



GABRIELE CASAGRANDE



3-7 July 2023 **STRING PHENO 2023**

Institute for Basic Science Daejeon, South Korea

Gravitino in nonlinear supergravity and the Swampland



Gravitino in nonlinear supergravity and the Swampland

GABRIELE CASAGRANDE

BONNEFOY, CASAGRANDE, DUDAS 2206.13451 & ongoing work with E. Dudas, M. Peloso and Q. Bonnefoy



3-7 July 2023 STRING PHENO 2023

Institute for Basic Science Daejeon, South Korea

I. Nonlinear supergravity

- Nonlinear SUSY and the goldstino
- Superfield constraints
- $\mathcal{N} = 1$ orthogonal SUGRA

I. Nonlinear supergravity

- Nonlinear SUSY and the goldstino
- Superfield constraints
- $\mathcal{N} = 1$ orthogonal SUGRA

II. Gravitino in FRW

- The sound speed c_s
- Potential superluminality: $c_s > 1$
- Unbounded particle production: $c_s = 0$

I. Nonlinear supergravity

- Nonlinear SUSY and the goldstino
- Superfield constraints
- $\mathcal{N} = 1$ orthogonal SUGRA

II. Gravitino in FRW

- The sound speed c_s
- Potential superluminality: $c_s > 1$
- Unbounded particle production: $c_s = 0$

III. Causality constraints in the rigid limit

- The equivalence theorem & positivity bounds
- Improved setup & full causality

I. Nonlinear supergravity

- Nonlinear SUSY and the goldstino
- Superfield constraints
- $\mathcal{N} = 1$ orthogonal SUGRA

II. Gravitino in FRW

- The sound speed c_s
- Potential superluminality: $c_s > 1$
- Unbounded particle production: $c_s = 0$

III. Causality constraints in the rigid limit

- The equivalence theorem & positivity bounds
- Improved setup & full causality

IV. <u>Remarks on $c_s = 0$ </u>

- EFT breakdown
- Backreaction & QFT in curved spacetime

- Spontaneous SUSY breaking:

$$\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle$$

 $|^2 \neq 0$

- Spontaneous SUSY breaking:

$$\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle$$

$|^2 \neq 0 \quad \longleftrightarrow \quad \langle F_i \rangle \neq 0 \quad \lor \quad \langle D_a \rangle \neq 0$

- Spontaneous SUSY breaking:

$$\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle$$

- F-term breaking:

 $\langle F \rangle \neq 0$

$|^2 \neq 0 \quad \longleftrightarrow \quad \langle F_i \rangle \neq 0 \quad \lor \quad \langle D_a \rangle \neq 0$

- Spontaneous SUSY breaking:

$$\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle$$

- F-term breaking: $\langle F \rangle \neq 0 \iff W = f \Phi + \dots$ linear term

$|^2 \neq 0 \quad \longleftrightarrow \quad \langle F_i \rangle \neq 0 \quad \lor \quad \langle D_a \rangle \neq 0$

Spontaneous -SUSY breaking:

$$\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle$$

$|^2 \neq 0 \quad \longleftrightarrow \quad \langle F_i \rangle \neq 0 \quad \lor \quad \langle D_a \rangle \neq 0$

- F-term breaking: $\langle F \rangle \neq 0 \iff W = f \Phi + \dots$ linear term SUSY breaking scale

- Spontaneous SUSY breaking: $\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle|^2$

- F-term breaking: $\langle F \rangle \neq 0 \iff W = f \Phi + \dots$ linear term
- The resulting spectrum contains a massless spin-1/2 state, the *goldstino*:

$$|^2 \neq 0 \quad \longleftrightarrow \quad \langle F_i \rangle \neq 0 \quad \lor \quad \langle D_a \rangle \neq 0$$

+ ... linear term SUSY breaking scale

 $G \equiv \frac{1}{f} \left[\langle F_i \rangle \psi^i + \langle D_a \rangle \lambda^a \right]$

- Spontaneous SUSY breaking: $\langle V \rangle = |\langle F_i \rangle|^2 + \frac{1}{2} |\langle D_a \rangle|^2$

- F-term breaking: $\langle F \rangle \neq 0 \iff W = f \Phi + \dots$ linear term
- The resulting spectrum contains a massless spin-1/2 state, the *goldstino*:
- The goldstino couplings are such that supersymmetry is realised nonlinearly: [VOLKOV, AKULOV '73]

$$|^2 \neq 0 \quad \longleftrightarrow \quad \langle F_i \rangle \neq 0 \quad \lor \quad \langle D_a \rangle \neq 0$$

+ ... linear term SUSY breaking scale

$$G \equiv \frac{1}{f} \left[\langle F_i \rangle \psi^i + \langle D_a \rangle \lambda^a \right]$$

$$\delta_{\epsilon}G = f\epsilon + \frac{i}{f} \left(G\sigma^{\mu}\bar{\epsilon} - \epsilon\sigma^{\mu}\bar{G}\right)\partial_{\mu}G$$

Nonlinear SUSY can be implemented by imposing *superfield constraints*:

Nonlinear SUSY can be implemented by imposing *superfield constraints*:

> $S^2 = 0$ *nilpotent constraint*:

Nonlinear SUSY can be implemented by imposing *superfield constraints*:

<u>nilpotent constraint</u>: $S^2 = 0 \longrightarrow S = \frac{G^2}{2F_S}$

Nonlinear SUSY can be implemented by imposing *superfield constraints*:

<u>nilpotent constraint</u>: $S^2 = 0 \longrightarrow s = \frac{G^2}{2F_S}$ SUSY breaking $\leftrightarrow G \equiv goldstino$

Nonlinear SUSY can be implemented by imposing *superfield constraints*:

Orthogonal constraint:

[KOMARGODSKI, SEIBERG '09]



 $\mathbf{S}\left(\mathbf{\Phi}-\bar{\mathbf{\Phi}}\right)=0$

Nonlinear SUSY can be implemented by imposing *superfield constraints*: -

Orthogonal constraint:

 $\mathfrak{Tm}\phi = \mathfrak{Tm}\phi(G, F_S, \mathfrak{Re}\phi),$



Nonlinear SUSY can be implemented by imposing *superfield constraints*:

Orthogonal constraint:

 $\mathfrak{Tm}\phi = \mathfrak{Tm}\phi(G, F_S, \mathfrak{Re}\phi),$

[KOMARGODSKI, SEIBERG '09]



EFT construction

Constraints in local supersymmetry: -

$$\begin{cases} \mathbf{S}^2 = 0\\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 \end{cases}$$

[FERRARA, KALLOSH, THALER, '15]



φ

+

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)



φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 \end{cases} \longrightarrow \begin{pmatrix} g_{\mu\nu}, \psi_{\mu} \end{pmatrix} + \end{cases}$$

SUGRA multiplet

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)



Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) & + \end{cases}$$

real scalar field

φ

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)



Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

minimal inflation in SUGRA

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0+φ (super-Higgs mechanism) real scalar field



 φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$



 φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

$$\mathcal{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^2 - 3g(\varphi)\right]$$

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\Phi) \mathbf{S} + g(\Phi)$

 $)^{2}] - i \bar{\Psi}_{\mu} \gamma^{\mu\nu\rho} \nabla_{\nu} \Psi_{\rho} - g(\varphi) \bar{\Psi}_{\mu} \gamma^{\mu\nu} \Psi_{\nu} + \mathcal{L}_{\text{torsion}}$



 φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

$$\mathscr{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^{2} - 3g(\varphi)\right]$$

Einstein-Hilbert term

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\Phi) \mathbf{S} + g(\Phi)$

 $)^{2}] - i \bar{\Psi}_{\mu} \gamma^{\mu\nu\rho} \nabla_{\nu} \Psi_{\rho} - g(\varphi) \bar{\Psi}_{\mu} \gamma^{\mu\nu} \Psi_{\nu} + \mathcal{L}_{\text{torsion}}$



 φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

 $\mathcal{L} = -\frac{R}{2} + \frac{1}{2} \partial_{\mu} \varphi \partial^{\mu} \varphi - \left[f(\varphi)^2 - 3g(\varphi)^2 \right] - i \bar{\Psi}_{\mu} \gamma^{\mu\nu\rho} \nabla_{\nu} \Psi_{\rho} - g(\varphi) \bar{\Psi}_{\mu} \gamma^{\mu\nu} \Psi_{\nu} + \mathcal{L}_{\text{torsion}}$

inflaton sector

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$



Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

 $\mathscr{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^{2} - 3g(\varphi)^{2}\right]$ $2 2^{\mu}$

scalar potential $V(\varphi)$

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$

$$^{2}] - i \bar{\Psi}_{\mu} \gamma^{\mu\nu\rho} \nabla_{\nu} \Psi_{\rho} - g(\varphi) \bar{\Psi}_{\mu} \gamma^{\mu\nu} \Psi_{\nu} + \mathscr{L}_{\text{torsion}}$$

+

 φ



 φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

$$\mathcal{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^2 - 3g(\varphi)\right]$$

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\Phi) \mathbf{S} + g(\Phi)$

 $)^{2}] - i \bar{\Psi}_{\mu} \gamma^{\mu\nu\rho} \nabla_{\nu} \Psi_{\rho} - g(\varphi) \bar{\Psi}_{\mu} \gamma^{\mu\nu} \Psi_{\nu} + \mathscr{L}_{\text{torsion}}$

gravitino sector



 φ

Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

$$\mathcal{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^2 - 3g(\varphi)\right]$$

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\Phi) \mathbf{S} + g(\Phi)$

 $)^{2} - i \bar{\Psi}_{\mu} \gamma^{\mu\nu\rho} \nabla_{\nu} \Psi_{\rho} - g(\varphi) \bar{\Psi}_{\mu} \gamma^{\mu\nu} \Psi_{\nu} + \mathcal{L}_{\text{torsion}}$

field-dependent mass parameter



Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

$$\mathscr{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^{2} - 3g(\varphi)^{2}\right] - i\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\nabla_{\nu}\Psi_{\rho} - g(\varphi)\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu} + \mathscr{L}_{\text{torsion}}$$

+

φ

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

& $W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$

torsion term



Constraints in local supersymmetry:

$$\begin{cases} \mathbf{S}^2 = 0 \\ \mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0 & \longrightarrow & \left(g_{\mu\nu}, \psi_{\mu} \right) \end{cases}$$

 $\longrightarrow \mathcal{N} = 1$ orthogonal Supergravity:

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4}\left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$

$$\mathscr{L} = -\frac{R}{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi - \left[f(\varphi)^{2} - 3g(\varphi)^{2}\right] - i\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\nabla_{\nu}\Psi_{\rho} - g(\varphi)\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu} + \mathscr{L}_{\text{torsion}}$$

+

φ

Nonlinear supergravity

[FERRARA, KALLOSH, THALER, '15]

Spectrum in the unitary gauge G = 0(super-Higgs mechanism)

&
$$W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$$

gravitino sector


- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA:

$$S = -\int d^4x \sqrt{-g} \left\{ i \right.$$

 $i\,\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\,\nabla_{\nu}\Psi_{\rho}+m\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu}\bigg\}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA:

$$S = -\int d^4x \sqrt{-g} \left\{ i \right.$$

- Ψ_{μ} contains 4 degrees of freedom:

 $i\,\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\,\nabla_{\nu}\Psi_{\rho}+m\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu}\bigg\}$

Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA: -

$$S = -\int d^4x \sqrt{-g} \left\{ i \right.$$

 $\pm \frac{3}{2}$ transverse components - Ψ_{μ} contains 4 degrees of freedom:

 $i\,\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\,\nabla_{\nu}\Psi_{\rho}+m\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu}\bigg\}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA:

$$S = -\int d^4x \sqrt{-g} \left\{ i \right.$$

- Ψ_{μ} contains 4 degrees of freedom: $\pm \frac{1}{2}$ transverse components $\pm \frac{1}{2}$ longitudinal components

 $i\,\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\nabla_{\nu}\Psi_{\rho}+m\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu}\bigg\}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA:

$$S = -\int d^4x \sqrt{-g} \left\{ i \right.$$

- Ψ_{μ} contains 4 degrees of freedom: $\begin{array}{c} \pm \frac{3}{2} \\ \pm \frac{1}{2} \end{array}$ transverse components $\begin{array}{c} \pm \frac{3}{2} \\ \pm \frac{1}{2} \end{array}$ longitudinal components $\begin{array}{c} \leftarrow \end{array}$ super-Higgs mechanism

 $i\,\bar{\Psi}_{\mu}\gamma^{\mu\nu\rho}\nabla_{\nu}\Psi_{\rho}+m\bar{\Psi}_{\mu}\gamma^{\mu\nu}\Psi_{\nu}\bigg\}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i \bar{\gamma}^{\mu} \partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i \bar{\gamma}^0 \partial_0 + i \left(C_R + i C_I \bar{\gamma}^0 \right) \bar{\gamma}^j \partial_j - am \right] \psi_{1/2} \right\}$$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0\right)\bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

transverse component $\mathscr{L}_{3/2}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0\right)\bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

longitudinal component $\mathscr{L}_{1/2}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0 \right) \bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

 \rightarrow The longitudinal gravitino is characterised by a non-trivial *sound speed*: [KALLOSH, KOFMAN, LINDE, VAN PROEYEN '00] [NILLES, PELOSO, SORBO, '01]

$$\omega_k^2 = c_s^2$$

 $k^{2} + a^{2}m^{2}$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0 \right) \bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

$$c_s^2 = C_R^2 + C_I^2 =$$

$$=\frac{(p-3m^2)^2+4\dot{m}^2}{(\rho+3m^2)^2}$$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I \bar{\gamma}^0 \right) \bar{\gamma}^j \partial_j - am \right] \psi_{1/2} \right\}$$

$$c_s^2 = C_R^2 + C_I^2 =$$

$$\frac{(p - 3m^2)^2 + 4\dot{m}^2}{(\rho + 3m^2)^2} \qquad \rho = \frac{1}{2}\dot{\phi}^2 + V(\phi)$$

$$p = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0 \right) \bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$
$$W = f(\mathbf{\Phi}) \mathbf{S} + g(\mathbf{\Phi})$$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0 \right) \bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2$$
$$W = f(\mathbf{\Phi}) \mathbf{S} + g(\mathbf{\Phi})$$
gravitino mass $m(\varphi) = g(\varphi)$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I\bar{\gamma}^0 \right) \bar{\gamma}^j\partial_j - am \right] \psi_{1/2} \right\}$$

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}}\right)^2 \implies c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)$$
$$W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$$

- Gravitino sector of $\mathcal{N} = 1$ orthogonal SUGRA — *FRW background* $g_{\mu\nu} = a^2(\tau)\eta_{\mu\nu}$:

$$S = \int d^4x \left\{ \bar{\psi}_{3/2} \left(i\bar{\gamma}^{\mu}\partial_{\mu} - am \right) \psi_{3/2} + \bar{\psi}_{1/2} \left[i\bar{\gamma}^0\partial_0 + i\left(C_R + iC_I \bar{\gamma}^0 \right) \bar{\gamma}^j \partial_j - am \right] \psi_{1/2} \right\}$$

 \rightarrow The longitudinal gravitino is characterised by a non-trivial *sound speed*: [KALLOSH, KOFMAN, LINDE, VAN PROEYEN '00] [NILLES, PELOSO, SORBO, '01]

$$K = \bar{\mathbf{S}}\mathbf{S} - \frac{1}{4} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right)^2 \implies c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)$$
$$W = f(\mathbf{\Phi})\mathbf{S} + g(\mathbf{\Phi})$$

The sound speed drives the dynamics of the longitudinal gravitino

$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

\longrightarrow *two major issues*:

$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

\rightarrow *two major issues*:

potential *overproduction* of gravitinos
 [HASEGAWA *et al.* '17] [KOLB, LONG, MCDONOUGH '21]



$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

\rightarrow *two major issues*:

potential *overproduction* of gravitinos
 [HASEGAWA *et al.* '17] [KOLB, LONG, MCDONOUGH '21]

 $c_s = 0$ vanishing sound speed



$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

\rightarrow *two major issues*:

potential *overproduction* of gravitinos
 [HASEGAWA *et al.* '17] [KOLB, LONG, MCDONOUGH '21]

 $c_s = 0$ vanishing sound speed

$$A_k(\tau) = \frac{\omega'_k}{\omega_k^2} = \frac{H}{m}$$
 adiabaticity parameter



$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

\rightarrow *two major issues*:

 potential *overproduction* of gravitinos [HASEGAWA *et al.* '17] [KOLB, LONG, MCDONOUGH '21]

 $c_s = 0$ vanishing sound speed $A_k(\tau) = \frac{\omega'_k}{\omega_k^2} = \frac{H}{m} \propto k$ gravitinos with arbitrary high momentumcan be produced



$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

two major issues:

- 1. potential *overproduction* of gravitinos [HASEGAWA et al. '17] [KOLB, LONG, MCDONOUGH '21]
- 2. potential *superluminal* propagation of gravitinos

 $f(\varphi)^2 \le 2g'(\varphi)^2 \implies c_s > 1$



$$c_s^2 = 1 - \frac{4\dot{\varphi}^2}{\left(\frac{1}{2}\dot{\varphi}^2 + f(\varphi)^2\right)^2} \left(\frac{1}{2}f(\varphi)^2 - g'(\varphi)^2\right)^2$$

two major issues:

- 1. potential *overproduction* of gravitinos [HASEGAWA et al. '17] [KOLB, LONG, MCDONOUGH '21]
- 2. potential *superluminal* propagation of gravitinos





- Subluminal conditions

Subluminal conditions

Positivity bounds on EFT operators

[ADAMS, ARKANI-HAMED, DUBOVSKY, NICOLIS, RATTAZZI '06]



Subluminal conditions

Positivity bounds on EFT operators

Equivalence theorem:

 $\psi_{\mu} \to \frac{1}{m} \partial_{\mu} G$

[FAYET '86] [CASALBUONI ET AL '88-'89] [ADAMS, ARKANI-HAMED, DUBOVSKY, NICOLIS, RATTAZZI '06]

At $E \gg m$, longitudinal gravitino \leftrightarrow goldstino



Subluminal conditions

Positivity bounds on EFT operators

Equivalence theorem:

[FAYET '86] [CASALBUONI ET AL '88-'89]



 $\psi_{1/2}$ couples with the SUSY-breaking scale *f* [ADAMS, ARKANI-HAMED, DUBOVSKY, NICOLIS, RATTAZZI '06]

At $E \gg m$, longitudinal gravitino \leftrightarrow goldstino

enhanced couplings in the limit $M_{\rm P} \rightarrow \infty$ with f fixed



Subluminal conditions

Positivity bounds on EFT operators

Equivalence theorem:

[FAYET '86] [CASALBUONI ET AL '88-'89]



 $\psi_{1/2}$ couples with the SUSY-breaking scale *f*

We can address the superluminality problem in terms of positivity bounds on the operators of the low-energy, SUSY goldstino theory.

[ADAMS, ARKANI-HAMED, DUBOVSKY, NICOLIS, RATTAZZI '06]

At $E \gg m$, longitudinal gravitino \leftrightarrow goldstino

enhanced couplings in the limit $M_{\rm P} \rightarrow \infty$ with f fixed



- Goldstino - real scalar SUSY lagrangian:

$$\begin{aligned} \mathscr{L} &= -f(\varphi)^2 + \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi + i \partial_\mu G \sigma^\mu \bar{G} + (\dots) \\ &+ \frac{1}{2f(\varphi)^2} \left(1 - \frac{2g'(\varphi)^2}{f(\varphi)^2} \right) \left(i \, \partial^\mu G \sigma^\nu \bar{G} \partial_\mu \varphi \right) \right) \end{aligned}$$

 $\bar{G}\partial_{\mu}\varphi\partial_{\nu}\varphi + h.c. - \frac{1}{4f(\varphi)^2} \left(1 - \frac{g'(\varphi)^2}{f(\varphi)^2}\right) \bar{G}^2 \Box G^2$

- Goldstino - real scalar SUSY lagrangian:

$$\mathcal{L} = -f(\varphi)^{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi + i\partial_{\mu}G\sigma^{\mu}\bar{G} + (\dots) \quad \text{higher-order operators} \\ + \frac{1}{2f(\varphi)^{2}}\left(1 - \frac{2g'(\varphi)^{2}}{f(\varphi)^{2}}\right)\left(i\partial^{\mu}G\sigma^{\nu}\bar{G}\partial_{\mu}\varphi\partial_{\nu}\varphi + h.c.\right) - \frac{1}{4f(\varphi)^{2}}\left(1 - \frac{g'(\varphi)^{2}}{f(\varphi)^{2}}\right)\bar{G}^{2} \Box G^{2}$$

- Goldstino - real scalar SUSY lagrangian:

$$\mathscr{L} = -f(\varphi)^{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi + i\partial_{\mu}G\sigma^{\mu}\bar{G} + (\dots) \quad \text{higher-order operators} \\ + \frac{1}{2f(\varphi)^{2}} \left(1 - \frac{2g'(\varphi)^{2}}{f(\varphi)^{2}}\right) \left(i\partial^{\mu}G\sigma^{\nu}\bar{G}\partial_{\mu}\varphi\partial_{\nu}\varphi + h \cdot c \cdot\right) - \frac{1}{4f(\varphi)^{2}} \left(1 - \frac{g'(\varphi)^{2}}{f(\varphi)^{2}}\right)\bar{G}^{2} \Box G^{2}$$

 $\rightarrow Positivity bounds:$ [DINE, FESTUCCIA, KOMARGODSKI '09] $\begin{cases}
1 - \frac{2g'^2}{f^2} \ge 0 \\
1 - \frac{g'^2}{f^2} \ge 0
\end{cases}$

Goldstino - real scalar SUSY lagrangian:

$$\mathscr{L} = -f(\varphi)^{2} + \frac{1}{2}\partial_{\mu}\varphi\partial^{\mu}\varphi + i\partial_{\mu}G\sigma^{\mu}\bar{G} + (\dots) \quad \text{higher-order operators} \\ + \frac{1}{2f(\varphi)^{2}} \left(1 - \frac{2g'(\varphi)^{2}}{f(\varphi)^{2}}\right) \left(i\partial^{\mu}G\sigma^{\nu}\bar{G}\partial_{\mu}\varphi\partial_{\nu}\varphi + h \cdot c \cdot\right) - \frac{1}{4f(\varphi)^{2}} \left(1 - \frac{g'(\varphi)^{2}}{f(\varphi)^{2}}\right)\bar{G}^{2} \Box G^{2}$$

 $\xrightarrow{} Positivity \ bounds:$ [DINE, FESTUCCIA, KOMARGODSKI '09] $\begin{cases} 1 - \frac{2g'^2}{f^2} \ge 0 \\ 1 - \frac{g'^2}{f^2} \ge 0 \end{cases} \implies f(\varphi) \ge 2g'(\varphi)^2 \ c_s < 1 \ condition \end{cases}$ it coincides with the $c_s < 1 \ c_s < 1 \ c_s$



- The possible longitudinal gravitino superlumit of positivity bounds

- The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory

- The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms
- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory

- The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

obstruction to a UV *completion of the theory*

- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory



The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

obstruction to a UV *completion of the theory*

The orthogonal constraint $\mathbf{S}(\mathbf{\Phi} - \mathbf{\bar{\Phi}}) = 0$ is not reliable

- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory



This is a consequence of the *decoupling of the auxiliary field* F_{ϕ} :

The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

obstruction to a UV *completion of the theory*

The orthogonal constraint $\mathbf{S} \left(\Phi - \overline{\Phi} \right) = 0$ is *not reliable*

- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory



This is a consequence of the *decoupling of the auxiliary field* F_{ϕ} : 1. it has no clear physical meaning

The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

obstruction to a UV *completion of the theory*

The orthogonal constraint $\mathbf{S}(\mathbf{\Phi} - \mathbf{\bar{\Phi}}) = 0$ is *not reliable*

- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory



- This is a consequence of the *decoupling of the auxiliary field* F_{ϕ} :
 - 1. it has no clear physical meaning
 - 2. it requires a UV higher-derivative operator

[DALL'AGATA, DUDAS, FARAKOS '16]

The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

obstruction to a UV *completion of the theory*

The orthogonal constraint $\mathbf{S}(\mathbf{\Phi} - \mathbf{\bar{\Phi}}) = 0$ is *not reliable*

- of positivity bounds
- Any non-trivial such bound should not come from a microscopic theory



The orthogonal constraints

- This is a consequence of the *decoupling of the auxiliary field* F_{ϕ} :
 - 1. it has no clear physical meaning
 - 2. it requires a UV higher-derivative operator

[DALL'AGATA, DUDAS, FARAKOS '16]

The possible longitudinal gravitino superluminal propagation has a low-energy realisation in terms

obstruction to a UV *completion of the theory*

raint
$$\mathbf{S}\left(\mathbf{\Phi}-\mathbf{\bar{\Phi}}\right)=0$$
 is not reliable

it may be badly defined from the very beginning

- We introduce a new setup:

- We introduce a new setup:

Generalised constraint: $\mathbf{S}\mathbf{S}\mathbf{Q}_L = \mathbf{0}$

[DALL'AGATA, DUDAS, FARAKOS '16]

- We introduce a new setup:

Generalised constraint: $< S\bar{S}Q_L = 0$



[DALL'AGATA, DUDAS, FARAKOS '16]

- We introduce a new setup:

Generalised constraint: $S\bar{S}Q_L = 0$



[DALL'AGATA, DUDAS, FARAKOS '16]

 $\mathbf{S}\bar{\mathbf{S}}\left(\mathbf{\Phi}-\bar{\mathbf{\Phi}}
ight)=0 \implies \Im m\phi = \Im m\phi(G,F_S,\mathfrak{R}e\phi,F_\phi),$

- We introduce a new setup:



[DALL'AGATA, DUDAS, FARAKOS '16]

- We introduce a new setup:



[DALL'AGATA, DUDAS, FARAKOS '16]

Only the physical components are removed

- We introduce a new setup:



The resulting theory has *automatically satisfied positivity bounds*

Only the physical components are removed

- We introduce a new setup:



The resulting theory has *automatically satisfied positivity bounds* - the sound speed results to be

$$c_s^2 = 1 - \frac{4f(\varphi)^2 \dot{\varphi}^2}{\left(\dot{\varphi}^2 + f(\varphi)^2 + g'(\varphi)^2\right)^2}$$

Only the physical components are removed

- We introduce a new setup:



The resulting theory has *automatically satisfied positivity bounds* - the sound speed results to be

$$c_s^2 = 1 - \frac{4f(\varphi)^2 \dot{\varphi}^2}{\left(\dot{\varphi}^2 + f(\varphi)^2 + g'(\varphi)^2\right)^2}$$

Only the physical components are removed

 $< 1 \ \forall \varphi \qquad right setup to build minimal \\ inflation in Supergravity$

- The sound speed can still vanish, leading to the unbounded gravitinos production

- The sound speed can still vanish, leading to the unbounded gravitinos production

Gravitino Swampland Conjecture: $c_s = 0 \in Swampland$ [KOLB, LONG, MCDONOUGH '21]

- The sound speed can still vanish, leading to the unbounded gravitinos production

$$c_s = 0 \quad \iff \quad g'(\varphi) = 0$$

&
$$\dot{\varphi} = f(\varphi)$$

- The sound speed can still vanish, leading to the unbounded gravitinos production

$$c_s = 0 \quad \iff \quad g'(\varphi) = 0$$

&
$$\dot{\varphi} = f(\varphi)$$

inflaton energy equal to the SUSY-breaking scale

- The sound speed can still vanish, leading to the unbounded gravitinos production

$$c_s = 0 \quad \iff \quad g'(\varphi) = 0$$

This dynamics is in contrast with the EFT range:

& $\dot{\varphi} = f(\varphi)$

inflaton energy equal to the SUSY-breaking scale



- The sound speed can still vanish, leading to the unbounded gravitinos production

$$c_s = 0 \quad \iff \quad g'(\varphi) = 0$$

This dynamics is in contrast with the EFT range: 1. the whole constraint setup relies on SUSY breaking

& $\dot{\varphi} = f(\varphi)$

inflaton energy equal to the SUSY-breaking scale



- The sound speed can still vanish, leading to the unbounded gravitinos production

$$c_s = 0 \quad \iff \quad g'(\varphi) = 0$$

This dynamics is in contrast with the EFT range: 1. the whole constraint setup relies on SUSY breaking 2. before reaching f, the theory starts to see the inflatino χ_{ϕ}

& $\dot{\varphi} = f(\varphi)$

inflaton energy equal to the SUSY-breaking scale



- The sound speed can still vanish, leading to the unbounded gravitinos production

$$c_s = 0 \quad \iff \quad g'(\varphi) = 0$$

This dynamics is in contrast with the EFT range:
 1. the whole constraint setup relies on SUSY breaking
 2. before reaching *f*, the theory starts to see the inflatino χ_φ
 the goldstino-inflatino mixing prevents the overproduction
 [Dudas, Garcia, MAMBRINI, OLIVE, PELOSO, VERNER '21]

& $\dot{\varphi} = f(\varphi)$

inflaton energy equal to the SUSY-breaking scale



- The sound speed can still vanish, leading to the unbounded gravitinos production

Backreaction effect

- The sound speed can still vanish, leading to the unbounded gravitinos production

Backreaction effect

 $G_{\mu\nu} = T_{\mu\nu}$

- The sound speed can still vanish, leading to the unbounded gravitinos production

Backreaction effect $G_{\mu\nu} = T_{\mu\nu} = T^{(\varphi)}_{\mu\nu}$ background

- The sound speed can still vanish, leading to the unbounded gravitinos production

Backreaction effect $G_{\mu\nu} = T_{\mu\nu} = T^{(\phi)}_{\mu\nu} + \langle T^{(\psi)}_{\mu\nu} \rangle$

gravitino contribution

- The sound speed can still vanish, leading to the unbounded gravitinos production

Backreaction effect $G_{\mu\nu} = T_{\mu\nu} = T^{(\phi)}_{\mu\nu} + \langle T^{(\psi)}_{\mu\nu} \rangle$

gravitino contribution

$$T^{\mu\nu} = i\bar{\Psi}_{\rho}\gamma^{\rho(\mu|\alpha} \left(\nabla_{\alpha}\Psi^{|\nu)} - \nabla^{|\nu\rangle}\Psi_{\alpha}\right) - m\bar{\Psi}$$
$$+i\nabla_{\rho} \left(\bar{\Psi}^{\rho}\gamma^{(\mu}\Psi^{\nu)}\right) + i\nabla^{(\nu} \left(\bar{\Psi}^{\mu)}\gamma^{\rho}\Psi_{\rho}\right)$$

 $\bar{\Psi}_{\rho}\gamma^{\rho(\mu}\Psi^{\nu)}+$ $+ig^{\mu\nu}\nabla_{\rho}(\bar{\Psi}_{\alpha}\gamma^{\alpha}\Psi^{\rho})$

- The sound speed can still vanish, leading to the unbounded gravitinos production

<u>Backreaction effect</u> $G_{\mu\nu} = T_{\mu\nu} = T_{\mu\nu}^{(\varphi)} + \langle T_{\mu\nu}^{(\psi)} \rangle \implies g_{\mu\nu} - g_{\mu\nu}$ gravitino contribution $T^{\mu\nu} = i\bar{\Psi}_{\rho}\gamma^{\rho(\mu|\alpha} \left(\nabla_{\alpha}\Psi^{|\nu)} - \nabla^{|\nu\rangle}\Psi_{\alpha}\right) - m\bar{\Psi}_{\rho}\gamma^{\rho(\mu}\Psi^{\nu)} +$ $+i\nabla_{\rho}\left(\bar{\Psi}^{\rho}\gamma^{(\mu}\Psi^{\nu)}\right)+i\nabla^{(\nu}\left(\bar{\Psi}^{\mu)}\gamma^{\rho}\Psi_{\rho}\right)+ig^{\mu\nu}\nabla_{\rho}\left(\bar{\Psi}_{\alpha}\gamma^{\alpha}\Psi^{\rho}\right)$

$$\rightarrow g'_{\mu\nu} = g_{\mu\nu} + \delta g_{\mu\nu}$$
 the metric changes

- The sound speed can still vanish, leading to the unbounded gravitinos production

<u>Backreaction effect</u> $G_{\mu\nu} = T_{\mu\nu} = T_{\mu\nu}^{(\varphi)} + \langle T_{\mu\nu}^{(\psi)} \rangle \implies g_{\mu\nu}$ gravitino contribution $T^{\mu\nu} = i\bar{\Psi}_{\rho}\gamma^{\rho(\mu|\alpha}\left(\nabla_{\alpha}\Psi^{|\nu)} - \nabla^{|\nu|}\Psi_{\alpha}\right) - m\bar{\Psi}_{\rho}z$ $+ i \nabla_{\rho} \left(\bar{\Psi}^{\rho} \gamma^{(\mu} \Psi^{\nu)} \right) + i \nabla^{(\nu} \left(\bar{\Psi}^{\mu)} \gamma^{\rho} \Psi_{\rho} \right)$

$$) + ig^{\mu\nu} \nabla_{\rho} (\bar{\Psi}_{\alpha} \gamma^{\alpha} \Psi^{\rho})$$



- We studied gravitino dynamics in supergravity theories with constrained superfields

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation



- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation
 - 2. unbounded gravitational production




- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:

1. potential superluminal propagation



- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation

superluminality

equivalence theorem

positivity bounds on the SUSY theory

obstruction to a UV *completion of the theory*

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $\mathbf{S}(\mathbf{\Phi} \bar{\mathbf{\Phi}}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation

superluminality

equivalence theorem

improved Supergravity models for inflation:

positivity bounds on the SUSY theory

obstruction to a UV *completion of the theory*

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation

equivalence superluminality theorem

improved Supergravity models for inflation:

$$\begin{cases} \bar{\mathbf{S}}\mathbf{S} \left(\mathbf{\Phi} - \bar{\mathbf{\Phi}} \right) = 0\\ \bar{\mathbf{S}}\mathbf{S} D_{\alpha} \mathbf{\Phi} = 0 \end{cases}$$

positivity bounds on the SUSY theory

obstruction to a UV *completion of the theory*

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation

$$\rightarrow$$
 superluminality $\leftarrow \frac{equivalence}{theorem}$

improved Supergravity models for inflation:

$$\begin{cases} \bar{\mathbf{S}}\mathbf{S} \left(\Phi - \bar{\Phi} \right) = 0 & \text{act only} \\ \bar{\mathbf{S}}\mathbf{S} D_{\alpha} \Phi = 0 & \text{well def} \end{cases}$$

positivity bounds on the SUSY theory

obstruction to a UV *completion of the theory*

on the physical DOF

fined in the UV

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 1. potential superluminal propagation

$$\rightarrow$$
 superluminality $\leftarrow \frac{equivalence}{theorem}$

improved Supergravity models for inflation:

$$\begin{cases} \bar{\mathbf{S}}\mathbf{S} \left(\Phi - \bar{\Phi} \right) = 0 & \text{act only} \\ \bar{\mathbf{S}}\mathbf{S} D_{\alpha} \Phi = 0 & \text{well def} \end{cases}$$

positivity bounds on the SUSY theory

obstruction to a UV *completion of the theory*

on the physical DOF

fined in the UV



- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 2. <u>unbounded gravitational production</u>



- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 2. <u>unbounded gravitational production</u>

 $c_s = 0$

 \rightarrow backreaction on the spacetime geometry $\langle T^{(\psi)}_{\mu\nu} \rangle \implies c_s$ changes

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $S(\Phi \overline{\Phi}) = 0$ suffer from two inconsistencies:
 - 2. <u>unbounded gravitational production</u>

$$\longrightarrow g'(\varphi) = 0 \& \dot{\varphi} = f(\varphi)$$

 $c_s = 0$

 \rightarrow backreaction on the spacetime geometry $\langle T^{(\psi)}_{\mu\nu} \rangle \implies c_s$ changes

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $\mathbf{S}(\mathbf{\Phi} \bar{\mathbf{\Phi}}) = 0$ suffer from two inconsistencies:
 - 2. <u>unbounded gravitational production</u>

$$\rightarrow g'(\varphi) = 0 \& \dot{\varphi} = f(\varphi)$$
 tens income

 $|c_s = 0|$

 \rightarrow backreaction on the spacetime geometry $\langle T_{\mu\nu}^{(\psi)} \rangle \implies c_s$ changes

sion with the SUSY-breaking scale

onsistency with the inflatino mass scale

- We studied gravitino dynamics in supergravity theories with constrained superfields
- Models built via the orthogonal constraint $\mathbf{S}(\mathbf{\Phi} \bar{\mathbf{\Phi}}) = 0$ suffer from two inconsistencies:
 - 2. <u>unbounded gravitational production</u>
 - backreaction on the spacetime geor

$$\rightarrow g'(\varphi) = 0 \& \dot{\varphi} = f(\varphi)$$
 tens incomes

consistency with the inflatino mass scale the fermion mixing prevents the overproduction

 $c_{s} = 0$

metry
$$\langle T^{(\psi)}_{\mu\nu} \rangle \implies c_s$$
 changes

sion with the SUSY-breaking scale

Thank you!