- \rightarrow lower than for the upgraded LHC MKI8D.

New concept of a helical beam screen

Each screen conductor has a form of a helix:



Photograph of the alumina tube with 21 screen conductors applied using Ag paint



For an **integer number** of turns:

$$V_i = \frac{dB}{dt} (d+R)L$$

- \rightarrow the same induced voltage on each conductor
- → this is **one-half** of the worst-case voltage induced on a straight conductor.

Significantly improved HV performance

Beam coupling impedance simulations

- Beam coupling impedance modelling (CST Studio Suite) to optimize:
 - Beam-induced heating
 - Power distribution
 - Impact on beam stability



Longitudinal impedance optimization





No heating issues are expected for the FCC-hh kicker magnets for nominal beam parameters.

Characterization of EM properties of ferrites

- Novel measurements of EM properties of ferrites as a function of **frequency** and **temperature** (short-circuit line method):
 - \rightarrow new input for beam impedance simulations.





Chmielińska, A. *et al.* "Measurements of electromagnetic properties of ferrites as a function of frequency and temperature", In *Journal of Physics: Conference Series*, vol. 1067. No. 8, p. 082018. IOP Publishing.

Field response of the kicker magnet simulations

- Field response modelling (coupling Opera-2D with Cadence PSpice) to achieve required:
 - Field rise time
 - Field flat-top quality



Impact of a beam screen upon kicker waveform (PSpice)



Proposed beam screen design satisfies the field rise time and the field flat-top quality requirements

Experimental benchmark

• **Coaxial-wire measurements** performed on the LHC MKI to validate numerical models (classical and resonant method)



Chmielińska, A. and Barnes, M. J. "Helical beam screen for the Future Circular Collider for hadrons injection kicker magnets." *Physical Review Accelerators and Beams* 23.4 (2020): 041002.

Measurements of beam coupling impedance

• **Coaxial-wire measurements** performed on the LHC MKI to validate numerical models (classical and resonant method)

Complete picture of the optimization of the beam screen design.

Consistent results in theoretical, numerical and experimental studies.

Chmielińska, A. and Barnes, M. J. "Helical beam screen for the Future Circular Collider for hadrons injection kicker magnets." *Physical Review Accelerators and Beams* 23.4 (2020): 041002.

Fellow in TE-MSC-TM (2020-2023)

Integration of numerical models with magnetic measurement data to model and characterize the magnetic field of superconducting magnets

- Design of experiments (test plan, specification of cycles)
- Data-driven updates of numerical models
- Magnetic-field prediction for machine operation (FiDeL)
 - Semi-empirical model decomposition
- Software design and implementation for data-analysis



Operational cycle

simulations in ROXIE

Cross-talk



Rotating-coil

PyMagAnalysis



https://gitlab.cern.ch/te-msc-mm/pymaganalysis

Analysis of simulations and measurements to predict the magnetic field for the operation of superconducting magnets

FiDeL model for machine operation (I)

• The **Field Description of the LHC (FiDeL)** is the magnetic field prediction system that provides a model of the transfer function of magnetic circuits installed in the machine in order to retrieve the current level given the required field:

 \rightarrow indispensable for LHC operation

- \rightarrow uses parametric equations to fit **magnetic measurements**
- → based on **physical decomposition** of effects that contribute to the total field in the aperture of the magnet
- → uses **several levels of approximations**, starting from a linear dependance on the current and adding correction terms (depending on the magnet type and function).

FiDeL model for machine operation (II)

1. Geometric: B

$$B_m^{\text{geo}} = \gamma_m I.$$

2. Magnetization:

$$B_m^{\text{mag}} = \mu_m |I| \left(\frac{I_{\text{inj}}}{|I|}\right)^{2-p_m} \left(\frac{I_c - |I|}{I_c - I_{\text{inj}}}\right)^{q_m} \left(\frac{T_{c0}^{1.7} - T^{1.7}}{T_{c0}^{1.7} - T_{\text{meas}}^{1.7}}\right)^{h_m}.$$

3. Saturation:

$$B_m^{\text{sat}} = \sum_{i=1}^N \sigma_m^i I\Sigma(I, S_m^i, I_{0m}^i, I_{\text{nom}}),$$

where
$$\Sigma(I, S, I_0, I_{nom}) = -\frac{1}{2} \left[1 + \operatorname{erf} \left(S\left(\frac{|I| - I_0}{I_{nom}} \right) \right) \right]$$
, $\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$.

The fitting parameters are implemented in the LHC control system.

PyMagAnalysis



https://gitlab.cern.ch/te-msc-mm/pymaganalysis

Analysis of nonlinear effects during luminosity calibration (I)

• MCBC/MCBY magnets are used to provide beam separation during van der Meer (vdM) scans.



 Nonlinearities and lack of repeatability during vdM scans triggered the necessity to review the hysteretic behaviour of MCBC/MCBY magnets.

Measurements in SM18



Extensive program of **standard** and **operational test cycles launched** to characterize the magnetic field and its **reproducibility from scan to scan** using the **rotating-coil technique**.

Analysis of measurements and simulations (I)



Developed numerical model to account for transitions between major magnetization branches

Benchmark of the proposed model

Comparison of the measured magnetization-branch transitions (markers) and the proposed extended FiDeL model for the MCBC magnet (dashed lines).

Predicted and measured impact of magnetic hysteresis on the transverse displacement



Chmielińska, A., et al. "Magnetization in superconducting corrector magnets and impact on luminosity–calibration scans in the Large Hadron Collider." *Eur. Phys. J. Plus* (2023) 138: 796.

Magnetic model for the MCBXF (I)



(main field B1, normal) Outer coil provides horizontal field (main field A1, skew).

Identifying the problem



The development of FiDeL model is particularly challenging due to angular dependence of saturation (and persistent current magnetization).

Proposed model (I)

Original definition:

 $B_m^{\text{sat}}(I_1) = \sum_{i=1}^N \sigma_m^i I_1 \Sigma (I_1, S_m^i, I_{0m}^i, I_{\text{nom}}), \text{ where: } \Sigma (I, S, I_0, I_{\text{nom}}) = -\frac{1}{2} \left[1 + \text{erf}\left(S\left(\frac{|I| - I_0}{I_{\text{nom}}}\right)\right) \right],$ m - multipole, $\sigma_m, S_m, I_{0m}, I_{\text{nom}}$ - fitting parameters.

Proposal:

$$B_m^{\text{sat}}(I_1) \mapsto B_m^{\text{sat}}(I_1, I_2)$$

$$A_m^{\text{sat}}(I_2) \mapsto A_m^{\text{sat}}(I_1, I_2)$$

$$I_1 - \text{inner current}$$

$$I_2 - \text{outer current}$$

 $B_{m}^{\text{sat}}(I_{1}, I_{2}) = \sum_{i=1}^{N} \sigma_{m}^{i} I_{1} G_{m}^{B}(I_{1}, I_{2}) \Sigma (I_{1} G_{m}^{B}(I_{1}, I_{2}), S_{m}^{i}, I_{0m}^{i}, I_{\text{nom}_1}),$ $A_{m}^{\text{sat}}(I_{1}, I_{2}) = \sum_{i=1}^{N} \sigma_{m}^{i} I_{2} G_{m}^{A}(I_{1}, I_{2}) \Sigma (I_{2} G_{m}^{A}(I_{1}, I_{2}), S_{m}^{i}, I_{0m}^{i}, I_{\text{nom}_2}),$

where $G_m(I_1, I_2)$ is a function of inner and outer current (determined experimentally).

Proposed model (II)

Input: *I*₁, *I*₂

$$G_{m}^{B}(I_{1},I_{2}) \qquad G_{m}^{B}(I_{1},I_{2})$$

$$B_{m}^{\text{sat}}(I_{1},I_{2}) = \sum_{i=1}^{N} \sigma_{m}^{i} I_{1} \left(\frac{a}{\alpha(I_{1},I_{2}) + \Delta} + b \right) \Sigma \left(I_{1} \left(\frac{a}{\alpha(I_{1},I_{2}) + \Delta} + b \right), S_{m}^{i}, I_{0m}^{i}, I_{\text{nom1}} \right),$$

$$G_{m}^{A}(I_{1},I_{2}) = \sum_{i=1}^{N} \sigma_{m}^{i} I_{2} \left(\frac{a}{\frac{\pi}{2} - \alpha(I_{1},I_{2}) + \Delta} + b \right) \Sigma \left(I_{2} \left(\frac{a}{\frac{\pi}{2} - \alpha(I_{1},I_{2}) + \Delta} + b \right), S_{m}^{i}, I_{0m}^{i}, I_{nom2} \right)$$

where: $\alpha(I_1, I_2) = \tan^{-1} \left(\frac{|I_1|/I_{\text{nom}1}}{|I_2|/I_{\text{nom}2}} \right)$

Modelling error

New model with angular dependence of saturation



The error is within required specification (±30 units for B1/A1)



Modelling error

New model with angular dependence of saturation

MCBXF saturation model (B1)

 $\alpha = 1.57 \text{ rad}, G(I_1, I_2) = 0.56$

With **PyMagAnalysis** framework we brought **the best available knowledge** from **simulations** and **measurements** to **understand the magnet behavior itself** and to provide **precise field predictions for operation**.



Thank you for your attention!

