

De Sitter Dialectics

Susha Parameswaran

University of Liverpool

String Pheno 2023, Daejeon

based on

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with Bruno Bento, Dibya Chakraborty and Ivonne Zavala
and work to appear with Joaquim Gomes and Ed Hardy

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Anti-thesis = metastable dS vacua are possible in QG?
Synthesis = transient dS vacua are possible in QG?

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Plan

- ▶ Some recent questions around KKLT/LVS constructions
- ▶ Weakly-warped LVS de Sitter solution
- ▶ Robustness against g_s and \mathcal{V} corrections
- ▶ Transient de Sitter from interacting Dark Sectors

Uplifting without runaways

KKLT '03; see also Bena, Dudas, Grana & Lüst '18; Lüst & Randall '22

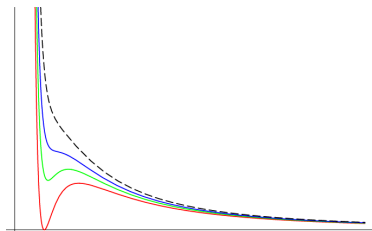


Figure from Angelova, Calo, Cicoli '09

- In KKLT/LVS, stabilisation via balance of classical, perturbative and non-perturbative effects and uplifting by e.g. $\overline{D3}$ -brane.

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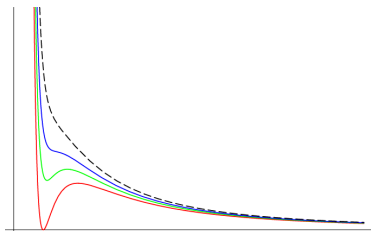


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- ▶ Uplifting must not destabilise volume modulus – place $\overline{D3}$ -brane at tip of strongly warped throat – large flux numbers
 $MK \sim \mathcal{O}(100) - \mathcal{O}(1000)$ – large contributions to tadpole, which must be cancelled.

Bena, Blaback, Graña & Lüst '20; Crinó, Quevedo & Valandro '20

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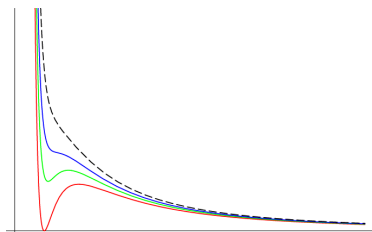


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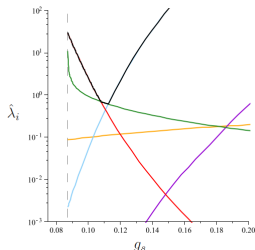
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- ▶ Alternative may be to have weak warping and large AW_0 in $W = W_0 + Ae^{-aT}$.

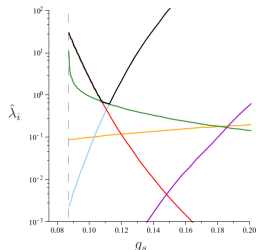
String-loop and curvature corrections

Junghans '22



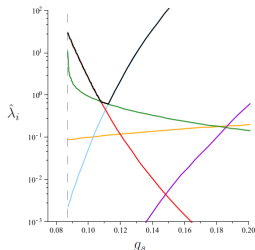
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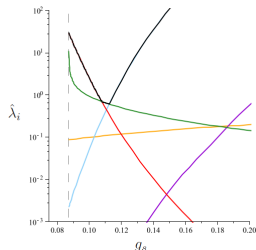
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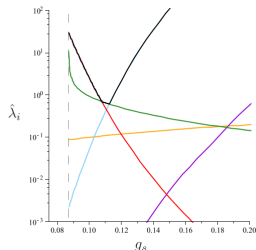
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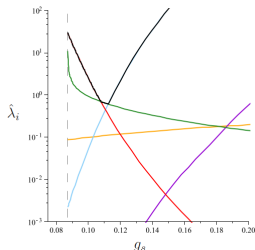
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- ▶ Weak warping and low MK may help to suppress α' corrections. ⁵

Weakly warped deformed conifold

Bento, Chakraborty, SLP, Zavala '23

- ▶ Type IIB flux compactification with warped deformed conifold region glued to compact CY:

$$ds^2 = h^{-1/2} g_{\mu\nu} dx^\mu dx^\nu + h^{1/2} c(x)^{1/2} g_{mn} dy^m dy^n$$

with warp factor $h \equiv 1 + \frac{e^{-4A_0(y)}}{c(x)}$ and $c(x) = \mathcal{V}^{2/3}$.

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- ▶ Near the conifold tip, with deformation parameter ϵ and $R_\epsilon^2 \sim \epsilon^{4/3}$:

$$ds_{10}^2 = h^{-1/2} g_{\mu\nu} dx^\mu dx^\nu + h^{1/2} c(x)^{1/2} \left(dr_0^2 + \frac{r_0^2}{8} d\Omega_{S^2}^2 + R_\epsilon^2 d\Omega_{S^3}^2 \right)$$

with $e^{-4A_0(\eta)} = 2^{2/3} \frac{(\alpha' g_s M)^2}{\epsilon^{8/3}} I(\eta)$ and $\frac{1}{(2\pi)^2 \alpha'} \int_{S^3} F_3 = M$.

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- ▶ For weak warping, physical radius at tip: $R_{S^3}^2 \approx \epsilon^{4/3} \mathcal{V}^{1/3}$, so sugra approx $\Rightarrow \epsilon^{4/3} \mathcal{V}^{1/3} \gg \ell_s^2$.

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$$\mathcal{K}(z, \bar{z}) \sim \left[|z|^2 \left(\log \frac{\Lambda_0^3}{|z|} + 1 \right) + \frac{9c'(g_s M)^2}{(2\pi)^4 \mathcal{V}^{2/3}} |z|^{2/3} \right]$$

Usually assumed that second term dominates, but large volume can suppress this warping contribution – we will assume regimes of parameter space where first term is also significant.

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GKP '01

$$W^{cs}/M_p^3 \sim \left[W_0 - \frac{M}{2\pi i} z \left(\log \frac{\Lambda_0^3}{z} + 1 \right) - i \frac{K}{g_s} z \right].$$

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- ▶ Antibrane uplift – *note suppression via small $|z|$:*

$$V_{D3} = c_{D3} \left(\frac{g_s^3}{8\pi} \right) \frac{2}{\mathcal{V}^2} \left\{ 1 + \frac{1}{(2\pi)^4} \frac{2}{c''} \frac{(g_s M)^2}{\mathcal{V}^{2/3} |z|^{4/3}} \right\}^{-1} M_p^4.$$

Deformation modulus stabilisation

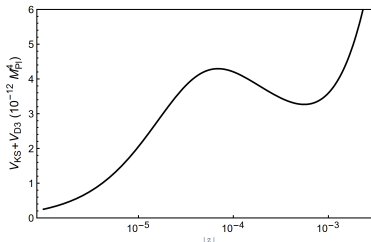
Find minimum in weak warping expansion $1/\beta$:

$$\beta \equiv \frac{\nu^{2/3} |z|^{4/3}}{(g_s M)^2} \frac{(2\pi)^4}{c'} \log \frac{\Lambda_0^3}{|z|}$$

with weak warping $\Leftrightarrow \beta \gg 1$:

$$|z|_{min} \approx |z|_{GKP} \cdot \exp \left\{ -c_{D3} \frac{4K}{3\pi^2 c' M |z_{GKP}|^{4/3} \nu^{2/3}} \right\}$$

\Rightarrow a small shift away from GKP background.



Parameters: $\Lambda_0 = 0.43$, $g_s = 0.17$, $M = 16$, $K = 2$, $\nu = 10^4 \Rightarrow$
solutions: $|z|_{min} \approx 5.57 \times 10^{-4}$ and $|z|_{max} \approx 6.89 \times 10^{-5}$.

Weakly-warped LVS dS solution

Embed in LVS scenario with $\mathcal{V} = \tau_b^{3/2} - \kappa_s \tau_s^{3/2}$:

$$\Delta\mathcal{K}/M_p^2 = -2 \log \left[\mathcal{V} + \frac{\xi}{2} \right] \text{ and } W = W^{CS} + A e^{-\frac{a}{g_s} T}.$$

Balasubramanian, Berglund, Conlon & Quevedo '05

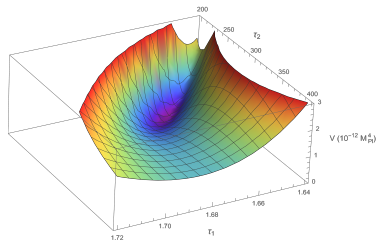
In weakly warped regime, $\beta \gg 1$, find dS solution:

$$\begin{aligned} \mathcal{V} &\approx \tau_b^{3/2} = \frac{3(a\tau_s - g_s)}{4a\tau_s - g_s} \cdot \frac{W_0 g_s \kappa_s \sqrt{\tau_s}}{aA} \cdot e^{\frac{a}{g_s} \tau_s}, \\ \tau_s^{3/2} &\approx \frac{\xi}{2\kappa_s} + \frac{g_s}{3a} + \frac{8\|\Omega\|^2 \mathcal{V}}{9g_s \kappa_s W_0^2} + \mathcal{O}\left(\frac{1}{\beta}\right), \\ \zeta_{min} &\approx \zeta_{GKP} \cdot \exp \left\{ -c_{D3} \frac{4K}{3\pi^2 c' M \zeta_{GKP}^{4/3} \mathcal{V}^{2/3}} \right\}. \end{aligned}$$

provided that:

$$\frac{16\|\Omega\|^2 a}{27\kappa_s} \left(\frac{2\kappa_s}{\xi} \right)^{1/3} < \frac{g_s^2 W_0^2}{\mathcal{V}} < c_{D3} \frac{8a\|\Omega\|^2}{9\kappa_s \sqrt{\tau_s}}.$$

A new dS solution?



W_0	σ	g_s	M	K	Λ_0	κ_s	χ	a	A
2000	0	0.17	16	2	0.43	$\frac{\sqrt{2}}{9}$	-280	$\frac{\pi}{3}$	870

τ_s	τ_b	ζ	\mathcal{V}	V_{crit}
1.80	239	4.17×10^{-4}	3684	1.70×10^{-13}

$m_1^2 \sim m_\zeta^2$	$m_2^2 \sim m_{\tau_s}^2$	m_3^2
9.23×10^{-5}	4.87×10^{-4}	4.32×10^{-11}

M_s	m_{KK}	$m_{3/2}$	M_s^W	m_{KK}^W	$m_{3/2}^W$
4.96×10^{-3}	1.26×10^{-3}	1.11×10^{-3}	2.87×10^{-3}	1.74×10^{-3}	6.40×10^{-4}

Subleading corrections

Bento, Chakraborty, SLP & Zavala '23;
see Junghans '22 for strongly warped case;
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- ▶ Compute $\Delta\mathcal{K}$, ΔW and $V(|z|, \tau_b, \tau_s)$; then $\langle \tau_b \rangle$, $\langle \tau_s \rangle$, $\langle |z| \rangle$, $\langle V \rangle$.
- ▶ Corrections to vevs:

$$\begin{aligned} V_{min} \sim & \alpha - 1 + \mathcal{O}(g_s) + \mathcal{O}(1/\beta) \\ & - C_s^{KK} \cdot \frac{g_s^2}{6\kappa_s} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-2/3} + C_s^{\log} \cdot \frac{a}{6} (11 + 12 \log \nu) + C_1^\xi \cdot \frac{a \log \nu}{\kappa_s} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} \\ & - C_2^\xi \cdot \frac{g_s}{3\hat{\xi}} - C_{flux} \cdot \frac{10a}{9\kappa_s} \frac{KM}{\nu^{2/3}} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} - C_F \cdot \frac{16a}{27\kappa_s} \frac{g_s W_0^2}{\nu^{2/3}} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} \\ & - C_b^{KK} \cdot \frac{8c' a (g_s M)^2}{9\kappa_s} \frac{g_s \zeta^{2/3}}{\nu^{1/3}} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} + C^{\text{con}} \cdot \mathcal{O}(1/\beta^3). \end{aligned}$$

$$\text{with } \nu = \frac{3g_s \kappa_s W_0}{8aA} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{1/3}.$$

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- ▶ Corrections to vevs:

$$\begin{aligned} V_{min} \sim & \alpha - 1 + \mathcal{O}(g_s) + \mathcal{O}(1/\beta) \\ & - C_s^{KK} \cdot \frac{g_s^2}{6\kappa_s} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-2/3} + C_s^{\log} \cdot \frac{a}{6} (11 + 12 \log \nu) + C_1^\xi \cdot \frac{a \log \nu}{\kappa_s} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} \\ & - C_2^\xi \cdot \frac{g_s}{3\hat{\xi}} - C_{flux} \cdot \frac{10a}{9\kappa_s} \frac{KM}{\mathcal{V}^{2/3}} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} - C_F \cdot \frac{16a}{27\kappa_s} \frac{g_s W_0^2}{\mathcal{V}^{2/3}} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} \\ & - C_b^{KK} \cdot \frac{8c' a (g_s M)^2}{9\kappa_s} \frac{g_s \zeta^{2/3}}{\mathcal{V}^{1/3}} \left(\frac{\hat{\xi}}{2\kappa_s} \right)^{-1/3} + C^{\text{con}} \cdot \mathcal{O}(1/\beta^3). \end{aligned}$$

$$\text{with } \nu = \frac{3g_s \kappa_s W_0}{8aA} \left(\hat{\xi}/(2\kappa_s) \right)^{1/3}.$$

- ▶ Consistent dS vacuum needs C to be small – we need to compute these...

Transient de Sitter via interacting Dark Sectors



Elephant in the Room by Banksy

Claim: interacting Dark Sector can source a transient dS with small field displacements and no fine-tuning between potential parameters or in initial conditions, consistently with string swampland conjectures.

Interacting Dark Sectors

Toy model - two interacting dark scalar fields:

$$\mathcal{L} = \frac{1}{2}g^{\mu\nu}\partial_\mu\phi\partial_\nu\phi + \frac{1}{2}g^{\mu\nu}\partial_\mu\psi\partial_\nu\psi + V(\phi, \psi) ,$$

with canonical kinetic terms and a scalar potential of the form:

$$V(\phi, \psi) = V(\phi) + \frac{1}{2}m_\psi^2\psi^2 + \frac{1}{2}\frac{m_{\text{int}}^2}{\Lambda^2}\phi^2\psi^2 .$$

and **Higgs-like hilltop** or **runaway potential** for ϕ :

$$V(\phi) = \rho_{\text{de}} \left(\left(\frac{\phi}{\Lambda} \right)^2 - 1 \right)^2 \quad \text{or} \quad V(\phi) = \rho_{\text{de}} e^{-\frac{\phi}{\Lambda}}$$

With $m_{\text{int}} = 0$ either ϕ or ψ could source slowly-rolling quintessence...
but only with **fine-tuning** to hilltop or **dangerous large field distances**.

but see e.g. Montero, Vafa & Valenzuela '22 for ideas on using large fields and light towers

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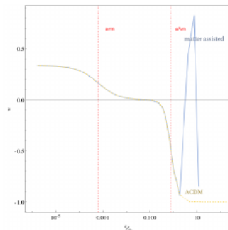
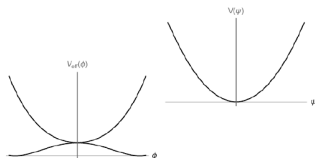
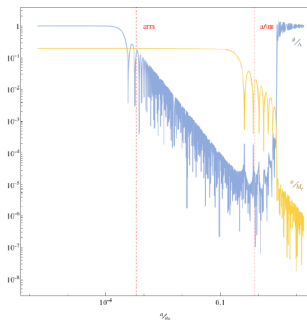
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With $m_{\text{int}} \neq 0$ and ψ behaving as DM, DR or subdominant DE - can stabilise ϕ near $\phi = 0$ to source observed DE as **transient dS!**

A transient de Sitter

Dvali & Kachru '03, Copeland & Rajantie '05, Axenides & Dimoulouos '04 'Locked Inflation/Dark Energy', Gomes, Hardy & SLP to appear

- ▶ Cosmological background $\psi(t) = \psi_0 e^{-3H(t-t_0)/2} \cos(m_\psi t)$; collection of scalar particles oscillating coherently.
- ▶ Background in $\frac{1}{2} \frac{m_{\text{int}}^2 \langle \psi^2 \rangle}{\Lambda^2} \phi^2$ creates false vacuum at $\phi = 0$ where $V_{\text{min}} = \rho_{\text{de}}$ until exit via parametric resonance – analytic understanding via Mathieu's equation.



$$\frac{\Lambda}{M_{\text{pl}}} = 0.1, \quad \frac{m_{\text{int}}}{H_0} = 10^5, \quad \frac{m_\psi}{H_0} = 10, \quad \frac{\psi_0}{M_{\text{pl}}} = 0.4, \quad \frac{\phi_0}{\Lambda} = 1.$$

- ▶ No tuning of initial conditions, no super-Planckian distances, a transient dS with no fine-tuning in Lagrangian parameters!

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- ▶ String model to be worked out... leaves the cc problem...

DM Assisted DE – parameter space

Gomes, Hardy & SLP to appear

